Frequency-Domain Multiplexed Readout for Superconducting Gamma-Ray Detectors

Jonathan G. Dreyer, Kam Arnold, Trevor M. Lanting, Matt A. Dobbs, Stephan Friedrich, Adrian T. Lee, and Helmuth G. Spieler

Abstract—We are developing a frequency-multiplexed readout for arrays of high-resolution Gamma detectors based on superconducting transition edge sensors (TESs). Each sensor is part of an LCR resonant circuit and is biased at an identifying carrier frequency. Several carrier signals are added and amplified with a single SQUID preamplifier at 4 K. Gamma absorption modulates the amplitude of the carrier, and demodulation at room temperature retrieves the initial temperature evolution of the sensor. This multiplexing system has originally been developed to read out large arrays of bolometers for cosmic microwave background studies. To accommodate the faster Gamma-ray signals, its demodulator bandwidth is being extended to ~20 kHz to allow reading out up to eight TESs with a detector bandwidth of ~10 kHz. Here we characterize the system noise performance and show how this multiplexing scheme can be adapted to read out arrays of superconducting Gamma-ray detectors.

Index Terms—Gamma ray detectors, frequency domain multiplexing, superconducting radiation detectors, transition edge sensors.

I. INTRODUCTION

Gamma ray spectrometers based on superconducting transition edge sensors offer an order of magnitude higher energy resolution than conventional detectors based on high-purity germanium crystals [1]. This can improve the precision of non-destructive isotope analysis by gamma spectroscopy in purity germanium crystals [1]. This can improve the precision of non-destructive isotope analysis by gamma spectroscopy in purity germanium crystals [1]. This can improve the precision of non-destructive isotope analysis by gamma spectroscopy in purity germanium crystals [1]. This can improve the precision of non-destructive isotope analysis by gamma spectroscopy in purity germanium crystals [1]. This can improve the precision of non-destructive isotope analysis by gamma spectroscopy in purity germanium crystals [1]. This can improve the precision of non-destructive isotope analysis by gamma spectroscopy in purity germanium crystals [1]. This can improve the precision of non-destructive isotope analysis by gamma spectroscopy in purity germanium crystals [1]. This can improve the precision of non-destructive isotope analysis by gamma spectroscopy in purity germanium crystals [1]. This can improve the precision of non-destructive isotope analysis by gamma spectroscopy in purity germanium crystals [1]. This can improve the precision of non-destructive isotope analysis by gamma spectroscopy in purity germanium crystals [1]. This can improve the precision of non-destructive isotope analysis by gamma spectroscopy in purity germanium crystals [1]. This can improve the precision of non-destructive isotope analysis by gamma spectroscopy in purity germanium crystals [1]. This can improve the precision of non-destructive isotope analysis by gamma spectroscopy in purity germanium crystals [1].

This work was supported by the U.S. Department of Energy, Office of Nonproliferation Research and Engineering NA-22 under grant LL035-DP, and by the NASA APRA program under grant NNG04WF071. Part of this work was performed under the auspices of the U.S. Department of Energy by University of California Lawrence Livermore National Laboratory under contract No. W-7405-Eng-48. J. G. Dreyer and S. Friedrich are with the Advanced Detector Group at the Lawrence Livermore National Laboratory, Livermore, CA 94550, USA (e-mail: friedrich1@llnl.gov; phone 925-423-1527; fax 925-424-5512). J. D. Dreyer is also with the Department of Nuclear Engineering at the University of California, Berkeley, CA 94720, USA.

This work was supported by the U.S. Department of Energy, Office of Nonproliferation Research and Engineering NA-22 under grant LL035-DP, and by the NASA APRA program under grant NNG04WF071. Part of this work was performed under the auspices of the U.S. Department of Energy by University of California Lawrence Livermore National Laboratory under contract No. W-7405-Eng-48. J. G. Dreyer and S. Friedrich are with the Advanced Detector Group at the Lawrence Livermore National Laboratory, Livermore, CA 94550, USA (e-mail: friedrich1@llnl.gov; phone 925-423-1527; fax 925-424-5512). J. D. Dreyer is also with the Department of Nuclear Engineering at the University of California, Berkeley, CA 94720, USA.

Each TES with resistance $R_{\text{TES}}$ is biased in an $LCR$ resonance circuit with a center frequency $\omega_0 = (LC)^{-1/2}$. The bandwidth $b = R_{\text{TES}}/L$ selected for each tuned circuit is set by the thermal time constant of the sensor. The spacing between the individual bias frequencies is chosen to minimize cross-talk.

Bias currents from multiple carrier sources are summed in a low-value bias resistor $R_{\text{bias}}$ to form the low-impedance voltage source needed to ensure constant voltage bias operation. The bias frequencies range from 300 to 700 kHz.
The amplitude of the carrier associated with each detector is modulated by changes in TES resistance. Since each detector is biased at a different frequency, the individual detector signals are uniquely associated with discrete bands in frequency space. These individual signals are then summed into a single SQUID preamplifier, preserving the amplitude modulation of each carrier. To reduce the dynamic range requirements for the SQUID, a nulling comb that is 180° out of phase with the bias carriers is applied to the SQUID. This suppresses the bias carrier while the modulated sensor signal is preserved for readout. Shunt feedback through $R_{\text{feedback}}$ establishes a low input impedance at the SQUID input. The output of the SQUID amplifier fed to a bank of demodulators, one for each bolometer, that translate the individual bolometer signals into original signal band. Onboard ADCs sample the output signals and the digitized signals are sent to a computer for subsequent signal processing.

### III. IMPLEMENTATION

Implementation of frequency-domain multiplexing for high-resolution Gamma detectors requires the modification of several components in the existing readout system originally designed for microwave bolometers. The major change is due to the higher sampling frequency required for the faster sensors used here. This requires changes to the firmware that controls the readout logic and the bandwidth of the post-detection anti-aliasing filter must be increased from 400 Hz to 20 kHz to match the higher sampling rate. The parameters of the sensors are also chosen to facilitate the required changes. These components are inexpensive and can be modified without significantly changing existing electronics.

The component values of the LC filter are set by the requirements of the readout electronics and TES detectors. These values must be optimized to offer sufficient signal bandwidth and limit cross talk between channels, while still remaining within the maximum frequency range of 1 MHz for multiplexing system. To achieve constant bandwidth each multiplexer branch operates with the same inductance and the capacitance is chosen to set the individual resonant frequencies. We have initially chosen inductances of $L = 16 \mu \text{H}$, which are available as chips that have already been well-characterized and proven to work well in multiplexing systems. The capacitance values are selected according to $C_i = 1/(4\pi f_i L_i)$ where $f_i$ is the center frequency of channel $i$ and $L_i$ is the inductance of 16\mu H. For a spacing of ~50 kHz between bias frequencies, the capacitor values range from 4 to 20 nF. NPO chip capacitors are used, whose temperature dependence has been characterized to 4.2 K. The capacitance linearly decreases to 85% of the room temperature value.

Electrothermal stability requires the LC filter bandwidth to be at least a factor six larger than the associated TES bandwidth, so for a detector time constant of 1 ms, the bandwidth must be greater than $6 \times 160$ Hz. Furthermore, a signal bandwidth of ~10 kHz is required to read out the TESs with high signal-to-noise ratio (Fig. 3). So the bandwidth of each tuned circuit must be increased to accommodate twice the signal bandwidth, as the signal is mirrored in two sidebands above and below the carrier. For an $L$ of 16\mu H, the total resistance (TES and parasitic resistances) in series with the LC circuit is required to be of the order 1 \Omega. Smaller inductors would allow operating TESs with lower resistances.

The Gamma-ray sensors used in the array are Mo/Cu multilayers on thin SiN membranes that provide the weak thermal link to the cold bath. A bulk Sn absorber attached to each sensor increases the detection efficiency of the individual pixels [2, 3]. A crucial feature of multilayer TESs is the ability to select the resistance of the sensor independently of the transition temperature (Fig 2). This characteristic allows for optimizing the resistance required for the readout electronics while still maintaining the desired transition temperature of ~120mK.

![Fig. 2. Critical temperature and square resistance of Mo/Cu multilayer TESs with a Mo/Cu ratio of 4:9 for different multilayer thicknesses. Note that the multilayer approach allows adjusting the sensor resistance while keeping the transition temperature constant.](image-url)
Multilayers with a thickness between 200 and 800 nm have resistances between 0.35 and 0.08 Ω, and we can increase the resistance into the 1 Ω range with a multilayer thickness of ~50 nm.

To first order the phonon noise current in TES Gamma-ray detectors is given by

\[ i_n \approx \sqrt{\frac{4k_BT^2G}{V_{bias}}} \]  

(1)

where \( T \) is the transition temperature of the sensors of 120 mK, \( G \) is the thermal conductance between absorber and TES, approximately 10 nW/K depending on design, and \( 1/V_{bias} \) is the responsivity, the conversion of phonon noise power into the input noise current. For \( V_{bias} \approx 1 \mu V \) we expect a sensor current noise on the order of ~100 pA/√Hz (Fig. 3).

IV. SYSTEM CHARACTERIZATION

To confirm that this frequency-domain multiplexing system can be used to read out large Mo/Cu Gamma-ray TES arrays, a setup consisting of eight 0.46 Ω resistors and LC filters with 16 μH inductors and capacitors ranging from ~5 to 18 nF at 300K was assembled (cf. Fig 1). The setup was placed in a test cryostat and cooled to a temperature of 1.8 K. A variable frequency bias was applied from 100 to 900 kHz, and the output response was recorded (Fig. 4). The measured center frequencies and the widths of the eight peaks correspond to the expected LC resonances to within one percent, taking into account the reduction of the capacitance by ~15% on cooldown. The response can be fine-tuned in future setups depending on the exact value of the TES resistance at the operating point.

The current noise of the multiplexer system was measured with the test resistors at 2.1 K. This noise level was calculated by measuring the spectrum of the output of the SQUID electronics with an HP4195 network analyzer, and referring it to the SQUID input coil as a current noise (Fig 5). The total noise between the filter resonances is measured to be ~5 pA/√Hz, as expected from the SQUID characteristics and the Johnson noise of the feedback resistor. On resonance, a single 0.46 Ω test resistor adds \((4k_BT/R)^{1/2} \approx 16 \text{ pA/√Hz}\) of noise. The measured average current noise on resonance of ~18 pA/√Hz is consistent with the quadrature sum of the resistor’s Johnson noise and the readout noise.

This noise floor of 5 pA/√Hz is more than adequate for TES gamma-ray sensors with a detector noise in the ~50 to 100 pA/√Hz range. Future decreases of the inductance of the bandpass filter and fabrication of ~1 Ω TESs will ensure that the detector bandwidth will cover the full range of frequencies where the signal-to-noise ratio exceeds unity in our gamma detectors. In addition, we are currently changing the firmware...
to increase the sampling rate of the A/D converter and the
demultiplexer to 50 kHz, and are adjusting the Nyquist filters
accordingly.

V. SUMMARY

We have characterized the response of our frequency-
multiplexed readout for superconducting transition edge
sensors. The response agrees with the design values, and the
readout noise floor of 5 pA/√Hz is well below the intrinsic
noise level of typical TES Gamma-ray detectors. We also
show that this frequency domain multiplexing scheme can be
adapted for the readout of arrays of TES gamma ray
spectrometers with moderate modifications. Such arrays will
provide increased detection efficiency without the significant
heat load associated with arrays of non-multiplexed sensors.
Therefore, this development will increase the sensitivity of
TES Gamma-ray spectrometers for nuclear analysis, and of
imaging spectrometers for high-energy astrophysics.

REFERENCES

[1] For an overview of the current status of cryogenic detector development,
[2] O. B. Drury, S. F. Terracol, S. Friedrich, Quantifying the benefits of
Solidi C5, 1468-1479 (2005)
Development for Nuclear Attribution and Non-Proliferation
Applications”, Conference Record, 2004 IEEE Nuclear Science
Symposium, Rome (2004)
Symposium, Puerto Rico (2005)
Macintosh, S. W. Nam, C. D. Reinstsema, L. R. Vale, and M. E. Huber,
[6] Lanting et al. “Frequency-domain multiplexed readout of transition-
edge sensor arrays with a superconducting quantum interference device”