# EXPERIMENTAL AND KINETIC MODELING STUDY OF METHANOL IGNITION AND OXIDATION AT HIGH PRESSURE

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## Abstract

A detailed chemical kinetic model for oxidation of CH<sub>3</sub>OH at high pressure and intermediate temperatures has been developed and validated experimentally. Ab initio calculations and RRKM analysis were used to obtain rate coefficients for key reactions of CH<sub>2</sub>OH and CH<sub>3</sub>O (dissociation, isomerization, reaction with O<sub>2</sub>). The experiments, involving CH<sub>3</sub>OH/O<sub>2</sub> mixtures diluted in N<sub>2</sub>, were carried out in a high pressure flow reactor at 600–900 K and 20-100 bar, varying the reaction stoichiometry from very lean to fuel-rich conditions. Under the conditions studied, the onset temperature for methanol oxidation was not dependent on the stoichiometry, while increasing pressure shifted the ignition temperature towards lower values. Model predictions of the present experimental results, as well as Rapid Compression Machine (RCM) data from the literature, were generally satisfactory. The governing reaction pathways have been outlined based on calculations with the kinetic model. Unlike what has been observed for unsaturated hydrocarbons, the oxidation

pathways for CH<sub>3</sub>OH under the investigated conditions were very similar to those prevailing at higher temperatures and lower pressures. At the high pressures, the modeling predictions for onset of reaction were particularly sensitive to the CH<sub>3</sub>OH+HO<sub>2</sub> $\Rightarrow$ CH<sub>2</sub>OH+H<sub>2</sub>O<sub>2</sub> reaction.

## Keywords

Methanol; oxidation; ignition; high pressure; flow reactor experiments; ab initio calculations

## 1. Introduction

Oxygenated organic compounds have been proposed as alternative fuels to diesel in order to improve the fuel properties and reduce particulates and  $NO_x$  emissions [1]. Methanol, which is the simplest alcohol, has been used in reformulated gasolines with significant success, and recently has also received attention for diesel engine applications. Its range of production sources and its chemical structure, with a high oxygen content and no C-C bonds, causes methanol to be one of the most important oxygenated additives. For these reasons, a number of studies have addressed the use of methanol-diesel blends in diesel engines [2-5]. Most of the authors conclude that the addition of methanol reduces the emissions of particulates and  $NO_x$ , even though some reported results seem contradictory.

Laboratory experiments allow more controlled conditions compared to tests in engines; methanol oxidation has been studied in flames [6-8], stirred reactors [9], static reactors [10, 11], flow reactors [12-15], Rapid Compression Machines [16], and shock tubes [17]. Of these studies, only three have been carried out at elevated pressures, i.e. 10 bar [9], 20 bar [13] and 30 bar [16], respectively. The aim of the present work is to investigate methanol oxidation at even higher pressures, which are of interest for combustion devices such as diesel engines. To achieve this, experiments were conducted

in a flow reactor at pressures of 20-100 bar and 600-900 K, as a function of stoichiometry. In addition, high pressure ignition delays, obtained in a Rapid Compression Machine [16], were taken into consideration. The results were analyzed in terms of an updated detailed chemical kinetic model. Selected reactions, potentially important at high pressure, were characterized from ab initio calculations over a wide range of pressure and temperature.

## 2. Experimental

Experiments have been conducted in a high pressure laminar flow reactor, described in detail elsewhere [18]. The quartz reactor, which has been designed to ensure a good plug approximation and a well-defined reaction volume (inside diameter 8 mm, outside diameter 10 mm, length 1545 mm), was enclosed in a stainless steel tube that acted as a pressure shell. A pressure control system delivered nitrogen to the shell-side of the reactor to obtain a pressure similar to that inside the reactor. The steel tube was placed horizontally in an oven with three individually controlled electrical heating elements that produced an isothermal reaction zone ( $\pm$ 5 K) of approximately 40 cm. The reactor temperature was monitored by a type K thermocouple ( $\pm$ 2.2 K) located inside a steel thermo pocket placed in the pressure shell. Reactant gases were fed into the system through high pressure differential mass flow controllers.

A liquid feeding system, described in greater detail elsewhere [19], was used to supply the methanol. The methanol was pressurized by an HPLC pump and its flow was controlled by a liquid mass flow controller. The liquid was evaporated at 250°C. All the reactants were mixed before entering the reactor. Gas cylinder concentration certification was  $\geq$  99.998% for O<sub>2</sub> and N<sub>2</sub>. The methanol, which was supplied by Sigma-Aldrich, had a purity of  $\geq$  99.8%. The reactor pressure was monitored upstream of the reactor by a differential pressure transducer and controlled by a pneumatic pressure valve positioned after the reactor. At the outlet of the reactor, the pressure was reduced down to atmospheric pressure prior to analysis. A gas chromatograph, equipped with TCD and FID detectors, was used to measure CH<sub>3</sub>OH, C<sub>2</sub>H<sub>5</sub>OH, CH<sub>4</sub>, O<sub>2</sub>, CO, and CO<sub>2</sub>. The uncertainty of the measurements was estimated to be  $\pm$ 5%.

## 3. Chemical kinetic model

The mechanism used in the present study is based on the atmospheric pressure methanol schemes by the authors [14, 15], drawing also on previous high pressure work on oxidation of CO and CH<sub>4</sub> [18, 20, 21]. The methanol subset has been updated with results from recent theoretical work [22-24] and from ab initio calculations and RRKM analysis made as part of the present work on reactions of CH<sub>2</sub>OH and CH<sub>3</sub>O (dissociation, isomerization, reaction with  $O_2$ ). Table 1 lists the reactions in the methanol subset; in the following, the numbering of the reactions corresponds to the numbers in this table. The full reaction mechanism is available as Supplemental Material.

## 3.1 Ab initio calculations

In the present work, the rate constants for the CH<sub>2</sub>OH/CH<sub>3</sub>O dissociation / isomerization system and for the CH<sub>3</sub>O+O<sub>2</sub> reaction have been calculated by ab initio methods. Because of the occurrence of significant spin-contamination, coupled-cluster approaches were employed. In brief, geometries and frequencies (scaled by 0.954) of reactants and transition states were derived with CCSD/6-311G(d,p) using the Gaussian 09 program [26], and are summarized in the Supplemental Material. For the reactions CH<sub>2</sub>OH = CH<sub>2</sub>O+H (R11), CH<sub>3</sub>O = CH<sub>2</sub>O+H (R19), and CH<sub>3</sub>O = CH<sub>2</sub>OH (R20), energies were obtained with the Molpro09 program [27] at the CCSD(T)/aug-cc-pVnZ levels of theory with n=3-5, and extrapolated as exp(-n) to the infinite basis set limit. For the larger CH<sub>3</sub>O+O<sub>2</sub>  $\Rightarrow$  CH<sub>2</sub>O+HO<sub>2</sub> (R27) system, energies were assessed via the CBS-QB3 method [28] which approximates this extrapolated limit. Table 2 lists the energies of the stationary points in the C/H<sub>3</sub>/O<sub>3</sub> system derived with coupled cluster theory.

Previous studies of the four reactions considered here provided preliminary structures for the transition state (TS) optimizations [29, 30]. The thermochemistry of methoxy and hydroxymethyl radicals is complicated by significant anharmonicity, coupling between normal modes, spin-orbit splitting and Jahn-Teller distortion, which have been analyzed in some detail by Ruscic et al. [31]. For CH<sub>3</sub>O we used the empirical electronic partition function based on splitting of the <sup>2</sup>E ground electronic state into  ${}^{2}E_{3/2}$  and  ${}^{2}E_{1/2}$  components, and a rotational symmetry number based on the average C3V symmetry, although a Jahn-Teller-distorted Cs symmetry was used for the energy calculations. We found this model reproduced the tabulated thermochemistry [31] reasonably well: the error in entropy is -0.7 J K<sup>-1</sup> mol<sup>-1</sup> at 298 K increasing to -4.1 J K<sup>-1</sup> mol<sup>-1</sup> at 1200 K, while the error in H(T)-H(0 K) varies from 0.2 to 2.5 kJ mol<sup>-1</sup> over this range of temperature. We treated CH<sub>2</sub>OH as having a hindered OH rotor separable from the other modes which were treated as harmonic oscillators. This is certainly incorrect [31] and neglects anharmonicity and coupling of this rotor to the other motions. The impact of the seriously incorrect *ab initio* vibrations computed at 451 and 711 cm<sup>-1</sup> instead of 234 and 420 cm<sup>-1</sup> is, fortuitously, to compensate for the neglected coupling with the result that the computed entropy is too small by the amount of 3.7 J K<sup>-1</sup> mol<sup>-1</sup> decreasing to 1.5 J K<sup>-1</sup> mol<sup>-1</sup> over 298 – 1200 K, and H(T)-H(0 K) is underestimated by 0.5 kJ mol<sup>-1</sup> at 298 K and overestimated by 1.1 kJ mol<sup>-1</sup> at 1200 K. These are modest errors.

RRKM and transition state theories as implemented within the Multiwell program [32] were applied to derive kinetics over 600-1200 K and 0.01-100 bar (see Table 1). The Multiwell program allows for quantum mechanical tunneling estimated from an Eckart potential. For the RRKM calculations, we assumed Lennard-Jones parameters of 417 K and 3.69 Å, taken from Chemkin, and an exponential down energy transfer model with  $\alpha = 350 \text{ cm}^{-1}$ .

The potential energy diagram summarized in Fig. 1 indicates that unimolecular decomposition of CH<sub>3</sub>O energetically favors CH<sub>2</sub>O formation over isomerization to CH<sub>2</sub>OH. In fact the isomerization also has a rather tight TS, with an A factor at 800 K and the high-pressure limit of  $2.4 \times 10^{13}$  s<sup>-1</sup> vs  $1.6 \times 10^{14}$  s<sup>-1</sup> for k<sub>19,∞</sub>. For this reason most unimolecular decomposition of CH<sub>2</sub>OH yields CH<sub>2</sub>O+H (A factor 7.8 x  $10^{13}$  s<sup>-1</sup>) rather than CH<sub>3</sub>O (A factor 5.6 x  $10^{12}$  s<sup>-1</sup>), even though their energies are similar. Any CH<sub>3</sub>O formed by isomerization of CH<sub>2</sub>OH initially has sufficient excess energy to further react on to CH<sub>2</sub>O+H. This means that isomerization is a minor channel and does not maintain equilibrium between the two CH<sub>3</sub>O isomers.

Figure 2 shows an Arrhenius plot for the reactions  $CH_2OH \Rightarrow CH_2O+H$  (R11),  $CH_3O \Rightarrow CH_2O+H$ (R19), and  $CH_3O \Rightarrow CH_2OH$  (R20). Of these steps, only the rate constant for  $CH_3O \Rightarrow CH_2O+H$ (R19) has been measured experimentally. Hippler et al. [33] derived a value for the high pressure limit of (R19) from experiments covering pressure and temperature ranges of 1–90 bar and 680–810 K. Their value for  $k_{19,\infty}$  is about a factor of two lower than the one determined in the present work. Presently, we are not certain about the source of this discrepancy. It may partly be attributed to the fact that they neglected tunneling. According to our calculations, tunneling increases  $k_{19,\infty}$  by 30-40% in the 600-800 K range, and it may also affect the falloff behavior.

For the  $CH_3O+O_2$  reaction (R27), an Arrhenius plot is shown in Fig. 3. While the low-temperature data for this reaction are largely consistent, the behavior at elevated temperatures has been controversial. From experiments, Wantuck et al. [42] reported a nonlinear Arrhenius behavior, with

the rate constant increasing strongly above 600 K. However, their results have been questioned due to possible interference by thermal dissociation of  $CH_3O$  (R19) at the higher temperatures. For this reason, Baulch et al. [34] chose to disregard the data from Wantuck et al. in their evaluation of  $k_{19}$ . It is interesting to note that our theoretical rate constant (denoted by the solid line in Fig. 3) supports the observation of Wantuck et al. of a nonlinear Arrhenius behavior, even though the temperature dependence is smaller than in their determination. Our rate constant is in good agreement with the literature data, but significantly higher than the recommendation of Baulch et al. [34] at elevated temperatures (Fig. 3). This has some implications for the modeling in the present work.

#### 3.2 Reaction mechanism

The methanol subset of the reaction mechanism is listed in Table 1. The main methanol consumption pathways under the present conditions are hydrogen abstraction reactions with H and OH. It is now well established that these steps mainly yield hydroxymethyl, while methoxy is only a minor product. For the CH<sub>3</sub>OH+H reaction,

$$CH_{3}OH+H \rightleftharpoons CH_{2}OH+H_{2} \tag{R2}$$

$$CH_3OH + H \rightleftharpoons CH_3O + H_2 \tag{R3}$$

the recent rate constants of Moses et al. [22] have been preferred, while for CH<sub>3</sub>OH+OH,

$$CH_{3}OH+OH \rightleftharpoons CH_{2}OH+H_{2}O \tag{R6}$$

$$CH_{3}OH+OH \Rightarrow CH_{3}O+H_{2}O \tag{R7}$$

we use the values recommended by Rasmussen et al. [15]. The reactions of CH<sub>3</sub>OH with HO<sub>2</sub>,

$$CH_3OH + HO_2 \rightleftharpoons CH_2OH + H_2O_2 \tag{R8}$$

$$CH_3OH + HO_2 \Rightarrow CH_3O + H_2O_2 \tag{R9}$$

and  $O_{2}$ ,

$$CH_3OH + O_2 \rightleftharpoons CH_2OH + HO_2 \tag{R10}$$

are important for ignition, but in the past kinetic models have had to rely on rough estimates for their rate constants. Two recent high-level theoretical studies offer rate constants for these steps with a strongly improved accuracy. For the  $CH_3OH+HO_2$  reaction, we adopt the results of Alecu and Truhlar [24], while the rate constant for  $CH_3OH+O_2$  is drawn from the work of Klippenstein et al. [23].

It is important to distinguish between  $CH_2OH$  and  $CH_3O$ , because the reactivity of these isomers is different. The  $CH_2OH$  isomer has a fairly high thermal stability and under fuel-lean conditions and not too high temperatures  $CH_2OH$  reacts preferably with  $O_2$  to form  $HO_2$ ,

$$CH_2OH+O_2 \rightleftharpoons CH_2O+HO_2$$
 (R18)

The CH<sub>3</sub>O isomer dissociates easily in a reaction that forms atomic hydrogen,

$$CH_3O \Rightarrow CH_2O+H$$
 (R19)

Since the HO<sub>2</sub> radical is comparatively unreactive, the partitioning among CH<sub>3</sub>O and CH<sub>2</sub>OH is important for the methanol oxidation rate. While the values of  $k_{18}$  and  $k_{19}$  are well established, the rate constants for dissociation of the hydroxymethyl radical, CH<sub>2</sub>OH  $\Rightarrow$  CH<sub>2</sub>O+H (R11), for the isomerization of methoxy to hydroxymethyl radical, CH<sub>3</sub>O  $\Rightarrow$  CH<sub>2</sub>OH (R20), and for methoxy reacting with oxygen, CH<sub>3</sub>O+O<sub>2</sub>  $\Rightarrow$  CH<sub>2</sub>O+HO<sub>2</sub> (R27), have been in dispute in the past, as discussed above. We consider the rate coefficients derived in the present work for these steps to be more reliable than data in the literature and recommend that they are used in kinetic modeling.

## 4. Results and Discussion

Oxidation of methanol has been studied in a flow reactor in the temperature range 600-900 K and in the pressure interval 20-100 bar as a function of excess air ratio ( $\lambda$ =0.25-16). Table 3 lists the conditions of the experiments. They were conducted under dilute conditions with the balance made up of N<sub>2</sub>. The constant flow of 5000 mL/min (STP) resulted in gas residence times in the isothermal zone being a function of the reaction temperature and pressure, t<sub>r</sub>(s)=65.85 ·P(bar)/T(K). The experimental results were analyzed in terms of the detailed chemical kinetic model described above. All model calculations were performed using SENKIN, the plug flow reactor code that runs in conjunction with the Chemkin-II library [43, 44].

Figures 4-6 compare experimental and predicted results for CH<sub>3</sub>OH, CO and CO<sub>2</sub>, obtained under reducing conditions ( $\lambda$ =0.25) as a function of temperature at pressures of 20, 50, and 100 bar, corresponding to sets 1-3 in Table 3. The scatter in the data for the methanol concentration below 700 K indicates that the liquid fuel feeding is not completely stable over time. However, the fluctuations are mostly within 10% of the nominal inlet concentration for CH<sub>3</sub>OH.

The results show that increasing pressure tends to lower the onset temperature for methanol oxidation, which changes from about 775 K at 20 bar to 700 K at 100 bar. It should be noted that the shift observed in the figures is partly caused by the increase in residence time at higher pressure (Table 3), but increased pressure also promotes ignition. At the higher temperatures, the increase in pressure serves to inhibit slightly the consumption of CH<sub>3</sub>OH. Methane and ethanol were also detected among the products at higher temperatures, but only in very small amounts. The modeling predictions are in satisfactory agreement with the experimental results, even though the CO concentrations are slightly overpredicted at the higher temperatures.

Figures 7 and 8 show results obtained at stoichiometric and oxidizing conditions, respectively, at 100 bar. Comparison of figures 6-8 shows that the temperature for onset of reaction, approximately 700 K, is almost independent of stoichiometry. As would be expected, the oxygen concentration has a pronounced effect on the consumption of CH<sub>3</sub>OH and formation of CO and CO<sub>2</sub>. While for fuel-rich conditions ( $\lambda$ =0.25), the exit CO concentration increases gradually with temperature, it reaches a peak value at around 750 K for stoichiometric ( $\lambda$ =1.0) and fuel-lean ( $\lambda$ =16) conditions. At this temperature, the methanol has been fully consumed and no longer competes for OH radicals. As the temperature increases further, the CO<sub>2</sub> concentration would be expected to increase as well due to the oxidation of CO to CO<sub>2</sub>. However, it is noteworthy that even under very lean conditions (Fig. 8), as much as 800 ppm of CO is found in the exit stream from the reactor. The model captures the effect of stoichiometry well and also predicts the fairly high CO level observed even at lean conditions. Actually, according to the model there is a small local minimum in CO at about 800 K under both stoichiometric and lean conditions (Figs. 7 and 8), but this cannot be confirmed from the experimental results.

The main reaction routes for methanol oxidation under the present conditions are shown in Fig. 9. Methanol is consumed by reaction with OH (R6/R7) and, under fuel-rich conditions, H (R2). These reactions yield hydroxymethyl and methoxy radicals, with the former radical being predominant. As expected [13-15], the hydroxymethyl radical mainly react with oxygen (R18), while methoxy undergoes thermal decomposition (R19). However, due to the larger value of  $k_{27}$  in the present work, compared to previous modeling studies, the reaction of CH<sub>3</sub>O with O<sub>2</sub> (R27) becomes competitive under lean conditions. Finally, formaldehyde formed in these reactions is converted to CO and CO<sub>2</sub> through the sequence CH<sub>2</sub>O  $\rightarrow$  HCO  $\rightarrow$  CO  $\rightarrow$  CO<sub>2</sub>. Some CO<sub>2</sub> is also produced directly from the formyl radical by the reaction with HO<sub>2</sub> under fuel-rich conditions at 20 bar, and from CO, with HOCO as an intermediate species, under stoichiometric and fuel-lean conditions. The main reaction pathways under the present conditions are consistent with those reported for atmospheric pressure [13-15]. Unlike unsaturated fuels such as ethylene [45], the high pressure does not introduce additional oxidation pathways for methanol because addition reactions of OH and  $O_2$  are not important.

A first order sensitivity analysis shows that the predicted onset temperature is sensitive to the same subset of reactions, independent of stoichiometry. The highest positive sensitivity coefficients are found for the important initiation steps,

$$CH_3OH + HO_2 \rightleftharpoons CH_2OH + H_2O_2 \tag{R8}$$

 $H_2O_2(+M) \Rightarrow OH+OH(+M)$ 

while initiation is inhibited by the terminating reaction

$$HO_2+HO_2 \Rightarrow H_2O_2+O_2$$

Also the reactions of CH<sub>3</sub>OH with OH (R6, R7) and O<sub>2</sub> (R10), as well as  $CH_2O + HO_2$ , show up in the sensitivity analysis.

To test if the present mechanism can predict the autoignition behavior of methanol/oxidizer mixtures at high pressures and moderate temperatures, we used it to simulate the autoignition delay times measured in a RCM by Kumar et al. [16]. For a meaningful comparison between measurements and simulations one must account for deviations from ideal behavior that occur during the experiments, such as heat loss, and it is necessary to take into account the entire compression history of the gas mixture in the simulations. Towards this end the specific volume of the assumed adiabatic core is used as input parameter in the simulations. Similar to the procedure described by Gersen and coworkers [46, 47], we determined the specific volume from the measured pressure in the period between compression and the moment after compression where substantial heat release begins (~1

ms prior the minimum in the pressure trace after compression), and extrapolate the time dependence in this fashion to the region in which substantial heat release begins. The entire derived history of the specific volume is used as input into the simulations. Again, the simulations are performed with the SENKIN code [43].

Figure 10 shows two pressure traces measured by Kumar et al. [16] in a stoichiometric CH<sub>3</sub>OH/O<sub>2</sub>/Ar mixture at different compressed peak pressures of  $P_c \sim 15$  bar (thin solid line) and  $P_c \sim 30$  bar (bold solid line) and compressed peak temperature conditions of  $T_c \sim 916$  K and  $T_c \sim 904$  K, respectively. Comparison of the measured and simulated pressure profiles (Fig. 10) shows that the ignition delay times are underpredicted at stoichiometric conditions; for the  $P_c \sim 30$  bar case the difference is about 25% and at  $P_c \sim 15$  bar the difference between the calculated and measured autoignition delay time is about ~40%. Figure 11 presents the measured and calculated ignition delay times of CH<sub>3</sub>OH/O<sub>2</sub>/Ar mixtures at  $P_c \sim 15$  bar and at an equivalence ratio of  $\varphi=0.25$  as function of the reciprocal temperature. As can be seen from Fig. 11, the calculations for the lean conditions are in good agreement with the measurements over the entire measured domain.

## 5. Conclusions

The methanol oxidation has been studied in a flow reactor in the temperature range 600-900 K and in the pressure interval 20-100 bar as a function of stoichiometry. While the temperature for onset of reaction is largely independent of stoichiometry, it decreases significantly with increasing pressure. The experimental results from the present work, as well as RCM ignition delays reported in literature, have been compared with an updated chemical kinetic model. Ab initio calculations and RRKM analysis were used to obtain rate coefficients for key reactions of  $CH_2OH$  and  $CH_3O$ (dissociation, isomerization, reaction with  $O_2$ ). Modeling predictions were generally in good agreement with the experimental results. Rate-of-production analysis and sensitivity analysis were used to identify the important reactions. The predicted onset of reaction was particularly sensitive to the reaction  $CH_3OH+HO_2 \rightleftharpoons CH_2OH+H_2O_2$ .

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Table 1. Selected reactions in the methanol oxidation subset.

	A (cm <sup>3</sup> /mol·s)	β	E <sub>a</sub> (cal/mol)	Source
1. CH <sub>3</sub> OH(+M) ⇒ CH <sub>3</sub> +OH(+M) Low-pressure limit: Troe parameters: 0.7656 1910 59.51 9374	2.1E18 2.6E49	-0.60 -8.80	92540 101500	[21]
2. $CH_3OH+H \Rightarrow CH_2OH+H_2$	6.6E04	2.73	4449	[22]
3. $CH_3OH+H \Rightarrow CH_3O+H_2$	4.1E04	2.66	9221	[22]
4. $CH_3OH+O = CH_2OH+OH$	2.1E13	0.00	5305	[21]
5. $CH_3OH+O = CH_3O+OH$	3.7E12	0.00	5305	[21]
6. $CH_3OH+OH \Rightarrow CH_2OH+H_2O$	1.5E08	1.44	113	[21]
7. $CH_3OH+OH \Rightarrow CH_3O+H_2O$	2.7E07	1.44	113	[21]
8. $CH_3OH+HO_2 \Rightarrow CH_2OH+H_2O_2$	3.5E-04	4.85	10346	[24]
9. $CH_3OH+HO_2 \Rightarrow CH_3O+H_2O_2$	1.5E-03	4.61	15928	[24]
10. $CH_3OH+O_2 \Rightarrow CH_2OH+HO_2$	3.6E05	2.27	42760	[23]
11. $CH_2OH \Rightarrow CH_2O+H$	3.3E32 8.3E41 2.3E36 5.7E34 3.4E31 1.7E27 4.2E07	-7.18 -9.58 -7.53 -6.64 -5.57 -4.21 1.86	39740 45730 44880 46430 45610 44130 36590	pw, 0.01 bar pw, 0.1 bar pw, 1 bar pw, 20 bar pw, 50 bar pw, 100 bar pw, k <sub><math>\infty</math></sub>
12. $CH_2OH+H \Rightarrow CH_2O+H_2$	1.4E13	0.00	0	[21]
13. $CH_2OH+H \Rightarrow CH_3+OH$	6.0E12	0.00	0	[21]
14. $CH_2OH+H(+M) \Rightarrow CH_3OH(+M)$ Low pressure limit: Troe parameters: 0.2500 210 1434 1.0E30	4.3E15 3.8E37	-0.80 -6.20	0 1333	[21]
15. $CH_2OH+O \Rightarrow CH_2O+OH$	6.6E13	0.00	-693	[21]
16. $CH_2OH+OH \Rightarrow CH_2O+H_2O$	2.4E13	0.00	0	[21]
17. $CH_2OH+HO_2 \Rightarrow CH_2O+H_2O_2$	1.2E13	0.00	0	[21]
18. $CH_2OH+O_2 \Rightarrow CH_2O+HO_2$ Duplicate	7.2E13 2.9E16	0.00 -1.50	3736 0	[21] [21]
19. $CH_3O \Rightarrow CH_2O+H$	1.8E31 2.9E20 7.7E115 5.4E111 5.6E87 4.7E93 1.1E10	-6.96 -3.64 -32.10 -30.46 -23.06 -24.61 1.23	26670 20140 63410 62690 52090 57200 24250	pw, 0.01 bar pw, 0.1 bar pw, 1 bar pw, 20 bar pw, 50 bar pw, 100 bar pw, k∞
20. $CH_3O \Rightarrow CH_2OH$	1.4E27 2.3E16 6.0E111 4.2E107 4.4E83 3.6E89 8.9E05	-5.99 -2.66 -31.12 -29.49 -22.09 -23.64 2.21	27390 20860 64130 63410 53620 57920 24970	pw, 0.01 bar pw, 0.1 bar pw, 1 bar pw, 20 bar pw, 50 bar pw, 100 bar pw, $k_{\infty}$
21. $CH_3O+H \Rightarrow CH_2O+H_2$	5.3E13	0.00	745	[21]
22. $CH_3O+H \Rightarrow CH_3+OH$	4.6E12	0.00	745	[21]
23. $CH_{3}O+H(+M) = CH_{3}OH(+M)$	2.4E12	0.50	50	[21]

Low pressure limit: Troe parameters: 0.7000 100 90000 10000	4.7E40	-7.44	140803	
24. $CH_3O+O \Rightarrow CH_2O+OH$	3.8E12	0.00	0	[21]
25. $CH_3O+OH \Rightarrow CH_2O+H_2O$	1.8E13	0.00	0	[21]
26. $CH_3O+HO_2 \Rightarrow CH_3OH+O_2$	1.4E11	0.00	0	[25]
27. $CH_3O+O_2 \Rightarrow CH_2O+HO_2$	4.8E-01	3.57	-1055	pw

pw: present work

**Table 2.** Energies of stationary points in the C/H<sub>3</sub>/O<sub>3</sub> system derived with coupled cluster theory.

aug-cc-pVTZ Au	aug-cc-pVQZ aug-cc-pVQZ	aug-cc-pV5Z aug-cc-pV5Z	CBS limit au	relative enthalpy at 0 K kJ/mol (a)
-114.8850816	-114.9138885	-114.922965	-114.9271403	0
-0.49982118	-0.49994832	-0.49999478	-0.500021534	
-114.3428434	-114.3723302	-114.3815812	-114.3858104	82.4
-114.8987541	-114.9286887	-114.9379937	-114.9421907	-37.6
-114.8367789	-114.8663888	-114.8756417	-114.8798475	103.0
-114.8330452	-114.8622654	-114.8714234	-114.8756039	124.0
-114.8269853	-114.8566007	-114.8658411	-114.8700318	127.0
-150.1407415	-150.1785442	-150.1909034	-150.1969069	
-150.7265278	-150.7653707	-150.7780254	-150.7841404	-119.8
	aug-cc-pVTZ Au -114.8850816 -0.49982118 -114.3428434 -114.8987541 -114.8367789 -114.8330452 -114.8269853 -150.1407415 -150.7265278	aug-cc-pVTZaug-cc-pVQZAuaug-cc-pVQZ-114.8850816-114.9138885-0.49982118-0.49994832-114.3428434-114.3723302-114.8987541-114.9286887-114.8367789-114.8663888-114.830452-114.8663888-114.8269853-114.8566007-150.1407415-150.1785442-150.7265278-150.7653707	aug-cc-pVTZaug-cc-pVQZaug-cc-pV5ZAuaug-cc-pVQZaug-cc-pV5Z-114.8850816-114.9138885-114.922965-0.49982118-0.49994832-0.49999478-114.3428434-114.3723302-114.3815812-114.8987541-114.9286887-114.9379937-114.8367789-114.8663888-114.8756417-114.8330452-114.8622654-114.8714234-114.8269853-114.8566007-114.8658411-150.1407415-150.1785442-150.1909034-150.7265278-150.7653707-150.7780254	aug-cc-pVTZaug-cc-pVQZaug-cc-pV5ZCBS limitAuaug-cc-pVQZaug-cc-pV5Zau-114.8850816-114.9138885-114.922965-114.9271403-0.49982118-0.49994832-0.49999478-0.500021534-114.3428434-114.3723302-114.3815812-114.3858104-114.8987541-114.9286887-114.9379937-114.9421907-114.8367789-114.8663888-114.8756417-114.8798475-114.8330452-114.8622654-114.8714234-114.8760398-114.8269853-114.8566007-114.8658411-114.8700318-150.1407415-150.1785442-150.1909034-150.1969069-150.7265278-150.7653707-150.7780254-150.7841404

(a) Includes zero-point energy at CCSD/6-311G(d,p) level scaled by 0.954.

(b) The TS for CH<sub>3</sub>O + O<sub>2</sub> → CH<sub>2</sub>O + HO<sub>2</sub> has an enthalpy of 13.1 kJ mol<sup>-1</sup> relative to reactants based on CBS-QB3 energies obtained with CCSD/6-311G(d,p) geometries and zero-poin energies.

**Table 3.** Matrix of experimental conditions. The experiments are conducted at constant mass flow, 5000 ml(STP)/min, in the temperature range 600-900 K and pressure interval 20-100 bar. The residence time is dependent on the reaction temperature and pressure.

Exp Pressure (bar)	CH <sub>3</sub> OH (ppm) (p	O <sub>2</sub>	2 NO <sub>x</sub> m) (ppm)	λ	t <sub>res</sub>	
		(ppm)			(s)	
Set 1	20	2606	904	0	0.23	1317/T
Set 2	50	2390	939	0	0.26	3293/T
Set 3	100	2370	951	0	0.27	6587/T
Set 4	100	3010	4500	0	1.00	6587/T
Set 5	100	2930	71300	0	16.2	6587/T



**Figure 1.** Relative enthalpies computed for 0 K based on CCSD(T) energies extrapolated to the complete basis set limit for the  $CH_3O/CH_2OH$  system.



**Figure 2.** Arrhenius plot for the reactions  $CH_2OH = CH_2O+H$  (R11),  $CH_3O = CH_2O+H$  (R19), and  $CH_3O = CH_2OH$  (R20) (lower). The (thin) solid lines denote the calculated high-pressure limit for each reaction, while pressure-specific rate constants are shown as dash-dotted lines (20 bar), short-dashed lines (50 bar), and long-dashed lines (100 bar). The thick solid line denotes the high-pressure limit derived by Hippler et al. [32] for (R19).



**Figure 3.** Arrhenius plot for the reaction CH<sub>3</sub>O+O<sub>2</sub>  $\rightleftharpoons$  CH<sub>2</sub>O+HO<sub>2</sub> (R27). Solid line: present work. Literature data: dashed line: CEC recommendation [33]; • Wiebe et al. [34], □ Alcock and Mile [35], ○ Barker et al. [36], ■ Kirsch and Parkes [37], ▲ Batt and Robinson [38], △ Cox et al. [39], ▼ Gutman et al., [40],  $\bigtriangledown$  Lorentz et al. [41], ◆ Zaslonko et al. [42], ◊ Wantuck et al. [43].



**Figure 4.** CH<sub>3</sub>OH, CO, CO<sub>2</sub>, and O<sub>2</sub> concentration profiles as a function of temperature for the conditions of set 1 in Table I (fuel-rich, 20 bar). Comparison between experimental data (symbols) and model predictions (lines).



**Figure 5.** CH<sub>3</sub>OH, CO, CO<sub>2</sub>, and O<sub>2</sub> concentration profiles as a function of temperature for the conditions of set 2 in Table I (fuel-rich, 50 bar). Comparison between experimental data (symbols) and model predictions (lines).



**Figure 6:** CH<sub>3</sub>OH, CO, CO<sub>2</sub>, and O<sub>2</sub> concentration profiles as a function of temperature for the conditions of set 3 in Table 3 (fuel-rich, 100 bar). Comparison between experimental data (symbols) and model predictions (lines).



**Figure 7:** CH<sub>3</sub>OH, CO, CO<sub>2</sub>, and O<sub>2</sub> concentration profiles as a function of temperature for the conditions of set 4 in Table 3 (stoichiometric, 100 bar). Comparison between experimental data (symbols) and model predictions (lines).



**Figure 8:**  $CH_3OH$ , CO, and  $CO_2$  concentration profiles as a function of temperature for the conditions of set 5 in Table 3 (lean, 100 bar). Comparison between experimental data (symbols) and model predictions (lines).



Figure 9. Reaction path diagram for methanol oxidation at the high pressure conditions of the present work.



Figure 10. Measured and modeled pressure traces for stoichiometric methanol/O<sub>2</sub>/Ar mixtures with  $O_2$ :Ar=1:3.76 at compressed peak pressures of  $P_c \sim 15$  bar (thin lines) and  $P_c \sim 30$  bar (solid lines) and compressed peak temperature conditions of  $T_c \sim 916$  K and  $T_c \sim 904$  K, respectively. The solid lines represent the measurements performed by Kumar et al. [16], and the dashed lines represent the simulations performed with the present mechanism.



Figure 11. Measured (dots) and calculated (line) autoignition delay times as function of the reciprocal temperature for methanol/ $O_2$ /Ar mixtures at  $\phi$ =0.25 and a pressure of  $P_c$ =15 bar. The measurements were performed by Kumar et al. [16].