

LA-UR-08- 8039

Approved for public release;  
distribution is unlimited.

*Title:* Superconducting Materials: What the Record Tells Us

*Author(s):* Z. Fisk, H. Ott, J.D. Thompson

*Intended for:* Philosophical Magazine



Los Alamos National Laboratory, an affirmative action/equal opportunity employer, is operated by the Los Alamos National Security, LLC for the National Nuclear Security Administration of the U.S. Department of Energy under contract DE-AC52-06NA25396. By acceptance of this article, the publisher recognizes that the U.S. Government retains a nonexclusive, royalty-free license to publish or reproduce the published form of this contribution, or to allow others to do so, for U.S. Government purposes. Los Alamos National Laboratory requests that the publisher identify this article as work performed under the auspices of the U.S. Department of Energy. Los Alamos National Laboratory strongly supports academic freedom and a researcher's right to publish; as an institution, however, the Laboratory does not endorse the viewpoint of a publication or guarantee its technical correctness.

## Superconducting Materials: what the record tells us

Z. Fisk<sup>a</sup>, H.-R. Ott<sup>b</sup>, J. D. Thompson<sup>c</sup>

<sup>a</sup>Department of Physics and Astronomy, University of California, Irvine, CA 92697

<sup>b</sup>D-Phys, ETH-Hönggerberg, 8093 Zürich, Switzerland

<sup>c</sup>MPA-10, K764, Los Alamos National Laboratory, Los Alamos, NM 87545

(in honor of James L. Smith's 65<sup>th</sup> birthday)

**Abstract:** We argue that the materials which appear to maximize the superconducting transition temperature can be regarded as living at the interface between chemistry and physics.

Although the phenomenon of superconductivity in common metals is mainly affecting the system of itinerant charge carriers in an electrically conducting solid, it is not completely independent from the atomic arrangement in the crystal lattice of the corresponding material. While the latter is dictated by minimizing chemical energies of the order of 1 eV or more, the electronic instability that provides the transition to the superconducting state involves much smaller energy gains. This difference is most obvious in the observation that at the superconducting transition, the thermal expansion coefficient and the compressibility of the solid, both mainly dictated by the bonding of the atoms in the crystal lattice, undergo some but rather insignificant changes, if compared with the resistivity and the magnetic susceptibility for instance. Likewise, the influence of external pressure on the critical temperature  $T_c$  is usually, with typical values of  $\partial T_c / \partial p$  of the order of  $10^{-2}$  K/kbar, rather modest. This is due to the fact that the degenerate Fermi gas of the conduction electrons is intrinsically under very high pressure and the lattice excitation spectrum usually does not change much with varying external pressure.

Apart from superconductivity, other electronic instabilities, such as charge- or spin density wave states have been identified. Some of them seem to be favoured by special symmetries of the crystal lattice, in particular if the atomic arrangements exhibit trends to lower dimensional features, such as stacking of weakly interacting planes or chains. Those are obviously resulting from aspects of chemical bonding, and, hence, these instabilities are dictated by chemistry as much as by physics. In some of the more recent developments of superconducting materials, it has been recognized that, also here, aspects of solid state chemistry play an increasingly important role and that instabilities as those mentioned above are much more closely related.

A striking aspect of superconductivity materials is the remarkable phase space they inhabit. From the alkali metals to the halides, some 50 plus of the elements are superconducting but for more than half of these the transition is observed only under external pressure. For some of them, e.g., Gallium, so called high pressure phases exhibit considerably higher  $T_c$ 's than the ambient pressure variety. Here, external pressure induces changes of the crystal structure, i.e., the atomic arrangement. Intermetallics, B-doped diamond, rare-earth antiferromagnetic compounds, alkali metal  $C_{60}$  fullerenes, organics, a variety of oxides, and recently Fe and Fe-pnictides all exhibit superconductivity. Superconductivity is everywhere but, nevertheless, sparse. For all the ternary intermetallics examined *inter alia*, it is surprising how few are superconductors.

So the central question in superconductivity and the search for new superconducting materials is whether there is anything common to the known superconductors?

Fröhlich's in retrospect inadequate theory of electron-phonon mediated superconductivity<sup>1</sup> immediately provoked the objection that the lattice would be unstable against the electron-phonon coupling strengths needed.<sup>2</sup> This consideration persisted as a limit to achievable  $T_c$ 's, i.e., avoiding intervening lattice instabilities, in the later successful BCS theory (e.g. Cohen and Anderson<sup>3</sup>). In fact among the materials with the highest known  $T_c$ 's at the time, the cubic A15 compounds, such instabilities were known. They appeared as the so-called martensitic transformations at  $T_m$ , just above  $T_c$  observed in  $V_3Si$  and  $Nb_3Sn$ ,<sup>4</sup> with  $T_c$ 's in the range of 17K. In  $V_3Si$  and  $Nb_3Sn$ ,  $T_c$  and  $T_m$  have opposite sign variation with pressure, in the former they approach each other at positive pressure, in the latter at negative pressure. Our viewpoint here is that we have a phase diagram for these superconductors where  $T_c$  appears to be approaching a maximum in the T-P phase diagram at the terminal point of another phase transition line, here  $T_m$ . For  $V_3Si$ ,  $T_m$  exceeds  $T_c$  by a few degrees Kelvin. The growing lattice distortion with decreasing temperature below  $T_m$  is abruptly terminated at the onset of superconductivity at  $T_c$ .<sup>5</sup>

For heavy Fermion materials, it is a generally held belief that all the superconductors are found in the vicinity of a magnetic quantum critical point where an antiferromagnetic ordering temperature has been driven to  $T = 0K$  (Fig. 1). What is observed in general is

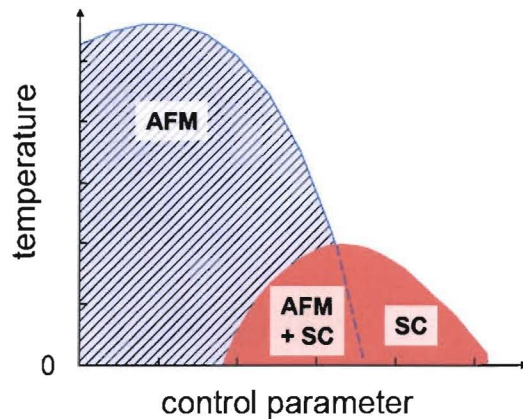


Figure 1. Schematic phase diagram for occurrence of superconductivity in heavy Fermion systems.

very close to this, as shown for example in Fig. 2. What we see is the antiferromagnetic line of phase transitions intersecting near the maximum observed  $T_c$  in the phase diagram. We point out the similarity to the earlier discussed electron-phonon phase diagram.

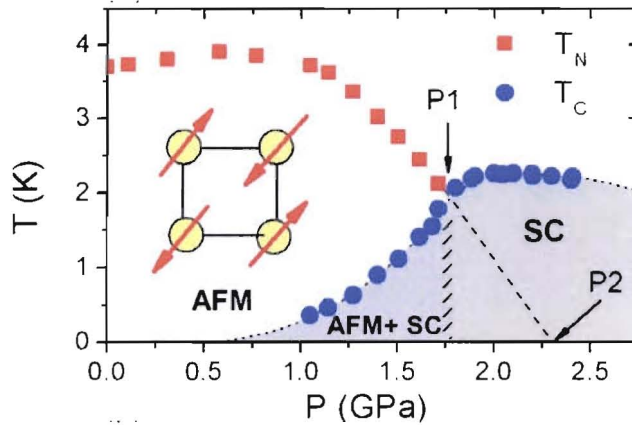


Figure 2. Pressure-temperature phase diagram of CeRhIn<sub>5</sub>.<sup>6</sup>

The generic features of high  $T_c$  cuprates are often discussed on the basis of the phase diagram shown in Fig. 3. Here the so-called pseudogap line intersects the boundary to superconductivity at maximum  $T_c$ , but some claim that the pseudogap is also seen in

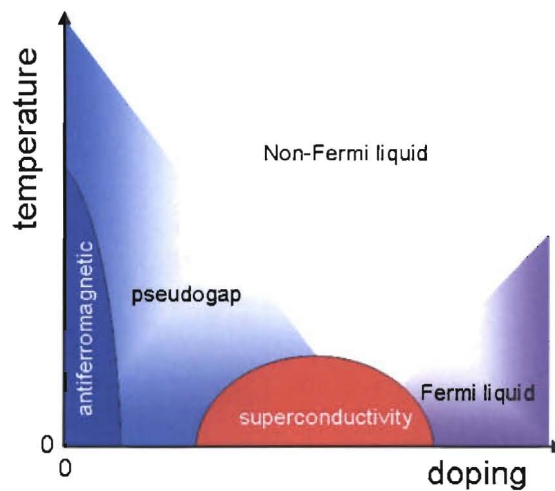


Figure 3. Generic phase diagram of high  $T_c$  cuprates as a function of doping.

tunneling beyond  $T_c^{\max}$ . While most experiments give no indication that the pseudogap line represents a true phase transition, polar Kerr rotation experiments suggest it might.<sup>7</sup> Nevertheless, it is a temperature below which a distinct and measurable change in the electronic properties develops. Again, the similarity to the earlier discussed phase diagrams is apparent with the pseudogap temperature intersecting a maximum in the  $T_c$  versus doping phase diagram.

We can also bring into this discussion the recently discovered Fe-pnictide superconductors.<sup>8</sup> In this set of materials we have two high-temperature transitions in the phase diagram, a structural and a spin-density wave transition. These appear to occur close to each other in temperature, sometimes coinciding and sometimes not exactly at the same temperature. Pressure is found to strongly suppress the transition temperature

for both these phases, and superconductivity does not appear to coexist with the spin density wave. The data at present do not allow one to make the claim that in the ideal case, the spin-density-wave transition will intersect a maximum in the  $T_c$  versus pressure curve for these materials, but the data are suggestive of this. We also mention the organics in which phase diagrams similar to those we have been discussing are found.

The simplest way to think about the phase-diagram similarity discussed above is that we have in all interesting cases where we are trying to maximize  $T_c$  a competing second phase, with superconductivity winning out when the ordering temperature of the competing phase is brought down to the  $T_c$  of the superconducting phase. But this does not explain why this happens in the vicinity of the maximum observed for the superconducting  $T_c$ . It is an old idea that a sufficient increase in the superconducting pairing interaction will lead to some instability limiting  $T_c$ . In the electron-phonon case for instance, it is a lattice distortion which relieves so to speak the tension arising from the large coupling. But the coincidence noted above suggests that it is more useful to think in different terms, namely that the fundamental instability is that of the Fermi surface and that perhaps this instability is related to the mechanism behind the nearby phase whose boundary intersects the maximum  $T_c$ . Again, experiments have often suggested the view that superconductivity is competing for Fermi surface with another phase. But our somewhat different viewpoint is that this competition results from a more fundamental instability which is telling us that the material is balanced between conflicting tendencies corresponding to quite different ground states. This “pairing” of phases is reminiscent of dualities used in discussions of other condensed matter phenomena, such as localized/delocalized, magnetic/non-magnetic, bonding/non-bonding. Another way to perhaps think about the dichotomy is in terms of real space versus momentum space condensation, similar in ways to the differing chemical and physical viewpoints of bonds versus bands. Cohen and Anderson<sup>3</sup> in their early BCS-based discussion of maximum  $T_c$  observed that the limit of large electron-phonon coupling is equivalent to something like a covalent bond. Our point here is simply to note the possibility that what may be limiting  $T_c$  is akin to some kind of localization. From this viewpoint, superconductivity is, then a phenomenon, when pushed to its limit, that lives at the intersect of chemistry and physics.

#### Acknowledgments

Z.F. thanks V. Zlatic and R. B. Laughlin for discussions on this topic and NSF-DMR-0710492 for support. Work at Los Alamos was performed under the auspices of the US DOE/Office of Science.

#### References

1. Fröhlich, H., Phys. Rev. **79**, 845 (1950).
2. According to Zürich folklore, it was Pauli who claimed that if the coupling to the lattice played a decisive role in superconductivity, a  $T_c$  above 30 K was unlikely because of the stability of the lattice. Unfortunately, there is no place in the literature which could be cited. At the same time, in the early 1950ties he suggested pressure experiments to test Fröhlich’s scheme.

3. Cohen, M. L., Anderson, P. W., *Superconductivity in d- and f-Band Metals*, ed. D. H. Douglass (American Institute of Physics, New York, 1972), p. 17
4. Chu, C. W., Testardi, L. R., Phys. Rev. Lett. **32**, 766 (1974); Lim, K. C., Thompson, J. D., Webb, G. W., Phys. Rev. B **27**, 2781 (1983).
5. Ott, H.R., Chandrasekhar, B.S., Seeber, B., Phys. Rev. B **31**, 2700 (1985).
6. Park, T. et al., Nature **440**, 65 (2006).
7. Xia, J. et al., Phys. Rev. Lett. **100**, 127002 (2008).
8. Kamihara, Y. Et al., J. Am. Chem. Soc. **130**, 3296 (2008).