*HENC Performance Evaluation
and Plutonium Calibration*

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HENC PERFORMANCE EVALUATION AND PLUTONIUM CALIBRATION

by

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ABSTRACT

We have designed a high-efficiency neutron counter (HENC) to increase the plutonium content in 200-L waste drums. The counter uses totals neutron counting, coincidence counting, and multiplicity counting to determine the plutonium mass. The HENC was developed as part of a Cooperative Research and Development Agreement between the Department of Energy and Canberra Industries. This report presents the results of the detector modifications, the performance tests, the add-a-source calibration, and the plutonium calibration at Los Alamos National Laboratory (TA-35) in 1996.

I. INTRODUCTION

This report gives a summary of the performance tests and calibrations that were performed on the high-efficiency neutron counter (HENC) that was designed as part of the Cooperative Research and Development Agreement (CRADA) between Canberra Industries and the Department of Energy (DOE). The design and fabrication of the HENC was completed in early 1996.

The initial characterization of the HENC was performed by Canberra in April 1996 and results are documented elsewhere.¹ Small design modifications were made at the Los Alamos National Laboratory (Los Alamos), and the system was calibrated for plutonium. The add-a-source (AS) matrix correction calibration was made to reflect the final moderator configuration.

II. MECHANICAL DESIGN

The HENC was fabricated out of high-density polyethylene (HDPE) with 10-cm-thick detector banks backed by 30-cm-thick HDPE shielding. The sample cavity size is 82 cm wide by 109 cm high by 77 cm deep. There are 113 ³He tubes in the system with uniform coverage on the four sides, the top, and the bottom. Figure 1 shows the complete system and Fig. 2 shows the sample cavity. There are doors on both the front and back of the system for the potential of having the drums pass through the system. Figure 3 shows the electronics cabinet with the computer, printer, and the CI2150 multiplicity electronics. The system uses 16 AMPTEK amplifiers.

A complete system description is provided in Canberra's technical manual.²

Fig. 1. Photograph of complete HENC.

Fig. 2. Sample cavity of HENC.

Fig. 3. Electronic cabinet for HENC showing the personal computer (PC) and the CI2150 multiplicity module.

III. PERFORMANCE TESTS

A. High Voltage Bias

The neutron counting rate vs high voltage was measured in Ref. 1 and the results are shown in Fig. 4. The operating voltage was set at 1720 V.

Fig. 4. Detector high voltage plateau for the HENC.

B. Neutron Life Time

The neutron life time was measured using a ^{252}Cf source and different gate lengths. The die-away time is smaller for short gate settings and larger for long gate settings. We operate the HENC at a gate of 128 μs , and the die-away time of 50 μs is representative of the gate times of 64 to 128 μs .

C. Gate Width

Long gate settings are desirable to get a large fraction of the neutrons in the coincidence gate. However, at higher counting rates, the long gate settings give a large accidental counting rate and this increases the statistical error on the results. For waste quantities of plutonium, the totals rates are low and we have selected a gate setting of 128 μs . For plutonium masses greater than 100 g, the appropriate gate setting would be 64 μs but this is outside the normal range for waste drums. The ratio of the doubles rate at a gate of 128 μs to 64 μs is 1.320. Figure 5 shows the relative statistical error vs gate setting for a 3-g plutonium sample (6% ^{240}Pu).

Fig. 5. Normalized measurement error vs gate setting for a 3-g plutonium sample.

D. Deadtime

The deadtime was measured using multiple ^{252}Cf sources at Canberra and at Los Alamos. The results agreed within the statistical errors and we have used the average values shown in Table I.

The deadtime correction is slightly different for the two software codes that can be used to operate the HENC. The Canberra code Neutron Assay System (NAS)³ uses the standard correction for the Reals (R) and the Totals (T) and it uses the multiplicity deadtime with linear corrections (c and d) for the doubles, and triples.¹

TABLE I. Performance Specification for the HENC.	
Parameter	Value
Neutron efficiency for ^{240}Pu	31%
Neutron die away-time	50 μs
Deadtime (a)	0.50 μs
(b $\times 10^{-6}$)	0.161 μs
Multiplicity deadtime	171 μs
correction coefficient c	0.147 μs
correction coefficient d	0.147 μs
Coincidence gate	128 μs
Predelay	3.0 μs
Doubles calibration coefficient	53.8 counts/s $\cdot$ g ^{240}Pu
Multiplication constant (ρ_0)	0.178

The Los Alamos software NCC⁴ uses the same corrections for T and R as NAS, but it uses the same corrections for the singles (S) and doubles (D) as were used for T and R. The multiplicity deadtime is only used for the triples rate together with the linear correction (d).

In practice, these differences in the deadtime corrections are completely negligible for waste quantities of plutonium.

E. Efficiency

The efficiency of the HENC was measured at Canberra giving the results of 29.5% for ²⁵²Cf. This efficiency was remeasured at Los Alamos after shipment giving an efficiency of 29.6% for ²⁵²Cf in the center of an empty drum. The volume-averaged efficiency was 30.2% because the counts at the perimeter of the drum are slightly higher than at the center. Subsequent corrections in the absolute yields of the ²⁵²Cf calibration sources has reduced the measured efficiencies by a few percent.

Initial measurements with polyethylene shavings in the drum gave a significant increase in efficiency indicating that the HENC had a wall thickness that was too thin for matrix independence. A series of measurements were performed to determine the optimum wall thickness. Figure 6 shows the measurement geometry that was used to evaluate the efficiency change with the polyethylene thickness at the sample. The HDPE thickness was increased and the counting rate for a small ²⁴⁰Pu sample was measured at each thickness. The results are shown in Fig. 7 when the rates have been normalized to the original as-built condition. We see that the efficiency increases up to a thickness of about 1.8 cm. The doubles rates peak before the singles rates because the thermalized neutrons at the sample location do not make it into the coincidence gate because of the long flight time to the ³He tubes.

The addition of polyethylene at the sample location has a different effect than the addition on the wall near the detector. On the detector wall, the HDPE moderates the direct neutrons from the sample as well as the reflected neutrons from the other interior walls. The HDPE at the sample location has less impact on the efficiency than the HDPE on the detector wall. Figure 8 shows the efficiency increase with the addition of the HDPE on the detector walls. With the addition of ~0.6 cm of HDPE on the four sides and top of the cavity walls, we see a 3% increase in the ²⁴⁰Pu counting efficiency. When we increase the thickness to ~1.2 cm, the efficiency is approximately the same as for the 0.6-cm increase so we used the 0.6-cm thickness to leave the detector slightly under moderated.

After the addition of 0.6 cm to the wall, we added HDPE at the sample location and we got the results shown in Fig. 9. With the HDPE on the wall, the sample HDPE thickness to get the peak efficiency is ~0.5 cm so the detector is still slightly under moderated.

In summary, the addition of 0.6 cm of HDPE to the wall increased the efficiency to 30% for ²⁵²Cf and 31% for ²⁴⁰Pu. The system also has less matrix sensitivity for low hydrogen loadings in waste drums.

Fig. 6. Geometry for adding HDPE around a ^{240}Pu sample to measure neutron moderation and absorption.

Fig. 7. Change in singles and doubles counting rate as a function of HDPE thickness around the plutonium sample.

Fig. 8. Normalized singles and doubles counting rate as a function of HDPE added to the sample cavity walls.

Fig. 9. Normalized singles and doubles counting rate as a function of HDPE added around the sample after 0.6 cm of HDPE was added to the detector walls.

F. Coincidence Gate Fractions

The fraction neutrons that fall within the doubles coincidence gate can be calculated or measured. The calculated value is

$$f_g = e^{\frac{-PD}{\tau}} \left(1 - e^{\frac{-G}{\tau}} \right),$$

where

PD = Predelay = 3.0 μ s

G = Gate = 128 μ s

τ = Die-away time $\approx$ 50 μ s.

The doubles gate fraction calculation gives

$$f_g = 0.869,$$

and the triples gate fraction is approximately the square of the doubles gate fraction.

$$f_g (\text{triples}) = (0.869)^2 = 0.755.$$

These fractions can be measured using a ^{252}Cf source together with the nuclear decay constants for ^{252}Cf . The measured gate doubles fraction was

$$f_g (\text{doubles}) \sim 0.70.$$

The gate fractions are not used for the normal AS applications, but they are used for passive multiplicity counting if plutonium standards are not available. The triples gate fraction is adjusted to give the correct ^{240}Pu mass when standards are available.

G. Multiplication Constant

For two-parameter passive coincidence counting, a multiplication constant is defined as

$$\rho_0 = \frac{R_0}{T_0} (1 + \alpha),$$

where

R_0 = Doubles rate for a nonmultiplying sample,

T_0 = Totals rate for a nonmultiplying sample, and

α = Ratio of (α, n) neutrons to spontaneous fission neutrons.

To measure ρ_0 , we used a small (0.705 g) ^{240}Pu sample to give

$$\rho_0 = \frac{37.87}{246.8}(1 + 0.142) = 0.178 .$$

H. Calibration Options

The calibration of the HENC is normally performed for the doubles rate; however, the singles and triples rates can also be used for the calibration. The singles rate gives the most sensitive response but it is subject to variations in α . The triples rate is less sensitive than the doubles for low mass assays.

IV. BACKGROUND MEASUREMENT

The neutron backgrounds in the HENC originate from nuclear materials stored nearby, such as waste drums and from cosmic-ray spallation reactions. The net results come exclusively from cosmic-ray spallation.

The HENC has 30-cm-thick HDPE on the outside of the 10-cm-thick detector banks and this effectively eliminates the room source neutrons. However, the 30 cm of shielding has a negligible reduction in the cosmic-ray coincidence neutrons. The spallation neutrons originating in the floor and room walls are eliminated; however, the HDPE shield creates additional spallation neutrons next to the detector banks.

The effectiveness of the HENC shield for room neutrons was measured in Ref. 1 by placing a ^{252}Cf source inside the sample cavity and then outside the cavity at different distances from the system. Table II shows the normalized singles rates as a function of distance. The doubles background was negligible.

TABLE II. Room Background Rate vs Distance to HENC for a ^{252}Cf Source.	
Distance to HENC (m)	Normalized Totals Rate
Inside	1.00
0.3	1.09×10^{-3}
0.91	0.69×10^{-3}
1.52	0.43×10^{-3}
2.29	0.29×10^{-3}
3.05	0.22×10^{-3}

One conclusion of Table II is that a drum containing ~10 g of plutonium placed ~2.3 m from the outside of HENC would give a singles count equivalent to < 3 mg of plutonium inside the HENC. The doubles rate from the drum's background would be negligible. If the totals rate is used to measure low mass samples, it is necessary to measure the totals background near the time of the sample measurement.

A. Cosmic-Ray Backgrounds

The dominant source of neutron background is the neutrons from cosmic-ray spallation. To investigate the effectiveness of the external HDPE shield on the cosmic rays, we placed a portable coincidence counter⁵ near the HENC and then measured the singles and doubles backgrounds as a function of HDPE shielding thickness.

Table III gives the results of the measurements where the detector background was first measured in the 12-m-deep basement of Building 27, TA-35 at Los Alamos. The detector was then moved adjacent to the HENC and the cosmic-ray background was measured.

TABLE III. Cosmic-Ray Background for INVS.			
Location	Shielding (cm)	S (counts/s)	D (counts/s)
Los Alamos, Bldg. 27 (12 m below surface)	0	1.72	0.013
	10	1.09	0.014
Los Alamos, Bldg. 27 (surface)	0	9.41	0.100
	10	4.88	0.100
	20	4.20	0.098

We conclude from the measurements that the singles and doubles rates increased by factors of 5.47 and 7.21, respectively, at the surface compared with the basement. The 20-cm-thick HDPE shielding decreased the singles background by a factor of 2.24 and the doubles background by only a few percent.

The cosmic-ray backgrounds for HENC at Meriden, CT were a factor of 3 and 4 less for singles and doubles, respectively, compared to Los Alamos.

B. Statistical Filter for Cosmic-Ray Rejection

The cosmic-ray spallation events have a higher average multiplicity than the spontaneous fission events. The detectability limit is a function of the neutron time-correlation background from cosmic rays, and we have reduced the background by eliminating the cosmic-ray spallation events that produce high-multiplicity neutrons by using a statistical filter. The cosmic-ray events can be counted as prompt charged-particle reactions in the detector tubes or as spallation-source neutrons that extend in time over the slowing-down time of the detector body. The predelay (3.0 μ s) eliminates the first category because they are short lived and the predelay⁶ vetoes them from the coincidence gate. The spallation neutrons fall within the coincidence gate but often with high multiplicity. We use the data collection software to isolate the high-multiplicity events and to eliminate them from the data averages. We are currently using statistical techniques to accomplish this.

Our statistical filter for background reduction consists of 2.4–3.0 σ rejection threshold from the average of multiple, short data intervals. The normal counting time for a drum is 600 s, and we divide this into 60 intervals of 10 s each. If any interval is more than ± 2.4 –3.0 σ out of the average, we reject that interval from the average. The software statistical quality control (QC)

filter⁴ at the 2.4σ rejection level decreased the doubles background by a factor of 1.35 and it gave a negligible reduction in the singles rate. This type of filter does not interfere with the routine assay of drums with significant plutonium loadings.

The HENC has such a large singles and doubles background rate at Los Alamos that the statistical filter has less benefit than at the lower counting rates. The average doubles background rate is ~ 2.5 counts/s so a 60-s counting interval contains ~ 138 counts. Thus, a cosmic-ray burst with high multiplicity is diluted by the 138 average counts. To help this problem, we decreased the counting intervals to 10 s so the average number of counts is $2.5 \times 10 = 25$ and more of the burst events can be eliminated. There is much less of a problem at sea level because the doubles rate of ~ 0.6 counts/s would result in an average of 6 counts in a 10-s interval and the burst events can be more effectively eliminated.

The triples rate is more effective at eliminating burst events because the average triples rate is about an order of magnitude less than the doubles rate so the cosmic-ray burst stands out above the background.

V. RESPONSE PROFILES

A. Axial Profiles

The vertical response profiles for the singles and doubles rates were measured for the HENC with the original wall thickness and again after the addition of 0.6 cm of HDPE to the upper sections of the sample cavity.

A ^{252}Cf source was measured as a function of distance above the bottom of a 200-L drum and the results are given in Fig. 10 for the original wall thickness. The source was positioned at a radius of 20 cm in a rotating drum to give a volume-averaged response. After the addition of the 0.6 cm of HDPE to the upper walls of the sample cavity, the response was more uniform as shown in Fig. 11.

Fig. 10. Vertical response profiles for singles and doubles rates for HENC prior to adding 0.6 cm HDPE to the cavity walls.

Fig. 11. Vertical response profiles for singles and doubles rates after adding 0.6 cm of HDPE to the upper sections of the HENC detector walls.

B. Radial Profiles

The radial response profiles were measured using a ^{252}Cf source at the midplane of a rotating 200-L drum. The normalized totals and reals rates are shown in Fig. 12. The rates are higher near the perimeter of the drum because of the higher counting efficiency near the sides of the HENC. The volume averaged response was used for the calibrations.

Fig. 12. Radial response profiles for singles and doubles for a ^{252}Cf source in a rotating 200-L drum.

VI. ADD-A-SOURCE MATRIX CORRECTIONS

The HENC was designed to obtain a maximum figure of merit. However, it is also necessary to make the system insensitive to matrix variations in the drum. The initial measurements with HENC showed that the system was under moderated so the additional 0.6 cm of HDPE was added to the walls. This addition changed the AS calibration that was performed in Ref. 1. The AS calibration was measured at Los Alamos for the modified unit.

The basis of the AS method is to measure the matrix perturbation to the counting rate from a small ^{252}Cf source (9×10^4 n/s) on the outside of the sample and to use the information to correct for the matrix perturbation on the inside of the sample. For the present case, we have positioned the AS neutron source at three positions on the side of the 200-L drum. Figure 13 shows a schematic diagram that illustrates the technique in which the AS neutrons originate from ^{252}Cf and that is stopped at three positions along the side of the drum. The plutonium neutrons originate from the drum's interior.

Fig. 13. Schematic diagram of the HENC AS positions for a rotating 200-L drum.

To correct for the matrix perturbation on the neutron signal, the AS method measures each drum both with and without the ^{252}Cf source on the outside of the drum. The measured quantities are

T_0, R_0 = Totals and reals rates from ^{252}Cf for an empty drum.

T, R = Totals and reals rates from a sample drum without ^{252}Cf .

$T(\text{Cf}), R(\text{Cf})$ = Totals and reals rates from a sample drum with the ^{252}Cf .

The net ^{252}Cf reals rate for the ^{252}Cf and a loaded sample drum is

$$R(\text{Cf}) - R = R(\text{net}) .$$

We use the ratio of the empty drum (after source decay correction) to the net loaded drum to make the matrix correction as follows:

$$\left(\frac{R_0 e^{-\lambda t}}{R(\text{net})} - 1 \right) = x ,$$

and the correction factor (CF) is defined as

$$CF = 1 + f(x) ,$$

where $f(x)$ is a polynomial function of x based on empirical measurements. The measured R for a drum is corrected to give

$$R(\text{corrected}) = R(\text{measured}) CF .$$

The functional relationship between the AS perturbation x and the volume-averaged sample perturbation $f(x)$ was determined empirically by measuring a large variety of matrix loadings with the AS.

A separate ^{240}Pu neutron source was counted at six positions in the drum to give a *volume averaged* matrix effect. The average of the six positions was then ratioed to the empty drum case to give the volume averaged perturbation

$$\left[\frac{R'_0(\text{empty vol})}{R'(\text{matrix vol})} - 1 \right] = y(\text{vol av perturbation}) ,$$

where

R'_0 = reals rates averaged over the volume of an empty drum and

R' = reals rates averaged over the volume of the drum with matrix material.

A. Add-a-Source Procedures

To obtain the volume-averaged matrix response, the small ^{240}Pu sample (0.700 g ^{240}Pu -eff.) was counted at six positions in the drum. The neutron source was positioned at two heights (20 cm and 40 cm), and three radial positions (12, 20, and 25 cm). At each position, the source was measured for 10 min in the passive mode.

The ^{252}Cf AS measurements were made at three positions along the outside of the drum. These positions are given in Table IV.

TABLE IV. AS Positions for the HENC.			
Position	Steps	Distance from Home (cm)	Distance from Bottom of Drum (cm)
1	2745	139.7	51.8
2	3079	156.7	34.8
3	3413	173.7	17.8

The AS positions did not continue to the top of the drum to avoid measuring air with partially filled drums.

B. Matrix Materials

The matrix materials that were used in the calibration are listed in Table V. The primary variable that requires the matrix correction is the hydrogen density. The other atoms such as carbon, nitrogen, oxygen, silicon, and aluminum have very little effect on the passive neutron measurement. Thermal-neutron poisons such as boron and chlorine reduce the totals rate but not the reals rate because the thermal neutrons in the drum are eliminated by the coincidence time gate.

Two AS measurement procedures were evaluated. They were

- (1) all detector banks operational and
- (2) the detector banks directly behind the AS position turned off.

The reason for procedure (2) was to enhance the AS perturbation by increasing the fraction of neutrons that interact with the drum.

TABLE V. Matrix Loadings in 200-L Drums Used for AS Calibration.	
Sample	
Empty drum	Empty metal drum weighing 20 kg
Polyethylene drum liner (9.1 kg)	5-mm thick polyethylene liner on inside wall of drum
Boron glass (173 kg)	Raschig rings
Concrete/rubble (220 kg)	High-density concrete block mixed with 10 kg paper
Sand (290 kg)	Dry sand
Iron metal (196 kg)	Iron pieces (~170 kg) mixed with aluminum (~26 kg)
Polyethylene shavings (7.2 kg)	Low-density polyethylene shavings
Polyethylene tubes (32.5 kg)	Chopped polyethylene tubes
Paper (38.4 kg)	Low-density paper
Paper (53.9 kg)	Medium-density paper
Paper (71 kg)	High-density paper
Wood (51 kg)	Wood pieces
Polyethylene beads/vermiculite (66 kg)	Polyethylene beads (41 kg) plus vermiculite (22 kg) plus borax (~31 kg)

C. Results with the Detector Bank Off

The results with the detector bank off (procedure 2) are given in Table VI. We see that the nonhydrogenous matrices such as sand and iron, scatter neutrons with negligible absorption. This scattering reflects the AS neutrons back into the detector bank that has been turned off. However, the scattering increases the volume-averaged response. Thus, the AS perturbation has the opposite sign as the plutonium volume-averaged response. To avoid this ambiguity, we did not use the bank-off procedure.

TABLE VI. Add-a-Source Calibration Data with the Adjacent Detector Bank Off.										
Sample	S av	D av	T av	So/S	Do/D	To/T	As S av	AS D av	AS So/S	AS Do/D
Empty drum	246.8	37.87	4.3	1	1	1	25 823	7420	1	1
Polyethylene drum liner (9.1 kg)	248.3	37.55	3.97	0.994	1.008	1.083	26 082	7400	0.99	1.003
Boron glass (173 kg)	237.4	38.54	4.24	1.04	0.982	1.014	23 130	6670	1.116	1.112
Concrete/rubble (219.5 kg)	257.8	38.49	4.22	0.957	0.984	1.019	24 741	6640	1.044	1.118
Sand (290 kg)	265.3	47.39	6.1	0.93	0.799	0.705	24 542	7166	1.052	1.035
Iron metal (196 kg)	259	44.45	6.2	0.953	0.852	0.694	25 291	7586	1.021	0.976
Polyethylene shavings (7.2 kg)	258.4	39.89	4.67	0.955	0.949	0.92	25 948	7170	0.995	1.035
Polyethylene tubes (32.5 kg)	231.6	26.84	2.44	1.066	1.411	1.76	23 010	5022	1.122	1.478
Paper (38.4 kg)	253.8	35.32	3.44	0.972	1.072	1.25	NA	NA	NA	NA
Paper (53.9 kg)	246.6	32.71	3.09	1.001	1.158	1.392	24 185	5790	1.068	1.282
Paper (71 kg)	228.1	25.84	2.35	1.082	1.466	1.83	22 872	5062	1.129	1.466
Wood (50 kg)	247.7	31.96	2.84	0.996	1.185	1.51	24 400	5867	1.058	1.265
Polyethylene beads/vermiculite (66 kg)	202.2	21.6	1.392	1.221	1.74	3.09	19 003	4434	1.36	1.669
Half polyethylene and half aluminum	100.7	4.58	0.133	2.45	8.27	32.33	NA	2900	NA	2.56

D. All Banks on Results

To obtain the correct scattering results for nonhydrogenous matrix materials, we have used the results with all detector banks turned on. The measurement results and the perturbation summary are given in Table VII.

The AS calibration curve are shown in Fig. 14 and the results with both the bank on and off is shown in Fig. 15. We see that the AS perturbation is increased by the bank-off procedure; however, the results in the negative domain for metals and sand are ambiguous.

The maximum volume-averaged plutonium perturbation was 0.76 for the drum containing 47 kg of polyethylene beads plus 19 kg of vermiculite and borax. This is a higher hydrogen density than can be found in normal combustible matrix materials.

The polynomial fit to the AS calibration data gave the following coefficient for the function:

$$y = a_0 + a_1x + a_2x^2 + a_3x^3 ,$$

where

$$\begin{aligned} a_0 &= -0.0131, \\ a_1 &= 1.378, \\ a_2 &= 1.916, \text{ and} \\ a_3 &= 5.86. \end{aligned}$$

TABLE VII. AS Calibration Data with all Detector Banks On.

Sample	S av	D av	T av	So/S	Do/D	To/T	As S av	AS D av	AS So/S	AS Do/D
Empty drum	246.8	37.87	4.3	1	1	1	29 820	10 030		1
Polyethylene drum liner (9.1 kg)	248.3	37.55	3.97	0.994	1.008	1.083	30 253	10 148	0.986	0.988
Boron glass (173 kg)	237.4	38.54	4.24	1.04	0.982	1.014	28 080	9 778	1.062	1.026
Concrete/rubble (220 kg)	257.8	38.49	4.22	0.957	0.984	1.019	30 047	9 925	0.992	1.013
Sand (290 kg)	265.3	47.39	6.1	0.93	0.799	0.705	30 220	10 985	0.987	0.913
Iron metal (196 kg)	259	44.45	6.2	0.953	0.852	0.694	29 940	10 696	0.996	0.94
Polyethylene shavings (7.2 kg)	258.4	39.89	4.67	0.955	0.949	0.92	30 397	10 055	0.981	0.997
Polyethylene tubes (32.5 kg)	231.6	26.84	2.44	1.066	1.411	1.76	27 798	7 540	1.073	1.33
Paper (38.4 kg)	253.8	35.32	3.44	0.972	1.072	1.25	29 798	9 278	1.001	1.081
Paper (53.9 kg)	246.6	32.71	3.09	1.001	1.158	1.392	29 389	8 500	1.015	1.181
Paper (71 kg)	228.1	25.84	2.35	1.082	1.466	1.83	27 858	7 711	1.07	1.3
Wood (51 kg)	247.7	31.96	2.84	0.996	1.185	1.51	29 203	8 594	1.021	1.167
Polyethylene beads/ vermiculite (66 kg)	202.2	21.6	1.392	1.221	1.753	3.09	25 650	6 942	1.163	1.448
Half polyethylene and half aluminum	100.7	4.58	0.133	2.45	8.27	32.33	25 991	7 066	1.147	1.419

Fig. 14. AS calibration function (all banks on) for the matrix materials listed in Table V. The horizontal axis is the AS perturbations and the vertical axis is the volume averaged ²⁴⁰Pu doubles rate perturbation.

Fig. 15. AS calibration data with all detector banks on and with one detector bank turned off.

These coefficients are based on a volume-averaged loading pattern. If the loading is concentrated in the central area of the drum as would be the case of a radial-averaged loading, the plutonium perturbation increases more rapidly with hydrogen concentration. For this radial-averaged loading case, the AS correction function coefficients are

$$\begin{aligned} a_0 &= -0.0224, \\ a_1 &= 1.401, \\ a_2 &= -2.041, \text{ and} \\ a_3 &= 8.493. \end{aligned}$$

These coefficients are entered in the software under the AS calibration. An empty drum was counted to obtain the AS ^{252}Cf reference rate. This rate was measured at TA-35, Bldg. 27 to be

$$R_0 = 10\,027 \text{ on } 11/1/96.$$

When the AS correction is applied to the volume-averaged plutonium rates, we get the corrected rates shown in Table VIII and Fig. 16. The corrected values are normalized to the empty drum case so that the empty drum calibration can be used. The average deviation of the corrected rates from the true values was 3.1%. The boron glass matrix (Raschig rings) has very little effect on the real rate. The heterogeneous mixture of concrete rubble and compacted paper rate agreed with the known value to within 2%.

TABLE VIII. AS Corrected Doubles Rate.

Sample	S av	D av	Do/D	AS Do/D	CF	DxCF	Norm. D
Empty drum	246.8	37.87	1	1	1	37.87	1.00
Polyethylene drum liner (9.1 kg)	248.3	37.55	1.0085	0.988	1	37.55	0.99
Boron glass (173 kg)	237.4	38.54	0.9826	1.026	1.005	38.5	1.02
Concrete/rubble (220 kg)	257.8	38.49	0.9839	1.013	1	38.49	1.02
Sand (290 kg)	265.3	47.39	0.7991	0.913	0.845	40.04	1.06
Iron metal (196 kg)	259	44.45	0.852	0.94	0.892	39.6	1.05
Polyethylene shavings (7.2 kg)	258.4	39.89	0.9494	0.997	1	39.89	1.05
Polyethylene tubes (32.5 kg)	231.6	26.84	1.411	1.33	1.437	38.57	1.02
Paper (38.4 kg)	253.8	35.32	1.0722	1.081	1.069	37.78	1.00
Paper (53.9 kg)	246.6	32.71	1.1577	1.181	1.19	38.92	1.03
Paper (71 kg)	228.1	25.84	1.4656	1.3	1.377	35.58	0.94
Wood (51 kg)	247.7	31.96	1.1849	1.167	1.172	37.46	0.99
Polyethylene beads/ vermiculite (66 kg)	202.2	21.6	1.7532	1.448	1.75	37.8	1.00
Average	244.85	35.27				38.31	1.01
St dev (counts/s)	±16.8	±7.36				±1.19	±0.0316
St dev (%)	±6.88	±20.9				±3.10	±3.12

Fig. 16. Normalized doubles rates with before and after the AS correction for the materials listed in Table V.

VII. HENC CALIBRATION

The HENC was calibrated on November 1–3, 1996, using NIS-5 standards. The counter walls were configured with the standard 0.25-in. HDPE liners and the measurements were performed inside an empty drum using positions to give a volume-averaged response.

A. Standards

The standards were a set of five MOX pellets and a small ^{240}Pu oxide sample. The MOX pellet standards were prepared by the GE Vallecitos Nuclear Center under contract number LC3-18688-1 in 1973 with traceability to GE fuel batch history and the National Bureau of Standards (NBS). The physical description of the MOX standards is listed in Table IX. The isotopics are listed in Table X for both the plutonium and uranium.

The ^{240}Pu enriched standard (FZC-158) was prepared at Los Alamos in 1978 and the isotopics and ^{240}Pu -eff mass are listed in Table XI. This sample contains less than 1 g of plutonium and the alpha value was 0.142 on October 26, 1996.

Because of the small fissile masses in the pellet and powder samples, the neutron multiplication is negligible for the waste calibration.

TABLE IX. MOX Pellets Specifications.

Standard	Dimensions (in.)			Mass (g)	Density (% TD)*	
	Length	Diameter	Annulus		Fuel	Bulk
A1-066	0.4143	0.4716	0.146	10.7721	91.67	82.88
A1-081	0.4133	0.4706	0.147	10.6510	91.25	82.35
A1-078	0.3925	0.4719	0.148	10.1886	91.41	82.41
A1-089	0.4067	0.4718	0.150	10.5145	91.60	82.34
A1-119	0.2508	0.1952	—	1.2452	92.13	92.13

*Theoretical density = 12.87 g/cc

TABLE X. Los Alamos MOX Pellets Specifications

		August 30, 1991											
		Wt % (Relative to total plutonium mass)											
ID	g Pu	²³⁸ Pu	²³⁹ Pu	²⁴⁰ Pu	²⁴¹ Pu	²⁴² Pu	²⁴¹ Am	P (mW/g Pu)	g ²³⁹ Pu -eff	g ²⁴⁰ Pu -eff	Alpha	MOX Ratio U/Pu	Enrich (%)
A1-066	0.1526	0.0856	88.4296	10.3053	0.8844	0.2951	1.8465	5.0603	0.9552	0.0168	0.9729	61.3413	0.7400
A1-081	0.5077	0.0856	88.4296	10.3053	0.8844	0.2951	1.8465	5.0603	1.2561	0.0559	0.9729	17.5092	0.7400
A1-078	0.8061	0.0856	88.4296	10.3053	0.8844	0.2951	1.8465	5.0603	1.4778	0.0888	0.9729	10.0922	0.7400
A1-089	0.2267	0.2354	77.8275	18.8824	1.8130	1.2417	3.8233	8.6038	0.9880	0.0489	0.8665	39.7135	0.7400
A1-119	0.2651	0.0519	87.1700	11.8495	0.7160	0.2126	1.4045	4.440	0.7709	0.0327	0.7521	3.1112	92.8100

$$^{240}\text{Pu-eff} = 2.52 \text{ }^{238}\text{Pu} + ^{240}\text{Pu} + 1.68 \text{ }^{242}\text{Pu}.$$

TABLE XI. FZC-158 Isotopics.	
Isotope	Wt %
	12/15/78
^{238}Pu	0.016
^{239}Pu	0.955
^{240}Pu	93.76
^{241}Pu	0.699
^{242}Pu	4.56
^{241}Am	0.001
$^{240\text{e}}\text{Pu}$	0.705 g

B. Calibration Procedure

The standards were positioned in the drum to give a volume-averaged response over the interior of the empty drum. The five MOX pellets were attached to a thin aluminum tube to give uniform vertical spacing between 10 and 60 cm from the bottom of the drum. The radius was set to 12 cm and the drum was rotated for a 60×30 -s counts. The drum was removed from the counter, the sample radius was increased to 20 cm, and the drum was counted for 60×30 -s counts. The drum was removed from the counter, the sample radius was increased to 25 cm, and the drum was counted for 60×30 -s counts. The average of the three positions was used to give a volume-averaged response.

This procedure was repeated with three MOX pellets on the tube and with the FZC-158 standard and the MOX pellets on the tube. In each case, the volume-averaged response was used.

C. Calibration Results

The calibration measurement results are listed in Table XII where the room background has been subtracted from the rates. The ^{252}Cf source CR-6 was counted at the same time in the center of the drum to give a reference rate for future normalizations. The AS reference rate is also given in Table XIII.

The HENC can be used in the singles mode (S), the doubles mode (D), and the triples mode (T). All three modes are calibrated and the data are in Table XII. Figure 17 shows the calibration curve for the normal doubles mode and the triples mode. The linear calibration is forced through the origin because the cosmic-ray and room backgrounds have been subtracted. The doubles calibration slope is 53.8 ± 0.2 counts/s $\cdot$ g ^{240}Pu -eff.

For high-density matrix materials such as iron, it is necessary to adjust the cosmic-ray background subtraction. The combustible matrix materials have approximately the same background as the empty drum.

TABLE XII. HENC Calibration Data.

Sample	²⁴⁰ Pu (g)	Time (s)	S (counts/s)	D (counts/s)	Sigma (counts/s)	T (counts/s)	Alpha (a)	S	D
								(1+a)	g ²⁴⁰ e
A1-066, 081, 078	0.161	60 x 30	98.93	8.863	0.06	0.971	1.013	49.14	55.05
A1-066-078,081, 089, 119	0.243	60 x 30	143.7	13.12	0.09	1.42	1.013	71.39	53.99
FZC-158	0.705	60 x 30	246.8	37.87	0.28	4.3	0.142	216.1	53.72
A1-066, 078, 081, 089, 119, FZC-158	0.948	60 x 30	390.5	50.98	0.31	5.41	0.365	286.1	53.78
²⁵² Cf, CR-6 Centered (11/1/96)	NA	20 x 30	2120	646.4	3.5	118.1	NA	NA	NA
PDP1-3.0 ^a	0.1833	40 x 30	158.5	9.8	0.1	1.25	0.787 ^b	NA	53.46
PDP1-3.0. + 0.5 ^a	0.2138	59 x 10	186.2	11.94	0.1	1.88	0.787 ^b	NA	55.84
PDP1-3.0 + 0.5 + 0.1 ^a	0.2199	513 x 10	189.4	11.97	0.1	1.37	0.787 ^b	NA	54.43

^aThe PDP standards were counted at two-thirds radius and half height to approximate a volume-averaged response.

^bThe effective α is ~1.80 because of the increase from the diatomaceous earth.

Fig. 17. Doubles and triples rate calibration for the plutonium samples listed in Tables X and IX.

The singles rate calibration is shown in Fig. 18 where the two types of standards give different slopes because of their different alpha values (alpha reaction neutrons/spontaneous fission neutrons). The MOX standards have an alpha of 1.013, whereas the alpha for the enriched ^{240}Pu is only 0.142. The alpha of ~ 1 is typical of low-burnup plutonium oxide and the calibration slope is approximately 600 counts/s $\cdot$ g ^{240}Pu -eff. Materials that contain beryllium, boron, fluorine, magnesium, chlorine, and sodium will have a higher alpha value and a higher singles calibration. When the singles calibration is divided by $(1 + \alpha)$, all of the oxide standards fall on a straight line as shown in Fig. 19.

The triples calibration is shown in Fig. 17 with the calibration slope of 5.83 counts/s $\cdot$ g ^{240}Pu -eff. The triples calibration can be useful when the plutonium mass gets above ~ 0.2 g ^{240}Pu -eff. The primary use of the triples calibration is to flag matrices that have localized hydrogenous shielding and nonuniform plutonium distributions in the shielding.

The doubles calibration is normally used for quantitative assay and the singles calibration is used to obtain a lower detectability limit. The 30 cm of external CH_2 shielding makes the HENC singles rate relatively insensitive to the variations in the room backgrounds.

Fig. 18. Singles rate calibration for the plutonium samples listed in Tables IX and XII.

Fig. 19. Singles rate calibration after correcting for the (α, n) reaction neutrons.

D. Matrix Corrections

All of the calibrations are performed relative to the empty 200-L drum case. For drums containing a significantly high loading of hydrogenous matrix materials, the measured response is corrected back to the empty drum calibration condition. The ^{252}Cf AS measurement is used for this purpose. The AS measurement is made on all drums and the empty drum that was used for the plutonium calibration is the reference condition with no correction. The AS correction is automatically made by the software that is used for routine assay.

If the plutonium is mixed with materials that have a high (α, n) reaction yield, the singles neutron emission is increased. Plutonium standards ($\alpha = 0.787$) were mixed into diatomaceous earth and the singles rate vs ^{240}Pu -eff. mass is shown in Fig. 20. We see that the yield increases from 600 to 865 counts/s.g ^{240}Pu -eff. This corresponds to an α of ~ 1.8 for the diatomaceous earth.

Fig. 20. Singles rate calibration for the plutonium standards mixed in diatomaceous earth.

VIII. MEASUREMENT PRECISION AND BACKGROUNDS

A. Precision

The measurement time required to reach a given statistical precision are shown in Fig. 21. The measured precisions are based on the standards listed in Table IX with an α of ~ 1.0 . Figure 22 shows the relationship of the net doubles rate to the statistical error. As the net doubles rate decreases below ~ 3 counts/s, the background rate starts to significantly increase the statistical error.

Fig. 21. Counting double error (STD) vs time for different plutonium mass values.

Fig. 22. Net doubles rate vs statistical error for different counting times.

B. Backgrounds

The cosmic-ray background varies with the barometric pressure. Figure 23 shows the variation of the singles and doubles rates with time over a five-day period. Both the singles and doubles rates track the inverse of the atmospheric pressure. When the atmospheric pressure is high, the incoming cosmic-ray dose is decreased by absorption in the atmosphere.

The precision for measuring the singles background is much better than the precision for the doubles and triples rates. The precision for the singles background is better than 1% in 10 min. Thus, we can use the singles rate from a short 10-min measurement to predict the doubles and triples rate for that time period. A long overnight measurement is used to establish the singles, doubles, and triples cosmic-ray background ratios. A typical combustible matrix was used to determine the background ratios shown in Table XIII.

Fig. 23. Variation in background rate with time and barometric pressure change.

TABLE XIII. Ratio Singles, Doubles, and Triples. ^a		
	NCC	NAS
Singles (S)	22.25	22.20
Doubles (D)	2.23	2.63
Triples (T)	0.437	0.64
D/S	0.1002	0.1185
T/S	0.0196	0.0288

^aThese rates and ratios will change for each new location for the HENC.

The change in the doubles background from different matrix materials is shown in Fig. 24. These measurements were made using the JCC-21 in a well shielded location. We see that the high density materials such as lead and cadmium give a significant increase in the background. The combustible matrices and concrete rubble have essentially the same background as an empty drum. The high hydrogen matrices such as polyethylene have a lower background than the empty drum case because the neutron absorption in the drum is greater than the cosmic-ray neutron production rate.

Fig. 24. Doubles background rate vs matrix loading for the JCC-21 in a well-shielded location.

C. Background Reduction by Statistical Filters

The cosmic-ray background neutrons have a higher multiplicity than ^{240}Pu spontaneous fission events. The background counting rates in the HENC are relatively high and this reduces the effectiveness of the statistical QC filter.

1. NAS Software. The NAS software applies the QC test to the singles and doubles rates but not to the triples rates. We performed experiments to study the benefit from using short time intervals and different sigma QC rejection tests. The 10-s time bins gave a slightly lower background than the 30-s time bins. The shorter time bins carry the disadvantage of the lost time (or deadtime) for the data transfer from the multiplicity module (CI2150) to the computer. The large amount of data in the multiplicity distributions require about 3.5 s for the data transfer at 9600 Baud. Thus the 10-s data collection intervals introduce an ~30% deadtime. Using two CI2150 modules would cut this in half.

Because of the high cosmic-ray background levels at Los Alamos, we used the 10-s time bins to give a better rejection of the large burst events. However, at sea-level locations the background is ~5 times lower and a time bin of 30 s would be more appropriate.

We performed measurements with 10-s bins to establish the best sigma level for the QC test. Figure 25 shows the doubles background rate as a function of the signal level for the rejection of an individual run. We selected a 2.4 sigma outlier limit to get a doubles background of 2.63 counts/s. The doubles background increased to ~3.0 with the QC test off.

The triples rate background is more sensitive to the QC test than the doubles rate. Figure 26 shows the triples background vs the sigma level for the NAS software. We see that the triples background is ~0.63 for a 2.4 sigma QC test. The triples background with the QC test off increases to ~1.3 counts/s.

It is necessary to set the QC outlier test high enough to pass most of the runs. Figure 27 shows the percent of runs passed as a function of the QC sigma level. The 2.4 sigma level passed ~92% of the background runs. A higher fraction of the runs are passed when a significant mass of plutonium is in the drum because the cosmic-ray background becomes negligible.

Fig. 25. Doubles background vs outlier sigma test level using NAS software and 10-s time bins.

Fig. 26. Triples background vs outlier sigma test level using NAS software and 10-s time bins.

Fig. 27. Percentage of runs passed vs sigma test level for 10-s time bins and NAS software.

2. Neutron Coincidence Counting Software. The QC tests for the Neutron Coincidence Counting (NCC) software are different than for the NAS software because the NCC code applies the QC test to the singles, doubles, and triples; whereas, the NAS does not use the triples rate in the test.

Long background runs were performed with the paper matrix drum (53.9 kg) to simulate typical combustible drum loadings. Time intervals of 10 s, 20 s, 30 s, 60 s, and 120 s were used with total run times of 2 to 12 h. Figure 28 shows the triples background rate vs the sigma level for different data time bins. We see that the triples rate drops to 0.44 counts/s for a 3 sigma QC test. This level is significantly below the 0.63 counts/s using the NAS software. The reason for the difference is that the NCC software performs the QC test on the triples rate.

The NCC doubles background vs the sigma level is shown in Fig. 29 and we see that the rate is 2.23 counts/s for a 3 sigma QC test. This can be compared with the 2.63 counts/s doubles background for the NAS software. The 10-s time bin is significantly better than the larger counting time bins. However, in lower background (doubles) locations, the time bins should be increased to ~30 s to reduce the data transfer time losses.

The percentage of runs passed vs the sigma test level are shown in Fig. 30. We see that ~90% of the runs are passed at the 3 sigma level and 10-s time bins.

For routine operation at Los Alamos, we recommend 10-s runs with a 3 sigma outlier test for NCC software and a 2.4 sigma test for NAS software.

Fig. 28. Triples background rate vs outlier sigma test level for different data collection time bins using NCC software.

Fig. 29. Doubles background rate vs outlier sigma test level for different data collection time bins using NCC software.

Fig. 30. Percentage of runs passed vs sigma level for different data time bins using NCC software.

D. Spallation Background Tests

The spallation background rate in the HENC is a function of the shielding location, cavity size, moderator materials, and construction materials.

A series of measurements were performed to evaluate the background contribution from the moderator and construction materials. Three different multiplicity counters were used for the study.

1. the HENC (LANL, TA-35, Bldg. 27)
2. the HCAS⁷ (hot cell AS system)
3. the portable neutron MOX scanner

The HCAS uses a composite moderator of polypropylene and Delrin with two rows of ³He tubes. The total volume of the composite moderator is ~0.34 m³, and the mass is ~340 kg. The HCAS is located in an underground hot cell with over an order of magnitude less cosmic-ray background than the HENC that is located at the surface level at Los Alamos. The MOX Scanner is a portable multiplicity counter that can be moved to the locations of the HENC and the HCAS to compare the cosmic-ray backgrounds. The MOX Scanner has an efficiency of 50% and a mass of ~60 kg. It can be placed by the side or inside both the HENC and the HCAS to sample the background levels at the two locations.

Table XIV gives the counting rates measured at the different locations in Bldg. 27, TA-35 at Los Alamos. The HENC background rates at Meriden, CT are listed for comparison. The

TABLE XIV. Background Comparison for HENC and HCAS.

Detector and Location	Time(s)	Totals	Doubles	Triples
<u>1st floor high bay</u>				
HENC empty bkg (QC off)	>100,000	23.8	2.97	4.1
HENC empty bkg (QC on)	>100,000	23.8	2.24	0.44
HENC FZC-158 (0.705 g ²⁴⁰ Pu) (QC on)	30 × 10	246.8	37.87	4.11
MOX Scanner by HENC (QC off)	2 × 100,000	9.21	0.11	NA
MOX Scanner by HENC (20 cm CH ₂ , QC off)	600 × 30	4.21	0.098	NA
MOX Scanner inside HENC (QC off)	3 × 10,000	3.69	0.075	NA
<u>2nd floor hot cell</u>				
HCAS empty bkg (QC off)	2000 × 10	45.3	1.175	0.581
HCAS empty bkg (QC on)	2000 × 10	45.2	0.978	0.023
HCAS FZC-158 (QC on)	30 × 10	235.2	30.26	2.73
MOX Scanner by HCAS back (QC off)	3 × 100,000	2.08	0.0061	NA
MOX Scanner by HCAS side (QC off)	2 × 10,000	7.53	0.0058	NA
MOX Scanner inside HCAS (QC off)	2 × 50,000	2.96	0.0039	NA
<u>3rd floor low bkg area</u>				
MOX Scanner (QC off)	2 × 1000	0.68	0.001	NA
<u>Meridian, CT</u>				
HENC empty bkg (QC off)	500 × 60	7.35	0.70	0.21
<u>Summary</u>		<u>Ratios</u>		
HENC/HCAS rate ratios for ²⁴⁰ Pu		1.049	1.26	1.51
HENC/HCAS sample cavity vol. ratio	1.293			
HENC/HCAS sample cavity surface ratio	1.174			
HENC/HCAS iron mass ratio = 75 kg/126 kg	0.60			
HENC/HCAS doubles iron bkg increase 1.68/2.3	0.73			
Normalized doubles bkg ratio HCAS/HENC (0.978/2.24) (1.26) (17.9) (0.73) (1.17) = (D bkg rates) (eff.) (location) (metal) (size)			8.4	

standard plutonium sample FZC-158 was counted in both counters to give a comparison of the efficiency. We see that the singles efficiency of the HENC and HCAS are similar but the doubles and triples efficiencies of the HENC are significantly higher than for HCAS because of the larger gate fractions for the HENC resulting from the shorter die-away times in the HENC.

The MOX Scanner was used to get a ratio of backgrounds between the hot cell and the surface. The measured doubles ratio for the two locations was ~18.

To estimate the relative cosmic-ray backgrounds for the HCAS and the HENC, it is necessary to make corrections for the intrinsic efficiency, the detector size, the location, and the iron inside the two systems.

Both the HCAS and the HENC contain much more iron than is desirable for low-background systems. We estimated that the HCAS has ~75 kg of iron inside the cavity and the HENC has ~126 kg of iron in the rotator and reflector. We used Fig. 24 to estimate a background increase of ~2.3 for the HCAS and ~1.7 for the HENC caused by the iron in both systems.

To compare the backgrounds in the two systems we need to combine the following ratios:

$$\frac{\text{HCAS}}{\text{HENC}} = \underbrace{\left[\frac{0.97}{2.24} \right]}_{\text{meas. bkg rates}} \times \underbrace{[1.26]}_{\text{doubles efficiency}} \times \underbrace{[18]}_{\text{doubles location}} \times \underbrace{\left[\frac{1.7}{2.3} \right]}_{\text{interior iron}} \times \underbrace{[1.17]}_{\text{cavity size}} = 8.4$$

We see that the normalized background rate in the HCAS is ~8 times higher than for HENC. Most of this increase probably comes from the increased spallation rate in the double row composite moderator materials. Compounding the problem is the background increase from the long transport distance for cosmic-ray source neutrons in the high carbon plastics.

E. HENC Backgrounds in 200-L Drums

We loaded 200-L drums with a series of matrix materials to determine the background levels for the HENC at the Los Alamos surface location. Table XV lists the measured backgrounds. We see that the high carbon graphite loading has a significantly higher doubles background than the pure polyethylene case. Figure 31 shows the doubles backgrounds, and Fig. 32 shows the triples backgrounds for different matrix materials.

For routine applications, a typical combustible matrix should be used to establish the background. A long overnight run is used to get accurate ratios for the S/D and S/T and subsequently short ~10-min runs can be used to get a current S rate. The S/D and S/T ratios (see Table XIII) are then used to calculate D and T, which can be entered manually into the background file in the software. For measurements that are near the detectability limit, the appropriate matrix loading should be used for the background subtraction.

TABLE XV. Cosmic-Ray Backgrounds as a Function of Matrix Loading.

Background matrix effects data				
Matrix	Mass (kg)	Singles (counts/s)	Doubles (counts/s)	Triples (counts/s)
Empty	0	23.67	2.61	0.551
Paper	53.9	22.58	2.23	0.437
Polyethylene beads/vermiculite	66	22.25	2.24	0.460
Polyethylene tubes	32.5	22.99	2.30	0.463
Polyethylene shavings	7.2	23.34	2.40	0.492
Wood	51	23.44	2.40	0.521
Boron glass	173	24.54	2.94	0.668
Concrete rubble	220	26.05	2.97	0.678
Graphite block	126.6	26.96	3.13	0.743
Sand	290	27.53	3.68	1.003
Iron	196	30.38	5.70	2.191

Fig. 31. Doubles background rates as a function of matrix type and mass in the 200-L drums.

Fig. 32. Triples background rates as a function of matrix type and mass in the drums.

IX. SUMMARY

When we combine the original MCNP design calculations with the subsequent evaluation, we can make the following recommendations.

1. Include the drum matrix effects in future optimization studies.
2. Increase the distance between the face of the HDPE and the ^3He tubes by 0.826 cm (0.325 in.). If the iron rollers remains in the bottom of HENC, remove one ^3He tube from the bottom detector bank.
3. Consider removing the drum rotator in HENC and making the AS a one stop location under the drum (same as WDAS-2 and -3).
4. Separate the electromechanical functions including the rollers, draw bridge, doors, and rotators from the data collection/analysis software to improve the reliability.
5. Remove software controls that block the use of the detector system when various switches and parameters are out of sequence or range. Just give warning flags in the printout.

6. The efficiency after adding ~0.6 cm of HDPE to the walls was 30.5% for ^{252}Cf and 32% for ^{240}Pu using the old ^{252}Cf reference source (CR-6). The ^{240}Pu efficiency decreases to 31% with the more recently calibrated ^{252}Cf sources.
7. In high background locations such as Los Alamos, use 10-s time bins; however, for lower background locations, use 30-s time bins.
8. For the Los Alamos locations, use 2.4 sigma QC test levels for NAS software and 3.0 sigma levels for NCC software.
9. Use the singles background from a ~10-min measurement and calculate the D and T rates from the known ratios to S. Enter the D and T background values manually into the software.
10. For plutonium masses above ~0.4 g ^{240}Pu , use the multiplicity results to supplement the AS results.
11. Refer to Appendix A for the calibration after moving the HENC to TA-54 at Los Alamos.
12. Appendix B presents the Performance Demonstration Program Round 2 test results.
13. Refer to Appendix B for the system parameters that are loaded in the software.

REFERENCES

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3. Model 5430 Neutron Assay Software Technical Manual, 5430-TEC, V1.0 (August 1996).
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6. M. S. Krick and J. E. Swansen, "Neutron Multiplicity and Multiplication Measurements," *Nucl. Instr. Meth.* **219** (1984) 384.
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APPENDIX A

HENC RECALIBRATION AT TA-54

The HENC was moved from TA-35 to TA-54 at Los Alamos for routine use in the measurement of waste drums. During the movement, the CH₂ sheets (0.6 cm thick) that were added to the sample cavity were moved from the inside of the sample cavity liner to the detector wall. This provided a small increase to the efficiency of ~2% as determined by counting the same ²⁵²Cf source (CR-6) and PDP plutonium standards at both locations.

To determine the plutonium calibration at the TA-54 location, a set of standards were measured. To obtain a radial averaged response, the standards were counted at three radial positions including the center, one-half radius, and a 25-cm radius. The sources were positioned at the mid height of an empty drum.

Table A-I lists the standards and the calibration data. The plutonium oxide was uniformly mixed into diatomaceous earth. The isotopics for the standards are given in Table A-II. The calculated α for the pure oxide was 0.787; however, the mixing with the diatomaceous earth increased the effective α to ~1.80. This results in a high singles calibration.

The doubles and triples calibration are shown in Fig. A-1 where the slopes are 56.35 and 5.73, respectively. The singles calibration is shown in Fig. A-2 where the slope is 886.1. Table A-III gives a comparison of the backgrounds and calibration constants at TA-35 and TA-54.

The calibration slopes are the same for both the NAS and NCC software codes. However, the backgrounds are slightly different for the two codes because of the differences in the QC test to reduce the cosmic-ray spallation events. The NCC code makes use of the triples rates, whereas, the NAS code does not.

The HENC system can be used with either software code and both codes are loaded in the computer.

TABLE A-I. HENC Calibration Data for PDP Standards at TA-54.

Sample	Time (s)	m _{Pu240 eff} (g)	Singles (counts/s)	Doubles (counts/s)	σ_{Doubles}	Triples (counts/s)	D/m _{Pu240 eff}
PDP 0.5 g + 0.1 g	2097	0.03612	32.4	0.40	0.22	0.447	59.38
PDP 3.0 g	3 × 600	0.1806	161.0	10.03	0.23	0.924	55.55
PDP 10.0 g	3 × 600	0.602	532.5	33.83	0.55	3.342	56.20
PDP 10 g + 3 g	3 × 600	0.7826	694.1	44.19	0.60	4.592	56.47
²⁵² Cf CR-6 (3/11/97) (centered in drum)	600	NA	2008.2	626.49	2.31	116.267	NA

TABLE A-II. Isotopics for LANL PDP Standards.	
Isotope	Weight (%)
²³⁸ Pu	0.0145
²³⁹ Pu	93.76
²⁴⁰ Pu	5.945
²⁴¹ Pu	0.223
²⁴² Pu	0.0559
²⁴¹ Am	0.0873
(as of 7/15/95)	

TABLE A-III. Background and Calibration Constants at TA-35 and TA-54.		
Item	TA-35	TA-54 ^a
<u>NCC</u>		
Singles bkg (empty drum)	23.65	28.0
Doubles bkg (empty drum)	2.61	3.0
Triples bkg (empty drum)	0.551	0.65
<u>Empty Drum</u>		
Singles calibration slope (PDP)	865	886
Doubles calibration slope	53.8	56.35
Triples calibration slope	5.83	5.73

^aThe CH₂ liner was moved to the detector slab wall.

Fig. A-1. Doubles and triples calibration at TA-54 using PDP standards. The data points correspond to the radial average positioning.

Fig. A-2. Singles calibration (radial averaged) for the diatomaceous earth standards at TA-54.

APPENDIX B

HENC CALIBRATION FOR PDP DRUMS

The two PDP test drums (PDP-002 and PDP-003) were measured using two software codes that have identical data collection and analysis algorithms. The first code was Canberra's neutron assay system (NAS) software that used an OS-2 operating system. The second code was the Los Alamos neutron coincidence counter (NCC) software that is Windows based.

For the PDP drum measurements, the codes were loaded with identical calibration and dead-time parameters. The PDP drums were measured November 20–21, 1996, using the NAS code and November 21–22, 1996, using the NCC code. The NCC measurement results are presented here.

The PDP drums were measured using 10 min of passive counting followed by 3 min of ^{252}Cf AS counting for making matrix corrections. The measurement data were taken in 10-s time bins to aid in the rejection of cosmic-ray spallation background neutrons. A statistical QC test was used to reject outlier data points resulting from the cosmic-ray events.

The isotopics were assumed to be 6% ^{240}Pu -eff. and the calibration constant of 53.8 counts/s g ^{240}Pu -eff. was used for all of the measurements. MOX pellets and a ^{240}Pu oxide powder sample were used for the calibration.

Table B-I gives the results of the six repeat measurements for each of the drums. The measurement routine was identical to what would be used for an unknown sample drum. The drum was completely discharged and reloaded for each of the cycles.

The average measurement results gave:

$$0.906 \pm 0.047 \text{ g Pu for PDP-002}$$

and

$$9.00 \pm 0.10 \text{ g Pu for PDP-003 ,}$$

where the uncertainties are 1 sigma value based only on counting statistics. The uncertainty in the matrix correction was ~5%.

A more complete analysis of the data using the singles, doubles, and triples ratio is given below. The agreement using the three different calibration curves shows that the results are accurate.

PDP DRUM ASSAY USING SINGLES, DOUBLES, AND TRIPLES

The HENC system gives the multiplicity data in addition to the normal coincidence data and the samples can be assayed using the singles and triples rates as well as the doubles rates. We calibrated the HENC for singles and triples counting. If we get agreement for the doubles (D) and triples (T) assays, we can conclude that the matrix correction is correct.

The singles rate is alpha dependent so we used the PDP1-3.0 standard to simulate the alpha increase in a diatomaceous earth standard. The neutron counting rate was 865 counts/s/g ^{240}Pu -eff. compared with 304 counts/s/g ^{240}Pu -eff. for a zero alpha sample. The isotopic based alpha (oxide) for PDP1-3.0 was 0.787 and the measured value of alpha was 1.85. For a benign "alpha matrix" such as an oxide, the expected count rate for 1 g of ^{240}Pu would be

$$S = 304 (1 + \alpha) = 304(1.787) = 543 \text{ counts/s ;}$$

however, the measured rate was

$$S = 304(1 + 1.85) = 865 \text{ counts/s.}$$

Thus, the diatomaceous earth increased the totals rate by a factor of 1.59.

The AS correction factor was measured using drums loaded with a wide range of matrix materials. The AS correction factor (CF) is small except for drums with a high hydrogen content. The AS correction is routinely made for the doubles rate and the singles CF and triples CF can be derived from the doubles correction as follows:

$$CF(\text{singles}) = CF(\text{doubles})^{0.18}$$

$$CF(\text{triples}) = CF(\text{doubles})^{1.8}$$

Table B-II gives the measurement results for PDP-002 and PDP-003 drums for S, D, and T counts and the results are displayed in Fig. B-1. The statistics for the average triples rate for the low mass sample was marginal at ~20%, so the 25% deviation from the true mass is not significant.

We see that all three assay results agree for both matrix loadings so we can conclude that our matrix correction was accurate and our alpha assumption was correct. For samples that have a high plutonium content, we can use the required agreement between the doubles and triples assay to make an accurate matrix correction.

TABLE B-I. PDP Drum Results Using NCC Software.

	Real (counts/s)	Reals sigma	AS Corr.	Pu (g)	Sigma (g)	%
PDP-002 (11/27/96)	2.511	0.125	1.409	1.083	0.056	5.2
	2.315	0.128	1.397	0.99	0.057	5.8
	1.961	0.121	1.34	0.803	0.051	6.3
	2.015	0.123	1.407	0.868	0.054	6.2
	2.117	0.123	1.413	0.915	0.054	5.9
	1.869	0.121	1.358	0.777	0.052	6.7
Av.	2.131			0.906	0.047	5.2
PDP-003 (11/26/96)	23.87	0.466	1.232	9.002	0.201	2.2
	24.55	0.454	1.226	9.21	0.203	2.2
	24.01	0.439	1.224	8.99	0.191	2.2
	23.78	0.435	1.199	8.72	0.195	2.2
	24.30	0.449	1.236	9.191	0.203	2.2
	23.67	0.441	1.223	8.855	0.201	2.2
Av.	24.03			9.00	0.010	0.9

TABLE B-II. PDP Drum Assay Using Singles, Doubles, and Triples Rates.										
Reals counts/s	AS corr.	Pu (g)	S counts/s	T count/s	AS corr. S	AS corr. T	S corr.	T corr.	g Pu via S	g Pu via T
PDP-002 (11/27/96)	2.511 2.315 1.961 2.015 2.117 1.869	1.083 0.99 0.803 0.868 0.915 0.777	46.33 47.14 45.95 45.29 45.61 45.25	0.448 0.267 0.11 0.331 0.183 0.091	(CF) * 0.18	(CF) * 1.8			865 counts/s/g	5.83 counts/s/g
Av. $^{240}\text{Pu-eff.} =$	2.131 1.3873	0.906 0.0551	45.928	0.2383	1.061	1.802	48.73	0.429	0.927 0.0563	1.21 0.0736
PDP-003 (11/26/96)	23.87 24.55 24.01 23.78 24.30 23.67	9.002 9.21 8.99 8.72 9.191 8.855	446.6 445.6 446.2 445.7 445.7 444.7	1.984 2.386 2.01 2.301 2.676 2.018						
Av. $^{240}\text{Pu-eff.} =$	24.03 1.2233	9.00 0.547	445.75	2.2292	1.037	1.437	462.3	3.205	8.78 0.534	9.04 0.55

PDP DRUM RESULTS FOR SINGLES, DOUBLES, AND TRIPLES

Fig. B-1. PDP cycle 2 test drum (002 and 003) results using S, D, and T calibrations.

APPENDIX C

HENC SOFTWARE PARAMETER SETTINGS

Prior to using the HENC for assay verifications, it is necessary to load the parameter settings in the software. After moving the HENC to TA-54, the parameters were loaded into the files with the values shown in the following software screens.

Select Method of Calculating Singles, Doubles and Triples Errors [X]

☐ Use theoretical standard deviation
☒ Use sample standard deviation

OK
 Cancel
 Help

Test Parameter Setup [X]

Accidentals/singles test rate limit (1/seconds)
 Accidentals/singles test precision limit (%)
 Accidentals/singles test outlier limit (sigma)
 Outlier test limit (sigma)
 Background doubles rate limit
 Background triples rate limit
 Measurement control chi-squared limit (%)
 Max # of checksum or accidentals/singles failures
 Declared - assay quality check limit (sigma)
 High voltage test limit (%)
 Max # runs for outlier test to be performed after every run
☒ Perform multiplicity checksum test

OK
 Cancel
 Help...

Archive Configuration [X]

Number of days before data file auto archive
 Number of days before data file auto delete
 Number of days before database data auto delete

OK
 Cancel
 Help...

Add-a-source Calibration

Item type: **Pu oxide** Print calibration parameters OK Cancel Help...

Curve type: **$D = a + b \cdot m + c \cdot m^2 + d \cdot m^3$**

Add-a-source Parameters		Calibration Curve Parameters	
Number of runs	4	a	0.000000E+000
a	-2.240000E-002	b	5.636000E+001
b	1.401000E+000	c	0.000000E+000
c	-2.041000E+000	d	0.000000E+000
d	8.493000E+000	Variance a	0.000000E+000
D0 reference date	02/13/97	Variance b	2.000000E-001
D0 average	9593.800	Variance c	0.000000E+000
D0 position 1	9619.000	Variance d	0.000000E+000
D0 position 2	9537.000	Covariance ab	0.000000E+000
D0 position 3	9623.000	Covariance ac	0.000000E+000
D0 position 4	0.000	Covariance ad	0.000000E+000
D0 position 5	0.000	Covariance bc	0.000000E+000
		Covariance bd	0.000000E+000
		Covariance cd	0.000000E+000
		Sigma x (%)	0.000000E+000

Detector Parameter Setup

Shift register type: **MSR4** Shift register serial port: **COM 1** OK Cancel Help...

Predelay (microseconds)	3.00	Multiplicity deadline (1e-9)	171.0000
Gate length (microseconds)	128.00	Deadtime coefficient A (1e-6)	0.5000
2nd gate length (microseconds)	64.00	Deadtime coefficient B (1e-12)	0.1610
High voltage	1680	Deadtime coefficient C (1e-7)	0.1470
Drift away time (microseconds)	50.0000	Doublet gate fraction	0.6660
Efficiency	0.3040	Triplet gate fraction	0.5150

Add-a-source Calibration [X]

Item type: **Pu oxide** [v] Print calibration parameters [OK] [Cancel] [Help...]

Curve type: **$D = a + b * m + c * m^2 + d * m^3$** [v]

Add-a-source Parameters

Number of runs	4
a	-2.240000E-002
b	1.401000E+000
c	-2.041000E+000
d	8.493000E+000
D0 reference date	02/13/97
D0 average	9593.000
D0 position 1	9619.000
D0 position 2	9537.000
D0 position 3	9623.000
D0 position 4	0.000
D0 position 5	0.000

Calibration Curve Parameters

a	0.000000E+000
b	5.636000E+001
c	0.000000E+000
d	0.000000E+000
Variance a	0.000000E+000
Variance b	2.000000E-001
Variance c	0.000000E+000
Variance d	0.000000E+000
Covariance ab	0.000000E+000
Covariance ac	0.000000E+000
Covariance ad	0.000000E+000
Covariance bc	0.000000E+000
Covariance bd	0.000000E+000
Covariance cd	0.000000E+000
Sigma x (%)	0.000000E+000

Detector Parameter Setup [X]

Shift register type: **MSR4** [v] Shift register serial port: **COM 1** [v] [OK] [Cancel] [Help...]

Predelay (microseconds)	3.00	Multiplicity deadline (1e-9)	171.0000
Gate length (microseconds)	128.00	Deadline coefficient A (1e-6)	0.5000
2nd gate length (microseconds)	64.00	Deadline coefficient B (1e-12)	0.1610
High voltage	1680	Deadline coefficient C (1e-7)	0.1470
Die away time (microseconds)	50.0000	Doubles gate fraction	0.6660
Efficiency	0.3040	Triples gate fraction	0.5150