**An Elliptic Problem Arising From the Unsteady
Transonic Small Disturbance Equations**

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Abstract

We prove a theorem on existence of a weak solution of the Dirichlet problem for a quasilinear elliptic equation with a degeneracy on one part of the boundary. The degeneracy is of a type ("Keldysh type") associated with singular behavior — blow-up of a derivative — at the boundary. We define an associated operator which is continuous, pseudo-monotone and coercive and show that a weak solution displaying singular behavior at the boundary exists.

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1 Introduction

The following Dirichlet boundary value problem for a quasilinear degenerate elliptic equation arises in studying self-similar solutions of hyperbolic conservation laws in two space dimensions. The equation is

$$\left((u - \rho)u_\rho + \frac{u}{2} \right)_\rho + u_{\eta\eta} = 0, \quad (1)$$

in the bounded domain Ω shown in Figure 1. The boundary data is given by $u = g$ on $\partial\Omega$, where $g = \rho_m$ on the degenerate boundary Γ_2 and on Γ_∞ , and $g = \tilde{g}$ on Γ_3 , where $\tilde{g} \geq \rho_m$ is continuous and symmetric with respect to the ρ -axis.

Equation (1) results from scaling variables ($\xi = x/t$, $\eta = y/t$, $\rho = \xi + \eta^2/4$) in the unsteady transonic small disturbance (TSD) equation, also called the 2-D Burgers equation or Zabolotskaya-Khoklov equation:

$$\begin{aligned} u_t + u u_x + v_y &= 0, \\ -v_x + u_y &= 0. \end{aligned} \quad (2)$$

Figure 1: Domain and Boundary Components of the Dirichlet Problem

Reasons for this scaling are discussed in [3]. The system of equations (2) has been derived and used by several authors (see [7], [14] and the references therein) to model problems in nonlinear acoustics and nonlinear geometrical optics, and also in unsteady transonic flow. Brio and Hunter in [2] use (2) as an asymptotic reduction of the Euler equations in order to study weak shock reflection. The system (2) retains the nonlinear wave structure of the gas dynamics equations in the form of a particularly simple nonlinearity, and may give a description of the transition between regular and Mach reflection for weak shocks. This phenomenon has been extensively studied numerically and experimentally (see, for example, [5], [15]); however, there are only a few related analytical results. We mention Morawetz's study of a full-potential-equation model, [10], which is in the same spirit as the TSD model. The recent paper of Brio and Hunter, [2], gives a geometrical analysis of equation (2), but does not consider questions which lead to problems like equation (1) and Figure 1; on the other hand, Tabak and Rosales in [16] use an alternative analysis which leads to a different sort of elliptic problem.

The observation we pursue in this paper is that the standard reduction of equation (2) to similarity variables leads to a degenerate elliptic equation

with a degeneracy which has not previously been studied (but see the recent paper of Choi, Lazer and McKenna, [4], which studies a degeneracy which is almost the same). An examination of the reduction is made elsewhere, [3], but we note here for completeness that the unusual degeneracy in the equation has nothing to do with the fact that (2) is not hyperbolic in time — this is an artifact of the weakly nonlinear geometric optics reduction of the compressible Euler equations; the degeneracy seems to characterize self-similar reductions of quasilinear hyperbolic equations in two or more space dimensions.

In this paper, we show the existence of weak solutions for a boundary-value problem for equation (1) obtained by making a set of reasonable (but slightly different) assumptions on the boundary values, g . We also simplify the shape of the domain in an inessential way. We are now in the process of applying this result to the original Mach reflection model. This involves matching the elliptic to the hyperbolic region via a free-boundary problem, and modifying the elliptic region to be an unbounded domain. However, the degenerate problem we have solved seems to be interesting in its own right, since it is a new type of degeneracy which occurs in a physically interesting problem.

The problem we study in this paper is obtained by letting $x = -\rho$ and $y = \eta$ so that (1) becomes

$$\begin{aligned} \left((u+x)u_x - \frac{u}{2} \right)_x + u_{yy} &= 0 \quad \text{in } \Omega, \\ u &= g \quad \text{on } \partial\Omega. \end{aligned} \tag{3}$$

We take Ω to be a rectangle, $0 \leq x \leq a$, $-b \leq y \leq b$, in the right half-plane, and the boundary data to satisfy $g = 0$ on $x = 0$, making this part of the boundary, Γ_2 , degenerate. We shall also assume $g \geq 0$ on the rest of the boundary; this, we shall show, leads to a comparison principle. We shall need more conditions on g ; see the Assumption in Section 4.1. We look for solutions of (3) which make the equation elliptic in the interior of Ω .

The degenerate characteristic form at the boundary corresponds to change of type in the flow from hyperbolic (supersonic) to elliptic (subsonic). In fact, the steady TSD equation also has this character, as well as the difficulty that the location in physical space of the sonic line is not known apriori. How the techniques of this paper would apply to the steady TSD equation is an open question.

Elliptic equations with degenerate characteristic forms have been much studied: [12] surveys the literature up to 1973, while Triebel, [17], extensively covers a number of linear degenerate equations and provides a detailed embedding theory of corresponding weighted Sobolev spaces. However, the prototype degeneracies in [17] apply uniformly to the entire second-order part of the operator. Both [12] and [17] study only linear operators. Linear techniques point to differences between Tricomi type equations (modeled by $xu_{yy} + u_{xx} = 0$) and Keldysh type equations, [8], (for example $u_{yy} + xu_{xx} = 0$) in the elliptic region, $x \geq 0$. For the linear Keldysh equation, unlike the Tricomi equation, it is necessary to impose a boundary condition at the degenerate boundary. However, solutions of the Keldysh equation are not differentiable at the degenerate boundary. This was noted by Kohn and Nirenberg, [9]. Because of this behavior, generalizations of Keldysh type equations are not covered by most results on degenerate elliptic equations. Keldysh's equation is closely related to a linearized version of equation (1). Current research in degenerate quasilinear equations focusses on models like the p -Laplacian and other geometric problems whose structure appears to be quite different from the problem at hand. Nonlinear equations of Keldysh type have not been studied at all, as far as we know.

The structure of the singularity of equation (3) is discussed in section 2, where we show that solutions exhibit a square root type singularity in the neighborhood of the degenerate boundary. Thus we are led to consider weak solutions of (3) in a Hilbert space with a weighted Sobolev norm.

As a motivation for the choice of the solution space, we study a linear equation associated to the original quasilinear problem. Existence of a weak solution can be proved, as in [12], using the classical Riesz theorem. This is the content of section 3.

In section 4 we define the weighted Sobolev solution space and the weak formulation for (3). To apply the Browder-Minty theorem, adapting the theory expounded in [13] to a degenerate operator, we need to show that the operator is continuous, pseudo-monotone and coercive. It turns out to be necessary to reformulate the problem so as to obtain a new operator which satisfies all the conditions for the Browder-Minty theorem. We verify the conditions in Section 5. Finally, using the techniques of sub- and super-solutions we show that a solution to the reformulated problem actually satisfies the original equation.

2 The Degenerate Boundary and the Comparison Principle

In this section, we study equation (3) with Dirichlet boundary conditions on the rectangle Ω . The vanishing of the characteristic form for equation (3) may induce singular behavior in solutions on Γ_2 , the part of the boundary on which the equation is degenerate. From the standard theory of elliptic equations we expect the solution to be smooth ($C^{2,\mu}$) everywhere in Ω away from Γ_2 . In addition, a comparison principle holds for elliptic solutions of (3), and we formulate this here. The type of singularity expected at the boundary Γ_2 is described first.

2.1 Singularities in the Solution

Let us assume that $u \in C^2(\Omega) \cap C(\bar{\Omega})$ is a classical solution of equation (3) in Ω , and that it satisfies equation (3) in a weak sense in $\bar{\Omega}$, with $u = 0$ on Γ_2 . We shall show that if u attains its minimum on Γ_2 , then u belongs to $C^\alpha(\bar{\Omega})$, $\alpha > 0$, and it must have growth of order $\sqrt{x}$ precisely at $x = 0$.

For this purpose, we write a solution u in the form $u(x, y) = x^\alpha v(x, y)$, with $\alpha > 0$, and $v(0, y) \neq 0$. Then equation (3) yields an equation for v :

$$\begin{aligned} & x^{2\alpha-2} \left\{ \alpha(2\alpha-1)v^2 + x(4\alpha vv_x) + x^2(v_x^2 + vv_{xx}) \right\} \\ & + x^{\alpha-1} \left\{ \alpha(\alpha - \frac{1}{2})v + x(v_{yy} + (\frac{1}{2} + 2\alpha)v_x) + x^2v_{xx} \right\} = 0 \end{aligned} \quad (4)$$

The relative order of terms depends on whether α is between 0 and 1, $\alpha = 1$, or $\alpha > 1$.

If $0 < \alpha < 1$, equation (4) implies that at $x = 0$, $\alpha(2\alpha - 1)v^2 = 0$. Since we have assumed that $v(0, y) \neq 0$, it follows that $\alpha = \frac{1}{2}$. Furthermore, in the case when $\alpha = \frac{1}{2}$, equation (4) implies that $v_x(0, y) = 0$.

If $\alpha = 1$, then at $x = 0$, $v(v + \frac{1}{2}) = 0$, so $v(0, y) = -\frac{1}{2}$. Thus $v(x, y) = -\frac{1}{2} + w(x, y)$, where $w(0, y) = 0$.

Finally, we consider the case when $\alpha > 1$. The leading-order term in (4) is then $\alpha(\alpha - \frac{1}{2})x^{\alpha-1}v$. Therefore, at $x = 0$, equation (4) implies that $v(0, y) = 0$, and so no such solutions exist.

Thus we have proved the following theorem:

Theorem 1 Let $u \in C^2(\Omega) \cap C(\bar{\Omega})$ be a solution of (1) in $\bar{\Omega}$, such that $u = 0$ on Γ_2 . Then, if u is not identically zero everywhere in $\bar{\Omega}$, either:

- (1) $u \in C^1(\bar{\Omega})$ and u vanishes exactly to first order at Γ_2 , or
- (2) $u \in C^\alpha(\bar{\Omega})$ for $\alpha \leq \frac{1}{2}$, and $u(x, y) = \sqrt{x} v(x, y)$, where $v(0, y) \neq 0$.

In both cases there are further restrictions on u . In case (1), $u_x(0, y) = -\frac{1}{2}$; in case (2), then $v_x(0, y) = 0$.

In particular, if $u \geq 0$ in Ω , then case (1) cannot hold, and we have:

Corollary 1 There is no solution of equation (3) which is C^1 in $\bar{\Omega}$ and attains its minimum on Γ_2 .

If case (2) holds, then $u + x$ is negative near Γ_2 unless $v > 0$ on Γ_2 . Hence, we have

Corollary 2 A nonnegative solution of (3) which is $C^\alpha(\bar{\Omega})$, $\alpha > 0$, and elliptic in Ω must have growth of order $\sqrt{x}$ precisely at $x = 0$ and must satisfy $u > 0$ near Γ_2 .

2.2 The Comparison Principle

We now prove a comparison theorem for the quasilinear equation (3). We prove that every function $v \in C(\bar{\Omega}) \cap C^2(\Omega)$, which satisfies the usual conditions for a sub- or super-solution of an elliptic problem, is in fact one, provided in addition it satisfies the condition $v_{xx} \leq 0$ in Ω . This leads to a unique elliptic solution in the class of functions u for which $u_{xx} \leq 0$ in Ω . Moreover, using this comparison principle, we obtain estimates for a weak solution in the neighborhood of the degenerate boundary.

More precisely, let $u \in C(\bar{\Omega}) \cap C^2(\Omega)$. We denote by $Q(u)$ the quasilinear operator associated with equation (3),

$$Q(u) \equiv \left((u + x)u_x - \frac{u}{2} \right)_x + u_{yy}. \quad (5)$$

Theorem 2 Let $u, v \in C(\bar{\Omega}) \cap C^2(\Omega)$. Suppose that u is a solution of (3) which is elliptic in Ω , and that $v_{xx} \leq 0$ in Ω . If $Q(v) \leq Q(u) = 0$ in Ω , and $u \leq v$ on $\partial\Omega$, then $u \leq v$ in Ω .

Proof: Let $w = u - v$. Then

$$Q(u) - Q(v) = (u + x)w_{xx} + w_{yy} + (u_x + v_x + \frac{1}{2})w_x + v_{xx}w \geq 0. \quad (6)$$

Denote by L the linear operator on w defined by (6). Then we have $Lw \geq 0$ in $\Omega^+ \equiv \{(x, y) \in \Omega | w(x, y) > 0\}$, and $w \leq 0$ on $\partial\Omega$. Since L is elliptic in Ω , with the coefficient v_{xx} of w less than zero, we can apply the maximum principle for linear equations elliptic in the interior of Ω (see [6, Theorem 3.1, Corollary 3.2]), to see that $w \leq 0$ in Ω , and therefore, $v \geq u$ in Ω . Similarly, one can show that a function z , such that $z_{xx} \leq 0$, $Q(u) \leq Q(z)$ in Ω , and $u \geq z$ on $\partial\Omega$, has the property that $z \leq u$ in Ω . ■

Corollary 3 *If $u, v \in C(\bar{\Omega}) \cap C^2(\Omega)$ are two solutions of (3), elliptic in Ω , and u_{xx} and $v_{xx} \leq 0$, then $u \equiv v$ in Ω .*

3 A Linear Model Equation

As a first step in understanding the behavior of a solution to the Dirichlet problem for the quasilinear equation (3), we study the linear equation obtained by replacing the coefficient $u + x$ of u_{xx} by a fixed function with the singularity we found in the last section. Let $w(x, y)$ be a bounded, measurable solution of (3) which satisfies

$$w + x \geq 0 \text{ in } \bar{\Omega}, \quad w > -x \text{ in } \Omega, \quad w = 0 \text{ on } \Gamma_2 = \{x = 0\};$$

for the calculation below, we choose $w(x, y) = \sqrt{x}$. The linear equation associated with (3) can be written in the form

$$\left((w + x)u_x - \frac{u}{2} \right)_x + u_{yy} + cu = cw, \quad c < 0. \quad (7)$$

Equation (7) is elliptic everywhere in $\bar{\Omega}$ except at the boundary Γ_2 , on which the characteristic form vanishes. Appropriate boundary conditions on Γ_2 are found by computing the *Fichera function* for (7); we summarize the procedure given in [12].

For a linear operator

$$L(u) = a^{kj}(x)u_{x_k x_j} + b^k(x)u_{x_k} + c(x)u, \quad (8)$$

with non-negative characteristic form

$$a^{kj}(x)\xi_k\xi_j \geq 0, \quad \forall \vec{\xi} = (\xi_1, \xi_2), \quad (9)$$

let Σ^0 denote the part of $\partial\Omega$ along which (9) vanishes, and let $\Sigma_3 = \partial\Omega - \Sigma^0$.

Definition 1 Let $\vec{n} = (n_1, n_2)$ be the inward normal to $\partial\Omega$. The function

$$b \equiv (b^k - a_{x_j}^{kj})n_k,$$

defined on Σ^0 , is called the **Fichera function**.

According to the sign of the Fichera function we divide Σ^0 into three parts, with Σ_2 the part along which $b < 0$. As studied in [12], a well-posed boundary value problem for the operator (8) or (7) is the following:

$$\begin{aligned} L(u) &= f \quad \text{in } \Omega \\ u &= g \quad \text{on } \Sigma_2 \cup \Sigma_3 \end{aligned}$$

where f is a given function in Ω , and g a given function on $\Sigma_2 \cup \Sigma_3$.

For equation (7), the Fichera function on Γ_2 has the constant value $-1/2$, so that imposing data (for u) along Γ_2 leads to a well-posed boundary value problem.

3.1 Existence of a Weak Solution in a Hilbert Space

In this section we study the existence of a weak solution for the linear degenerate elliptic boundary value problem (7) with

$$u = g \quad \text{on } \partial\Omega, \quad (10)$$

where w is a bounded, measurable function in $\bar{\Omega}$ and $c < 0$.

We formulate this problem as a Dirichlet problem with zero boundary data in a standard way: let $v = u - g$, where $g \in C^2(\Omega) \cap C(\bar{\Omega})$ is an extension of the boundary data to the whole domain such that $g \geq 0$ in $\bar{\Omega}$ and g never exceeds its maximum on the boundary. Then, existence of a solution of (7) and (10) is equivalent to solving

$$\begin{aligned} \left((w+x)v_x - \frac{v}{2} \right)_x + v_{yy} + cv &= f \quad \text{in } \Omega, \\ v &= 0 \quad \text{on } \partial\Omega. \end{aligned} \quad (11)$$

We define a weak solution of problem (11) as follows.

Let v be a function in $C_0^2(\overline{\Omega}) = \{v \in C^2(\overline{\Omega}) | v = 0 \text{ on } \partial\Omega\}$, and z be in $C_0^1(\overline{\Omega})$. Multiplying the PDE in (11) by a test function $z \in C_0^1(\overline{\Omega})$ and integrating by parts we obtain

$$-\int_{\Omega} \left((w+x)v_x - \frac{v}{2} \right) z_x + v_y z_y - cvz = \int_{\Omega} fu, \quad \forall z \in C_0^1(\overline{\Omega}).$$

For functions in $C_0^1(\overline{\Omega})$ we introduce the scalar product

$$(v, z)_{\mathcal{H}} \equiv \int_{\Omega} (w+x)v_x z_x + v_y z_y + vz.$$

Define $\mathcal{H}$ to be the completion of $C_0^1(\overline{\Omega})$ in the norm defined by this scalar product.

For $v, z \in C_0^1$ we study the bilinear form

$$B(v, z) \equiv -\int_{\Omega} \left((w+x)v_x - \frac{v}{2} \right) z_x + v_y z_y - cvz.$$

Proposition 1 For each fixed $z \in C_0^1$, the operator $v \mapsto B(v, z)$ is a linear, bounded functional on $v \in \mathcal{H}$.

Proof: It is enough to show that for all $v, z \in C_0^1$, $B(v, z)$ satisfies the inequality

$$|B(v, z)| \leq M \sqrt{\int_{\Omega} z_x^2 + z_y^2 + z^2} \|v\|_{\mathcal{H}}, \quad (12)$$

where M depends only on the coefficients of the operator B . Inequality (12) follows from the Cauchy-Schwartz inequality in $(L^2)^3$. That is,

$$\begin{aligned} |B(v, z)| &= \left| \int_{\Omega} \left((w+x)v_x - \frac{v}{2} \right) z_x + v_y z_y - cvz \right| \\ &\leq \left(\int_{\Omega} \left((w+x)v_x - \frac{v}{2} \right)^2 + v_y^2 + c^2 v^2 \right)^{\frac{1}{2}} \left(\int_{\Omega} z_x^2 + z_y^2 + z^2 \right)^{\frac{1}{2}} \\ &= M \sqrt{\int_{\Omega} z_x^2 + z_y^2 + z^2} \|v\|_{\mathcal{H}} \end{aligned}$$

where $M = \max\{\sqrt{(w+x)}, 1, |c| + \frac{1}{2}\}$. From this inequality we conclude that $B(v, z)$ can be extended to $v \in \mathcal{H}$, and that for each fixed $z \in C_0^1$, $B(v, z)$ is a bounded, linear functional on $\mathcal{H}$. ■

Definition 2 Let $f \in L^2(\Omega)$. A function $v \in \mathcal{H}$ is called a weak solution of (11) if for all $u \in C_0^1$ the following equation holds:

$$\int_{\Omega} fu = B(v, u) \quad (13)$$

Theorem 3 Let $f \in L^2(\Omega)$ and $c < 0$. Then there exists a weak solution $v \in \mathcal{H}$ of the problem (11).

The existence of a weak solution of (11) in the Hilbert space $\mathcal{H}$ is a consequence of a more general result proved by Oleinik and Radkevič [12, Theorem 1.4.1, page 29]. Note that these authors require more smoothness in w than we have; but it is used only to straighten the degenerate part of the boundary. For the problem we consider here, we do not need derivatives of w at Γ_2 .

4 The Weighted Sobolev Space Setting

In this section we formulate the problem of finding a weak solution of the Dirichlet problem for equation (3) in a weighted Sobolev space defined below. The extension from the linear case is nontrivial because of the lack of smoothness of u at Γ_2 . We get around this difficulty by defining a new operator with a fixed singularity at Γ_2 . In Section 5 we show that the new operator is continuous, coercive, and pseudo-monotone, and use the Browder-Minty theorem to show the existence of a weak solution. Then, using the techniques of sub- and super- solutions, we show that a solution to the redefined problem actually satisfies the original problem (3).

4.1 Modification of the Problem

We now make an assumption on the boundary data g which is necessary for redefining the operator as outlined above.

Assumption: The boundary data g can be extended to a $C^2(\Omega)$ function everywhere in Ω with the properties $g \geq 0$, $g_{xx} < 0$ and $\epsilon\sqrt{x} \leq g \leq (1/\epsilon)\sqrt{x}$ in Ω , for some $\epsilon > 0$. The maximum of g is attained on $\partial\Omega$, and we can write $g(x, y) = \sqrt{x}\bar{g}(x, y)$, where $\bar{g}$ is smooth.

Note that this is stronger than the conditions in problem (1), and is inconsistent with the behavior of g on Γ_∞ in Figure 1. We shall consider modifications in future work.

To use the monotone operator framework, we recast the problem as a Dirichlet problem with zero boundary data, by the substitution $u = v + g$. We obtain in place of (3)

$$\begin{aligned} -\left((v + g + x)v_x + \left(g_x - \frac{1}{2}\right)v\right)_x - v_{yy} &= f \quad \text{in } \Omega \\ v &= 0 \quad \text{on } \partial\Omega. \end{aligned} \quad (14)$$

Here

$$f = g_{yy} + ((g + x)g_x)_x - \frac{1}{2}g_x. \quad (15)$$

The weak formulation of (14) is as follows. Let $v \in C^2(\Omega)$ satisfy (14) and let w be in $C_0^1(\bar{\Omega})$. Multiplying the PDE in (14) by the test function w and integrating by parts one obtains

$$B(v, w) \equiv \int_{\Omega} \left((v + g + x)v_x + \left(g_x - \frac{1}{2}\right)v \right) w_x + v_y w_y = \int_{\Omega} f w. \quad (16)$$

For functions v, w in $C_0^1(\bar{\Omega})$ we introduce the scalar product

$$(v, w)_{\mathcal{H}_\alpha} \equiv \int_{\Omega} x^\alpha v_x w_x + v_y w_y + v w. \quad (17)$$

Define the Hilbert space $\mathcal{H}_\alpha$ to be the completion of $C_0^1(\bar{\Omega})$ in the norm defined by (17). We have a form of the Poincaré inequality which makes this norm equivalent to

$$\|u\|^2 \equiv \int_{\Omega} x^\alpha v_x^2 + v_y^2. \quad (18)$$

Proposition 2 *Let $0 < \alpha < 2$. There exists a constant $C > 0$ such that for every $v \in \mathcal{H}_\alpha$ the following inequality holds:*

$$\int_{\Omega} v^2 dx dy \leq C \int_{\Omega} x^\alpha v_x^2 dx dy$$

Proof: We may assume $v \in C_0^1(\bar{\Omega})$. An integration by parts and the Cauchy-Schwartz inequality imply

$$\begin{aligned} \int_{\Omega} v^2 dx dy &= - \int_{\Omega} x \frac{\partial v^2}{\partial x} dx dy \\ &= -2 \int_{\Omega} x^{1-\alpha/2} v x^{\alpha/2} \frac{\partial v}{\partial x} dx dy \\ &\leq 2 \max_{\Omega} x^{1-\alpha/2} \left(\int_{\Omega} v^2 \right)^{\frac{1}{2}} \left(\int_{\Omega} x^{\alpha} v_x^2 \right)^{\frac{1}{2}}. \end{aligned}$$

Therefore,

$$\int_{\Omega} v^2 dx dy \leq C \int_{\Omega} x^{\alpha} v_x^2 dx dy,$$

where $C = (2 \max_{\Omega} x^{1-\alpha/2})^2$. ■

As a consequence, we have:

Proposition 3 *The norms defined by (18) and (17) are equivalent.*

For the remainder of this paper we shall use (18) for the norm, and will modify (17) appropriately for the corresponding scalar product.

Definition 3 *Let $f \in L^2(\Omega)$. A function $v \in \mathcal{H}_{\alpha}$ is called a weak solution of (14) if (16) holds for every $w \in C_0^1(\bar{\Omega})$.*

We anticipate a couple of difficulties in proving existence of a weak solution using (16). First, if u is not greater than $-x$ in Ω , the equation will no longer be elliptic. As we shall see, the value of α for which we get an existence theorem is $\alpha = \frac{1}{2}$; however for no value of α is (16) defined for all v, w in $\mathcal{H}_{\alpha}$, continuous on $\mathcal{H}_{\alpha} \times \mathcal{H}_{\alpha}$ and coercive, even for $v \geq -g$. Assuming that we are looking for an elliptic solution u of equation (3), then, from the comparison principle of Theorem 2, we expect u to be squeezed between sub- and super-solutions given by $\delta\sqrt{x}$, $M\sqrt{x}$, respectively, for any $\delta \leq \epsilon$ and $M \geq (1/\epsilon)$. This is equivalent to saying that $v + g$ in problem (14) exhibits square-root behavior on Γ_2 . Therefore, we redefine $B(v, w)$ by introducing a cut-off function for the coefficient $v + g$ in such a way that the modified coefficient exhibits square-root behavior on Γ_2 . We define the new operator $B_g(v, w)$ as follows.

Let $\epsilon > 0$ be the constant which bounds g in the Assumption. We define

$$\tilde{u} = \begin{cases} v + g & \text{if } \epsilon\sqrt{x} \leq v + g \leq (1/\epsilon)\sqrt{x} \\ \epsilon\sqrt{x} & \text{if } v + g < \epsilon\sqrt{x} \\ (1/\epsilon)\sqrt{x} & \text{if } (1/\epsilon)\sqrt{x} < v + g \end{cases} \quad (19)$$

Now define

$$B_g(v, w) = \int_{\Omega} \left((\tilde{u} + x) v_x + \left(g_x - \frac{1}{2}\right)v \right) w_x + v_y w_y. \quad (20)$$

We shall prove the existence of a solution in $\mathcal{H}_{\alpha}$, $\alpha = \frac{1}{2}$, for

$$B_g(v, w) = \int_{\Omega} f w. \quad (21)$$

Then we shall show that the solution we have found satisfies $\epsilon\sqrt{x} \leq v + g \leq (1/\epsilon)\sqrt{x}$. Therefore, $B_g(v, w)$ coincides with $B(v, w)$ for this solution, and v satisfies the original problem.

We have introduced a new difficulty: B_g depends on g , as does f in (21). We handle this by supposing f to be an arbitrary function, in (21), (not related to g by (15)), and then letting f be the function defined by (15) to obtain the existence of a weak solution for problem (14).

Notice that for each fixed $v \in \mathcal{H}_{\alpha}$, $B_g(v, w)$ is a linear functional in $w \in C_0^1$. We show that $w \mapsto B_g(v, w)$ is bounded and can be extended to a bounded, linear functional on $\mathcal{H}_{\alpha}$, for $\alpha = \frac{1}{2}$. In order to do this, we first investigate some properties of the space $\mathcal{H}_{\alpha}$.

4.2 Embeddings of the Hilbert Space $\mathcal{H}_{\alpha}$

Let $(\Omega, \mathcal{B}, \mu)$ be the measure space defined on Ω with product measure μ given by

$$d\mu = \frac{1}{x^{\beta}} dx dy.$$

The value of β will be specified later. We denote by $L^2(\Omega, \beta)$ the space

$$L^2(\Omega, \beta) \equiv \left\{ u \mid u \text{ measurable and } \int_{\Omega} |u|^2 d\mu < \infty \right\}.$$

Clearly, $L^2(\Omega, \beta)$ is a Hilbert space with scalar product

$$(u, v)_{L^2(\Omega, \beta)} \equiv \int_{\Omega} u v d\mu = \int_{\Omega} u v \frac{dx}{x^{\beta}} dy.$$

Theorem 4 *The space $\mathcal{H}_\alpha$, for each $0 < \alpha < 1$, is continuously embedded in $L^2(\Omega, \beta)$, for any $\beta \leq 2 - \alpha$.*

Proof: Let $u \in C_0^1(\overline{\Omega})$. If $\beta \neq 1$, an integration by parts and the Cauchy-Schwartz inequality yield

$$\begin{aligned}
\|u\|_{L^2(\Omega, \beta)}^2 &= \int_{\Omega} \frac{u^2}{x^\beta} dx dy \\
&= -\frac{1}{1-\beta} \int_{\Omega} x^{1-\beta} \frac{\partial u^2}{\partial x} dx dy \\
&= -\frac{2}{1-\beta} \int_{\Omega} x^{1-\beta} u \frac{\partial u}{\partial x} dx dy \\
&= -\frac{2}{1-\beta} \int_{\Omega} x^\delta \frac{u}{x^{\beta/2}} x^{1-\beta/2-\delta} \frac{\partial u}{\partial x} dx dy \\
&\leq \frac{2}{|1-\beta|} \max_{\Omega} x^\delta \left(\int_{\Omega} \frac{u^2}{x^\beta} dx dy \right)^{\frac{1}{2}} \left(\int_{\Omega} x^{2-\beta-2\delta} u_x^2 \right)^{\frac{1}{2}} \\
&\leq C \|u\|_{L^2(\Omega, \beta)} \|u\|_{\mathcal{H}_\alpha}
\end{aligned}$$

where $\delta \geq 0$, and $\alpha = 2 - \beta - 2\delta$. Therefore, if $\beta \leq 2 - \alpha$, there exists a constant $C \equiv \frac{2}{|1-\beta|} \max_{\Omega} x^{1-(\alpha+\beta)/2}$ such that

$$\|u\|_{L^2(\Omega, \beta)} \leq C \|u\|_{\mathcal{H}_\alpha}. \quad (22)$$

If $\beta = 1$, integration by parts introduces a logarithmic factor, and the rest of the argument is the same. Since $C_0^1(\overline{\Omega})$ is dense in $\mathcal{H}_\alpha$, inequality (22) holds for all functions u in $\mathcal{H}_\alpha$. ■

For future purposes, we shall need the fact that the embeddings are actually compact. The following theorem is proved in the Appendix.

Theorem 5 *The space $\mathcal{H}_\alpha$, $0 < \alpha < 1$, is compactly embedded in $L^2(\Omega, \beta)$ for all $\beta < 2 - \alpha$.*

4.3 Weak Formulation of the Modified Problem

We now return to the problem of section 4.1. As described at the end of that section, we replace (14) by

$$\begin{aligned} - \left((\tilde{u} + x)v_x + \left(g_x - \frac{1}{2}\right)v \right)_x - v_{yy} &= f, \quad \text{in } \Omega \\ v &= 0, \quad \text{on } \partial\Omega, \end{aligned} \quad (23)$$

where f is now not considered to depend on g . Equation (20) defines the bivariate form associated with the weak formulation of this problem.

In this section we show that B_g is bounded on $\mathcal{H}_{\frac{1}{2}} \times \mathcal{H}_{\frac{1}{2}}$, then we define a nonlinear mapping $T : \mathcal{H}_{\frac{1}{2}} \rightarrow \mathcal{H}_{\frac{1}{2}}^*$ associated with B_g (T also depends on g). In Section 5, we establish that this mapping is continuous, pseudo-monotone and coercive, and then use a version of the Browder-Minty theorem to obtain the existence of a weak solution to (21) in $\mathcal{H}_{\frac{1}{2}}$.

To define T we first prove

Theorem 6 *For each $v \in \mathcal{H}_{\frac{1}{2}}(\Omega)$, and $w \in C_0^1(\Omega)$, the following inequality holds:*

$$|B_g(v, w)| \leq \tilde{C} \|v\|_{\mathcal{H}_{\frac{1}{2}}} \|w\|_{\mathcal{H}_{\frac{1}{2}}}, \quad (24)$$

where $\tilde{C}$ depends on g .

Proof: We apply the Cauchy-Schwartz inequality to $B_g(v, w)$ and use the embedding of $\mathcal{H}_{\frac{1}{2}}$ in $L^2(\Omega, \beta)$ with $\beta = 3/2$ in the following way:

$$\begin{aligned} |B_g(v, w)| &= \left| \int_{\Omega} \sqrt{\tilde{u} + x} v_x \sqrt{\tilde{u} + x} w_x + \left(g_x - \frac{1}{2}\right) \frac{v}{x^{1/4}} x^{1/4} w_x + v_y w_y \right| \\ &\leq \left(\int_{\Omega} (\tilde{u} + x) v_x^2 + \left(g_x - \frac{1}{2}\right)^2 \frac{v^2}{\sqrt{x}} + v_y^2 \right)^{\frac{1}{2}} \left(\int_{\Omega} (\tilde{u} + x) w_x^2 + \sqrt{x} w_x^2 + w_y^2 \right)^{\frac{1}{2}} \end{aligned} \quad (25)$$

We estimate each integral in the first factor of the product separately. To estimate the second integral we use the assumption that there exists a smooth function $\bar{g}$ with $g(x, y) = \sqrt{x} \bar{g}(x, y)$ and $\epsilon \leq \bar{g}(x, y) \leq 1/\epsilon$. Using this property of g we obtain

$$|g_x| \leq Cx^{-1/2} + D. \quad (26)$$

Therefore, we have

$$\int_{\Omega} (g_x - \frac{1}{2})^2 \frac{v^2}{\sqrt{x}} \leq C^2 \int_{\Omega} \frac{v^2}{x^{3/2}} + E \int_{\Omega} \frac{v^2}{x} + F \int_{\Omega} \frac{v^2}{\sqrt{x}}.$$

Now, from the embedding of $\mathcal{H}_{\frac{1}{2}}$ in $L^2(\Omega, \beta)$, for $\beta \leq 3/2$, we obtain

$$\int_{\Omega} (g_x - \frac{1}{2})^2 \frac{v^2}{\sqrt{x}} \leq K \|v\|_{\mathcal{H}_{\frac{1}{2}}}^2$$

where K depends on g .

We estimate the first integral in both factors of (25) using the bound on $\tilde{u}$:

$$\begin{aligned} \int_{\Omega} (\tilde{u} + x) v_x^2 &\leq \int_{\Omega} ((1/\epsilon)\sqrt{x} + x) v_x^2 \\ &\leq (1/\epsilon) \max_{\Omega} \{1 + \epsilon\sqrt{x}\} \int_{\Omega} \sqrt{x} v_x^2 \\ &\leq C_1(\epsilon) \|v\|_{\mathcal{H}_{\frac{1}{2}}}^2 \end{aligned}$$

The second factor of (25) is then bounded:

$$\int_{\Omega} (\tilde{u} + x) w_x^2 + \sqrt{x} w_x^2 + w_y^2 \leq C_2(\epsilon) \|w\|_{\mathcal{H}_{\frac{1}{2}}}^2.$$

Combining the last three estimates we see that there is a constant $\tilde{C}(g) > 0$ such that (24) holds. ■

Corollary 4 *The operator $w \mapsto B_g(v, w)$, defined on $w \in C_0^1(\overline{\Omega})$, is a bounded, linear functional for each fixed $v \in \mathcal{H}_{\frac{1}{2}}(\Omega)$. Moreover, it can be extended to a bounded, linear functional on $\mathcal{H}_{\frac{1}{2}}(\Omega)$.*

Corollary 5 *There exists a mapping*

$$\mathcal{H}_{\frac{1}{2}} \ni v \mapsto T(v) \in \mathcal{H}_{\frac{1}{2}}^*$$

such that

$$B_g(v, w) = \langle T(v), w \rangle,$$

for all v in $\mathcal{H}_{\frac{1}{2}}$, $w \in \mathcal{H}_{\frac{1}{2}}$. Furthermore, the nonlinear mapping is bounded.

Remark The space $\mathcal{H}_{\frac{1}{2}}^*$ is the negative norm weighted Sobolev space dual to $\mathcal{H}_{\frac{1}{2}}$ under the pairing

$$\langle v, w \rangle = \int_{\Omega} vw \, dx \, dy.$$

A weak solution of (21) is a function $v \in \mathcal{H}_{\frac{1}{2}}$ such that

$$B_g(v, w) = \langle T(v), w \rangle = \langle f, w \rangle, \quad \forall w \in \mathcal{H}_{\frac{1}{2}}.$$

To prove existence of a solution, we use a version of the classical Browder-Minty theorem (Theorem 7, below) which gives the existence of a weak solution if T is continuous, coercive, and pseudo-monotone.

We recall some definitions from [13].

Definition 4 Let X be a real, reflexive Banach space. We say that the mapping $T : X \rightarrow X^*$ is **coercive** if

$$\frac{\langle T(u), u \rangle}{\|u\|} \rightarrow \infty \quad \text{as } \|u\| \rightarrow \infty.$$

Definition 5 Let X be a real, reflexive Banach space. An operator $T : X \rightarrow X^*$ is called **pseudo-monotone** if T is bounded and if whenever

$$u_n \rightharpoonup u \quad \text{in } X$$

and

$$\limsup_{n \rightarrow \infty} \langle T(u_n), u_n - u \rangle \leq 0,$$

it follows that

$$\liminf_{n \rightarrow \infty} \langle T(u_n), u_n - v \rangle \geq \langle T(u), u - v \rangle \quad \forall v \in X.$$

Theorem 7 (The Browder-Minty Theorem; [13, Theorem 9.57, page 367]) Let X be a real, reflexive Banach space and suppose that $T : X \rightarrow X^*$ is continuous, coercive and pseudo-monotone. Then for every $f \in X^*$ there exists a solution $u \in X$ of the equation

$$T(u) = f.$$

In the next section we prove that the operator T defined in Corollary 5 of Theorem 6 is continuous, pseudo-monotone and coercive.

5 Existence of a Weak Solution

In this section we verify the properties of T , apply Theorem 7, and then show that the solution solves the original problem.

5.1 Properties of T

We show first that T is continuous by showing that each coefficient of T , as a function of $(x^{-3/4}u(x,y), x^{1/4}u_x(x,y), u_y(x,y)) \in (L^2)^3$, is continuous and bounded, and can be considered as a Nemytskii operator from $(L^2)^3$ to its dual. The continuity of Nemytskii operators (see [11], [13]), then implies the continuity of T . The following argument is adapted from [11] and [13] to the weighted Sobolev space we are dealing with here. We recall the definition of a Nemytskii operator.

Definition 6 For $(\xi, \omega) \in \Omega \times \mathbf{R}^m$, we say that a function $f(\xi, \omega)$ satisfies the Caratheodory conditions if the mapping $\omega \mapsto f(\xi, \omega)$ is continuous for almost every $\xi \in \Omega$, and $\xi \mapsto f(\xi, \omega)$ is measurable for every $\omega \in \mathbf{R}^m$. For such an f and any function $\omega : \Omega \rightarrow \mathbf{R}^m$,

$$a(\omega)(\xi) \equiv f(\xi, \omega(\xi)). \quad (27)$$

is called a Nemytskii operator.

Theorem 8 ([11],[13]) Let

$$\Omega \times \mathbf{R}^m \ni (\xi, \omega) \mapsto f(\xi, \omega) \in \mathbf{R}^m$$

satisfy the Caratheodory conditions, and

$$|f(\xi, \omega)| \leq C |\omega|^{p-1} + h(\xi),$$

where $p \in (1, \infty)$, $h \in L^q$, $(\frac{1}{p} + \frac{1}{q} = 1)$. Then the Nemytskii operator defined by (27) is a bounded and continuous map from L^p to L^q .

We now apply the calculus of Nemytskii operators to prove the continuity of T .

Theorem 9 The operator T is continuous.

Proof: To study the coefficients of T we write $\langle T(v), w \rangle$ as

$$\begin{aligned} \langle T(v), w \rangle &= \int_{\Omega} \frac{(\tilde{u} + x)x^{1/4}v_x}{\sqrt{x}} x^{1/4}w_x \\ &\quad + \int_{\Omega} (g_x - \frac{1}{2}) \frac{v}{x^{3/4}} \sqrt{x} x^{1/4}w_x + \int_{\Omega} v_y w_y. \end{aligned}$$

Since $v \in \mathcal{H}_{\frac{1}{2}}$ implies $v \in L^2(\Omega, \beta)$ with $\beta = 3/2$, we know that $v/(x^{3/4}) \in L^2$. Therefore, we may consider the coefficients of T as functions from $(L^2)^3$ to L^2 by defining the symbol vector

$$\omega = (\omega_1, \omega_2, \omega_3) = (\frac{v}{x^{3/4}}, x^{1/4}v_x, v_y).$$

Then the coefficients of the symbol of T become

$$\begin{aligned} a_1(x, y, \omega) &= \frac{(\tilde{\omega}_1 + x)\omega_2}{\sqrt{x}}, \\ a_2(x, y, \omega) &= (g_x - \frac{1}{2})\sqrt{x}\omega_1, \\ a_3(x, y, \omega) &= \omega_3. \end{aligned}$$

Here, $\tilde{\omega}_1$ denotes the function $\tilde{\omega}_1 \equiv \omega_1 x^{3/4} + g$, which satisfies the bounds

$$\epsilon \sqrt{x} \leq \tilde{\omega}_1 \leq (1/\epsilon) \sqrt{x}.$$

Now, all three functions a_1 , a_2 , and a_3 are continuous for almost every $(x, y) \in \Omega$, and measurable for every $\omega \in \mathbf{R}^3$. Moreover,

$$\begin{aligned} |a_1(x, y, \omega)| &\leq \frac{(1/\epsilon)\sqrt{x} + x}{\sqrt{x}} |\omega_2| \leq C_1 |\omega_2|, \\ |a_2(x, y, \omega)| &\leq \left| g_x - \frac{1}{2} \right| \sqrt{x} |\omega_1| \leq (\frac{K}{\sqrt{x}} + D) \sqrt{x} |\omega_1| \leq C_2 |\omega_1|, \\ |a_3(x, y, \omega)| &= |\omega_3|. \end{aligned}$$

This shows that the conditions of Theorem 8 hold with $p = 2$, and thus the symbol is continuous on a weighted L^2 space and the corresponding map T is continuous from $\mathcal{H}_{\frac{1}{2}}$ to $\mathcal{H}_{\frac{1}{2}}^*$. ■

To verify pseudo-monotonicity we use the approach presented in [13]; we show that the operator T is of calculus of variations type, which implies that T is pseudo-monotone. We state the definition and theorem from [13, pages 367-368].

Definition 7 An operator $T : X \rightarrow X^*$ is said to be of calculus of variations type if it is bounded, and it has the representation $T(v) = \hat{T}(v, v)$ where the mapping from $X \times X$ to X^* given by $(v, z) \mapsto \hat{T}(v, z)$ satisfies the following properties:

1. The mappings $z \mapsto \hat{T}(v, z)$ and $v \mapsto \hat{T}(v, z)$ are bounded and continuous.
2. For each $v \in X$,

$$\langle \hat{T}(v, v) - \hat{T}(v, z), v - z \rangle \geq 0, \quad \forall z \in X. \quad (28)$$

3. If $v_n \rightharpoonup v$ in X , and $\langle \hat{T}(v_n, v_n) - \hat{T}(v_n, v), v_n - v \rangle \rightarrow 0$, then for every $z \in X$

$$\hat{T}(v_n, z) \rightharpoonup \hat{T}(v, z) \quad \text{in } X^*.$$

4. If $v_n \rightharpoonup v$ in X , and $\hat{T}(v_n, z) \rightarrow \psi$ in X^* , then

$$\langle \hat{T}(v_n, z), v_n \rangle \rightarrow \langle \psi, v \rangle.$$

Theorem 10 ([13]) If T is of calculus of variations type, then T is pseudo-monotone.

To show that T is of calculus of variations type, we define $\overline{B}_g(v, z, w)$ by

$$\overline{B}_g(v, z, w) \equiv \int_{\Omega} \left((\tilde{u} + x) z_x + \left(g_x - \frac{1}{2} \right) z \right) w_x + z_y w_y.$$

Using the same argument as in Theorem 6, we can show that there is a mapping $(v, z) \mapsto \hat{T}(v, z)$, from $\mathcal{H}_{\frac{1}{2}} \times \mathcal{H}_{\frac{1}{2}}$ to $\mathcal{H}_{\frac{1}{2}}^*$, such that

$$\overline{B}_g(v, z, w) = \langle \hat{T}(v, z), w \rangle \quad \text{for all } w \in \mathcal{H}_{\frac{1}{2}}, \quad (29)$$

and $T(v) = \hat{T}(v, v)$. Note that $\overline{B}$ is linear in its second and third arguments, while its dependence on v is bounded (by the cut-off in $\tilde{u}$) and does not involve derivatives. Thus it is straightforward to verify the conditions of Definition 7, as we now do.

We need the following proposition.

Proposition 4 *If $v_n \rightarrow v$ in $L^2(\Omega, \beta)$, then $\tilde{u}_n \rightarrow \tilde{u}$ in $L^2(\Omega, \beta)$, where $\tilde{u}_n$ and $\tilde{u}$ are defined as in (19).*

Proof: Let $\Omega^\epsilon = \{x \in \Omega \mid v + g < \epsilon \sqrt{x}\}$, $\Omega^{1/\epsilon} = \{x \in \Omega \mid v + g > (1/\epsilon) \sqrt{x}\}$, and let Ω^{v+g} denote the subset of Ω on which $\tilde{u} = v + g$. Similarly, define the sets Ω_n^ϵ , $\Omega_n^{1/\epsilon}$, and $\Omega_n^{v_n+g}$, corresponding to v_n . Now, we write the $L^2(\Omega, \beta)$ -norm of $\tilde{u}_n - \tilde{u}$ as follows:

$$\begin{aligned} \|\tilde{u}_n - \tilde{u}\|^2 &= \int_{\Omega_n^\epsilon \cap \Omega^\epsilon} (\tilde{u}_n - \tilde{u})^2 d\mu + \int_{\Omega_n^{1/\epsilon} \cap \Omega^{1/\epsilon}} (\tilde{u}_n - \tilde{u})^2 d\mu \\ &\quad + \int_{\Omega_n^{v_n+g} \cap \Omega^{v+g}} (\tilde{u}_n - \tilde{u})^2 d\mu + \int_{\Omega^C} (\tilde{u}_n - \tilde{u})^2 d\mu, \end{aligned}$$

where Ω^C denotes the subset of Ω on which $\tilde{u}_n$ and $\tilde{u}$ are defined using different formulas in definition (19). Since $v_n \rightarrow v$ in $L^2(\Omega, \beta)$, the measure of the set Ω^C tends to 0. Also, since $\tilde{u}_n$ and $\tilde{u}$ are bounded on Ω^C the integral tends to zero, that is,

$$\int_{\Omega^C} (\tilde{u}_n - \tilde{u})^2 d\mu \leq \int_{\Omega^C} (|\tilde{u}_n| + |\tilde{u}|)^2 d\mu \leq K \mu\{\Omega^C\} \rightarrow 0.$$

Now, since $v_n \rightarrow v$ in $L^2(\Omega, \beta)$, we have

$$\int_{\Omega_n^{v_n+g} \cap \Omega^{v+g}} (\tilde{u}_n - \tilde{u})^2 d\mu = \int_{\Omega_n^{v_n+g} \cap \Omega^{v+g}} (v_n - v)^2 d\mu \rightarrow 0.$$

Finally, since the integrands in the remaining integrals are identically 0, we see that $\|\tilde{u}_n - \tilde{u}\| \rightarrow 0$. $\blacksquare$

Now we can prove

Theorem 11 *The operator T is of calculus of variations type.*

Proof: Boundedness of T follows from Theorem 6; we write $T(v) = \hat{T}(v, v)$, where $\hat{T}(v, z)$ is given by (29). We show that the operator $\hat{T}(v, z)$ satisfies the conditions of Definition 7:

1. Boundedness of the mapping $z \mapsto \hat{T}(v, z)$ follows in the same way as the boundedness of $B(v, w)$, proved in Theorem 6. Continuity follows since $z \mapsto \hat{T}(v, z)$ is linear. Boundedness of the mapping $v \mapsto \hat{T}(v, z)$ also follows from Theorem 6. Using the same arguments as in Theorem 9 to prove continuity of T , we can show that $v \mapsto \hat{T}(v, z)$ is a Nemytskii operator, and, therefore,

is continuous.

2. Now we prove that inequality (28) holds. Recall that we have extended the boundary data g to a function on Ω with the property that $g_{xx} \leq 0$. Therefore, we have

$$\begin{aligned} & \langle \hat{T}(v, v) - \hat{T}(v, z), v - z \rangle \\ &= \int_{\Omega} (\tilde{u} + x)(v_x - z_x)^2 + (g_x - \frac{1}{2})(v - z)(v - z)_x + (v_y - z_y)^2 \quad (30) \end{aligned}$$

We estimate the middle term in (30) by integrating by parts and taking into account that v and z vanish on $\partial\Omega$, to obtain

$$\int_{\Omega} (g_x - \frac{1}{2})(v - z)(v - z)_x = - \int_{\Omega} \frac{1}{2} g_{xx} (v - z)^2 \geq 0. \quad (31)$$

Since $v - z \in L^2(\Omega, \beta)$ for $\beta = 3/2$, and $g_{xx} \leq 0$, this integral is bounded and non-negative. Therefore, (28) follows from (30) and (31).

3. Let $v_n \rightharpoonup v$ in $\mathcal{H}_{\frac{1}{2}}$. By the compactness of the embedding of $\mathcal{H}_{\frac{1}{2}}(\Omega)$ in $L^2(\Omega, \beta)$, for $\beta < 3/2$, we have $\tilde{u}_n \rightarrow \tilde{u}$ (strongly) in $L^2(\Omega, \beta)$. Therefore, even without the hypothesis $\langle \hat{T}(v_n, v_n) - \hat{T}(v_n, v), v_n - v \rangle \rightarrow 0$ we obtain $\overline{B}_g(v_n, z, w) \rightarrow \overline{B}(v, z, w)$, or equivalently,

$$\hat{T}(v_n, z) \rightharpoonup \hat{T}(v, z) \quad \text{in } \mathcal{H}_{\frac{1}{2}}^*.$$

4. Let $v_n \rightharpoonup v$ in $\mathcal{H}_{\frac{1}{2}}$, and $\hat{T}(v_n, z) \rightharpoonup \psi$ in $\mathcal{H}_{\frac{1}{2}}^*$. We need to show that

$$\langle \hat{T}(v_n, z), v_n \rangle = \langle \hat{T}(v_n, z), v \rangle + \langle \hat{T}(v_n, z), v_n - v \rangle \rightarrow \langle \psi, v \rangle.$$

It suffices to show that

$$\langle \hat{T}(v_n, z), v_n - v \rangle \rightarrow 0.$$

Now,

$$\begin{aligned} \langle \hat{T}(v_n, z), v_n - v \rangle &= \int_{\Omega} \left((\tilde{u}_n + x)z_x + (g_x - \frac{1}{2})z \right) (v_n - v)_x + z_y(v_n - v)_y \\ &= \int_{\Omega} \sqrt{x}h_n(x, y) z_x (v_n - v)_x + \int_{\Omega} (g_x - \frac{1}{2}) z (v_n - v)_x \\ &\quad + \int_{\Omega} z_y (v_n - v)_y, \end{aligned}$$

where h_n is a bounded function, since $\tilde{u}_n \leq (1/\epsilon)\sqrt{x}$. Now Proposition 4, with $\beta = 1$, shows that $h_n \rightarrow h$ in $L^2(\Omega)$, and this, along with weak convergence of v_n in $\mathcal{H}_{\frac{1}{2}}$, shows that the first integral converges to zero; similarly, the third integral converges to zero by weak convergence of v_n . In order to estimate the second integral, we write

$$\int_{\Omega} (g_x - \frac{1}{2}) z (v_n - v)_x = \int_{\Omega} (g_x - \frac{1}{2}) x^{3/4} \frac{z}{x^{3/4}} (v_n - v)_x \quad (32)$$

and again use (26); since $z/x^{3/4} \in L^2$ and $|g_x|$ is bounded by $x^{-1/2}$, the right side of (32) converges to zero, as a consequence of the weak convergence of v_n to v .

We have proved that the mapping $\hat{T}$ is of calculus of variations type. ■

By Theorem 10 we now have

Corollary 6 *The operator T is pseudo-monotone.*

We have now shown that the mapping T is continuous and pseudo-monotone. In order to satisfy all the hypotheses of Theorem 7 we have to check that T is coercive.

Theorem 12 *The operator T is coercive.*

Proof: We shall show that there is a constant $M > 0$ such that

$$|B_g(v, v)| \geq M \|v\|_{\mathcal{H}_{\frac{1}{2}}}^2.$$

Recall that

$$B_g(v, v) = \int_{\Omega} (\tilde{u} + x) v_x^2 + (g_x - \frac{1}{2}) v v_x + v_y^2. \quad (33)$$

We estimate the second integral using the assumptions on g . Integrating by parts and using $v = 0$ on the boundary, we obtain

$$\int_{\Omega} (g_x - \frac{1}{2}) v v_x = - \int_{\Omega} \frac{1}{2} g_{xx} v^2 \geq 0.$$

This integral is bounded since $v \in L^2(\Omega, \beta)$, for $\beta = 3/2$. The first integral in (33) can be estimated using the fact that $\tilde{u}$ is bounded from below by $\epsilon\sqrt{x}$. Therefore, there exists a $C_1 > 0$ such that

$$\int_{\Omega} (\tilde{u} + x) v_x^2 \geq C_1 \int_{\Omega} \sqrt{x} v_x^2.$$

Combining the last two estimates we obtain

$$\int_{\Omega} (\tilde{u} + x) v_x^2 + (g_x - \frac{1}{2}) v v_x \geq C_1 \int_{\Omega} \sqrt{x} v_x^2.$$

Therefore, we have

$$|B_g(v, v)| \geq M \int_{\Omega} \sqrt{x} v_x^2 + v_y^2 = M \|v\|_{\mathcal{H}_{\frac{1}{2}}}^2,$$

where $M = \min\{C_1, 1\} > 0$. ■

By Theorem 9, Corollary 6, and Theorem 12, the operator T defined in Corollary 5 of Theorem 6 is continuous, pseudo-monotone and coercive. Therefore, Theorem 7 gives a weak solution v in the space $\mathcal{H}_{\frac{1}{2}}$.

Theorem 13 *For every $f \in \mathcal{H}_{\frac{1}{2}}^*$, there exists a function $v \in \mathcal{H}_{\frac{1}{2}}$, such that*

$$B_g(v, w) = \int_{\Omega} f w, \quad \text{for all } w \in \mathcal{H}_{\frac{1}{2}}. \quad (34)$$

In particular, we may now suppose f is given by (15), then the solution of (21), which we have found, is the solution we want, save for the possibility that B_g does not coincide with B for our weak solution v . We rule this out in the next section.

5.2 Bounds on the Weak Solution

In this section we show that a weak solution v of the problem defined by the bivariate form B_g satisfies the bounds

$$\epsilon \sqrt{x} \leq v + g \leq (1/\epsilon) \sqrt{x}.$$

In order to show this we construct sub- and super-solutions. We shall show that for every $\delta < \epsilon$ and $M > 1/\epsilon$, the functions $\delta\sqrt{x} - g$ and $M\sqrt{x} - g$ are sub- and super-solutions respectively to (23), where f satisfies (15). This is equivalent to showing that $\delta\sqrt{x}$ and $M\sqrt{x}$ are sub- and super-solutions to

$$\left((\tilde{u} + x) u_x - \frac{1}{2} u \right)_x + u_{yy} = 0 \quad \text{in } \Omega, \quad (35)$$

$$u = g \quad \text{in } \partial\Omega. \quad (36)$$

Here, $\tilde{u}$ is defined by (19), where we set $v + g = u$.

Theorem 14 *The functions $\delta\sqrt{x}$ and $M\sqrt{x}$ are sub- and super-solutions respectively to problem (35), (36).*

Proof: If $u < \epsilon\sqrt{x}$ or $u > (1/\epsilon)\sqrt{x}$, then the equation (35) satisfied by u is linear. Let $w = M\sqrt{x}$; then a calculation shows that w satisfies (35) for any M . Since $M\sqrt{x} > g$ on the boundary, w is a super-solution of problem (35), (36). Therefore, any solution u of the problem satisfies $u \leq (1/\epsilon)\sqrt{x}$, by the comparison principle, Theorem 2. Similarly, we can show that $\delta\sqrt{x}$, for $\delta \leq \epsilon$, is a sub-solution, and therefore $u \geq \epsilon\sqrt{x}$. ■

Corollary 7 *(Existence of a solution to (14)) The weak solution v of Theorem 13 satisfies*

$$\epsilon\sqrt{x} \leq v + g \leq (1/\epsilon)\sqrt{x}.$$

and is a solution of

$$B(v, w) = \int_{\Omega} f w, \quad \forall w \in \mathcal{H}_{\frac{1}{2}},$$

where the bivariate form B is defined by (16) and f by (15).

6 Conclusions

We have proved existence of a weak solution in a weighted Sobolev space $\mathcal{H}_{\frac{1}{2}}$ for problem (3), under assumptions on g . By the classical theory of elliptic equations, this weak solution is actually C^2 everywhere in Ω . Although for uniformly elliptic operators the solution is smooth all the way up to the boundary, we have shown that the degeneracy in this equation induces singular behavior of solutions along the degenerate boundary; in this case, a square-root type singularity.

The growth assumptions on the boundary data g are forced by the need to obtain a coercive operator on the weighted Sobolev space. We believe these assumptions are consistent with the actual behavior of the solution to the original problem. We consider this issue in work now in progress, [3].

Some interesting questions remain open. Our immediate objective is to apply this result to the motivating problem of self-similar solutions to the unsteady TSD equation. To this end, we would like to weaken the hypotheses on g ; alternatively, we could attempt to formulate the free-boundary problem

which connects the elliptic and hyperbolic regions in a way that respects these hypotheses, and show that this problem has a solution. In any case, but especially if we adopt this approach, the question of uniqueness remains open.

The theory in this paper also may be extended to unbounded domains Ω : an appropriate extension for our problem would be to replace the rectangle by a quarter-plane. In this case, the y -axis will remain the degenerate boundary. The boundary condition on the x -axis might continue to contain upper and lower square-root type bounds on its growth near $x = 0$, but a new far-field condition would be needed.

Beyond extensions to the TSD equation, the approach in this paper might be useful in studying other degenerate elliptic equations of Keldysh type. Our interest was focussed on the model problem at hand, but the dependence of the singularity on dimension and on the coefficients of the operator could be studied asymptotically, as in Section 2, and corresponding theorems formulated.

A Appendix: Compactness of the Embedding

Here we assemble the ingredients for the proof of compactness in Theorem 5 in Section 4.2.

The following lemma is a version of the Arzelà-Ascoli theorem for the spaces $L^2(\Omega, \beta)$.

Lemma 1 (*Adams, [1]*) *A bounded subset $K \subset L^2(\Omega, \beta)$ is precompact in $L^2(\Omega, \beta)$ if and only if for every $\epsilon > 0$, there exists a $\delta > 0$, and a subset $G \subset\subset \Omega$ such that for every $u \in K$ and every $h \in \mathbb{R}^2$ with $|h| < \delta$*

$$\int_{\Omega} |u(x+h) - u(x)|^2 d\mu < \epsilon^2, \quad \text{and} \\ \int_{\Omega-G} |u(x)|^2 d\mu < \epsilon^2.$$

Using this Lemma, the following Theorem is also proved in [1, page 33]:

Theorem 15 Let $K \subset L^2(\Omega, \beta)$ be a bounded subset of $L^2(\Omega, \beta)$. Suppose that there exists a sequence $\{\Omega_j\}$ of subdomains of Ω having the following properties:

1. For each j , $\Omega_j \subset \Omega_{j+1}$,
2. For each j the set of restrictions to Ω_j of the functions in K is precompact in $L^2(\Omega_j, \beta)$,
3. For every $\epsilon > 0$ there exists j such that

$$\int_{\Omega - \Omega_j} |u|^2 d\mu < \epsilon, \quad \forall u \in K. \quad (37)$$

Then K is precompact in $L^2(\Omega, \beta)$.

Using this Theorem we show that the embeddings from Theorem 4 are compact.

Proof of Theorem 5: Continuity of the embedding is proved in Theorem 4. In order to prove the compactness of the embedding we use Theorem 15 as follows. Let K be a bounded set in $\mathcal{H}_\alpha$. Then, there exists a constant D such that for each $u \in K$,

$$\|u\|_{\mathcal{H}_\alpha} \leq D.$$

From the continuity of the embedding, it follows that K is bounded in $L^2(\Omega, \beta)$.

Let $\Omega_j \subset \Omega$ be the subset of Ω such that $x \geq 1/j$ for $(x, y) \in \Omega_j$. Obviously, $\Omega_j \subset \Omega_{j+1}$.

Denote by $\mathcal{H}_\alpha(\Omega_j)$ the space of restrictions of functions in $\mathcal{H}_\alpha(\Omega)$ to the set Ω_j . Define $H^1(\Omega_j)$, $L^2(\Omega_j)$, and $L^2(\Omega_j, \beta)$ the same way. (Here H^1 is the standard Sobolev space.) The norms in $H^1(\Omega_j)$ and $\mathcal{H}_\alpha(\Omega_j)$ are equivalent. The same is true for the norms in $L^2(\Omega_j)$ and $L^2(\Omega_j, \beta)$, for every measure μ for which $\beta < 2 - \alpha$. Since $H^1(\Omega_j)$ is compactly embedded in $L^2(\Omega_j)$, it follows that $\mathcal{H}_\alpha(\Omega_j)$ is also compactly embedded in $L^2(\Omega_j, \beta)$, $\beta < 2 - \alpha$.

Therefore, it suffices to show that for every $\epsilon > 0$ there exists j , such that (37) holds. Let $\epsilon > 0$. Assume for the moment that $\beta \neq 1$. Take j such that

$$\epsilon = \frac{4}{|1 - \beta|^2} D^2 \left(\frac{1}{j} \right)^{2 - \alpha - \beta}. \quad (38)$$

Then, using an estimate similar to that in the proof of Theorem 4 with

$\delta = 1 - (\alpha + \beta)/2 > 0$, we obtain:

$$\begin{aligned} \int_{\Omega-\Omega_j} |u|^2 d\mu &\leq \frac{2}{|1-\beta|} \left(\frac{1}{j}\right)^{1-(\alpha+\beta)/2} \left(\int_{\Omega-\Omega_j} \frac{u^2}{x^\beta} dx dy\right)^{\frac{1}{2}} \left(\int_{\Omega-\Omega_j} x^\alpha u_x^2\right)^{\frac{1}{2}} \\ &\leq \frac{2}{|1-\beta|} D \left(\frac{1}{j}\right)^{1-(\alpha+\beta)/2} \left(\int_{\Omega-\Omega_j} |u|^2 d\mu\right)^{\frac{1}{2}} \end{aligned}$$

Therefore,

$$\int_{\Omega-\Omega_j} |u|^2 d\mu \leq \frac{4}{|1-\beta|^2} D^2 \left(\frac{1}{j}\right)^{2-\alpha-\beta} = \epsilon.$$

as required. If $\beta = 1$, one modifies the choice of ϵ in (38). ■

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A discussion with Joe McKenna at the Fields Institute led to an independent investigation of a problem similar to the boundary value problem for equation (3): Choi, Lazer and McKenna, [4], consider a family of equations which includes a singular term of the form $uu_{xx} + u_{yy}$ but a different lower order term. They use a comparison principle and a priori estimates to bound the solution and then use compactness arguments to show existence in a bounded, convex domain with a smooth boundary. The solution they obtain has the same type of square root singularity.

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