

# Intense Terahertz Pulses from SLAC Electron Beams using Coherent Transition Radiation

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SLAC has two electron accelerators, the Linac Coherent Light Source (LCLS) and the Facility for Advanced Accelerator Experimental Tests (FACET), providing high-charge, high-peak-current, femtosecond electron bunches. These characteristics are ideal for generating intense broadband terahertz (THz) pulses via coherent transition radiation. For LCLS and FACET respectively, the THz pulse duration is typically 20 and 80 fs RMS and can be tuned via the electron bunch duration; emission spectra span 3 to 30 THz and 0.5 THz to 5 THz; and the energy in a quasi-half-cycle THz pulse is 0.2 and 0.6 mJ. The peak electric field at a THz focus has reached 4.4 GV/m (0.44 V/Å) at LCLS. This paper presents measurements of the terahertz pulses and preliminary observations of nonlinear materials response.

## I. INTRODUCTION

The terahertz (THz) region of the electromagnetic spectrum (roughly from 100 GHz to 10 THz) has long been used as a characterization tool, but in recent years the development of new, higher intensity sources, driven by both lasers and electron beams, has created broad interest in using THz light for chemical and biological imaging, and opened up a new regime in which THz

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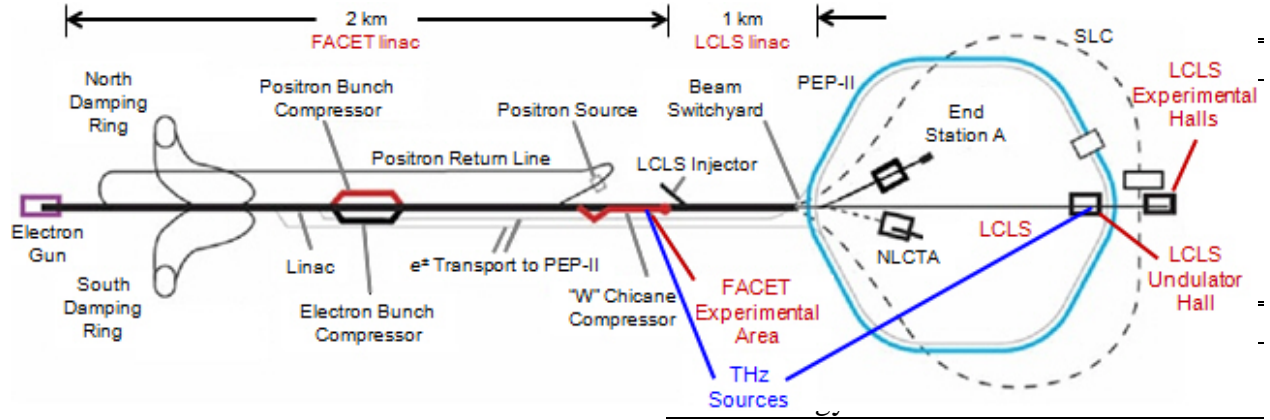
fields can be used to drive nonlinear responses in matter. Other articles in this Special Edition discuss a range of sources and applications.

Spectroscopic applications prefer a source with a narrow frequency span: a long pulse or a continuous source that is tunable over the THz band. Time-resolved applications prefer short pulses in the single-cycle limit, with correspondingly wide bandwidths, to quickly “kick” a system into an excited state; after a delay of femtoseconds to picoseconds, the response is probed with, for example, an optical pulse from a laser.

A relativistic electron bunch in a linear accelerator (linac) can be compressed during acceleration to a length as short as a few fs, depending on the bunch charge. For wavelengths longer than the bunch length, the emitted fields of all the electrons in the bunch add coherently, and the total radiated power scales with the square of the number of electrons. The electron bunches described here have charges of 100 to 3000 pC compressed to durations of 10 to 200 fs RMS, and so are well suited for the generation of intense THz pulses. The electromagnetic fields traveling with the beam can be converted to radiating fields in several ways: by bending the beam in a strong magnetic field (coherent synchrotron radiation and coherent edge radiation, CSR and CER), by passing the beam through a small aperture (coherent diffraction radiation, CDR), or by passing the beam from vacuum into a medium (coherent transition radiation, CTR). Several accelerator-based sources of THz radiation using these approaches have been constructed.<sup>1- 11</sup>

In CTR from a thin metallic foil, surface currents that match boundary conditions on the foil produce the fields of a collapsing dipole as the electrons approach, and an expanding dipole as they leave, generating an impulse of dipole radiation on both sides.<sup>12</sup> The temporal shape has a long-wavelength cutoff due to the finite size of the foil and apertures. Without these low frequencies, the shortest pulse resembles a “quasi-half-cycle,” with a high-amplitude central pulse of one polarity and two longer but weaker lobes of opposite polarity before and after. For highly relativistic electrons, the radiation emitted on exit from the foil travels in the beam direction, while the entrance radiation goes in the direction of specular reflection from the foil surface. By tilting the foil at 45° to the electron beam, this entrance radiation can be quickly extracted.

SLAC National Accelerator Laboratory has two such THz sources, at FACET and at LCLS. The former uses the first 2 km of the 3-km SLAC linear accelerator, while the latter uses the third km (Figure 1). Both sources use CTR, which offers a small and well defined source point (unlike



**Figure 1.** The 3-km SLAC linac, divided into FACET and LCLS linacs. THz sources are in the FACET experimental area and in the LCLS undulator hall.

CSR) and the absence of a short-wavelength cutoff from the hole in CDR. A multi-GeV beam experiences little scattering when passing through a thin ( $\leq 10 \mu\text{m}$ ) foil.

## A. FACET

FACET, the Facility for Advanced Accelerator Experimental Tests,<sup>13</sup> provides 3.2-nC electron (and, in future, positron) bunches at 20 to 23 GeV, compressed to lengths as short as 67 fs RMS, corresponding to a peak current of over 20 kA (Table 1). The beam is tightly focused at the interaction point (IP), a long optical table for user experiments (Figure 2). Advanced acceleration concepts are a major emphasis, notably electron acceleration in wakefields in plasmas<sup>14</sup> and also in dielectrics.<sup>15</sup> Both depend on the longitudinal bunch profile, which is measured by a new transverse deflecting cavity and by two THz experiments. The first studies Smith-Purcell radiation.<sup>16</sup> The second, which is presented here, uses CTR.

**Table 1.** FACET parameters

		Typical	Best	Design	
Electron energy		20.35	21.1	23	GeV
Charge per bunch		2500–2900	3200	3200	pC
Repetition rate		10	10	30	Hz
Bunch length	$\sigma_z$	25–30	<25	20	$\mu\text{m}$
	$\sigma_t$	83–100	<83	67	fs
Size at IP	$\sigma_x$	35	20	20	$\mu\text{m}$
	$\sigma_y$	35	23	20	$\mu\text{m}$
		Standard Optics		Double Waist	
Size at upstream	$\sigma_x$	1200		317	$\mu\text{m}$
THz foil	$\sigma_y$	6		36	$\mu\text{m}$

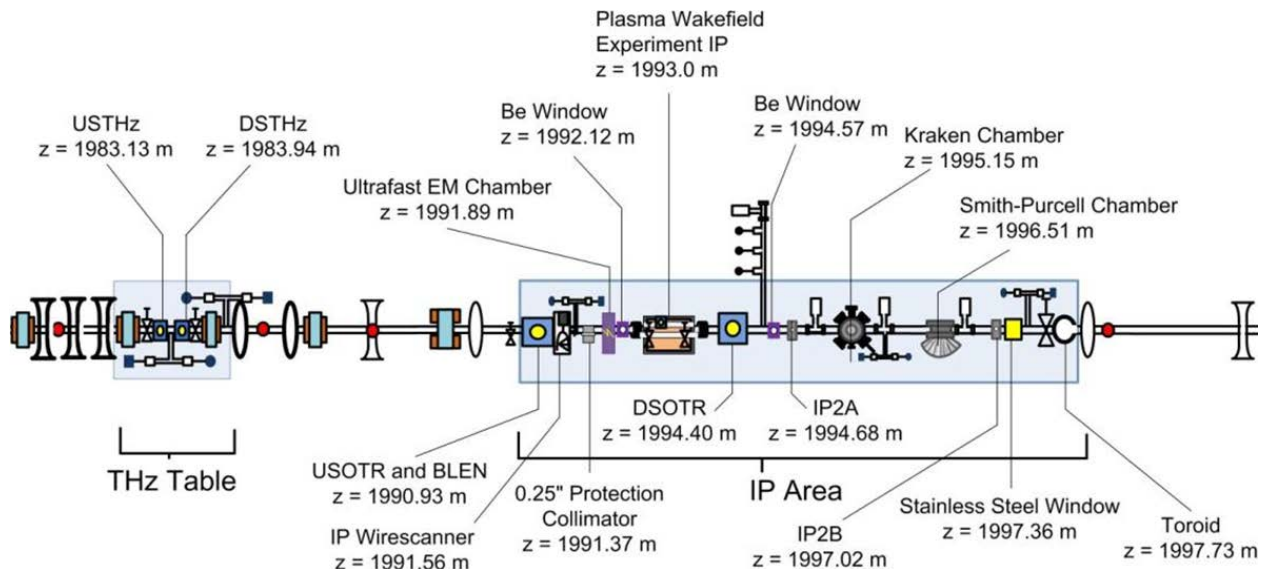
## B. LCLS

LCLS, the Linac Coherent Light Source, is the world's first hard x-ray free-electron laser (FEL).<sup>17</sup> Self-amplified spontaneous emission along more than 100 m of undulator produces coherent x-ray pulses. The electron energy is tuned over a broad energy range (Table 2) to provide a wide range of photon energies, down to wavelengths as short as 1.3 Å. Recent progress in x-ray self-seeding<sup>18</sup> offers nearly Fourier-transform-limited pulses of hard x-rays.

## II. THz AT FACET

### A. Setup

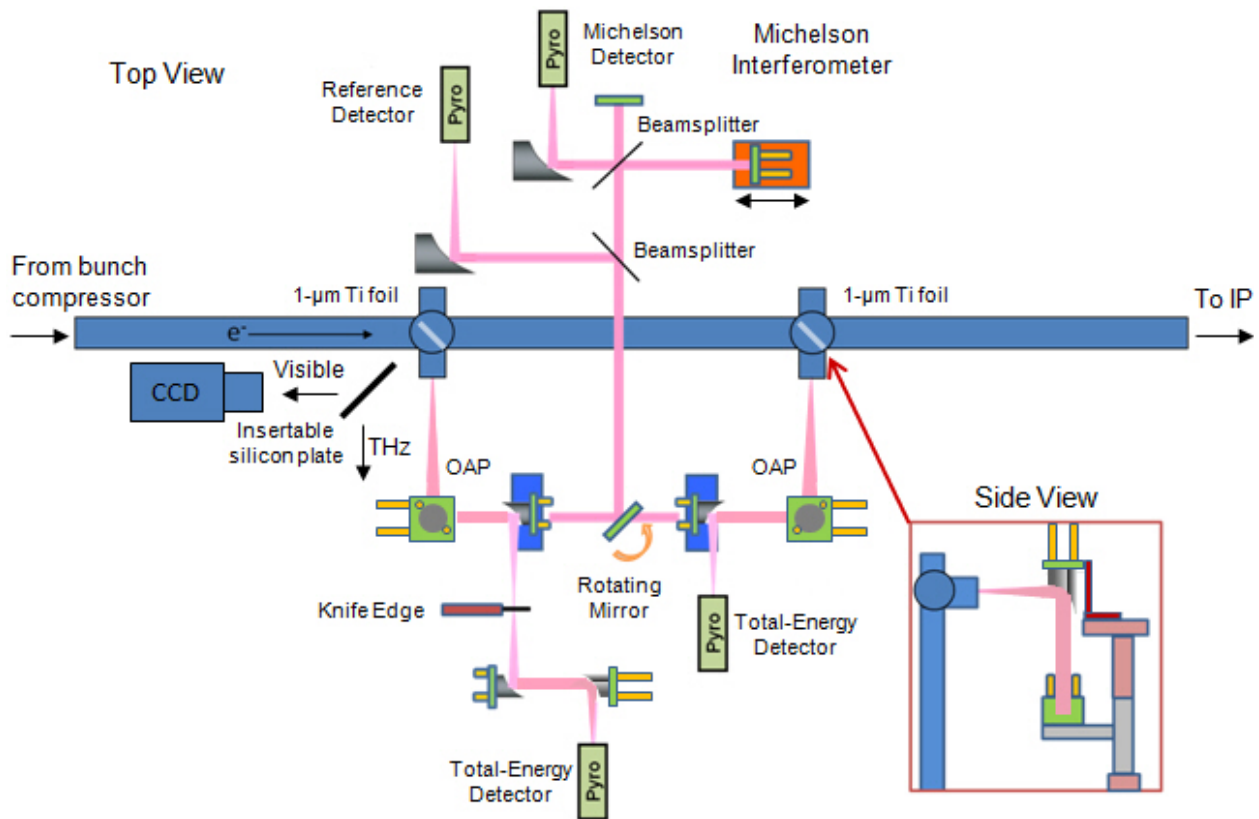
To provide sufficient space for optics while avoiding interference with experiments on the IP table, the THz source at FACET is on a separate optical table, 10 m upstream of the IP (Figure 2). The THz radiator is a 1- $\mu\text{m}$ -thick titanium foil rotated about the vertical axis by 45°. As electrons enter the foil, the THz they emit exits horizontally at 90° to the electron direction. The radiation passes through a 25-mm-diameter, 0.25-mm-thick, diamond window and is collimated by an off-axis parabolic mirror (OAP). Instruments on the table (Figure 3) characterize the THz pulses. A pyroelectric joulemeter gives the energy, knife-edge scans measure the size at the focus of another OAP, and a Michelson interferometer provides the spectrum. The entire tabletop is enclosed (Figure 4) and continually purged with dry air to avoid the strong absorption of THz by water vapor.



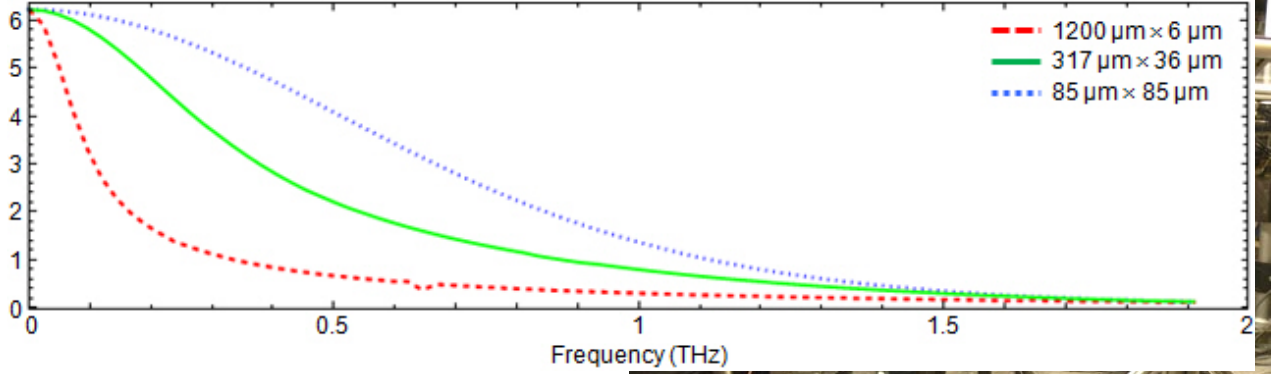
**Figure 2.** The THz and IP tables of FACET

The THz table is located after the final bunch compressor but in the middle of the sequence of quadrupole magnets that make the final focus at the IP. With standard IP optics, the nominal beam size at the upstream THz foil (Table 1) is highly elliptical, with  $\sigma_x = 1.2$  mm and  $\sigma_y = 6$   $\mu\text{m}$  at the upstream foil, and 1.5 mm by 13  $\mu\text{m}$  at the downstream foil. The emitted THz energy is reduced for wavelengths shorter than this transverse size, since the fields of the particles interfere destructively. Retuning the quadrupole magnets for a more circular spot at the THz foil gives a shallow waist of 317  $\mu\text{m}$  by 36  $\mu\text{m}$  while retaining the small focus at the IP. The measurements reported here use this “double-waist” configuration (Table 1).

The 1- $\mu\text{m}$  foil thickness is a trade-off between the effect of multiple Coulomb scattering on beam quality and the skin depth, 0.33  $\mu\text{m}$  at 1 THz, which reduces reflectivity and so the CTR yield below 1 THz. Titanium, a light and strong metal with a high melting temperature, was expected to be resistant to beam damage. However, the beam has punctured several of these foils on the IP table, near the minimum beam radius, and one foil on the THz table. The damage does not appear to be driven thermally but instead by the peak field, for reasons elaborated in Sec. III A. Other foil materials will be tested in the 2013 FACET run.



**Figure 3.** Layout of the FACET THz table.



**Figure 6.** Calculations comparing the spectrum (dotted blue line) to the double waist (solid green line) and to an 85- $\mu\text{m}$  waist (dashed red line).

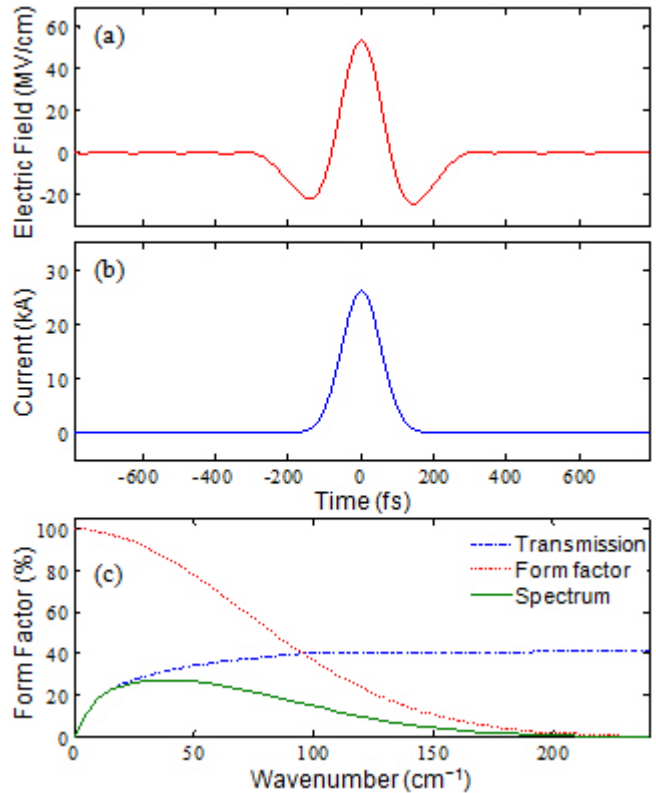
## B. Simulations

The THz power and spectrum radiated by CTR was calculated with a code<sup>19,20</sup> developed at SLAC. The foil is modeled as a perfectly conducting boundary with a unity reflection coefficient.

Figure 5 plots the calculated electric field of a 3-nC electron bunch, at an energy of 23 GeV, with a 50-fs length and a 10- $\mu\text{m}$  diameter. For this case, the peak current is roughly 27 kA and the maximum electric field reaches 0.6 V/Å in the  $\sim 150$  fs long quasi-half cycle pulse. The power spectrum peaks at 1.2 THz ( $40 \text{ cm}^{-1}$ ), with significant content up to 6 THz ( $200 \text{ cm}^{-1}$ ). The energy per pulse is 13 mJ. The bunch dimensions used were significantly smaller than those available later in the experiment, but the high fields and pulse energy show FACET's potential as a THz source.

Because the experiment initially used the standard optics of Table 1, with a ribbon-like beam, another simulation studied the effect on the CTR of the ribbon compared to a circular

**Figure 4.** FACET THz table and dry-air enclosure.



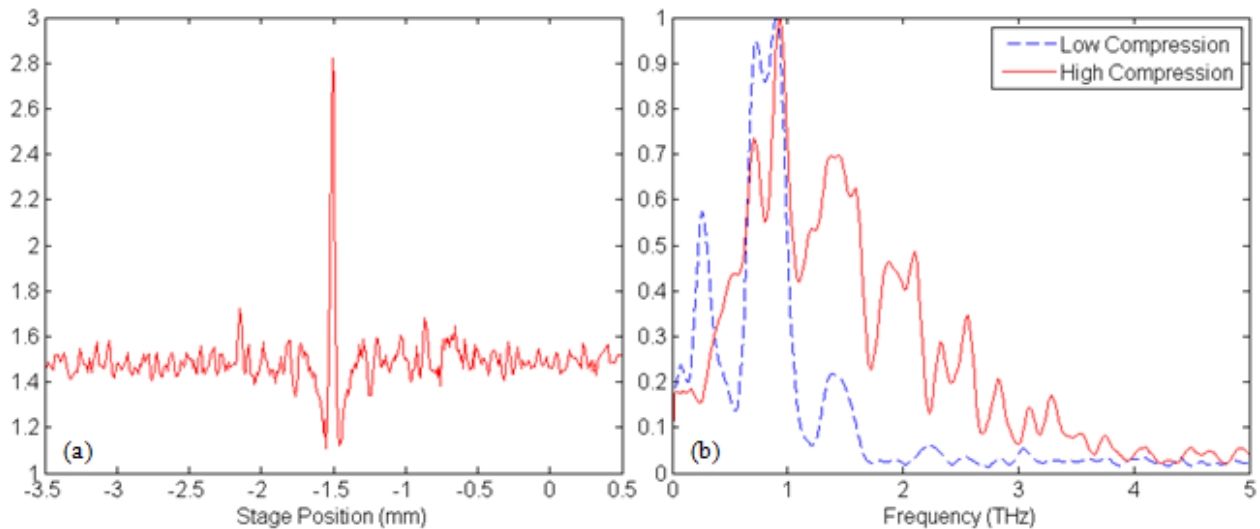
**Figure 5.** Calculated (a) CTR electric field, (b) electron-bunch peak current, and (c) CTR power spectrum, using the parameters discussed in the text.

beam. This required a different code<sup>21</sup> that follows Ter-Mikaelian.<sup>12</sup> Figure 6 shows the CTR radiation spectrum, normalized to the square of the charge and integrated over solid angle. The dashed curve corresponds to the elliptical spot of standard optics (Table 1,  $\sigma_x = 1.2$  mm,  $\sigma_y = 6$   $\mu\text{m}$ ), the solid line to the double waist ( $\sigma_x = 317$   $\mu\text{m}$ ,  $\sigma_y = 36$   $\mu\text{m}$ ), and the dotted curve to an 85- $\mu\text{m}$  circular beam. With double-waist optics, the power at frequencies near 0.5 THz increases by more than a factor of 2, and the circular beam provides an even greater improvement.

### C. Measurements

Figure 7(a) shows a typical Michelson interferogram of THz emitted by a highly compressed FACET electron bunch. The pyroelectric detector readings are averaged over 10 shots and normalized to readings from a reference pyro detector right before the interferometer (Figure 3). The 4-mm scan has a resolution of 37 GHz in the frequency domain.

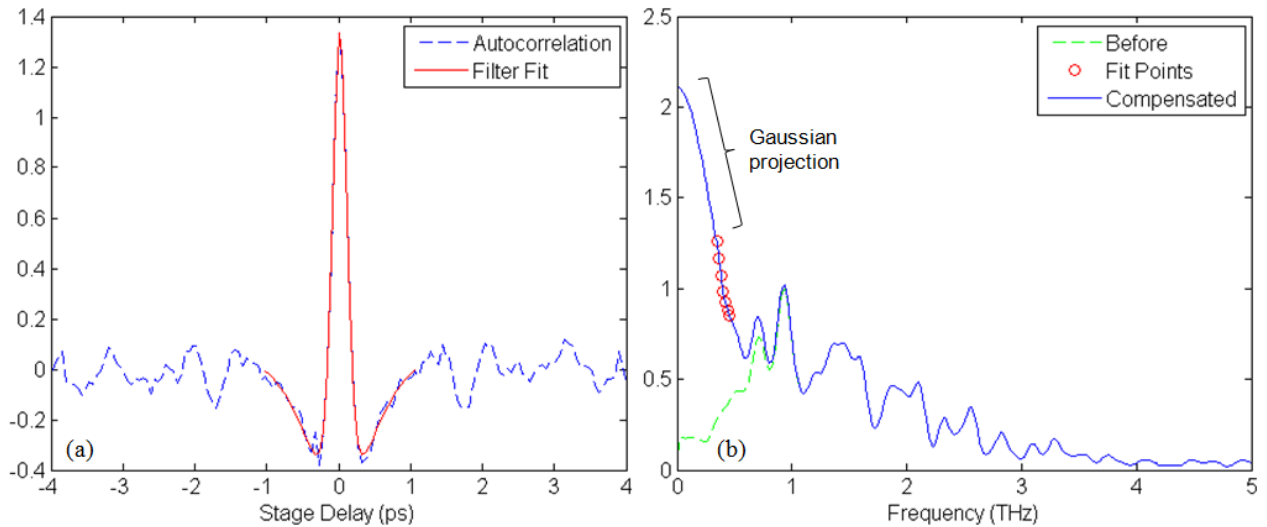
Fourier transforms of the scans give power spectra. Figure 7(b) shows the change in spectral content for high and low bunch compression. Strong water absorption lines near 1.3 THz are not present, due to the low humidity inside the enclosure. We attribute the 0.24-THz modulation in both spectra to Fabry-Pérot interference in reflections from the layers of the pyroelectric detector. The



**Figure 7.** (a) Michelson interferogram of a CTR THz pulse from a highly compressed electron bunch. The points are normalized to the reference detector. (b) Normalized power spectra of THz pulses from bunches with low and high compression.

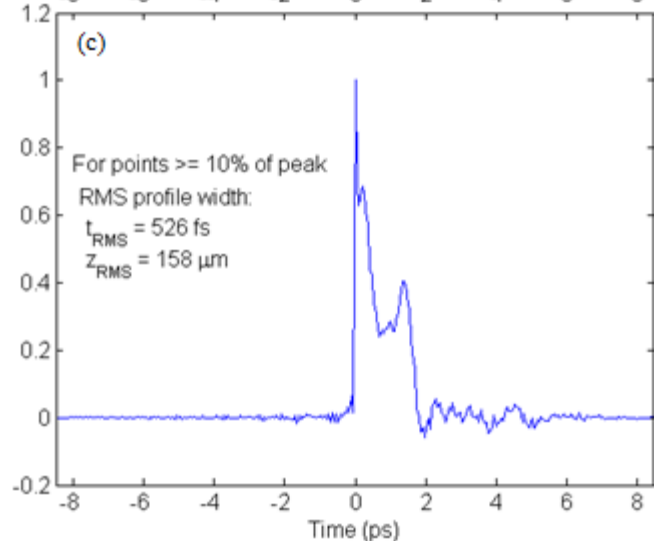
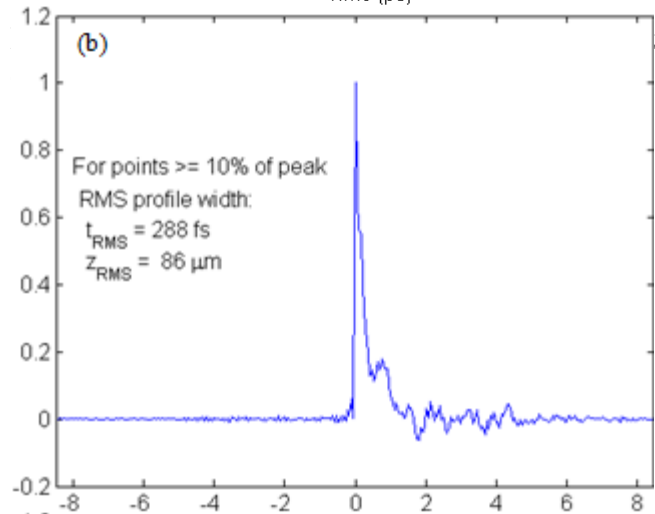
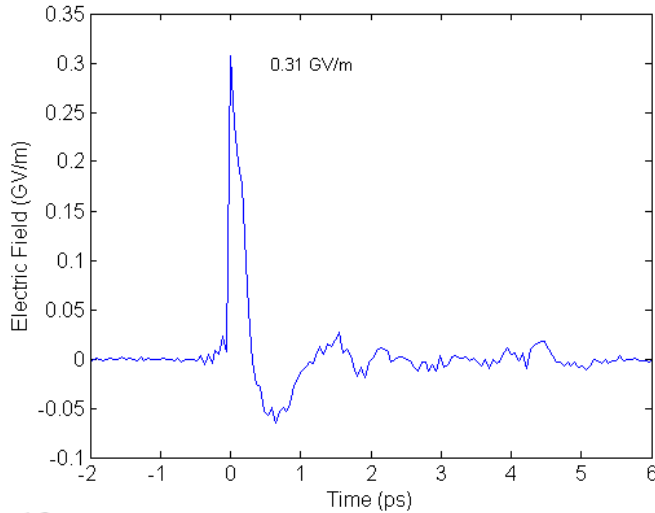
spectra were smoothed to reduce this effect.

The spectrum can be used to reconstruct the temporal profile of a bunch. Although the wakefield of a relativistic bunch follows the bunch shape, the fields radiated by the foil are filtered by the finite beampipe and window, and so lose long wavelengths. To compensate the spectrum for this loss, we model the attenuation of the field as a Gaussian high-pass filter,<sup>22</sup>  $1 - \exp(-f^2/f_0^2)$ , which leads to negative lobes on either side of the main peak of the interferogram. With data from a highly compressed beam, which has a nearly Gaussian profile, we fit the center of the autocorrelation in the time domain (Figure 8(a)) with a Gaussian beam and filter to find the filter constant  $f_0$ . For FACET,  $f_0$  lies between 0.35 and 0.45 THz, depending on the adjustment of the optics. We then discard the assumption of a Gaussian beam but retain  $f_0$ . The power spectrum is compensated by dividing by the square of the filter down to frequencies near  $f_0$ , where the correction starts to diverge and where little signal is available. From this point down to zero frequency, the correction is completed by fitting a Gaussian peaking at  $f = 0$  (Figure 8(b)). Empirically, it is found that the reconstruction does not depend strongly on the detailed shape at the lowest frequencies. Then the Kramers-Kronig relations allow reconstruction of the electron-bunch profile from the corrected spectrum,<sup>23</sup> as shown in Figure 9.

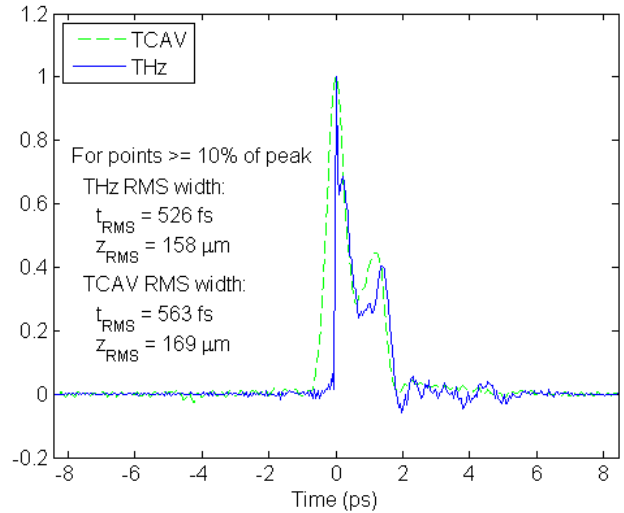


**Figure 8.** (a) Central region of the autocorrelation of Figure 7(a), after subtracting the mean, with a fit about the peak to find the high-pass filter. (b) The high-compression spectrum (dashed line) of Figure 7(b) is compensated (solid) for the high-pass filter, down to the leftmost circled point, below which a Gaussian fit to all the circled points extends the curve to  $f = 0$ .





**Figure 9.** Reconstruction of bunches at (a) high, (b) medium, and (c) low compression. A typical “double horn” structure develops as the compression is reduced.



**Figure 11.** A comparison of the THz profile reconstruction of Figure 9(c) with a measurement using a transverse deflecting cavity.

If this spectral correction is not applied, the Kramers-Kronig procedure reconstructs the electric field of the THz pulse on the optical table. Figure 10 shows an example using intermediate compression, in which measurements of the pulse energy (350  $\mu$ J and the transverse beam sizes (1.36 and 1.09 mm RMS, using knife-edge scans across a THz focus) were made along with the autocorrelation. The calculation gives a field peaking at 0.31 GV/m, corresponding to a magnetic field of 1.0 T. A reconstruction of this bunch gives an RMS length of 130 fs, double the length of the shortest FACET bunches. These should yield twice the energy and field.

The Kramers-Kronig algorithm does not resolve time reversal and generally places the main peak first; moreover, the entire tail

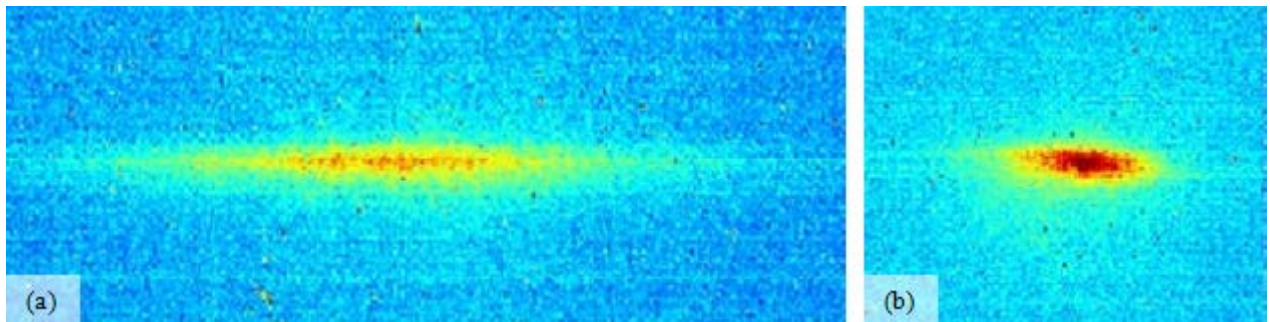
usually appears afterward, even for a symmetric single-cycle pulse like that in Figure 5(a). However, the bunch asymmetry in Fig 9 agrees well with measurements (Figure 11) using a newly commissioned X-band transverse deflecting cavity (TCAV).<sup>24</sup>

With the double-waist electron optics discussed in sections II A and B and in Table 1, a shallow waist with a more circular electron beam is formed at the upstream THz foil. Figure 12 compares optical-transition-radiation images of the transverse spot on the foil with the standard and double-waist configurations; the size scales as expected. The THz energy per pulse increased by 33%, from 0.47 to 0.63 mJ, and the spectrum shows more high-frequency content, in qualitative agreement with Figure 6. Quadrupole settings have been calculated to produce even smaller sizes, including the circular spot of Figure 6, when the constraint of a small size at the IP is lifted. Dedicated beam time to test these optical configurations has been requested during the 2013 run in order to determine the largest transverse size consistent with maximum CTR.

### III. THz at LCLS

#### A. Setup

The LCLS THz source is located in the Undulator Hall, 31 m downstream of the undulator. Measurements have been carried out at electron energies of 14 GeV with bunch charges from 150 to 350 pC. A beryllium foil tilted downward by 45° sends THz through a 25-mm-diameter, 0.25-mm-thick diamond window to optics on a table below. These include characterization measurements similar to those at FACET, plus a 20-fs Ti:sapphire oscillator, for THz-pump and laser-probe studies. The reflection and transmission of a variety of samples can be measured as they are exposed to intense THz pulses. Again the tabletop is enclosed, both for a dry-air purge and for laser safety.



**Figure 12.** The FACET electron beam on the upstream THz foil, imaged with the same scale using optical transition radiation, for (a) standard and (b) double-waist electron optics.

Initial measurements used a 2- $\mu\text{m}$ -thick beryllium foil, a metal chosen both for low scattering of the electron beam and for weak absorption of FEL x rays that also pass through the beampipe. The skin depth for Be is 95 nm at 1 THz. This foil was broken after several months when a vacuum-valve controller malfunctioned. To avoid the toxicity of Be particulates, aluminum was substituted. Maintaining low x-ray absorption required a 500-nm thickness, which is still much larger than the 85-nm skin depth at 1 THz.

The electron beam punctured three Al foils, even when the undulators were pulled away from the beam so that no FEL x rays were produced. In a systematic study with a constant charge of 150 pC and a low 10-Hz repetition rate, the bunch was gradually compressed. No damage was seen until the bunch length was reduced to 100 fs FWHM. Then a video camera detected light emitted from the beam position on the foil, and within seconds a small hole ( $\sim 1$  mm) was visible. We concluded that damage depending on compression at the femtosecond timescale cannot be thermal but is likely a mechanism driven by the peak field of the bunch, such as a surface discharge. Material selection by thermal and mechanical properties is not sufficient. We resumed work with Be, but with a 10- $\mu\text{m}$ -thickness and a new holder that parks the Be when not in use inside a volume with low conductance to the beampipe. This design slows a sudden pressure rise and equalizes the pressure on both sides of the foil. Be has proven to be resistant to beam damage, even at full charge (350 pC) and compression (50 fs FWHM).

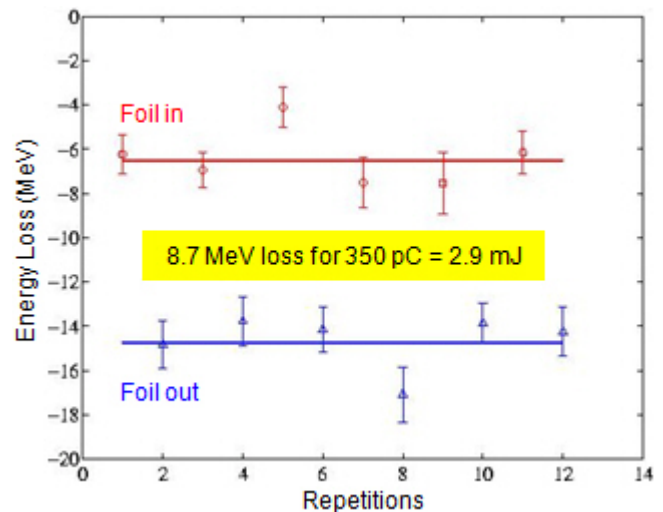
## B. Simulations

Simulations of the radiated field strength and temporal profile are similar to those described above for FACET. The spatio-temporal electron beam profile is obtained by a simulation of the entire LCLS accelerator using Impact-T and Elegant codes.<sup>25</sup> The initial electric field at the Be foil is transformed into a frequency-domain profile and propagated using Laguerre-Gaussian modes through the optical setup, including diffraction effects from the finite size of the foil. These simulations calculate the full extracted spatio-temporal THz field at the focus of the optical setup, predicting a radially polarized beam with zero intensity on axis, reflecting the relativistic Coulomb field of the electron bunch. For the highest charge and shortest bunch measured (350 pC, 50 fs FWHM), at the focus of an OAP with an effective focal length of 150 mm ( $\text{NA} = 0.1$ ), the model predicts a 400- $\mu\text{m}$  FWHM transverse profile and a peak field 100  $\mu\text{m}$  from the axis of 7 GV/m (0.7 V/Å). Integrating over the temporal and spatial profile gives a radiated energy of 560  $\mu\text{J}$ .

### C. Measurements

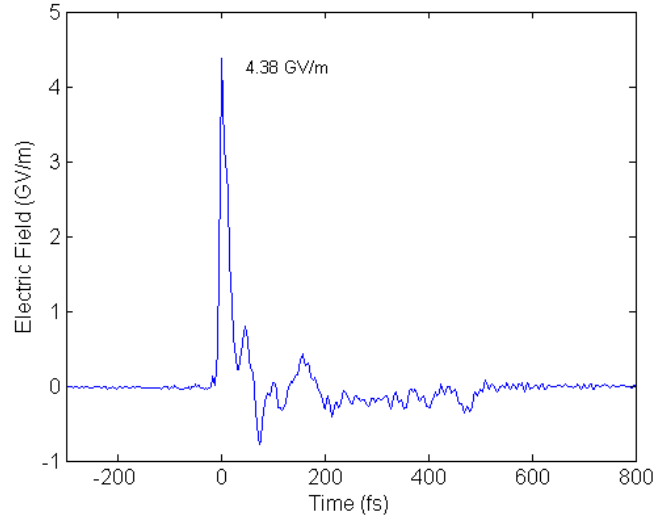
Total energy measurements of the radiated THz fields are carried out using several methods to check for consistency. A thermopile power meter with a graphite absorber and a pyroelectric detector (Gentec) were used initially, with maximum recorded energy of  $\sim 200 \mu\text{J}/\text{pulse}$ . The energy may be higher, since Gentec later provided preliminary calibration correction factors of 2.1 to 4 at 2.52 THz, depending on the detector type. Similarly, the graphite in the thermopile is not a perfect THz absorber.

The total radiated THz energy can also be obtained by measuring the energy loss of the electrons due to CTR, an approach not subject to uncertainties about the broadband calibration of the detectors. The bend at the electron beam dump was designed as an energy spectrometer, allowing us to compare the deflection with the foil radiator in and out of the beam. The electron beam is estimated to lose only  $30 \mu\text{J}$  in a  $2\text{-}\mu\text{m}$  Be foil tilted at  $45^\circ$ , and  $150 \mu\text{J}$  in a  $10\text{-}\mu\text{m}$  foil.<sup>26</sup> Figure 13 shows measurements with the  $2\text{-}\mu\text{m}$  Be foil indicating an 8.7-MeV loss per electron, corresponding to a 2.9-mJ loss for 350 pC. Half of this THz is radiated in the forward direction as the beam exits the foil. The radiation on entrance heads downward, but 30% is lost by reflection from the high-index diamond window, so that about 1 mJ would be expected on the table. Additional losses, particularly absorption by water vapor (before the enclosure was installed) reduce these fields by another factor of  $\sim 2$ . This approach provides a complimentary estimate of the radiated fields. Given the uncertainties, the simulations and the energy loss are in reasonable agreement with the detector measurements.

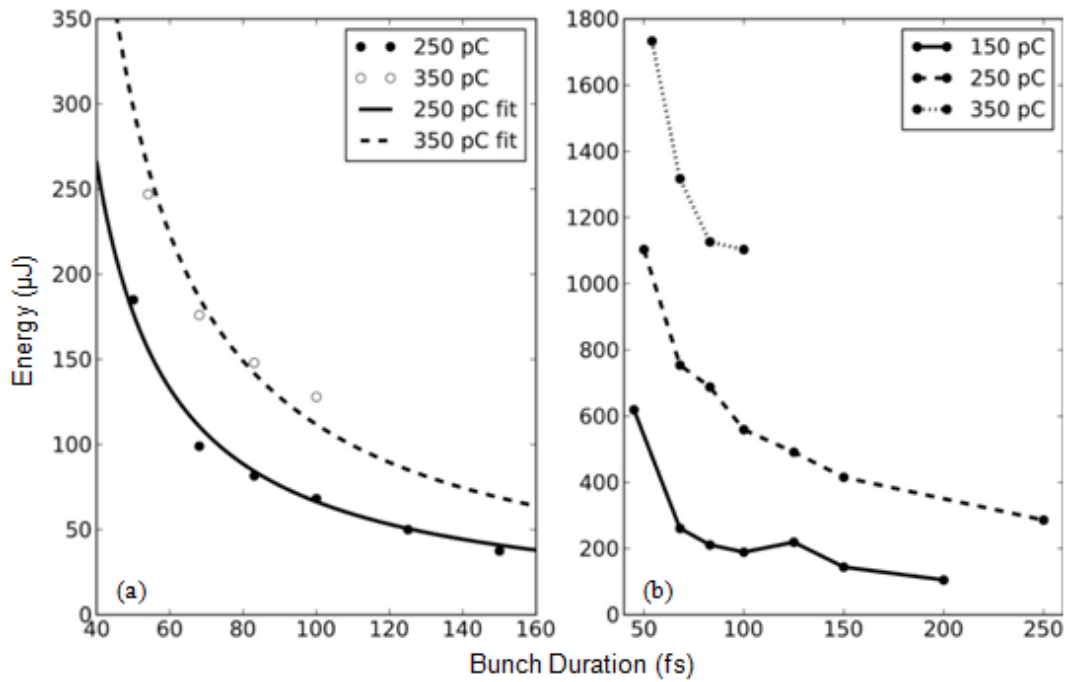


**Figure 13.** Electron-beam energy loss due to inserting the  $2\text{-}\mu\text{m}$  Be foil.

For a charge  $q$  and pulse duration  $t$ , the radiated THz field is expected to scale with the time-dependent peak current  $q/t$ , and the intensity with  $q^2/t^2$ . Integrating over time, the radiated energy then scales with  $q^2/t$ . Figure 14 shows measurements which roughly confirm this scaling law, using (a) a pyroelectric detector at a THz focus, and (b) the electron energy loss. This scaling cannot be precise: the bunch shape changes with compression, becoming more peaked and so increasing emission faster than  $1/t$ . Also, the nominal bunch length in the figure was taken from a monitor measuring filtered infrared from a bend at the end of the final compressor. It was calibrated with a transverse deflecting cavity, but the bandpass filter makes it somewhat inaccurate for the shortest bunches. The best fit to the above scaling offsets the bunch length by approximately 20 fs.



**Figure 15.** Reconstruction of the focused THz field in LCLS.



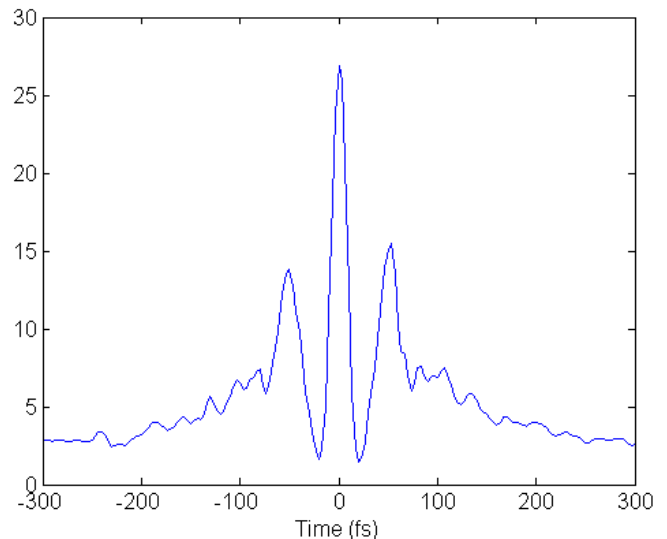
**Figure 14.** Dependence of THz energy on charge and FWHM bunch duration, measured by (a) a pyroelectric joulemeter at a THz focus, and (b) the electron-beam energy loss.

As with FACET, the Kramers-Kronig procedure can reconstruct the temporal profile of the electric field from an autocorrelation. Using a measured autocorrelation of a 350-pC, 70-fs FWHM bunch and a typical energy of 180  $\mu$ J from Figure 14, the reconstruction in Figure 15 calculates a peak field of 4.4 GV/m at the focus on the optical table. This field is much stronger than the FACET field in Figure 10 since this energy is contained in a shorter bunch and a smaller focal spot, 200  $\mu$ m RMS. The smaller source size and shorter bunch result in a spectrum with higher frequencies, peaking at 10 THz,<sup>11</sup> and so allow a tighter focus.

Further evidence for the high field strength is provided by nonlinear autocorrelations. The field (linear) autocorrelation shown in Figure 7(a) provides a direct measurement of the spectrum of the radiated field. The peak-to-baseline ratio for such scans is limited to 2, with a perfect overlap of the fields in space and time on the pyroelectric detector. In other autocorrelations, two THz pulses were incident collinearly on a silicon photodiode, which provides an intrinsically nonlinear response since its 1.1-eV bandgap is well above the THz photon energy.<sup>27</sup> Figure 16 plots the diode's photocurrent as a function of the relative time delay between the two THz fields for a bunch with intermediate compression (90 fs FWHM for 350 pC). The scan shows the 3-peak autocorrelation expected for a double-horn distribution,<sup>11</sup> but with a high peak-to-baseline ratio of 9 that indicates the photodiode's strongly nonlinear response to the high fields.

#### IV. ACKNOWLEDGMENTS

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**Figure 16.** A nonlinear autocorrelation measured with a silicon photodiode.

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