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Nondestructive Assay Tests of High-Efficiency Neutron Counter (HENC) for Waste Assay and Possible Diversion Scenario

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Edited by Jeff Skiby, Group CIC-1 Prepared by Karen Griggs, Group NIS-5

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Nondestructive Assay Tests of High-Efficiency Neutron Counter (HENC) for Waste Assay and Possible Diversion Scenario

D. R. Mayo H. O. Menlove J. M. Pecos

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by

D. R. Mayo, H. O. Menlove, and J. M. Pecos

ABSTRACT

An advanced passive neutron counter, the high-efficiency neutron counter (HENC), has been used to measure plutonium content in 200-L waste drums. The HENC was designed with the ²⁵²Cf "add-a-source" (AS) feature to improve accuracy over a wide range of waste matrix materials. The current implementation allows for passive neutron coincidence counting, AS analysis, and multiplicity analysis. Passive neutron assay of "typical" waste containers is intrinsically more accurate than active neutron techniques because of the penetrability of the spontaneous fission neutrons originating from within the waste matrix. In addition, the HENC is designed as a slightly undermoderated detector to be less sensitive to low loading of hydrogen-bearing matrices. The following paper considers the applicability of three different nondestructive assay methods for analysis of waste drums and the flagging of possible diversions in waste drums. The ²⁵²Cf AS method, multiplicity counting, and a bounded-parameter multiplicity analysis are presented with areas of applicability.

I. INTRODUCTION

The measurement of plutonium content in waste containers is required prior to longterm storage or shipment. The high-efficiency neutron counter (HENC) was developed as part of a cooperative research and development agreement (CRADA) between the Department of Energy (DOE) and Canberra Industries to measure the plutonium content in 200-L waste drums.¹ The initial characterization of the detector and performance evaluation is described elsewhere.²

The HENC consists of high-density polyethylene (HDPE) with 10-cm detector banks backed by 30-cm-thick HDPE shielding. The sample cavity is designed to accommodate up to 200-L drums. The 200-L sample drums are brought into the detector via rollers onto a rotating stand. The detector has doors on opposite sides for ease of loading and to allow for placement in an industrial setting. A full description of the system can be found in the HENC Hardware Manual.³

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The HENC uses totals, coincidence, multiplicity, and the ²⁵²Cf "add-a-source" (AS) method to determine plutonium mass in a sample. The passive system records the multiplicities measured from the waste drums and calculates the single (S), double (D), and triple (T) counting rates. The passive neutron counting relies on the spontaneous fission neutrons from the sample. The signal level, however, of spontaneous fission neutrons is low, limiting the sensitivity of the system. The multiplicity method uses the S, D, and T rates to solve the point model equations and determine the plutonium mass. The AS feature uses an external ²⁵²Cf source to interrogate the sample. The dependence of neutron coincidence counting on the waste matrix is minimized by the use of the AS correction. For all of the methods considered, it is necessary to know the plutonium isotopics to convert measured ²⁴⁰Pu effective mass to total plutonium.

Three methods of nondestructive assay (NDA) for 200-L waste drums are considered with regards to applications in waste assay, performance demonstration program (PDP) trials, and diversion scenarios. A brief discussion of the AS, multiplicity, and bounded-parameter multiplicity analysis NDA methods are given, followed by a presentation of the results.

A. Add-A-Source

The AS method⁴ measures the matrix perturbation to the counting rate from a small 252 Cf source ($\approx 3 \times 10^4$ n/s) on the outside of the sample. This information is used to correct for the matrix perturbation inside the sample.

The HENC, as shown in Fig. 1, positions the ²⁵²Cf source at three locations along the side of the drum. The drum rotates to give an averaged effect. The sample matrix has two primary effects on the neutrons. The neutrons are reduced in energy by scattering reactions and low-energy neutrons are captured by the materials. As shown in Fig. 2, hydrogen-bearing materials tend to decrease the observed count rate, while metal-bearing materials tend to enhance the count rate.



Fig. 1. HENC conceptual diagram with the ²⁵²CF AS feature.



Fig. 2. Effects of matrix loading on measured double count rate from FZC-158 for the HENC.

Matrices of glass, concrete/rubble, sand, iron, various densities of polyethylene, paper, and wood were used to calculate the AS perturbation. The AS perturbation is the measured change in the D rate of ²⁵²Cf source exterior to the drum due to the presence of the waste drum and its contents. The AS perturbation corresponds to a correction factor resulting in a "corrected doubles rate." The corrected doubles rate, D_c , is related to the ²⁴⁰Pu effective mass. The mass determination relies on calibration of the detector in the passive mode. One possible weakness of the AS method is in the assumption that the plutonium is uniformly distributed in the drum.

B. Multiplicity Analysis

The analysis of the drums using multiplicity counting assumes the point model⁵ and sets the neutron self multiplication to 1 (M = 1). The point model relates the observed S, D, and T rates to the ²⁴⁰Pu effective mass (m_{Pu240e}), the detector efficiency (ϵ), the neutron self multiplication (M), and the (α , *n*) to spontaneous fission ratio (α).

Typically, one measures an average efficiency with a calibrated standard and calculates the other three parameters. For the cases of waste, it is assumed that the self-multiplication is small (M \approx 1). This allows for solution of ε , α , and m_{Pu240e}, where the matrix effects are taken into account by the change in detector efficiency.

The multiplicity method has an advantage in identifying diversions as it relies on the S, D, and T rates (which go as $\varepsilon:\varepsilon^2:\varepsilon^3$) to calculate the m_{Pu240e} , whereas the AS method relies only on the corrected D rate of the sample. Changes in the detector efficiency are due to matrix effects and are shown by changes in the ratios of the S:D:T counting rates, thus resulting in an indicator of shielding. The efficiency measured by multiplicity counting has the advantage in that it directly measures the detection efficiency of each neutron escaping the drum. A sample having localized shielding and/or nonuniform plutonium distributions are not a problem for the multiplicity method. The primary limitation is that low plutonium mass samples, which is typical of waste, will result in poor counting statistics unless long counting times are used.

C. Bounded Parameter Multiplicity Analysis

The limitations of the AS and multiplicity methods for detection of diversion lead to the consideration of an alternative method, the bounded-parameter multiplicity analysis (BPMA). The BPMA method does not use the point model equations with M = 1, but incorporates the calibration curves that relate m_{Pu240e} to D and T coincidence rates. The calibration curves for S, D, and T to m_{Pu240e} for the HENC were measured using plutonium standards including mixed oxide and PuO₂ powder samples. The calibration² gives a fit of the data for D and T rates to m_{Pu240e} . The neutron multiplication was negligible for the calibration standards used (i.e., M \approx 1). The BPMA generates a D_c using the D and T counting rates, which is related to m_{Pu240e} ($m_{Pu240e} \propto D_c$). The correction factor is of the form

$$CF = 1 + f\left(\frac{D}{T}\right),$$

which results in a lower error/uncertainty for the mass than the multiplicity equation results for mass.

D. Backgrounds

For all of the analysis methods used with passive neutron counting, it is necessary to measure background levels with similarly loaded 55-gal. drums because the contents of the drum affect the background level to be subtracted. A metal-containing drum will increase background levels whereas hydrocarbons decrease the background, both affecting the D rate and inducing a bias. For the HENC, Table I contains the data measured for various matrices and Figs. 3 and 4 compare the results of D and T background effects respectively.

Table I. Cosmic-Ray Backgrounds as Measured in HENC as a Function								
of 200-L Drum Matrix Loading ²								
Matrix	Matrix Mass	Singles	Doubles	Triples				
	(kg)	(counts/s)	(counts/s)	(counts/s)				
Empty Drum	0.0	23.67	2.61	0.551				
Paper	53.9	22.58	2.23	0.437				
Polyethylene Beads with Vermiculite	66.0	22.25	2.24	0.460				
Polyethylene Tubes	32.5	22.99	2.30	0.463				
Polyethylene Shavings	7.2	23.34	2.40	0.492				
Wood	51.0	23.44	2.40	0.521				
Boron Glass	173.0	24.54	2.94	0.668				
Concrete Rubble	220.0	26.05	2.97	0.678				
Graphite Block	126.6	26.96	3.13	0.743				
Sand	290.0	27.53	3.68	1.003				
Iron	196.0	30.38	5.70	2.191				



Fig. 3. Doubles background rate for the HENC as a function of matrix type and mass in the 200-L drum.



Fig. 4. Triples background rates for the HENC as a function of matrix type and mass in the 200-L drums.

E. Applications of Methods

The three analysis methods are compared for applications to waste assay, including the PDP cycles, and specifically the possible diversion of material from any location where appreciable amounts of special nuclear material, especially plutonium, are accessed by workers. The HENC with AS has the capability of flagging local shielding by comparing the resulting AS measured D ratio to an empty drum at the three points along the outside of the drum. Large fluctuations in the ratios at the three points would indicate a possibly stratified matrix and possibly local shielding. The multiplicity analysis is best suited for determining local shielding of a single source, but suffers from poor counting statistics with low mass samples. The BPMA produces results with smaller uncertainty than multiplicity for the same measurement times, when using the D and T rates to correct for matrix effects.

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II. ANALYSIS METHODS

Data were recorded for ²⁵²Cf and plutonium samples for various matrix loadings to arrive at a volume and/or radial averaged result. S, D, and T rates from the multiplicity register, along with the D rates of the AS at three vertical positions, were recorded for each setup. Additional data were recorded with plutonium standards which consisted of the plutonium oxide powders mixed in diatomaceous earth. The data were taken using the CANBERRA WM3100 Series/HENC Passive Neutron Coincidence Drum Counter with the Neutron Coincidence Counting software, Windows NCC. The HENC Detector Parameters are listed in Table II. The data were then analyzed using the three methods described in the previous sections and detailed below. The ²⁵²Cf data, in addition to measuring the detector parameters, were used as a reference calibration source by being placed in the center of the detector.

Table II. HENC Detector	Parameters ²
Detector Parameter	Value
High Voltage	1720 V
Pre-Delay	3.0 µs
Neutron Die-Away	50 µs
Coincidence Gate Length	128 µs
Deadtime (a)	0.50 μs
(b x 10 ⁻⁶)	0.161 µs
Multiplicity Deadtime	0.147 μs
Neutron Efficiency: ²⁵² Cf	30%
²⁴⁰ Pu	31%

The passive neutron coincidence calibration data for the HENC are listed in Table III, and the resulting calibration curves relating D and T to m_{Pn240e} are shown in Fig. 5.

Table III. HENC	Passive	Neutron	Counting	Calibratio	n Data ²				
	m _{Pu240e}	Т	S	D	Т	Alpha			
Sample	(g)	(s)	(counts/s)	(counts/s)	(counts/s)	(α)			
A1-066, 081, 078	0.161	60 x 30	98.9	8.86	0.971	1.013			
A1-066-078, 081, 089, 119	0.243	60 x 30	143.7	13.12	1.42	1.013			
FZC-158	0.705	60 x 30	246.8	37.87	4.3	0.142			
A1-066, 078,081, 089, 119, FZC-158	0.948	60 x 30	390.5	50.98	5.41	0.365			
²⁵² Cf, CR-6 Centered (96/11/01)	N.A.	20 x 30	2120.0	646.40	118.1	NA			
PDP1-3.0 ^a	0.1833	40 x 30	158.5	9.80	1.25	0.787 ^b			
$PDP1-3.0 + 0.5^{a}$	0.2138	59 x 30	186.2	11.94	1.88	0.787 ^b			
$PDP1-3.0 + 0.5 + 0.1^{a}$	0.2199	513 x 10	189.4	11.97	1.37	0.787 ^b			
^a The PDP Standards were counted at two-	^a The PDP Standards were counted at two-thirds radius and half height to approximate volume-averaged response.								
^b The effective α is ≈ 1.80 because of the	increase fro	om diatoma	ceous earth.						



Fig. 5. D and T counting rate calibration of HENC for plutonium samples listed in Table II.

A. Add-A-Source

The AS analysis uses a ²⁵²Cf source external to the waste drum to correct for the effects of the matrix on the neutrons generated in the drum. The HENC measures the D rate due to the ²⁵²Cf source at three positions along the exterior of the drum and calculates an average (D). A ratio of the average D rate to the D rate measured for an empty drum (D_o) is used to correct the passive doubles. The matrices used to generate the calibration are described in Table IV. The effect on the average D of the external ²⁵²Cf source are related to the effect on neutrons passing from inside the matrix to the detector. A plutonium source is placed inside the test matrix materials at known locations to generate a volume or radially averaged matrix perturbation for the neutrons. The resulting data is presented in Table V.

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Table IV. Ma	atrix Loadings of	200-L Drums Used for HENC Add-a-Source Calibration
Sample Drum Loading	Mass of Matrix (kg)	Description
Empty Drum	0.0	Empty metal drum
Polyethylene Drum Liner	9.1	5-mm-thick polyethylene liner on inside wall of drum
Boron Glass	173.0	Raschig rings
Concrete/Rubble	220.0	High-density concrete block mixed with 10 kg paper
Sand	290.0	Dry sand
Iron Metal	196.0	Iron pieces (≈ 170 kg) mixed with aluminum (≈ 26 kg)
Polyethylene Shavings	7.2	Low-density polyethylene shavings
Polyethylene Tubes	32.5	Chopped polyethylene tubes
Paper I	38.4	Low-density paper
Paper II	53.9	Medium-density paper
Paper III	71.0	High-density paper
Wood	51.0	Wood pieces
Polyethylene Beads/ Vermiculite	66.0	Polyethylene beads (41 kg) plus vermiculite (22 kg) plus borax (\approx 31 kg)

Table	IV.	Matrix	Loadings	of	200-L	Drums	Used	for	HENC	Add-a-Source
					Calibr	ation				

	Table V. Add-a-Source Calibration Data ²								
Sample	S _{av}	D _{av}	T_{av}	S _° /S	D₀/D	T₀/T	AS D _{av}	AS D _o /D	
Empty Drum	246.8	37.87	4.11	1.000	1.000	1.000	10030	1.000	
Poly Drum Liner	248.3	37,55	3.97	0.994	1.009	1.035	10148	0.988	
Boron Glass	237.4	38.87	4.36	1.040	0.974	0.943	9778	1.026	
Concrete/Rubble	257.8	38,85	4.35	0.957	0.975	0.945	9925	1.011	
Sand	265.3	48.45	6.55	0.930	0.782	0.627	10985	0.913	
Iron Metal	259.0	47.54	7.84	0.953	0.797	0.524	10696	0.938	
Poly Shavings	258.4	39.68	4.61	0.955	0.954	0.892	10055	0.998	
Poly Tubes	231.6	26.53	2.35	1.066	1.427	1.749	7540	1.330	
Paper I	253.8	35.09	3.36	0.972	1.079	1.223	9278	1.081	
Paper II	246.6	32.33	2.98	1.001	1.171	1.379	8500	1.180	
Paper III	228.1	25.33	2.20	1.082	1.495	1.868	7711	1.301	
Wood	247.7	31.75	2.81	0.996	1.193	1.463	8594	1.167	
Poly Beads/ Vermiculite	202.2	21.23	1.30	1.221	1.784	3.162	6942	1.445	
Half Poly/ Half Aluminum	100.7	4.58	0.13	2.451	8.269	31.615	7066	1.419	
^a The D and T va	The D and T values have been corrected for matrix induced backgrounds								

The volume-averaged plutonium D perturbation (Δy_D) was approximated with a thirdorder polynomial and is given by

$$\Delta y_D = a_0 + a_1 \delta + a_2 \delta^2 + a_3 \delta^3. \tag{1}$$

The fit coefficients are given by

 $a_0 = -0.0131,$ $a_1 = 1.378,$ $a_2 = -1.916,$ and $a_3 = 5.86;$

and the AS perturbation is represented as $\delta = (D_0/D) - 1$, where D_0 is the measured D rate of the ²⁵²Cf source with an empty drum and D is the measured D rate with the sample drum.

The resulting relationship is a single curve from AS perturbation (δ) to plutonium D (Δy_D) or mass perturbation, as shown in Fig. 6. As implemented, an additional constraint is placed upon the perturbations, for $|\delta| \le 0.02$ then $\Delta y_D \equiv 0$, to handle the fluctuations near zero correction.



Fig. 6. AS calibration function for the matrix materials listed in Table V. The horizontal axis is the perturbation of the external AS and the vertical is the volume-averaged perturbation of the 240 Pu effective D rate.

The AS corrected D rate (D₂) is given by

$$D_c = CF \times D$$
, where $CF = 1 + \Delta y_D$. (2)

The m_{Pu240e} is then calculated using a linear calibration curve of m_{Pu240e} vs D and is given by

$$D_c = am_{Pu240e},\tag{3}$$

where *a* is given in Fig. 5 as 53.8 counts/s•g-Pu240e and the results of the calibration given in Table VI. The spread in the D is reduced from 21% to 3%. The m_{Pu240e} is calculated using (Eq. 3) and is in good agreement with the mass of FZC – 158 (0.705 – g).

Tabl	Table VI. Add-a-Source Corrected D Rate ²								
Sample	D _{avg}	D _o /D	AS D _o /D	CF	D _c	m _{Pu240e}			
Empty Drum	37.87	1.000	1.000	1.000	37.870	0.704			
Poly Drum Liner	37.55	1.009	0.988	1.000	37.550	0.698			
Boron Glass	38.87	0.974	1.026	1.005	39.064	0.726			
Concrete/Rubble	38.85	0.975	1.011	1.000	38.850	0.722			
Sand	48.45	0.782	0.913	0.845	40.940	0.761			
Iron Metal	47.54	0.797	0.938	0.892	42.406	0.788			
Poly Shavings	39.68	0.954	0.998	1.000	39.680	0.738			
Poly Tubes	26.53	1.427	1.330	1.437	38.124	0.709			
Paper I	35.09	1.079	1.081	1.069	37.511	0.697			
Paper II	32.33	1.171	1.180	1.190	38.473	0.715			
Paper III	25.33	1.495	1.301	1.377	34.879	0.648			
Wood	31.75	1.193	1.167	1.172	37.211	0.692			
Poly Beads/ Vermiculite	21.23	1.784	1.445	1.750	37.153	0.691			
Mean	35.47				38.44	0.714			
Std. Dev. (counts/s)	8.03				1.87	0.035			
Std. Dev. (%)	22.65				4.87	4.867			

B. Multiplicity Analysis

The multiplicity analysis, in this case, assumes the point model equations with M = 1 (i.e., nonmultiplying sample with low levels of plutonium and/or dispersed throughout volume). Typically the detector efficiency is fixed and a solution is found for the self-multiplication (M), the (α ,n) to spontaneous fission rate (α), and the m_{Pu240e}. The multiplicity analysis relies on point model equations⁶ for the passive case and are given by

$$S = F \varepsilon v_{s1}(1+\alpha),$$

$$D = F \varepsilon^{2} \frac{f_{d}}{2} M^{2} \bigg[v_{s2} + \bigg(\frac{M-1}{v_{i1}-1} \bigg) v_{s1}(1+\alpha) v_{i2} \bigg], \text{ and}$$

$$T = F \varepsilon^{3} \frac{f_{t}}{6} M^{3} \bigg\{ v_{s3} + \bigg(\frac{M-1}{v_{i1}-1} \bigg) [3v_{s2}v_{i2} + v_{s1}(1+\alpha)v_{i3}] + 3 \bigg(\frac{M-1}{v_{i1}-1} \bigg) v_{s1}(1+\alpha) v_{i2}^{2} \bigg\};$$
(4)

where

S	=	the singles rate
D	=	the doubles rate,
Т	=	the triples rate,
m _{Pu240e}	=	the effective ²⁴⁰ Pu mass,
F_{\perp}	=	the plutonium fission rate, $m_{Pu240e} \times 479$,
ε	=	the neutron detection efficiency,
f_d	=	the fraction of neutron events in the doubles gate, $f_d = e^{-PD/\tau_d} (1 - e^{-G/\tau_d})$,
f_t	=	the fraction of neutron events in the triples gate $f_i \approx (f_d)^2$,
τ	=	the detector die-away time,
PD	=	the multiplicity circuit predelay,
G	=	the multiplicity electronics correlation event gate width,
М	=	the sample leakage multiplication,
α	=	the ratio of (α, n) neutrons to fission neutrons,
V_{si}	=	the i th moment of the spontaneous fission multiplicity distribution, and
V _{ii}	-	the i th moment of the induced fission multiplicity distribution.

The S, D, and T rates are the quantities measured by the shift register and multiplicity circuits. The fraction of doubles (f_d) , and triples (f_p) in the gate are typically measured using a ²⁵²Cf source as $\alpha = 0$ and M = 1 for ²⁵²Cf. The detector efficiency can be easily measured using the ²⁵²Cf source, but the energy spectra of the emitted neutrons are different from the plutonium resulting in a slightly different efficiency. Measurement of the detector efficiency with a plutonium sample requires that the sample be very well characterized. The moments of the spontaneous and induced fission distributions are given in Table VII. The detector dieaway (τ), predelay (PD), and coincidence gate (G) are a function of the detector and the resulting parameters are given in Table II.

T	Table VII. Measured Prompt Fission Moments ⁷								
Probability	²³⁸ Pu	²³⁹ Pu	²⁴⁰ Pu	²⁴¹ Pu	²⁴² Pu	²⁴¹ Am	²⁵² Cf		
Distribution									
ν _{s1}	2.21	2.16	2.156	2.25	2.145	3.22	3.757		
ν _{s2}	3.957		3.825		3.794		11.962		
ν _{s3}	5.596		5.336		5.317		31.812		
ν_{i1}	2.9	2.879	2.8	2.931	2.81	3.09	4.06		
ν _{i2}		6.773		6.9946					
ν _{i3}		12.63		13.2396					

The standards used in the calibration were chosen such that M = 1, reducing the equations to

$$S = F \varepsilon v_{s1}(1+\alpha),$$

$$D = F \varepsilon^2 \frac{f_d}{2} v_{s2}, \text{ and}$$

$$T = F \varepsilon^3 \frac{f_t}{6} v_{s3},$$

(5)

which can be solved for the ε , α , and m_{Pu240e} . The fraction of coincident events in the gate parameters, f_d and f_t , are calculated from the measurement of known plutonium samples.

The fractions were calculated using the FZC-158 plutonium sample in an empty drum and the volume-averaged point. The volume-averaged point occurs at the half height and two-thirds radius of the 200-L drum. The isotopics of FZC-158 are given in Table VIII. Given α and m_{Pu240e} , we can calculate ϵ , f_d , and f_t which are given by

$$\varepsilon = S/Fv_{s1}(1+\alpha),$$

$$f_d = 2\frac{D}{S}\frac{v_{s1}(1+\alpha)}{v_{s2}}, \text{ and}$$

$$f_t = 3\frac{T}{D}\frac{v_{s2}}{v_{s3}}\frac{f_d}{\varepsilon}.$$
(6)

Table VI	II. FZC-1	58 Isotopics								
and α										
FZC-158	12/16/82	03/07/97								
^{Pu} 238	0.016	0.008								
Pu239	0.955	0.969								
^{Pu} 240	93.76	94.381								
^{Pu} 241	0.699	0.007								
^{Pu} 242	4.56	4.637								
²⁴¹ Am	0.001	0.624								
²⁴⁰ Pu-eff	0.705	0.698								
α	0.130	0.145								

The resulting gate fractions are $f_d = 0.666$ and $f_t = 0.515$ (Note: this is not the usually assumed relationship of $f_t = f_d^2$). The efficiency was $\varepsilon = 0.298$ with the polyethylene moderator of the interior of the sample cavity, and it changed to $\varepsilon = 0.304$ when the polyethylene was moved behind the steel of the sample cavity. This differs slightly from the efficiency measured using the ²⁵²Cf source because their fission neutron energy spectra differ. Upon moving the detector to the installation site, the resulting efficiencies were 30% for ²⁵²Cf and 31% for plutonium.

The multiplicity approach given by Eq. (3) results in three equations and four unknowns ({ $m_{Pu240e}, \varepsilon, \alpha, M$ }). Typically, the average efficiency of the detector is measured and becomes a fixed parameter for the point model to be soluble for ({ m_{Pu240e}, α, M }). Since it is now assumed that the self-multiplication is given by M = 1, the efficiency becomes a free parameter and the solution set results in ({ $m_{Pu240e}, \varepsilon, \alpha$ }). The detection efficiency changes with respect to location and distribution of the sources within the drum. The efficiency is additionally altered by the addition of matrix materials which decrease the signal (hydrocarbons) or increase the signal (metals) as shown if Figs. 2, 3, and 4. A change in the detector's efficiency gives an indication of the matrix material within the drum.

The point model equations reduce to solve for ε , α , and m_{Pu240e} are given by

$$\varepsilon = 3 \frac{T}{D} \frac{f_d}{f_t} \frac{v_{s2}}{v_{s3}},$$

$$\alpha = \frac{3}{2} \frac{ST}{D^2} \frac{f_d^2}{f_t} \frac{v_{s2}^2}{v_{s1} v_{s3}} - 1, \text{ and}$$

$$m_{Pu240} = \frac{F}{479} = \left(\frac{1}{479}\right) \frac{2}{9} \frac{D^3}{T^2} \frac{f_t^2}{f_d^3} \frac{v_{s3}^2}{v_{s2}^3}.$$
(7)

The uncertainty in the mass m_{Pu240} is given by

$$\sigma_{m}/m_{Pu240e} = \sqrt{\left(3\frac{\sigma_{D}}{D}\right)^{2} + \left(2\frac{\sigma_{T}}{T}\right)^{2} + 2\left(\frac{3}{D}\right)\left(\frac{-2}{T}\right)\operatorname{cov}(D,T).}$$
(8)

In order to estimate the uncertainty without inverting the error matrices, the covariance terms are dropped to compare dependencies. Typically, the cross terms (covariance) has a 5% - 40% effect.⁸ Alternatively, if it is assumed that the correlation between the observable parameters are small then the uncertainty reduces to

$$\sigma_{m/m_{Pu240e}} \approx \sqrt{\left(3\frac{\sigma_{D}}{D}\right)^{2} + \left(2\frac{\sigma_{T}}{T}\right)^{2}} \propto \frac{2\sigma_{T}}{T}$$
 (9)

If the cov(D,T) > 0, then the error is reduced in the mass. As the full error propagation has not been implemented in software for this case, multiple runs for a given sample are used to estimate the uncertainty of the measurement.

C. Bounded-Parameter Multiplicity Analysis

BPMA uses the D and T rates to generate a passive D correction factor. A D correction factor is generated using the relationship between D and m_{Pu240e} and T and m_{Pu240e} as given in Table III and Fig. 5.

The passive D and T rates can be corrected for matrix effects and are related to mass by a "correction factor" and linear correspondence between a corrected rate to m_{Pu240e} . The mass and rates are represented by

$$D_c = CF)_D D = a_1 m_{Pu240e}$$
 and
 $T_c = CF)_T T = a_2 m_{Pu240e}$, (10)

where $CF)_D$ and $CF)_T$ are the matrix correction factors for D and T, and a_1 and a_2 are the calibration slopes for the D and T related to mass of ²⁴⁰Pu. The above reduces to

$$\frac{CF)_D}{CF)_T} \frac{D}{T} = \frac{a_1}{a_2} . \tag{11}$$

Measurements made to calibrate the HENC² for typical matrix loadings of drums are given in Tables III and IX. One can generate the perturbation in the D (Δy_D) and T (Δy_T) as a function of AS perturbation (δ), as in Fig. 9.

The form of the D and T correction factors are assumed to be

$$(12)$$

 (12)
 $(CF)_T = 1 + \Delta y_T$.

Table IX. Perturbation of HENC Measured D and T								
Sample	D	Т	δ	D/T	Δy_d	Δy_1		
Poly Tubes	26.53	2.35	0.330	11.29	0.427	0.749		
Paper I	35.09	3.36	0.081	10.44	0.079	0.223		
Paper II	32.33	2.98	0.180	10.85	0.171	0.379		
Paper III	25.33	2.20	0.301	11.51	0.495	0.868		
Wood	31.75	2.81	0.167	11.30	0.193	0.463		



Fig. 9. Perturbation of $D(\Delta y_D)$ and $T(\Delta y_T)$ counting rate as a function of AS perturbation for the HENC.

A much steeper correction to the T rate is observed, which is a result of the T being a function of ε^3 and D of ε^2 . The relationship between the D and T perturbation is shown in Fig. 10. Although a higher polynomial may be used, the quality of the data does not warrant more than a linear relationship.

$$\Delta y_T = A \cdot \Delta y_D + B, \tag{13}$$

where A and B are the constants from the fit. The D perturbation (Δy_D) can be expressed as a function of the D and T rates.



Fig. 10. Perturbation of T rate to D rate for HENC.

Using Eqs. (11-13), the D perturbation as a function of D and T is given by

$$\frac{CF)_D}{CF)_T} = \frac{a_1}{a_2} \frac{T}{D},$$

$$\frac{1 + \Delta y_D}{1 + A\Delta y_D + B} = \frac{a_1}{a_2} \frac{T}{D}, \text{ and}$$

$$\Delta y_D = \frac{\frac{a_1}{a_2} (1 + B)T - D}{D - \frac{a_1}{a_2} A T}.$$
(14)

For the HENC,² $a_1 = 53.8$ and $a_2 = 5.83$. If an empty drum is considered, then by definition $\Delta y_D = \Delta y_T = 0$, which implies that B = 0. If, however, Eq. (14) is set to zero as should be the case for an empty drum ($\Delta y_D = 0$), then we have

$$B = \frac{D}{T} \frac{a_2}{a_1} - 1 = \left(\frac{37.87}{4.11}\right) \left(\frac{5.83}{58.3}\right) - 1 \approx -0.0786 \pm 0.0920 , \qquad (15)$$

which, within statistics is 0. As shown in Fig. 10, the resulting fit for B = 0 is $\Delta y_D = 1.8414 \ \Delta y_D$.

The BPMA D correction factor CF)_D can be rewritten, resulting in the m_{Pu240e} being given by

$$m_{P_{u}240e} = \frac{D_{C}}{a_{1}} = \frac{(1+\Delta y_{D})D}{a_{1}} = \frac{(1-A)DT}{(a_{2}D-a_{1}AT)}.$$
 (16)

The uncertainty in the mass is given approximately by

$$\sigma_{m} / m_{Pu240e} = \left[\left(\frac{\partial m_{Pu240e}}{\partial D} \frac{\sigma_{D}}{m_{Pu240e}} \right)^{2} + \left(\frac{\partial m_{Pu240e}}{\partial T} \frac{\sigma_{T}}{m_{Pu240e}} \right)^{2} + 2 \left(\frac{\partial m_{Pu240e}}{\partial D} \frac{\sigma_{D}}{m_{Pu240e}} \right) \left(\frac{\partial m_{Pu240e}}{\partial T} \frac{\sigma_{T}}{m_{Pu240e}} \right) \cos(D,T) \right]^{\frac{1}{2}}.$$
(17)

Once again, the cov(D,T) is dropped for comparison, resulting in

$$\sigma_{m} / m_{Pu240e} \approx \sqrt{\left[\left(\frac{-a_{1}AT}{a_{2}D - a_{1}AT}\right)\frac{\sigma_{D}}{D}\right]^{2} + \left[\left(\frac{a_{2}D}{a_{2}D - a_{1}AT}\right)\frac{\sigma_{T}}{T}\right]^{2}},$$

$$= \frac{1}{|a_{2}D - a_{1}AT|} \sqrt{\left[\left(a_{1}AT\right)\frac{\sigma_{D}}{D}\right]^{2} + \left[\left(a_{2}D\right)\frac{\sigma_{T}}{T}\right]^{2}}.$$
(18)

The uncertainty in the mass calculated by the BPMA method is less than that of the multiplicity analysis as the D perturbation is a small fraction of the correction factor.

D. Comparison of Multiplicity and BPMA Uncertainties

Consider two samples with masses of 69 g and 95 g of plutonium. The 69-g and 95-g cases have a D = 330.6 and T = 34.43, and D = 206.3 T = 20.59, respectively. Then using Eqs. (9) and (18) as shown in Fig. 11 for a fixed uncertainty of 1% in the D rate (i.e., $\sigma_{D/D} = 0.01$), the multiplicity error scales as $\frac{2\sigma_T}{T}$. The BPMA results in a lower uncertainty in the mass and scales as $\frac{3}{2}\frac{\sigma_T}{T}$.



Fig. 11. Comparison of uncertainties in the mass (σ_m/m) for a 69-g and 95-g sample as a function of the uncertainty in the T counting rate (σ_{r}/T). The multiplicity result matches 2 σ_{r}/T .

III. RESULTS

The three analysis methods (AS, multiplicity, BPMA) are examined using the calibration data set. The results of the analyses are given in Table X and displayed in Fig. 12. The AS method does best over the entire range as expected. The multiplicity analysis does poorly for the iron metal and sand cases and overcorrects for the lighter loadings of paper and wood. The attenuation of the matrix shows up in the change in detector efficiency as well as in the mass. The BPMA seems to fluctuate from the AS result in the same direction as the multiplicity analysis, but the fluctuations are damped due to the lesser dependence on the T rate.

Table X. Assay Results for Calibration										
Matrix	Mass (kg)	Doubles (counts/s)	I	Add-a-Son	ırce	Multiplicity		BPMA		
			δ	D _c	m _{Pu240e}	eff	m _{Pu240e}	Δy_D	D _c	m _{Pu240e}
Empty Drum	reference	37.87	0.000	37.389	0.695	0.299	0.700	-0.002	37.802	0.703
Polyethylene Drum Liner	9.1 kg	37.55	-0.012	36.575	0.680	0.291	0.731	0.031	38.698	0.719
Boron Glass	173 kg	38.87	0.026	39.421	0.733	0.309	0.673	-0.038	37.386	0.695
Concrete/Rubble	220 kg	38.85	0.013	38.894	0.723	0.308	0.675	-0.036	37.453	0.696
Sand	290 kg	48.45	-0.087	42.325	0.787	0.373	0.577	-0.191	39.201	0.729
Iron Metal	196 kg	47.54	-0.060	43.446	0.808	0.454	0.380	-0.290	33.777	0.628
Polyethylene Shavings	7.2 kg	39.68	-0.003	39.044	0.726	0.320	0.639	-0.074	36.729	0.683
Polyethylene Tubes	32.5 kg	26.53	0.330	37.864	0.704	0.244	0.734	0.359	36.041	0.670
Paper I	38.4kg	35.09	0.081	37.427	0.696	0.264	0.833	0.186	41.615	0.774
Paper II	53.9 kg	32.33	0.181	37.775	0.702	0.254	0.830	0.267	40.943	0.761
Paper III	71 kg	25.33	0.300	34.421	0.640	0.239	0.731	0.418	35.914	0.668
Wood	51 kg	31.75	0.167	36.589	0.680	0.244	0.882	0.364	43.297	0.805
Polyethylene Beads/ Vermiculite	66 kg	21.23	0.448	38.577	0.717	0.169	1.231	10.552	245.277	4.559
Half Polyethylene and Half Aluminum		4.58	0.419	7.791	0.145	0.080	1.182	-1.445	-2.039	0.038



Fig. 12. Comparison of results from AS, multiplicity analysis, and BPMA.

A. PDP Cycle 2

The HENC with the Los Alamos NCC software was tested on two PDP drums in the PDP cycle 2.⁹ The passive calibration, AS calibration, and deadtime parameters used for the cycle are given in Figs. 5 and 6, and Table II, respectively. Each of the PDP drums were measured in passive mode for 10 min followed by a 3-min run of the ²⁵²Cf AS for matrix corrections. Six sets of data were taken for each barrel. The data were taken in 10-s time intervals to minimize the effects of cosmic-ray spallation events² in the background. The higher multiplicities are typical of cosmic-ray events and are rejected by the quality control tests in the software. Table XI lists the data for both of the drums measured, PDP-002 and PDP-003, as well as the AS analysis performed by the code.

Table XI. HENC PDP Cycle 2 Data with AS Results											
	Singles	Doubles	Triples	Add-a-Source							
PDP-002	(counts/s)	(counts/s)	(counts/s)	CF	m_{Pu240e}	m _{Pu}					
1	46.33	2.511	0.448	1.409	0.066	1.083					
2	47.14	2.315	0.267	1.397	0.060	0.990					
3	45.95	1.961	0.110	1.340	0.049	0.804					
4	45.29	2.015	0.331	1.407	0.053	0.868					
5	45.61	2.117	0.183	1.413	0.056	0.916					
6	45.25	1.869	0.091	1.358	0.047	0.777					
Mean	45.93	2.131	0.238	1.387	0.055	0.906					
Std. Dev.	0.72	0.241	0.138	0.031	0.007	0.116					
Std. Dev. [%]	1.57	11.289	57.716	2.212	12.765	12.765					
	Singles	Doubles	Triples	Add-a-Source							
PDP-003	(counts/s)	(counts/s)	(counts/s)	CF	m _{Pu240e}	m _{Pu}					
1	446.6	23.87	1.984	1.232	0.547	9.002					
2	445.6	24.55	2.386	1.226	0.559	9.214					
3	446.2	24.01	2.010	1.224	0.546	8.996					
4	445.7	23.78	2.301	1.199	0.530	8.728					
5	445.7	24.30	2.676	1.236	0.558	9.194					
6	444.7	23.67	2.018	1.223	0.538	8.862					
Mean	445.8	24.03	2.229	1.223	0.546	8.999					
Std. Dev.	0.64	0.34	0.276	0.013	0.011	0.188					
Std. Dev. [%]	0.14	1.39	12.402	1.057	2.087	2.087					

The isotopic composition of the samples were assumed to be ~6% 240 Pu, and the passive calibration constant of 53.8 counts/s•g - Pu240e was used for all measurements. The HENC measurement resulted in

 0.906 ± 0.047 g plutonium for PDP-002

and

 9.00 ± 0.077 g plutonium for PDP-003,

where the uncertainties are the one standard deviation values based solely on counting statistics. The uncertainty in the matrix correction was $\sim 5\%$.⁹

An alternative way to calculate the uncertainty, which would include systematic as well as statistical uncertainties, is to calculate the mean and standard deviation of the six measurements. In this case, the HENC reports

 0.906 ± 0.116 g plutonium for PDP-002

and

9.00 ± 0.19 g plutonium for PDP-003.

This method allows for the intercomparison with the two other methods. Table XII contains the multiplicity and BPMA results. The results of AS, multiplicity analysis, and BPMA are shown in Fig. 13.

	Table XII.	Multiplicit	y and B	PMA Res	ults for	PDP Cycle	2
	Multiplicity			BPMA			
PDP-002	eff	m _{Pu240c}	m _{Pu}	$\Delta y_{\rm D}$	CF	m _{Pu240c}	m _{Pu}
1	0.492	0.0171631	0.283	-0.318	0.682	0.032	0.524
2	0.318	0.0378653	0.624	-0.067	0.933	0.040	0.661
3	0.155	0.1355995	2.233	-10.303	-9.303	-0.339	-5.585
4	0.453	0.0162472	0.268	-0.288	0.712	0.027	0.439
5	0.238	0.0616412	1.015	0.431	1.431	0.056	0.928
6	0.134	0.1715362	2.825	-3.190	-2.190	-0.076	-1.253
Mean	0.298	0.073	1.208	-2.289	-1.289	-0.043	-0.714
Std. Dev.	0.150	0.065	1.076			0.152	2.509
Std. Dev. [%]	50.372	89.054	89.054			-351.208	-351.208
	Multiplicity			BPMA			
PDP-003	eff	m _{Pu240e}	m _{Pu}	$\Delta y_{\rm D}$	CF	m _{Pu240e}	m _{Pu}
1	0.229	0.7517693	12.381	0.565	1.565	0.694	11.435
2	0.268	0.5654891	9.313	0.158	1.158	0.529	8.705
3	0.231	0.7454097	12.276	0.538	1.538	0.687	11.306
4	0.267	0.5526026	9.101	0.166	1.166	0.515	8.489
5	0.303	0.4359706	7.180	-0.019	0.981	0.443	7.300
6	0.235	0.7085379	11.669	0.475	1.475	0.649	10.689
Mean	0.255	0.627	10.320	0.314	1.314	0.586	9.654
Std. Dev.	0.029	0.128	2.110			0.1043766	1.719
Std. Dev. [%]	11.501	20.443	20.443			17.805631	17.806

Both the multiplicity and BPMA do poorly for the low-mass relatively high-attenuating ($\delta \approx 0.4$) PDP-002 case. This is directly related to large fluctuations in the T count rate, on which, the AS does not rely. Note that the results from each method are consistent for PDP-002 and PDP-003. The large fluctuations in the T rate leads to a large D/T ratio, which is outside the range of the BPMA.



Fig. 13. Comparison of HENC results using the AS, multiplicity analysis, and BPMA for PDP cycle 2.

The multiplicity and BPMA method perform better for the PDP-003 drum ($\delta \approx 0.2$) which has a stronger neutron source and a lower degree of fluctuations in the measured T rate. The uncertainty in the BPMA method is less than the multiplicity method as expected from the discussion in Section II.

B. PDP Cycle 3

Measurements for the PDP cycle 3 followed the same procedure as cycle 2 for the drums PDP-004 and PDP-003. The resulting data are given in Table XIII. The HENC with AS reports

 95.11 ± 1.19 g plutonium for PDP-004

and

 69.32 ± 0.72 g plutonium for PDP-003.

Both drums were fairly nonperturbative for neutrons ($\delta \approx 0.1$ –0.2).

T	able XIII.	HENC PD	P Cycle 3	Data with A	S Results	
	Singles	Doubles	Triples	Add-a-Source		
PDP-004	(counts/s)	(counts/s)	(counts/s)	CF	m _{Pu240e}	m _{Pu}
1	5640	333.7	35.91	1.057	6.156	95.11
2	5640	333.8	37.99	1.072	6.245	96.49
3	5640	332.0	36.56	1.077	6.240	96.42
4	5645	330.1	32.74	1.065	6.135	94.80
5	5639	328.3	29.03	1.066	6.108	94.37
6	5639	325.7	34.33	1.064	6.048	93.45
Mean	5641	330.6	34.43	6.401	6.155	95.11
Std. Dev.	2	3.2	3.21	0.007	0.077	1.19
Std. Dev. [%]	0.04	0.969	9.316	0.108	1.247	1.247
	Singles	Doubles	Triples	Add-a-Source		
PDP-003	(counts/s)	(counts/s)	(counts/s)	CF	m _{Pu240e}	m _{Pu}
1	3966	205.0	21.05	1.243	4.447	68.71
2	3965	209.0	20.95	1.238	4.516	69.77
3	3970	207.8	19.78	1.248	4.526	69.93
4	3974	205.4	21.53	1.237	4.434	68.51
5	3967	206.6	20.05	1.260	4.543	70.20
6	3971	204.0	20.20	1.251	4.454	68.82
Mean	3969	206.3	20.59	7.477	4.487	69.32
Std. Dev.	3	1.9	0.68	0.009	0.047	0.72
Std. Dev. [%]	0.09	0.90	3.311	0.116	1.043	1.043

The multiplicity and BPMA are given in Table XIV with a comparison of the three methods in Fig. 14. The BPMA and AS do not overlap within one standard deviation for the PDP-003, the multiplicity analysis easily overlaps with both. This is indicative of the refinement needed in the data, as shown in Fig. 10, for calculating the BPMA equations .

Tab	le XIV. M	ultiplicity	and BPN	MA Resu	ılts for	PDP Cyc	le 3
	Multiplicity		BPMA				
PDP-004	eff	m _{Pu240c}	m_{Pu}	Δy	CF	m _{Pu240e}	m _{Pu}
1	0.297	6.270	96.87	0.008	1.008	5.873	90.74
2	0.314	5.607	86.63	-0.054	0.946	5.512	85.17
3	0.303	5.957	92.04	-0.019	0.981	5.686	87.86
4	0.273	7.301	112.81	0.124	1.124	6.473	100.02
5	0.244	9.135	141.15	0.366	1.366	7.827	120.94
6	0.290	6.378	98.55	0.035	1.035	5.880	90.86
Mean	0.287	6.775	104.68	0.077	1.077	6.209	95.93
Std. Dev.	0.025	1.288	19.90			0.857	13.23
Std. Dev. [%]	8.747	19.012	19.012			13.796	13.796
	X 6 1.1 11 1.						
	Multiplicity		-	BPMA	<u>an</u>		
PDP-003	eff	m _{Pu240e}	m _{Pu}	Δy	CF	m _{Pu240e}	m _{Pu}
1	0.283	4.230	65.36	0.070	1.070	3.829	59.17
2	0.276	4.526	69.93	0.107	1.107	4.036	62.37
3	0.262	4.990	77.10	0.197	1.197	4.341	67.07
4	0.289	4.067	62.85	0.042	1.042	3.735	57.71
5	0.267	4.773	73.75	0.161	1.161	4.186	64.67
6	0.273	4.527	69.95	0.126	1.126	4.010	61.96
Mean	0.275	4.519	69.82	0.117	1.117	4.023	62.16
Std. Dev.	0.010	0.339	5.23			0.223	3.44
Std. Dev. [%]	3.560	7.495	7.50			5.539	5.539



Fig. 14. Comparison of HENC results using the AS, muliplicity analysis, and BPMA for PDP cycle 3.

IV. REGIONS OF APPLICABILITY

Each of the three assay techniques (AS, multiplicity analysis, BPMA) have various regions of applicability, but are achievable with one instrument, such as the HENC. The AS method relies only on measuring the D rate and the AS perturbation measurement. The AS technique covers a wide range of matrix drums, but becomes unreliable for drums with a high hydrogen content and segregated plutonium. However, "typical" waste matrix doses have a high hydrogen mass fraction. The HENC can achieve 1% statistical uncertainty in the D count rate for a 10-g plutonium sample with a 10-min passive run and 3-min ²⁵²Cf AS run. The lower detectability limit depends on the matrix and background rate. The detectability limit at Los Alamos is 20 mg ²⁴⁰Pu and, at sea level, drops to less than 10 mg ²⁴⁰Pu.

The multiplicity assay does not require the construction of a wide range of matrix drums to calibrate the system. This method requires accurate determination of the fraction in the gates of D and T. It requires reasonable statistics in the T count rate for low uncertainties in the measurement. This becomes a problem of highly attenuating matrices (i.e, high hydrogen content) or low sample masses, thus resulting in longer and possibly unreasonable run times

for assay. A 10-g sample resulted in a measurement with a multiplicity precision of 20% for a 10-min measurement.

The BMPA reduces the effects of poor counting statistics found in the multiplicity assay and results in lower uncertainties than multiplicity. It does not, however, approach the small statistical uncertainty levels of the AS technique.

None of the three techniques are very accurate for the extremely attenuating cases and small source strengths. An additional constraint on the system's ability is the construction of a drum similar to those measured for proper characterization of background effects of the matrix for samples with plutonium loadings near the detectability limit.

V. DIVERSION SCENARIO

A safeguards concern is the possibility of a diversion of plutonium from dismantlement or decommissioning processes. For example, an insider purposely shields the material within a barrel. The supposition is that waste drums are from high-plutonium areas and are modified to pass the required tests to be certified as waste to be shipped to a less-guarded facility. The signal reduction may be achieved by local shielding. The source in question is placed within a moderator to reduce the S or D rates. The AS technique, as implemented in the HENC with three positions, allows for indication of stratification layers or local shielding. A local shield would be indicated by a larger absorption of neutrons at a given level. Presence would be shown depending on source location compared to shield.

The first series of tests, reported in Table XV, investigate the case of a plutonium source shielded by polyethylene (CH₂). Single runs were recorded. The source was not at the volume-averaged point (half height and two-thirds radius), so the resulting data for the empty drum may be higher or lower depending on position. The FZC-158 plutonium source, which is ~0.70 g²⁴⁰Pu (~10 g plutonium), was placed inside of a CH₂ cylinder, and additional masses of polyethylene (poly) bead and borated poly blocks were added.

The AS method produced no correction for the lone CH_2 cylindrical shield and consistently underestimates the plutonium mass for all cases. The multiplicity analysis and BPMA analyses are consistent and failed in regions where $D/T \ge 15$. This is due to the extreme scatter in the measured T rate for multiplicity analysis. The BPMA fails in this region due to the ratio of D/T being outside the calibrated range.

Table XV. Diversion Scenario Test 1 of HENC									
Matrix Single		Doubles	Triples	Triples		Add-a-Source			
	(counts/s)	(counts/s)	(counts/s)	δ	D _c	m _{Pu240c}			
Empty drum	243.0	37.40	3.80	0.000	36.93	0.686			
CH ₂ Cylinder	242.0	29.12	2.38	0.002	28.82	0.536			
$CH_2Cyl. + 1$ -gal. beads	208.4	18.93	1.46	0.020	19.09	0.355			
$CH_2Cyl. + poly blocks$	183.1	13.72	0.71	0.045	14.17	0.263			
$CH_2 Cyl. + more poly$ blocks	136.8	7.34	0.24	0.157	8.38	0.156			
$CH_2Cyl. + 22$ -kg beads	127.1	6.01	0.09	0.161	6.89	0.128			
Matrix		Multiplicity		BPMA					
		efficiency	m_{Pu240e}	$\Delta y_{\rm D}$	D _c	m _{Pu240c}			
Empty drum		0.280	0.788	0.0858637	40.611303	0.7548569			
CH ₂ Cylinder		0.225	0.948	0.6321043	47.526878	0.8833992			
$CH_2Cyl. + 1$ -gal. beads	CH_2 Cyl. + 1-gal. beads			0.9281579	36.500029	0.6784392			
$CH_2Cyl. + poly blocks$	0.143	0.143 1.115 -4.33061		-45.696065	-0.8493692				
CH_2 Cyl. + more poly blocks		0.090	1.481	-1.5767002	-4.2329796	-0.0786799			
$CH_2 Cyl. + 22$ -kg beads		0.041	5.963	-1.1536457	-0.9234107	-0.0171638			

A second set of data were taken to systematically explore local shielding. These were performed with the highly attenuating matrices of the Poly Beads/Vermiculite/Borax (PBV) mixture, Poly Beads, and Borated Poly Blocks. The data from these runs are presented in Table XVI.

Table XVI. Diversion Scenario Test 2 Data from HENC										
Sample	Source Position	Singles	Doubles	Triples	delta	Pos 1	Pos 2	Pos 3		
		(counts/s)	(counts/s)	(counts/s)						
Empty Drum	Two-thirds R, half H	245.156	37.354	3.891	0.006	0.996	1.007	1.016		
Empty Drum	Two-thirds R, half H	245.622	37.473	4.025	0.004	1.000	1.015	0.996		
Empty Drum	Two-thirds R, half H	246.026	37.667	4.177	-0.003	0.994	0.999	0.998		
Poly Beads/Verm	Top - Center	206.737	23.033	1.955	0.443	1.438	1.469	1.423		
Poly Beads/Verm	Top - Center	208.926	22.674	1.831						
Poly Beads/Verm	Center - 5 in. down	183.937	15.420	1.089						
Poly Beads/Verm	Center - 5 in. down	182.847	15.180	1.008						
Poly Beads /Verm	1/2 R, 5 in. down	196.496	18.240	1.303						
Poly Beads /Verm	Edge, 6 in. down	224.096	27.887	2.678	0.476	1.487	1.487	1.455		
Poly Beads	Top - Center	179.660	15.807	1.103	0.489	1.238	1.404	2.020		
Poly Beads	Top - Center	184.759	1 6.97 7	1.180						
Poly Block	In block	193.270	22.488	1.973	0.030	1.025	1.022	1.043		
3 Poly Blocks	In block	118.993	8.244	0.422			-			
3 Poly Blocks	In block	118.707	8.333	0.449	0.090	1.052	1.079	1.144		
13 Poly Blocks	Center of blocks	31.257	0.535	-0.021	0.251	1.128	1.207	1.465		
13 Poly Blocks	Center of blocks	31.363	0.611	0.021	0.256	1.130	1.230	1.448		
13 Poly Blocks	Outer block	82.650	4.142	0.186	0.250	1.134	1.205	1.453		

The repeat runs were summed, and the three analysis methods were performed on the data with results given in Table XVII. The PBV data show that the location of the source affects the results. If the source is placed on top or the edge of the material, the AS overcorrects. Placing the source partially into the materials results in an underestimation of the mass. The multiplicity analysis gives results which are within 10% and detection efficiencies that change by eight percentage points. The BPMA method gives results which are consistently higher than the multiplicity analysis and fail for PBV-III. The large discrepancies can be attributed to the cases for which D/T approach 15. The AS and multiplicity analysis do reasonably well for the poly beads. The BPMA method is getting close to the D/T limit of 15.

			Table	XVII. Di	version S	Scenario Te	st Resu	lts		
Matrix	I	Doubles	Add-a-Source			Multiplicity		BPMA		
	<u>I</u> C	counts/s]	0	D _c	m _{Pu240c}	eff	m _{Pu240c}	$\Delta y_{\rm D}$	D _C	m _{Pu240c}
Empty Drun	1	37.498	0.002	37.498	0.697	0.296	0.707	0.012	37.93	0.705
PBV-I		22.854	0.443	41.045	0.763	0.228	0.726	1.213	40.23	0.748
PBV-II		15.420	0.443	27.694	0.515	0.195	0.673	1.741	42.26	0.786
PBV-III		15.180	0.443	27.263	0.507	0.183	0.749	3.017	60.97	1.133
PBV-IV		18.240	0.443	32.759	0.609	0.197	0.778	1.593	47.30	0.879
PBV-V		27.887	0.476	54.184	1.007	0.265	0.658	0.180	32.91	0.612
Poly Beads		16.392	0.489	32.899	0.612	0.192	0.735	1.949	48.35	0.899
Poly Block		22.488	0.030	22.893	0.426	0.242	0.636	0.388	31.21	0.580
3 Poly Block	s	8.289	0.090	8.910	0.166	0.145	0.655	-5.005	-33.24	-0.618
13 PB-I		0.573	0.254	0.728	0.014	-0.007	0.094	-1.229	-0.15	-0.003
13 РВ-П		4.142	0.250	5.240	0.097	0.124	0.447	-2.472	-6.10	-0.113

Note: PBV-I: Top Center, PBV-II: Center 5 in. down, PBV-III: Center 6 in. down, PBV-IV: half R, PBV-V: Edge filled half way

The PBV drum was replaced with a drum half full with polyethylene beads. The AS results in a mass which is within 15%, the multiplicity within 10%. The BPMA fails as it is past the D/T limit of 15. The AS method does show the effects of the half-full barrel in the three positions. Although the average is 1.489 for the three positions, the lowest position on the barrel (position 3) gives a ratio of ≈ 2 , which indicates a highly attenuating material. Location of the source relative to the material would greatly affect the measured value. The AS perturbation is $\delta = 1$ which is outside the calibration range of the AS technique.

The last measurements considered different configurations of the FZC-158 source and boron-loaded polyethylene blocks (¹⁰B captures thermal neutrons). The blocks are $5 \text{ cm} \times 10 \text{ cm} \times 20 \text{ cm}$ in size. One of the blocks had a hole machined out to accommodate the FZC-158 source to be centrally located. The first setup considers the source and the single block. The AS method underestimates the mass due to the relatively small size of the block. The AS and multiplicity analysis both result in lower mass, which is not only compounded by the shielding, but also the location of the source. The BPMA is within 20% of the correct value. A poly block is placed underneath and on top of the source containing the block.

For this three-block configuration, the AS underestimates the mass by a factor of 6. The multiplicity analysis results in a mass which is within 10% but at a reduced efficiency. The BPMA fails for this and the additional cases.

Thirteen blocks of poly are stacked within the sample drum to about two-thirds height. The source block is placed in the center of the drum and the center of the three layers. Although all three methods fail to give the correct mass, the three together give a useful picture as to the possible effects of shielding. An additional "flag" for anomalous drums may be generated from the D/T ratio. From Eq. (5),

$$\frac{D}{T} = 3 \left(\frac{f_{\rm D}}{f_{\rm T}} \right) \left(\frac{v_{s2}}{v_{s3}} \right) \left(\frac{1}{\varepsilon} \right) \propto \frac{1}{\varepsilon} .$$
(19)

As the amount of shielding increases, the efficiency for detection decreases, causing an increase in the D/T ratio.

VI. CONCLUSIONS

The purpose of this work was to consider three analysis techniques for applications to waste and diversion of material. The data used in the analysis were taken using the HENC with the ²⁵²Cf AS feature. All three methods of analysis bring an additional piece to the assay of an "unknown" barrel.

A. Waste

The data and results presented in the previous sections point to the AS feature as the overwhelming candidate for the measurement of drums containing a wide range of matrices. The AS technique is limited only by the fact that it assumes that the plutonium is not all concentrated at the center or perimeter of the drum. The AS relies on creating matrix drums to calibrate over the desired range of perturbations. The multiplicity analysis adds to the information on the drum. It may be applicable in determining whether a sample is undergoing self-multiplication, a nonuniform source distribution, or a high (α ,n) rate. The BPMA method has reduced errors in some regions compared to the multiplicity method. Additional comparisons are needed to map the possible applications, especially for cases with a higher (α ,n) rate.

B. Diversion Scenarios

The AS method fails to measure mass properly if all of the plutonium is concentrated in a localized neutron shield. In some cases the difference in the AS perturbation at the three positions may be used to flag possible diversions or problem drums. The position of the source and the attenuating matrix are not discernible with this technique. The source may or may not be within the matrix in question.

Multiplicity analysis can deal with these cases, but requires counting time which may be prohibitive. In addition to the mass reported by the technique, efficiency and α are given.

For some cases, the BPMA analysis resulted in an uncertainty less than that produced by multiplicity analysis for the equivalent counting time. The D/T ratio may be more applicable and easier to implement than the BPMA method for flagging a possible diversion. The ratio is proportional to the efficiency for detecting the neutrons from their point of origination. The BPMA method, however, requires further study, and additional calibration data sets would improve the relationship. All three measurements in conjunction would provide a series of software tests pointing to barrels which need additional analysis or closer scrutiny.

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