Pressure and Fluid Saturation Prediction in a Multicomponent Reservoir, using Combined Seismic and Electromagnetic Imaging

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ABSTRACT

This paper presents a method for combining seismic and electromagnetic measurements to predict changes in water saturation, pressure, and CO₂ gas/oil ratio in a reservoir undergoing CO₂ flood. Crosswell seismic and electromagnetic data sets taken before and during CO₂ flooding of an oil reservoir are inverted to produce crosswell images of the change in compressional velocity, shear velocity, and electrical conductivity during a CO₂ injection pilot study. A rock properties model is developed using measured log porosity, fluid saturations, pressure, temperature, bulk density, sonic velocity, and electrical conductivity. The parameters of the rock properties model are found by an L1-norm simplex minimization of predicted and observed differences in compressional velocity and density. A separate minimization, using Archie's law, provides parameters for modeling the relations between water saturation, porosity, and the electrical conductivity. The rock-properties model is used to generate relationships between changes in geophysical parameters and changes in reservoir parameters. Electrical conductivity changes are directly mapped to changes in water saturation; estimated changes in water saturation are used along with the observed changes in shear wave velocity to predict changes in reservoir pressure. The estimation of the spatial extent and amount of CO₂ relies on first removing the effects of the water saturation and pressure changes from the observed compressional velocity changes, producing a residual compressional velocity change. This velocity change is then interpreted in terms of increases in the CO₂ /oil ratio. Resulting images of the CO₂/oil ratio show CO₂-rich zones that are well correlated to the location of injection perforations, with the size of these zones also correlating to the amount of injected CO₂. The images produced by this

process are better correlated to the location and amount of injected CO₂ than are any of the individual images of change in geophysical parameters.

INTRODUCTION

Crosswell seismic and electromagnetic technology has developed over the past two decades to provide high spatial resolution images of the seismic velocities (P and S) and electrical conductivity of the interwell region. The majority of effort, as measured by the topics of published and presented work, has concentrated on developing and improving algorithms for estimating the geophysical parameters themselves (Newman, 1995; Lazaratos et al., 1995; Wilt et al., 1995; Nemeth et al., 1997; Goudswaard et al. 1998 to In most applications where nongeophysical parameters, such as list but a few). temperature during a steam flood (Lee et al., 1995) or CO₂ saturations during CO₂ flood (Harris et al., 1995; Wang et al., 1998) are the object of the crosswell survey, correlations between the geophysical parameters, e.g., velocity or electrical conductivity, and the desired reservoir parameter are derived and used to infer the distribution of reservoir parameters from the distribution of the geophysical parameters. The output from the survey is still most commonly a cross section of velocity, electrical conductivity or the time-lapse change of these parameters, which is then interpreted in terms of its implications for the distribution and/or change of the parameter of interest (temperature, CO_2 saturation, etc.).

The simple extension of interpreting the geophysical parameters themselves is to use relationships between geophysical and reservoir parameters (e.g., a regression fit between velocity and temperature) to convert a geophysical parameter to a reservoir parameter image. This approach can be used successfully in relatively simple reservoir systems with a minimum of fluid components and/or spatial variations in other controlling parameters (such as porosity, pressure, and temperature). However, in many settings the geophysical parameters depend on a number of reservoir parameters that are variable in both space and time. In particular, porosity, pressure, water, and gas saturation strongly influence seismic velocity. Electrical conductivity can generally be described as a function of porosity, water saturation, and fluid conductivity (Archie, 1942), although clay content may also need to be considered. As we will show, in a complex reservoir fluid system, the spatial distribution of the time-lapse change in geophysical parameters, such as velocity, can vary significantly from the spatial distribution of the time-lapse change in a desired reservoir parameter, such as CO₂ saturation in oil. This difference results from the dependence of the geophysical parameters on more than one reservoir parameter (such as pressure and water saturation). These multiple dependencies must be sorted out before a picture of any single reservoir parameter can be obtained.

It has become common practice to use time-lapse changes in compressional and shear impedance mapped at the top of a reservoir. These changes are used to calculate timelapse changes in effective pressure and water saturation within a reservoir without significant gas saturation (Landro, 2001). However, in systems where natural gas is present in significant concentrations or where gas in the form of CO_2 is introduced, quantitative prediction of pressure and fluid saturation changes becomes problematic because of trade-offs in the effects of the multiple reservoir parameters on the mapped geophysical parameters. The situation is further complicated if the objective is to monitor CO_2 injection into a reservoir already containing natural gas (in addition to oil and water).

The objective of the work described in this paper is to demonstrate a methodology of combining time-lapse changes in electric conductivity and compressional- and shear-wave velocity with a detailed rock-properties model, to produce quantitative estimates of the change in reservoir pressure, water saturation, and CO₂/oil ratio.

THE FIELD EXPERIMENT

Crosswell seismic tomography and electromagnetic imaging have been demonstrated in separate applications over the last decade. The SEG special issue 'Crosswell Methods' (Rector, 1995) contains several papers on the application of crosswell seismic tomography specifically for thermal process monitoring and several others on crosswell EM monitoring of water floods. Wilt et al. (1995) report on the application of crosswell EM in water flood monitoring.

In the fall of 2000 and spring of 2001 we conducted crosswell seismic and electromagnetic (EM) measurements in the Lost Hills oil field in southern California during a CO_2 injection pilot study by Chevron Petroleum Co. The objective of the pilot study was to demonstrate enhanced oil recovery resulting from CO_2 injection. We used this opportunity to study geophysical imaging of the reservoir during CO_2 injection.

The portion of the Lost Hills field where this experiment took place has been undergoing water flood since 1995. The CO_2 pilot covers four injection wells and surrounding producers. Figure 1 shows the well placement in the affected portion of the

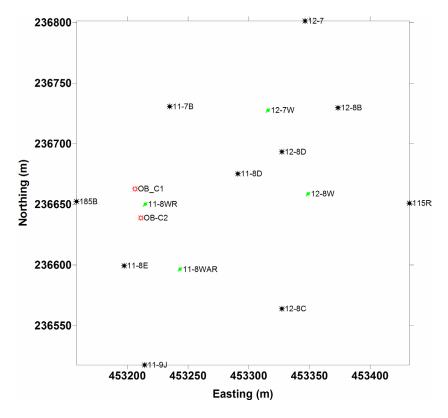


Figure 1. Area of the Lost Hills field affected by CO_2 injection. Four water injectors (shown in green) were converted to CO_2 injection in September 2000. The crosswell experiments took place between observation wells OB-C1 and OB-C2 (shown in red). A flow simulation production history match was done on the portion of the field covered by this figure.

field. Observation wells, OB-C1 and OB-C2, were drilled for the pilot and were fiberglass-cased to enable the use of crosswell EM. The nearby CO_2 injector (11-8WR) is located 20 feet out of the crosswell-imaging plane. These injection wells were hydraulically fractured to increase injectivity into the low-permeability diatomite reservoir. In some cases, downhole pressures were increased above the lithostatic

pressure, which may have induced fracturing above the desired injection interval. If the fracture did indeed extend above the desired interval, much of the injected CO_2 would likely not sweep its intended target, but rather move into the higher section.

The baseline crosswell seismic and EM surveys were conducted in September 2000, just prior to the beginning of CO_2 injection. A second EM survey was conducted in mid April 2001, and a second seismic survey was conducted in May 2001. In addition to the crosswell surveys, the two observation wells OB-C1 and OB-C2 were relogged for electrical resistivity in January 2001.

A ROCK-PROPERTIES MODEL

The reservoir parameters that have a dominant affect on geophysical parameters are porosity, pore pressure, effective pressure (lithostatic-pressure minus pore-pressure), fluid saturation, and the amount of dissolved hydrocarbon gas or CO₂ in oil. Pressure has a significant effect at Lost Hills because it is a shallow reservoir in soft rock. Converting geophysical images of the interwell region to reservoir parameters requires a rock-properties model relating the geophysical parameters to the reservoir parameters. We sought a model that would be able to predict observed velocity, density and electrical conductivity from observed pressure, porosity, and fluid saturations. Table 1 gives all the symbol definitions used in this paper. Laboratory measurements of the dry-frame moduli and grain density of the diatomite reservoir rock were unavailable, so to compute the seismic velocity we used the Hertz-Mindlin contact theory for the effective bulk (K_{dry}) and shear (G_{dry}) moduli of a dry, dense, random pack of spherical grains given by the following expressions:

$$K_{dry} = \left[\frac{l^2 \left(1 - \phi_0\right)^2 G_{grain}^2}{18\pi^2 \left(1 - v\right)^2} \cdot P_{eff}\right]^{1/3}$$
(1)

$$G_{dry} = \frac{5 - 4\nu}{5(2 - \nu)} \left[\frac{3l^2 (1 - \phi_0)^2 G_{grain}^2}{2\pi^2 (1 - \nu)^2} P_{eff} \right]^{1/3},$$
(2)

where ϕ_0 is the critical porosity (the porosity above which the grains become a liquid suspension), P_{eff} is the effective pressure, ν is the grain Poisson's ratio, G_{grain} is the grain shear modulus and l is the average number of other grains each grain contacts. Equations (1) and (2) describe the effective dry-frame moduli at the critical porosity ϕ_0 . The modified Hashin-Shtrikman lower bounds (Hashin and Shtrikman, 1963) given by Dvorkin and Nur (1996),

$$K_{eff} = \left[\frac{\phi/\phi_0}{K_{dry} + 4/3G_{dry}} + \frac{1 - \phi/\phi_0}{K_{grain} + 4/3G_{dry}}\right]^{-1} - 4/3G_{dry}$$
(3)

$$G_{eff} = \left[\frac{\phi/\phi_{0}}{G_{dry} + \frac{G_{dry}}{6}\left(\frac{9K_{dry} + 8G_{dry}}{K_{dry} + 2G_{dry}}\right)} + \frac{1 - \phi/\phi_{0}}{G + \frac{G_{dry}}{6}\left(\frac{9K_{dry} + 8G_{dry}}{K_{dry} + 2G_{dry}}\right)}\right]^{-1}$$
(4)
$$G_{dry}\left(9K_{dry} + 8G_{dry}\right)$$

$$\frac{G_{dry}}{6} \left(\frac{9K_{dry} + 8G_{dry}}{K_{dry} + 2G_{dry}} \right)$$

are used to model the dry frame moduli (K_{eff} and G_{eff}) at porosity ϕ , where K_{grain} is the grain bulk modulus.

The bulk modulus of the fluid saturated rock (K_{sat}) is modeled by Gassmann's equation (Gassmann, 1951):

$$K_{sat} = K_{grain} \left[\frac{\phi \cdot K_{eff} - \left[\frac{(1+\phi) \cdot K_{fluid} \cdot K_{eff}}{K_{grain}} \right] + K_{fluid}}{(1+\phi) \cdot K_{fluid} + \phi \cdot K_{grain} - \left[\frac{K_{fluid} \cdot K_{eff}}{K_{grain}} \right]} \right] , \qquad (5)$$

where K_{fluid} is the aggregate bulk modulus of the fluids filling the pore space. The bulk shear modulus of the fluid saturated rock is assumed to equal that of the dry rock.

The possible fluids filling the pore space are oil, brine, hydrocarbon gas, and CO₂. A common approach for calculating K_{fluid} is to use Wood's mixing formula (Wood, 1955):

$$1/K_{fluid} = S_w/K_{brine} + S_{oil}/K_{oil} + S_{hcg}/K_{hcg} + S_{co2}/K_{co2}, \qquad (6)$$

where the water saturation (S_w) , oil saturation (S_{oil}) , hydrocarbon gas saturation (S_{hcg}) and CO₂ saturation (S_{co2}) sum to 1.0. The bulk moduli of brine, oil, hydrocarbon gas, and CO₂ are K_{brine} , K_{oil} , K_{hcg} , and K_{CO2} , respectively. We will discuss this method of calculating K_{fluid} at the end of this section. The bulk density is given by a simple mixing law

$$\rho_{bulk} = (1 - \phi) \rho_{grain} + \phi (1 - S_{hcg} - S_w - S_{CO_2}) \rho_{oil} + \\
\phi S_w \rho_{brine} + \phi S_{hcg} \rho_{hcg} + \phi S_{CO_2} \rho_{CO_2}$$
(7)

where ρ_{grain} , ρ_{oil} , ρ_{brine} , ρ_{hcg} , and ρ_{C0_2} are the grain, oil, brine, hydrocarbon gas, and

CO₂ densities, respectively, as a function of pressure and temperature.

The fluid bulk moduli $K_{brine}, K_{oil}, K_{hcg}$ and densities ρ_{brine} , ρ_{oil} , ρ_{hcg} of the brine, oil, and hydrocarbon gas respectively are computed using relations from Betzel and Wang (1992). The bulk modulus and density of CO₂, K_{CO_2} and ρ_{CO_2} , respectively, as well as the bulk moduli and densities of CO₂-hydrocarbon gas mixtures, are modeled using relations from Magee and Howley (1994).

The bulk electrical conductivity (σ_{bulk}) of the reservoir rock is modeled using Archie's (1942) relationship

$$\sigma_{bulk} = \sigma_{brine} \cdot \phi^m \cdot S_w^n \quad , \tag{8}$$

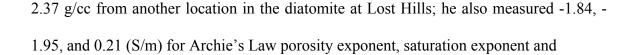
where $\sigma_{\rm brine}$ is the fluid conductivity, and *m* and *n* are numbers usually between 1 and 3.

The model parameters in Equations (1) through (7) were found by using a simplex algorithm to minimize L_1 given by Equation (9).

$$L_{1} = \sum_{1}^{N} \left(V_{p}^{obs} - V_{p}^{calc} \right) + \sum_{1}^{N} \left(\rho^{obs} - \rho^{calc} \right),$$
(9)

where V_p^{obs} , V_p^{calc} , ρ^{obs} , and ρ^{calc} are the sonic log compressional velocity, model calculated sonic compressional velocity, log density, and model calculated density, respectively. The units used in defining L₁ were m/s and Kg/m³, so that the velocity and density had approximately equal numerical magnitude, and hence equal weight in the value of L₁. Because the observation wells used in the crosswell surveys did not have full logging suites (no sonic logs), the nearest well (1,000 ft away) with a full suite of logs was used. Electrical parameters in Equation (8) were determined by a regression using the OB-C1 σ , ϕ , and S_w logs. Predicted V_p, ρ , and 1/ σ compared to the observed logs are shown in Figure 2, with the model parameters determined from the regressions listed in Table 2.

Parameters listed in bold type in Table 2 (critical porosity, oil API gravity, brine salinity, and temperature) were held fixed in the regression. These values, with the exception of critical porosity, came from direct measurement. Although we are not interested in the model parameters *per se* (we are only interested in the model's ability to predict V_p , V_s , and ρ , given reservoir parameters), note that their values are quite realistic. The gas density G is very close to that of methane. Estimated shear modulus and grain density of the diatomite grains is very close to the values of 18 (GPa) and 2.3 (g/cc) estimated by Wang (2001). Bilodeau (1995) measured an average grain density of



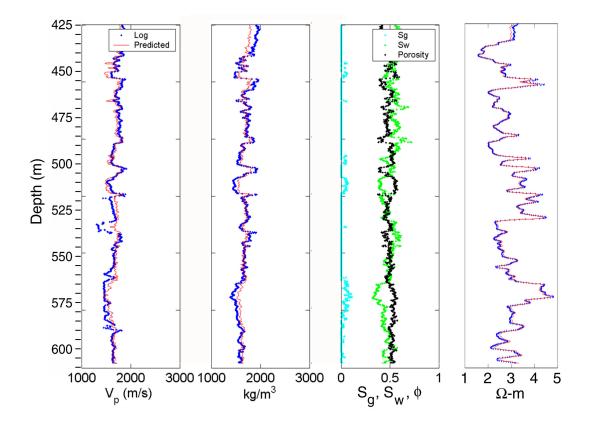


Figure 2. Rock properties model uses logged porosity (black), water saturation (green) and gas saturation (light blue) as inputs in a multi-parameter regression to predict the velocity (left panel), density (second from left panel) and electrical resistivity (right panel). Measured velocity, density, and resistivity are shown in blue; model predicted values are shown in red.

fluid conductivity, respectively, on the same samples. A value of critical porosity was determined by a set of minimizations of Equation (9) where ϕ_0 was varied between 0.5 and 0.7, all of which reached essentially the same value of L₁. The value of ϕ_0 was

chosen that resulted in values of G_{grain} and ρ_{grain} that were closes to those estimated by Wang (2001).

In Table 1, only one parameter, "gas correction", is listed under the Gassmann fluid substitution column. In addition, the Gassmann formula uses the dry-frame modulus as well as the fluid bulk moduli derived from the Batzle and Wang (1992) relations. However, we found that to fit the observed velocity in areas where the gas saturation was non-zero, the gas effect had to be reduced. The overestimation of the gas effect on fluid bulk modulus by the Wood's mixing law, Equation (6), has been observed by Brie et al. (1995). A better match between predicted and observed velocity could be achieved by a simple correction to the gas term in Equation (6), yielding a modified equation

$$K_{fluid} = 1.0 \left[\frac{G_c * (S_{hcg} / K_{hcg}) + (S_w / K_{brine}) +}{((1 - S_{hcg} - S_w - S_{C02}) / K_{oil}) + (S_{CO2} / K_{CO2})} \right]$$
(10)

where G_c is the gas correction listed in Table 1.

The pressure prediction capability of the model was validated by comparison to measurements made by Wang (2001) on core samples of diatomite from Lost Hills. Figure 3 shows the measured compressional velocity for vertical and horizontal propagation. These measurements show a horizontal-to-vertical velocity anisotropy of 1.047 that varies slightly as a function of pressure. We will come back to the velocity anisotropy when we consider the velocity inversion of the crosswell

data.

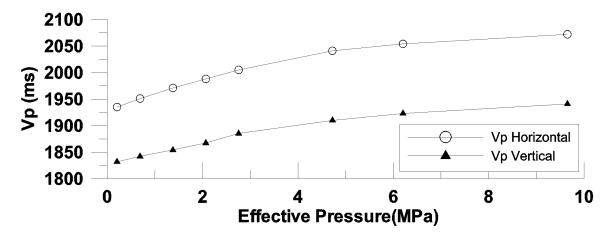


Figure 3 Vertical and horizontal compressional velocity as a function of effective pressure measured on Lost Hills Diatomite core by Wang (2001). Core was saturated with 19 API oil and 200,000 ppm brine (50-50 ratio) at 22.7 C.

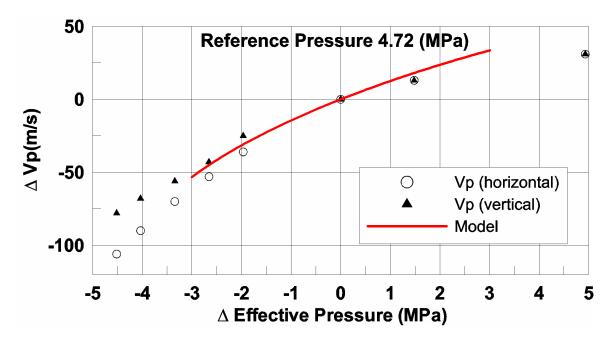


Figure 4. *Predicted velocity change as a function of change in effective pressure compared to laboratory measurements on Lost Hills diatomite core samples.*

Figure 4 presents the data from Figure 3 recast as velocity changes as a function of pressure changes at a reference pressure of 4.7 MPa, the average effective pressure in the reservoir at the start of CO_2 injection. For expected decreases in effective pressure (increases in pore pressure) in the range 0 to 3 MPa from the reference pressure, the rock properties model predictions are within a few percent of the lab measurements vertical velocity. The rock-properties model is derived from log sonic measurements dominated by vertical propagation along the borehole, so the correspondence to the vertical core measurements is expected. For changes in effective pressure above the reference pressure, the lab measurements show a change in the slope of the curve, with the quality of the fit between model and lab data decreasing. The difference in this region is probably associated with pore crushing in the lab samples not accounted for in the rock-properties model.

The estimates of the time-lapse changes in geophysical parameters derived from inversion of the observed geophysical data, as described in the following sections, are used with the rock-properties model described by Equations (1)–(5), (7), (8), and (10), with constants listed in Table 2, to calculate time-lapse changes in reservoir parameters.

INTEGRATED TIME-LAPSE GEOPHYSICAL IMAGES

The algorithms, assumptions, starting models, and amount of incorporated *a priori* information all greatly affect the velocity and conductivity models resulting from inversion. Inversions of the individual data sets done separately, without any mechanism for linking the models, produces images of V_p , V_s , and σ with little spatial correlation.

Since we assume that the changes in reservoir parameters affect all of the geophysical parameters (albeit in different ways), we expect a certain degree of spatial correlation between changes in the different geophysical parameters. This assumption acts as a constraint on the possible solutions. In this experiment, sonic logs were not run in OB-C1 or OB-C2, but conductivity logs were run in both wells. The strategy we adopted to maximize the spatial correlation between velocity and conductivity images was to begin with the EM data, where the greatest amount of *a priori* information existed, and then use the conductivity image to produce a starting V_p model, followed by producing a starting V_s model from the final V_p model. Conductivity logs were used to build the starting conductivity model for the EM inversion. The EM inversion algorithm is described by Newman (1995). We chose to use the conjugate gradient algorithm of Jackson and Tweeton (1996) for the travel-time tomography because the final model is sensitive to the initial model and is perturbed from the starting model only as much as needed to fit the observed data. Both EM and seismic inversions models were parameterized by 3 m cells.

EM inversion for the data at initial conditions (late August 2000 before CO₂ injection) was started from a model built by laterally interpolating the conductivity logs between the OB-C1 and OB-C2 wells. The final inversion model from this data was then used as the starting model for the inversion of the April 2001 data. The difference of the two inversions provides the time-lapse change in conductivity shown in Figure 5c. A high degree of correlation exists between the permeability log from the injector and the areas where the largest decrease in conductivity occurs. The correlation between high permeability and large changes in conductivity (water saturation) is expected.

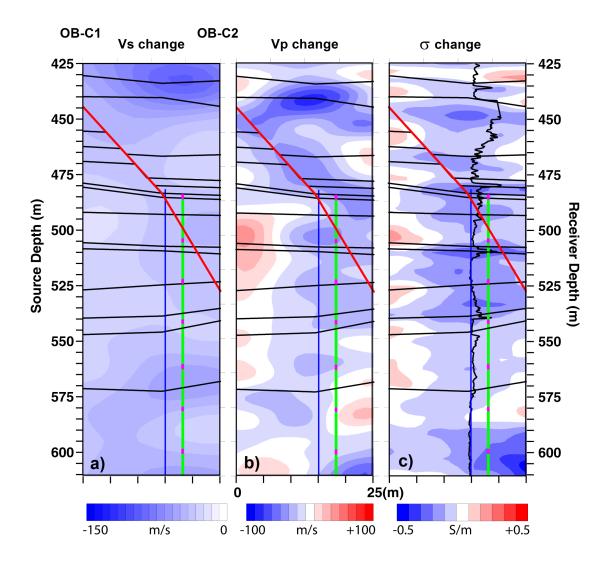


Figure 5. Time-lapse changes in (a) shear velocity, (b) compressional velocity and (c) electrical conductivity. The EM images were used to construct starting models for the V_p inversions; the resulting V_p images were used to construct starting models for the V_s inversions. Major unit boundaries are shown as black sub-horizontal lines, estimated location of previous water injection fracture is shown as a vertical blue line, estimated location of the CO₂ injection fracture is shown as a vertical green line, perforation intervals for CO₂ injection are shown as magenta dots, and the mapped location of a fault zone is shown as the red diagonal line. The permeability log in the out-of-plane CO₂ injection well (11-8WR) is shown in black on panel (c).

Next, the conductivity models from the two inversions were converted to compressional velocity. Values of ϕ , P_{eff}, P_{pore}, and S_{heg}, based on averages from the log data, were used with regression derived parameters (Table 2) to calculate V_p and σ as a function of S_w, using Equations (1)-(5), (7), (8), and (10). A linear regression between the calculated V_p and σ was done to provide a function for converting σ to V_p. The converted σ models were then used as initial models in the inversion of the V_p travel-time data to produce the change in V_p shown in Figure 5b. In addition to a decrease in V_p in the region around the estimated locations of the old water and new CO₂ injection fracture locations, there are decreases in V_p that align with the upper section of the mapped fault, implying that P_{pore} increases along the upper section of the fault. Since there are few conductivity changes associated with the fault, these results indicate that pressure changes occur along the fault zone without significant changes in water saturation at the time of the experiment.

The largest σ , V_p , and V_s changes occur in a region bordered by the old water injection fracture and the new CO₂-injection fracture. The water injection was ongoing for more than six years and likely produced a high-permeability damage zone that has been intersected by the newer CO₂ fracture. We speculate that this has produced a relatively high permeability zone in the region between and surrounding the two ideal fracture locations. Both the conductivity and V_p change sections (Figure 5c and 5b) show an increase in conductivity and V_p near the OB-C1 and OB-C2 wells. This is caused by an increase in water saturation, as shown in the relogging of the wells in January 2001. Water moving outward and away from the high permeability injection zone as CO₂ is injected causes a "rind" of increased S_w surrounding the volume affected by CO_2 . The volume of rock affected by CO_2 injection will have reduced water content as either CO_2 fills the pore space or oil absorbs CO_2 and swells, expelling water. This volume will have a surrounding "rind" of increased water saturation.

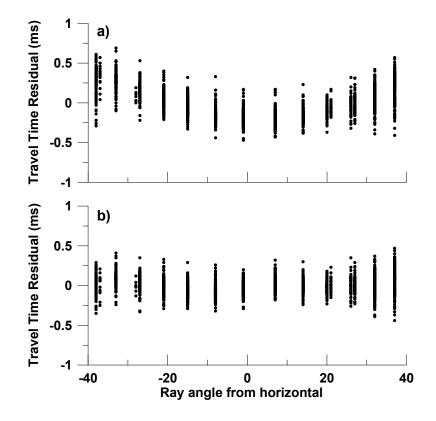


Figure 6. Travel time residual (observed – calculated) vs. ray angle from horizontal. Panel (a): homogeneous halfspace starting model with no anisotropy or dip of the velocity field. Panel (b): homogeneous halfspace starting model with $V_{horizontal}/V_{vertical} = 1.05$ and symmetry axis 7 degrees from vertical.

The algorithm (Jackson and Tweeton, 1996) used to produce the velocity tomograms shown in Figure 5 allows setting a constant velocity anisotropy and a constant dip of the anisotropy symmetry axis for the entire cross section. In a series of tomographic inversions, values of the horizontal/vertical velocity and the dip of the symmetry axis were varied between 0.9 and 1.1 and -10 to +10 degrees, respectively. The final values of 1.05 and 7 degrees from vertical (respectively) used in Figure 5 produced the flattesttravel-time misfit-versus-ray angle scatter plot with the minimum RMS data misfit. Figure 6a shows the travel-time residual plot for a V_p model without anisotropy, and Figure 6b shows the residual for the final V_p model shown in Figure 5b. The horizontal-to-vertical velocity ratio of 1.05 from the crosswell seismic tomograms compares remarkably well with the value of 1.047 from core measurements shown earlier in Figure 3. In addition, the structural dip of the reservoir units in the plane of the crosswell experiment is 7 degrees.

The starting models for the V_s inversions were converted from the final V_p sections using a V_p/V_s ratio derived from the rock properties model. The final V_s models were differenced to produce the change in V_s section shown in Figure 5a. The V_s change section is much smoother than either the conductivity or V_p change sections. This results partially from the lower frequency content in the shear-wave data. Shear-wave data were acquired using an orbital vibrator source with a center frequency of 500 Hz, whereas the compressional wave data were acquired using a piezoelectric source with a center frequency of 2,000 Hz. The V_s change section is also smoother because V_s is relatively insensitive to changes in water saturation (which have high spatial variability) and more sensitive to pressure changes (which have much lower spatial variability). Even with the smoother spatial changes in V_s we see a correlation with V_p and conductivity changes. In particular, the zone along the fault shows a decrease in V_s , lending support to our interpretation that pore pressure is increasing along the fault zone.

THE EFFECTS OF GAS ON SEISMIC VELOCITY AND DENSITY

The goal is to predict changes (Δ) in reservoir pressure, fluid saturations, and the amount of absorbed CO₂ in the oil as the CO₂ flood proceeds. We assume that the porosity remains constant over the time of the experiment. To use the rock-properties model to predict changes in reservoir parameters from changes in geophysical parameters, we must define certain values for reference parameters with respect to which the changes will be computed. In particular, reference water saturation (S_w) and porosity (ϕ) of 0.5 and 0.52, respectively, are taken from the averages in the OB-C1 well over the reservoir interval prior to CO₂ injection. The reference pore pressure (P_{pore}) is taken from a history-matched flow simulation model at the beginning of CO₂ injection. The reference effective pressure (P_{eff}) on the rock frame for seismic velocity calculations is calculated from the integrated density log minus P_{pore}. We will consider the sensitivity of our predictions to values of the reference parameters below.

Both hydrocarbon gas and CO₂ in the reservoir affect the seismic velocities through three possible mechanisms:

- (1) by directly changing the bulk modulus of the composite fluid in the pore space as gas saturation changes (Equation 10).
- (2) By changing the bulk modulus of the oil as the amount of dissolved gas changes.
- (3) By changing the bulk density of the rock.

Equation (11), from Batzle and Wang (1992), gives the maximum amount of gas that can dissolve in oil expressed as a gas/oil ratio (R_G^{max}) as a function of pore pressure (P_{pore}), temperature in degrees Celsius (T), oil API gravity (API), and gas gravity (G_{grav}):

$$R_G^{\text{max}} = 2.03 G_{grav} \left[P_{pore} \exp(0.02878API - 0.00377T)^{1.205} \right]^{1.205} (11)$$

The gas/oil ratio is the volume ratio of liberated gas to remaining oil at atmospheric pressure and 15.6° C. Batzle and Wang (1992) also provide formulas for computing the velocity and density of oils with dissolved gas, which we have used in our calculations.

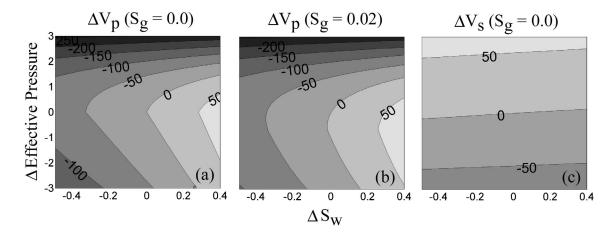


Figure 7. Change in velocity (m/s) as a function of change in effective pressure and water saturation at reference values of $S_w=0.5$, $S_{hcg}=0.0$, $\phi = 0.52$ and $P_{eff}=4.7MPa$. Panel (a) ΔV_p ($S_{hcg} = 0.0$) (b) ΔV_p ($S_{hcg} = 0.02$) (c) ΔV_s ($S_{hcg}=0.0$). The oil contains the maximum amount of dissolved hydrocarbon gas as a function of pressure for the parameters of the rock properties model given in Table 1.

An increase in the amount of dissolved gas in the oil, as measured by R_G, decreases both the bulk modulus and density of the oil. The bulk modulus is reduced more than the density, resulting in a decrease in the compressional velocity of the oil. Figures 7a and 7c show the calculated ΔV_p and ΔV_s using oil with the maximum amount of dissolved hydrocarbon gas as functions of ΔP and ΔS_w , at a reference point (reservoir just prior to CO₂ injection) where S_w , S_{hcg} , ϕ , and P_{eff} are equal to 0.5, 0.0, 0.52, and 4.7 (MPa), respectively. When S_{hcg} is non-zero and free gas exists, the behavior of ΔV_p with ΔP and ΔS_w changes markedly. Figure 7b shows ΔV_p for the same reference values as Figure 7a, but with $S_{hcg} = 0.02$. Equation (11) is used to compute the maximum amount of dissolved gas as a function of pressure. As Ppore increases above the reference pressure, R_{G}^{\max} increases, and we assume that *in situ* gas will dissolve into the oil up to R_{G}^{\max} . As the pressure decreases below the reference pore pressure, R_G^{max} decreases, and gas will come out of solution, thereby increasing Shcg above its reference value. This behavior is shown in Figure 8. At the reference pressure $S_{hcg} = 0.02$, as P_{pore} increases (- ΔP_{eff}), gas dissolves in the oil and S_{hcg} decreases until ΔP_{eff} reaches -0.6 MPa, when all of the gas has dissolved in the oil. If P_{pore} decreases (+ ΔP_{eff}), gas comes out of the oil and S_{hcg} increases. This increase in S_{hcg} with $+\Delta P_{eff}$ accounts for the sharp gradients in ΔV_p seen in the upper portion of Figure 7a and 7b.

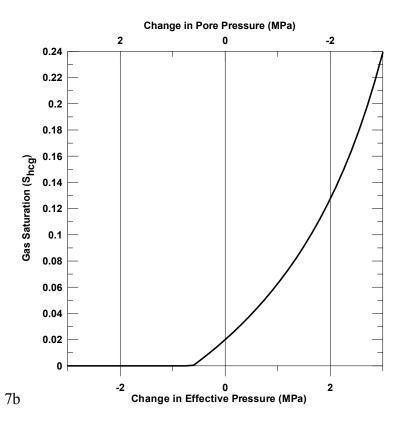


Figure 8. Change in S_{hcg} as a function of change in effective pressure for the model calculations shown in Figure 12. Reference $S_w=0.5$, $S_{hcg}=0.02$, $\phi = 0.52$ and P=4.7MPa.

Although developed for hydrocarbon gas in oil, Equation (11) can be used to predict R_G^{max} for CO₂ if the appropriate gas gravity is used. Chung et al. (1988) present experimental results for CO₂ solubility in 22 API gravity oil at 66.7° C over a range of pressures. The predicted values of R_G^{max} from Equation (11), using G_{grav} =1.51 for CO₂, are within 1% of the measured values over the range of P_{pore} found in the Lost Hills reservoir, between 800 and 1,500 psi. Based on this comparison, Equation (11) is appropriate for both hydrocarbon gas and CO₂. Data from Chung et al. (1988) also show that the amount of CO₂ that will dissolve in oil, at the relatively low temperatures and pressures in our experiment, is independent of the amount of hydrocarbon gas already

dissolved in the oil. We will make use of this fact later in our interpretation of the observed velocity changes.

The dissolution of gas into oil as P_{pore} increases produces two opposite effects on the composite fluid bulk modulus (Equation (10)) and hence the bulk velocity of the rock. An increase in P_{pore} causes R_{G}^{max} to increase, allowing more gas to dissolve in the oil, lowering K_{oil} while at the same time reducing S_{hcg} . From Equation (10), we see that a decrease in K_{oil} and in S_{hcg} acts in opposition on K_{fluid} . In addition, an increase in P_{pore} (decrease in P_{eff}) reduces K_{dry} and G_{dry} , which reduces the bulk velocity of the rock. When S_{hcg} is small, decreasing S_{hcg} to zero increases the bulk velocity of the rock more than lowering K_{oil} (by dissolving gas) decreases it. In contrast, decreasing P_{pore} increases K_{dry} which, by itself, would increase V_p . However, this effect is outweighed by the decrease in V_p caused by the increase in S_{hcg} as hydrocarbon gas comes out of solution from the oil. The net effect on V_p is seen in the upper halves of Figure 7a and 7b, where V_p decreases as P_{eff} increases (P_{pore} decreases).

PREDICTING TIME-LAPSE CHANGES IN RESERVOIR PARAMETERS

Before describing the process we have followed to estimate changes in fluid saturations, including *in situ* fluids and introduced CO_2 , we acknowledge that the multitude of possible interactions between changes in pressure, hydrocarbon gas, and CO_2 , as well as the effects on the oil from dissolved gas components, is too large to be uniquely determined from our geophysical measurements. We propose a procedure that makes use of a number of (what we consider to be) reasonable and most probable assumptions to estimate the change in CO_2 gas/oil ratio, ΔR_{CO2} , and CO_2 saturation, ΔS_{CO2} . The most critical assumption, supported by field reservoir engineers and operations staff (Perri, 2001), is that introduced CO_2 will dissolve in oil almost immediately after injection. Thus, we treat changes in the CO_2 gas/oil ratio as the primary mechanism for velocity reduction after changes in S_w and P have been accounted for.

EM data provide an independent estimate of ΔS_w . Electrical conductivity (σ) is a much simpler function of reservoir parameters than is velocity and can be described by Archie's law (Archie, 1942). Assuming ϕ is constant, $\Delta \sigma$ is only a function of ΔS_w and $\Delta \sigma_{\text{brine}}$. Because a water flood had been in effect for over 6 years at the start of CO₂ injection, we assume σ_{brine} has reached equilibrium between injected and native water and does not change. Therefore, conductivity changes are interpreted solely in terms of water saturation changes.

The process of converting the geophysical Δ images to Δ reservoir parameters begins with predicting ΔS_w between the wells from the $\Delta \sigma$ image, assuming that ϕ and σ_{brine} are constant. The predicted ΔS_w is used with the observed ΔV_s and the relation illustrated in

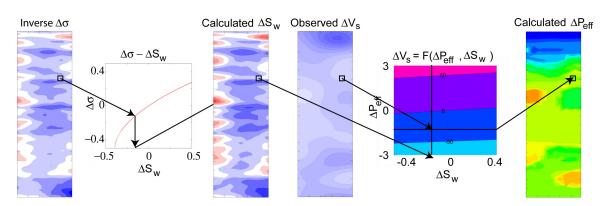


Figure 7(c) to predict ΔP_{eff} . The process is illustrated schematically in Figure 9. At

Figure 9: Schematic diagram of process of estimating ΔS_w from inverse $\Delta \sigma$ and then using the estimated ΔS_w with the observed ΔV_s to estimate ΔP_{eff}

this point, the predicted ΔS_w and ΔP_{eff} sections have required only the assumption that $\Delta \sigma_{brine}$ does not change appreciably. Going beyond this point to use the observed ΔV_p with the predicted ΔS_w and ΔP to predict changes in CO₂ saturation, along with changes in absorbed gases, requires more assumptions and becomes riskier.

The predicted ΔS_w and ΔP_{eff} are used to calculate the ΔV_p that results from ΔS_w and ΔP_{eff} alone, assuming S_{hcg} =0. Over the majority of the image plane, ΔS_w is negative, with the exception of small zones in the rind of water saturation which increase (as noted earlier). Predicted ΔP_{eff} is negative over the entire interwell section, thus producing a - ΔV_p . The residual change in velocity (ΔV_R) is defined by Equation (12):

$$\Delta V_R = \Delta V_p^{obs} - \Delta V_p^{calc} , \qquad (12)$$

where ΔV_p^{obs} is the observed change in V_p and ΔV_p^{calc} is the calculated change in V_p . We expect the injected CO₂ to decrease V_p in excess of the effects of ΔS_w and ΔP_{eff} by dissolving CO₂ in oil and possibly producing $S_{co2} > 0$. Figure 10 schematically represents the process of calculating ΔV_R and using this with the rock properties model to estimate ΔR_{CO2} .

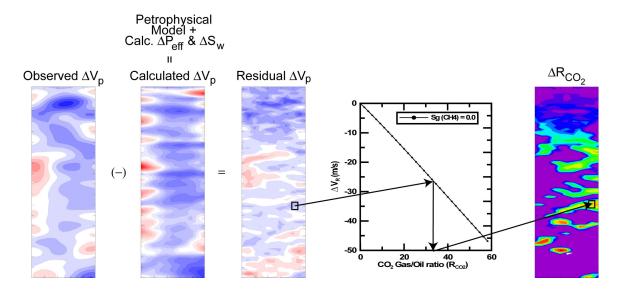


Figure 10: Schematic representation of combination of ΔV_R and rock properties model to predict ΔR_{CO2} .

On the other hand, a $+\Delta V_R$ can result if the assumption of no *in situ* hydrocarbon gas, S_{hcg}=0, is incorrect. This effect can be seen by comparing Figure 7a to 7b, where the presence of hydrocarbon gas reduces the change in V_p associated with a given ΔS_w and ΔP . As a consequence, a calculated ΔV_{p} , assuming S_{hcg} = 0 when S_{hcg} > 0, yields a calculated ΔV_p that is too large and hence a $+\Delta V_R$. However, if *in situ* hydrocarbon gas is present and has been accounted for in the calculation of ΔV_R , $+\Delta V_R$ can result if a P_{pore} increase causes hydrocarbon gas to dissolve in the oil, in which case S_{hcg} is reduced. The OB-C1 log shows the presence of hydrocarbon gas over certain intervals within the reservoir. We noted a strong correlation between depth intervals with a non-zero S_{hcg} and $+\Delta V_R$. Therefore, a two-step process was used to calculate ΔV_R . The first pass used $S_{hcg}=0$ as described. Next, sections of the image with $+\Delta V_R$ were recalculated assuming $S_{hcg} = 0.02$ (the average non-zero S_{hcg} in the reservoir interval). After the second pass calculation of ΔV_R , many of the areas that had $+\Delta V_R$ after the first pass calculations became negative, as ΔV_p^{calc} became less negative.

There are thus three regions of the ΔV_R section between the wells to interpret: (1) $S_{hcg} = 0$ and $\Delta V_R < 0$, (2) $S_{hcg} > 0$ and $\Delta V_R < 0$, and (3) $S_{hcg} > 0$ and $\Delta V_R > 0$. Regions of the crosswell section corresponding to $S_{hcg} = 0$ and $\Delta V_R < 0$ require an assumption about the partitioning of effects of free CO₂ and CO₂ dissolved in oil on $-\Delta V_R$. We chose to allow the maximum increase in R_{CO2} , as given by Equation (11), for the given ΔP_{pore} and ΔS_w . If the $+\Delta R_{CO2}$ does not completely account for the $-\Delta V_R$, then ΔS_{CO2} was calculated to account for the rest. For regions where $S_{hcg} > 0$ and $\Delta V_R < 0$, we assumed that the $+\Delta P_{pore}$ caused by injection would drive as much of the initial S_{hcg} into the oil as possible, followed by the same assumption about the partitioning of the $-\Delta V_R$ between $+\Delta R_{CO2}$ and ΔS_{CO2} as above. Regions where $S_{hcg} > 0$ and $\Delta V_R > 0$ were converted to $-\Delta S_{hcg}$.

Analysis shows that the oil is fully saturated with hydrocarbon gas at the ambient reservoir pressure. Therefore, we assume that the starting point for CO_2 absorption is oil with R_{hcg} at its maximum value for the given P_{pore} and T. As noted earlier, R_{CO2} and R_{hcg}

are essentially independent, so that the oil can absorb the amount of hydrocarbon gas and CO_2 up to their respective R_G^{max} , indicated by Equation (11). Because we lack an equation for calculating K_{oil} with two separate dissolved gases, we have assumed that Equation (13) is an adequate approximation of the bulk modulus of the oil:

$$K_{oil}^{co2+hcg} = K_{oil}^{dead} + \Delta K_{oil}^{co2} + \Delta K_{oil}^{hcg}$$
(13)

 K_{oil}^{dead} is the oil bulk modulus without any gas, $\Delta K_{oil}^{co^2}$ is the difference between K_{oil}^{dead} and the oil bulk modulus with CO₂ dissolved, and ΔK_{oil}^{hcg} is the difference between K_{oil}^{dead} and the oil bulk modulus with hydrocarbon gas dissolved.

Calculating $+\Delta R_{CO2}$ for regions where $\Delta V_R < 0$ is a simple linear interpolation between observed $-\Delta V_R$ and calculated $-\Delta V_R$ for a range of $+\Delta R_{CO2}$. If R_{CO2} reaches the maximum given by Equation (11), then the remaining observed $-\Delta V_R$ is used in a linear interpolation between calculated $-\Delta V_R$ over a range of $+\Delta S_{CO2}$ to calculate S_{CO2} .

The linear relation between ΔV_R and R_{CO2} is shown by the dotted line ($S_{hcg} = 0$) in Figure 11. In Figure 11, R_{CO2} increases from 0 to R_{CO2}^{max} from test number 1 to 34. From test number 34 to 41, S_{CO2} increases from 0 to 0.02, simulating the effect of progressively adding CO₂ that first dissolves in oil. After R_{CO2}^{max} is reached, CO₂ goes into the gas phase. Figure 9 also illustrates the effects of incorrectly assigning *in situ* hydrocarbon gas saturation. If $S_{hcg} > 0$ when $S_{hcg}=0.0$ is assumed, the estimated R_{CO2} will be low. On the other hand, if $S_{hcg}=0.0$ when $S_{hcg} > 0$ is assumed, the estimated R_{CO2} will be high. The error introduced by an incorrect S_{hcg} of 0.02 is approximately 15%.

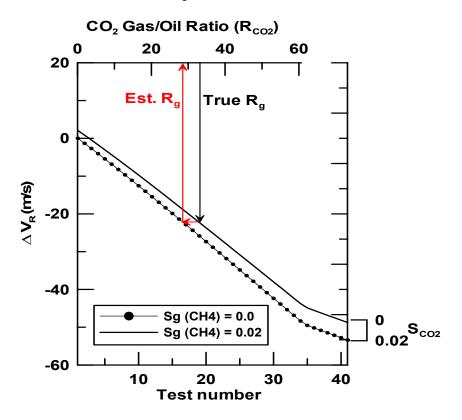


Figure 11. ΔV_R as a function of R_{CO2} with and without hydrocarbon gas saturation, solid line is $S_{hcg}=0.02$, dotted line is $S_{hcg}=0.0$. The presence of S_{hcg} causes ΔV_R to be less negative than if $S_{hcg}=0.0$.

Using this model, we have assessed the errors caused by incorrect values of the reference parameters S_w , ϕ , and $P_{ref.}$ The error expressed as a percent of the true value is plotted in Figure 12 for the same R_{CO2} values used in Figure 11. A 15% perturbation of the true reference values were used, which we feel covers the expected variation in these

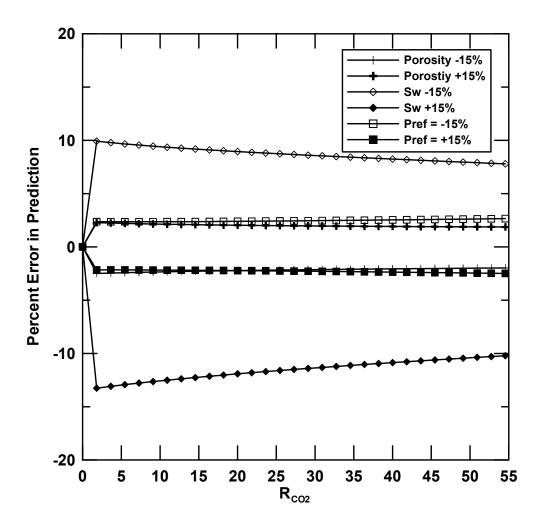
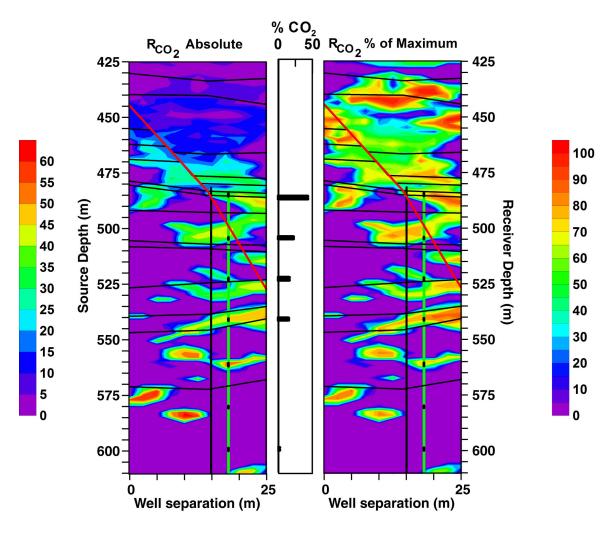


Figure 12. Error in predicted R_{CO2} as a percentage of the true value. Reference values of porosity (ϕ), water saturation (S_w) and effective pressure (P_{ref}) are in error by +-15%. The 15% range covers the expected variation in these parameters over the inter-well section.

parameters over the interwell section. The error response as a function of R_{CO2} is approximately symmetric for positive and negative perturbations in the reference parameters used. The assumed S_w has the largest effect, followed by the assumed effective pressure, with the assumed porosity having the smallest effect. Overall, the estimated R_{CO2} is most sensitive to S_{hcg} , since an error of 0.02 in S_{hcg} causes a comparable error of 15% in S_w , but S_{hcg} may vary by more that 0.02.

Figure 13 shows the calculated absolute R_{CO2} (left side) and R_{CO2} expressed as a percent of R_{CO2}^{max} (right side) generated from the geophysical parameter changes shown in Figure 5, using the two-step process described above. Effective pressure from a historymatched flow simulation model and integrated density log at the beginning of CO₂ injection was used as the reference pressure. The predicted R_{CO2} never reached R_{CO2}^{max} , so no S_{CO2} was needed to account for remaining $-\Delta V_R$. The predicted R_{CO2} shows a strong correlation with the location of perforation intervals (shown as black dots on the green CO_2 hydro-fracture line) that account for the majority of the injected CO_2 . The percentage of injected CO_2 going into each perforation in the 11-8WR well (Figure 1) is plotted in the center of Figure 13 and shows that the upper four perforations account for 95% of all the CO₂. Almost 50% of the CO₂ goes into the uppermost perforation. The location of this perforation corresponds to the large $+R_{CO2}$ associated with the fault zone and region above, indicating loss of substantial CO₂ into the upper portions of the reservoir. The second, third, and fourth perforations from the top account for roughly



another 45% of injected CO₂, with each perforation aligning with a laminar zone of

Figure 13. Predicted CO_2/oil ratio (R_{CO2}). Left side shows absolute R_{CO2} , right side shows R_{CO2} as a percent of the maximum value for the given pressure and temperature. Major unit boundaries are shown as black sub-horizontal lines, estimated location of the previous water injection fracture is shown as a vertical black line, estimated location of the CO_2 injection fracture is shown as a vertical green line, perforation intervals for CO_2 injection are shown as black dots on top of the CO_2 injection fracture, and the mapped location of a fault zone is shown as a red diagonal line.

 $+R_{CO2}$. The only poor correlation between injected CO₂ and predicted $+R_{CO2}$ occurs at the perforation at a depth of 1,850 ft. At this depth, a laminar $+R_{CO2}$ zone aligns with a perforation, but the injectivity log indicates little injected CO₂. A possible explanation for this zone of increased CO₂ is the down-dip CO₂ injector 12-7W. This injector lies along the same hydraulic fracture azimuth as the 11-8WR (Figure 1) and shows considerable CO₂ injection into the geologic unit that intersects our image plane at 1,850 ft depth.

The upper section of the fault (left side), where geologic-unit boundaries are offset, correlates with an increase in R_{CO2} , whereas the lower section (right side), where no displacements are mapped, does not. This is consistent with an increased permeability along portions of the fault that have significant movement compared to portions that do not. We interpret this image as indicating that CO_2 from the uppermost perforation is moving up dip along the fault zone and leaking into the high-permeability units above.

The image of R_{CO2} shown in Figure 13 has apparent higher vertical resolution of increased CO₂ zones compared to the geophysical anomalies shown in Figure 5. While there are zones of $-\Delta V_p$ associated with the same perforation intervals correlated with $+R_{CO2}$, there are additional areas of $-\Delta V_p$ above and below that do not correspond to $+R_{CO2}$. Because V_s is insensitive to the fluid substitutions (Figure 7c), we do not expect to see a correlation between ΔV_s and CO₂, either in the gas phase or dissolved in oil. Electrical conductivity changes will be related to changes in oil saturation through the change in S_w; these conductivity changes would also show a correlation to the displacement of water by oil, which may or may not be oil with dissolved CO_2 in it. Thus, although the $\Delta\sigma$ image (Figure 5) is correlated with the ΔV_p image, it also does not correlate with the injection intervals nearly as well as the derived R_{CO2} image of Figure 13. Overall, the R_{CO2} image has higher correlation with the injection intervals than the geophysical-change images and is also more horizontally stratified, as is the permeability structure of the formation.

CONCLUSIONS

We have used a rock-properties model, based on a close packing of spherical grains in conjunction with Gassmann's equation, to simulate the relationships between reservoir parameters of the Lost Hills diatomite and seismic compressional and shear velocities. A volumetric mixing law models bulk density. Parameters of the rock-properties model are derived by a simultaneous fitting of compressional velocity and density logs, using a simplex L1-norm minimization, given the observed porosity and fluid-saturation logs as well as measured pressure, temperature, and oil properties. Although the spherical grain model may not ideally represent the microscopic structure of the diatomite, the model accurately predicts the bulk seismic velocities and densities as a function of the fluid saturations, pressure, and porosity, as measured by log data and measurements made on core samples.

Calculations using the derived rock-properties model show that the rock bulk shear velocity primarily depends on pressure changes, with the effects of water saturation changes on shear velocity being of second order. Calculations also show that the

presence of even a small amount of hydrocarbon gas strongly affects the relationships between V_p and the reservoir parameters. The influence of gas on compressional velocity makes it impossible to separate the effects of changes in hydrocarbon gas saturation, CO_2 gas saturation, and the effects on the oil caused by dissolved CO_2 on V_p without additional independent information. Crosswell EM data was used to provide estimates of changes in electrical conductivity that are directly related to changes in water saturation, thus providing an estimate of the change in water saturation that is independent from the seismic data.

To predict quantitatively the location and amount of CO_2 in the crosswell image plane, the change of P-wave velocity is decomposed into the part that can be predicted by the estimated changes in water saturation and pressure and the part predictable by a change in CO_2 content. The process relies on the assumption that the CO_2 will first dissolve in the oil and will only enter the gas phase after the oil has absorbed the maximum amount of CO_2 possible for the *in situ* pressure and temperature conditions. Using this procedure, we have demonstrated that by combining seismically derived changes in compressional and shear velocity with EM-derived changes in electrical conductivity, estimates of pressure change, water saturation change, and CO_2 gas/oil ratio can be made in a complex reservoir containing oil, water, hydrocarbon gas, and injected CO_2 . The resulting predicted CO_2 /oil ratio, R_{CO2} , is better correlated with logged unit boundaries than are any of the images of changes in geophysical parameters. The size of the predicted CO_2 -rich zones correlate with the amount of CO_2 that enters the formation through each perforation. The predicted ΔR_{CO2} images indicate that a significant portion of the injected CO_2 is filling the upper portions of the section above the intended injection interval. These conclusions are validated by CO_2 injectivity measurements made in the 11-8WR Well.

While we have tried to produce quantitative estimates of the CO_2 in place by estimating the CO_2 /oil ratio, the values of this ratio depend on our assumptions about the partitioning of CO_2 between oil and gas phases. In addition, the assumed values of *in situ* hydrocarbon gas affect the estimates of the CO_2 /oil ratio, so that the absolute values of our estimates may be in error. The main advantage of the approach described in this paper is the decoupling of the effects of pressure and water saturation changes from those caused by CO_2 . This produces the improved spatial correlation between the estimated CO_2 /oil ratio and the CO_2 injectivity logs when compared to the geophysical change images.

This analysis relies on many assumptions that were required because the project was not originally designed to use this methodology. In future applications, the number of assumptions could be substantially reduced by design. In particular, considerable benefit could be drawn from repeat logging of the wells with a full suite of logs. This would provide control points for the ΔP , ΔS_w , ΔS_g , ΔV_p , ΔV_s , and $\Delta \sigma$, all of which would serve to greatly constrain the problem. Log measurements of the geophysical parameters would provide information for better starting models, with constraints on the velocity, density, and electrical conductivity at the well locations. Additionally, measurements of S_{CO2} and the amount of CO₂ dissolved in the oil would provide a basis for determining the partitioning of the residual velocity between the two, as well as eliminate the need to assume that all of the CO_2 dissolves in the oil before CO_2 gas is evoked as a mechanism of velocity change.

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REFERENCES

Archie, G. E., 1942, The electrical resistivity log as an aid in determining some reservoir characteristics: Trans. Am. Inst. Mech. Eng., **146**, 54-62.

Batzle, M. and Wang, Z., 1992, Seismic properties of pore fluids: Geophysics, **57**, 1396-1408.

Bilodeau, B. J., 1995, Determining water saturation in diatomite using wireline logs, Lost Hills Field, California: SPE29653, in Proceedings of the Western Regional SPE Meeting, Bakersfield, CA, 369-382.

Brie, A., Pampuri, F., Marsala, A. F., and Meazza, O., 1995, Shear sonic interpretation in gas-bearing sands: SPE30595, in Proceeding of the SPE Annual Technical Conference, Dallas, Tx, 701-710.

Chung, F. T. H., Jones, R. A., and Burchfield, T. E., 1988, Recovery of viscous oil under high pressure by CO₂ displacement: A Laboratory Study: SPE 17588, in Proceedings of the SPE Annual Technical Conference, Tianjin, China, 401-410.

Dvorkin, J., and Nur, A., 1996, Elasticity of high-porosity sandstones: Theory for two North Sea datasets: Geophysics, **61**, 1363-1370.

Gassmann, F., 1951, Elasticity of porous media: Uber die Elastizitat poroser Medien: Vierteljahrsschrift der Naturforschenden Gesselschaft in Zurich, Heft 1.

Goudswaard, J. C. M., ten Kroode, A. P. E., Snieder, R. K. and Verdel, A. R., 1998, Detection of lateral velocity contrasts by crosswell traveltime tomography: Geophysics, **63**, 523-533.

Hashin, Z., and Shtrikman, S., 1963, A variational approach to the elastic behavior of multiphase materials: J. Mech. Phys. Solids, **11**, 127–140.

Harris, J. M., Nolen-Hoeksema, R. C., Langan, R. T., Van Schaack, M., Lazaratos, S. K., and Rector, J. W., 1995, High-resolution crosswell imaging in a west Texas carbonate reservoir: Part 1- Project summary and interpretation: Geophysics, **60**, 667-681.

Jackson, M. J., and Tweeton, D. R., 1996, 3DTOM: Three-Dimensional Geophysical Tomography: US Dept. of the Interior, Bureau of Mines, Report of Investigation 9617.

Landro, M., 2001, Discrimination between pressure and fluid saturation changes from time-lapse seismic data: Geophysics, **66**, 836-844.

Lazaratos, S. K., Harris, J. M., Rector J. W. and Van Schaack, M., 1995, High-resolution crosswell imaging of a West Texas carbonate reservoir: Part 4- Reflection imaging: Geophysics, **60**, 702-711.

Magee, J. W., and Howley, J. A., 1994, A predictive model for the thermophysical properties of carbon dioxide rich mixtures: Research Report RR-136, Gas Processors Association, Tulsa, OK.

Nemeth, T., Normark, E., and Qin, F., 1997, Dynamic smoothing in crosswell traveltime tomography: Geophysics, **62**, 168-176.

Newman, G. A., 1995, Crosswell electromagnetic inversion using integral and differential equations: Geophysics, **60**, 899-911.

Perri, P., 2001, Personal communication.

Rector W. J. ed., 1995, Special issue: Crosswell methods: Geophysics, 60., No. 3.

Wang, Z., 2001, Personal communication.

Wang, Z., Cates, M. E. and Langan, R. T., 1998, Seismic monitoring of a CO₂ flood in a carbonate reservoir: A rock physics study: Geophysics, **63**, 1604-1617.

Wilt, M. J., Alumbaugh, D. L., Morrison, H. F., Becker, A., Lee, K. H., and Deszcz-Pan,M., 1995, Crosswell electromagnetic tomography: System design considerations and fieldresults: Geophysics, 60, 871-885.

Wood, A. W., 1955, A Textbook of Sound: The MacMillan Co., New York, 360pp.

TABLES

 Table 1: Definition of symbols.

Parameter Name	Symbol	Parameter Name	Symbol
Dry frame bulk modulus	K _{dry}	CO ₂ bulk modulus	K _{CO2}
Dry frame shear modulus	G _{dry}	Composite fluid bulk modulus	K _{fluid}
Effective dry frame bulk modulus	K _{eff}		
Effective dry frame shear modulus	G _{eff}	Grain density	ρ _{grain}
Saturated rock bulk modulus	K _{sat}	Oil density	Poil
Saturated rock shear modulus	G _{sat}	Brine density	ρ _{brine}
Grain bulk modulus	K _{grain}	CO ₂ density	ρсо2
Grain shear modulus	G _{grain}	Gas correction	G _c
Grain Poisson ratio	ν	Hydrocarbon gas density	Phcl
Critical porosity	ϕ_0	Rock bulk electrical conductivity	$\sigma_{\it bulk}$
Rock porosity	ϕ	Brine electrical conductivity	$\sigma_{\it brine}$
Number of grain contacts	l	Archie porosity exponent	m
Water (brine) saturation	Sw	Archie saturation exponent	n
Oil saturation	So	Pore pressure	P _{pore}
Hydrocarbon gas saturation	S _{hcg}	Effective Pressure	P _{eff}
CO ₂ saturation	S _{CO2}	Compressional velocity	V _p
Brine bulk modulus	K _{brine}	Shear velocity	Vs
Oil bulk modulus	K _{oil}	Gas/Oil ratio	R_{g}
Hydrocarbon gas bulk modulus	K _{hcg}	Gas gravity	G _{grav}

Table 2: Rock-properties model parameters by model constituent. Model parameters fixed in the regression of well log data are shown in bold type. Asterisk (*) indicates that lithostatic pressure was calculated as a function of depth using the integrated density log, and pore pressure was taken as hydrostatic. Effective Pressure = Lithostatic – Hydrostatic Pressure.

		Dry Frame	Gassmann's Fluid		Electrical	Regression	
Parameter Name	Symbol	Modulus	Substitution	Fluids	Conductivity	Values	
Grain Shear							
Modulus	G _{grain}	х				17.84 (Gpa)	
Grain Poisson Ratio	v	x				0.107	
Grain Density	ρgrain	x	Dry Rock			2.358 (g/cc)	
# of contacts/grain	l	х	К			3.68	
Effective Pressure	P_{eff}	х				*	
Critical Porosity	ϕ_0	x				0.55	
Pore Pressure	P_{pore}			x		*	
Oil API gravity	API			х		21.7	
Gas gravity	G _{grav}		Fluid K	x		0.585	
Brine Salinity	S			x		0.023 (PPM/10 ⁶)	
Temperature	Т			х		42 (C)	
Fluid conductivity	$\sigma_{\it brine}$				x	0.23 (S/m)	
Porosity exponent	т				х	-1.66	
Saturation exponent	n				х	-1.85	
Gas correction	Gc		X			0.0068943	
*	Litostatic pressure from integrated density logs, <i>P_{eff}</i> = Litho - Hydro static						