

Control and Elimination of Cracking of AlGaN Using Low-Temperature AlGaN

Interlayers

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Abstract

We demonstrate that the insertion of low-temperature (LT) AlGaN interlayers is effective in reducing mismatch-induced tensile stress and suppressing the formation of cracks during growth of AlGaN directly upon GaN epilayers. Stress evolution and relaxation is monitored using an *in-situ* optical stress sensor. The combination of *in-situ* and *ex-situ* characterization techniques enables us to determine the degree of pseudomorphism in the interlayers. It is observed that the elastic tensile mismatch between AlGaN and GaN is mediated by the relaxation of interlayers; the use of interlayers offers tunability in the in-plane lattice parameters.

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Most of the III-nitride devices to date have been grown on thick (2~4 μm) GaN layers nucleated on heavily mismatched substrates (such as Al_2O_3 or SiC). The thick GaN layer leads to improved structural quality and serves as a (nearly relaxed) pseudo-substrate on which active layers consisting of AlGaInN heterostructures are grown. Unlike zincblende semiconductors, misfit dislocations do not easily afford stress relief during strained heteroepitaxial growth of the (0001)-oriented, wurtzite, III-nitrides. Difficulty of misfit-dislocation glide in wurtzite crystals in response to a biaxial strain field is due to the absence of effective slip systems¹ and is manifested by a very large critical thickness for both AlGaIn and GaInN epilayers grown on GaN.² We have shown that relaxation of tensilely stressed AlGaIn (on GaN) proceeds through the propagation of incipient cracks, which in turn activates the slip systems for misfit dislocations around the crack tips.³ Even though III-nitride devices are known for their relative insensitivity to a high density of dislocations,⁴ cracking due to the tensile mismatch between AlGaIn and the GaN pseudo-substrate tends to severely deteriorate device performance.⁵ Various types of *interlayers*, generally referred to as a low-temperature (LT) grown layer inserted between two layers grown at regular, high temperatures (HT), have been reported to influence the structural properties of the III-Nitrides. Nakamura *et al.*⁶ reported the use of an InGaIn layer before the growth of AlGaIn cladding layers (in purple laser diode structures) to avoid cracking. Iwaya *et al.*⁷ reported the use of low-temperature (LT) deposited AlN and GaN interlayers for the reduction of threading dislocations. Very recently, Amano *et al.*⁸ demonstrated that the introduction of a LT AlN interlayer on a high-temperature (HT) grown GaN epilayer facilitates subsequent crack-free growth of HT AlGaIn. Since the LT interlayers (typically 100~300Å thick) are embedded in thick

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(> 1 μm) heterostructures, it is very difficult to extract information concerning the strain-state in the interlayers. In this work, we investigate the strain relaxation of $\sim 150\text{-\AA}$ -thick LT AlN, GaN, and ternary-alloy interlayers inserted into HT-GaN/AlGaN heterostructures. LT-interlayer strain was quantified using *in-situ* film-stress measurements in combination with *ex-situ* characterization of alloy composition. We found that the in-plane lattice parameter of the partially strain-relaxed LT interlayer can be directly tuned to control and compensate growth stress during subsequent growth of AlGaN heterostructures of various compositions.

Five samples (labeled A-E herein) were grown by metal-organic vapor-phase epitaxy (MOVPE) on (0001) Al_2O_3 substrates. The growth sequence for these samples was: (I) nucleation of 200-\AA -thick GaN at $550\text{ }^\circ\text{C}$, (II) growth of a $1.0\text{-}\mu\text{m}$ -thick GaN pseudo-substrate at $1050\text{ }^\circ\text{C}$, (III) followed by a 150-\AA -thick LT-AlGaN interlayer deposited at $700\text{ }^\circ\text{C}$, and (IV) capped with a $0.9\text{-}\mu\text{m}$ -thick $\text{Al}_{0.2}\text{Ga}_{0.8}\text{N}$ layer again grown at $1050\text{ }^\circ\text{C}$ (all layers Si-doped, $n\sim 2\times 10^{18}\text{ cm}^{-3}$). The schematics of layering and growth sequence are shown in Figure 1. The composition of the LT $\text{Al}_x\text{Ga}_{1-x}\text{N}$ interlayer for samples A-E was intentionally varied from $x_{LT} = 0$ (GaN) to $x_{LT} = 1$ (AlN), respectively; the samples are otherwise similar. In addition to using the binary AlN interlayers as was reported previously⁸, we extended our investigation into ternary AlGaN interlayers based on two considerations: 1) AlN is normally insulating and may not be compatible with devices requiring vertical current transport, and 2) the employment of AlGaN should render direct tunability in the in-plane lattice parameter. Our growth procedures for GaN on sapphire and a description of the MOVPE reactor appear elsewhere.⁹ The LT AlGaN interlayers and HT AlGaN layers of present interest were grown at 40 Torr. NH_3 and H_2

flows were set at 2.5 and 5 l/min, respectively. Trimethylgallium and trimethylaluminum were employed as metal-organic precursors.

The Al composition of the thin LT interlayers (x_{LT}) was determined by Rutherford backscattering spectrometry (RBS)¹⁰ using a 2.5 MeV $^4\text{He}^+$ ion beam. The RBS was performed on a separate set of 400-Å-thick LT-AlGaN calibration samples grown in the same manner as samples A-E, but without HT-AlGaN capping layers. The composition of the much thicker HT AlGaN layers (x_{HT}) of samples A-E was directly measured using high-resolution x-ray diffractometry (XRD). Symmetric (0002) and asymmetric (20-24) reflections were used to determine both in-plane (a) and out-of-plane (c) lattice parameters. Once the strained a and c lattice parameters were measured, corresponding unstrained lattice parameters, elastic strains, and Al compositions were computed from elasticity theory¹¹ under the assumption that Vegard's law holds for the unstrained-lattice parameters¹² and elastic constants.¹³ The resulting RBS (for x_{LT}) and XRD (for x_{HT}) compositions are listed in Table I.

Real time *in-situ* stress monitoring based on wafer-curvature measurements was performed with a multi-beam optical stress sensor (MOSS)¹⁴ modified for use on our reactor. To determine the wafer curvature, the divergence of multiple initially parallel laser beams is measured on a CCD camera after reflection of the beams from the film/substrate surface. Changes in wafer curvature induce a proportional change in the beam spacing on the camera. This technique provides a direct measurement of the stress-thickness product during MOVPE.¹⁵ Slopes of the (stress) \times (thickness) traces versus time during deposition cycles can be converted to instantaneous stress once the time scale is

converted to thickness scale using growth rates derived from *in-situ* reflectance measurements.¹⁵

The evolution of growth stress was measured by MOSS and the results are shown in Fig. 2. Similar to our previous observations,³ the growth of GaN on LT GaN buffers (Region II in Fig. 2) exhibited a characteristic residual tensile stress (σ_{GaN}) that varies between 0.35 and 0.56 GPa (see Table I). A much wider difference, on the other hand, was observed in the variation of the initial growth stress ($\sigma_{HT-AlGaN}$) during growth of HT AlGaN layers (Region IV in Figure 2) on the interlayers. Very steep positive slopes, accompanied by successive step-wise stress relaxation, were seen in Samples A and B. As the Al concentration in the LT AlGaN interlayers was increased from A (GaN) to E (AlN), the magnitude of the tensile stress of the HT AlGaN gradually decreased. The sample grown using a binary LT AlN interlayer, Sample E, exhibited a negative slope indicative of an initial compressive stress.

Surface-cracking of samples A-E was examined using a differential interference-contrast Nomarski microscope. We also calculated the critical thickness for cracking (h_g) of the HT AlGaN using the initial stress ($\sigma_{HT-AlGaN}$) produced by growth on the partially relaxed LT interlayers on HT GaN. The results for h_g given by Griffith's criterion for crack propagation³ are compared with the Nomarski micrographs in Figure 3. Recalling that the HT Al_{0.2}Ga_{0.8}N is 0.9 μ m thick, we find consistent results. As expected, the LT GaN interlayer, which produces no relaxation of HT-AlGaN coherency stress, has a critical thickness for fracture much smaller than the HT-layer thickness. The result is the highest degree of cracking as seen in Fig 3(a). Addition of intermediate amounts of Al ($x_{LT}=0.35-0.42$) to the LT interlayer reduces the HT-AlGaN coherency stress and

increases $h_g \sim$ fourfold. Cracking is still observed in Figs. 3(b) and 3(c) since the HT-AlGaN thickness is still twice of h_g . Increasing the LT-interlayer composition to $x_{LT}=0.62$ raises h_g by \sim thirtyfold -- now far above the HT-AlGaN thickness. Examining Fig. 3(d), the HT-Al_{0.2}Ga_{0.8}N growth is appropriately crack-free. Finally, as the LT interlayer becomes pure AlN, the HT AlGaN moves into compression; crack propagation can no longer be driven by composition-induced stress at any layer thickness as seen in Fig. 3(e).

A key parameter in employing the interlayer scheme is the degree of pseudomorphism of the LT interlayers. Information of the in-plane lattice constants of the interlayers immediately *before* the nucleation of HT AlGaN layers is required to determine the degree of coherency. This lattice information is replicated during initial and presumably coherent growth of HT AlGaN layers and can be "captured" with the availability of the *in-situ* stress sensor before relaxation takes place; the HT AlGaN test layers thus serve as "probes" to yield the lattice parameters of the underlying interlayers. The methodology of extracting the degree of lattice strain relaxation of the interlayers is presented in the following. In brief, the use of XRD and RBS provide the information of the Al concentrations in the HT AlGaN (x_{HT}) and LT AlGaN (x_{LT}) interlayers, respectively. The slopes measured via MOSS during the initial growth of HT AlGaN layers and during HT GaN yielded instantaneous growth stresses, $\sigma_{HT-AlGaN}$ and σ_{GaN} , respectively. The measured stresses were converted to strains ($\epsilon_{HT-AlGaN}$ and ϵ_{GaN}) with information of the biaxial bulk modulus ($\epsilon = \sigma/M$, $M = 449.6 \text{ GPa}^{15}$). Assuming an initial coherent growth of HT AlGaN test layers on the LT interlayers, the in-plane lattice constant of the LT AlGaN interlayers can be expressed as $a_{LT-AlGaN} = a_{HT-AlGaN} = (1 + \epsilon_{HT-$

$a_{\text{AlGaN}} a_0(x_{\text{HT}})$, where $a_0(x_{\text{HT}})$ is the in-plane lattice constant of the free-standing AlGaN. With the knowledge of $a_{\text{LT-AlGaN}}$, $a_0(x_{\text{LT}})$, and $a_{\text{GaN}} = (1 + \epsilon_{\text{GaN}}) a_0(\text{GaN})$, the fraction of the strain relief can be approximated to be $[a_{\text{GaN}} - a_{\text{LT-AlGaN}}]/[a_{\text{GaN}} - a_0(x_{\text{LT}})]$. All the parameters are summarized in Table I.

The fact that $a_{\text{HT-AlGaN}}$ for Sample A (the sample with a LT GaN interlayer), which was derived from a combination of stress ($\sigma_{\text{HT-AlGaN}}$) during AlGaN growth and Al concentration (x_{HT}) from XRD, approaches a_{GaN} lends strong support to both the assumption of an initial coherent HT AlGaN growth and the validity of current treatment combining *in-situ* and *ex-situ* characterizations. It is clear from Table I (columns $a_{\text{LT-AlGaN}}$ and a_{GaN}) that crack suppression through stress engineering is based on the ability to vary the in-plane lattice parameters of the LT interlayers. Under the growth conditions employed in this work, the deposition of LT $\text{Al}_x\text{Ga}_{1-x}\text{N}$ interlayers produces templates with in-plane lattice constants equal to those of free-standing $\text{Al}_y\text{Ga}_{1-y}\text{N}$ with $y \sim x/3$. The observed partial lattice relaxation in the LT interlayers was possibly accomplished through the nucleation and gliding of edge-dislocations on the basal (0001) plane³ between LT interlayers and HT-GaN. The geometrical constraints for nucleation and gliding of misfit dislocations during planar, coherent heteroepitaxy in wurtzite systems¹ are likely lifted due to the alteration of microstructural properties. Atomic force microscopy performed on as-grown interlayers revealed the presence of non-planar discrete and/or interconnected islands. Such a change in microstructure causes a highly inhomogeneous, non-biaxial distribution of the strain field, which makes the admission of misfit dislocations energetically favorable. Alternatively, strain relaxation of the LT interlayers might occur through initial cracking and subsequent dislocating; the final,

crack-free HT AlGa_N layers (Samples A and B) would then be the result of lateral growth over the cracks. We note that in the cases of crack-free growth on the AlGa_N interlayers (Al>60%, Samples A and B), the structural quality of the 0.9 μm AlGa_N layers was as good as the underlying Ga_N pseudo-substrates as indicated by the rocking curve linewidths using (0004) diffraction (around 320 arcsec for both AlGa_N and Ga_N).

In conclusion, we found that the use of LT AlGa_N interlayers is effective in “re-defining” the in-plane lattice parameter through strain relaxation. The interlayers serve to mediate the elastic tensile mismatch between the adjacent layers and extend the cracking critical thickness. It is anticipated that the interlayer scheme will greatly increase the flexibility in the heteroepitaxy of AlGa_N-based devices.

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Figure Captions

Figure 1. (a) Schematic drawing of the layer structure employed in this work, and (b) schematic drawing of the temperature profile during a typical growth of the LT GaN buffer, HT GaN, LT AlGa_N interlayer, and HT AlGa_N test layers.

Figure 2. (stress) \times (thickness) versus time during growth of Samples A-E. Dotted lines serve to demarcate the beginning and end of the HT GaN and HT AlGa_N growths. (Automatic power control based on the intensity of reflected beams as feedback was employed during this work. One should focus on the general slopes during HT GaN and HT AlGa_N growth and ignore the finer oscillations, which are artifacts possibly due to the power dependence of the laser beam profile.)

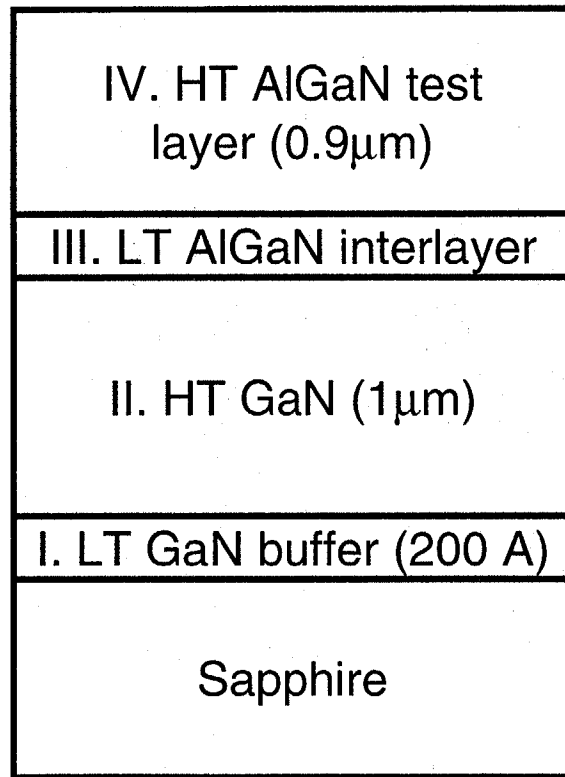
Figure 3. Nomarski surface micrograph ($\sim 200\times$) of Samples A-E. a) $x_{LT} = 0.00$ (GaN); b) $x_{LT} = 0.34$; c) $x_{LT} = 0.42$; d) $x_{LT} = 0.62$; e) $x_{LT} = 1.00$ (AlN).

Table Caption

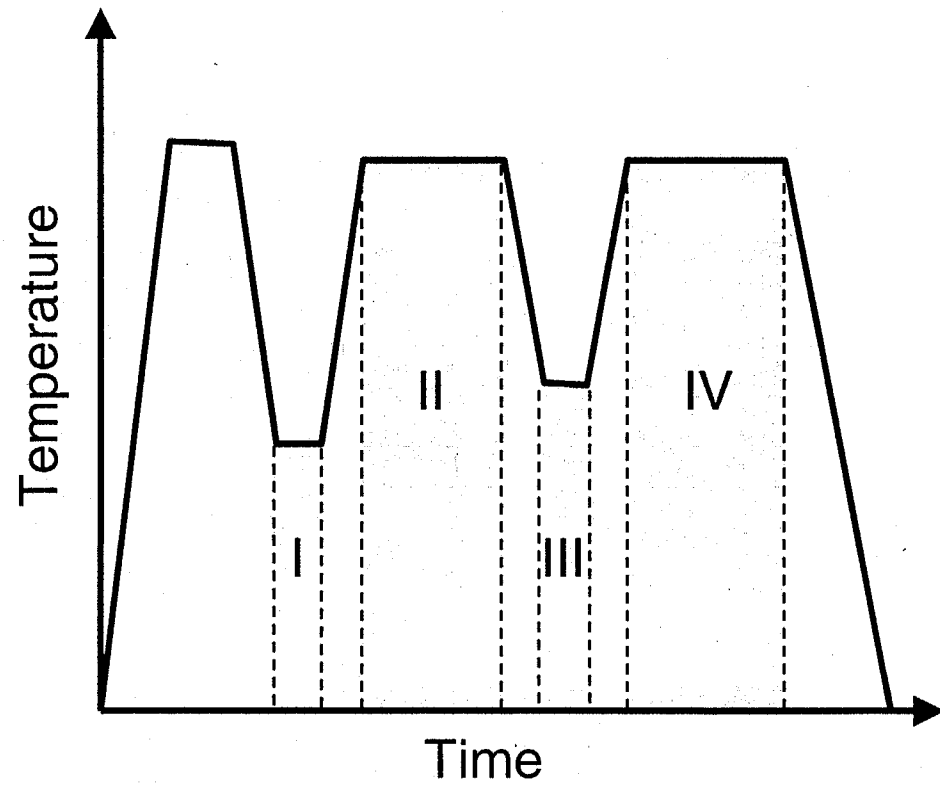
Table I. Summary of measured and calculated parameters for Samples A-E. (Fraction of strain relief in Sample A is not applicable since a LT GaN interlayer was grown on GaN.)

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(a)



(b)

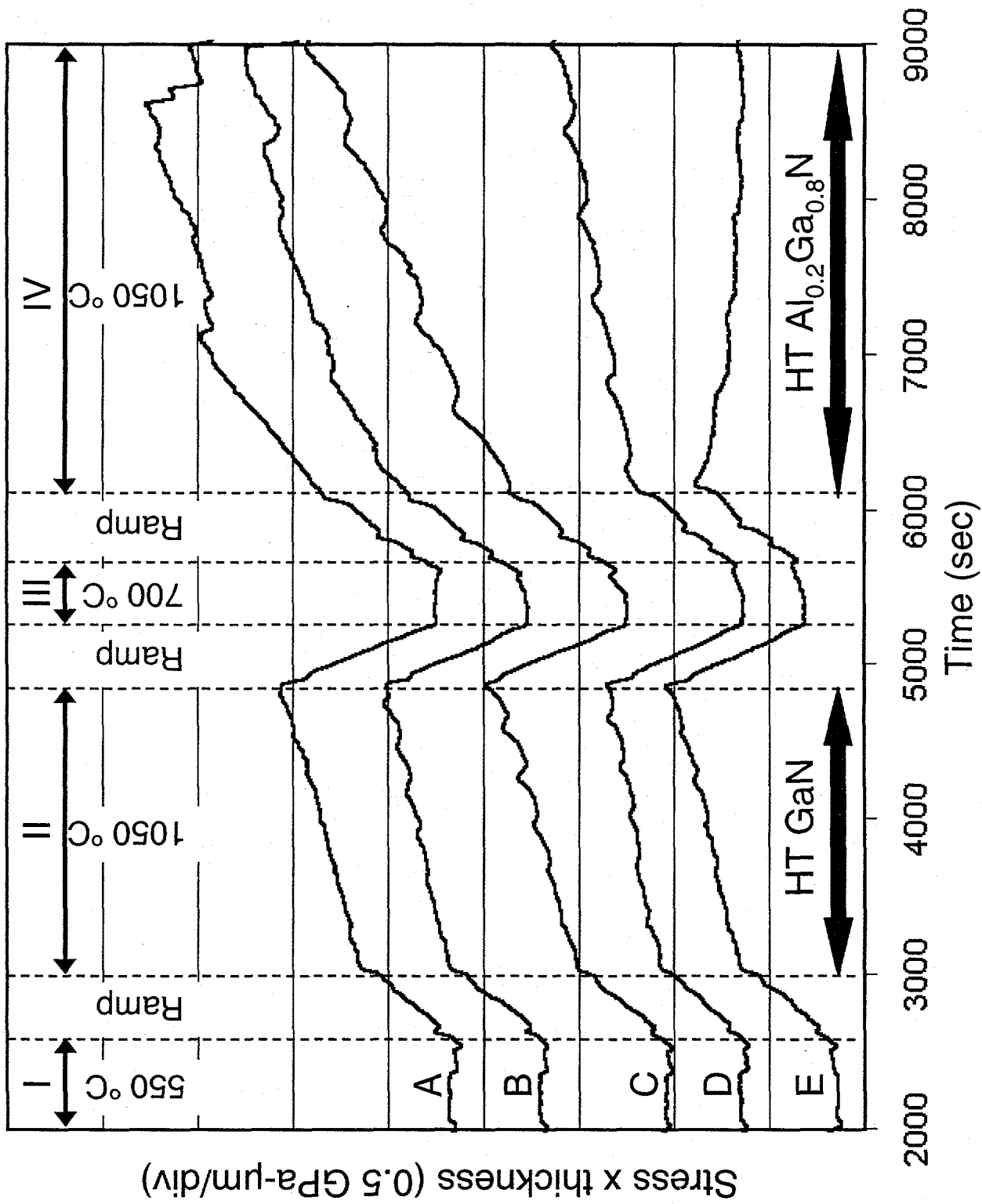


Fig 2. Han et al.

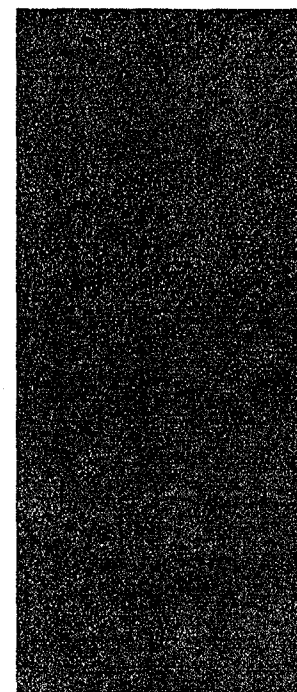
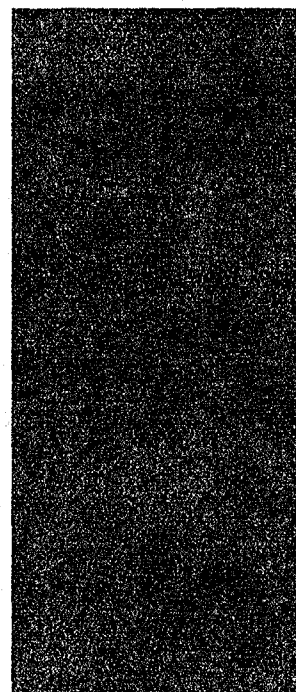
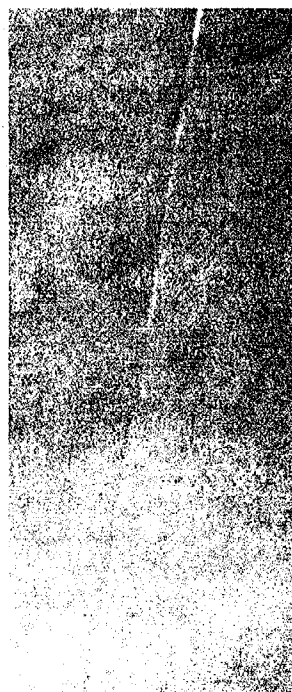
$h_g=0.10 \mu\text{m}$

$h_g=0.43 \mu\text{m}$

$h_g=0.38 \mu\text{m}$

$h_g=3.0 \mu\text{m}$

$h_g=\infty$



(a) $x_{LT}=0$

(b) $x_{LT}=0.34$

(c) $x_{LT}=0.42$

(d) $x_{LT}=0.62$

(e) $x_{LT}=1$

Fig 3. Han et al.

Sample	x_{HT}	$\sigma_{HT-AlGaN}$ (GPa)	$\epsilon_{HT-AlGaN}$ (%)	$a_{HT-AlGaN}$ $a_{LT-AlGaN}$ (Å)	σ_{GaN} (GPa)	ϵ_{GaN} (%)	a_{GaN} (Å)	x_{LT}	$a_0(x_{LT})$ (Å)	Fraction of strain relief
A	0.201	2.730	0.607	3.1930	0.448	0.100	3.1922	0.00	3.1890	NA
B	0.203	1.310	0.291	3.1828	0.393	0.087	3.1918	0.34	3.1631	0.314
C	0.201	1.400	0.311	3.1836	0.567	0.126	3.1930	0.42	3.1570	0.262
D	0.200	0.501	0.111	3.1773	0.355	0.079	3.1915	0.62	3.1418	0.286
E	0.200	-0.746	-0.166	3.1685	0.459	0.102	3.1923	1.00	3.1129	0.299