

Composition of the Earth's Inner Core from High-pressure Sound Velocity Measurements in Fe-Ni-Si alloys

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14 Abstract

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16 We performed room-temperature sound velocity and density measurements on a 17 polycrystalline alloy, Fe<sub>0.89</sub>Ni<sub>0.04</sub>Si<sub>0.07</sub>, in the hcp phase up to 108 GPa. Over the 18 investigated pressure range the aggregate compressional sound velocity is ~ 9% higher 19 than in pure iron at the same density. The measured aggregate compressional (V<sub>P</sub>) and 20 shear (V<sub>S</sub>) sound velocities, extrapolated to core densities and corrected for anharmonic 21 temperature effects, are compared with seismic profiles. Our results provide constraints 22 on the silicon abundance in the core, suggesting a model that simultaneously matches the 23 primary seismic observables, density, P-wave and S-wave velocities, for an inner core 24 containing 4 to 5 wt% of Ni and 1 to 2 wt% of Si. 25

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29 The study of seismic wave propagation and normal mode oscillation are, without doubt, 30 two of the most direct probes of the Earth's interior, providing a remote sensing method to 31 obtain sound velocities and density. However, to derive an accurate compositional model, 32 these seismic observations have to be combined with laboratory-based experiments 33 constraining the density and elastic properties of highly compressed minerals. Based on shock 34 wave measurements, Birch proposed in the early 1950's that the Earth's core was mainly 35 composed of iron alloyed with nickel and some "light element(s)", required to account for the 36 observed density difference between the core and pure iron [Birch, 1952]. Since then, a large 37 number of experimental and theoretical studies addressed the nature and concentration of 38 these light elements (see review by Poirier [Poirier, 1994]), with several elements proposed. 39 Geo- and cosmochemical arguments suggest that most likely candidates are sulfur [Williams 40 and Jeanloz, 1990; Sherman, 1991; Sherman, 1995; Sherman, 1997; Li et al., 2001; Vočadlo, 41 2007; Morard et al., 2008; Côté et al., 2008a], oxygen [Sherman, 1991; Sherman, 1995; 42 Stixrude et al., 1997; Rubie et al., 2004; Badro et al., 2007; Asahara et al., 2007; Corgne et al., 43 2009], and silicon [Sherman, 1997; Dobson et al., 2003; Vočadlo, 2007; Badro et al., 2007; 44 Côté et al., 2008a; Côté et al., 2008b; Asanuma et al., 2008]. Very recently, high-pressure 45 measurements of sound velocity and density of several iron compounds (FeO, FeSi, FeS, and 46 FeS<sub>2</sub>) in conjunction with data on pure iron, pointed out inconsistencies when considering 47 sulfur as the only major light element, and proposed an Earth's inner core composed of iron 48 alloyed with silicon (2.3 wt%) and traces oxygen (0.1 wt%) [Badro et al., 2007]. However, 49 these results are based on three primary assumptions:



a) a linear dependence of compressional sound velocity (V<sub>P</sub>) on density ("Birch's law");

b) the P-wave velocity of the alloy is equal to a geometrical average of that of the minor
compound and the metal (*i.e.* ideal mixing behaviour);

53 c) the inclusion of up to 5 wt% Ni has a negligible effect on sound velocity.

54 Further, a clear limitation of this model is that only the aggregate compressional sound 55 velocity was taken into account, whereas the largest discrepancy between the seismological 56 observations and the results from molecular dynamics simulations and diamond anvil 57 experiments is observed for the shear wave velocity [Deuss, 2008].

58 Therefore, in order to validate the overall proposed model of the core composition, as 59 well as to address the legitimacy and the possible limits of the adopted approximations, we 60 undertook an experimental investigation of (Fe,Ni,Si) alloy samples. Specifically, we carried 61 out sound velocity and density measurements on polycrystalline samples containing 4.3 wt% 62 of Ni and 3.7 wt% of Si, compressed in diamond anvil cell (DAC) to megabar pressures. 63 Details of sample synthesis and characterization, as well as of the inelastic x-ray scattering 64 (IXS) measurements are reported in the next section, while the obtained results are illustrated 65 and discussed in section 3. Our main conclusions are summarized in section 4.

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67 **2. Experimental details** 

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Polycrystalline homogeneous samples of silicon bearing iron-nickel alloy have been prepared at high pressure and high temperature. Partial oxidation of silicon metal is a critical issue in the synthesis of silicon-rich iron alloy. The experiment has been conducted at superliquidus conditions to segregate from a SiO<sub>2</sub> glass a (Fe,Ni,Si) metallic blob free of SiO<sub>2</sub> inclusions. The starting material consisted in homogenized mixture of high purity metallic and oxide powders of Fe, Si, Ni and SiO<sub>2</sub>, with a 60 wt. % metallic portion relative to oxide. Piston cylinder experiment was carried out at 10 kbars and 1850°C at the Lawrence Livermore National Laboratory, using a standard 1/2" BaCO<sub>3</sub> pressure cell assembly, with a graphite furnace and a MgO capsule. The recovered metallic blob of about 1 mm diameter was analyzed with an electron probe micro-analyzer operating at 20 kV and 50 nA. Multiple analyses as well as backscattered electron images show homogenous (Fe,Ni,Si) alloy composition without quench textures at least at the scale of imaging and analytical resolution. Silicon and nickel concentrations are 3.7 wt% and 4.3 wt%, respectively.

Compacted pellets of about 90  $\mu$ m diameter and 20  $\mu$ m thick were loaded in DACs equipped with Re gasket, using 300  $\mu$ m flat anvils and neon as pressure transmitting medium for measurements up to 50 GPa, and 150/300  $\mu$ m beveled anvils with no pressure transmitting medium for higher pressures. Pressures were determined by ruby fluorescence and, most importantly, the densities ( $\rho$ ) were directly obtained from diffraction measurements and crosschecked with the equation of state previously measured on the same samples [Fiquet et al., 2008].

89 Inelastic x-ray scattering measurements were carried out on the ID28 beamline at the 90 European Synchrotron Radiation Facility, using the Si(8,8,8) instrument configuration, which 91 provides the best compromise between flux and energy resolution (5.5 meV full width half 92 maximum (FWHM)) for polycrystalline samples compressed in DAC. Spectra have been 93 collected in transmission geometry, with the x-ray beam impinging on the sample through the 94 diamonds, along the main compressional axis of the cell, and hence probing exchange 95 momenta q almost perpendicular to the cell-axis. The transverse dimensions of the focused xray beam of 30 x 90  $\mu$ m<sup>2</sup> (horizontal X vertical, FWHM) were further reduced by slits on the 96 vertical direction. Momentum resolution was set to 0.25 nm<sup>-1</sup>. Further details of the 97 98 experimental setup can be found elsewhere [Krisch, 2003; Antonangeli et al., 2004a; 99 Antonangeli et al., 2005]. By scanning the scattering angle at the elastic energy (*i.e.* q-scan at 100  $\Delta E=0$ ) we collected the [100], [002] and [101] reflections out of our sample, with the 101 momentum resolution of 0.06 nm<sup>-1</sup> set by slits in front of the analyzer.

102 We collected data in the hcp phase, at 27, 37 and 47 GPa on quasi-hydrostatically 103 compressed samples, and at 32, 68 and 108 GPa on non-hydrostatically compressed samples. 104 At each investigated pressure point, we mapped the longitudinal acoustic phonon dispersion throughout the entire first Brillouin zone collecting 8-9 spectra in the range 3.5-12 nm<sup>-1</sup> 105 106 (Figure 1). The energy positions of the phonons were extracted by fitting a set of Lorentzian 107 functions convolved with the experimental resolution function to the IXS spectra, utilizing a standard  $\chi^2$  minimization routine. We then derived the aggregate compressional sound 108 velocity V<sub>P</sub> from a sine fit (Born-von Karman lattice-dynamics theory limited to 1<sup>st</sup> neighbor 109 110 interaction) to the phonon dispersion [Antonangeli et al., 2004a; Antonangeli et al., 2005], 111 with error bars between  $\pm 2$  and  $\pm 3\%$  (Figure 1). Combining our measurements of V<sub>P</sub> and  $\rho$ 112 with the values of bulk modulus, K [Figuet et al., 2008], we also obtained the aggregate shear sound velocities  $V_S$  from the relation  $V_S = [3/4 (V_P^2 - K/\rho)]^{1/2}$ , although with larger uncertainty 113 114  $(\pm 4-6\%)$  due to error propagation.

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- 116 **3. Results and discussion**
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The measured compressional sound velocities are plotted as a function of density in Figure 2, together with values for pure iron [Fiquet et al., 2001; Antonangeli et al., 2004a] and Fe<sub>0.78</sub>Ni<sub>0.22</sub> alloy [Kantor et al., 2007]. To provide an analysis unbiased by systematic differences resulting from different techniques, pressure scales or approximations in the equation of state, and thus to be able to resolve variations as low as a few percent, we have only considered data obtained by IXS for conditions where the density was directly measured. While no systematic offsets can be observed between data for pure iron and iron-nickel alloy, our velocity measurements for  $Fe_{0.89}Ni_{0.04}Si_{0.07}$  are systematically higher, as highlighted by the linear fits to the experimental data (Figure 2). Over the investigated pressure range,  $Fe_{0.89}Ni_{0.04}Si_{0.07}$  is approximately 9% faster than pure iron at the same density. We also note that our experimental results compare favourably with calculations on  $Fe_{0.9375}Si_{0.0625}$ [Tsuchiya and Fujibuchi, 2009] (Figure 2), further stressing that the increase in the sound velocity is solely due to the silicon incorporation with no observable effect due to nickel.

131 To compare directly with seismic models however, our results need to be extrapolated to 132 core conditions. Within the experimental uncertainties, all the datasets exhibit a linear 133 dependence of V<sub>P</sub> with density (*i.e.* follow "Birch's law"), as also observed in several other 134 high-pressure experimental and theoretical studies [Vočadlo, 2007; Badro et al., 2007; Kantor 135 et al., 2007; Antonangeli et al., 2008; Tsuchiya et Fujibuchi, 2009]. Hence, as a first 136 approximation (see discussion later on), within a quasi-harmonic limit, we assume a linear 137 dependence of the compressional sound velocities with density, irrespective of the specific 138 pressure and temperature conditions.

139 Also, since we are interested in the isotropic aggregate properties, we carefully checked 140 our data for preferential alignment and elastic anisotropy. In the case of non-hydrostatically 141 compressed iron, angular dependence of  $V_P$  has been documented starting above ~80 GPa 142 [Antonangeli et al., 2004a], as a combined effect of the deformation-related development of 143 preferred orientation and of the intrinsic single-crystalline elastic anisotropy. Preferential 144 alignment of the c-axis along the main compression axis of the cell has been observed for 145 several hcp metals when compressed uniaxially [Wenk et al., 2001; Merkel et al., 2004; 146 Merkel et al., 2006]. Hence, we compared the relative intensities of the [100], [002] and [101] 147 reflections and the c/a ratio obtained under hydrostatic conditions using Ne as pressure 148 transmitting medium with those obtained under non-hydrostatic conditions. While all the 149 diffraction patterns collected from hydrostatically compressed samples show no significant 150 variation on the intensities with increasing pressure. However, the intensity of the [002] 151 reflection exhibits a strong reduction upon compression and almost vanishes at pressures exceeding 30 GPa, for samples loaded without pressure transmitting medium. Such behavior 152 153 is expected for the utilized diffraction geometry as a consequence of the progressive 154 alignment of the crystalline c-axis with the main compression axis of the cell. However, the 155 volumes and the values for c/a derived solely from the [100] and [101] reflections did not 156 significantly differ from those obtained under quasi-hydrostatic compression up to  $\sim 90$  GPa. 157 Conversely, the volume derived at 108 GPa is quite different from expected according to the 158 quasi-hydrostatic equation of state [Fiquet et al., 2008] and the c/a ratio displays a large 159 deviation. Deviation of individual d-spacings from the values expected in the limit of 160 hydrostatic compression is a direct consequence of the presence of a deviatoric stress within the cell. For elastically isotropic materials, the strain  $\epsilon_{hkl} = (d_{hkl}-d^{hydro}_{hkl}) / d^{hydro}_{hkl}$  for all the 161 162 lattice planes is the same. However, this is not the case in the presence of elastic anisotropy 163 [see for instance Sing et al., 1998; Merkel et al., 2009]. Thus, we can conclude that all our 164 non-hydrostatic pressure points ( $P \sim 32$ , 68 and 108 GPa) exhibit some texture, but a 165 significant deviation from a radially averaged sound velocity is observed only at 108 GPa. To 166 support this conclusion, we note that as a direct consequence of the developed preferential 167 alignment and the predominant contribution of the slow-velocity basal plane [Antonangeli et 168 al., 2004a; Antonangeli et al., 2004b; Antonangeli et al., 2006; Mao et al., 2008], the  $V_P$  value 169 at the highest density (highest pressure) is somewhat lower than the linear trend, although still 170 within uncertainties (Figure 2).

Due to the deviation from radial average properties at the highest pressures, we only considered data up to 68 GPa for extrapolation of  $V_P$  and  $V_S$  to core conditions. Our results are reported in Figure 3 along with linear regressions and the density evolution proposed for pure Fe [Badro et al., 2007], and compared with the seismic velocity profile from the radial 175 PREM model for the inner core [Dziewonski and Anderson, 1981]. As already suggested 176 [Fiquet et al., 2001; Badro et al., 2007],  $V_P$  for pure hcp-Fe is somewhat lower than the 177 PREM; this study shows that adding 3.7 wt% Si yields a velocity that is too fast relative to 178 PREM (Figure 3). If we consider linear mixing of pure Fe and Fe<sub>0.89</sub>Ni<sub>0.04</sub>Si<sub>0.07</sub> (assuming an 179 ideal two-component solid solution), we simultaneously match the PREM values of  $V_P$  and  $\rho$ 180 for an alloy with 1.2 wt% of Si (gray dashes in Figure 3).

181 Considering V<sub>s</sub> significantly increases the complexity of the problem. Both Fe and  $Fe_{0.89}Ni_{0.04}Si_{0.07}$  exhibit values of shear velocities<sup>1</sup> significantly higher than PREM (Figure 3). 182 183 The mismatch for  $Fe_{0.89}Ni_{0.04}Si_{0.07}$  is even larger than for pure iron. Indeed, inclusion of 184 silicon leads to a much larger increase in  $V_S$  than in  $V_P$ , due to a smaller pressure derivative of 185 the bulk modulus (K') that is only partially balanced by a larger value of K<sub>0</sub>. As a result, 186  $Fe_{0.89}Ni_{0.04}Si_{0.07}$  has a larger pressure derivative for V<sub>S</sub> than pure Fe, which manifests itself as 187 a considerable divergence between the data sets when extrapolated to core densities. Using 188 this model, it is clear that we cannot simultaneously solve for  $V_P$ ,  $V_S$  and  $\rho$  of the core by 189 simply varying the amounts of Ni and Si or pressure.

190 At core temperatures (4000-6000 K) anharmonic effects are expected. These might be 191 particularly relevant to V<sub>s</sub>, as already pointed by computational studies [Laio et al., 2000; 192 Steinle-Neumann et al., 2001], that suggest significantly lower values of V<sub>s</sub> as temperature 193 approaches the melting point. As such, we applied temperature corrections to our ambient 194 temperature results. Very recent *ab initio* calculations on hcp iron [Vočadlo et al., 2009] 195 suggest about 9.5% and about 35% softening for  $V_P$  and  $V_S$  respectively, between 0 and 5000 196 K. Such a temperature increase yields a density reduction about 4%. Since we are interested 197 in only the anharmonic high temperature effects (we are comparing data at the same density), 198 we corrected these computational results according to the measured density dependence of

<sup>&</sup>lt;sup>1</sup> Extrapolated values of V<sub>S</sub> are obtained assuming the Birch's law to be valid for both V<sub>P</sub> and  $V_{\phi}=(K/\rho)^{1/2}$ . The resulting V<sub>S</sub> exhibits a sub-linear dependence on density.

199	sound velocities to compensate for the density variation, obtaining 4% softening of $V_{\text{P}}$ and
200	30% softening of $V_S$ at 5000 K and 13000 Kg/m <sup>3</sup> . The thermal softening of $V_P$ at constant
201	density requires an increase in the Si-content to ~ $1.5 \text{ wt\%}$ in order to match the seismic
202	observations. Most importantly, this composition appears to provide a simultaneous solution
203	for both $V_P$ and $V_S$ , consistent with PREM values (light gray dots in Figure 3).
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206	4. Conclusions
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208	Our new sound velocity and density measurements on hcp $Fe_{0.89}Ni_{0.04}Si_{0.07}$ polycrystalline

samples compressed in DAC to 108 GPa yield values of  $V_P$  that are about 9% higher than in pure iron (at the same density). As it has already been shown than alloying with Ni does not significantly influence the elastic properties of Fe, at least up to a concentration of 22 at% [Kantor et al., 2007], we can reasonably assume that the increase in the sound velocity is solely due to the presence of Si. This conclusion is further supported by comparison of our measurements with computational results on Fe<sub>0.9375</sub>Si<sub>0.0625</sub> [Tsuchiya and Fujibuchi, 2009].

215 Extrapolation to core densities and comparison with seismic velocity and density profiles 216 from PREM allow us to constrain compositional models of the core. Our results suggest an 217 inner core composition containing 4-5 wt% of Ni and 1-2 wt% of Si. The exact amount of Si 218 might vary depending upon the temperature corrections (here we used values calculated for 219 pure Fe), or if other elements are also present in the inner core. The approximations in our 220 model and the relatively large uncertainties, especially on V<sub>S</sub>, do not allow us to definitively 221 rule out other compositional models, and other light elements might be present in amount 222 below our present limit of sensitivity (~1 wt%). It should be noted that our conclusions 223 pertain strictly to solid Fe-alloys and hence the inner core. Elements such as oxygen, that

224 may be incompatible in solid relative to liquid Fe-alloys are expected to reside mainly in the 225 outer core (see for instance [Alfe et al., 2002; Badro et al., 2007]) and cannot be adequately 226 constrained here. However, our proposed core composition is consistent with existing 227 experimental data and, for the first time, simultaneously matches all three geophysical observables ( $V_P$ ,  $V_S$  and  $\rho$ ). Combined with Si partition coefficients between liquid and solid 228 229 Fe we can estimate the Si concentration of the liquid outer core and, hence, obtain a Si 230 concentration for the entire core. If we assume a molar partition coefficient for silicon between the liquid and solid phase of iron of  $1.2 \le D^{\text{Liq/Sol}} \le 1.9$  [Alfe et al., 2002], we obtain 231 232 a core composition with Si ranging from 1.2 to 4 wt%, for an inner core containing 1 to 2 233 wt% of Si. This result is at the lower range of those from core-formation and core-mantle 234 interaction models that often call for large amount of Si in the core (e.g. 7.3 wt% [Allègre et 235 al, 1995], 10.3 wt% [Javoy, 1995], 5-7 wt% [Wade and Wood, 2005]).

236 Several mechanisms have been proposed to explain the low shear velocity in the core, 237 including fluid inclusions [Singh et al., 2000; Vočadlo, 2007], viscoelastic relaxation [Jackson 238 et al., 2000] or the presence of randomly oriented anisotropic "patches" [Calvet and Margerin, 239 2008]. According to our present results, none of these is strictly necessary, and V<sub>P</sub> and V<sub>S</sub> can 240 be matched, although with some uncertainties due to the discussed approximations, by only 241 considering the effect of alloving silicon at the few wt% level with reasonable high 242 temperature anharmonic corrections inferred form recent theoretical calculations [Vočadlo et 243 al., 2007]. However, the above-mentioned possibilities [Singh et al., 2000; Vočadlo, 2007; 244 Jackson et al., 2000; Calvet and Margerin, 2008] become relevant to reconciliation of 245 observed seismic wave attenuation. In addition, seismic anisotropy, its variation with depth, 246 as well core hemisphericity, require higher complexity than the simple radial model discussed 247 here.

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## **383 Figure Captions**

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## Figure 1: Longitudinal acoustic phonon dispersion of Fe<sub>0.89</sub>Ni<sub>0.04</sub>Si<sub>0.07</sub> at ambient temperature and pressures of 27, 37, 47, 68 and 108 GPa (bottom to top). For clarity the dispersion at 32 GPa is not plotted. The lines are sine fits to the experimental data.

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391 **Figure 2:** Aggregate compressional sound velocity as a function of density. Circles: 392 Fe<sub>0.89</sub>Ni<sub>0.04</sub>Si<sub>0.07</sub>; squares: Fe [Figuet et al., 2001; Antonangeli et al., 2004a]; triangles: 393 Fe<sub>0.78</sub>Ni<sub>0.22</sub> [Kantor et al., 2007]; open hexagons: computational results on Fe<sub>0.9375</sub>Si<sub>0.0625</sub> 394 [Tsuchiya and Fujibuchi, 2009]. The displayed error bars of the velocities result from the 395 experimental uncertainties, the statistical error of the fit, and the finite-q resolution of the 396 spectrometer. The uncertainties in the densities are smaller than the symbols. Lines are linear 397 regressions to the experimental data (solid - Fe<sub>0.89</sub>Ni<sub>0.04</sub>Si<sub>0.07</sub>; dotted - Fe; dashed -398 Fe<sub>0.78</sub>Ni<sub>0.22</sub>).

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401 **Figure 3:** Aggregate compressional  $(V_P)$  and shear  $(V_S)$  sound velocities and density 402 extrapolations (considering only data up to 68 GPa). Circles: IXS data on Fe<sub>0.89</sub>Ni<sub>0.04</sub>Si<sub>0.07</sub>; 403 diamonds: PREM [Dziewonski and Anderson, 1981]. Solid lines - Fe<sub>0.89</sub>Ni<sub>0.04</sub>Si<sub>0.07</sub>; dotted lines – pure Fe [Badro et al., 2007]. V<sub>S</sub> has been derived combining the measured V<sub>P</sub> with 404 405 literature equation of state (Fe<sub>0.89</sub>Ni<sub>0.04</sub>Si<sub>0.07</sub>: [Figuet et al., 2008]; Fe: [Dewaele et al., 2006]). 406 Gray dashes are estimated values of  $V_P$  and  $V_S$  for Fe<sub>0.936</sub>Ni<sub>0.040</sub>Si<sub>0.024</sub> (Si content ~1.2 wt%), 407 neglecting temperature corrections. Light gray dots are estimated values of  $V_P$  and  $V_S$  for  $Fe_{0.93}Ni_{0.04}Si_{0.03}$  (Si content ~1.5 wt%) at 13000 Kg/m<sup>3</sup> and 5000 K. For clarity uncertainties 408 409 in our estimations are not reported.

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