A proposed high-power UV industrial demonstration laser at CEBAF


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ABSTRACT

The Laser Processing Consortium, a collaboration of industries, universities, and the Continuous Electron Beam Accelerator Facility (CEBAF) in Newport News, Virginia, has proposed building a demonstration industrial processing laser for surface treatment and micro-machining. The laser is a free-electron laser (FEL) with average power output exceeding 1 kW in the ultraviolet (UV). The design calls for a novel driver accelerator that recovers most of the energy of the exhaust electron beam to produce laser light with good wall-plug efficiency. The laser and accelerator design use technologies that are scalable to much higher power. We will describe the critical design issues in the laser such as the stability, power handling, and losses of the optical resonator, and the quality, power, and reliability of the electron beam. We will also describe the calculated laser performance. Finally progress to date on accelerator development and resonator modeling will be reported.

Key words: free-electron laser, UV laser, industrial laser, picosecond pulses

1. Introduction

In 1993 a collaboration of industries, universities, and the Continuous Electron Beam Accelerator Facility (CEBAF) in Newport News, Virginia, formed the Laser Processing Consortium with the purpose of using technologies developed at CEBAF to enhance the capabilities of U.S. industry. The Laser Processing Consortium's ultimate objective is to enhance U.S. industry's capability to profitably tailor the surface characteristics of polymers, composites, ceramics, and metals in an environmentally safe manner. To provide this capability, the consortium intends to develop a cost-effective, high-volume surface-processing system based on a high-average-power, wavelength-tunable free-electron laser (FEL). A companion paper gives examples of applications of the FEL [1].

Providing for industrial applications cost-effectively requires an average laser power of 50 to 100 kW delivered in 1 psec pulses at a cost below $0.01/kJ. The laser must be continuously pulsed and should be tunable through wide portions of the UV and IR spectra where surface absorption is strong. The consortium proposed to develop this industrial capacity in two phases. Phase 1 entails construction of a UV/IR FEL demonstrator (hereafter called the UV Demo). The UV Demo, pictured in Figure 1, will provide 1 kW of average laser power with a continuously-pulsed, mode-locked time structure, at wavelengths between 190 and 300 nm, with an option for IR operation between 3 and 20 μm. It will use technology that can be scaled up to meet industrial requirements. The UV Demo will enable validation of this technology, and industry will use this laser to develop and demonstrate processes of commercial interest. Experience gained from the UV Demo will be the basis for the Phase 2 laser. Phase 2 entails production of the full-scale 100 kW device and transfer of the technology to industry.

2. UV FEL Capabilities

An FEL works by extracting power from an electron beam via the mechanism of stimulated emission of bremsstrahlung in the sinusoidally varying magnetic field of a wiggle magnet [2]. Several features of the FEL combine to make it an attractive candidate for commercializing light-based manufacturing:
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- **Average Power:** Although not yet demonstrated, there is good reason to believe that FELs are capable of very high average power. The UV Demo will provide average powers in the kilowatt range. While it is straightforward to upgrade an FEL to higher power levels without increasing the cost substantially (decreasing the cost per kilojoule), existing sources in the UV increase in capital cost superlinearly with average power, so the cost per kilojoule remains high for high power systems.

- **Tunability:** Free-electron lasers are easily tunable by varying either the field strength in the wiggler or the electron beam energy according to the formula

\[
\lambda_\epsilon = \frac{\lambda_w}{2\gamma^2} \left(1 + K^2\right)
\]

\[
K = \frac{\lambda_w eB}{2\pi mc^2},
\]

in which \(B\) is the rms magnetic field of the wiggler, \(\lambda_\epsilon\) is the wavelength of laser radiation, \(\lambda_w\) is the wiggler wavelength, and \(\gamma\) is the energy of the electrons divided by the rest mass energy of the electron \(mc^2\).

The UV Demo will provide light that is tunable across most of the UV spectrum (190–300 nm) and, as an option, the IR spectrum (3 to 20 μm). As a bonus, the entire visible spectrum will be accessible (350–750 nm). This light will have all of the characteristics of high-quality laser emission: narrow bandwidth (typically <0.1%), spatial coherence (1–2 times the diffraction limit), and linear polarization.

- **Temporal Structure:** Existing lasers are either CW at low intensity or have pulses that are much too long to enable efficient surface processing. By contrast, FELs typically generate very short pulses (~1 psec) that are ideally suited to rapid thermal annealing or ablation of near-surface regions [3]. This pulse length matches the time scales of surface molecular rearrangements and vibrations. In ablation applications using high-power excimer lasers, pulse lengths exceed 10 nsec, and these pulses are long enough to interact with gas-phase ejecta with an associated loss of surface-interaction efficiency. The use of picosecond pulses avoids this problem.

![Figure 1. UV Demo schematic. The FEL extracts energy from electrons accelerated in two passes through an energy-recovering superconducting radio-frequency (SRF) driver linear accelerator (linac). A 10 MeV electron beam from the injector attains 105 MeV in the first of two passes of acceleration through the linac. After returning to the injection point clockwise through the low-energy recirculator line (at the center of the machine, between the linac and the wigglers), the beam attains 200 MeV energy by the end of its second acceleration pass. Then it is steered through an FEL wiggler, where it yields about 0.5% of its power as laser light. Gain is produced in the wiggler via a coupling of the wiggler motion to the transverse field of the laser light. The exhaust electron beam then decelerates in two energy-recovery passes through the linac so that most of its energy can be recycled. Finally, the remaining electrons are absorbed when, decelerated to 10 MeV, they reach a cooled, shielded copper beam dump.](image-url)
- **Efficiency and Cost:** The cost per unit energy of light delivered from conventional UV sources is presently too high (of order $0.10$/kJ) for profitable industrial applications. For demonstration and development purposes, the UV Demo will provide light at about $1$/kJ. The cost of a free-electron laser does not scale strongly with the output power, however. Goals for the Phase 2 UV FEL are operation at 10% wall-plug efficiency and a cost of delivered light below $0.01$/kJ.

![Figure 2. Power output versus wavelength calculated for the UV Demo and for the IR FEL option. The short-wavelength fall-off and kinks in the curve near 200 and 300 nm are due to changes in mirror technologies. The fall-off at long wavelengths is due to a lower electron beam energy leading to less available electron beam power.](image)

The highest-priority considerations for the Phase 1 device are efficient, CW operation; low development cost; short development schedule; scalability; and high reliability, availability, and maintainability. All these will be even more important in the Phase 2 device. The Laser Processing Consortium accounted for all four considerations in formulating its development plan for the UV Demo.

One essential aspect of an industrial laser is energy recovery. During lasing, typically <1% of the electrons’ energy is converted to light. It is, however, possible to recover most of the energy in the spent beam and improve the overall efficiency. In the UV demo the electron beam is returned to the driver accelerator where it is decelerated, and most of its energy is converted back to radio-frequency (RF) power at the constituent cavities’ resonant frequency. This recovered RF power is used to accelerate other electrons. The decelerated electrons are then deposited in a cooled, shielded beam dump. Energy recovery greatly reduces the accelerator’s RF power requirements, waste heat, and radiation.

To enhance efficiency further, one can use superconducting radio-frequency (SRF) cavities. The very low wall loss in these cavities enables efficient CW operation. By contrast, room-temperature accelerators using copper cavities generally operate in a pulsed mode, producing typically ~10 μsec of beam several times a second. Accordingly, an SRF accelerator is essential to high average power and energy efficiency in CW acceleration and deceleration of the electron beam. Moreover, SRF technology is well suited for transporting high-current beams while preserving high beam quality at the entrance to the FEL wiggler, and thereby ensuring efficient conversion of electron beam energy into light. Finally, the technology can be scaled up for use in Phase 2.

### 3. UV Demo Design

The user requirements and the design goals of the UV demo laser are shown in Table 1 [4]. The design goals exceed the user requirements to provide a safety margin. To minimize cost and schedule, existing designs or commercially available components are incorporated into the UV Demo wherever possible. Many of these existing designs derive from CEBAF’s 4 GeV recirculating SRF accelerator, which began operation for nuclear physics research in 1995. The injector consists of a 500 keV DC
photocathode gun driven by a commercially available Nd:YLF laser, followed by a copper cavity to bunch the beam and a CEBAF-type SRF quarter-cryomodule operating at 1500 MHz to provide acceleration to 10 MeV. The driver accelerator uses CEBAF SRF cavities operating at 1500 MHz and producing 32 MV per eight-cavity cryomodule. The UV Demo also incorporates a commercially available wiggler and copies of CEBAF’s RF system, control system, and safety system.

Table 1. User requirements and design goals for the UV Demo

<table>
<thead>
<tr>
<th>Wavelength range</th>
<th>User Requirements</th>
<th>Design Goal</th>
</tr>
</thead>
<tbody>
<tr>
<td>UV</td>
<td>190–300 nm</td>
<td>160–1000 nm</td>
</tr>
<tr>
<td>IR</td>
<td>3000–20,000 nm</td>
<td>2500–25,000 nm</td>
</tr>
<tr>
<td>Power</td>
<td>1 kW</td>
<td>1.7 kW</td>
</tr>
<tr>
<td>Optical beam quality</td>
<td>&lt; 2x diffraction limit</td>
<td>1.2 x diffraction limit</td>
</tr>
<tr>
<td>Bandwidth</td>
<td>Within 4x Fourier limit</td>
<td>Fourier limited</td>
</tr>
<tr>
<td>Stability</td>
<td></td>
<td></td>
</tr>
<tr>
<td>Temporal</td>
<td>&lt; 25%</td>
<td>&lt;10% ΔP/P in 30 min.</td>
</tr>
<tr>
<td>Operational</td>
<td>&lt; 1 trip every 10 min.</td>
<td>&lt;1 trip every hour</td>
</tr>
<tr>
<td>Wavelength</td>
<td>&lt; 10^2 Δλ/λ</td>
<td>&lt;2×10^-3 Δλ/λ</td>
</tr>
<tr>
<td>Pointing</td>
<td>&lt; 10 μrad</td>
<td>&lt;1 μrad</td>
</tr>
<tr>
<td>Runtime</td>
<td>&gt; 2.5 hours (gun limited)</td>
<td>Continuous, 1000 hours</td>
</tr>
<tr>
<td>Energy recovery</td>
<td>&gt; 90%</td>
<td>95%</td>
</tr>
</tbody>
</table>

The use of CEBAF cryomodules also provides confidence that the reliability, availability, and maintainability of the UV Demo will be high, based on CEBAF’s experience with building, installing, and operating the 42 1/4 cryomodules in the existing accelerator [5]. A related requirement is that electron beam impingement on the walls of the device be kept low (<5 μA at >25 MeV) to reduce radiation damage, radiation-shielding requirements, and spurious electronic noise. Low beam loss—which is an attribute of the intrinsic large aperture of SRF cavities—also contributes to safe hands-on maintenance of the device. Otherwise, over time, radioactivation of components due to beam impingement could conceivably become a problem.

As noted above, the wavelength of the FEL is determined by the wiggler wavelength, the electron beam energy and the wiggler strength, represented by the wiggler parameter $K$. Significant laser gain requires $K \approx 1$. Although it is certainly possible to produce a wiggler period as short as 1 cm with $K \approx 1$, the resulting wiggler gap is quite small and the available electron beam power is reduced due to the lower electron beam energy. We have therefore chosen a wiggler with a period of 3.3 cm which allows a 1 cm clear vertical aperture for the electron beam and optical mode. This leads to the need for a 200 MeV electron beam to generate a 190 nm wavelength.

Experience with other FELs (confirmed by detailed modeling) suggests that the conversion efficiency of the FEL is $\sim \left(4N_w \right)^{-1}$, where $N_w$ is the number of wiggler periods. To obtain sufficient gain at UV wavelengths requires $N_w \approx 70$, corresponding to 0.35% conversion of electron beam energy to light. Allowing for the significant (~50% net) losses in commercially available optics at the shortest wavelengths, the electron beam power must be 1 MW to obtain a laser power of 1.75 kW. Thus, the average beam current must be $I_{ave} \approx 5$ mA.

It is possible to estimate limits on the electron beam's energy spread and transverse emittance as well. The total phase slip in the wiggler is $N_w$. If the energy spread is greater than $\left(5N_w \right)^{-1}$, then the electrons will dephase by a full optical period with respect to the optical field while traversing the wiggler. To avoid significant reduction in gain, one must therefore keep the fractional rms energy spread within $2 \times 10^{-3}$. To ensure transverse overlap of the electron beam with the optical mode, the normalized beam emittance must be $\gamma \Delta \lambda/(4\pi) \pi$-mm-mrad. This is a difficult condition to meet and it cannot be grossly violated, but a factor-of-two violation does not unduly degrade the beam-mode coupling. This means the normalized transverse emittance ($\varepsilon_t$) must be within approximately 12 $\pi$-mm-mrad.
We have found several areas that introduce risk into the accelerator design.

- One must achieve high brightness with a continuously pulsed, high-charge electron beam.
- Emittance growth in such a machine must be carefully controlled.
- An accelerator with a net synchrotron tune of one-half is inherently unstable and must have active feedback control [6].
- Beam losses in a high-power machine must be kept very small both for machine protection and to ensure RF stability.
- Mirror coatings in the deep UV are lossy. This leads to problems with poor output coupling and mirror distortion.

We also considered other limits that might affect FEL performance. In all other cases, either the FEL performance was insensitive to the parameter or the parameter was easily achieved.

The sensitivity of the laser performance with respect to these parameters is apparent by way of the following scaling relationship for the laser power:

\[ P_{\text{laser}} = \frac{E_b I_{\text{avg}}}{4N_w} \cdot \frac{E_b \lambda_e}{4\pi e_\text{m} c^2} \left( 1 - \frac{3\delta}{G_{ss}} \right) \]

where \( \delta \) represents the non-output-coupling losses (absorption and scatter from the mirrors, and diffraction from apertures), and \( G_{ss} \) is the laser gain at small signal. This relationship includes several simplifying approximations, yet detailed modeling shows that it is a good qualitative representation of the scaling of laser power with parameters associated with the electron beam and photon optics.

The laser power is clearly seen to be most sensitive to the electron beam energy \( E_b \). The average current, normalized transverse emittance, and mirror losses are also critical design parameters. The dependence on the energy spread is hidden in \( G_{ss} \) and is relatively weak. The dependence on \( N_w \) is also relatively weak because \( G_{ss} \) varies linearly with \( N_w \), but the efficiency is inversely dependent on \( N_w \), so the net power dependence is fairly flat. The dependence on peak current is stronger because \( G_{ss} \) also varies linearly with peak current, and so the peak current should be as high as possible. Modeling indicates that the electron beam bunches can be made short enough to produce a peak current \( I_{pk} = 270 \text{ A} \).

One critical parameter not reflected in the scaling relationship is mirror deformation. Detailed modeling of various short-wavelength lasers indicates that laser power is not significantly degraded as long as the mirrors do not deform by more than 1/5 of an optical wave. Detailed analysis of the optical resonator for the UV Demo remains to be done, but the specification is believed to be similar or more relaxed. This is still an extremely tight tolerance at 200 nm, especially with high power loading, and active figure control may be required.

The critical electron beam design values chosen for the conceptual design of the UV Demo after detailed analysis are listed in Table 2 with respect to a nominal laser wavelength \( \lambda_e = 200 \text{ nm} \). Sensitivities of the laser gain and output power with respect to variations in these critical parameters about their design values (\( \lambda_e = 200 \text{ nm} \)) are illustrated in Table 3. Our best estimates of the largest possible degradation are assumed in calculating the performance degradation in the laser. We have found some evidence that coherent synchrotron radiation in the bending magnets of the driver accelerator may lead to a growth in the emittance larger than the estimate given in Table 3. This phenomenon is poorly understood, but since the emittance is a critical parameter, we are trying to better understand the effect and come up with both better estimates of its size and ways to alleviate the problem [7].

As noted previously, energy recovery is an essential aspect of the UV Demo design and must be demonstrated for the Phase 2 device to be realizable. Another important aspect is multipass acceleration (recirculation). With recirculation, the machine generates the required beam energy at the wiggler with a shorter accelerator. Since costs of the cryomodule and associated RF powers dominate the cost of the accelerator, a multipass machine is considerably less expensive than a single-pass machine. To produce a 200-MeV beam with a 32-MeV CEBAF cryomodules preceded by a 10 MeV injector, there are several
options: three passes with two modules, two passes with three modules, six passes with one module, or one pass with six modules. The last option is too costly, and a six-pass machine was deemed too complex. The choice between two passes versus three passes was slightly in favor of the two-pass machine based on reduced complexity and better preservation of beam quality. Thus, the conceptual design of the UV Demo consists of a 10 MeV injector feeding into a two-pass recirculating accelerator comprising three CEBAF cryomodules.

Table 2. Critical electron beam design values for the UV Demo

<table>
<thead>
<tr>
<th>Parameter</th>
<th>Value</th>
</tr>
</thead>
<tbody>
<tr>
<td>Electron beam energy</td>
<td>$E_b = 200$ MeV</td>
</tr>
<tr>
<td>Average current</td>
<td>$I_{\text{avg}} = 5$ mA</td>
</tr>
<tr>
<td>Peak current</td>
<td>$I_{\text{pk}} = 270$ A</td>
</tr>
<tr>
<td>Normalized transverse emittance</td>
<td>$\epsilon_n = 11 \pi$-mm-mrad</td>
</tr>
<tr>
<td>Non-output-coupling loss</td>
<td>$\delta = 8%$</td>
</tr>
<tr>
<td>Mirror deformation</td>
<td>40 nm</td>
</tr>
</tbody>
</table>

Table 3. Sensitivities of laser gain and power versus degradation in critical parameters for 200 nm operation

<table>
<thead>
<tr>
<th>Parameter</th>
<th>Possible degradation</th>
<th>Impact on Power and Gain</th>
</tr>
</thead>
<tbody>
<tr>
<td></td>
<td></td>
<td>Power</td>
</tr>
<tr>
<td>Energy</td>
<td>-10%</td>
<td>32%</td>
</tr>
<tr>
<td>Transverse emittance</td>
<td>+25%</td>
<td>38%</td>
</tr>
<tr>
<td>Energy spread</td>
<td>+33%</td>
<td>12%</td>
</tr>
<tr>
<td>Peak current</td>
<td>-25%</td>
<td>15%</td>
</tr>
<tr>
<td>Optical cavity losses</td>
<td>+25%</td>
<td>13%</td>
</tr>
</tbody>
</table>

A user facility will be constructed on the CEBAF site to house the UV Demo and provide laboratories in which to demonstrate and develop laser-processing applications. The UV Demo will be housed below grade level on the first floor, a “rigid-box” concrete structure that will “float” on compacted soil for stability. The laboratory space on the second floor, a steel-framed structure including areas for housing RF power, electronics, and controls, can be configured into six or more user areas served by optical beamlines from the FEL below. There is substantial room for laser-applications work, with low-conductivity cooling water, dry gas, electrical services, chemical hoods, and exhaust. Quick vacuum interconnects will permit enclosure of the optical beam and transport to any of the areas, and remotely insertable mirrors will allow delivery of the beam to any station desired. All areas will be networked to the computer systems controlling the laser. Shielding between floors will reduce radiation hazards to acceptably low levels typical of laboratory work.

4. Critical Technical Issues

During the design process we identified several critical technical issues. The design was iterated so that the reduction in performance due to each technical risk area was similar. In this way the risk is evenly spread among all possible risk factors. In each case there are means to recover from shortfalls in performance. The development plan arrived at for building the UV Demo is focused on mitigating the critical technical issues.

4.1 Injector

No injector exists which will provide the required combination of continuous average current and emittance. Space-charge effects in the high-current beam are important. If necessary, the injector can be operated at twice the planned pulse-repetition rate and half the charge per bunch to reduce space-charge effects. If cathode lifetime becomes a limiting factor, an alternative cathode can be used.

An injector test stand is already under construction at CEBAF to test the injector design [8]. A shielded enclosure has been modified to accept the injector. Several subsystems such as the photocathode drive laser and the high-voltage power supply have been constructed and tested. Components for the HV gun are almost complete and assembly has begun. The quarter-
cryomodule and RF systems are scheduled to be completed by spring of 1996. Prototypes of non-standard components for this quarter-cryomodule have been fabricated and successfully tested. The performance of CEBAF cavities in the presence of high beam loading will be measured during injector tests. The emittance and charge-density profile of the beam will be carefully measured, and the results will be correlated with analytic theory and numerical simulations.

4.2 Energy Recovery

The basic principles of energy recovery have been demonstrated experimentally [9], but energy recovery at high average beam current with an FEL increasing the energy spread of the exhaust beam has not been demonstrated. Transport must be nearly lossless from ~199 MeV down to the beam dump at 10 MeV. Recirculating, energy-recovering accelerators exhibit instabilities that arise from fluctuations of the cavity fields. Energy changes can cause beam loss on apertures, or, when coupled to energy dependent timing, phase oscillations. Both effects change the beam induced voltage in the cavities and can lead to unstable variations of the accelerating field. Stability analysis for small perturbations from equilibrium has been performed [6] and the threshold current for the CEBAF FEL has been determined to be 300 μA. The model has been extended to include amplitude and phase feedback, with the transfer function of the feedback presently modeled as a low-pass filter. It was found that, for small variations, modest gain frequencies, well within CEBAF’s RF control system capability, are required to stabilize the system.

Instabilities may arise, and the RF control system may need to incorporate additional fast-feedback capability to damp them. An alternative lattice has been developed which has small net beam loading, and it could be incorporated should the baseline control system prove insufficiently responsive to transients. It has the disadvantage of reduced energy acceptance in the first recirculation pass, but it should be absolutely stable. Beam loss must be kept low to reduce radiation damage, shielding requirements, spurious electronic noise, and radioactivation of components. To mitigate beam loss, a few specific apertures could be increased, and local shielding and beam scrapers could be inserted at strategic locations around the machine.

Until actual construction and commissioning of the UV Demo, analysis of energy recovery will continue. Modeling of the system performed to date incorporates known effects, including ongoing analysis of transients related to turn-on of the FEL.

4.3 Optical Cavity

To handle the high average power density in the UV Demo with reasonable mode quality, we plan to use a retro-reflecting, re-imaged, ring resonator (or R5 resonator) [10]. The R5 resonator is a negative-branch, unstable resonator with scraper output coupling. This design has the advantage of using all reflective surfaces. The lowest-order mode of this resonator has a top-hat profile at the scraper and an Airy profile at the wiggler (the Airy profile is the two-dimensional Fourier transform of the uniform disk). Simulations both with a Beer’s law gain profile (the gain was exponential in the distance along the wiggler and had a Gaussian transverse profile with the same size as the electron beam) and with the FEL 3D simulation code FELEX [11] show that the presence of the wiggler bore and a gain medium rounds the top-hat profile and reduces the output coupling for a given magnification. The gain of this mode is less than the gain of a Gaussian mode. Gaussian mode simulations at 200 nm predict a gain of 125% while FELEX predicts approximately 100%. The Gaussian simulation predicts a peak power output of 130 MW while the FELEX simulation predicts a power output of 70 MW. Given a duty cycle of 1.8×10^{-5}, the average power is greater than 1300 W. An interesting feature of the FELEX simulations is that the power output is insensitive to the electron beam focusing, contrary to expectations.

A serious disadvantage of the R5 resonator design is the need for six reflections in each round trip. Though mirrors are quite adequate for wavelengths longer than 250 nm, the reflectivity of mirrors in the deep UV is at best 98–99% for normal incidence and is usually only about 92% for a 54° p-plane reflection. We plan to work with the University of Arizona Optical Science Center to develop deep UV optics with higher p-plane reflectivity. If this does not work, we will use a roof-top reflector with all s-plane reflections, which is quite adequate for the kilowatt demonstrator but is not scalable to the 100 kW device.

To test such issues as depolarization, net reflectivity, length control, and figure control, we are building a full-size mockup of the R5 resonator using mirrors with reflective coatings at 633 nm. This will allow us to test the basic design of the resonator before building the final device.

High-power laser operation will result in a significant heat load (~50 W/cm²) on the mirrors, making it difficult to limit mirror deformation. Moreover, losses in commercially available UV optical coatings increase significantly below 225 nm. Improved optical coatings for the shortest wavelengths are being developed in collaboration with industrial and university labora-
ories that have established reputations in deep-UV optics. An alternative optical cavity which uses commercially available coatings has been developed as a backup. This cavity fulfills the requirements for the UV Demo, but it has limited scalability to higher power. To mitigate mirror deformation, the substrates can be cooled to liquid-nitrogen temperatures where the coefficient of expansion approaches zero. Alternatively, the cavity can be made longer.

The optical system requires development of an advanced mounting, mirror substrates, and optical controls and diagnostics. An optical test stand will be constructed early in the project and subsequently used to validate the various components of the optical cavity. A special optical test bed at the National Institute of Standards and Technology will be used for absorption and scattering measurements of optical coatings prior to their installation on the UV Demo.

5. Conclusion

In summary, industry has significant need for the development of a cost-effective, high-power UV/IR FEL for surface-processing applications. In response, the Laser Processing Consortium formed, and it has proposed a systematic program culminating in the desired machine. The program should proceed in two phases, Phase 1 culminating in the 1 kW UV Demo, and Phase 2 culminating in the full-scale UV/IR FEL. No show-stoppers are foreseen. The proposed plan is a responsible development path that should allow technical approaches to be demonstrated at 1 kW before the scale-up to an industrial 50–100 kW system.

Optimization studies for the UV Demo have produced a robust design that meets the minimum requirements of the users. Several potential problems have been extensively studied and found soluble. We are continuing the study of the emittance growth of short-pulse electron beams, which might be worsened by coherent synchrotron emission [7], and the study of the overall phase and gradient stability of energy-recovered systems. The lattice is also being optimized to reduce delivery time and cost.

We have recently received approval for a kilowatt demonstrator in the infrared. This device is to be built in support of a program of ship defense sponsored by the Office of Naval Research. It is a useful first step in the Phase I plan since the injector, energy recovery, and emittance preservation can all be tested in the planned device for one-third the cost of the Phase I device. The IR FEL will use the same injector as the UV FEL but will only use a single acceleration pass through a single cryomodule. The final energy will therefore be 42 MeV instead of 200 MeV. The peak current will also be at least a factor of five lower than in the UV device, and the emittance constraint is slightly relaxed. The wiggler will be a small fixed-gap commercial wiggler with a wiggler period of approximately 2 cm. This will provide 1 kW of power near 2 μm. The optical resonator will be a nearly concentric design with dielectric mirrors.

6. Acknowledgments

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