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ANALYTICAL METHODS FOR DISCRIMINATING STARDUST IN AEROGEL CAPTURE MEDIA

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ABSTRACT

Comet 81P/Wild 2’s serendipitous orbit change to the inner solar system in 1974 offered researchers a rare opportunity to sample cometary material from the Kuiper belt, a repository of material left over from solar system formation ~4.6 Gyr ago. NASA’s Stardust mission intercepted the comet in January 2004 and returned with material collected from its tail in January 2006. The cometary material, consisting of particles ranging from 10 microns down to <2 nm, was collected in aerogel, a very low density (~3 mg/cm³) silica foam, to minimize the effects of deceleration from 6.1 km/s. The entire deceleration track is extracted from the aerogel block as a pyramidal shape known as a keystone which can be mapped using x-ray fluorescence prior to extraction of terminal or intermediate particles for other analyses. One goal of the track mapping is to determine the bulk composition of the cometary material returned. Unfortunately, although the aerogel is predominantly SiO₂, there are sufficient quantities of trace elements similar to those expected in the cometary material to require sophisticated discrimination techniques in order to decide whether a fluorescence map pixel contains only aerogel or both aerogel and cometary material. We have developed a dual threshold analysis approach for better distinguishing cometary material from aerogel contaminants and have applied it to five Stardust impact tracks and terminal particles. Here, we present aspects of the dual threshold approach and demonstrate its impact on track composition for one track.

INTRODUCTION

Stardust, a NASA Discovery class mission, returned to Earth in January, 2006 with dust captured from the tail of Comet 81P/Wild 2 [1,2]. Microscopic dust particles were captured at a relative velocity of 6.1 km/s in low-density silica aerogel tiles and adjacent aluminum foils [3]. These Stardust samples are unique in our collections of extraterrestrial materials: They are the first samples returned from a known parent body that originated in the Kuiper belt beyond the gas giants. Prior to 1974, Wild 2 orbited at heliocentric distances beyond Jupiter, but a gravitational encounter with that planet sent it into a short-period orbit in the inner solar system and allowed the Stardust spacecraft to intercept the comet’s coma within the orbit of Mars. Since Comet Wild 2 has only been in its new orbit for 6 revolutions it is likely to be little-altered and a good representative
of the material in the Kuiper belt. It is expected that most small bodies in the Kuiper belt have experienced little thermal or aqueous processing since formation and are therefore relatively unaltered since the solar system formed \(~4.6\) billion years ago [4].

The focus of the synchrotron work is the elemental composition of the Stardust particles using x-ray fluorescence. In parallel with our studies are studies by other teams using a wide variety of instrumentation including x-ray micro-diffraction, electron microscopy, infra-red spectroscopy, organic abundance, and isotopic analysis. In addition to the efforts at the Stanford facility research teams at 5 other synchrotron facilities around the world were also involved in the preliminary examination team. The results of this examination were published in a series of papers in a special issue of Science [2,5,6,7,8,9,10]. Since presenting the results of our preliminary examination, we have attempted to refine our methods for the challenging problem of properly identifying which pixels in our fluorescence maps truly contain cometary material and which only contain aerogel. As will be discussed below, this is complicated by the presence in the aerogel of elements other than silicon and oxygen which can mimic the cometary material. This refined analysis has also been submitted for publication [11]. The focus of this work is to discuss in more detail the techniques used to distinguish cometary material from the surrounding background of aerogel plus contaminants. While specific to the issue of Wild 2 sample return analysis, the methods here are generally applicable to those situations where a fluorescence signal of interest is intermixed with known background signals.

**EXPERIMENTAL**

We studied five Stardust deceleration tracks during Stardust Preliminary Examination. All were extracted from the same aerogel tile, Cell 44, in the form of individual aerogel “keystones” [12], shown as micrographs in Fig. 1 at approximately similar scale. Each keystone contains a single deceleration track, a cavity in the silica aerogel formed by the hypervelocity capture of a piece of cometary material (impactor). Some impactors fell apart on impact, generating multiple sub-tracks from a single entry hole, and some remained relatively intact. For both types, measurable cometary material is present along the track. The track lengths vary from approximately 0.3 to 3.3 mm and from conical, “carrot” shapes to more bulbous shapes. Details of the tracks can be found in the supplemental online material of [5] and in [11].

Tracks were mapped using a hard x-ray scanning fluorescence microprobe on the wiggler Beam Line 6-2 at the Stanford Synchrotron Radiation Laboratory. The microprobe uses Kirkpatrick-Baez final focusing optics and adjustable virtual source slits for a beam size of 2 \(\times\) 2 microns\(^2\) with \(10^9\) photons/second (ph/s). For efficient mapping of the deceleration track, the focused spot size was increased to 6 \(\times\) 15 microns\(^2\) with \(2 \times 10^{10}\) ph/s. Si(111) monochromator crystals selected 14 keV incident x-rays for easy access to the K absorption edges of elements from Si through Br. The microprobe is equipped with a Leica optical microscope and aluminum-coated low-Z mirror for x-ray line-of-sight viewing and positioning of the sample in the x-ray beam (Fig. 2).
Keystones were translated in the beam in steps matched to the FWHM of the focused spot size and full fluorescence spectra were collected with dwell times of 30 seconds/pixel. Terminal particles and other particles along tracks were located primarily by high Fe fluorescence and scatter count rates. An ultra-clean Si(Li) detector with ~150 eV resolution (Mn Kα) was used to collect the fluorescent x-rays in a geometry perpendicular to the incident beam in the plane of the storage ring. An ion chamber before the sample and a PIN diode after were used to monitor the incident beam intensity. A helium shower covering the x-ray – sample interaction region reduces air absorption, ozone generation and Ar fluorescence. Recent improvements in that shower have essentially eliminated argon fluorescence from the spectra. (Fig. 2).

Spectra were fit using PyMca, a program developed by Dr. Armando Sole at the European Synchrotron Radiation Facility (ESRF) and available for public use [13]. Whole track masses were derived from the difference between the sum of the on-track spectra containing cometary material and the normalized aerogel matrix background spectrum. The details of how this separation is achieved are discussed below. Two reference standards were employed in quantifying element masses: a thin (200 nm) Fe
film acted as an absolute reference standard, and a USGS basaltic glass microprobe standard NKT-1G thin section (200 nm) was used to determine the energy-dependent correction for elements below Fe [14]. The NKT-1G basaltic glass standard contains ~9.5 wt% Fe, so is especially useful for measuring self-absorption of low-Z fluorescence x-rays (for example, sulfur) and corrects for detector response. What this standard does not address is the variation in self-absorption due to cometary particle size. Without specific knowledge of the cometary particle, the self-absorption can only be approximated. Thus, self-absorption is small for elements above Ca, but considerably more uncertain for lower-Z elements. The error in Fe mass is dominated by uncertainty in the thickness of the Fe standard and is at most a few percent.

ANALYSIS

One of the simplest and most powerful techniques is to use a region of interest (ROI) technique to identify where certain elements are present. One such map is presented in Fig. 3, showing the concentration of Fe along the deceleration track, with an optical image of the same track above it. The horizontal scale for the two images is the same, but the vertical scale of the fluorescence map is enlarged to more easily see the variability along the track. The warmth of the color in the map is related to the intensity of the Fe signal from that pixel; the data are scaled as the log of the intensity so that range of the map is not dominated by the very intense terminal particle. The distinction between the midnight and royal blue pixels on the periphery of the track are discussed below. There is strong evidence that Fe is deposited throughout the deceleration track, although in this case the majority of the Fe mass is at the terminal particle. There are other deceleration tracks where the majority of the Fe mass is distributed along the track, leaving only a small part of the total at the terminal particle.

![Optical Microscope Image of Track 12](https://via.placeholder.com/150)

The midnight blue is for aerogel only pixels, the royal blue is for the in-between pixels and warmer colors are where cometary material is found.

There are several challenges in calculating a total cometary mass in a deceleration track, most of them fundamentally related to the nature of the aerogel capture medium. Aerogel is a silica “foam”. This fundamentally limits our ability to measure the presence of silicates in the cometary material. Where one would normally use the Si fluorescence signal as a reference for comparison to other kinds of cometary or meteoritic material, in this case, the Si fluorescence is dominated by the surrounding aerogel, especially as it has been shown that there is compressed aerogel surrounding the terminal particle. Typical aerogel keystone thicknesses are 100-500 microns, so silicates in the cometary material
are completely masked by the surrounding aerogel. Analysis is further complicated by the presence of contamination within the aerogel, including Fe, Zn, Cu, Cr, Ti; many elements of interest in the comet are also potentially present in the aerogel. The surrounding aerogel is also effective in masking Mg and Al fluorescence that may be emanating from the cometary material, making those elements hard to quantify. Self-absorption of low-Z elements (< Ca, e.g.) increases the error bars on estimates of those masses.

We have tried several methods for identifying pixels that include cometary material. In the end, we have chosen the Fe fluorescence concentration in each pixel as the best metric. Another method, using the total scattered (elastic) signal from a pixel, was tested and found to be similar to but less reliable than using iron. We have chosen a dual threshold method: If the Fe concentration in a pixel was greater than some threshold, then it was considered to contain both cometary material and aerogel, if less than some lower threshold, than it contains only aerogel, and if the Fe concentration is between those two thresholds, it is discarded. The resulting spectra can be seen in Fig. 4, with the

![Graph](image)

**Fig. 4.** Summed fluorescence spectra from the track map of track 12. The red curve is from those pixels which include both cometary and aerogel, the blue curve is from those that only contain aerogel, and the black curve is from those where we are uncertain. The blue curve has been scaled to the red curve, but the black curve is unscaled.

spectra plotted as the log of the fluorescence intensity vs. the photon energy in eV. The red curve is the sum of those spectra thought to contain both cometary material and aerogel, the blue curve is the sum of those pixels which only contain aerogel, and the black curve the spectra for which we are undecided. Note that the blue curve has been
normalized to have the same Si and scatter peak amplitude as the red curve, but the black curve has not been scaled. To give a sense of the scaling factors, the red curve represents almost 29000 seconds of collected data, the background just under 12000 seconds of data and the undecided spectrum is 5500 seconds of data. To see where the undecided pixels are on the map, compare the midnight-blue areas of Fig. 3 to the royal-blue areas. It is immediately clear that the undecided pixels are where one would expect to see them: on the periphery of the track, where the transition occurs from clearly cometary plus background to clearly background.

The reason it is important not to include these undecided spectra in either of the other two is because the cometary material is calculated by subtracting the results of analyzing the background spectrum from the results of analyzing the cometary spectrum. If cometary material is mistakenly included in the background spectrum, then one loses twice: cometary material missing from that spectrum is then subtracted from the remaining cometary material. If one is too free in including pixels without any cometary material into that spectrum, then those elements with only trace concentrations will be effectively buried by the statistical noise of the additional pixels with only aerogel in them. The challenge is to select the thresholds so that one is sure that the background spectrum does not contain any cometary material, and that as little material is in the undecided spectrum as possible.
To explore the effect of having this undecided bin for some of the pixels, consider Fig. 5, where instead of separating the undecided pixels, they are instead added to the background spectrum, effectively going to a single threshold method. One can see that, although the changes are small, there clearly are differences between the blue and magenta spectra. This is because the scatter/Fe ratio is different for the two, with more Fe, Ni, Zn, etc present in the magenta (relative to the scatter peak) than in the true background spectrum. In practice, the lower threshold is determined by watching metrics such as the Fe/scatter ratio as the threshold is changed, picking as the appropriate level that point where there is a significant change in that ratio. To more clearly see that the Fe and Ni (and especially Zn) ratios to scatter peak are different between the background and undecided spectra, Fig. 6 shows the cometary, background and undecided spectra, now all scaled so that they have the same Si and scatter peak amplitudes.

![Graph](image)

In this figure the very different ratios for Fe-scatter and especially Zn-scatter are clear. There is significantly more cometary material in these pixels than in the surrounding background pixels. Adding these pixels only has a small effect on the calculated Fe mass, but a very large effect on the lower-mass elements such as Zn. (Remember, these figures show the log of the fluorescence intensity). The results of analyzing these two approaches can be shown in Table 1, which compares the calculated masses of Track 12 using the two methods.
Table 1. Element abundances by mass for Stardust impact Track 12 by single and dual threshold analysis.

<table>
<thead>
<tr>
<th></th>
<th>single threshold mass (g)</th>
<th>dual threshold mass (g)</th>
<th>fractional change</th>
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</thead>
<tbody>
<tr>
<td>Cl</td>
<td>3.05E-11</td>
<td>5.78E-11</td>
<td>1.89</td>
</tr>
<tr>
<td>K</td>
<td>6.19E-12</td>
<td>1.27E-11</td>
<td>2.04</td>
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<tr>
<td>Ca</td>
<td>2.07E-13</td>
<td></td>
<td></td>
</tr>
<tr>
<td>Cr</td>
<td>5.92E-12</td>
<td>8.51E-12</td>
<td>1.44</td>
</tr>
<tr>
<td>Mn</td>
<td>4.37E-14</td>
<td>6.17E-14</td>
<td>1.41</td>
</tr>
<tr>
<td>Fe</td>
<td>2.89E-10</td>
<td>2.98E-10</td>
<td>1.03</td>
</tr>
<tr>
<td>Ni</td>
<td>2.48E-11</td>
<td>2.58E-11</td>
<td>1.04</td>
</tr>
<tr>
<td>Cu</td>
<td>1.52E-13</td>
<td>3.10E-13</td>
<td>2.04</td>
</tr>
<tr>
<td>Zn</td>
<td>2.15E-13</td>
<td>8.83E-13</td>
<td>4.11</td>
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<tr>
<td>Ga</td>
<td>2.12E-14</td>
<td></td>
<td></td>
</tr>
<tr>
<td>As</td>
<td>1.39E-13</td>
<td>2.48E-13</td>
<td>1.78</td>
</tr>
<tr>
<td>Se</td>
<td>2.09E-13</td>
<td>2.05E-13</td>
<td>0.98</td>
</tr>
<tr>
<td>total</td>
<td>3.57E-10</td>
<td>4.04E-10</td>
<td>1.13</td>
</tr>
</tbody>
</table>

* Measured masses above minimum detection limits are reported here.

While the total mass only changes marginally, the changes in some elements (notably Zn, has been significant (>4). The total mass is dominated by the 3% increase in the Fe mass, a small but significant change. One curious outcome of this analysis is that the sulfur signal does not rise above that measured in the background in either single or dual threshold analyses. Since collecting these data we have become suspicious that volatile elements may have a larger distribution perpendicular to the track than would be expected from looking at the Fe distribution. This conjecture will be tested in the near future.

CONCLUSIONS
The analysis of Stardust deceleration tracks has been an exciting project for our team and has raised a set of questions which are potentially common to a variety of fluorescence mapping projects. If there is an element of interest distributed in a matrix which includes low levels of contamination of that same element, how does one successfully isolate the material of interest? We have presented one such solution here, where one accepts that a limited number of pixels are “too close to call” and the consequences of guessing wrong can be significant, especially for trace elements which have not been used to make the segregation between signal and background. One can create guidelines of how one chooses to set the two thresholds, but we have yet to identify a rigorous method for their selection.

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REFERENCES


