RESEARCH MEMORANDUM

WIND-TUNNEL TESTS AT LOW SPEED OF SWEPT AND YAWED WINGS HAVING VARIOUS PLAN FORMS

By

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ERRATA

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THERMODYNAMIC CHARTS FOR THE COMPUTATION OF FUEL QUANTITY REQUIRED FOR CONSTANT-PRESSURE COMBUSTION WITH DILUENTS
By Donald Bogart, David Okrent, and L. Richard Turner
July 1948

The page-size version of figure 6 was incorrectly prepared for reproduction; the corresponding large-size chart enclosed in the back of the report was, however, prepared correctly.

The ordinate label of figure 13 was incorrectly listed nitrogen-air ratio instead of oxygen-air ratio.

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Figure 6. - Fuel-air ratio increment \( \Delta f \) for addition to fuel-air mixture of water-alcohol mixtures containing more than 75-percent water by weight. 
\[
\Delta f = K_2 \Delta f = (f^* + 5g_{f^*} + 5g_{f^*} + \delta_{f^*}).
\]
Figure 13. - Fuel-air-ratio increment $\Delta f$ for addition of liquid oxygen to fuel-air mixture. $\Delta f = K_m K_i f''$. 
WIND-TUNNEL TESTS AT LOW SPEED OF SWEP'T AND YAWED
WINGS HAVING VARIOUS PLAN FORMS

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SUMMARY

Wind-tunnel tests of an exploratory nature have been made at low speed of various small-scale models of swept-back, swept-forward, and yawed wings. The tests covered changes in aspect ratio, taper ratio, and tip shape. Some data were obtained with high-lift devices on swept-back wings and with ailerons on swept-forward wings. The data have been briefly analyzed and some comparisons have been made with the available theory.

The results of the tests and the analyses indicated that the values of lift-curve slope and effective dihedral of swept wings can be computed with a reasonable degree of accuracy in the low-lift-coefficient range by means of existing theories.

In general, reducing the aspect ratio and the ratio of root chord to tip chord resulted in increases in drag and effective dihedral and increased the longitudinal stability near the stall. Cutting off the tip of a swept-back wing normal to the leading edge reduced the effective dihedral at low lift coefficients and gave a slight reduction in the drag at high lift coefficients. Sweeping forward a part of the outer panel of a swept-back wing improved the longitudinal stability and decreased the effective dihedral but also slightly decreased the maximum lift coefficient and increased the drag at high lift coefficients. The use of high-lift devices at either the leading edge or the trailing edge of swept-back wings increased the lift-drag ratio and the effective dihedral at high lift coefficients. An increase in the ratio of root chord to tip chord for swept-forward wings gave decreases in aileron rolling-moment effectiveness that were greater than the values computed for unswept wings.

INTRODUCTION

Much interest in the use of highly swept wings has arisen since the theory of reference 1 indicated the increases in flight critical
Mach number that could be obtained by the use of sweep. The effects of sweep on the low-speed characteristics of wings have long been recognized and theory (reference 2) indicates that the effects may be rather large. Some experimental data on untapered swept-back wings are provided in reference 3. The present paper reports tests made on various swept and yawed wings as an extension of the work of reference 3 to include the additional effects of taper ratio and sweepforward and to provide data for comparison with the theory of reference 2.

COEFFICIENTS AND SYMBOLS

The results of the tests are presented as standard NACA coefficients of forces and moments which are referred in all cases to the quarter-chord point of the mean aerodynamic chord of the model tested. The data for the swept-wing tests are referred to the stability axes (fig. 1(a)), and the data for the yawed-wing tests are referred to the stability axes and to the wind axes (fig. 1(b)).

For the stability axes the coefficients and symbols are defined as follows:

- $C_L$ lift coefficient ($\frac{\text{Lift}}{\rho S}$ where Lift = $-Z$)
- $C_{L_{\text{max}}}$ maximum lift coefficient
- $C_n$ yawing-moment coefficient ($\frac{N}{\rho SB}$)
- $C_X$ longitudinal-force coefficient ($\frac{X}{\rho S}$)
- $C_Y$ lateral-force coefficient ($\frac{Y}{\rho S}$)
- $C_l$ rolling-moment coefficient ($\frac{L}{\rho SB}$)
- $C_m$ pitching-moment coefficient ($\frac{M}{\rho Sc}$)
- $X$ force along X-axis, pounds
- $Y$ force along Y-axis, pounds
force along Z-axis, pounds
rolling moment about X-axis, pound-feet
pitching moment about Y-axis, pound-feet
yawing moment about Z-axis, pound-feet

For the wind axes the coefficients and symbols are defined as follows:

\[ C_D = \frac{\text{drag coefficient}}{q S} \]

force along X-axis, pounds
force along Y-axis, pounds
force along Z-axis, pounds
rolling moment about X-axis, pound-feet
pitching moment about Y-axis, pound-feet
yawing moment about Z-axis, pound-feet

Other symbols are defined as follows:

aspect ratio \( \left( \frac{b^2}{S} \right) \)
free-stream dynamic pressure, pounds per square foot \( \left( \frac{\rho v^2}{2} \right) \)
wing area
airfoil section chord, measured in flight direction
wing mean aerodynamic chord \( \left( \frac{2}{S} \int_0^{b/2} c^2 \, dy \right) \)
wing span
distance along wing span
$V$ air velocity, feet per second
$
\rho$
mass density of air, slugs per cubic foot
$
\alpha$
angle of attack of chord line in stability-axis $XZ$-plane, degrees
$
\alpha'$
angle of attack of chord line in wind-axis $X'Z$-plane, degrees
$
\psi$
angle of yaw, degrees
$
\Lambda$
angle of sweep of airfoil leading edge, positive for sweepback, degrees
$
\Lambda_{\frac{3}{4}}$
angle of sweep of quarter-chord line, positive for sweepback, degrees
$
\Gamma$
angle of dihedral, degrees
$
\lambda$
taper ratio $\frac{\text{Root chord}}{\text{Tip chord}}$
$
\delta_f$
flap deflection, measured in flight direction, degrees
$
\delta_a$
aileron deflection, measured in flight direction, degrees
$
\eta_0$
aerodynamic-center location, percent mean aerodynamic chord

Subscripts:

$L_0$
conditions for zero lift

Symbols used as subscripts denote partial derivatives of coefficients with respect to angle of yaw, angle of attack, flap deflection, aileron deflection, and lift coefficient. For example,

$$
(C_{L\psi})_{C_L} = \frac{\partial}{\partial C_L} \left( \frac{\partial C_{L\psi}}{\partial \psi} \right)
$$
MODELS

The models, which were mahogany wings used in previous investigations in the Langley 7- by 10-foot tunnel, are illustrated in figures 2 and 3. The models having conventional taper were of NACA 23012 airfoil section in planes parallel to the original planes of symmetry. The untapered models were of NACA 0012 and NACA 0015 airfoil section in planes normal to the leading edges. The model having inverse taper had low-drag-type airfoil sections, the ordinates of which are given in table I. The wing tips were faired on only the inverse-taper model. The full-span split flap tested on one of the untapered swept-back models was of \( \frac{1}{16} \)-inch steel and had a chord equal to 25 percent of the wing chord. The half-span split flap tested on the inverse-taper model was of \( \frac{1}{8} \)-inch Masonite and had chords equal to 20 percent of the airfoil section chord. The nose flap (or slat) tested on the inverse-taper model was of NACA 22 airfoil section (reference 4) in a plane normal to its leading edge and had a constant chord equal to \( 8\frac{1}{2} \) percent of the average chord of the part of the wing \( (0.36\text{ to } 0.95) \) over which the flap (or slat) was located.

TESTS AND RESULTS

Test Conditions

The tests were made in the Langley 7- by 10-foot tunnel at dynamic pressures of 16.37 and 9.21 pounds per square foot, which correspond to airspeeds of about 80 and 60 miles per hour, respectively. The test Reynolds numbers (fig. 4) ranged from 620,000 to 1,250,000, the value depending on the dynamic pressure and on the mean aerodynamic chord of the model tested. Because of the turbulence factor of 1.6 for the tunnel, the effective Reynolds numbers (for maximum lift coefficients) ranged from 992,000 to 2,000,000 (fig. 4).

Corrections

Data for only the inverse-taper model have been corrected for tares caused by the model support strut. No tare data were obtained for the other wing models because experience has shown that for the
The data are presented in figures 6 to 43 in three general groups—force-test data, tuft sketches, and comparison plots—and are indexed in table II.

THEORETICAL RELATIONSHIPS

The basic theory for swept and yawed wings as developed by Betz (reference 2) is based on the concept that only the component of velocity normal to the wing leading edge determines the chordwise pressure distribution. Among the simplifying assumptions made by Betz are: The spanwise load distribution is rectangular, the two semispans of a swept wing may be considered independently as yawed wings, and the wing is swept by first setting the panels at an angle of attack and then sweeping the wing in such a manner that the leading edges of the panels remain in a horizontal plane. The last assumption, since it introduces a geometric dihedral, primarily affects the rolling moments, and, since maintaining the panel leading edges in a plane is not a practical arrangement, a series of equations was developed from Betz's work without such an assumption.

The normal-component-of-velocity concept and the assumptions of independent semispans and rectangular span loading, however, were retained in the development of the following equations, which are not all used in the present paper but are presented for future reference:

**Yawed wings**:

\[
C_{L\alpha} = \left( C_{L\alpha} \right)_{\psi=0} \cos^2 \psi 
\]

(1)

\[
C_{L\alpha'} = \left( C_{L\alpha} \right)_{\psi=0} \cos \psi = \left( C_{L\alpha} \right)_{\psi=0} \cos \psi
\]

(2)

\[
C_{L\sigma} = \left( C_{L\sigma} \right)_{\psi=0} \cos^2 \psi
\]

(3)
Swept wings without flaps or camber:

\[ C_{L_\alpha} = \left( C_{L_\alpha} \right)_{\Delta=\psi=0} \cos \Delta \cos^2 \psi \]  \hspace{1cm} (4)

\[ C_{L_\alpha'} = \left( C_{L_\alpha} \right)_{\Delta=\psi=0} \cos \Delta \cos \psi \]  \hspace{1cm} (5)

\[ C_l = \frac{1}{4} C_L \tan \Delta \tan \psi + \frac{1}{4} C_{L_\alpha} \tan \Gamma \tan \psi \]  \hspace{1cm} (6)

\[ C_{L_\psi} \approx 0.004 \dot{\psi} \left( C_L \tan \Delta + C_{L_\alpha} \tan \Gamma \right) \]  \hspace{1cm} (7)

\[ C_{\delta a} = \left( C_{\delta a} \right)_{\Delta=\psi=0} \cos \Delta \cos^2 \psi \]  \hspace{1cm} (8)

Swept wings with full-span flaps or camber:

\[ C_{L_\delta_\alpha} = \left( C_{L_\delta_\alpha} \right)_{\Delta=\psi=0} \cos \Delta \cos^2 \psi \]  \hspace{1cm} (flaps) \hspace{1cm} (9)

\[ \left( C_L \right)_{\alpha=0} = \left( C_L \right)_{\alpha=\Delta=\psi=0} \cos^2 \Delta \cos^2 \psi \]  \hspace{1cm} (camber) \hspace{1cm} (10)

\[ \alpha_{L_0} = \left( \alpha_{L_0} \right)_{\Delta=\psi=0} \cos \Delta \]  \hspace{1cm} (11)

\[ C_l = \frac{1}{2} \left( C_L \right)_{\alpha=0} \tan \Delta \tan \psi - \frac{1}{4} C_{L_\alpha} \tan \Delta \tan \psi \sin \alpha_{L_0} \]

\[ + \frac{1}{4} C_{L_\alpha} \tan \Gamma \tan \psi + \frac{1}{4} \left[ C_L - \left( C_L \right)_{\alpha=0} \right] \tan \Delta \tan \psi \]  \hspace{1cm} (12)
Page 8, equations (6), (7), and (12) should read

\[ C_L = \frac{1}{4} C_L \tan \Delta \tan \psi + \frac{57.3}{4} C_{L\alpha} \tan \Gamma \tan \psi \]  \hspace{1cm} (6)

\[ C_{L\psi} \approx 0.0044 \left( C_L \tan \Delta + 57.3 \ C_{L\alpha} \tan \Gamma \right) \]  \hspace{1cm} (7)

\[ C_L = \frac{1}{2} (C_L)_{\alpha=0} \tan \Delta \tan \psi + \frac{57.3}{4} C_{L\alpha} \tan \Gamma \tan \psi \]

\[ + \frac{1}{4} \left[ C_L - (C_L)_{\alpha=0} \right] \tan \Delta \tan \psi \]  \hspace{1cm} (12)
\[ C_\psi \approx 0.0037 \left( C_L \right)_{\alpha=0} \tan \Lambda - 0.0044 C_{L\alpha} \tan \Lambda \sin \alpha_{I0} \]

\[ + 0.0044 C_{L\alpha} \tan \Gamma + 0.0044 \left[ C_L - (C_L)_{\alpha=0} \right] \tan \Lambda \]  \hspace{0.5cm} (13)

or since

\[ C_{L\alpha} \sin \alpha_{I0} \approx (C_L)_{\alpha=0} \frac{1}{57.3} \]  \hspace{0.5cm} (14)

\[ C_\psi \approx 0.0044 \left[ (C_L)_{\alpha=0} \tan \Lambda + C_{L\alpha} \tan \Gamma + C_L \tan \Lambda \right] \]  \hspace{0.5cm} (15)

Equations (1) to (15) take no account of aspect ratio and taper ratio. For lift and aileron effectiveness these factors may be accounted for approximately in several ways as follows:

1. By use of standard corrections with the aspect ratio and taper ratio based on an unswept wing having the same panels as the swept wing (reference 3).

2. By use of charts developed by Mutterperl (reference 6) which give the span loading and total lift of swept-back wings calculated by a method based on Weighardt's extension to lifting-line theory (reference 7).


For effective dihedral, in order to account for aspect ratio and taper ratio, the following items may be noted:

1. Equations (7), (13), and (15) actually provide only increments in \( C_{L\psi} \) caused by sweep and dihedral.

2. The basic values of \( C_{L\psi} \) may be obtained from Weissinger (reference 9) by using the values of aspect ratio and taper ratio actually existing on the swept wings.
Page 9, equations (13), (14), and (15) should read

\[ C_{l,\psi} \approx \frac{\frac{1}{2}(C_L)_{\alpha=0} \tan \Lambda + \frac{1}{4} C_{L\alpha} \tan \Gamma + \frac{1}{4}[C_L - (C_L)_{\alpha=0}] \tan \Lambda}{57.3} \]  

(13)

\[ C_{l,\psi} \approx 0.0087(C_L)_{\alpha=0} \tan \Lambda + \frac{1}{4} C_{L\alpha} \tan \Gamma + 0.0044C_L \tan \Lambda - 0.0044(C_L)_{\alpha=0} \tan \Lambda \]  

(14)

\[ C_{l,\psi} \approx 0.0044 \left[ (C_L)_{\alpha=0} \tan \Lambda + 57.3 C_{L\alpha} \tan \Gamma + C_L \tan \Lambda \right] \]  

(15)
DISCUSSION

Longitudinal Stability of Swept Wings

Effect of aspect ratio.- As has been shown in references 3 and 10, the pitching-moment curves become increasingly nonlinear as the sweep angle is increased and tend to become unstable near the stall. Decreasing the aspect ratio generally reduces the nonlinearity and tends to make the pitching-moment curve stable near the stall. (See figs. 6, 7, 9, and 36 for example.) The data for all the wings included in the present investigation, both swept-back and swept-forward, agree very well with the summary chart of reference 10 as to the effects of sweep angle and aspect ratio on the pitching-moment characteristics near the stall. As shown in figure 36, increases in aspect ratio moved the aerodynamic center at low lift coefficients slightly back for the unswept and swept-forward wings and slightly forward for the swept-back wings.

Effect of taper ratio.- In agreement with the data of reference 10, the present investigation showed little or no effect of taper on the pitching-moment characteristics near the stall for swept-back wings. (See figs. 13 and 14.) For swept-forward wings, however, increasing the ratio of root chord to tip chord provided a slight stabilizing effect on the pitching-moment curve near the stall. (See figs. 26 to 28.) Increases in the ratio of root chord to tip chord moved the aerodynamic center at low lift coefficients back for swept-back wings, very little for unswept wings, and forward for swept-forward wings. (See fig. 37.)

Effect of high-lift devices.- The use of a full-span split flap at the trailing edge or of a spoiler extending from the nose on an untapered 60° swept-back wing (figs. 7, 8, and 38) had little effect on the pitching-moment curve except for a change in trim produced by the trailing-edge flap. For the inverse-taper swept-back wing (figs. 14 and 38) the use of a half-span center-section split flap at the trailing edge and a half-span tip slat or flap at the leading edge - either separately or in combination - delayed the excessive stability at high lift coefficients and had little effect on the stability at low lift coefficients. All combinations produced some change in trim, and in the order of increasing the negative value of $C_{m}$ at $C_{L} = 0$ the devices are: leading-edge slat, trailing-edge flap, trailing-edge flap and leading-edge slat, trailing-edge flap and leading-edge flap, and leading-edge flap.
Effect of tip modification.- Cutting off the tip normal to the leading edge on an untapered 60° swept-back wing had little effect on the nonlinearity of the pitching-moment curve or on the stability near the stall (figs. 6 and 10) but did move the aerodynamic center back at low lift coefficients (fig. 39). When the outer 40 percent of the wing panels was swept forward, however, the pitching-moment curve became nearly linear and indicated stability near the stall. (See figs. 6, 11, and 39.)

Effective Dihedral of Swept Wings

Effect of aspect ratio.- For unswept wings the slope of the curve of $C_\psi$ against $C_L$ is increased positively as the aspect ratio is decreased. (See reference 9 and fig. 36.) The same effect is shown in figure 36 for untapered swept-back wings. Although insufficient data are available to show directly the effects of aspect ratio on $(C_\psi/C_L)$ for swept-forward wings, the agreement between experiment and calculation shown in the section entitled "Comparison with Theory" supports the argument that aspect-ratio effects on $(C_\psi/C_L)$ are independent of sweep. The maximum value of $C_\psi$ for the swept-back wings (fig. 36) was increased slightly as the aspect ratio was reduced.

Effect of taper ratio.- According to the calculations of Weissinger (reference 9) an increase in the ratio of root chord to tip chord should give a reduction in the positive value of $(C_\psi/C_L)$. That this result is true is indicated by the data of figure 37 for both swept-back and swept-forward wings. The apparent discrepancy for the unswept and for the approximately unswept wings (fig. 37) is attributable to the fact that the tapered wing built with a straight trailing edge had enough sweepback to counteract the small taper-ratio effect. For swept-back wings, increases in the ratio of root chord to tip chord apparently increased the maximum positive value of $C_\psi$ and the lift coefficient at which this maximum value occurred.

Effect of high-lift devices.- The data of figure 38 show that the use of high-lift devices can greatly increase the maximum values
of $C_\psi$ obtained with swept-back wings. The use of a full-span
split flap at the trailing edge of an untapered wing having 60°
sweepback gave an increment in the value of $C_\psi$ at $C_L = 0$,
an increment in the maximum value of $C_\psi$, and an increment in
the value of $C_L$ at which the maximum value of $C_\psi$ occurred.

For the inverse-taper swept-back wing, a half-span center-section
split flap at the trailing edge produced practically no change in
the value of $C_\psi$ at $C_L = 0$, probably because at $C_L = 0$ the
wing tips were carrying a negative load; this load in turn produced
a negative value of $C_\psi$ to counteract the positive increment
provided by the flap. The use of the flap did, however, extend the
curve of $C_\psi$ enough to produce an appreciable increase in the
maximum value of $C_\psi$ and in the lift coefficient at which the
maximum value of $C_\psi$ occurred. For the inverse-taper swept-back
wing the use of the half-span tip-section leading-edge slat
(or flap) - either alone or in combination with the trailing-edge
flap - resulted in little change in the value of $C_\psi$ at $C_L = 0$.
but did increase the maximum value of $C_\psi$ and the lift coefficient
at which the maximum value occurred, probably because the leading-
edge devices improved the flow over the tips at high lift coefficients.
The use of full-span and half-span tip-section nose spoilers extending
forward from the chord plane on the 60° swept-back wing apparently
improved the flow conditions over the wing outer panel and slightly
increased the maximum value of $C_\psi$.

Effect of tip modification: Cutting off the tip normal to the
leading edge on an untapered 60° swept-back wing reduced the slope
of the curve of $C_\psi$ against $C_L$ at low lift coefficients but did
not change the maximum value of $C_{L\psi}$. Sweeping forward the outer 40 percent of the span, however, markedly reduced both $C_{L\psi}/C_{L}$ and the maximum value of $C_{L\psi}$. (See fig. 39.)

**Induced Drag, Maximum Lift, and Stalling of Swept Wings**

**Effect of aspect ratio.** Curves in figures 19 and 36 indicate the effect of aspect ratio on the induced drag, the maximum lift, and the stalling characteristics for unswept straight wings. Reducing the aspect ratio from 6 to 3 increases the drag, since the induced drag varies inversely as the aspect ratio. A reduction in $C_{L_{\text{max}}}$ occurs as the aspect ratio is decreased although the stall angle is higher for the lower aspect ratio.

Wings swept back $30^\circ$ (fig. 15) show generally the same effect as unswept straight wings. When the aspect ratio is reduced from 5.2 to 4.5, an increase in drag and a reduction in $C_{L_{\text{max}}}$ occur. Wings swept back $60^\circ$ (figs. 6, 7, 9, and 36) also show an increase in drag as the aspect ratio is reduced in the lower lift-coefficient range, but at higher lift coefficients the drag of the wing with the smaller aspect ratio is less than that of the wing with the higher aspect ratio. The same effect was obtained in tests of $60^\circ$ swept-back wings in the Langley 300 MPH 7- by 10-foot tunnel (reference 11). The higher drag of the wing with the larger aspect ratio is probably caused by the spanwise flow toward the tips of swept-back wings; this flow results in a thickening of the boundary layer and causes separated flow over the wing. This condition apparently becomes more aggravated at the higher sweep angles as the span is increased and results in a drag increment large enough to offset any decrease in induced drag caused by increasing the aspect ratio.

Aspect-ratio changes have a normal effect on swept-forward wings, as seen in figures 25 and 34. The effect is similar to that for unswept and for $30^\circ$ swept-back wings, but the increase in drag and the loss in $C_{L_{\text{max}}}$ with decreases in aspect ratio appear larger for the swept-forward wings.
Effect of taper ratio.- For unswept wings figure 37 shows that an increase of taper reduced the induced drag, but the apparent increase in $C_{l_{max}}$ for the wing with a taper ratio of 3.0 is probably a false effect since the tapered wings are cambered (NACA 23012) airfoil sections whereas the untapered wing is uncambered. Comparison of the tapered-wing data with data on a rectangular NACA 23012 airfoil section (reference 12) shows no effect of taper on $C_{l_{max}}$. As the wings are swept either forward or back the favorable effect of increased ratio of root chord to tip chord in reducing the induced drag becomes quite large.

Tuft studies of the swept-back wings (fig. 35) indicate that the stall pattern is similar to that observed on other swept-back wings at low Reynolds numbers. At moderate lift coefficients a region of disturbed flow occurs on the leading edge; then the tip stalls and the stall moves progressively toward the center section. Changes in taper did not appreciably affect the general pattern of the stall.

Effect of high-lift devices.- The use of full-span split flaps on the trailing edge of an untapered $60^\circ$ swept-back wing (fig. 7) increased $C_{l_{max}}$ only slightly but did reduce the angle of attack for $C_{l_{max}}$. The drag was increased over most of the lift-coefficient range and became less than for the plain wing only slightly below $C_{l_{max}}$. The full-span nose spoiler tested on the $60^\circ$ swept-back wing (fig. 8) gave a slightly larger increment of $C_{l_{max}}$ than did the split flap but indicated no change in the stall angle. The drag was increased up to a lift coefficient of about 0.6 but was less than the drag of the plain wing above $C_{l} = 0.6$.

Deflecting a half-span split flap on the trailing edge of a $37.5^\circ$ swept-back wing (fig. 14) or adding either a leading-edge slat or flap on the tip increased $C_{l_{max}}$. Deflecting the flap increased the drag up to a lift coefficient of 0.65 and then gave less drag than the plain wing up to $C_{l_{max}}$. The addition of either the leading-edge slat or flap further reduced the drag from a lift coefficient of 0.65 up to $C_{l_{max}}$. The addition of either the leading-edge slat or flap with the trailing-edge flap undeflected reduced the drag in the higher lift range by an amount about equal to that caused by deflecting the trailing-edge flap alone. Deflecting the
split flap had little effect on the stall pattern but use of the tip slat considerably delayed the stall at the wing tip (figs. 35(c) and 35(a)).

Estimates based on aileron data (fig. 30) were made to determine the effectiveness of a split flap on the tip of swept-forward wings. The increment of lift at \( \alpha = 0 \) for the half-span split flap on the tip of a 45° swept-forward wing was slightly greater than that for an inboard half-span split flap on a 45° swept-back wing (reference 3) and almost twice as great as that for an outboard half-span split flap on a 45° swept-back wing (references 3 and 13). Little difference was noted in the increment of \( C_{l_{\text{max}}} \) provided by the split flap on swept-forward and swept-back wings.

**Effect of tip modification.** Cutting off the tip of a swept-back wing normal to the leading edge caused a reduction in drag from a lift coefficient of 0.50 up to maximum lift since the taper ratio was effectively increased (fig. 39). Sweeping the outer 40 percent of the wing forward increased the drag from a lift coefficient of 0.80 to \( C_{l_{\text{max}}} \) and slightly reduced \( C_{l_{\text{max}}} \) probably because of the increased interference between the swept-forward and the swept-back panels.

**Aileron Effectiveness for Swept-Forward Wings**

Data for two 45° swept-forward wings of taper ratio 1.0 and 4.0 equipped with half-span split-flap-type 0.20c ailerons deflected on the left wing only are presented in figures 30 and 33.

Comparisons which accounted for the relative effectiveness of plain and split flaps (reference 13) indicate that the aileron effectiveness \( C_{l_{\text{a}}} \) at a lift coefficient of 0.2 for the 45° untapered swept-forward wing was about 10 percent greater than the value that would be obtained for the 45° untapered swept-back wing of reference 3. This result is probably caused by the thinner boundary layer and the less turbulent flow existing on the tips of swept-forward wings.

The data showed that the loss in aileron rolling-moment effectiveness resulting from increased taper was greater for the swept-forward wing than the loss indicated for unswept wings in reference 14.
COMPARISON WITH THEORY

Yawed-Wing Lift-Curve Slope

The tests of the yawed wings were made primarily to provide a relatively quick preliminary check on Betz's concept of the effect of yaw on the lift-curve slope (reference 2). As shown by figure 40 the data for the NACA 0012 wing of aspect ratio 6 agreed almost exactly with the cosine law. Tests of an NACA 0012 wing of aspect ratio 3, however, showed less effect of yaw on \( C_{L\alpha} \) than is indicated by the cosine law. In an effort to explain the discrepancy, tests were made of two flat plates having aspect ratios of 3, one rectangular and one of infinite taper. As shown by figure 40 the infinite-taper model showed more effect of yaw than the cosine law and the rectangular plate showed less effect. Additional tests of a flat plate having an aspect ratio of 1.27 showed an increase rather than a decrease in \( C_{L\alpha} \) as the model was yawed. These results may be partly explained by the fact that as a rectangle is yawed the span normal to the air-stream direction - and thus the aspect ratio - increases for part of the yaw range. The amount of increase and the angles of yaw over which this increase appears are functions of the aspect ratio and the taper of the basic model. Corrections applied on this basis indicate that all the data would group about the curve for the infinite-taper plate having an aspect ratio of 3. The resulting curve showed a slightly greater effect of yaw than is indicated by the cosine law.

Swept-Wing Lift-Curve Slope

The data of reference 3 indicate that in the computation of the lift-curve slope of swept wings the cosine law is valid provided the aspect ratio used is that of an unswept wing having the same panels as the swept wing. On this basis and by use of the lifting-surface-theory equation for the lift-curve slope (reference 15) figure 41 was derived. By use of figure 41 and a value of 0.099 for the section lift-curve slope the values of \( C_{L\alpha} \) were computed for all the swept-wing tests. The measured and the computed values of \( C_{L\alpha} \) are shown in figure 42. The agreement is reasonably good but indicates, as did the yawed-wing data, that the cosine law does not indicate quite enough drop in \( C_{L\alpha} \) as \( \Lambda \) is increased.
Swept-Wing Effective Dihedral

In the calculation of the effective dihedral the same procedure was followed as in reference 3 except that the aspect ratio and taper ratio as well as the sweep were accounted for by obtaining

\[
(C_L^* \psi_{C_L})_{\lambda=0}
\]

from the following formula of Weissinger (reference 9):

\[
57.3 \frac{\partial C_L^*}{\partial \psi \partial C_L} = 57.3 C_L^* = \frac{2K}{A} \left[ \frac{1 + 0.29(\lambda - 1)}{\lambda + 1} \right] - 0.10 \tag{16}
\]

Reference 9 states that the constant \( K \) is indeterminate but depends on the wing-tip shape and is probably of the order of magnitude of unity for square-cut tips. The data for the NACA 0012 airfoils having aspect ratios of 3 and 6 were used to evaluate \( K \) and a value of 0.78 was obtained.

The values of \( (C_L^* \psi_{C_L}) \) for the models tested in the present investigation were computed by using \( K = 0.78 \) and equations (15) and (16). Figure 43 shows the remarkably close agreement obtained between the measured and the computed values.

CONCLUSIONS

The results of low-speed tests in the Langley 7- by 10-foot tunnel of several small-scale models of yawed and swept wings indicated the following conclusions:

1. The lift-curve slope and the effective dihedral for swept wings can be computed with a reasonable degree of accuracy in the low lift-coefficient range by means of existing theories.

2. In general, reducing the aspect ratio and the ratio of root chord to tip chord produced increases in drag and effective dihedral and slightly increased the longitudinal stability near the stall.

3. Cutting off the tip of a swept-back wing normal to the leading edge reduced the effective dihedral at low lift coefficients and gave a slight reduction in the drag at high lift coefficients.
Page 17, equation (16) should read

\[ 57.3 \frac{\partial^2 C_L}{\partial \psi \partial \theta_L} = 57.3 (C_L \psi)_C L = \frac{2K}{A} \left[ \frac{1 + 0.15(\lambda - 1)}{\lambda + 1} \right] - 0.10 \]  

(16)
RESTRICTED

ERRATA

NACA RM No. L7D23

WIND–TUNNEL TESTS AT LOW SPEED OF SWEPT AND YAWED
WINGS HAVING VARIOUS PLAN FORMS
By Paul E. Purser and M. Learoy Spearman

May 22, 1947

An error has been found in Weissinger's formula (reference 9) which has been corrected in the following reference:


Page 17, equation (16) should read

\[
57.3 \frac{\partial^2 C_L}{\partial \psi \partial C_L} = 57.3 (C_{\psi L})_0 = 0.5 \left\{ \frac{2K \left[ \frac{1 + 0.15(\lambda - 1)}{1 + \lambda} \right]}{A} - 0.10 \right\} \quad (16)
\]

The factor 0.5 converts the formula from terms of semispans as used by Weissinger to spans. With the formula in the corrected form the factor K should be changed from 0.78 to 1.51 (see p. 17, lines 11 and 13). The foregoing equation (16) supersedes the correction to equation (16) in the previous errata of this paper.

RESTRICTED
4. Sweeping forward a part of the outer panel of a swept-back wing improved the longitudinal stability and decreased the effective dihedral but also increased the drag at high lift coefficients and slightly decreased the maximum lift coefficient.

5. The use of either leading-edge or trailing-edge high-lift devices on swept-back wings increased the lift-drag ratio and the effective dihedral at high lift coefficients.

6. An increase in the ratio of root chord to tip chord on a swept-forward wing caused decreases in aileron rolling-moment effectiveness that were greater than the losses computed for unswept wings.

Langley Memorial Aeronautical Laboratory
National Advisory Committee for Aeronautics
Langley Field, Va.
REFERENCES


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NATIONAL ADVISORY COMMITTEE FOR AERONAUTICS
### Table II

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#### Tuft sketches

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#### Comparison figures

- Effect of aspect ratio
- Effect of taper ratio
- Effect of high-lift devices
- Effect of tip modification
- Yawed-wing lift-curve slope
- Lift-curve slope for swept wings
- Comparison of measured and computed lift-curve slopes for swept wings
- Comparison of measured and computed values of effective dihedral for swept wings

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Figure 1. - Systems of axes used. Positive values of forces, moments, and angles are indicated by arrows.
Figure 2. - Plan forms and dimensions of wing models.
Figure 2 - Continued.
Figure 2 cont.

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\[ \frac{\alpha}{4} = 37.5^\circ \]
\[ \alpha = 3 \]
\[ \lambda = 2.04 \]
\[ b = 38.84 \text{ in.} \]
\[ c^1 = 13.17 \text{ in.} \]
\[ c_R = 17.07 \text{ in.} \]
\[ c_T = 8.35 \text{ in.} \]
\[ S = 490 \text{ sq. in.} \]
\[ \bar{x} = 10.65 \text{ in.} \]
\[ \bar{y} = 0.25 \text{ in.} \]

\[ \frac{\alpha}{4} = 37.5^\circ \]
\[ \alpha = 3 \]
\[ \lambda = 0.617 \]
\[ b = 87 \text{ in.} \]
\[ c^1 = 15.61 \text{ in.} \]
\[ c_R = 11.875 \text{ in.} \]
\[ c_T = 19.33 \text{ in.} \]
\[ S = 721.44 \text{ sq. in.} \]
\[ \bar{x} = 12.42 \]

\[ \frac{\alpha}{4} = 30^\circ \]
\[ \alpha = 5.8 \]
\[ \lambda = 1 \]
\[ b = 60 \text{ in.} \]
\[ c^1 = 12.54 \text{ in.} \]
\[ c_R = 11.54 \text{ in.} \]
\[ c_T = 11.54 \text{ in.} \]
\[ S = 692.88 \text{ sq. in.} \]
\[ \bar{x} = 11.548 \text{ in.} \]
\[ \bar{y} = 10.393 \text{ in.} \]

Figure 2.- Continued.
Figure 2.- Continued.
Figure 2.- Continued.
Fig. 2 cont.

Models 21 and 22

Model 21

\[ \frac{A_0}{\lambda} = 30^\circ \]

A

5.2

4.5

\lambda

1

1

b

60 in.

52 in.

a

11.54 in.

11.54 in.

\beta

11.54 in.

11.54 in.

S

602.84 sq. in.

600 sq. in.

\overline{c}

9.774 in.

4.639 in.

Model 22

\[ \frac{A_0}{\lambda} = 30^\circ \]

A

4.5

4.5

\lambda

1

1

b

52 in.

52 in.

a

11.54 in.

11.54 in.

\beta

11.54 in.

11.54 in.

S

600 sq. in.

600 sq. in.

\overline{c}

4.639 in.

4.639 in.

Wind

Model 23

A

3.59

3.59

\lambda

1

1

b

41.5 in.

41.5 in.

a

11.54 in.

11.54 in.

\beta

11.54 in.

11.54 in.

S

480 sq. in.

480 sq. in.

\overline{c}

5.63 in.

5.63 in.

Model 24

A

3.61

3.61

\lambda

2.85

2.85

b

46.6 in.

46.6 in.

a

13.9 in.

13.9 in.

\beta

19.1 in.

19.1 in.

S

600.9 sq. in.

600.9 sq. in.

\overline{c}

1.57 in.

1.57 in.

\overline{E}

0.3 in.

0.3 in.

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Figure 2. - Continued.
Figure 2. - Continued.
Fig. 2 conc.

\[ A_{c/4} = 45^\circ \]
\[ A = 2.1 \]
\[ \lambda = 3.28 \]
\[ b = 36.62 \text{ in.} \]
\[ c_t = 18.51 \text{ in.} \]
\[ c_R = 26.38 \text{ in.} \]
\[ c_T = 6.78 \text{ in.} \]
\[ s = 607 \text{ sq. in.} \]
\[ = 0.97 \text{ in.} \]
\[ r = 0.31 \text{ in.} \]

Model 28

Models 29 and 30

Figure 2.- Concluded.

\[ A_{c/4} = 60^\circ \]
\[ A = 3 \]
\[ \lambda = 1 \]
\[ b = 60 \text{ in.} \]
\[ e_t = 20 \text{ in.} \]
\[ e_R = 20 \text{ in.} \]
\[ e_T = 20 \text{ in.} \]
\[ s = 1200 \text{ sq. in.} \]
\[ = 20.365 \text{ in.} \]
\[ r = 7.99 \text{ in.} \]

Model 29

Model 30

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(a) 45° swept-forward wing.

Figure 3.- Swept wings mounted in test section of Langley 7- by 10-foot tunnel. Front view.
(b) $60^\circ$ swept-back wing with $60^\circ$ swept-forward outer panels.

Figure 3.- Concluded.
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\( \lambda = 0.617 \); low-drag-type section with faired tips and various high-lift devices; model 9.
Figure 15: Aerodynamic characteristics of swept-back wings. For models 10 and 11.

\[ \frac{\alpha}{\lambda} = 30^\circ, A = 5.2 \text{ and } 4.5, \lambda = 1; \text{NACA 0015 airfoil section}; \]

\[ \text{Angle of attack, } \alpha, \text{ deg} \]

\[ \begin{array}{ccccccc}
0 & 4 & 8 & 12 & 16 & 20
\end{array} \]

\[ \text{Lift coefficient, } C_l \]

\[ \frac{A}{4} = 43 \]

\[ \text{Pitching-moment coefficient, } C_m \]

\[ 0 \]

\[ \text{Longitudinal-force coefficient, } C_x \]

\[ \begin{array}{ccccccc}
0 & 0.08 & 0.16 & 0.20 & 0.24
\end{array} \]

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Figure 16.- Aerodynamic characteristics of a swept-back wing. \( \Lambda \frac{c}{4} = 14^\circ; A = 6; \lambda = 3; \) NACA 23012 airfoil section; model 12.
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\[ \frac{\Lambda_c}{4} = 60^\circ; \ A = 6; \ \lambda = 5; \ NACA \ 23012 \ airfoil \ section; \ model \ 13. \]
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\( A/4 = 0.6; A = 6; N = 1 \); NACA 0012 airfoil section; model 14.

Angle of attack, \( \alpha \), deg

Pitching-moment coefficient, \( C_m \)

Longitudinal-force coefficient, \( C_x \)
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\[ \Lambda_{c/4} = 0^\circ; A = 6; \lambda = 1; \] NACA 0012 airfoil section; model 14.

(a) Stability axes.
(b) Wind axes.

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(b) Wind axes.

Figure 21. - Concluded.
(a) Stability axes.

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$\Lambda_{c/4} = 0^\circ$; $A = 3$; $\lambda = 1$; flat-plate section; model 18.
(b) Wind axes.

Figure 22. Concluded.
(b) Wind axes.

Figure 23.— Concluded.
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\( \Lambda_{c/4} = 0^o \); \( A = 3 \); \( \lambda = \infty \); flat-plate section; model 19.
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$
\frac{A_c}{L} = 0^\circ; \ A = 1.27; \ \lambda = 1; \ \text{flat-plate section; model 20.}$
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\[ \Lambda_c/4 = -30^\circ; \ A = 5.2 \text{ and } 4.5; \ \lambda = 1; \ \text{NACA 0015 airfoil section; models 21 and 22.} \]
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Figure 30. - Aerodynamic characteristics of a swept-forward wing. $\Lambda c/\Lambda = -45^\circ$; $A = 2.1$; $\lambda = 1$; NACA 0012 airfoil section with a 0.20c split-flap-type aileron on left wing only; model 26.
Figure 31. - Aerodynamic characteristics of a swept-forward wing. 
\( \Lambda_{C/4} = -46.6^\circ; A = 2.1; \lambda = 2.5; \) NACA 23012 airfoil section; model 27.
Figure 32. - Aerodynamic characteristics of a swept-forward wing. $\Delta_c/4 = -45^\circ$; $A = 2.1$; $\lambda = 3.88$; NACA 23012 airfoil section; model 28.
Figure 33. - Aerodynamic characteristics of a swept-forward wing. $\Lambda c/4 = -45^\circ$; $A = 2.1$; $\lambda = 3.88$; NACA 23012 airfoil section with a 0.20c split-flap-type aileron on left wing only; model 28.
Figure 34. - Aerodynamic characteristics of swept-forward wings. \( \Lambda_{c/4} = -60^\circ \); \( A = 3 \) and 1.5; \( \lambda = 1 \); NACA 0015 airfoil section; models 29 and 30.
(a) $\Delta c/4 = 60^\circ$; $A = 1.5$; $\lambda = 1$; NACA 0012 airfoil section; model 2.

Figure 35.- Tuft studies of various wing plan forms.
(b) $\Delta_c/4 = 56^\circ$; $A = 2.1$; $\lambda = 2.5$; NACA 23012 airfoil section; model 7.

Figure 35.- Continued.
(c) $A_c/4 = 37.5^\circ$; $A = 3$; $\lambda = 0.617$; low-drag-type airfoil section; model 9.

Figure 35.- Continued.
(d) $\Delta c/4 = 37.5^\circ$; $A = 3$; $\lambda = 0.617$; low-drag-type airfoil section with tip slat; model 9.

Figure 35.- Continued.
(e) \( \Delta_c/4 = 8^\circ; A = 8; \lambda = 5; \) NACA 23012 airfoil section; model 13.

Figure 35. - Continued.
(f) $\Delta_{c/4} = -46.6^\circ$; $A = 2.1$; $\lambda = 2.5$; NACA 23012 airfoil section; model 27.

Figure 35.- Concluded.
Figure 36.— Effect of aspect ratio on aerodynamic characteristics of various untapered wings. Models 1, 2, 14, 15, 16, 17, 29, and 30.
Figure 37. - Effect of taper ratio on aerodynamic characteristics of various wings. Models 1, 2, 7, 8, 9, 12, 13, 14, 18, 19, 23, 24, 25, 26, 27, and 28.
Figure 37.— Concluded.
Figure 38. - Effect of high-lift devices on aerodynamic characteristics of swept-back wings. Models 2 and 9.
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Figure 40. - Variation of lift-curve slope with angle of yaw for various wings.
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Figure 42. - Comparison of measured and computed values of $C_{L\alpha}$. Flagged symbols denote swept-forward wings.
Figure 43. - Comparison of measured and computed values of \( \left( \frac{C_{ly}}{C_L} \right) \). Flagged symbols denote swept-forward wings.