Evaluation of the DRAGON Code for VHTR Design Analysis

prepared by
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Abstract

This letter report summarizes three activities that were undertaken in FY 2005 to gather information on the DRAGON code and to perform limited evaluations of the code performance when used in the analysis of the Very High Temperature Reactor (VHTR) designs. These activities include:

- Use of the code to model the fuel elements of the helium-cooled and liquid-salt-cooled VHTR designs. Results were compared to those from another deterministic lattice code (WIMS8) and a Monte Carlo code (MCNP).

- The preliminary assessment of the nuclear data library currently used with the code and libraries that have been provided by the IAEA WIMS-D4 Library Update Project (WLUP).

- DRAGON workshop held to discuss the code capabilities for modeling the VHTR.
1. INTRODUCTION

A previous assessment of deterministic lattice physics codes revealed there are only a few codes that allow the detailed treatment of the double heterogeneity effect inherent in the Very High Temperature Reactor (VHTR) design at the fuel element (assembly) level. [1] Two such codes, DRAGON and WIMS8, [2,3] are available for use at ANL. The WIMS8 code is a proprietary code that has a yearly fee. The access to the WIMS8 source code is also quite expensive. On the other hand, the DRAGON code is quite attractive to the Gen IV program because it is open software that could be additionally developed to meet the specific needs of the program with only limited incremental cost. The code was developed by researchers at the Ecole Polytechnique de Montreal, Canada and is currently maintained by this institute.

The DRAGON code has a collection of models for simulating the neutronic behavior of a unit cell or a fuel lattice in a nuclear reactor. [2] The typical functionalities found in most modern lattice codes are contained in DRAGON. These include interpolation of microscopic cross sections supplied by means of standard libraries; resonance self-shielding calculations in multidimensional geometries; multigroup and multidimensional neutron flux calculations which can take into account neutron leakage; transport-transport or transport-diffusion equivalence calculations; and modules for editing condensed and homogenized nuclear properties for reactor calculations.

The current version of the code contains three algorithms for the solution of the integral transport equation, ranging from a simple collision probability method coupled with the interface current method to the full collision probability method. The code also performs isotopic depletion calculations. The code user must however supply cross sections in one of the following standard formats: DRAGLIB, MATXS (TRANSX-CTR), WIMSD4, and WIMS-AECL. Macroscopic cross sections can also be read by DRAGON via the input data stream. At ANL, the 69- and 172-group cross section libraries created in WIMSD4-format by the Reduced Enrichment for Research and Test Reactors (RERTR) project are used with the DRAGON code. The depletion chains and types of fission products to be tracked by the code are obtained from the cross section library used. For the DRAGON calculations, cross section data for the heavy
nuclides are tabulated at different temperatures (2-4 points) and all the heavy nuclides contained in the library are treated as resonance materials.

An attractive feature of the DRAGON code is its ability to treat particulate fuel in a graphite matrix in a full-assembly calculation. This capability has been used for modeling the fuel assemblies of block-type, high-temperature gas-cooled thermal reactors and the pebble elements in alternative pebble-bed concepts.

An evaluation of the DRAGON lattice capability for use in the design of the Very High Temperature Reactor (VHTR) system was initiated this year (FY 2005) under the Gen IV program. The activities performed this year can be classified into the following three areas:

- DRAGON code was used to model the fuel elements of the helium-cooled and liquid-salt-cooled VHTR designs. Results were compared to those from another deterministic lattice code (WIMS8) and a Monte Carlo code (MCNP).

- The nuclear data library currently used with the code was evaluated by comparing the code performance against those obtained when libraries made available by the IAEA WIMS-D4 Library Update Project (WLUP) are used. This step was also an evaluation of the WLUP libraries to determine their potential for use in future DRAGON calculations.

- A DRAGON workshop was held to discuss the code capabilities for modeling the VHTR.

These three activities are summarized in Sections 2, 3, and 4, respectively.
2. EVALUATION OF DRAGON CODE FOR GRAPHITE-MODERATED REACTORS

2.1 DRAGON VHTR Deterministic Lattice Models

The DRAGON code does not have an explicit model for representing the prismatic-block VHTR fuel element. In the DRAGON full-assembly model that has been developed at Argonne, the fuel element is formed by a collection of pin-cell sized hexagons. Each pin-cell contains the fuel compact and its surrounding block graphite. When all the fuel and coolant-hole pin-cells are represented, the block graphite content is not totally accounted for and therefore an extra ring of pin-cell sized hexagons is used to represent the remaining graphite. The number density of the graphite in these peripheral cells is modified to preserve the graphite content of the assembly block. Due to the use of the pin-cell sized hexagons, the DRAGON assembly model has jagged boundaries, not the flat boundaries of the hexagonal block.

2.2 Assessment of DRAGON for VHTR and GT-MHR Burner Fuel Elements

The DRAGON code has been used for the modeling of the VHTR/Next Generation Nuclear Plant (NGNP). [4] In that study, two kinds of prismatic fuel elements were considered: fuel elements of the NGNP and GT-MHR Minor Actinide Burner (GT/AD-MHR) designs. Both design concepts are alternative designs of the General Atomics (GA) Gas Turbine-Modular Helios Reactor (GT-MHR). The NGNP is being developed under the Gen-IV program to provide very high temperature gas for both electricity and hydrogen production. The GT/AD-MHR was proposed by GA for deep burning of the minor actinides. The assembly dimensions, materials, and composition data for these designs can be found in Reference 4.

The performance of the DRAGON and WIMS8 codes was evaluated by comparing results from the codes to those from the high fidelity MCNP4C code. The MCNP4C calculations for the fuel element were performed using the ENDF/B-VI based nuclear data library distributed with the code. The calculations were for the cold state (294 K). A million neutron histories are utilized in each MCNP4C calculation, with the first 50,000 histories ignored to allow convergence of the fission source before averaging the $k_e$. 

5
The results for fuel element problems are summarized in Tables 1 and 2. Both NGNP (UC$_{0.5}$O$_{1.5}$ fuel) and GT/AD-MHR results (with TRUO$_{1.7}$ and UO$_2$ fuels) are reported in the tables. Note that the difference between cases is defined using $\Delta \rho = \Delta k/k_1 k_2$. The results in Table 1 show that eigenvalues predicted with the deterministic DRAGON and WIMS8 codes generally agree well with the reference values obtained from MCNP4C calculations with stochastic distribution models. However, WIMS8 consistently underestimates the assembly eigenvalue for all the cases considered, while DRAGON overestimates it. For the GT/AD-MHR UO$_2$ assembly, DRAGON and WIMS8 show excellent agreements with MCNP4C. The DRAGON and WIMS8 errors are increased in the GT/AD-MHR TRU and VHTR/NGNP fuel element problems that have much larger resonance absorption. This suggests a potential need to improve the resonance absorption treatments in DRAGON and WIMS8.

### Table 1. Comparison of Deterministic and Monte Carlo Eigenvalue Results for NGNP and GT/AD-MHR Assembly Problems.

<table>
<thead>
<tr>
<th>Computation model</th>
<th>NGNP, UC$<em>{0.5}$O$</em>{1.5}$</th>
<th>GT-MHR, UO$_2$</th>
<th>GT-MHR, TRUO$_{1.7}$</th>
</tr>
</thead>
<tbody>
<tr>
<td>Stochastic Average</td>
<td>1.53280 ± 0.00082</td>
<td>1.57335 ± 0.00040</td>
<td>1.25838 ± 0.00040</td>
</tr>
<tr>
<td>DRAGON</td>
<td>1.54393 (-470)*</td>
<td>1.57565 (93)</td>
<td>1.26794 (599)</td>
</tr>
<tr>
<td>WIMS8</td>
<td>1.52993 (-122)</td>
<td>1.57121 (-86)</td>
<td>1.25326 (-325)</td>
</tr>
</tbody>
</table>

* Difference between deterministic (DRAGON or WIMS8) and stochastic (MCNP4C) average eigenvalues in pcm.

The double heterogeneity effect for the fuel assemblies have been calculated using the MCNP4C, WIMS8, and DRAGON codes. The double heterogeneity effect measures the difference in the assembly eigenvalue calculations performed with detailed heterogeneous compact model (coated fuel particle layers and compact graphite matrix are explicitly represented) and a homogenized model in which the various material components of the compact have been volume weighted to give a single material composition. For consistent comparison of WIMS8 calculations, the calculation with homogeneous compact model was performed using a two-step procedure employed for the heterogeneous compact model; [1] the super-cell
calculation was done using the PIJ module and the subsequent assembly calculation using the CACTUS module. It is noted that the CACTUS module has an option for treating directly such a singly homogeneous-compact lattice problem.

Results for the pin-cell and assembly calculations showed that the double heterogeneity effect in an assembly problem is slightly smaller than that of a corresponding unit pin-cell problem, but it is still significant. The reduced resonance absorption resulting from additional moderation provided by the graphite outside fuel cells decrease the double heterogeneity effect in assembly problems. The double heterogeneity effects for the assembly are summarized in Table 2. Compared to neutron capture in U-238 of the GT/AD-MHR UO2 assembly, the capture in U-238 of the NGNP assembly is 2.6 times larger due to a higher packing fraction of fuel particles and a bigger size of fuel kernel. (The amount of U-238 in the NGNP assembly is six times larger than that of GT/AD-MHR UO2 assembly.) As a result, the double heterogeneity effect is about two times larger in the NGNP assembly. The double heterogeneity effect is quite large in the TRU fuel assembly because of the large amount of Pu-239 and Pu-240.

<table>
<thead>
<tr>
<th>Computation Model</th>
<th>NGNP, UC_{0.5}O_{1.5}</th>
<th>GT-MHR, UO_{2}</th>
<th>GT-MHR, TRUO_{1.7}</th>
</tr>
</thead>
<tbody>
<tr>
<td>MCNP4C</td>
<td>2.3</td>
<td>1.4</td>
<td>13.1</td>
</tr>
<tr>
<td>DRAGON</td>
<td>2.9</td>
<td>1.5</td>
<td>13.1</td>
</tr>
<tr>
<td>WIMS8</td>
<td>2.7</td>
<td>1.4</td>
<td>12.9</td>
</tr>
</tbody>
</table>

Comparison of the fuel element flux spectra calculated by the codes has also been done and results for two cases are summarized in Figures 1 and 2, for the cases using the detailed heterogeneity model. Generally, the spectra calculated by the deterministic codes are very similar to the results of the MCNP4C calculations. A slight over-estimation of the thermal neutron flux peak is observed for the DRAGON results. The reason for this has not been determined and would require additional study.
Figure 1. GT/AD-MHR TRU Spectrum calculated by MCNP4C, DRAGON and WIMS8.

Figure 2. VHTR/NGNP Assembly Spectrum calculated by MCNP4C, DRAGON and WIMS8.
2.3 Evaluation of the DRAGON Code for Liquid-Salt-Cooled Very High Temperature Reactor (LS-VHTR) Analysis

The study presented in Section 2.2 was performed to evaluate the performance of the DRAGON and WIMS8 code models for the analysis of the helium-cooled VHTR. The study demonstrated that the codes are capable of calculating lattice parameters of interest very accurately compared to Monte Carlo reference solutions in the range of interest. While the helium-cooled VHTR and the liquid-salt cooled VHTR employ the same assembly design, the different characteristics of the coolants suggest the need for further verification of the performance of the lattice codes.

The performance of the deterministic lattice codes DRAGON and WIMS8 for analyzing the LS-VHTR fuel element has been evaluated by comparing results from the codes to those from the Monte Carlo code MCNP4C. [5] The MCNP4C calculations for the fuel assembly were performed using the ENDF/B-VI nuclear data library distributed with the code. In the DRAGON and WIMS8 calculations, the 172 group transport equations were solved. Information on the fuel element design can be found in Reference 5. In the study, the packing fraction, uranium enrichment and coolant-hole diameter are 25%, 10% and 0.953 cm respectively. A lithium enrichment value of 99.995% (Li-7 content) is used. [5] Results from the study are summarized in Tables 3 and 4. In addition to the traditional element eigenvalue, an estimate of the double heterogeneity effect was also obtained. This effect was determined by taking the difference in the multiplication factor ($k_{inf}$) values derived from calculations using the spatially heterogeneous and homogeneous compact models. In the heterogeneous model the coated fuel particles are explicitly represented in the compact. On the other hand, in the homogeneous model, the compositions of the coated fuel particles are smeared with those of the graphite matrix using volume weighting. In the latter model, the self-shielding effect of the fuel at the particle level is not represented.

Table 3 summarizes the eigenvalues calculated by the DRAGON and WIMS8 lattice codes and the MCNP4C Monte-Carlo code. The DRAGON and WIMS8 codes calculated fuel element $k_{inf}$ values that are within 100 pcm and 200 pcm, respectively, of those from MCNP4C calculations. (The reactivity differences were computed using $\Delta \rho = \Delta k/k_1 k_2$.) Components of
these differences come from the different nuclear data files used in the calculations (e.g., JEF2.2 for WIMS8). These differences are similar to those observed in Section 2.2.

Table 3. Comparison of $k_{\text{inf}}$ Values for Fuel Element at Cold State.

<table>
<thead>
<tr>
<th>Compact model</th>
<th>Li Enrichment</th>
<th>Code</th>
<th>Eigenvalue</th>
<th>Difference from reference, pcm $\Delta \rho$</th>
</tr>
</thead>
<tbody>
<tr>
<td>Heterogeneous</td>
<td>99.995 %</td>
<td>MCNP4C</td>
<td>1.53861 ± 0.00067</td>
<td>Reference</td>
</tr>
<tr>
<td></td>
<td></td>
<td>WIMS8</td>
<td>1.53756</td>
<td>-44</td>
</tr>
<tr>
<td></td>
<td></td>
<td>DRAGON</td>
<td>1.53832</td>
<td>-12</td>
</tr>
<tr>
<td>Homogeneous</td>
<td>99.995 %</td>
<td>MCNP4C</td>
<td>1.47886 ± 0.00075</td>
<td>Reference</td>
</tr>
<tr>
<td></td>
<td></td>
<td>WIMS8</td>
<td>1.47681</td>
<td>-94</td>
</tr>
<tr>
<td></td>
<td></td>
<td>DRAGON</td>
<td>1.48216</td>
<td>151</td>
</tr>
</tbody>
</table>

The deterministic and Monte Carlo codes predicted very similar values for the double heterogeneity effect in the LS-VHTR fuel element (difference in $k_{\text{inf}}$ from calculations using the smeared and explicit heterogeneous fuel compact models). The MCNP4C calculation predicted a value of $2.6\%\Delta \rho$ (see Table 4) and the two lattice codes give a deviation of about $0.1\%\Delta \rho$ for these cases; these are surprisingly small differences. For comparison, an earlier MCNP4C calculation for the helium-cooled VHTR fuel element gave a value of $2.3\%\Delta \rho$ for the double heterogeneous effect.

Table 4. Double Heterogeneity Effect ($\Delta \rho$) for LS-VHTR.

<table>
<thead>
<tr>
<th>Code</th>
<th>Lithium Enrichment (%Li-7 Content)</th>
<th>Double Heterogeneity Effect, $%\Delta \rho$</th>
</tr>
</thead>
<tbody>
<tr>
<td>MCNP4C</td>
<td>99.995</td>
<td>2.6</td>
</tr>
<tr>
<td>WIMS8</td>
<td>99.995</td>
<td>2.7</td>
</tr>
<tr>
<td>DRAGON</td>
<td>99.995</td>
<td>2.5</td>
</tr>
</tbody>
</table>

Also in the same LS-VHTR study, [5] the coolant void reactivity coefficient (CVRC) was estimated as a function of burnup for the reference fuel-element design, using the deterministic lattice codes DRAGON and WIMS8. The burnup-dependent CVRC were obtained by branch calculations at the burnup points of interest using full voiding of the coolant to derive the 100% voiding reactivity effect. The result is then normalized to the reactivity change per percent voiding; Monte Carlo calculations were not available for these state. For the fuel element CVRC study, an assembly fuel enrichment of 15% and fuel packing fraction of 25% were used.
Results of the study are summarized in Table 5. Noticeable differences are observed in the CVRC values calculated by the DRAGON and WIMS8 codes. The differences between the two codes were also found to become larger with burnup. The reason for these differences has not been completely investigated in the current study. It is however probably due to differences in basic cross section data (JEF2.2 for WIMS8 and ENDF/B-VI for DRAGON) and differences in material compositions at burnup points arising from different solution methodology and code data. The material composition differences with burnup have been attributed to the differences in the energy released per fission values used in the two codes. This effect is investigated in Section 3.2.

It is noted that the magnitude of the differences in the CVRC is actually small, the largest difference being about 600 pcm over 100% voiding for the case with lithium enrichment of 99.99% and about 300 pcm over 100% voiding for the case with lithium enrichment of 99.995%. The results also show that the magnitude of the positive CVRC is small, being less than $1 per 100% voiding (assuming a delayed neutron fraction value of 600 to 700 pcm). Further evaluations of the differences with burnup, however, need to be performed.

<table>
<thead>
<tr>
<th>Burnup (GWd/t)</th>
<th>WIMS8</th>
<th>DRAGON</th>
<th>WIMS8</th>
<th>DRAGON</th>
</tr>
</thead>
<tbody>
<tr>
<td>0.0</td>
<td>3.6</td>
<td>5.4</td>
<td>0.2</td>
<td>1.5</td>
</tr>
<tr>
<td>15.0</td>
<td>3.0</td>
<td>2.9</td>
<td>-0.4</td>
<td>-0.2</td>
</tr>
<tr>
<td>30.0</td>
<td>1.6</td>
<td>0.4</td>
<td>-1.8</td>
<td>-2.4</td>
</tr>
<tr>
<td>45.0</td>
<td>0.6</td>
<td>-2.0</td>
<td>-3.0</td>
<td>-4.2</td>
</tr>
<tr>
<td>60.0</td>
<td>-0.2</td>
<td>-4.2</td>
<td>-4.0</td>
<td>-6.1</td>
</tr>
<tr>
<td>90.0</td>
<td>-1.0</td>
<td>-7.1</td>
<td>-5.3</td>
<td>-8.6</td>
</tr>
</tbody>
</table>

Additionally, the results indicate that the DRAGON and WIMS8 codes give similar trends for the CVRC. It is observed that the coefficient is positive at the beginning of cycle and decreases with increasing burnup. At a burnup between 15 and 20 GWd/t, the fuel element CVRC becomes negative in the case using a lithium enrichment of 99.995%.
3. EVALUATION OF DRAGON LIBRARIES

The DRAGON code can read several standard library formats such as DRAGLIB, MATXS, WIMS-D, and WIMS-AECL. Macroscopic cross sections can also be read via the input data stream. [2] The 69- and 172-group cross section libraries used with the DRAGON code at ANL have been based on the WIMS-D4 format. These ANL libraries will be referred to as WIMS-ANL in this section. Previous clean comparisons of the DRAGON code results (using the WIMS-ANL library) to other lattice and Monte Carlo codes have shown that DRAGON has promise as a lattice capability for analyzing VHTRs (see for example Sections 2.2 and 2.3). However, more detailed evaluations have revealed the need for additional assessment of the neutron cross sections that are typically used with it at ANL (no data library is distributed with the DRAGON code); see the comparison of the CVRC variation with burnup in Section 2.3.

The WIMS-ANL library is based on the WIMSD4 cross-section data-library format that is one of the popular ones used in several lattice codes (DRAGON, WIMS-D4, etc). This library needs to be updated as new revisions to the base nuclear data file (ENDF/B in the U.S. case) are released. The validation of the WIMS-ANL file requires great effort and alternative approaches are being investigated to minimizing this effort at ANL. At the same time, WIMS-format libraries based on recent evaluated nuclear data files have been created for general use under the IAEA WIMS-D Library Update Project (WLUP) since 1990. Several libraries based on this format are currently open to the public through an IAEA website (http://www-nds.iaea.org/wimsd/downloads.htm). It was decided at ANL to obtain these WLUP files, exercise (assess) them to determine their pedigree and then make a determination of their applicability to reactor physics calculations for the VHTR design. A study was therefore conducted to evaluate the performance of the WIMS libraries available from ANL and the IAEA WLUP by intercomparison of DRAGON code results, and for some cases comparison to MCNP4C (Monte Carlo) results. The results of that study are summarized in this section.

3.1 General Descriptions of WIMS-D Format Libraries

The IAEA supported WIMS-D Library Update Project (WLUP) had the goal of providing libraries compatible with the WIMS-D family of codes or equivalent lattice codes. This would
enable the use of the most recent evaluated nuclear data files (ENDF/B-VI.8, JEF2.2, JENDL3.2, etc) for both thermal research and power reactor calculations. Table 6 summarizes the WIMS-D format libraries available for use in the current study and sources of the evaluated nuclear data files. In this work, most of the DRAGON calculations were performed using the 69- and 172-group WIMS-ANL and IAEA recommended libraries.

<table>
<thead>
<tr>
<th>Library name</th>
<th>Groups</th>
<th>Description</th>
</tr>
</thead>
<tbody>
<tr>
<td>WIMS-ANL</td>
<td>69</td>
<td>ANL developed based on ENDF/B-VI data</td>
</tr>
<tr>
<td></td>
<td>172</td>
<td></td>
</tr>
<tr>
<td>IAEA</td>
<td>69</td>
<td>Recommended by IAEA-WLUP participants</td>
</tr>
<tr>
<td></td>
<td>172</td>
<td></td>
</tr>
<tr>
<td>IAEA-ENDFB6</td>
<td>69</td>
<td>Library based on ENDF/B-VI evaluated nuclear data</td>
</tr>
<tr>
<td></td>
<td>172</td>
<td></td>
</tr>
<tr>
<td>IAEA-JEF2.2</td>
<td>69</td>
<td>Library based on JEF2.2 evaluated nuclear data</td>
</tr>
<tr>
<td></td>
<td>172</td>
<td></td>
</tr>
<tr>
<td>IAEA-JENDL3</td>
<td>69</td>
<td>Library based on JENDL3.2 evaluated nuclear data</td>
</tr>
<tr>
<td></td>
<td>172</td>
<td></td>
</tr>
</tbody>
</table>

Some similarities and differences exist between the nuclear-data library files. Both the WIMS-ANL and IAEA recommended libraries have been created using the WIMSR module of the NJOY code. Generally, most of the cross sections types (transport, absorption, $P_1$ transport-corrected scattering matrix, resonance integrals, potential cross sections, etc) included in the libraries are the same. However, the total number of fission products, energy released per fission values, and actinide chains (in particular, Am-241, U-233, U-238, etc) are different.

The resonance parameters of the WIMS-D format library are created assuming that the intermediate resonance approximation and equivalence theorem are used in the lattice code employing the library. In the library, the resonance integrals are tabulated as functions of temperature and background cross section ($\sigma_b$). The intermediate parameters (i.e., Goldstein-Cohen $\lambda$ factor) are stored for each resonance absorber and resonance group. [6]

Since some of the fission products have a significant impact on the nuclear performance of the core, they must be represented explicitly. Although individually unimportant, others have a cumulative effect that cannot be neglected; usually they are lumped together and called pseudo
fission product. In the IAEA WLUP libraries, 56 nuclides are explicitly represented and 79 nuclides are lumped into a pseudo fission product. The criteria for selecting the fission products that are represented explicitly in the WLUP are provided in Reference 6. Thirty-four fission products are explicitly represented in the WIMS-ANL library. Table 7 compares the fission yields for the primary fissionable nuclides (U-235 and Pu-239) contained in the IAEA WLUP and WIMS-ANL libraries. Values from the WIMS8 library are also included for comparison. In this table, the “IAEA” library denotes the WIMS-D format library recommended by the WLUP participants. Note that the cross sections for nuclides contained in the WIMS8 library were created using JEF2.2 nuclear data. [6] Although about 100 fission products are explicitly represented in the WIMS8 library, only the nuclides appearing in the IAEA library are compared in Table 7.

In the ANL library, many more fission products are ignored compared to the IAEA libraries because the fission yield values for these nuclides are negligibly small and their reactivity impacts have been found to be small in core calculations. Except for the lumped fission product, the fission yield values of the IAEA and WIMS-ANL libraries are quite similar because both fission yield values have been taken from the ENDF/B-VI nuclear data file. However, the fission yield values of the WIMS8 library are slightly different from those of the other two because of the different sources.

The effective energy released per fission is required for burnup calculations. In the IAEA WLUP WIMS-D format libraries, the effective energy released per fission (simply Q-value) includes the contributions from fission, from the kinetic energy of incident neutrons, and from the decay of the capture products:

\[ Q^i = Q_f^i + Q_c^i + Q_n^i, \]

where \( Q_f \) is the energy released in fission excluding the energy carried away by neutrinos and the kinetic energy of incident neutrons, \( Q_c \) is the contribution from the decay of the capture products, and \( Q_n \) is the kinetic energy of the incident neutrons. Index \( i \) refers to a fissionable nuclide in all cases. Among the three contributions, the \( Q_f \) is dominant.
Table 7. Explicitly Represented Fission Products and Fission Yields.

<table>
<thead>
<tr>
<th></th>
<th>IAEA</th>
<th>WIMS-D Format Library</th>
<th>ANL</th>
<th>WIMS8 library</th>
</tr>
</thead>
<tbody>
<tr>
<td></td>
<td>U235</td>
<td>Pu239</td>
<td>U235</td>
<td>Pu239</td>
</tr>
<tr>
<td>Krypton-83</td>
<td>0.005</td>
<td>0.003</td>
<td>0.005</td>
<td>0.003</td>
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<tr>
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<td>0.048</td>
<td>0.065</td>
<td>0.048</td>
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<tr>
<td>Technetium-99</td>
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<td>0.062</td>
<td>0.061</td>
<td>0.062</td>
</tr>
<tr>
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<td>0.060</td>
<td>0.052</td>
<td>0.060</td>
</tr>
<tr>
<td>Ruthenium-103</td>
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<td>0.030</td>
<td>0.070</td>
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<tr>
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<td>N/A</td>
</tr>
<tr>
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<td>N/A</td>
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<tr>
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<td>0.003</td>
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<tr>
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<tr>
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<td>0.000</td>
</tr>
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<tr>
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<td>0.039</td>
<td>0.030</td>
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<td>0.011</td>
<td>0.009</td>
</tr>
<tr>
<td>Promethium-147*</td>
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<td>0.011</td>
<td>0.012</td>
<td>0.011</td>
</tr>
<tr>
<td>Promethium-148</td>
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<td>0.000</td>
<td>0.000</td>
</tr>
<tr>
<td>Promethium-148*</td>
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<td>0.000</td>
<td>0.000</td>
<td>0.000</td>
</tr>
<tr>
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<td>0.011</td>
<td>0.012</td>
<td>0.011</td>
<td>0.012</td>
</tr>
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<td>Samarium-147</td>
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<td>0.002</td>
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<td>N/A</td>
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<td>0.000</td>
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<tr>
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<td>N/A</td>
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<td>Dysprosium-160</td>
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<td>0.000</td>
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<td>N/A</td>
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<tr>
<td>Dysprosium-161</td>
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<td>0.000</td>
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<td>N/A</td>
</tr>
<tr>
<td>Dysprosium-162</td>
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<td>N/A</td>
<td>N/A</td>
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<td>Dysprosium-163</td>
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<td>N/A</td>
<td>N/A</td>
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<td>Dysprosium-164</td>
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<td>N/A</td>
<td>N/A</td>
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<tr>
<td>Holmium-165</td>
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<tr>
<td>Erbium-166</td>
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<td>N/A</td>
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<td>Erbium-167</td>
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<td>N/A</td>
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<tr>
<td>Lumped FP</td>
<td>1.069</td>
<td>0.867</td>
<td>0.304</td>
<td>0.329</td>
</tr>
</tbody>
</table>

*Decays to meta-stable Promethium-148
In Table 8, the Q-values of the different libraries are compared. For the comparison, the Q-values used in the WIMS8 library are included in this table. The Q-values of the IAEA and WIMS8 libraries are quite similar, while those of the WIMS-ANL library are consistently smaller than the others. This indicates that the Q-values of the WIMS-ANL library do not include the minor contributions (such as capture or kinetic energy of incident neutrons, or both). The impact of the different Q-values will be discussed in Section 3.2.

Table 8. Comparison of Q-Values (MeV per Fission).

<table>
<thead>
<tr>
<th>Library</th>
<th>ANL WIMS-D</th>
<th>IAEA WIMS-D</th>
<th>WIMS8</th>
</tr>
</thead>
<tbody>
<tr>
<td>Th232</td>
<td>188.5</td>
<td>192.8</td>
<td>192.8</td>
</tr>
<tr>
<td>Pa233</td>
<td>189.8</td>
<td>196.6</td>
<td>196.6</td>
</tr>
<tr>
<td>U233</td>
<td>191.0</td>
<td>199.1</td>
<td>199.1</td>
</tr>
<tr>
<td>U234</td>
<td>190.3</td>
<td>199.2</td>
<td>199.2</td>
</tr>
<tr>
<td>U235</td>
<td>193.7</td>
<td>201.7</td>
<td>201.7</td>
</tr>
<tr>
<td>U236</td>
<td>194.5</td>
<td>201.8</td>
<td>201.8</td>
</tr>
<tr>
<td>U237</td>
<td>200.0</td>
<td>202.0</td>
<td>202.0</td>
</tr>
<tr>
<td>U238</td>
<td>198.1</td>
<td>205.0</td>
<td>205.0</td>
</tr>
<tr>
<td>Np237</td>
<td>196.4</td>
<td>205.1</td>
<td>205.1</td>
</tr>
<tr>
<td>Np239</td>
<td>190.0</td>
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<td>208.0</td>
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<tr>
<td>Pu238</td>
<td>197.4</td>
<td>208.5</td>
<td>208.5</td>
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<tr>
<td>Pu239</td>
<td>199.9</td>
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<td>210.0</td>
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<tr>
<td>Pu240</td>
<td>199.5</td>
<td>210.4</td>
<td>210.4</td>
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<tr>
<td>Pu241</td>
<td>202.0</td>
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<tr>
<td>Pu242</td>
<td>201.6</td>
<td>213.1</td>
<td>213.1</td>
</tr>
<tr>
<td>Am241</td>
<td>202.3</td>
<td>214.7</td>
<td>214.7</td>
</tr>
<tr>
<td>Am242m</td>
<td>202.3</td>
<td>215.8</td>
<td>215.8</td>
</tr>
<tr>
<td>Am243</td>
<td>203.6</td>
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<td>Cm242</td>
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<td>Cm244</td>
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<tr>
<td>Cm245</td>
<td>200.0</td>
<td></td>
<td>221.1</td>
</tr>
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</table>

Figure 3 illustrates the actinide chains contained in the IAEA WIMS-D format libraries. The WIMS-D format does not contain the \((n,2n)\) or \((n,3n)\) cross sections, but the total reaction rate is preserved by correcting the absorption cross section. However, the neglect of these reactions might result in the underestimation of some nuclide densities (for examples, U-237, Pu-238, etc). To compensate the underestimation, two pseudo nuclides are introduced in the U-233 and U-238 chains in Figure 3. In WIMS-D format libraries, these pseudo nuclides are treated as
fission products of U-233 and U-238 and decay instantly to U-232 and U-237, respectively. The fission yield values for these pseudo nuclides are obtained from the ratio of \((n,2n)/(n,f)\) reactions.

There are two states of Am-242: stable (Am-242g) and meta-stable (Am-242m). The ground state Am-242 decays to Cm-242 (beta decay with 16.02 hours half-life) and Pu-242 (electron capture with 4.2 ms half-life). The branching ratio to Am-242m is dependent on the neutron spectrum, but it is about 14% in a typical thermal spectrum. [3] In the IAEA libraries, a simplification is introduced to represent the Am-241 and higher actinide chains because only one nuclide can be specified as the capture product in the WIMS-D format.

The meta-stable isotope (Am-242m) in the IAEA libraries is specified as a capture product explicitly with a branching ratio \(B_{Am2m}\), while the ground state isotope \(^{242}\text{Am}\) in Figure 3) is treated as a fission product with a fission yield value of \((c/f) \times (1-B_{Am2m})\). Here, \(c/f\) denotes the capture-to-fission ratio of Am-241. In Figure 3, \(^{242}\text{Am}\) decays to Cm-242, or goes to Am-243 after neutron capture, and the electron capture (i.e., decay to Pu-242) of Am-242g is ignored.

**Figure 3. Actinide Chains of IAEA WLUP Libraries. [6]**
Similar simplification is observed in the ANL library representation of the Am-241 and higher actinide chains. The meta-stable isotope (Am-242m) is specified as the capture product with a branching ratio $B_{Am^{2m}}$, but the ground state isotope is skipped and Am-241 decays to Cm-242 directly with $(1 - B_{Am^{2m}})$ branching ratio. Thus, ANL library ignores the production of Pu-242 from the electron capture of Am-242.

In the WIMS8 actinide chains, Am-242 is bypassed because its half-life is very short and Am-241 goes to Am-242m, Cm-242, and Pu-242 directly with appropriate branching ratios. For a typical thermal system, these branching ratios are 14%, 71%, and 15%, respectively. This indicates that about 15% of Pu-242 production from the Am-241 is ignored in both WIMS-D format libraries.

3.2 Impact of Q-Values

The comparison of the DRAGON and WIMS8 calculated CVRC variation with burnup in Section 2.3 shows the increasing deviation of the code results for this parameter with burnup. This trend has previously been observed for other parameters (even for the fuel element eigenvalue, $k_{inf}$). It has also been observed that when a constant flux assumption is made in the DRAGON and WIMS8 code runs, the time evolutions of the $k_{inf}$ from the two codes are similar. Additionally, since the initial $k_{inf}$ is typically similar for the two codes, these trends suggest the need to evaluate the depletion chain data (fission product yields, effective energy released per fission [Q-values], decay constants, etc.). The Q-values are important to obtaining accurate flux levels in lattice and whole-core neutronics calculations; the flux level has a direct effect on the evolution of the fuel mass composition over depletion time, which in turn could affect the core reactivity.

In the ANL utilization of the DRAGON code, the depletion chain data are taken directly from the WIMS-ANL cross-section library used with the code. It has however been observed that different code libraries might have different Q-values. Table 8 is a comparison of the Q-values contained in the WIMS-ANL library (69- or 172-group) to those in the IAEA WIMS-D and WIMS8 libraries. It is observed that the latter libraries have identical Q-values, while the WIMS-ANL library values are different.
In order to evaluate the impact of the differences in the Q-values of Table 8, we have developed a procedure to modify manually the depletion chain data during a DRAGON run. This allows overriding the data contained in the cross section (binary-file) library. This procedure uses the “descdepl” option of the “LIB:” module of DRAGON. To first test this user defined depletion chain algorithm, the same data contained in the cross-section library has been collected and used in the DRAGON input deck. The results (eigenvalue and nuclide number densities) obtained with the code were then compared with those obtained using the depletion chain data in the cross-section library. If the algorithm is working well and the depletion data (i.e., fission product yields, decay constants, Q-values, etc) of the input deck are identical to the data in the library, the same results should be obtained. For this study, the WIMS-ANL library was used as the reference library and a homogeneous cell problem (“TCWD08-burnup of homogeneous cell”) of the DRAGON code package (manual) [2] was utilized.

![Graph showing comparison of eigenvalues and burnup](image)

**Figure 4. Comparison of Eigenvalues from Different Implementations of Depletion Chains.**

Figure 4 shows the eigenvalues and differences between the two depletion cases. The result labeled “WIMSD4 depletion chain” denotes the case using DRAGON with the depletion chains specified in the library, while “Manual depletion chain” denotes the case using DRAGON with the depletion chains specified in the input deck. The eigenvalues ($k_{inf}$) obtained for the cases
are quite similar for the whole depletion period (maximum difference less than 3 pcm Δk). In addition, nearly identical number densities were observed. The results indicate that the user specified depletion chains for the DRAGON code are correctly specified and employed as input to the code.

To evaluate the impact of the different isotopic Q-values on the VHTR fuel element results, a depletion problem for the VHTR has been employed. The VHTR fuel element geometry is identical to that used in Section 2.2; the fuel kernel size is 350 µm and the fuel packing fraction in the graphite matrix is 0.289. The uranium enrichment of the fuel is 10% U-235. The VHTR fuel element eigenvalue letdown curves obtained with the DRAGON code using different libraries (and Q-values) are compared in Figure 5. It is observed that the results obtained with the IAEA WIMS-D libraries (using the same Q-values) are quite similar. However, as expected, that obtained using the WIMS-ANL library is different (lower) than the others, with the difference increasing with burnup. The smaller Q-values of the WIMS-ANL library compared to the others gives a higher flux level; thus the depletion rates of the fissile nuclides increase (or the $k_{inf}$ decreases) with the depletion calculation.

![Figure 5. Depletion Results with Original Q-Values of Libraries.](image)

When the Q-values of the WIMS-ANL library are replaced manually with those from the WIMS8 library, closer agreements in the depletion results are observed. This is shown in
Figure 6. Given the level of code and data qualification that has been done for the WIMS8 library in the U.K. program, it has been decided to use the WIMS8 library Q-values in future DRAGON depletion calculations.

![Figure 6. Depletion Results with Modified WIMS-ANL Q-Values.](image)

3.3 Performance of WIMS-D Libraries

Another potential source of differences in the results obtained with different libraries is the cross section data contained in the libraries. Three numerical benchmark problems have been used to evaluate the differences in results obtained using the WIMS-ANL libraries and the IAEA WLUP libraries. The first two problems are infinite cell benchmarks, provided by the IAEA WLUP for testing the library implementation on different computers. These problems are designated ZPR7-a2 and BNW-Pu-a1. [6] These cases are both light water moderated lattice problems, but with different fuel compositions. The ZPR7-a2 problem is for a low-enriched uranium oxide fuel (uranium enrichment of 3.04%), while BNW-Pu-a1 is weapons-grade MOX fuel. Both cases are hexagonal lattice cell problems.

Calculation results for the infinite cell problems are summarized in Table 9. The MCNP4C calculations listed on the table have been used as reference solutions. Generally, the results obtained using the DRAGON code with the WIMS-ANL library are closer to the reference values (less than 400 pcm $\Delta \rho$), compared to those obtained with the IAEA WLUP
library (about 600 – 110 pcm). The MCNP4C and DRAGON/WIMS-ANL calculations used ENDF/B-VI based data, while as noted earlier, most of the IAEA WLUP library data are derived from JEF2.2 data files. Given the simple nature of the problems and the relatively large differences between the reference and DRAGON/WLUP-WIMS solutions, the differences have been attributed to either (1) differences between the ENDF/B-VI and JEF2.2 data or (2) differences in the cross section generation path (e.g., different assumptions in the NJOY code input used in the library creation process).

Table 9. Comparison of Eigenvalues for Infinite Cell Benchmarks.

<table>
<thead>
<tr>
<th>Problem</th>
<th>Fuel</th>
<th>Library</th>
<th>Code</th>
<th>Eigenvalue*</th>
</tr>
</thead>
<tbody>
<tr>
<td>ZPR7-a2</td>
<td>UO₂</td>
<td>ENDF/B-VI</td>
<td>MCNP4C</td>
<td>1.26732 ± 0.0006</td>
</tr>
<tr>
<td></td>
<td></td>
<td>ANL</td>
<td>DRAGON</td>
<td>1.27108 (233)</td>
</tr>
<tr>
<td></td>
<td></td>
<td>IAEA</td>
<td>DRAGON</td>
<td>1.28499 (1085)</td>
</tr>
<tr>
<td>BNW-Pu-a1</td>
<td>MOX</td>
<td>ENDF/B-VI</td>
<td>MCNP</td>
<td>1.35787 ± 0.0006</td>
</tr>
<tr>
<td></td>
<td></td>
<td>ANL</td>
<td>DRAGON</td>
<td>1.35064 (-394)</td>
</tr>
<tr>
<td></td>
<td></td>
<td>IAEA</td>
<td>DRAGON</td>
<td>1.34677 (-607)</td>
</tr>
</tbody>
</table>

*Value in parenthesis is difference between DRAGON and MCNP4C results in pcm $\Delta\rho$.

The third problem is for a VHTR fuel element; the same fuel element used for the study reported in Section 2.2. Again, reference solutions have been generated using the MCNP4C code. Additionally, results obtained with the WIMS8 code have been included for comparison.

Results of the calculations are summarized in Table 10. Eigenvalues obtained using the 69- and 172-group libraries are included. The following trends in results are observed. The DRAGON calculations show that the results are dependent on the number of energy groups. The results obtained with the 69-group cross section libraries are generally worse than those with the 172-group libraries, when compared to the reference MCNP4C solution. The use of the 172-group libraries results in less than 700 pcm over-prediction of the eigenvalue, while the 69-group libraries give a much higher difference (an additional ~300 pcm higher than the 172-group library results).

It is also observed that the WIMS8 calculation using JEF2.2 data gave the smallest difference compared to the MCNP4C reference solution; it under-predicts the eigenvalue by 122 pcm. This low difference may have arisen from the cancellation of errors between the code.
solution effect and the library effect. Note again that the DRAGON calculations with the WLUP library (IAEA recommended library and the IAEA-JEF2.2 library) are predominantly based on JEF2.2 data. The use of these libraries result in the over-prediction of the eigenvalue and have higher differences relative to the MCNP4C results (an additionally 100 – 200 pcm) than those obtained with the WIMS-ANL library.

The magnitudes of some of these differences are not satisfactory and additional evaluations of the code and nuclear-data library are required.

Table 10. VHTR Assembly Eigenvalues with Different Libraries.

<table>
<thead>
<tr>
<th>Code</th>
<th>Library</th>
<th>Group</th>
<th>Eigenvalue</th>
</tr>
</thead>
<tbody>
<tr>
<td>MCNP4C</td>
<td>ENDF/B-VI</td>
<td>Continuous</td>
<td>1.53280 ± 0.00082</td>
</tr>
<tr>
<td>WIMS8</td>
<td>WIMS8 library</td>
<td>172</td>
<td>1.52993 (-122)</td>
</tr>
<tr>
<td>DRAGON</td>
<td>ANL WIMS-D library</td>
<td>172</td>
<td>1.54393 (470)</td>
</tr>
<tr>
<td></td>
<td>IAEA recommended library</td>
<td></td>
<td>1.54850 (661)</td>
</tr>
<tr>
<td></td>
<td>IAEA-JEF2.2 library</td>
<td></td>
<td>1.54814 (646)</td>
</tr>
<tr>
<td></td>
<td>IAEA-JENDL3 library</td>
<td></td>
<td>1.54689 (594)</td>
</tr>
<tr>
<td></td>
<td>ANL WIMS-D library</td>
<td>69</td>
<td>1.55198 (806)</td>
</tr>
<tr>
<td></td>
<td>IAEA recommended library</td>
<td></td>
<td>1.55639 (989)</td>
</tr>
<tr>
<td></td>
<td>IAEA-JEF2.2 library</td>
<td></td>
<td>1.55627 (984)</td>
</tr>
<tr>
<td></td>
<td>IAEA-JENDL3 library</td>
<td></td>
<td>1.55461 (915)</td>
</tr>
</tbody>
</table>

*Value in parenthesis is difference between DRAGON and MCNP4C results in pcm Δρ.
4. DRAGON CODE WORKSHOP

A workshop on the DRAGON code was organized in FY 2005 in order to familiarize some participants of the Gen IV and AFCI programs that have not used DRAGON with the code and to provide a refresher course for some of the other code users. This workshop was held at ANL, February 28 and March 1, 2005, and was attended by ~20 people from ANL, INL, ORNL, BNL, and KAERI (Korea). The format of the workshop was a technical exchange on the code. One of the code authors, Dr. Guy Marleau (Ecole Polytechnique de Montreal), gave a series of presentations. There were also opportunities for fruitful discussions between him and the participants, particularly concerning the application of the code for VHTR analysis.

The topics covered during the workshop included general code theory and features, cross section libraries to be used with the code, user instructions, and sample decks/problems. There were also discussions on the double heterogeneity and resonance treatments for coated fuel particles, transmutation and decay chains, modularity of code for staged calculations, and code performance compared to those of higher fidelity models/codes and experiments. Additional items included discussions on detailed code documentation and descriptions of advanced models and capabilities not in public domain (e.g., method of characteristics solution, homogenization and de-homogenization, parallel code version). A preliminary assessment of code deficiencies and necessary modifications that would make the code attractive for VHTR applications was also pursued.

There were positive feedbacks from the workshop participants. The code showed promise as one to be considered for use in the Gen IV program for the analysis of thermal systems. The code theories are also applicable to fast spectrum systems, but significant modifications would have to be made to the models used for resonance and slowing-down treatments. It is being planned to invite the DRAGON code developers to give another seminar in FY 2006 focusing more on the resonance and homogenizations treatments for the particulate fuels used in the advanced high temperature reactor systems.

The Workshop Agenda, list of participants, and the presentation material from the workshop are provided in Appendix A.
5. CONCLUSIONS AND RECOMMENDATIONS

A preliminary evaluation of the open-source lattice code DRAGON has been done to determine its performance in the analysis of the Very High Temperature Reactor (VHTR) fuel element designs. The activities performed under this project in FY 2005 were divided into three parts:

- DRAGON code was used to model the VHTR fuel elements of the helium-cooled and liquid-salt-cooled VHTR designs. Results were compared to those from another deterministic lattice code (WIMS8) and Monte Carlo code (MCNP). These results showed the promise of the DRAGON code. Some relatively large differences when the code results are compared to MCNP4C and WIMS8 results suggest that additional developmental work and validation might be necessary for bring the code to a level that it can be used routinely for VHTR analysis with great confidence.

- The nuclear data library currently used with the code was evaluated by comparing the code performance against those obtained when data libraries obtained from the IAEA WIMS-D4 Library Update Project (WLUP) are used. This step was also an evaluation of the WLUP libraries to investigate their potential for use in future DRAGON calculations. Differences were observed in the results obtained using the different libraries that could mostly be attributed to the data contained in the cross section libraries. Additional evaluations of these differences and their sources are recommended.

- A DRAGON workshop was held to discuss the code capabilities for modeling the VHTR. A code author visited ANL to deliver a two-day seminar on the DRAGON code to a group of reactor physics specialists from U.S. national laboratories. Future collaboration with the DRAGON code team is recommended to support the development of the code in order to provide features that make the code both accurate and more user-friendly for VHTR applications.
ACKNOWLEDGMENTS

This ANL work was supported by the U.S. Department of Energy, Office of Nuclear Energy, Science and Technology (VHTR element of the Generation IV Program) under Contract number W-31-109-Eng-38.

The contributions of Jessica Matheron, a 2005 summer intern in the Nuclear Engineering Division at ANL, are appreciated. She participated in the evaluation of the IAEA WIMS-D Library Update Project (WLUP) libraries.
REFERENCES


APPENDIX A

AGENDA, LIST OF PARTICIPANTS, AND PRESENTATION MATERIAL FOR THE DRAGON WORKSHOP

Argonne National Laboratory

February 28 and March 1, 2005
Agenda for DRAGON Workshop

Presentations by Dr. Guy Marleau, Ecole Polytechnique de Montreal
Monday February 28 – Tuesday March 1, 2005
Argonne National Laboratory
Argonne, Illinois 60439, U.S.A.

Monday:

9:00-12:00 1. Presentation of DRAGON 3.04.
1.1 General features.
1.2 Cross section libraries.
1.3 Geometry and tracking.
1.4 Resonance self-shielding.
1.5 Homogenization and condensation.
1.6 Depletion calculations.
14:00-17:00 2. Status of DRAGON.
2.1 Verification.
2.2 Validation.
2.3 Documentation.
2.4 Using DRAGON.

Tuesday:

9:00-12:00 3  DRAGON today and tomorrow.
3.1 Advanced features in DRAGON.
3.2 Current improvements in version 3.04 revision R.
3.3 Limitations in DRAGON.
3.4 What to expect in version 3.05.
3.5 Towards DRAGON 4.00.
13:00-14:00 4  Discussion.
4.1 Features required for Generation IV reactors analysis.
4.2 Additional user’s feedback.

End of Workshop

14:00-16:00 5  Meeting with ANL-NE Division Management.
# List of Participants for DRAGON Workshop

*Presentations by Dr. Guy Marleau, Ecole Polytechnique de Montreal*

Monday February 28 – Tuesday March 1, 2005  
Argonne National Laboratory  
Argonne, Illinois 60439, U.S.A.

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</tr>
</tbody>
</table>
Description of DRAGON 3.04

G. Marleau

Institut de génie nucléaire
École Polytechnique de Montréal
Contents:

1. General features.
2. Cross-section libraries.
3. Geometry and tracking.
4. Resonance self-shielding.
5. Homogenization and condensation.
6. Depletion calculations.
General features

DRAGON is a lattice code:

- It should be used to compute the neutron flux distribution inside assemblies of heterogeneous cells.
- It should use the computed neutron flux distribution to produce the homogenized and condensed nuclear properties that will be used in finite reactor calculations.
- It should solve the isotopic depletion equations and predict the concentration of each isotope in the cell as a function of time.
General features

CANDU example (reactor and cell):

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General features

CANDU example (reactivity devices):

- Symmetrized bottom z-plane
- Symmetrized central z-plane
- Top z-plane
- Central z-plane
- Bottom z-plane
- Homogenization cells
- Surface of symmetry
- Cooled cell
- Voided cell

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CANDU core in 2-D:

*Annular model – Color by region type*
- Void
- Reflectors
- Fuel

*Pin Model – Color by region type*
- Void
- Reflectors
- Fuel
- Tubes
- Moderator

Reflector limits in the Cartesian model

\( \beta = 0 \)
General features

CANDU reactor:

- 79 different physical regions in each fuel cell.
- 380 fuel cells in the 2–D reactor plane.
- 12 planes of fuel in the 3–D reactor.
- 3–D fuel bundles are not uniform along the axis (end-regions).

- Reactivity and shutdown devices placed at various location in the core.
- Over 700,000 different physical regions in the core.
Prismatic modular reactor example (reactor and assemblies):
General features

2-D assembly calculations using DRAGON:

- Assume assembly is uniform and infinite in the axial direction.
- Use reflective or periodic boundary conditions at assembly limits.
- Extract multigroup cell properties from microscopic cross section library.
- Self-shield the cross section for resonant isotopes.
General features

2–D assembly calculations using DRAGON (continue):

- Perform fine mesh multigroup transport calculation inside 2–D assembly.
- Perform isotopic depletion calculations.
- Generate burnup dependent fewgroup (multigroup) homogenized (or partly homogenized) assembly properties.
General features

3-D supercell (assembly) calculations using DRAGON:

- Assume reactor is made up of an infinite number of identical supercell.
- Use reflective or periodic boundary conditions at supercell limits.
- Extract multigroup, partly homogenized assembly properties from 2-D assembly calculation (EDI: database).
- Perform fine mesh multigroup transport calculation inside 3-D supercell.
- Generate fewgroup homogeneous properties for 3-D supercell.
General features

2-D or 3-D reactor calculations using DRAGON:

- Use burnup dependent fewgroups properties from assembly calculation (EDI: database).
- Use vacuum boundary conditions at reactor limits.
- Perform fine mesh multigroup transport calculation inside 2-D or 3-D reactor model.
- Generate fission power distribution in the reactor.
Using Version 3.04 of DRAGON, one can:

- Solve the neutron transport equation using the CP method in 1, 2 and 3–D geometry.
- Solve the neutron transport equation using the method of cyclic characteristics in 2–D geometry.
- Solve the isotopic depletion equations for isotopic concentration.
General features

One can also:

- Process multigroup microscopic cross section libraries including self-shielding.
- Process multigroup macroscopic cross section libraries generated by DRAGON or other codes (WIMS-AECL and HELIOS have been coupled to DRAGON for CANDU applications).
- Generate the fewgroup/homogenized information required for finite reactor calculations (DRAGON EDI: format and DONJON CPO: and FBM: formats are supported).
The multigroup static neutron transport equation:

\[
\left[ \vec{\Omega} \cdot \vec{\nabla} + \Sigma^g (\vec{r}) + D^g (\vec{r}) B^2 \right] \Phi^g (\vec{r}, \vec{\Omega}) = Q^g_s (\vec{r}, \vec{\Omega}) + \frac{1}{k} Q^g_f (\vec{r}) + Q^g_e (\vec{r})
\]

Scattering source \( Q^g_s (\vec{r}, \vec{\Omega}) \):

\[
Q^g_s (\vec{r}, \vec{\Omega}) = \sum_{h=1}^{G} \int d^2 \Omega' \Sigma^g_s (\vec{r}, \vec{\Omega} \rightarrow \vec{\Omega}') \Phi^h (\vec{r}, \vec{\Omega}')
\]

Fission source \( Q^g_f (\vec{r}) \):

\[
Q^g_f (\vec{r}) = \chi^g \sum_{h=1}^{G} \nu \Sigma^h_f (\vec{r}) \int d^2 \Omega' \Phi^h (\vec{r}, \vec{\Omega}')
\]
The multigroup macroscopic cross sections are defined as:

\[
\Sigma_{m,x}^g (\vec{r}, \vec{\Omega}) = \sum_I N_{m,I}(\vec{r}) \sigma_{I,x}^g (\vec{\Omega})
\]

for mixture \( m \) that contains the isotopes \( I \).

The multigroup microscopic cross sections are extracted from a cross-section library (many formats supported).

The multigroup microscopic cross sections can be self-shielded.
Transport equation in the absence of external sources \( Q_e^g(\vec{r}) = 0 \) is an eigenvalue problem:

1. \( k \) eigenvalue with imposed leakage \( (D^g(\vec{r})B^2 \text{ fixed}) \):
   \( k \) indicates how the fission rate should be modified to make the system critical (reach a non-trivial solution to the transport equation).

2. Buckling eigenvalue with imposed \( k \):
   \( D^g(\vec{r})B^2 \) represent the amount of leakage required to make the system critical.
Boundary conditions are required to close the system:

- **Albedo conditions:**

  \[
  \Phi_-(\vec{r}_S, \vec{\Omega} - 2(\vec{N}_S \cdot \vec{\Omega})) = \beta(\vec{r}_S, \vec{\Omega}) \Phi_+(\vec{r}_S, \vec{\Omega})
  \]

- \(\beta(\vec{r}_S, \vec{\Omega}) = 0\) for vacuum (void) BC.
- \(\beta(\vec{r}_S, \vec{\Omega}) = 1\) for reflective BC.
The static multigroup neutron transport equation (integral form) for the angular flux:

\[
\phi^g(\vec{r}, \Omega) = e^{-\tau^g(R_S)} \Phi^g(\vec{r}', \Omega) \Theta(\vec{r}, \vec{r}', \Omega) + \int_0^R e^{-\tau^g(R')} Q^g(\vec{r}', \Omega) \Theta(\vec{r}, \vec{r}', \Omega) dR'
\]

with \( R = |\vec{r} - \vec{r}'| \) and:

\[
\Theta(\vec{r}, \vec{r}', \Omega) = \begin{cases} 
1 & \text{for } \vec{r}' = \vec{r} - R\hat{\Omega} \\
0 & \text{otherwise}
\end{cases}
\]
Example of a neutron path for the integral transport equation:
General features

- Optical path $\tau^g(R)$:

$$\tau^g(R) = \int_{0}^{R} \Sigma^g(\vec{r} - R'\vec{\Omega})dR'$$

- Neutron source $Q^g(\vec{r}, \vec{\Omega})$ includes scattering, fission and fixed sources:

$$Q^g(\vec{r}, \vec{\Omega}) = Q^g_s(\vec{r}, \vec{\Omega}) + \frac{1}{k} Q^g_f(\vec{r}) + Q^g_e(\vec{r})$$

- $\Phi^g_{-}(\vec{r}'_S, \vec{\Omega})$ is incoming angular flux at surface $\vec{r}'_S$. 
Boundary conditions are required to close the system:

- For cell or assembly isolated in space one uses vacuum BC:

  \[ \Phi^g_\Omega (\vec{r}_S, \vec{\Omega}) = 0 \]
For cell or assembly repeating to infinity albedo conditions or periodic BC are considered:

\[
\Phi_-(\vec{r}_S, \vec{\Omega} - 2(\vec{N}_S \cdot \vec{\Omega})) = \beta(\vec{r}_S, \vec{\Omega})\Phi_+(\vec{r}_S, \vec{\Omega})
\]

\[
\Phi_-(\vec{r}'_S, \vec{\Omega}) = \Phi_+(\vec{r}_S, \vec{\Omega})
\]

with \(\Phi^g_+(\vec{r}'_S, \vec{\Omega})\) the outgoing angular flux at surface \(\vec{r}_S\).
General features

Approaches used in DRAGON to solve the multigroup neutron transport equation:

- Collision probability method.
- Method of characteristics.
General features

Approximations for CP and MOC:

- Divide domain into \( N_V \) regions of volume \( V_i \) with constant and isotropic sources and cross sections:

\[
\Sigma^g(\vec{r}) = \Sigma^g_j \quad \text{for} \quad \vec{r} \in V_j
\]

\[
Q^g(\vec{r}, \Omega) = \frac{1}{4\pi} q^g(\vec{r}) = q^g_j \quad \text{for} \quad \vec{r} \in V_j
\]

- Divide the external boundary \( S \) into \( N_S \) surfaces of area \( S_\alpha \) with constant angular flux:

\[
\Phi^g_-(\vec{r}_S, \Omega) = \frac{1}{4\pi} \phi^g_\alpha,-
\]
For CP, integrate transport equation over regions and surfaces:

\[
\phi_i^g = \frac{1}{V_i} \int_{V_i} d^3r \int d^2\Omega \phi^g(\vec{r},\vec{\Omega})
\]

\[
\phi_{+,\alpha}^g = \frac{1}{S_\alpha} \int_{S_\alpha} d^2r \int d^2\Omega (\vec{\Omega} \cdot \vec{N}_-) \phi^g(\vec{r}_S,\vec{\Omega})
\]

The CP transport equations are:

\[
\phi_i^g = \sum_{\alpha=1}^{N_S} p_{i\alpha}^g \phi_{-,\alpha}^g + \sum_{j=1}^{N_V} p_{ij}^g q_j^g
\]

\[
\phi_{+,\alpha}^g = \sum_{\beta=1}^{N_S} p_{\alpha\beta}^g \phi_{-,\beta}^g + \sum_{j=1}^{N_V} p_{\alpha j}^g q_j^g
\]

ANL DRAGON workshop
Four types of probabilities:

\[ \tilde{p}_{i,j}^g = V_i p_{i,j}^g = \int_{V_i} \int_{V_j} \frac{e^{-\tau^g(R)}}{4\pi R^2} \Theta_i \Theta_j d^3r' d^3r \]

\[ \tilde{p}_{i,\alpha}^g = V_i p_{i,\alpha}^g = \int_{V_i} \int_{S_\alpha} \frac{e^{-\tau^g(R_S)}}{4\pi R_S^2} (\vec{\Omega} \cdot \vec{N}_-) \Theta_i \Theta_\alpha d^3r' d^2r \]

\[ \tilde{p}_{\alpha,i}^g = \frac{S_\alpha}{4} p_{\alpha,i}^g = \int_{S_\alpha} \int_{V_i} \frac{e^{-\tau^g(R)}}{4\pi R^2} (\vec{\Omega} \cdot \vec{N}_+) \Theta_\alpha \Theta_i d^2r' d^3r \]

\[ \tilde{p}_{\alpha,\beta}^g = \frac{S_\alpha}{4} p_{\alpha,\beta}^g = \int_{S_\alpha} \int_{S_\beta} \frac{e^{-\tau^g(R_S)}}{4\pi R_S^2} (\vec{\Omega} \cdot \vec{N}_-)(\vec{\Omega} \cdot \vec{N}_+) \times \Theta_\alpha \Theta_\beta d^2r d^2r' \]

These will be evaluated numerically.
For the method of characteristics we will use again the integral transport equation and assume:

- Divide domain into $N_V$ regions of volume $V_i$ with constant and isotropic sources and cross sections:

$$\Sigma^g(\vec{r}) = \Sigma_j^g \quad \text{for } \vec{r} \in V_j$$

$$Q^g(\vec{r}, \Omega) = \frac{1}{4\pi} q^g(\vec{r}) = q_j^g \quad \text{for } \vec{r} \in V_j$$

- Divide the external boundary $S$ into $N_S$ surfaces of area $S_\alpha$ with constant angular flux:

$$\Phi^g_-(\vec{r}_S, \Omega) = \frac{1}{4\pi} \phi^g_{\alpha,-}$$
General features

- The sources $q_i(\vec{\rho}_k)$ are constant on each line segment $\vec{r}' = \vec{r}' + t\vec{\Omega}$ crossing the domain in direction $\vec{\Omega}$.

- The incoming flux on a point $\vec{r}_S$ on surface $\alpha$ is assumed known and given by $\psi_{\alpha,-}(\vec{\rho}_k, \vec{\Omega})$. 
The average flux inside each region will be given by:

\[ V_i \phi_i = \int \frac{d^2 \Omega}{4\pi} \int d^2 \rho \sum_{j=1}^{i-1} q_j(\vec{\rho}) \left( \frac{1 - e^{-\tau_j(\vec{\rho})}}{\Sigma_j} \right) \]

\[ \times \left[ e^{-\sum_{i=j+1}^{i-1} \tau_i(\vec{\rho})} \right] \left( \frac{1 - e^{-\tau_i(\vec{\rho})}}{\Sigma_i} \right) \]

\[ + \int \frac{d^2 \Omega}{4\pi} \int d^2 \rho q_i(\vec{\rho}) \left( \frac{\tau_i(\vec{\rho}) - 1 + e^{-\tau_i(\vec{\rho})}}{\Sigma_i^2} \right) \]

The \( \frac{d^2 \Omega}{4\pi} \) and \( \int d^2 \rho \) for MOC are evaluated numerically using the same procedure as for CP method.
Isotopic depletion equations:

\[
\frac{dN_I(t)}{dt} = \sum_J S_I^J(t) - N_I(t) \left( \int dE \sigma_{I,a}(E) \phi(E,t) + \lambda_I \right)
\]

- $\lambda_I$ is the decay rate of $I$.
- $\sigma_{I,a}(E)$ is the neutron absorption cross section for $I$.
- $S_I(t)$ is the production rate for isotope $I$ from isotope $J$.
- Initial concentrations $N_I(t_0)$ required.
Example of simplified decay chains for $^{235}\text{U}$ and $^{238}\text{U}$:

- $^{235}\text{U}$ to $^{236}\text{U}$ via $(n,\gamma)$ reaction with 86.67 barns
  - $^{236}\text{U}$ undergoes fission with 506.8 barns
- $^{238}\text{U}$ to $^{239}\text{U}$ via $(n,\gamma)$ reaction with 2.241 barns
  - $^{239}\text{U}$ decays to $^{239}\text{Np}$ via $\beta^-$ with (23.45 minutes) and then to $^{239}\text{Pu}$ via $\beta^-$ with (2.3565 hours)
    - $^{239}\text{Np}$ undergoes $(n,\gamma)$ reaction with 32.79 barns and then decays to $^{239}\text{Pu}$ via $(n,\gamma)$ with 274.2 barns
  - $^{239}\text{Pu}$ undergoes fission with 698.3 barns
General features

- Decay chain:
  - Available on some libraries.
  - Can be provided manually.
- Flux normalization options available:
  - Constant flux burnup.
  - Constant power burnup.
  - Out of core depletion.
- Includes an isotope saturation model.
Cross-section libraries

Two types of multigroup cross-section database can be accessed by DRAGON:

MACROLIB:

- Contains the multigroup macroscopic cross-section associated with each mixture.
- Generally created by MAC: or EDI: modules.
- Required when macroscopic cross section are used (CP calculations module, flux solution module, editing module, etc.).
Cross-section libraries

MICROLIB:

- Contains the multigroup microscopic cross-section associated with each isotope at a fixed temperature and dilution (self-shielded cross sections).
- Contains the depletion chain.
- Includes a MACROLIB.
- Generally created by LIB: or EDI: modules.
- Required when microscopic cross section are used (depletion calculations, self-shielding calculation, etc).
Cross-section libraries

Minimum cross-section requirements for each mixture \( m \) in a MACROLIB:

- The multigroup total cross section \( \Sigma^g_m \).
- The isotropic component of the multigroup scattering cross section \( \Sigma^{h\rightarrow g}_{m,s,0} \) defined as:

\[
\Sigma^{h\rightarrow g}_{m,s,0} = \int_{4\pi} \, d^2\Omega^2 \Sigma^{h\rightarrow g}_{m,s} (\vec{\Omega}' \rightarrow \vec{\Omega}) P_0(\vec{\Omega}' \cdot \vec{\Omega})
\]
Cross-section libraries

For multiplicative problems one also needs:

- The product of the average neutron emitted per fission with the multigroup fission cross section $\nu \Sigma^{g}_{m,f}$.

- The multigroup fission spectrum $\chi^{g}_{m}$.

For a fixed source problem (multiplicative or not):

- The external neutron source $q^{g}_{m}$. 
Cross-section libraries

Additional cross sections:

- The linearly isotropic component of the multigroup scattering cross section $\Sigma_{m,s,1}^{h\rightarrow g}$.

$$\Sigma_{m,s,1}^{h\rightarrow g} = \int \frac{d^2\Omega'^2}{4\pi} \Sigma_{m,s}^{h\rightarrow g}(\vec{\Omega}' \rightarrow \vec{\Omega}) P_1(\vec{\Omega}' \cdot \vec{\Omega})$$

- Required only if $B_1$ leakage method is used.
The transport correction $\Sigma_{m,tc}^g$.

The transport calculations can be performed using transport corrected total ($\tilde{\Sigma}_m^g$) and scattering ($\tilde{\Sigma}_{m,s,0}^{h\rightarrow g}$) cross sections:

$$\tilde{\Sigma}_m^g = \Sigma_m^g - \Sigma_{m,tc}^g$$

$$\tilde{\Sigma}_{m,s,0}^{h\rightarrow g} = \Sigma_{m,s,0}^{h\rightarrow g} - \delta^h \Sigma_{m,tc}^g$$

Takes partially into account the linearly anisotropic scattering contributions.
Cross-section libraries

For contribution of multi-neutron production reactions such as:

\[ X_A^Z + n_1^0 \rightarrow X_{A-1}^Z + 2n_1^0 + \gamma \]

The scattering cross section (microscopic and macroscopic) are corrected to take into account this effect:

\[ \tilde{\sigma}^{h \rightarrow g}_{I,s,0} = \sigma^{h \rightarrow g}_{I,s,0} + 2\delta^{gh}_{I,(n,2n)} \sigma^{g}_{I,(n,2n)} \]

\[ \tilde{\Sigma}^{h \rightarrow g}_{m,s,0} = \Sigma^{h \rightarrow g}_{m,s,0} + 2\delta^{gh}_{m,(n,2n)} \Sigma^{g}_{m,(n,2n)} \]
MACROLIB can be created:
- From the input file using the MAC: module.
- From a GOXS file using the MAC: module.
- From another MACROLIB using the MAC: module.
- From a basic microscopic cross section library using the LIB: module.
- From a MICROLIB using the LIB: module.
- From MICROLIB or a MACROLIB using the EDI: module.
Cross-section libraries

MACROLIB can be updated using:

- The \texttt{MAC}: module for direct cross section perturbation.
- The \texttt{LIB}: module for perturbation of the isotopic contents of a mixture.
- The \texttt{EVO}: module when the isotopic composition of the mixture has changed.
Basic microscopic libraries that can be accessed in DRAGON:

- WIMS–AECL format.
- MATXS format.
- WIMD-D4 format.
- DRAGLIB format.
WIMS-AECL microscopic cross section libraries:
- Direct access format only.
- Automatic depletion chains.
- Winfrith 69 group library.
- ENDF/B-V based 89 group library.
- ENDF/B-VI based 89 group library.
Cross-section libraries

WIMSD4 microscopic cross section libraries:
- Sequential binary format.
- Automatic depletion chains.
- WIMSKAL 69 group library.
- ARGONNE ENDF/B-VI based 69 and 172 group libraries.
- IAEA WLUP Libraries 69 and 172 group libraries.

http://www-nds.iaea.org/wimsd/downloads.htm
Cross-section libraries

MATXS microscopic cross section libraries:
- Sequential binary format.
- No depletion chains.
- ENDF/B-V based 69 group library in NJOY-89 MATXS format.
- ENDF/B-V based 70 group library in NJOY-91 MATXS format.
DRAGLIB microscopic cross section libraries:

- Direct access format (XSM file).
- Automatic depletion chains.
- Winfrith 69 group library (a version of the library generated from the WIMS-AECL 69 groups library).
Cross-section libraries

Contents of the MICROLIB:

- An isotope on the MICROLIB refers to an isotope on the basic microscopic cross section library at a specified temperature and dilution.

- For each isotope, the LIB: module extracts the cross section from the basic library, interpolates in temperature and dilution and stores the information on the MICROLIB.

- The LIB: module initially assumes that the isotopes are at infinite dilution.

- The SHI: can be used to update the dilution cross section for resonance self-shielding based on the improved Stamm’ler method.
Depletion chain

- The MICROLIB can also contain a depletion chain.
- This chain can be extracted from the basic cross section library (WIMS-AECL, WIMS-D4 and DRAGLIB).
- The chain can also be defined manually in the LIB module.
- When the chain is defined manually, it must be consistent with the isotopic cross section information available on the MICROLIB.
Cross-section libraries

MICROLIB can be created:

- From a basic microscopic cross section library using the LIB: module.
- From another MICROLIB using the LIB: module.
- From another MICROLIB using the EDI: module.
MICROLIB can be updated using:

- **The LIB**: module for manual perturbation of the isotopic contents of a mixture, dilution or temperature of an isotope.
- **The EVO**: module when the isotopic composition of the mixture has changed.
- **The SHI**: for self-shielding calculation.
Approximations for CP methods:

- Divide domain into $N_V$ regions of volume $V_i$ with constant and isotropic sources and cross sections:

\[
\Sigma^g(\vec{r}) = \Sigma^g_j \quad \text{for} \quad \vec{r} \in V_j
\]

\[
Q^g(\vec{r}, \Omega) = \frac{1}{4\pi} q^g(\vec{r}) = q^g_j \quad \text{for} \quad \vec{r} \in V_j
\]

- Divide the external boundary $S$ into $N_S$ surfaces of area $S_\alpha$ with constant angular flux:

\[
\Phi^g_-(\vec{r}_S, \Omega) = \frac{1}{4\pi} \phi^g_{\alpha,-}
\]
Integrate transport equation over regions and surfaces:

\[
\phi_{i}^{g} = \frac{1}{V_{i}} \int_{V_{i}} d^{3}r \phi^{g}(\vec{r})
\]

\[
\phi_{+,\alpha}^{g} = \frac{1}{S_{\alpha}} \int_{S_{\alpha}} d^{2}r \phi^{g}(\vec{r}_{S})
\]

The CP transport equations are:

\[
\phi_{i}^{g} = \sum_{\alpha=1}^{N_{S}} p_{i\alpha}^{g} \phi_{-,\alpha}^{g} + \sum_{j=1}^{N_{V}} p_{ij}^{g} q_{j}^{g}
\]

\[
\phi_{+,\alpha}^{g} = \sum_{\beta=1}^{N_{S}} p_{\alpha\beta}^{g} \phi_{-,\beta}^{g} + \sum_{j=1}^{N_{V}} p_{\alpha j}^{g} q_{j}^{g}
\]
Four types of probabilities:

\[ \tilde{p}^g_{ij} = V_i \tilde{p}^g_{ij} = \int_{V_i} \int_{V_j} \frac{e^{-\tau^g(R)}}{4\pi R^2} \Theta_i \Theta_j d^3r' d^3r \]

\[ \tilde{p}^g_{i\alpha} = V_i \tilde{p}^g_{i\alpha} = \int_{V_i} \int_{S_{\alpha}} \frac{e^{-\tau^g(R_S)}}{4\pi R^2_S} (\vec{\Omega} \cdot \vec{N}_-) \Theta_i \Theta_{\alpha} d^3r' d^2r \]

\[ \tilde{p}^g_{\alpha i} = \frac{S_\alpha}{4} \tilde{p}^g_{\alpha i} = \int_{S_{\alpha}} \int_{V_i} \frac{e^{-\tau^g(R)}}{4\pi R^2} (\vec{\Omega} \cdot \vec{N}_+) \Theta_{\alpha} \Theta_i d^2r' d^3r \]

\[ \tilde{p}^g_{\alpha\beta} = \frac{S_\alpha}{4} \tilde{p}^g_{\alpha\beta} = \int_{S_{\alpha}} \int_{S_{\beta}} \frac{e^{-\tau^g(R_S)}}{4\pi R^2_S} (\vec{\Omega} \cdot \vec{N}_-) (\vec{\Omega} \cdot \vec{N}_+) \times \Theta_{\alpha} \Theta_{\beta} d^2r d^2r' \]
Geometry and tracking

Symmetry relations:

\[ V_i p^g_{ij} = V_j p^g_{ji} \]
\[ 4 V_i p^g_{i\alpha} = S_\alpha p^g_{\alpha i} \]
\[ S_\alpha p^g_{\alpha \beta} = S_\beta p^g_{\beta \alpha} \]

Conservation properties:

\[ \sum_{\alpha=1}^{N_\alpha} p^g_{i\alpha} + \sum_{j=1}^{N_j} p^g_{ij} \Sigma^g_j = 1 \]
\[ \sum_{\beta=1}^{N_\beta} p^g_{\alpha \beta} + \sum_{i=1}^{N_i} p^g_{\alpha i} \Sigma^g_i = 1 \]
Geometry and tracking

Surface flux approximation:
- Incoming angular flux on outer surfaces assumed to be independent of $\vec{\Omega}$.

Comments
- Outgoing angular flux on outer surfaces integrated over $\vec{\Omega}$.
- Angular flux not used at region interfaces.
- Approximation for incoming angular flux exact for vacuum BC.
- Approximation for incoming angular flux leads to large errors in surface flux when flux is not isotropic.
Geometry and tracking

Illustration of approximate BC:

Reflection

Translation

isotropic

specular

isotropic
Geometry and tracking

Recommendations to reduce errors due to approximate use of BC for a fixed direction:

- No special treatment for 2 vacuum BC.
- Unfold cell once for 1 vacuum and 1 reflection BC.
- Multiply unfold cell for 2 reflection or periodic cell, apply approximate BC on final surfaces and consider results in cell located far from these surfaces.
Geometry and tracking

Example of cell unfolding in direction $X$:

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</tbody>
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Geometry and tracking

Simplifying CP equations using approximate BC:

- Assume:

\[ \vec{J}_g = A^g \vec{J}_g \]

- Final transport equation:

\[ \vec{\phi}^g = P_{c,vv}^g \vec{Q}^g \]

with the complete collision probability matrix \( P_{c,vv}^g \):

\[ P_{c,vv}^g = \left( P_{vv}^g + P_{vs}^g \left( (A^g)^{-1} - P_{ss}^g \right)^{-1} P_{sv}^g \right) \]
Numerical CP evaluation:

\[ \tilde{p}_{ij}^g = \int dx' \int dy' \int dR' \int_{4\pi} d^2\Omega \int_{R_i - \frac{1}{2}}^{R_i + \frac{1}{2}} dR 
\times e^{-\tau^g(R)} \Theta_i \Theta_j \]
General 3–D geometry for collision probability integration:
Projection of 3–D geometry on a 2–D plane:
After integration over $R'$ and $R$, one obtains:

$$
\tilde{p}^g_{ij} = \frac{1}{4\pi \Sigma_i \Sigma_j} \int_0^{4\pi} d^2 \Omega \int dx' \int dy' \\
\times \left[ 1 - \exp \left( -\tau^g_{i-\frac{1}{2},i+\frac{1}{2}} \right) \right] \exp \left( -\tau^g_{i+\frac{1}{2},j-\frac{1}{2}} \right) \\
\times \left[ 1 - \exp \left( -\tau^g_{j-\frac{1}{2},j+\frac{1}{2}} \right) \right] \Theta_i \Theta_j
$$

with:

$$
\tau^g_{i\pm\frac{1}{2},j\pm\frac{1}{2}} = \Sigma_i \left( R_{i+\frac{1}{2}} - R_{i\pm\frac{1}{2}} \right) + \tau_{i+\frac{1}{2},j-\frac{1}{2}} \\
+ \Sigma_j \left( R_{j+\frac{1}{2}} - R_{j-\frac{1}{2}} \right)
$$
For 3–D geometry, the final integrations are evaluated numerically:

\[
\tilde{p}_{ij}^g = \frac{1}{4\pi \sum_i \sum_j} \sum_{\alpha=1}^{N_\Omega} W_\Omega \sum_{m=1}^{N_\rho} \sum_{n=1}^{N_\rho} \delta^2 \\
\times \left[ 1 - \exp\left(-\tau_{i-\frac{1}{2},i+\frac{1}{2}}^g\right) \right] \exp\left(-\tau_{i+\frac{1}{2},j-\frac{1}{2}}^g\right) \\
\times \left[ 1 - \exp\left(-\tau_{j-\frac{1}{2},j+\frac{1}{2}}^g\right) \right] \Theta_i \Theta_j
\]

and a 4–D numerical quadrature remains.

For 2–D geometry the integration over the polar angle leads to Bickley-Nayler functions and a 2–D numerical quadrature remains.
For the method of characteristics the flux is evaluated using:

\[
V_i \phi_i = \sum_{\alpha=1}^{N_\Omega} W_\Omega \sum_{m=1}^{N_\rho} \sum_{n=1}^{N_\rho} \delta^2 \sum_{j=1}^{i-1} q_j \left( \frac{1 - e^{-\tau_j(\bar{\rho}_{\alpha,m,n})}}{\Sigma_j} \right)
\]

\[
\left[ e^{-\sum_{l=j+1}^{i-1} \tau_l(\bar{\rho}_{\alpha,m,n})} \right] \left( \frac{1 - e^{-\tau_i(\bar{\rho}_{\alpha,m,n})}}{\Sigma_i} \right)
\]

\[
+ \sum_{\alpha=1}^{N_\Omega} W_\Omega \sum_{m=1}^{N_\rho} \sum_{n=1}^{N_\rho} \delta^2 q_i \left( \frac{\tau_i(\bar{\rho}_{\alpha,m,n}) - 1 + e^{-\tau_i(\bar{\rho}_{\alpha,m,n})}}{\Sigma_i^2} \right)
\]

The numerical quadrature used is identical to that considered for the CP method.
Geometry and tracking

Quadrature for Cartesian 3–D geometry:
Angular quadrature for 3-D Cartesian geometry:

- $S_n$ type $\text{EQ}_N$ quadrature with $4N_\Omega(N_\Omega + 2)/8$ angular directions $\vec{\Omega}_1,i$, $\vec{\Omega}_2,i$, $\vec{\Omega}_3,i$ and $\vec{\Omega}_4,i$.

- Global quadrature weight $W_\Omega = 2/(N_\Omega(N_\Omega + 2))$.

- Tracking in upper half sphere only.
Comments on Cartesian surface quadrature in 3–D:

- 3–D different planes selected for each direction.
Geometry and tracking

Spatial quadrature for Cartesian 3–D geometry:

1. Locate the \((x'_c, y'_c)\) center of the geometry.
2. Compute radius \(h_+\) of smallest sphere containing the geometry.
3. From track density \(\rho\) compute:

\[
N_\rho = (2\sqrt{\rho h_+}) + 1 \quad \delta = \frac{2h_+}{N_\rho}
\]

4. Integration line \(l_{mn}\) passes through:

\[
x'_m = x'_c + \left(\frac{2m - 1}{2}\right) \delta \quad y'_n = y'_c - h_+ + \left(\frac{2n - 1}{2}\right) \delta
\]
Angular quadrature for 2-D Cartesian geometry:

- Order $N$ trapezoidal quadrature for azimuth $\varphi$.
- Tracking for $0 \leq \varphi \leq \pi$.
- For standard tracking, use Bickley-Nayler functions for polar ($\theta$) quadrature.
- For cyclic tracking use order $\frac{(N + 1)}{2}$ Gauss-Legendre quadrature for polar ($\theta$) angle.
CP quadrature for Cartesian 2–D geometry:
Geometry and tracking

Angular quadrature for 3-D Hexagonal geometry:

- Order $N$ trapezoidal quadrature for azimuth $\phi$.
- Order $(N + 1)/2$ Gauss-Legendre quadrature for polar angle.
- Tracking in upper half sphere only.
Example of valid hexagonal 3–D geometry:
Angular quadrature for hexagonal 3–D geometry:
Geometry and tracking

Angular quadrature for 2-D Hexagonal geometry:

- Order $N$ trapezoidal quadrature for azimuth.
- Tracking for $0 \leq \varphi \leq \pi$.
- Use Bickley-Nayler functions for polar quadrature.
Angular quadrature for hexagonal 2–D geometry:
Comments on angular quadrature:

- Tracking in the lower half sphere (3–D) or circle (2–D) is not considered because of symmetry of CP matrices.

- Try to select as many angles as possible (neutrons travel on a straight line).

- Try to select as high a density as possible (to maximize the number of lines crossing each region).

- For 3–D Cartesian cases, the angular quadrature is currently limited to $N_\Omega = 16$ (144 angles).
Geometry and tracking

DRAGON tracking example:
DRAGON tracking of a line $l_{m,n}$:

1. Follow the tracking line as it travels through the cell.
2. Identify each region $i$ and surface $s$ uniquely.
3. Compute distance $\tilde{l}_{i,m,n}$ the neutron travels in each region.
4. Associate region flux number $I$ with set of spatial regions $i$.
5. Associate surface flux number $S$ with sets of surfaces $s$.
Post treatment of tracking file:

- For each direction, normalize tracks using:

\[ l_{i,m,n} = \left( \frac{V_i}{\tilde{V}_i} \right) \tilde{l}_{i,m,n} \]

where:

\[ \tilde{V}_i = \frac{W_p}{3} \sum_{m=1}^{N_p} \sum_{n=1}^{N_p} \tilde{l}_{i,m,n} \]
Geometry and tracking

Comments on storage requirements for tracking:

- Maximum number of line segments tracks $d_t$:
  
  $$d_t < 6N(N_\Omega(N_\Omega + 2))\rho_p h_+^2$$

- For $N = 1000$, $h_+ = 50$, $\rho_p = 20$ t/cm$^2$ and $N_\Omega = 8$:
  
  $$d_t < 20 \times 10^9$$

  tracks segments.
Multigroup cross section versus continuous energy:

Total cross section as a fonction of energy for $^{235}$U

- Pointwise
- Multigroup
Self-Shielding

Multigroup cross section versus continuous energy:

Total cross section as a function of energy for $^{235}$U

ANL DRAGON workshop

Description of DRAGON 3.04 – 84/110
Self-Shielding

Multigroup condensation procedure is selected to preserve the reaction rates (resonance integrals):

$$R_{I,x}^g(\sigma_0) = \int_{E_g}^{E_{g-1}} \sigma_{I,x}(E) \phi(E, \sigma_0)$$

Here, an approximated flux distribution $\phi(E, \sigma_0)$ is used corresponding to:

- An Infinite cell calculation for a binary mixture of hydrogen and isotope $I$.

The resonance integrals are tabulated as a function of the dilution $\sigma_0$ cross section and stored in the microscopic cross section library.
Self-Shielding

Multigroup cross sections are computed using:

\[ \sigma_{I,x}^g(\sigma_e) = \frac{R_{I,x}^g(\sigma_e)}{\phi^g(\sigma_e)} \]

Here:

- The effective dilution \( \sigma_e \) is approximated using the self-shielding process.
- \( R_{I,x}^g(\sigma_e) \) are extracted from the library and interpolated in dilution.
- \( \phi^g(\sigma_e) \) can be extracted from library of computed approximately.
Resonance self-shielding:

- Required to evaluate the equivalent dilution cross section $\sigma_e$ at which $R_{I,x}^g$ will be evaluated.

- Dilution cross section will be dependent on space, and mixture considered.
  
  - Cell is finite and heterogeneous.
  - Mixture not necessarily a binary mixture of hydrogen and isotope $I$.

- Use equivalence principle:
  
  - Two systems with the same value of $\sigma_e$ will have the same self-shielded cross section.
Self-Shielding

Improved Stamm’ler method:

1. Compute fuel-to-fuel collision probability \( P_{\text{ff}}(\sigma_e) \) in exact geometry for arbitrary dilutions \( \sigma_e \).

2. Approximate \( P_{\text{ff}}(\sigma_e) \) by a rational series in \( \sigma_e \) and invert series \([\sigma_e(P_{\text{ff}})]\).

3. Start iteration by using \( \tilde{\sigma}_e \) an initial dilution.

4. Evaluate self-shielded cross sections \( \sigma_{g,I,x}^{\tilde{\sigma}_e} \) and compute exact \( P_{\text{ff}} \).

5. Evaluate new approximation \( \tilde{\sigma}_e \) using rational series.

6. Check for convergence and exit if convergence reached otherwise return to 4.
Self-Shielding

Self-shielded flux evaluation:

- Interpolate in table if $\phi^g(\sigma_0)$ is tabulated as a function of dilution.

- Approximate using the conservation relation in an homogeneous media:

$$\phi^g(\sigma_e) = 1 - \frac{1}{N\sigma_e} \left[ R^g_I(\sigma_e) - \sum_h \frac{U^h}{U^g} R^h_{I,s}^g(\sigma_e) \right]$$
Approximate using the conservation relation in an heterogeneous media:

\[
\phi^g(\sigma_e) = 1 - \frac{1}{E(\sigma_e)} \left[ R^g_I(\sigma_e) - \sum_h \frac{U^h}{U^g} R^{h \rightarrow g}_{I,s}(\sigma_e) \right]
\]

where \( E(\sigma_e) \) is the escape matrix defined as:

\[
E(\sigma_e) = \frac{1}{P_{ff}(\sigma_e)} - N\sigma^g_{I,x}(\sigma_e)
\]

This last approximation is known as the Livolant-Jeanpierre normalization.
The DRAGON editing module **EDI**: performs the following tasks:

- Condense the multigroup microscopic and macroscopic cross sections.
- Homogenize the multigroup microscopic and macroscopic cross sections.
- Prints the reaction rates and the homogenize/condense cross sections.
- Save this information for further used in the form of a MACROLIB or a MICROLIB.
Condensation and homogenization based on the following assumptions:

- Reaction rates are physically meaningful and should be preserved by the condensation/homogenization procedure.

\[ R_i = \sum_g V_i \phi_i^g \Sigma_i^g = V_i \phi_i \Sigma_i \]
\[ R^g = \sum_i V_i \phi_i^g \Sigma_i^g = V \phi^g \Sigma^g \]

- The eigenvalue is physically meaningful and should be preserved by the condensation/homogenization procedure.
Homogenization/Condensation

Condensation:

- Condensed transport equation (macrogroup $K$):

\[
\vec{\Omega} \cdot \vec{\nabla} \sum_{g \in G_K} \Phi^g(\vec{r}, \vec{\Omega}) + \sum_{g \in G_K} \Sigma^g(\vec{r}) \Phi^g(\vec{r}, \vec{\Omega}) = \\
\sum_{g \in G_K} \left[ Q^g_s(\vec{r}, \vec{\Omega}) + \frac{1}{k} Q^g_f(\vec{r}) \right]
\]

- Few group version of the same equation:

\[
\vec{\Omega} \cdot \vec{\nabla} \Phi^K(\vec{r}, \vec{\Omega}) + \Sigma^K(\vec{r}) \Phi^K(\vec{r}, \vec{\Omega}) = \left[ Q^K_s(\vec{r}, \vec{\Omega}) + \frac{1}{k} Q^K_f(\vec{r}, \vec{\Omega}) \right]
\]

It should reproduce condensed multigroup results.
The condensation procedure that satisfies our requirements:

\[ \phi^K_i = \sum_{g \in G_K} \phi^g_i \]

\[ \Sigma^K_i = \frac{1}{\phi^K_i} \sum_{g \in G_K} \sum^g_i \phi^g_i \]

\[ \Sigma^{L \rightarrow K}_{s,i} = \frac{1}{\phi^L_i} \sum_{h \in G_L} \sum_{g \in G_K} \sum^h \rightarrow g \phi^h_i \]

\[ \chi^K_i = \sum_{g \in G_K} \chi^g_i \]

\[ \nu \Sigma^K_{f,i} = \frac{1}{\phi^K_i} \sum_{g \in G_K} \nu \sum^h \rightarrow g \phi^g_i \]
Homogenization/Condensation

Full cell homogenization:

- Homogenized transport equation (no leakage):
  \[
  \sum_i \Sigma_i V_i \phi_i^g = \sum_i V_i [Q_{s,i}^g + \frac{1}{k} Q_{f,i}^g]
  \]

- Homogeneous transport equation (no leakage):
  \[
  \hat{\Sigma}^g V \hat{\phi}^g = V [\hat{Q}_{s,i}^g + \frac{1}{k} \hat{Q}_{f,i}^g]
  \]

- Homogenized and homogeneous transport equations are identical if one selects a flux-volume homogenization technique.
Flux-volume homogenization technique:

\[ \hat{\phi}^g = \frac{1}{V} \sum_i V_i \phi_i^g \]

\[ \hat{\Sigma}^g = \frac{1}{V \hat{\phi}^g} \sum_i V_i \Sigma_i^g \phi_i^g \]

\[ \hat{\Sigma}_{s}^{h \rightarrow g} = \frac{1}{V \hat{\phi}^h} \sum_i V_i \Sigma_{s,i}^{h \rightarrow g} \phi_i^h \]

\[ \hat{\nu} \hat{\Sigma}_f^g = \frac{1}{V \hat{\phi}^g} \sum_i V_i \nu \Sigma_{f,i}^g \phi_i^g \]

\[ \hat{\chi}^g = \frac{1}{V \sum_h \hat{\nu} \hat{\Sigma}_f^h \phi^g} \sum_i \chi_i^g V_i \sum_h \nu \Sigma_{f,i}^g \phi_i^h \]
Flux-volume homogenization fails if:

- The cell is finite (a cell with leakage):

\[ \sum_{j=1}^{N_j} p_{i,j}^{g} \sum_{j}^{g} \neq 1 \]

- Partial cell homogenization is considered.
Partial homogenization:

- The $N$ region transport equation homogenized over $M$ regions:

$$\sum_{i \in M_I} V_i \Sigma_i^g \phi_i^g = \sum_{i \in M_I} \sum_J \sum_{j \in M_J} p_{ji}^g (\Sigma^g) [Q_{s,i}^g + \frac{1}{k} Q_{f,i}^g]$$

- The $M$ region transport equation:

$$V_I \hat{\Sigma}_I^g \phi_I^g = \sum_J \hat{P}_{JI}^g (\hat{\Sigma}^g) [Q_{s,J}^g + \frac{1}{k} Q_{f,J}^g]$$

- $\hat{P}_{JI}^g (\hat{\Sigma}^g)$ are CP computed using homogenized cross sections.
We need:

\[ \sum_{i \in M_I} V_i \Sigma_i^g \phi_i^g = V_I \Sigma_I^g \phi_I^g \]

and:

\[ \sum_{J} \hat{P}_{JI}^g (\hat{\Sigma}^g)[Q_{s,J}^g + \frac{1}{k} Q_{f,J}^g] = \]

\[ \sum_{i \in M_I} \sum_{J} \sum_{j \in M_J} p_{ji}^g (\Sigma^g)[Q_{s,i}^g + \frac{1}{k} Q_{f,i}^g] \]

to be simultaneously true.
The flux-volume homogenization method is not longer adequate because:

- There is no simple relation between $\hat{p}_{ji}^g (\hat{\Sigma}_I^g)$ and $p_{ji}^g (\Sigma_I^g)$.

The alternative here is to use a non-linear process:

- Consider a flux-volume homogenization for $\phi_I^g$ and $\Sigma_I^g$.
- Redefine the homogeneous flux $\hat{\phi}_I^g$ and cross sections $\hat{\Sigma}_I^g$ as follows:

$$\hat{\phi}_I^g = \frac{1}{\mu_I^g} \phi_I^g \quad \hat{\Sigma}_I^g = \mu_I^g \Sigma_I^g$$
Determine the SPH factors $\mu^g_I$ such a way that:

$$\sum_J \hat{p}_{JI}^g (\hat{\Sigma}^g) \left[ Q_{s,J}^g + \frac{1}{k} Q_{f,J}^g \right] =$$

$$\sum_{i \in M_I} \sum_J \sum_{j \in M_J} p_{ji}^g (\Sigma^g) \left[ Q_{s,i}^g + \frac{1}{k} Q_{f,i}^g \right]$$

is true.

This definition ensures that:

$$\sum_{i \in M_I} V_i \Sigma_i^g \phi_i^g = V_I \Sigma_I^g \phi_I^g = V_I \tilde{\Sigma}_I^g \tilde{\phi}_I^g$$
Homogenization/Condensation

Microscopic cross section:

- The macroscopic cross section associated with region \( m \) is:

\[
\Sigma^g_m = \sum_I \Sigma^g_{I,m}
\]

- \( \Sigma^g_{I,m} \) is the contribution of isotope \( I \) to \( \Sigma^g_m \):

\[
\Sigma^g_{I,m} = N_{I,m} \sigma^g_I
\]

- \( N_{I,m} \), the concentration of isotope \( I \) in region \( m \).

- The homogenization and condensation procedure described above remain valid for \( \Sigma^g_{I,m} \).
The concentration of isotope $I$ in homogenized region $M$ is:

$$N_{I,M} = \frac{1}{V_M} \sum_{m \in M} N_{I,m} V_m$$

The equivalent homogenized microscopic cross section is therefore:

$$\hat{\sigma}_{I,M}^g = \frac{\mu_M^g}{N_{I,M} V_M \phi_M^g} \sum_{m \in M} N_{I,m} V_m \sigma_I^g \phi_m^g$$

The microscopic cross sections now become dependent on the spatial position.
The isotopic depletion equations:

\[
\frac{dN_I(t)}{dt} = \sum_J S^J_I(t) - N_I(t) \left( \int dE \sigma_{I,a}(E) \phi(E, t) + \lambda_I \right)
\]

- \(\lambda_I\) is the decay rate of \(I\).
- \(\sigma_{I,a}(E)\) is the neutron absorption cross section for \(I\).
- \(S^J_I(t)\) is the production rate for isotope \(I\) from isotope \(J\).
- Initial concentrations \(N_I(t_0)\) required.
Isotopic Depletion

Example of simplified decay chains for $^{235}\text{U}$ and $^{238}\text{U}$.

- $^{235}\text{U} \xrightarrow{(n,\gamma)} 86.67$ barns $^{236}\text{U}$
- $^{236}\text{U}$ $\xrightarrow{\text{fission}}$ $506.8$ barns $\rightarrow$ fission 506.8 barns
- $^{238}\text{U} \xrightarrow{(n,\gamma)} 2.241$ barns $^{239}\text{U}$
- $^{239}\text{U} \xrightarrow{\beta^- (23.45 \text{ minutes})} ^{239}\text{Np}$
- $^{239}\text{Np} \xrightarrow{\beta^- (2.3565 \text{ hours})} ^{239}\text{Pu}$
- $^{238}\text{Pu} \xrightarrow{(n,\gamma)} 274.2$ barns
- $^{239}\text{Pu} \xrightarrow{\text{fission}}$ $698.3$ barns

Fertile isotope

Fissile isotope

ANL DRAGON workshop
Approaches used in DRAGON to solve the isotopic depletion equations:

- Use a quasi-static approach assuming the flux vary slowly over macro time steps (linear variation).
- Solve depletion equation using a Runge-Kutta or Kaps-Renthrop algorithm inside macro time steps with linear flux variations.
Isotopic Depletion

Flux normalization options available:

- Constant flux burnup.
- Constant power burnup.
- Out of core depletion.
Isotopic Depletion

Fuel depletion calculations:

- Depletion is controlled using the time variable.
- Selection of depletion time step is important.
- Includes an isotope saturation model.

For out of core depletion calculations:

- A single time step is sufficient since flux vanishes.

For in core depletion calculations:

- Start with short $\Delta t$.
- Increase $\Delta t$ using a log scale to follow spatial and energy variation in flux as a function of time.
Burnup at fixed power uses an inner iterative process (inside EVO:).

- Energy and spatial flux distribution are assumed constant.
- Assume integrated flux level has a linear time dependence.
- Iterate internally between final flux level and isotopic concentrations to ensure coherence between initial and final power production.
For in core depletion calculations:

- Burnup at fixed power and flux levels can also use an outer iterative process.
  - When energy and spatial flux distribution change are strongly affected by depletion.
  - Iterate externally between flux calculations (FLU:) and isotopic concentrations (EVO:).
  - This ensures that the final isotopic concentrations are coherent with the final energy and spatial flux distribution.
Status of DRAGON

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École Polytechnique de Montréal
Contents:

1. Verification.
2. Validation.
3. Documentation.
4. Using DRAGON.
DRAGON verification exercise:

1. The EXCELT: Module.
2. The ASM: Module.
3. The FLU: Module.
4. The EDI: Module
Verification

DRAGON initial verification:

- Select the DRAGON modules that are critical to 3–D CANDU applications.
- Select a series of routines in these modules for line by line verification.
- Select routines or blocks of routines for numerical verification.
  - Self-consistency tests.
  - Verification using an independent method or program.

Here I will concentrate on the numerical verification.
The EXCEL: Module:

- To analyze an already defined geometry.
- To assign to each region and surface an index and to compute the volume and surface area.
- Using quadrature parameters, to generate integration lines.
- Verify the integration lines and perform normalization if required.
Verification

The TRACKING data structure also contains a description of the geometry:

- Compare contents of TRACKING and GEOMETRY data structure.
- Use PSP: module to illustrate 2–D geometry:
Use Mathematica to illustrate 2–D and 3–D geometries using the information provided on the tracking data structure:

A Mathematica interface was written to access the TRACKING data structure.
Surface area and region volume evaluation:

- Comparison of DRAGON region volumes with explicit calculations.
- Comparison of DRAGON volume for equivalent 2–D and 3–D geometries:
Binary tracking file:

- Numerical evaluation of volume and comparison with exact values:

\[ V_i = h \sum_k l_{k,i} \]

- \( h \) is spatial integration weight.
- \( l_{k,i} \) is path length of line \( k \) in region \( i \).
### 2–D tracking volume errors:

<table>
<thead>
<tr>
<th>Integration parameters</th>
<th>Errors (%)</th>
</tr>
</thead>
<tbody>
<tr>
<td>Angular</td>
<td>Spatial</td>
</tr>
<tr>
<td>3</td>
<td>1.0</td>
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<tr>
<td>3</td>
<td>10.0</td>
</tr>
<tr>
<td>3</td>
<td>100.0</td>
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<tr>
<td>31</td>
<td>1.0</td>
</tr>
<tr>
<td>31</td>
<td>10.0</td>
</tr>
<tr>
<td>31</td>
<td>100.0</td>
</tr>
<tr>
<td>91</td>
<td>1.0</td>
</tr>
<tr>
<td>91</td>
<td>10.0</td>
</tr>
<tr>
<td>91</td>
<td>100.0</td>
</tr>
</tbody>
</table>
### 3–D tracking volume errors:

<table>
<thead>
<tr>
<th>Integration parameters</th>
<th>Errors (%)</th>
<th></th>
<th></th>
</tr>
</thead>
<tbody>
<tr>
<td></td>
<td>Maximum</td>
<td>Average</td>
<td></td>
</tr>
<tr>
<td>Angular</td>
<td>Spatial</td>
<td>Nb. lines</td>
<td></td>
</tr>
<tr>
<td>2</td>
<td>1.0</td>
<td>40820</td>
<td>6.61</td>
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<tr>
<td>2</td>
<td>5.0</td>
<td>203660</td>
<td>1.49</td>
</tr>
<tr>
<td>2</td>
<td>25.0</td>
<td>1021544</td>
<td>0.29</td>
</tr>
<tr>
<td>4</td>
<td>1.0</td>
<td>110644</td>
<td>6.61</td>
</tr>
<tr>
<td>4</td>
<td>5.0</td>
<td>555118</td>
<td>1.49</td>
</tr>
<tr>
<td>4</td>
<td>25.0</td>
<td>2775748</td>
<td>0.20</td>
</tr>
<tr>
<td>8</td>
<td>1.0</td>
<td>359432</td>
<td>5.75</td>
</tr>
<tr>
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<td>5.0</td>
<td>1793640</td>
<td>0.98</td>
</tr>
<tr>
<td>8</td>
<td>25.0</td>
<td>8982529</td>
<td>0.37</td>
</tr>
</tbody>
</table>
Verification

Extract integration line from binary tracking fine and plot over geometry using Matlab in 2–D:
Extract integration line from binary tracking fine and plot over geometry using Mathematica in 3-D:

- 12540 lines generated.
- Only a few lines plotted.
- 16/178 surfaces not touched by lines.
2–D projection of 3–D tracking lines using Matlab:
Verification

The ASM: Module:

- To compute the collision, leakage and escape probability matrix using a set of integration lines.
- To verify the CP conservation laws and perform CP normalization if required.
- To take into account the boundary conditions (approximately or exactly).
- To take into account within-group scattering.
CP conservation relations:

\[
\frac{4}{S_\beta} \left( \sum_{\alpha=1}^{N_\alpha} \frac{S_\alpha}{4} p_{\alpha\beta} + \sum_{i=1}^{N_i} \Sigma_i V_i p_{i\beta} \right) = R_\beta
\]

\[
\frac{1}{V_j} \left( \sum_{\alpha=1}^{N_\alpha} \frac{S_\alpha}{4} p_{\alpha j} + \sum_{i=1}^{N_i} \Sigma_i V_i p_{ij} \right) = R_j
\]

- If conservation relations satisfied: \( R_\beta = R_j = 1 \).
- Error on conservation relations:

\[
\epsilon_{\beta/j} = 100 \times \left( \frac{\|1 - R_{\beta/j}\|}{\|R_{\beta/j}\|} \right)
\]
Error on conservation relations:

- Computed directly in DRAGON.
- Before CP normalization $\epsilon_{\beta/j}^B$.
- After CP normalization $\epsilon_{\beta/j}^A$.

<table>
<thead>
<tr>
<th>Parameter</th>
<th>T2D</th>
<th>T3D</th>
</tr>
</thead>
<tbody>
<tr>
<td>$\epsilon_{\beta}^B$</td>
<td>0.23</td>
<td>4.36</td>
</tr>
<tr>
<td>$\epsilon_{\jmath}^B$</td>
<td>0.000003</td>
<td>0.000008</td>
</tr>
<tr>
<td>$\epsilon_{\beta}^A$</td>
<td>0.00004</td>
<td>0.00002</td>
</tr>
<tr>
<td>$\epsilon_{\jmath}^A$</td>
<td>0.00003</td>
<td>0.00003</td>
</tr>
</tbody>
</table>
Complete CP matrix:

\[
P_{c, vv}^g = \left( P_{vv}^g + P_{vs}^g \left( A^{-1} - P_{ss}^g \right)^{-1} P_{sv}^g \right)
\]

where \( A \) is the transmission/reflection matrix.

Scattering modified CP matrix:

\[
W_{vv}^g = \left( I - P_{c, vv}^g \Sigma_s^{g \rightarrow g} \right)^{-1} P_{c, vv}^g
\]

DRAGON evaluation is single precision.

Matlab evaluation is double precision.
Verification

Errors:

- Complete CP matrix:

\[ \epsilon_c = \left( \frac{\| P_{c,vv}^D - P_{c,vv}^M \|}{\| P_{c,vv}^M \|} \right) < 1.0 \times 10^{-6} \]

- Scattering modified CP matrix:

\[ \epsilon_W = \left( \frac{\| W_{vv}^D - W_{vv}^M \|}{\| W_{vv}^M \|} \right) < 3.0 \times 10^{-6} \]

Consistent with single precision nature of DRAGON.
The FLU: Module:

- To solve the collision probability equations for the flux (case with external source).

- To evaluate the problem eigenvalue (case without external source).

- To simulate leakage in an infinite cell and to evaluate the diffusion coefficients.
Scattering modified CP solution:

- Sources:
  \[ Q_{m,i,k}^g = F_{i}^g + \sum_{h=1, h \neq g}^{G} \sum_{s,i}^{h \rightarrow g} \phi_{i,k}^h - 1 \]

- Direct flux solution:
  \[ \phi_{i,k}^g = \sum_{j=1}^{N} W_{ij}^g Q_{j,k}^g \]

- Flux after rebalancing \( \alpha^g \):
  \[ \tilde{\phi}_{i,k}^g = \alpha^g \phi_{i,k}^g \]
### Verification

**DRAGON/Matlab comparison for T3D:**

<table>
<thead>
<tr>
<th>Iteration</th>
<th>Group</th>
<th>Matlab-DRAGON ($\times 10^7$)</th>
</tr>
</thead>
<tbody>
<tr>
<td></td>
<td></td>
<td>$Q$</td>
</tr>
<tr>
<td>$k$</td>
<td>$g$</td>
<td>$\phi$</td>
</tr>
<tr>
<td>1</td>
<td>1</td>
<td>1.0636</td>
</tr>
<tr>
<td>1</td>
<td>2</td>
<td>1.6990</td>
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<tr>
<td>2</td>
<td>1</td>
<td>1.0934</td>
</tr>
<tr>
<td>2</td>
<td>2</td>
<td>1.3440</td>
</tr>
</tbody>
</table>

<table>
<thead>
<tr>
<th></th>
<th>$\hat{\phi}$</th>
</tr>
</thead>
<tbody>
<tr>
<td>1.0636</td>
<td>2.4653</td>
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<tr>
<td>1.6990</td>
<td>2.1748</td>
</tr>
<tr>
<td>1.0934</td>
<td>2.9636</td>
</tr>
<tr>
<td>1.3440</td>
<td>2.1592</td>
</tr>
<tr>
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<tr>
<td></td>
<td>2.8039</td>
</tr>
<tr>
<td></td>
<td>18.1880</td>
</tr>
<tr>
<td></td>
<td>5.7258</td>
</tr>
</tbody>
</table>

Consistent with single precision nature of DRAGON.
The EDI: Module:

- To perform region homogenization of microscopic and macroscopic cross sections.
- To perform group condensation of microscopic and macroscopic cross sections.
- To compare the results from two different executions including:
  - The evaluation of the errors in reaction rates and output cross sections from two different executions.
  - The evaluation of incremental cross sections.
Homogenization and condensation procedure:

- **Flux:**
  \[
  \phi^H = \frac{1}{V_I} \sum_{i \in I} \sum_{h \in H} \phi_i^h V_i
  \]

- **Total cross section:**
  \[
  \Sigma^H_I = \frac{1}{V_I \phi^H} \sum_{i \in I} \sum_{h \in H} \phi_i^h V_i \Sigma^h_i
  \]

- **Scattering cross section:**
  \[
  \Sigma^{H \rightarrow G}_{I,s} = \frac{1}{V_I \phi^H} \sum_{i \in I} \sum_{h \in H} \sum_{g \in G} \phi_i^h V_i \Sigma^{h \rightarrow g}_{i,s}
  \]
### Verification

**DRAGON/Matlab comparison for T3D:**

<table>
<thead>
<tr>
<th>$(\times 10^7)$</th>
<th>Case 1</th>
<th>Case 2</th>
<th>Case 3</th>
</tr>
</thead>
<tbody>
<tr>
<td>$\delta \Sigma^1$</td>
<td>0.0401</td>
<td>-0.5698</td>
<td>-0.1446</td>
</tr>
<tr>
<td>$\delta \Sigma^2$</td>
<td>-0.2713</td>
<td>-0.4730</td>
<td>-0.0075</td>
</tr>
<tr>
<td>$\delta \Sigma^2_{1\rightarrow1}$</td>
<td>-0.2232</td>
<td></td>
<td></td>
</tr>
<tr>
<td>$\delta \Sigma^1_{1\rightarrow1}$</td>
<td>0.0311</td>
<td></td>
<td></td>
</tr>
<tr>
<td>$\delta \Sigma^2_{2\rightarrow2}$</td>
<td>-0.6877</td>
<td></td>
<td></td>
</tr>
<tr>
<td>$\delta \Sigma^1_{2\rightarrow2}$</td>
<td>0.1251</td>
<td></td>
<td></td>
</tr>
</tbody>
</table>

- **Case 1** is for full homogenization/no condensation.
- **Case 2** is for no homogenization/full condensation.
- **Case 3** is for full homogenization/full condensation.
Conclusions of the verification exercise:

- No major error detected in DRAGON using the verification process.

- Output file does not always contain adequate information. This lead to the revision of some output routine to produce more significant information.

- Contents of the data structures not always correctly documented. This lead to revisions in the DRAGON documentation.
Most of the major bugs in DRAGON were resolved using DRAGON self-verification procedures.

Some bugs were resolved using DRAGON over-programmation (in EXCELT and EXCELL modules for example).

Most of the errors remaining in DRAGON are detected by users.
Early DRAGON validation against experimental results:

- Experimental database provided by Chalk River Laboratory.
- Limited to 2–D zero-burnup transport calculations.
- Contains experimental information for a large number of important lattice parameters.
- Compiled for a variety of fuel and moderator types.
Validation

Evaluation of the performance of DRAGON by comparison with benchmark results:

1. NEACRP benchmark BWR lattice cell problems.
2. MOX Fueled VVER Cell Benchmarks.
3. 2–D and 3–D VENUS-2 MOX benchmark.
Validation

DRAGON validation against experimental results:

- Uses WIMS–AECL format Winfrith microscopic cross section library.
- Input data provided in the form of WIMS–AECL input data files.
- Only most reliable or default options of DRAGON considered.
- Cluster geometries in an hexagonal or square cell.
- Radial and axial buckling provided.
Validation

DRAGON Calculation Options:

- The explicit geometry used for self shielding.
- A discretized geometry used for solving the transport equation.
- Moderator, coolant and fuel regions divided into sub-regions of equal thickness not exceeding 0.4 cm.
- Test material are also self-shielded if required.
- Microscopic reaction rates for each test isotope in each region generated once flux distribution is known.
Isotropic collision probability technique (standard $P_{ij}$ method) in full 2–D geometry.

Geometry analyzed for each test case and best compromise between precision and calculation time for tracking selected.

Homogeneous $B_1$ leakage model with total buckling considered $B^2 = B_r^2 + B_z^2$.

Heterogeneous $B_1$ leakage model with radial and axial buckling also considered.
Validation

Validation Parameters:

- Spatial dependence of flux (region $i$):

$$\rho_i = V_{\text{fuel}} \frac{\sum g \sigma_{a,I}^g \phi_i^g}{\sum_{k \in \text{fuel}} V_k \sum g \sigma_{a,I}^g \phi_k^g}$$

- Select isotope $I$ such that $\sigma_{a,I}^g$ behaves as $1/\nu$ ($^{55}\text{Mn}$ or $^{63}\text{Cu}$ for example).
Validation

Neutron energy spectrum (region $i$):

$$N_i^{I/J} = \sum_g \frac{\sigma^g_{a/f,I} \phi^g_i}{\sum_g \sigma^g_{a/f,J} \phi^g_i}$$

Select isotope $I$ and $J$ such that:

$$\frac{\sigma^g_{a/f,I}}{\sigma^g_{a/f,J}} = \alpha (1 + \beta \delta_{gh})$$

to enhance artificially group $h$. Consider:

- Ratio of $^{239}$Pu to $^{235}$U fission.
- Ratio of $^{176}$Lu to $^{55}$Mn absorption.
- Ratio of $^{115}$In to $^{55}$Mn absorption.
Conversion ratio for cell with Uranium fuel:

\[ C_{i}^{\text{fuel}} = \frac{\int v \sigma_{a}^{238}U(v)\phi_{i}(v)dv}{\int v \sigma_{f}^{235}U(v)\phi_{i}(v)dv} \]

Reactivity of cell.

Void effect.

Fast fission ratio.
Validation

- Up to 52 parameters can be compared.
- All the parameters are not necessarily available for all experiments.
- The results are presented in the form of absolute errors or absolute error.
Validation

Results for $k_{\text{eff}}$:

- $B_1^{\text{hom}}$: systematic error of $1.48 \pm 3.63$.
- $B_1^{\text{het}}$: systematic error of $-1.0 \pm 2.83$. 
Relative errors on $k_{\text{eff}}$:

<table>
<thead>
<tr>
<th>Fuel</th>
<th>Coolant</th>
<th>N</th>
<th>$B_1^{\text{hom}}$ $E^r \pm \sigma^r$</th>
<th>$B_1^{\text{het}}$ $E^r \pm \sigma^r$</th>
</tr>
</thead>
<tbody>
<tr>
<td>All</td>
<td>All</td>
<td>197</td>
<td>3.64±7.26</td>
<td>1.15±5.16</td>
</tr>
<tr>
<td></td>
<td>D$_2$O</td>
<td>61</td>
<td>4.40±7.73</td>
<td>1.04±5.42</td>
</tr>
<tr>
<td></td>
<td>Void</td>
<td>75</td>
<td>3.63±7.82</td>
<td>0.88±5.56</td>
</tr>
<tr>
<td></td>
<td>Organic</td>
<td>39</td>
<td>3.97±6.96</td>
<td>2.34±4.99</td>
</tr>
<tr>
<td></td>
<td>H$_2$O</td>
<td>20</td>
<td>0.90±2.94</td>
<td>0.26±2.75</td>
</tr>
</tbody>
</table>
### Relative errors on $k_{\text{eff}}$

<table>
<thead>
<tr>
<th>Fuel</th>
<th>Coolant</th>
<th>N</th>
<th>$B_1^{\text{hom}}$</th>
<th>$B_1^{\text{het}}$</th>
</tr>
</thead>
<tbody>
<tr>
<td></td>
<td></td>
<td></td>
<td>$E^r \pm \sigma^r$</td>
<td>$E^r \pm \sigma^r$</td>
</tr>
<tr>
<td>UO$_2$</td>
<td>All</td>
<td>96</td>
<td>7.21$\pm$8.80</td>
<td>3.43$\pm$6.08</td>
</tr>
<tr>
<td></td>
<td>D$_2$O</td>
<td>35</td>
<td>7.86$\pm$8.66</td>
<td>3.24$\pm$6.05</td>
</tr>
<tr>
<td></td>
<td>Void</td>
<td>39</td>
<td>7.34$\pm$9.14</td>
<td>3.39$\pm$6.23</td>
</tr>
<tr>
<td></td>
<td>Organic</td>
<td>17</td>
<td>7.50$\pm$9.32</td>
<td>4.83$\pm$6.56</td>
</tr>
<tr>
<td></td>
<td>H$_2$O</td>
<td>5</td>
<td>0.63$\pm$2.32</td>
<td>0.37$\pm$2.35</td>
</tr>
</tbody>
</table>
### Validation

Relative errors on $k_{\text{eff}}$:

<table>
<thead>
<tr>
<th>Fuel</th>
<th>Coolant</th>
<th>N</th>
<th>$B_1^{\text{hom}}$</th>
<th>$B_1^{\text{het}}$</th>
</tr>
</thead>
<tbody>
<tr>
<td></td>
<td></td>
<td></td>
<td>$E^r \pm \sigma^r$</td>
<td>$E^r \pm \sigma^r$</td>
</tr>
<tr>
<td>PuUO$_2$</td>
<td>All</td>
<td>57</td>
<td>-0.35 ± 2.36</td>
<td>-1.39 ± 2.50</td>
</tr>
<tr>
<td></td>
<td>D$_2$O</td>
<td>13</td>
<td>-0.08 ± 1.38</td>
<td>-1.47 ± 1.94</td>
</tr>
<tr>
<td></td>
<td>Void</td>
<td>23</td>
<td>-0.65 ± 2.95</td>
<td>-1.80 ± 3.18</td>
</tr>
<tr>
<td></td>
<td>Organic</td>
<td>10</td>
<td>-0.06 ± 1.48</td>
<td>-0.61 ± 1.60</td>
</tr>
<tr>
<td></td>
<td>H$_2$O</td>
<td>9</td>
<td>-0.69 ± 2.86</td>
<td>-1.39 ± 2.36</td>
</tr>
</tbody>
</table>
Results for void effect:

- $B_1^{\text{hom}}$: systematic error of $-0.68 \pm 1.05$.
- $B_1^{\text{het}}$: systematic error of $-0.5 \pm 0.89$.
- $B_1^{\text{het}}$ slightly better than $B_1^{\text{hom}}$. 
Relative errors on void effect:

<table>
<thead>
<tr>
<th>Fuel</th>
<th>N</th>
<th>$B_1^{\text{hom}}$</th>
<th>$B_1^{\text{het}}$</th>
</tr>
</thead>
<tbody>
<tr>
<td></td>
<td></td>
<td>$E^r \pm \sigma^r$</td>
<td>$E^r \pm \sigma^r$</td>
</tr>
<tr>
<td>All</td>
<td>55</td>
<td>-0.68±1.05</td>
<td>-0.50±0.89</td>
</tr>
<tr>
<td>UO$_2$ (N)</td>
<td>28</td>
<td>-1.31±1.10</td>
<td>-0.95±0.99</td>
</tr>
<tr>
<td>UO$_2$ (D)</td>
<td>3</td>
<td>0.18±0.49</td>
<td>0.18±0.49</td>
</tr>
<tr>
<td>U-Metal</td>
<td>6</td>
<td>-0.06±0.60</td>
<td>-0.04±0.60</td>
</tr>
<tr>
<td>UC</td>
<td>5</td>
<td>-0.56±0.23</td>
<td>-0.49±0.21</td>
</tr>
<tr>
<td>U$_3$Si</td>
<td>2</td>
<td>-0.32±0.31</td>
<td>-0.26±0.27</td>
</tr>
<tr>
<td>PuUO$_2$</td>
<td>11</td>
<td>0.20±0.17</td>
<td>0.16±0.24</td>
</tr>
</tbody>
</table>
Validation

Results for neutron spatial density ratio:

- Results nearly independent of leakage model.
- For most cases $|E^r| < 1$ in fuel.
- Error is more important at cell edge.
### Validation

#### Relative errors for neutron density inner fuel pin:

<table>
<thead>
<tr>
<th>Fuel</th>
<th>Coolant</th>
<th>N</th>
<th>$B_1^{\text{hom}}$</th>
<th>$B_1^{\text{het}}$</th>
</tr>
</thead>
<tbody>
<tr>
<td></td>
<td></td>
<td></td>
<td>$E^r \pm \sigma^r$</td>
<td>$E^r \pm \sigma^r$</td>
</tr>
<tr>
<td>All</td>
<td>All</td>
<td>75</td>
<td>0.25±0.65</td>
<td>0.27±0.63</td>
</tr>
<tr>
<td></td>
<td>D$_2$O</td>
<td>20</td>
<td>0.50±0.55</td>
<td>0.50±0.55</td>
</tr>
<tr>
<td></td>
<td>Void</td>
<td>26</td>
<td>0.38±0.49</td>
<td>0.37±0.49</td>
</tr>
<tr>
<td></td>
<td>Organic</td>
<td>20</td>
<td>0.00±0.82</td>
<td>0.07±0.78</td>
</tr>
<tr>
<td></td>
<td>H$_2$O</td>
<td>9</td>
<td>-0.11±0.62</td>
<td>-0.05±0.59</td>
</tr>
</tbody>
</table>
Relative errors for neutron density cell edge:

<table>
<thead>
<tr>
<th>Fuel</th>
<th>Coolant</th>
<th>N</th>
<th>$B_1^{\text{hom}}$ $E^r \pm \sigma^r$</th>
<th>$B_1^{\text{het}}$ $E^r \pm \sigma^r$</th>
</tr>
</thead>
<tbody>
<tr>
<td>All</td>
<td>All</td>
<td>77</td>
<td>0.26±1.37</td>
<td>0.25±1.34</td>
</tr>
<tr>
<td>All</td>
<td>D$_2$O</td>
<td>20</td>
<td>0.32±0.85</td>
<td>0.38±0.85</td>
</tr>
<tr>
<td>All</td>
<td>Void</td>
<td>28</td>
<td>0.19±1.12</td>
<td>0.35±1.10</td>
</tr>
<tr>
<td>All</td>
<td>Organic</td>
<td>20</td>
<td>1.00±1.41</td>
<td>0.78±1.33</td>
</tr>
<tr>
<td>All</td>
<td>H$_2$O</td>
<td>9</td>
<td>-1.27±1.77</td>
<td>-1.52±1.66</td>
</tr>
</tbody>
</table>
Validation

Results for neutron energy distribution:

- Results nearly independent of leakage model.
- $|E^r| < 1$ in fuel.
- $|E^r| < 1$ at cell edge except for H$_2$O cooled cells.
Relative errors for fuel average $^{239}\text{Pu}/^{235}\text{U}$ fissions:

<table>
<thead>
<tr>
<th>Fuel</th>
<th>Coolant</th>
<th>N</th>
<th>$B_{1}^{\text{hom}}$</th>
<th>$B_{1}^{\text{het}}$</th>
</tr>
</thead>
<tbody>
<tr>
<td>All</td>
<td>All</td>
<td>4</td>
<td>-0.12 ± 1.34</td>
<td>-0.10 ± 1.34</td>
</tr>
<tr>
<td>All</td>
<td>D$_2$O</td>
<td>10</td>
<td>0.29 ± 1.69</td>
<td>0.30 ± 1.69</td>
</tr>
<tr>
<td>All</td>
<td>Void</td>
<td>14</td>
<td>-0.07 ± 1.20</td>
<td>-0.08 ± 1.20</td>
</tr>
<tr>
<td>All</td>
<td>Organic</td>
<td>7</td>
<td>0.03 ± 1.66</td>
<td>0.09 ± 1.66</td>
</tr>
<tr>
<td>All</td>
<td>H$_2$O</td>
<td>11</td>
<td>-0.62 ± 0.88</td>
<td>-0.59 ± 0.90</td>
</tr>
</tbody>
</table>
Relative errors for cell edge $^{239}$Pu/$^{235}$U fissions:

<table>
<thead>
<tr>
<th>Fuel</th>
<th>Coolant</th>
<th>N</th>
<th>$B_1^{\text{hom}}$ $E^r \pm \sigma^r$</th>
<th>$B_1^{\text{het}}$ $E^r \pm \sigma^r$</th>
</tr>
</thead>
<tbody>
<tr>
<td>All</td>
<td>All</td>
<td>6</td>
<td>-0.75$\pm$1.32</td>
<td>-0.74$\pm$1.32</td>
</tr>
<tr>
<td></td>
<td>D$_2$O</td>
<td>2</td>
<td>-0.09$\pm$0.05</td>
<td>-0.10$\pm$0.05</td>
</tr>
<tr>
<td>All</td>
<td>Void</td>
<td>2</td>
<td>-0.91$\pm$1.07</td>
<td>-0.91$\pm$1.07</td>
</tr>
<tr>
<td></td>
<td>Organic</td>
<td>0</td>
<td>-1.24$\pm$2.49</td>
<td>-1.21$\pm$2.51</td>
</tr>
<tr>
<td></td>
<td>H$_2$O</td>
<td>2</td>
<td></td>
<td></td>
</tr>
</tbody>
</table>
Validation

Conclusions of early validation attempt:

- Results are relatively good for most cases.
- Major problem are for relative conversion ratio in individual fuel pins (Probably due to $^{238}$U self-shielding).
- Gain from $B_1^{\text{hom}}$ to $B_1^{\text{het}}$ not considerable except possibly for $k_{\text{eff}}$ and void effect.
- Other libraries still remain to be tested.
Validation

MOX Fueled VVER Cell Benchmarks.
- Hexagonal cell.
- Various mixture composition including 2 types of MOX fuel.
- Local parameters variations.
- Burnup dependent.

Results available:
- ENDF/B-V based SAS2H calculations.
- ENDF/B-VI based HELIOS calculations.
Geometry:

- Hexagonal cell with $l = 1.275$ cm lattice pitch.
- Central fuel region with Zirconium fuel clad and light water moderator.

Neutron leakage:

- $k_\infty$ for cell without leakage.
- $k_{\text{eff}}$ for cell with imposed leakage (0.003 cm$^{-2}$).
- $k_0$ for critical cell with buckling search.
Validation

8 different fuel compositions including:

- Variant V1: 4.4% enriched LEU fuel.
- Variant V2: MOX from weapon grade plutonium.
- Variant V3: Spent 4.4% LEU fuel with fission products.
- Variant V10: MOX from reactor grade plutonium.
Validation

Local parameters for various operating conditions:

- State S1 for normal operating conditions: Fuel at 1027 K; Cladding+moderator at 579 K; $^{135}$Xe and $^{149}$Sm at saturation; 600 ppm of boron in moderator.
- State S3: S1 without boron.
- State S4: S1 without $^{135}$Xe and $^{149}$Sm.
- State S5: S4 with fuel+cladding+moderator at 579 K.
- State S6: S4 with fuel+cladding+moderator at 300 K.
Validation

Depletion calculations:
- All local parameters in state S1.
- Fuel variants V1 with $P = 41.9109$ kW/kg.
- Fuel variants V2 with $P = 41.7837$ kW/kg.
- Fuel variants V10 with $P = 34.0078$ kW/kg.
- Maximum burnup step of 2 MWd/kg.
- Final burnup of 60 MWd/kg.
Validation

Libraries for DRAGON:

- ENDF/B-V and ENDF/B-VI based 89 group WIMS-AECL library.
- No curium isotopes in ENDF/B-V library.
- No self-shielding for $^{241}$Pu, $^{242}$Pu and $^{241}$Am in ENDF/B-VI library.

2–D hexagonal pin-cell calculation for DRAGON using CP method.
Validation

Overall comparison:

\[ \Delta k^{\text{Code}}(mk) = 1000 \times \left( \frac{1}{k^\text{HELIOS}} - \frac{1}{k^\text{Code}} \right) \]

with:

- Code=DRGE5 represents DRAGON with ENDF/B-V WIMS-AECL library.
- Code=DRGE6 represents DRAGON with ENDF/B-VI WIMS-AECL library.
## Validation

$\Delta k_0$ (mk):

<table>
<thead>
<tr>
<th>Variant</th>
<th>Code</th>
<th>S1</th>
<th>S3</th>
<th>S4</th>
<th>S5</th>
<th>S6</th>
</tr>
</thead>
<tbody>
<tr>
<td></td>
<td></td>
<td></td>
<td></td>
<td></td>
<td></td>
<td></td>
</tr>
<tr>
<td>V1</td>
<td>SAS2H</td>
<td>-10.3</td>
<td>-10.1</td>
<td>-9.6</td>
<td>-8.5</td>
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</tr>
<tr>
<td></td>
<td>DRGE5</td>
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<td>-9.4</td>
<td>-8.2</td>
<td>-4.9</td>
</tr>
<tr>
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<td>DRGE6</td>
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<td>-2.5</td>
<td>-2.4</td>
<td>-2.6</td>
</tr>
<tr>
<td>V2</td>
<td>SAS2H</td>
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<td>-9.4</td>
<td>-9.0</td>
<td>-7.6</td>
<td>-3.7</td>
</tr>
<tr>
<td></td>
<td>DRGE5</td>
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<td>-6.3</td>
<td>-5.8</td>
<td>-4.8</td>
<td>-1.9</td>
</tr>
<tr>
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<td>DRGE6</td>
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<td>-0.7</td>
<td>-0.3</td>
<td>-0.3</td>
<td>-1.0</td>
</tr>
<tr>
<td>V3</td>
<td>SAS2H</td>
<td>-11.5</td>
<td>-11.3</td>
<td>-10.8</td>
<td>-9.7</td>
<td>-5.1</td>
</tr>
<tr>
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<td>DRGE5</td>
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<td>-4.5</td>
<td>-3.5</td>
<td>-3.2</td>
<td>-0.9</td>
</tr>
<tr>
<td></td>
<td>DRGE6</td>
<td>-3.9</td>
<td>-4.1</td>
<td>-3.3</td>
<td>-4.1</td>
<td>-4.2</td>
</tr>
<tr>
<td>V10</td>
<td>SAS2H</td>
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<td>-16.4</td>
<td>-16.0</td>
<td>-15.6</td>
<td>-10.9</td>
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<tr>
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<td>-22.3</td>
<td>-22.1</td>
<td>-23.1</td>
<td>-17.1</td>
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<tr>
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<td>DRGE6</td>
<td>-8.9</td>
<td>-9.0</td>
<td>-8.7</td>
<td>-10.9</td>
<td>-10.8</td>
</tr>
</tbody>
</table>
Comments:

- SAS2H average error -12.7 mk.
- DRGE5 average error -10.3 mk.
- DRGE6 average error -3.7 mk.
- Main differences between HELIOS and DRGE6:

  - due to the absence of resonance tables for some actinides (fuel variants V3 and V10).
  - $\Delta k_{\text{eff}} - \Delta k_0 \approx 4.0$ mk for V2 which indicates that $^{239}\text{Pu}$ has a large effect on leakage.
Reactivity effects comparison:

\[ \delta \rho_0^{\text{Code}}(S_i - S_j) = 100 \times \left( \frac{\rho_0^{\text{Code}}(S_i - S_j) - \rho_0^{\text{HELIOS}}(S_i - S_j)}{\rho_0^{\text{HELIOS}}(S_i - S_j)} \right) \]

\[ \rho_0(S_i - S_j) = 1000 \times \left( \frac{1}{k_0(S_i)} - \frac{1}{k_0(S_i)} \right) \]

- S3-S1 → reactivity worth of 600 ppm of boron.
- S4-S1 → reactivity worth of saturating fission products.
- S5-S4 → the Doppler effect in the fuel.
- S6-S5 → uniform temperature increase in the cell including the effect of coolant density and boron concentration reduction.
### Validation

\[ \delta \rho_0 \, (\%):\]

<table>
<thead>
<tr>
<th>Variant</th>
<th>Code</th>
<th>State</th>
</tr>
</thead>
<tbody>
<tr>
<td></td>
<td></td>
<td>S3-S1</td>
</tr>
<tr>
<td>V1</td>
<td>SAS2H</td>
<td>0.6</td>
</tr>
<tr>
<td></td>
<td>DRGE5</td>
<td>-1.8</td>
</tr>
<tr>
<td></td>
<td>DRGE6</td>
<td>-2.5</td>
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<td>V2</td>
<td>SAS2H</td>
<td>0.2</td>
</tr>
<tr>
<td></td>
<td>DRGE5</td>
<td>-0.2</td>
</tr>
<tr>
<td></td>
<td>DRGE6</td>
<td>-1.8</td>
</tr>
<tr>
<td>V3</td>
<td>SAS2H</td>
<td>0.5</td>
</tr>
<tr>
<td></td>
<td>DRGE5</td>
<td>-0.5</td>
</tr>
<tr>
<td></td>
<td>DRGE6</td>
<td>-0.7</td>
</tr>
<tr>
<td>V10</td>
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<tr>
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<td>1.1</td>
</tr>
<tr>
<td></td>
<td>DRGE6</td>
<td>-1.2</td>
</tr>
</tbody>
</table>
Validation

Comments:

- Reactivity worth of boron predicted to within 4.7 %.
- Reactivity worth of Xe and Sm predicted to within 6.2 %.
- Reactivity worth of fuel Doppler effect predicted to within 18 %.
- Reactivity worth of uniform fuel temperature change predicted to within 37 %.
Depletion for fuel variant V2:

![Graph showing depletion for fuel variant V2 with markers for HELIOS, SAS2H, DRGE5, and DRGE6.](image)

- **$k_0$** as a function of **Burnup (MWd/kg)**
Depletion for fuel variant V10:

![Graph showing depletion for fuel variant V10](image-url)
$^{235}\text{U}$ concentration for fuel V2 ($10^{-5}$ barns$^{-1}$ cm$^{-1}$):

<table>
<thead>
<tr>
<th>B</th>
<th>HELIOS</th>
<th>SAS2H</th>
<th>DRGE5</th>
<th>DRGE6</th>
</tr>
</thead>
<tbody>
<tr>
<td>0.0</td>
<td>3.84</td>
<td>3.84</td>
<td>3.84</td>
<td>3.84</td>
</tr>
<tr>
<td>10.0</td>
<td>3.17</td>
<td>3.17</td>
<td>3.17</td>
<td>3.18</td>
</tr>
<tr>
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<td>2.56</td>
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</tr>
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<td>2.01</td>
<td>2.01</td>
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</tr>
<tr>
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<td>1.17</td>
<td>1.15</td>
<td>1.15</td>
</tr>
<tr>
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<td>0.824</td>
<td>0.869</td>
<td>0.83</td>
<td>0.835</td>
</tr>
</tbody>
</table>
\(^{239}\)Pu concentration for fuel V2 \((10^{-4} \text{ barns}\, \text{cm}^{-1})\):

<table>
<thead>
<tr>
<th>B</th>
<th>HELIOS</th>
<th>SAS2H</th>
<th>DRGE5</th>
<th>DRGE6</th>
</tr>
</thead>
<tbody>
<tr>
<td>0.0</td>
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<td>6.59</td>
<td>6.59</td>
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<tr>
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<td>5.16</td>
<td>5.16</td>
<td>5.15</td>
</tr>
<tr>
<td>20.0</td>
<td>3.97</td>
<td>4.08</td>
<td>4.06</td>
<td>4.03</td>
</tr>
<tr>
<td>30.0</td>
<td>3.13</td>
<td>3.30</td>
<td>3.24</td>
<td>3.20</td>
</tr>
<tr>
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<td>2.76</td>
<td>2.68</td>
<td>2.62</td>
</tr>
<tr>
<td>50.0</td>
<td>2.17</td>
<td>2.42</td>
<td>2.39</td>
<td>2.23</td>
</tr>
<tr>
<td>60.0</td>
<td>1.93</td>
<td>2.20</td>
<td>2.05</td>
<td>1.97</td>
</tr>
</tbody>
</table>
Validation

$^{135}$Xe concentration for fuel V2 ($10^{-8}$ barns$^{-1}$cm$^{-1}$):

<table>
<thead>
<tr>
<th>B</th>
<th>HELIOS</th>
<th>SAS2H</th>
<th>DRGE5</th>
<th>DRGE6</th>
</tr>
</thead>
<tbody>
<tr>
<td>0.0</td>
<td>0.94</td>
<td>0.946</td>
<td>0.946</td>
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</tr>
<tr>
<td>10.0</td>
<td>1.57</td>
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<td>1.53</td>
<td>1.39</td>
</tr>
<tr>
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<td>1.39</td>
<td>1.26</td>
</tr>
<tr>
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<td>1.29</td>
<td>1.31</td>
<td>1.27</td>
<td>1.15</td>
</tr>
<tr>
<td>40.0</td>
<td>1.18</td>
<td>1.22</td>
<td>1.16</td>
<td>1.05</td>
</tr>
<tr>
<td>50.0</td>
<td>1.08</td>
<td>1.14</td>
<td>1.07</td>
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<tr>
<td>60.0</td>
<td>1.01</td>
<td>1.02</td>
<td>0.996</td>
<td>0.889</td>
</tr>
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</table>
149Sm concentration for fuel V2 ($10^{-7}$ barns$^{-1}$cm$^{-1}$):

<table>
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<th>B</th>
<th>HELIOS</th>
<th>SAS2H</th>
<th>DRGE5</th>
<th>DRGE6</th>
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<tbody>
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</tr>
<tr>
<td>10.0</td>
<td>1.79</td>
<td>1.74</td>
<td>1.76</td>
<td>1.88</td>
</tr>
<tr>
<td>20.0</td>
<td>1.75</td>
<td>1.74</td>
<td>1.74</td>
<td>1.88</td>
</tr>
<tr>
<td>30.0</td>
<td>1.66</td>
<td>1.68</td>
<td>1.66</td>
<td>1.77</td>
</tr>
<tr>
<td>40.0</td>
<td>1.56</td>
<td>1.60</td>
<td>1.56</td>
<td>1.63</td>
</tr>
<tr>
<td>50.0</td>
<td>1.46</td>
<td>1.52</td>
<td>1.46</td>
<td>1.49</td>
</tr>
<tr>
<td>60.0</td>
<td>1.37</td>
<td>1.31</td>
<td>1.36</td>
<td>1.37</td>
</tr>
</tbody>
</table>
Validation

Conclusions:

- ENDF/B-VI based DRAGON results are consistent with HELIOS.
- Main problem is for fuel variants that contain isotopes without resonance tables in the library.
- For depletion calculations, there seems to be a problem with the $^{135}$Xe and $^{149}$Sm concentration at high burnup.
General documents for DRAGON 3.04B:

- User’s guide (IGE-174).
- A guide describing the contents of the DRAGON data structures (IGE-232).
- Geometry numbering in DRAGON (IGE-233).
- The first part of the DRAGON theory manual (IGE-236).
- Programmer’s guide (IGE-251).
GANLIB and CLE-2000 reports:

- The GAN generalized driver (IGE-158).
Restricted DRAGON report:


Other documentations:

- A few reports comparing DRAGON with benchmarks.
- A large number of articles describing:
  - The methods implemented in DRAGON.
  - Comparison of DRAGON with benchmarks.
  - Comparison of DRAGON with other codes.
  - Use of DRAGON for modeling different problems.
Using DRAGON

Contents:

- The DRAGON input file.
- The DRAGON modules.
- The DRAGON data structures.
- Flow chart in DRAGON input decks.
Using DRAGON

Input file format:

- 72 columns.
- free format instruction.
- comments can be inserted in the input deck using:
  - * in first column.
  - ! anywhere.
- end of input for a command or instruction is indicated by:
  - ;
A typical input file for DRAGON contains instructions to:

- Define the modules that will be used.
- Define the data structures that are required and specify their format.
- Define local variables required to control the execution.
- Define the local procedures required for execution.
- Import the files required for the execution.
- Describe the sequence of modules to execute with their input data.
- Clean up, export files and finish the execution.
The DRAGON cross section generation modules:

- **MAC**: module used to generate or modify a DRAGON MACROLIB.

- **LIB**: module used to generate or modify a DRAGON MICROLIB. It can read different types of microscopic cross-section libraries. Cross sections are interpolated in temperature and dilution. It can access microscopic cross-section libraries in the formats DRAGLIB, MATXS, WIMS–D4 and WIMS–AECL. This module also produces a MACROLIB that is compatible with that generated with the MAC: module.

The spatial location of the mixtures defined in MAC: and LIB: will be defined using the GEO: module.
The DRAGON geometry analysis modules:

- **GEO**: module used to generate or modify a geometry.
- **JPMT**: the tracking module for approximate $J_{\pm}$ CP calculations.
- **SYBILT**: the tracking module for approximate Interface Current CP calculations.
- **EXCELT**: the tracking module for the standard 2–D and 3–D CP calculations and for 2–D MOC calculations. This is the most precise and general tracking procedure in DRAGON.
- **BIVACT**: a 2–D diffusion tracking module used for homogenization purpose.
The CP flux calculation modules:

- **ASM**: module which uses the tracking information to generate a multigroup response or collision probability matrix.

- **FLU**: module which uses the multigroup response or collision probability matrix to solve the transport equation for the flux.
The alternate CP evaluation module:

- **EXCELL**: module combines the tracking module
  **EXCEL T**: and the assembly module **ASM**: thereby avoiding the generation of a binary tracking file for 3–D CP calculations.

The MOCC flux calculation module:

- **MOCC**: module to solve the transport equation for the flux using the cyclic method of characteristics.
Additional MICROLIB processing modules:

- **SHI**: used to perform resonance self-shielding calculations.
- **EVO**: used for isotopic depletion calculations.
- **INFO**: a module to compute the density and isotopic contents of heavy or light water and the isotopic contents of UO$_2$ or ThUO$_2$ fuels.
The main reaction rates processing modules:

- **EDI**: editing module for group condensation and spatial homogenization of cross sections based on reaction rates conservation. The direct flux/volume homogenization technique is programmed as well at the SPH method for homogenization by equivalence.

- **CPO**: module for building the standard burnup dependant reactor database used in DONJON.

- **CFC**: module for building the local parameter burnup dependent reactor database based on feedback coefficients that can be used in DONJON.
Some utility modules:

- **MRG**: module to pre-homogenize a geometry tracked using the module **EXCELT**.

- **PSP**: module to generate PostScript images for 2–D geometries which can be tracked using the module **EXCELT**.
Data structures

The formats of the data structures permitted are:

- **LINKED_LIST**: data structure stored in memory.
- **XSM_FILE** data structure stored on direct-access file.
- **SEQ_BINARY** data structure stored on sequential binary file.
- **SEQ_BINARY** data structure stored on sequential ASCII file.

The name of a data structure is case sensitive.
Data structures

**LINKED_LIST**: data structure stored in memory:

- Generally used to transfer information between modules (exception: the binary tracking file).
- Should be avoided for ASMPI data structure (because one may rapidly run into memory limitations).
- Should not be used if one wants to keep the information in the data structure once the execution is completed (for restart for example)
- It is currently the only format permitted for the TRACKING data structure.
Data structures

**XSM_FILE**: data structure stored on direct-access file:

- Can be used to transfer information between modules (exception: the binary tracking file and the TRACKING data structure).
- Should be used for ASMPI data structure (unless the number of regions and group in the problem are small).
- Should be used if one wants to keep the information stores the data structure once the execution is completed (for restart for example).
- Can be used to create a backup of a LINKED_LIST on a direct-access file.
- The contents of this file is system dependent.
**Data structures**

**SEQ_BINARY**: sequential binary file:
- Mainly used to store the binary tracking file.
- Can be used to create a backup of a `LINKED_LIST` or a `XSM_FILE` file on a sequential binary file.
- The contents of this file is system dependent.

**SEQ_ASCII**: sequential binary file:
- Generally used to create a backup of a `LINKED_LIST` or a `XSM_FILE` file on a sequential ASCII file.
- The contents of this file is system independent.
- This file can be viewed online using a simple text editor.
A MACROLIB contains group ordered macroscopic cross-sections library:

- It can be a stand-alone structure or it can be included into a larger structure, such as a MICROLIB or an EDITION structure.
- It must be stored either in a linked list or an XSM file.
- It can be created by the MAC:, LIB: and EDI: modules.
- It can also be modified by the MAC:, LIB:, SHI: and EVO: modules.
A MICROLIB contains a microscopic and a macroscopic cross-section library:

- It always contains a MACROLIB substructure.
- It can be a stand-alone structure or included into a larger structure, such as an EDITION structure.
- It must be stored either in a linked list or an XSM file.
- It can be created by the LIB: and EDI: modules.
- It can also be modified by the LIB:, MAC:, SHI: and EVO: modules.
A GEOMETRY a contains "geometry" description:

- It can be a stand-alone structure or included into a larger structure, such as another GEOMETRY structure.
- It must be stored either in a linked list or an XSM file.
- It can be created and modified by the GEO: module.
Data structures

A TRACKING contains the general geometry description, the result of the geometry analysis and the main tracking parameters:

- It is a stand-alone structure.
- It must be stored on a linked list.
- It can be created by the JPMT:, SYBILT:, EXCEL and EXCELLE: modules.
- It can be modified by the MRG: modules.
Data structures

A ASMPIJ contains the multigroup response and collision probability matrices:

- It is a stand-alone structure.
- It must be stored on a linked list or an XSM file.
- It is created by the ASM: or EXCELL: modules.
Data structures

A FLUXUNK contains the multigroup fluxes and the problem eigenvalue:

- It is a stand-alone structure.
- It must be stored on a linked list or an XSM file.
- It is created by the FLU: and MOCC: modules.
A \textit{EDITION} contains condensed condensed and/or merged microscopic and macroscopic cross-sections libraries:

- It is a stand-alone structure but can contain MACROLIB and MICROLIB substructure.
- It must be stored on a linked list or an XSM file.
- It is created by the \texttt{EDI} module.
A BURNUP contains time dependent isotopic concentration:

- It is a stand-alone structure.
- It must be store on a linked list or an XSM file.
- It is created by the EVO: module.
A CPO contains a simple reactor database:

- It is a stand-alone structure.
- It must be store on a linked list or an XSM file.
- It is created by the CPO module.
A `FBMXSDB` contains a reactor database with feedback coefficients:

- It is a stand-alone structure.
- It must be stored on a linked list or an XSM file.
- It is created by the `CFC` module.
Local variable types:

- INTEGER
- REAL
- DOUBLE
- STRING 72 character long string variables enclosed between " "
- LOGICAL logical variable resulting from a comparison of defined using $True_L$ or $False_L$

The name of a variable is case sensitive.
Operations on variables:

- Arithmetic operations in RPN: +, -, *, / and **:
  \[ \text{delta} := b \ 2 \ ** \ 4 \ . \ a \ c \ * \ * \ - \]

- Unary operations: COS, SQRT, ABS, NOT and LN:
  \[ \text{delta} := \text{delta} \ \text{SQRT} \]

- Relational operations: <, >, <> and <=:
  \[ \text{condition} := a \ b \ <= \]

- Mixed mode operations: R_TO_I, I_TO_R, etc.:
  \[ \text{real} := \text{integer} \ \text{I_TO_R} ; \]
Assigning or evaluating variables:

```
REAL (variable names) := (value) ;
EVALUATE (variable names) := (value) ;
ECHO (variable names) ;
```

Using variable as input/output data to a module:

- `<< a >>` the variable `a` is used as input data.
- `>> a <<` the variable `a` is used as output data.
**Working with CLE-2000**

**IF/THEN/ELSE statement:**

```
IF (condition)  THEN  (statements)
ELSEIF (condition) THEN  (statements)
ELSE  (statements)
ENDIF  
```

**WHILE and REPEAT statements:**

```
WHILE (condition)  DO  |  REPEAT  
   (modif condition)     (modif condition)  
   (statements)          (statements)   
ENDWHILE  ;  |  UNTIL (condition)  ;
```
Input deck procedures:

- used to segment large input deck into sub-deck.
- the main input deck can be seen as a FORTRAN program while the sub-decks can be viewed as a subroutine.
- like in a subroutine the sub-decks can be programmed to use input variables.
Example of an input deck procedure (File PWRPinCell.c2m):

*----
* Procedure Geometry in file PWRPinCell.c2m
* Voltrak TrackFile := PWRPinCell ::
*  <<LocPitch>> <<LocRadius>> <<LocSplitR>>
*  >>MaxReg<< ;
* Procedure parameters
*----
PARAMETER Voltrak TrackFile ::
   ::: LINKED_LIST Voltrak ;
   ::: SEQ_BINARY TrackFile ;
REAL LocPitch LocRadius ;
INTEGER LocSplitR ;
   :: >>LocPitch<< >>LocRadius<< >>LocSplitR<< ;
*----
* Local information and procedure description
*----

MODULE GEO: EXCELT: DELETE: END: ;
LINKED_LIST LocalGeo ;
INTEGER MaxRegion := Mesh 1 + ;
LocalGeo := GEO: :: CARCEL 1
   X- REFL 0.0 <<LocPitch>> X+ REFL
   Y- REFL 0.0 <<LocPitch>> Y+ REFL
   RADIUS 0.0 <<LocRadius>> MIX 1 2 SPLITR <<LocSplitR>> ;
Voltrk TrackFile := EXCELT: LocalGeo ::
   MAXR <<MaxRegion>>
   TRAK TISO 20 100.0 ;
LocalGeo := DELETE: LocalGeo ;
:: <<MaxRegion>> ;
END: ;
Calling the procedure PWRPinCell (File PWRPinCell.c2m):

*----
* Calling procedure PWRPinCell.c2m
* Voltrak TrackFile := PWRPinCell ::
*  <<LocPitch>> <<LocRadius>> <<LocSplitR>> >>MaxReg<< ;
*----
LINKED_LIST Voltrak ;
SEQ_BINARY TrackFile ;
REAL Pitch Radius := 1.2 0.4 ;
INTEGER Mesh := 5 ;
INTEGER MaxReg ;
PROCEDURE PWRPinCell ;
VolTrk Trackfile := PWRPinCell ::
   <<Pitch>> <<Radius>> <<Mesh>> >>MacReg<< ;
ECHO "MaxReg =" MaxReg ;
VolTrk Trackfile := DELETE: VolTrk Trackfile ;
END: ;
Sequence of module for a straightforward DRAGON execution:

\[
\text{GEO:} \rightarrow \text{EXCELT:} \rightarrow \text{ASM:} \rightarrow \text{FLU:} \rightarrow \text{EDI:}
\]

\[
\text{SHI:} \rightarrow \text{LIB:} \rightarrow \text{MAC:}
\]

(geometry) (tracking info) (CP matrices) (fluxes)

(macroscopic XS)

(microscopic XS)
Sequence of module for a burnup cycle in DRAGON:

(tracking info)  \[\rightarrow\] ASM: \[\rightarrow\] FLU: \[\rightarrow\] EVO:

(macroscopic XS)

(microscopic XS)

(CP matrices)  \[\rightarrow\] (fluxes)
Multiple level DRAGON execution with homogenization:

\[ \text{GEO:} \rightarrow \text{EXCELT:} \rightarrow \text{ASM:} \rightarrow \text{FLU:} \rightarrow \text{EDI:} \]

\[ \text{(geometry 1)} \quad \text{(tracking info)} \quad \text{(CP matrices)} \quad \text{(fluxes)} \quad \text{(macroscopic XS)} \]

\[ \text{GEO:} \rightarrow \text{EXCELT:} \rightarrow \text{ASM:} \rightarrow \text{FLU:} \rightarrow \text{EDI:} \]

\[ \text{(geometry 2)} \quad \text{(tracking info)} \quad \text{(CP matrices)} \quad \text{(fluxes)} \quad \text{(macroscopic XS)} \]
Geometry validation and results analysis using the \textbf{PSP}: module:

- \textbf{PSP}: generate PostScript representation of 2D geometry.

- Can use a \texttt{GEOMETRY} or \texttt{TRACKING} data structure for geometry description.

- Limited to simple 2-D valid EXCELT type geometry.

- Geometry can be colored by mixture number, region number of flux level (when \texttt{FLUXUNK} data structure is provided).
Tracking validation:

- Select quadrature densities and angular quadrature level.
- Verify volume and surface integration errors provided in output file (EDIT 2 in EXCEL: module).
DRAGON today and tomorrow

G. Marleau

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École Polytechnique de Montréal
Contents:

1. Advanced features.
2. Improvements in version 3.04 release R.
3. Limitations in DRAGON.
4. What to expect in version 3.05.
5. Towards DRAGON version 4.0.
Advance features

1. Cross-section perturbations.
2. Restart capabilities.
3. Double heterogeneity.
5. Generating flux map in DRAGON.
**Advance features**

Cross-section perturbations using the LIB: module:

- **MACROLIB** data structure defined using the LIB: module can be perturbed in various ways:
  - Change the concentration of a single isotope in a mixture.
  - Change the density of a mixture.
  - Change the concentration of each fuel isotope using a BURNUP data structure.
Example for burnup dependent fuel Doppler effect or CVR:

- Define a first MACROLIB with fuel at normal operating temperature (RefLib).
- Define a second MACROLIB with fuel at perturbed temperature (PerLib).
- Solve transport and depletion equation using RefLib and save burnup dependent isotopic concentrations ConcIso.
- Update PerLib at a specific time step using ConcIso.

```
PerLib := LIB: PerLib ConcIso :: BURN <<time>> ;
```
Advance features

Restart capabilities and MACROLIB combination using the MAC: module:

- The data structures generated by the various modules can all be saved (:=):
  - In a direct access XSM format for later use on the same system.
  - In a sequential ASCII format for later use on the same or a different system.

- The data structures can be exchanged between input files:
  - Create a MACROLIB during one execution.
  - Use this MACROLIB in various other execution.
Example of CANDU incremental cross section evaluation:

a) Create burnup dependent macroscopic cross section libraries:

- Perform a time dependent CANDU 2–D cell calculation.
- At selected time step (0, mid and full burnup), use EDI: to generate homogenized macroscopic cross sections and save.
- Export the EDITING structure in FuelEdi.
- Create using LIB: a MACROLIB that contain the structure materials and export in STRCMAC.
b) Create a MACROLIB for 3-D transport calculation:

- Import in RDmac using MAC: the homogenized fuel properties at a specific burnup form FuelEdi.
- Import in RDmac using MAC: the cross section associated with structure materials stored in STRCMAC.
- Export RDmac.
c) Perform 3-D transport calculation:

- Import RDmac.
- Perform a 3–D CANDU supercell calculation.
- Evaluate the incremental cross sections using the EDI: module.
Example:

RDmac := UTL: RDmac ::
    STEP UP ‘FUEL1’ ;
MacroLib := MAC: RDmac ::
    NMIX 17
    MIX  1 1 MIX  2 2
    MIX  3 3 MIX  4 4
;
RDmac := UTL: RDmac ::
    STEP DOWN STEP UP ‘Structure’ ;
MacroLib := MAC: MacroLib RDmac ::
    NMIX 17
    MIX 11 1 MIX 12 1 MIX 13 1 MIX 14 1
    MIX 15 1 MIX 16 1 MIX 17 1
;
Geometry with double heterogeneity in DRAGON:

- The macro geometry is defined in terms of homogeneous mixtures and mixtures with structure (micro geometry).
- The micro geometry may be spherical (grains or pellets) or cylindrical (CANDU clusters).
Advance features

Double heterogeneity option in DRAGON:

1. Available only via the SYBILT: module.

2. All the processing is performed in the ASM: module via the sub-module SYBILP:

3. Similar to the technique that A. Hébert implemented in the APOLLO-1 code as part of his Ph.D. work.
Double heterogeneity treatment in DRAGON:

1. Compute the CP matrix for each micro geometry contributing to a DH mixture.

2. Use these CP matrix to generate an equivalent cross section for the DH mixtures.

3. Compute the CP matrix for the macro geometry.

4. Distribute the macro-geometry CP contribution in DH region on micro geometry.
Non-uniform meshes in DRAGON:

\[
N = 19 \times 16 \times 16 = 4864 \text{ regions, } d_a = 100 \text{ Mb/groups for uniform mesh.}
\]

\[
N \approx 1500 \text{ regions, } d_a = 9 \text{ Mb/groups for non-uniform mesh.}
\]
Generation of non uniform mesh in DRAGON using the **MRG**: module:

1. Define a geometry `RefGeo` with a uniform mesh.

2. Track this geometry with **EXCEL**: and generate the `VOLTRK` data structure `RefVol` and tracking files `RefTrk`.

3. Use the **MRG**: module to pre-homogenize `RefVol` and `RefTrk` into `MrgVol` and `MrgTrk`.

4. Use the **ASM**: module to compute the CP based on `MrgVol` and `MrgTrk`.
Advance features

Using the MRG: module:

```plaintext
MrgVol MrgTrk := MRG: MrgVol MrgTrk ::
REGI  1  1  2  3  4  5  6  7  8
   9 10 11 12 13 14 15 16 17 17
   1  1  2 18 19 20 21 22 23
  24 25 26 27 28 29 30 16 17 17
  31 31 32 33 34 35 36 37 38
  39 40 41 42 43 44 45 46 47 47
  31 31 32 48 49 50 51 52 53
  54 55 56 57 58 59 60 46 47 47
  61 62 63 64 65 66 67 68 69
  70 71 72 73 74 75 76 77 78 78
...
;
PIJ := ASM: MacLib MrgVol MrgTrk ;
```
A colored map by mixture number of a DRAGON BWR geometry:
Advance features

The **PSP** module works for most 2–D Geometry that can be processed by the **EXCEL** module to:

- Produce a colored map by mixture number in the geometry.
- Produce a colored map by region number in the geometry.
- Produce a colored map for the multigroup flux.

For validating 3–D geometry, proceed by projection of the 3–D geometry on various 2–D planes.
Advance features

A colored map for the multigroup flux in a BWR geometry:

Legend
Color by Flux in group 1
1.13E+01+i*5.78E-02 ≤ Flux(i) ≤ 1.13E+01+i*5.78E-02

ANL DRAGON workshop
The **PSP**: module can process:

- An explicit DRAGON geometry. In this case the module **PSP**: calls the geometry analysis part of the **EXCEL**: module.

- A **VOLTRK** data structure produced by the **EXCEL**: module.

- A **VOLTRK** data structure produced by the **MRG**: module.
Advance features

A colored map by mixture number of a prismatic modular assembly (hexagonal cluster):
DRAGON version 3.04R

Updates in:
- GANLIB.
- UTILIB.
- DRAGON.

New modules:
- DRAGON.
Update to GABLIB/CLE-2000:

- String variable with blanks are now permitted as parameters of procedures.
- String variable length is now preserved.
- Mixed mode operations \texttt{I\_TO\_S} implemented.
Example of I_TO_S use:

```plaintext
STRING Refname := "Test" ;
STRING Dirname ;
INTEGER Step := 0 ;
WHILE Step 50 < DO
    EVALUATE Step := Step + 1 ;
    EVALUATE Dirname := Refname Step I_TO_S + ;
    ...
    ...
    EditDS := EDI: EditDS Flux TmpMicLib TrackingF ::
    COND 0.625 MERG COMP SAVE ON <<Dirname>> ;
ENDWHILE ;
```
Update to UTILIB:

- PSP drawing routines for strike and fill now combined with \texttt{gsave} to reduce dimensions of PostScript file.

- \texttt{XDRSET} and \texttt{XDRDBL} routine rationalized and generalized to other types of variable.

- \texttt{XDRCST} routine created to rationalized use of constants in DRAGON. (Avogadro number, neutron mass, etc.).
Update to DRAGON:
- The GEO: Module.
- The SYBILT: Module.
- The EXCELT: Module.
- The ASM: Module.
- The EXCELL: Module.
- The FLU: Module.
- The EDI: Module.
- The EVO: Module.
The \texttt{GEO}: module:

- Boundary condition \texttt{SSYM} is similar to \texttt{SYME} but applied at the boundary of the cell or assembly.

The \texttt{SYBILT}: module:

- Large symmetric slab problems produced CP that were not symmetric.
- Both forward and backward tracking directions are now considered.
The `EXCelm` module:

- The symmetry options can now be applied to isolated cells.
- The geometry analysis process is now controlled by a single routine and completely decoupled from the tracking process:
  - This was required to extend the capacity of the `PSP`, `EDI` modules.
  - It was also required to simplify the maintenance of the `EXCelm` and `EXCelle` modules.
For symmetric assemblies, the tracking is now automatically symmetrized:

- The symmetric tracking option was activated manually in 3.04B.
- Only some symmetries were considered.
- This lead to a reduction of a factor of 2 to 4 in size for some tracking file.

In addition:

- Tracks were being lost on Linux for 2-D and 3-D hexagonal assembly.
- Volume computation errors were found for some cluster geometries.
The **ASM**: module:

- Space requirements for the execution of this module is now provided.

- It is now possible to compute the CP matrix while reading the tracking file only once.

The **EXCELL**: module:

- The CP normalization technique did not work for explicitly voided regions.

- The renormalization technique were reprogrammed from the original routine in **ASM**: and the results cross checked.
The **FLU**: module:

- Adjoint and generalized adjoint flux algorithm is now available.
- For imposed leakage on can select imposed $dB^2$ or imposed $B^2$ with computed $d$.

The **EDI**: module:

- Homogenization per geometry option is now implemented.
- The **MACROLIB** and **MICROLIB** generated by the **EDI**: module are now complete and can be used for burnup calculation in **DRAGON**.
The **EVO**: module:

- The contents of the **BURNUP** data structure was incomplete.

- The isotopes names are now stored on the **BURNUP** data structure (required by the **LIB**: module for some executions).
New modules in DRAGON:

- The SAD: Module.
- The PER: Module.
- The HST: Module.
The **SAD**: module:

- Compute the generalized adjoint sources required for generalized perturbation theory methods.

The **PER**: module:

- Perform generalized perturbation theory calculations.
The \texttt{HST} module:

- For history based diffusion/transport reactor calculations.
- It generates a \texttt{HISTORY} data structure:
  - It is a stand-alone structure.
  - It must be stored on a linked list or an XSM file.
  - It contains the local and global parameter associated with a series of fuel cells including their current isotopic contents.
Limitations in DRAGON

1. Geometry.
2. Memory.
3. Other.
Limitations in DRAGON

2-D Cartesian assembly in DRAGON:

- Cartesian mesh in each direction must extend to the whole geometry.
- A single cylinder per cell.
- Cylinders are by default centered in the cell but can be displaced using the OFFCENTER option.
- The Cartesian and annular mesh are superposed. Mixture must therefore be specified in $R$, then in $x, y$.
- Mixtures are specified even in location that do not exists.
Limitations in DRAGON

3-D Cartesian assembly in DRAGON:

- Similar to 2–D assembly.
- Cylinders in different directions cannot intersect.
- Cylinders must extend to the whole geometry.
- Cylinders are by default centered in the cell. They can be displaced using the OFFCENTER option.
Limitations in DRAGON

A Simple 3–D Cell:
Limitations in DRAGON

Mixture specification for simple 3–D cell:

```
TMPGEO := GEO: :: CAR3D 1 1 1
CELL FC1B
  X- REFL X+ REFL Y- REFL Y+ REFL Z- REFL Z+ REFL
::: FC1B := GEO: CARCELZ 2 3 2 1
  MESHX <<MXLP>> <<MXL>> <<C14CN>> <<MXYD>>
  MESHY <<MYLP>> <<CALLCN>> <<PYLP>>
  MESHZ 0.0 49.5
  OFFCENTER <<FC1XD>> <<FC1YD>>
  RADIUS 0.0 <<RF2>> <<RCT>>
  MIX 1 2 3
  11 12 13
  21 22 23
  31 32 33
  41 42 43
  51 52 53 ; ;
```

Note: mixtures 1, 2, 31 and 32 not used.
Limitations in DRAGON

2–D Hexagonal assembly in DRAGON:

- Either a single hexagonal cell or an assembly of hexagonal cell.
- Global boundaries of hexagonal assembly are jagged:
Limitations in DRAGON

2-D hexagonal assembly in DRAGON (continue):

- A single cylinder per cell.
- Cylinders are centered in the cell (no possibility to displace the center in this case).
- Sectors are possible (but not documented).

3-D hexagonal assembly in DRAGON:

- Similar to 2-D assembly.
- Cylinder completely fill the $z$ direction.
- Sectors are possible (but not documented).
Limitations in DRAGON

Cluster in DRAGON:

- Not available in 3–D.
- A single cell permitted in 2–D.
- Pin clusters must be located inside the outer annular ring.
- Outer boundary can be Cartesian, hexagonal or annular.
Limitations in DRAGON

Exact boundary conditions:

- VOID: applied at the explicit boundary of the cell or assembly.
- SYME: applied at the center of the cells closest to the explicit assembly boundary specified.
- DIAG: applied at the center of the cells closest to the explicit assembly boundary specified.
- SSYM: applied at the explicit boundary of the cell or assembly.
Approximate boundary conditions:

- **REFL**: applied at the explicit boundary of the cell or assembly. Mirror like reflections not available in 3–D.

- **TRAN**: applied at the explicit boundary of the cell or assembly. Mirror like translations not available in 3–D.

- **ALBE**: applied at the explicit boundary of the cell or assembly.
Limitations in DRAGON

Region identification for cells and assemblies:

Back

<table>
<thead>
<tr>
<th>6</th>
<th>9</th>
<th>10</th>
</tr>
</thead>
<tbody>
<tr>
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Front

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<tr>
<td>11</td>
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<td>15</td>
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Back

<table>
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<th>10</th>
<th>18</th>
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<tr>
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Front

<table>
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<tr>
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<th>16</th>
<th>20</th>
</tr>
</thead>
<tbody>
<tr>
<td>3</td>
<td>14</td>
<td>19</td>
</tr>
</tbody>
</table>

one cell

three cells assembly
Region identification for single cell:
- radially outward in a cell.
- from lower to upper $x$ location in a cell.
- from lower to upper $y$ location in a cell.
- from lower to upper $z$ location in a cell.
## Limitations in DRAGON

### Region identification for one cell problem:

<table>
<thead>
<tr>
<th>PLANE - 2</th>
<th></th>
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</thead>
<tbody>
<tr>
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<td>16 (1) 19 (1) 20</td>
</tr>
<tr>
<td>ABSENT (1) 18 ABSENT</td>
<td></td>
</tr>
<tr>
<td>ABSENT (1) 17 ABSENT</td>
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</table>

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</thead>
<tbody>
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<tr>
<td>ABSENT (1) 3 ABSENT</td>
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<tr>
<td>ABSENT (1) 13 ABSENT</td>
<td></td>
</tr>
<tr>
<td>ABSENT (1) 12 ABSENT</td>
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</table>

<table>
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</thead>
<tbody>
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<td>6 (1) 9 (1) 10</td>
</tr>
<tr>
<td>ABSENT (1) 8 ABSENT</td>
<td></td>
</tr>
<tr>
<td>ABSENT (1) 7 ABSENT</td>
<td></td>
</tr>
</tbody>
</table>
Region identification for assembly of cells:

- Inside each cell as above.
- From lower to upper $x$ cell location in the assembly.
- From lower to upper $y$ cell location in the assembly.
- From lower to upper $z$ cell location in the assembly.
Region identification for three cells assembly:

<table>
<thead>
<tr>
<th>Plane - 2</th>
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<th>(1) 20</th>
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<table>
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<th>(1) 13</th>
<th>(1) 19</th>
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<th>(1) 10</th>
<th>(1) 18</th>
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<tbody>
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<td>(1) 9</td>
<td>ABSENT</td>
<td></td>
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<tr>
<td>ABSENT</td>
<td>(1) 8</td>
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</table>

<table>
<thead>
<tr>
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<th>(1) 1</th>
<th>(1) 7</th>
<th>(1) 17</th>
</tr>
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<tbody>
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<td>ABSENT</td>
<td></td>
</tr>
<tr>
<td>ABSENT</td>
<td>(1) 5</td>
<td>ABSENT</td>
<td></td>
</tr>
</tbody>
</table>
Example of a simple 3–D problem.

- Total volume 1 liter \((V = 10^3 \text{ cm}^3)\).
- Central fissile region is red \((V = 27 \text{ cm}^3)\).
- Strong absorber is green \((V = 1 \text{ cm}^3)\).
Limitations in DRAGON

Example of mesh for simple 3–D problem:

Note: $x - y$ and $x - z$ planes are identical.
Limitations in DRAGON

Select transport calculation mesh:

- Use symmetry whenever possible.
- For regions with strong absorber and fission:
  - Try to avoid using approximate boundary conditions.
  - Use fine mesh discretization.
- For moderator region:
  - Fine to coarse mesh discretization may be required.
Select CP quadrature:

- As many angles as possible.
- Tracking density must be as dense as possible.
- Integration lines are renormalized using ratio of approximate to exact volumes.
Limitations in DRAGON

Storage requirements for tracking file:

- Size of tracking file linear in $N$:
  \[ d_t \propto \rho h_+^2 N \]

  for $h_+ = 50$ cm, $\rho = 20$ t/cm$^2$ and $N = 1000$ regions:
  \[ d_t = 600 \text{ Mb.} \]

- Tracking file limited to disk storage available.

- Use **EXCELL** module when disk storage is insufficient for 3-D (2-D not permitted) calculations since tracking file is not saved in this case.
Storage requirements for CP matrices:

- Number of elements per groups is $N^2$.
- Memory space required for execution is about $5N^2$.
- Total disk space required for storage of $G$ group CP is $GN^2$.
- For $G = 89$ groups and $N = 1000$ regions: $d_a = 356$ Mb.
- Store ASMPIJ data structure in XSM_FILE rather than LINKED_LIST.
Collision probability normalization:

- Collision, leakage and escape probability computed independently.
- Integration procedure preserves symmetry relations.
- Conservation relations generally not satisfied. For example:

\[
\sum_{\alpha=1}^{N_{\alpha}} \frac{S_{\alpha}}{4} p_{\alpha j} + \sum_{i=1}^{N_{i}} \Sigma_{i} V_{i} p_{ij} - V_{j} \neq 0
\]

- Conservation relations restored using HELIOS type normalization (default option).
Expectations for DRAGON 3.05

Major update to the DRAGON geometries:

- The **NXT**: tracking module will contain the new **EXCELT**: tracking option.
- It will be able to analyze more complex geometries in 2–D and 3–D including clusters.
What we have now for 2–D Cartesian assemblies:

Legend
Color by Region
1 2 3 4 5 6 7 8 9 10 11 12 13 14 15 16 17 18
Expectations for DRAGON 3.05

What we have now for 3–D Cartesian assemblies:
What we have now for 2–D cluster cells:
What we are in the process of implementing for 2–D assemblies with Cartesian boundaries:

This will also be available in 3–D geometry.
What we are in the process of implementing for 2–D assemblies with annular boundaries:

This will also be available in 3–D geometry.
What we are in the process of implementing for 2-D assemblies with clusters.

Legend
Color by Mixture

(a) Cluster with pin rotation (pins located using RPIN)
(b) Cluster without pin rotation (pins located using CPIN)
What we are in the process of implementing for 3–D assemblies with clusters.
Towards DRAGON version 4.0

2. Self-shielding options.
3. Method of characteristics.
Towards DRAGON version 4.0

Improvements in microscopic cross section library analysis:

- Upgrade of the decay chains treatment to take into account more reactions and to include energy production from other sources than fission (mainly LIB and EVO modules).

- Update the EDI module to process delayed neutron data.

- Upgrade the DRAGLIB format and generate of a DRAGLIB file from ENDF/B-VI using the DRAGR module of NJOY (mainly LIB:).

- Upgrade the LIB: module in such a way that it can process subgroup data for self-shielding.
Towards DRAGON version 4.0

Improvements in the self-shielding options:

- Upgrade the $\text{SHI} :$ module to permit distributed self-shielding (Nordheim model).
- Develop the $\text{SUBG} :$ module for self-shielding by subgroup.
Method of characteristics:

- Four versions currently in competition in developmental version of DRAGON:
  1. MOCC method for cyclic tracking (available in 2–D only for mirror like BC)
  2. G. Wu and R. Roy method of characteristics developed for 2–D and 3–D tracking with white BC.
  4. R. LeTellier and A. Hébert method of characteristics (available in 2–D only for white BC).

- Rationalize and implement a unique version.