

# Measurements of Time-Dependent $CP$ Asymmetries in $b \rightarrow s$ Penguin Dominated Hadronic $B$ Decays at $B_A B_{AR}$

Pietro Biassoni  
(On behalf of the  $B_A B_{AR}$  Collaboration)

*Università degli Studi and INFN Milano, via Celoria 16, I-20133 Milano, Italy*

**Abstract.** We report measurements of Time-Dependent  $CP$  asymmetries in several  $b \rightarrow s$  penguin dominated hadronic  $B$  decays, where New Physics contributions may appear. We find no significant discrepancies with respect to the Standard Model expectations.

**Keywords:** Charmless Hadronic  $B$  decays, Time-Dependent  $CP$  Violation,  $\sin 2\beta$  Measurement.

**PACS:** 13.25.Hw, 12.15.Hh, 11.30.Er, 13.66.Bc, 14.40.Cs, 13.25.Gv, 13.25.Jx, 13.20.Jf.

## INTRODUCTION

The measurement of  $CP$  violation in  $B$  meson decays provides crucial tests of the Standard Model (SM) and of the Cabibbo-Kobayashi-Maskawa (CKM) mechanism [1].

CKM-suppressed  $b \rightarrow q\bar{q}s$  ( $q = u, d, s$ ) processes are dominated by a single loop (penguin) amplitude, that, assuming penguin dominance and neglecting higher order contributions, is expected to have the same phase  $\beta$  of the CKM-favored  $b \rightarrow c\bar{c}s$  transition [2]. In many extensions of the SM new heavy particles may appear in the loop [3], giving rise to deviations from this expectation. These deviations are expected to be channel dependent. The measurement of the phase difference between  $B^0 \rightarrow K^*(892)^+\pi^-$  and  $\bar{B}^0 \rightarrow K^*(892)^-\pi^+$  can be used to constrain the CKM parameters in the  $(\bar{\rho}, \bar{\eta})$  plane [4].

## TIME-DEPENDENT DECAY RATES

The CKM phase  $\beta$  is accessible experimentally through the interference between the decay of mixed and unmixed  $B$  meson into a  $CP$  eigenstate. This interference is observable through the time evolution of the decay.

In the studies reported in this presentation, one  $B^0$  from  $Y(4S) \rightarrow B^0\bar{B}^0$  is reconstructed in  $\eta'K_S^0$ ,  $\eta'K_L^0$ ,  $\omega K_S^0$ , or  $K_S^0K_S^0K_S^0$   $CP$  eigenstate, or in  $\pi^+\pi^-K_S^0$  or  $K^+K^-K_S^0$  non- $CP$  eigenstate final state ( $B_{sig}$ ), and its vertex fitted using all charged daughter tracks. In  $K_S^0K_S^0K_S^0$  mode, where no charged track is present at  $B^0$  meson decay vertex,  $B_{sig}$  vertex is identified using the  $K_S^0$  reconstructed flight directions and the knowledge of the average interaction point [5]. From the remaining particles in the event we reconstruct the decay vertex of the other  $B$  meson ( $B_{tag}$ ) and identify its flavor, through the analysis of the decay product of  $B_{tag}$  [6].

The distribution of the difference  $\Delta t \equiv t_{CP} - t_{\text{tag}}$  of the proper decay times of  $B$  mesons into  $CP$ -eigenstate final states is given by

$$f(\Delta t) = \frac{e^{-|\Delta t|/\tau}}{4\tau} \{1 \pm [-\eta_f S_f \sin(\Delta m_d \Delta t) - C_f \cos(\Delta m_d \Delta t)]\} \quad (1)$$

where  $\eta_f$  is the  $CP$  eigenvalue of the final state  $f$  and  $\tau$  is the  $B^0$  meson lifetime. The upper (lower) sign denotes a decay accompanied by a  $B^0$  ( $\bar{B}^0$ ) tag, and  $\Delta m_d$  is the mixing frequency.

For three body non- $CP$ -eigenstate final state, the  $CP$ -violating parameters are a function of the position over the Dalitz Plot (DP). In this case Eq. (1) is written as

$$f(\Delta t) = \frac{e^{-|\Delta t|/\tau}}{4\tau} \{ |A|^2 + |\bar{A}|^2 \pm [\eta_f 2\text{Im}[\bar{A}A^*] \sin(\Delta m_d \Delta t) - (|A|^2 - |\bar{A}|^2) \cos(\Delta m_d \Delta t)] \}. \quad (2)$$

Let the decay  $B^0 \rightarrow X_1 X_2 X_3$  proceed through  $N$  intermediate states: the amplitude  $A$  depends only on the Mandelstam invariants  $s_{12}$  and  $s_{23}$ , and in the isobar approximation is

$$A(s_{12}, s_{23}) = \sum_{j=1}^N |c_j| e^{-i\phi_j} R_j(m_j) X_L(|\vec{p}^*| r') X_L(|\vec{q}| r) T_j(L, \vec{p}, \vec{q}) \quad (3)$$

where  $c_j$  and  $\phi_j$  are the relative magnitude and phase of the decay mode  $j$ ,  $R_j(m)$  is the lineshape term,  $X_L$  are Blatt-Weisskopf barrier factors [7],  $T_j$  is the angular distribution,  $\vec{p}$  ( $\vec{q}$ ) is the momentum of the prompt particle (one of the resonance daughters),  $L$  is the orbital angular momentum between  $\vec{p}$  and the resonance momentum, and asterisk denotes  $B$  rest frame. For a decay into a quasi-two-body  $CP$  eigenstate, one can extract the parameters  $\beta_{eff} = \frac{1}{2} \arg(c_k \bar{c}_k^*)$  and  $\mathcal{A}_{ch}(k) = [|\bar{c}_k|^2 - |c_k|^2] / [|\bar{c}_k|^2 + |c_k|^2]$ . For a decay into quasi-two-body non- $CP$  eigenstate, we measure the charge asymmetry and the phase between the two conjugate states  $\Delta\Phi(k) = \arg(c_k \bar{c}_k^*)$ .

A nonzero value of the parameter  $C_f$  or  $\mathcal{A}_{ch}$  would indicate direct  $CP$  violation. In these modes we expect  $-\eta_f S_f \equiv -\eta_f \sin 2\beta_{eff} \approx \sin 2\beta$ . Deviations  $\Delta S_f = S_f - \sin 2\beta$  from this expectation may appear even within the SM [8, 9], and are estimated in several theoretical approaches [8, 10].

## ANALYSIS TECHNIQUE

Analyses presented here are based on a sample of  $465 \times 10^6$   $B\bar{B}$  pairs ( $383 \times 10^6$  for  $B^0 \rightarrow K_S^0 \pi^+ \pi^-$ ), collected at a center-of-mass energy equal to the mass of the  $\Upsilon(4S)$  resonance at the PEP-II asymmetric  $e^+e^-$  collider, at the SLAC National Accelerator Laboratory, and recorded by the  $BABAR$  detector [11]. The  $B$  meson is reconstructed into the above-mentioned  $CP$  eigenstates. The  $B$  meson is kinematically characterized by the variables  $\Delta E \equiv E_B - \frac{1}{2}\sqrt{s}$  and  $m_{ES} \equiv \sqrt{s/4 - |\vec{p}_B|^2}$ , where  $(E_B, \vec{p}_B)$  is the  $B$  four-momentum vector expressed in  $\Upsilon(4S)$  rest frame.

Background arises primarily from random combinations of particles in  $e^+e^- \rightarrow q\bar{q}$  events ( $q = u, d, s, c$ ). We suppress this background with requirements on the event shape variables and on the energy, invariant mass and particle identification signature of the decay products. All events are required to have  $|\Delta t| < 20$  ps and  $\sigma_{\Delta t} < 2.5$  ps.

For each mode, results are obtained from an extended maximum likelihood fit with input variables  $\Delta E$ ,  $m_{\text{ES}}$ ,  $\Delta t$ , and the output of a multivariate discriminant combining different event shape variables. In  $\omega K_S^0$  decay we also use  $\omega$  mass and angular variables into the fit.  $K_L^0$  momentum is determined using a  $B$  mass constraint, hence  $m_{\text{ES}}$  is fully correlated to  $\Delta E$ , and is not used into the fit in  $\eta' K_L^0$  modes. The likelihood for a given event is the sum of the signal, continuum and the  $B$ -background components, weighted by their respective event yields. In  $K_S^0 \pi^+ \pi^-$  and  $K_S^0 K^+ K^-$  modes, a time-dependent DP analysis is performed. The DP model includes  $f_0(980)$ ,  $\rho^0(770)$ ,  $K^{*\pm}(892)$ ,  $(K\pi)_0^{*\pm}$ ,  $f_2(1240)$ ,  $f_x(1300)$ ,  $\chi_{c0}$  ( $f_0(980)$ ,  $\phi(1020)$ ,  $X(1550)$ ,  $f_2(1270)$ ,  $\chi_{c0}$ ,  $D^\pm$ ,  $D_s^\pm$ ) and non resonant component for  $K_S^0 \pi^+ \pi^-$  ( $K_S^0 K^+ K^-$ ) decay mode. In  $K_S^0 K^+ K^-$  analysis, the fit is first performed on the whole DP, and then in the low (high) mass region  $m_{K^+K^-} < 1.1$  GeV/ $c^2$  ( $m_{K^+K^-} > 1.1$  GeV/ $c^2$ ), fixing all the parameters to the values found in the whole DP fit, except the ones involving the  $f_0(980)$  ( $\phi(1020)$ ) resonance.

## RESULTS

In Table 1 and 2 we report the results for  $CP$ -violating parameters in analyses of the decay of a  $B^0$  meson into a  $CP$  eigenstates and a three body non- $CP$  eigenstates final state (DP analyses), respectively [12]. Results for  $K_S^0 K^+ K^-$  and  $K_S^0 K_S^0 K_S^0$  are preliminary.

**TABLE 1.** Results of analyses of  $b \rightarrow s$  decays into  $CP$  eigenstates. For each decay mode we report  $-\eta_f S_f$  and  $C_f$ . The first error is statistical, the second systematic.

Decay Mode	$-\eta_f S_f$	$C_f$
$\eta' K^0$	$0.57 \pm 0.08 \pm 0.02$	$-0.08 \pm 0.06 \pm 0.02$
$\omega K_S^0$	$0.55^{+0.26}_{-0.29} \pm 0.02$	$-0.52^{+0.22}_{-0.20} \pm 0.03$
$K_S^0 K_S^0 K_S^0$	$0.90^{+0.20+0.04}_{-0.18-0.03}$	$-0.16 \pm 0.17 \pm 0.03$

In  $K_S^0 \pi^+ \pi^-$  and  $K_S^0 K^+ K^-$  low mass region, the likelihood function has two minima. In  $B^0 \rightarrow f_0(980) K_S^0$  with  $f_0(980) \rightarrow K^+ K^-$ , the second solution is disfavored by the result from  $f_0(980) \rightarrow \pi^+ \pi^-$ . In  $K_S^0 \pi^+ \pi^-$  analysis we measure  $\mathcal{A}_{ch}(K^*(892)^+ \pi^-) = 0.20 \pm 0.10 \pm 0.02$ , where the first (second) error is statistical (systematic). We also exclude  $-137^\circ < \Delta\Phi(K^*(892)^+ \pi^-) < -5^\circ$  at 95% confidence level.

## CONCLUSIONS

We have reported the results of measurements of  $CP$ -violating parameters in several  $b \rightarrow s$  hadronic  $B$  meson decays. All the results are consistent with the SM. Results are

in agreement with and supersede previous *BABAR* measurements.

**TABLE 2.** Results of DP  $b \rightarrow s$  analyses. For each decay mode we report  $\beta_{eff}$ , and  $\mathcal{A}_{ch}$ , for both solutions. The first error is statistical, the second systematic.

Decay Mode	Solution I		Solution II	
	$\beta_{eff}$ ( $^\circ$ )	$\mathcal{A}_{ch}$	$\beta_{eff}$ ( $^\circ$ )	$\mathcal{A}_{ch}$
$K_S^0 \pi^+ \pi^-$				
$f_0(980)K_S^0$	$36.0 \pm 9.8 \pm 3.0$	$-0.08 \pm 0.19 \pm 0.05$	$56.2 \pm 10.4 \pm 3.0$	$-0.23 \pm 0.19 \pm 0.05$
$\rho^0(770)K_S^0$	$10.2 \pm 8.9 \pm 3.6$	$0.05 \pm 0.26 \pm 0.10$	$33.4 \pm 10.4 \pm 3.6$	$0.14 \pm 0.26 \pm 0.10$
$K_S^0 K^+ K^-$				
Whole DP	$25.2 \pm 4.0 \pm 1.1$	$0.03 \pm 0.07 \pm 0.02$	–	–
High Mass	$29.8 \pm 4.6 \pm 1.7$	$0.05 \pm 0.09 \pm 0.04$	–	–
$\phi K_S^0$	$7.4 \pm 7.4 \pm 1.1$	$0.14 \pm 0.19 \pm 0.02$	$8.0 \pm 8.0 \pm 1.1$	$0.13 \pm 0.18 \pm 0.02$
$f_0(980)K_S^0$	$8.6 \pm 7.4 \pm 1.7$	$0.01 \pm 0.26 \pm 0.07$	$197.1 \pm 10.9 \pm 1.7$	$-0.49 \pm 0.25 \pm 0.07$

## ACKNOWLEDGMENTS

I'd like to thank all my *BABAR* colleagues for their support and in particular Fernando Palombo and Alfio Lazzaro.

## REFERENCES

1. N. Cabibbo, Phys. Rev. Lett. **10**, 531 (1963); M. Kobayashi and T. Maskawa, Prog. Theor. Phys. **49**, 652 (1973).
2. Belle Collaboration, K.-F. Chen *et al.*, Phys. Rev. Lett. **98**, 031802 (2007); *BABAR* Collaboration, B. Aubert *et al.*, Phys. Rev. D **79**, 072009 (2009).
3. Y. Grossman and M. P. Worah, Phys. Lett. B **395**, 241 (1997); D. Atwood and A. Soni, Phys. Lett. B **405**, 150 (1997); M. Ciuchini *et al.*, Phys. Rev. Lett. **79**, 978 (1997).
4. H. J. Lipkin *et al.*, Phys. Rev. D **44**, 1454 (1991); N. G. Deshpande, N. Shina, and R. Shina, Phys. Rev. Lett. **90**, 061802 (2003); M. Ciuchini, M. Pierini, and L. Silvestrini, Phys. Rev. D **74**, 051301(R) (2006); M. Gronau *et al.*, Phys. Rev. D **75**, 014002 (2007).
5. *BABAR* Collaboration, B. Aubert *et al.*, Phys. Rev. Lett. **93**, 131805 (2004).
6. *BABAR* Collaboration, B. Aubert *et al.*, Phys. Rev. Lett. **94**, 161803 (2005).
7. J. Blatt and V. E. Weisskopf, *Theoretical Nuclear Physics* (J. Wiley (New York), 1952).
8. Y. Grossman *et al.*, Phys. Rev. D **68**, 015004 (2003); C.-W. Chiang, M. Gronau, and J. L. Rosner, Phys. Rev. D **68**, 074012 (2003); M. Gronau, J. L. Rosner, and J. Zupan, Phys. Lett. B **596**, 107 (2004); M. Beneke and M. Neubert, Nucl. Phys. B **675**, 333 (2003).
9. D. London and A. Soni, Phys. Lett. B **407**, 61 (1997).
10. M. Beneke, Phys. Lett. B **620**, 143 (2005); H. Y. Cheng, C-K. Chua, and A. Soni, Phys. Rev. D **72**, 014006 (2005), Phys. Rev. D **71**, 014030 (2005); S. Fajfer, T. N. Pham, and A. Prapotnik-Brdnik, Phys. Rev. D **72**, 114001 (2005); A. R. Williamson and J. Zupan, Phys. Rev. D **74**, 014003 (2006); M. Gronau, J. L. Rosner, and J. Zupan, Phys. Rev. D **74**, 093003 (2006).
11. *BABAR* Collaboration, B. Aubert *et al.*, Nucl. Instrum. Methods Phys. Res., Sect. A **479**, 1 (2002).
12. *BABAR* Collaboration, B. Aubert *et al.*, Phys. Rev. D **79**, 052003 (2009). *BABAR* Collaboration, B. Aubert *et al.*, arXiv:0905.3615v1 [hep-ex], submitted to Phys. Rev. D. *BABAR* Collaboration, B. Aubert *et al.*, arXiv:0808.0700v2 [hep-ex].