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OF TESTS OF 24 AIRFOILS AT HIGH MACH NUMBERS

(Derived from Interrupted Work at Guidonia, Italy
in the 1.31- by 1.74-Foot High-Speed Tunnel)

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Two-dimensional data for 24 airfoil sections tested in the 1.31- by 1.74-foot high-speed tunnel at Guidonia, Italy, are presented. The test Mach numbers ranged from 0.40 to 0.94 and the test Reynolds numbers from 340,000 to 420,000. The results indicate that thickness ratio is the dominating shape parameter at very high Mach numbers and that important aerodynamic advantages are to be gained by using the thinnest possible sections.

The results of preliminary tests made to investigate the effects of jet boundaries, Reynolds number, and humidity at very high speeds are also presented. It was found that the jet-boundary effects became very large at high Mach numbers when models large with respect to the tunnel height were used. In the absence of suitable correction factors for large models it was considered essential to use models small enough to make the jet-boundary effects negligible. It was indicated that the data presented for the 24 airfoils tested are essentially free from jet-boundary and humidity effects.

INTRODUCTION

The rapid increase in airplane speeds during the past 5 years has greatly accentuated the need for experimental data in the subsonic Mach number range above 0.7. Experimental aerodynamic data in this speed range, however, are still very scarce. There are two principal reasons for the lack of data. First, the experimental equipment required to obtain data at high
speeds on models of significant size is extremely costly to construct and operate. Second, the problems of technique involved in obtaining these speeds are very complex and are not yet fully understood. The tunnel-wall-effect phenomena occurring at very high Mach numbers with the presence of shock waves become so complex that there seems little hope at present of obtaining corrections for these effects by analytical methods.

The principal purpose of this report is to present aerodynamic data for 19 related airfoils and for 5 miscellaneous airfoils at Mach numbers in the range 0.40 to 0.94. The data were obtained on models of 1.575-inch and 1.965-inch chord in the 1.31- by 1.74-foot high-speed tunnel at Guidonia, Italy. Before the presentation of the test results, a description is given of the equipment used and the findings of preliminary tests made in an attempt to develop a suitable testing technique and to determine the isolated effects of such experimental variables as Reynolds number, ratio of the size of the model to the size of the tunnel, and humidity.

The results presented herein represent the completed part of a broad high-speed research program at Guidonia, which was interrupted by the war.

I. EFFECTS OF REYNOLDS NUMBER, JET BOUNDARIES, AND HUMIDITY IN TESTS OF AIRFOILS AT HIGH SPEEDS

A systematic study of the effects of Reynolds number, air-stream boundaries, and humidity at high speeds was made prior to the main part of the present investigation. It is not certain, of course, that these are the only factors affecting the results, but they are considered the most important.

WIND TUNNEL

All the tests were made in the high-speed tunnel at Guidonia (reference 1), a single-return tunnel that could function at a pressure below atmospheric. The pressure in the test section of the tunnel could be varied from 1.0 atmosphere to 0.1 atmosphere. The tunnel had a
system of refrigeration by which the temperature at low speeds could be held constant at as low a value as 15°C Centigrade. The temperature of the air as it left the compressor was very variable, depending on the velocity and the pressure of the jet.

The tunnel was powered by a 3000-horsepower fourteen-stage axial-flow compressor, which could produce a velocity ranging from 0.4 to 2.9 times the speed of sound when one minimum rectangular section of the jet 1.31 by 1.74 feet in size was used. In tests at subsonic speeds the test section of the jet was kept constant at these dimensions.

The jet was enclosed between two straight, parallel side walls, which were perpendicular to the axis of the model. The jet was not restrained by top and bottom walls. (See fig. 1.) The effuser A-A was shaped in such a way as to give a uniform flow at the plane a-a. This uniform flow was attained in a series of preliminary tests by increasing the length of the parallel-sided effuser until satisfactory flow distribution was obtained.

The diffuser B-B was placed in a position to give uniform flow and to eliminate the vibrations that tended to occur. With the diffuser shape and location finally determined, the velocity was constant along the plane b-b in the test section of the tunnel even at the highest speeds. By varying the position and the dimensions of the diffuser a stable and uniform flow could be obtained even in the Mach number range approaching and exceeding the speed of sound (Mach numbers of 0.9 to 1.2). The present test program included measurements made at Mach numbers up to 0.94. Information on the shape and location of the diffuser has been lost; therefore, the exact dimensions of this setup are not available.

The velocity and the Mach number were determined from a tunnel calibration based on measurements of the total pressure in the large section of the tunnel ahead of the entrance cone and on measurements of the static pressure at the wall near the exit of the entrance cone. In order to check the velocity measured in this manner, pitot-static tubes were installed at the top and bottom of the jet just downstream of the exit of the entrance cone. These tubes gave a qualitative indication of the jet-boundary interference effects. When the velocities measured by these tubes were appreciably different from the velocity indicated by the entrance-cone pressure.
ealibration, it was usually found that the interference effects were so large that they appreciably altered the aerodynamic characteristics of the test models. No data were taken when this condition existed.

**EFFECT OF REYNOLDS NUMBER AND AIR-STREAM BOUNDARIES**

Experimental methods. - In the study of Reynolds number effects at high speed, preliminary tests were made first on cylinders and spheres of various dimensions (reference 2). An analogous series of preliminary tests was then made for airfoils. Models of airfoils of constant profile but of varying chord were tested.

For the study of the effect of the air-stream boundaries, tests were made with varying ratio of model chord to tunnel height over a range of Mach numbers. The ratios used were: 0.0755, 0.0952, 0.113, and 0.151. The Reynolds number at each Mach number was held approximately constant by varying the density.

Test models. - A profile was chosen having an arc for the upper surface and a straight line for the lower surface because this profile could be exactly reproduced in various sizes. The upper surface could be made by use of a lathe and the lower surface could be formed by use of a shaper. The leading edge and the trailing edge were sharp. The maximum thickness chosen was 8 percent, and the profile was designated C-8 (fig. 2). Four models were constructed with such a profile; three with chords of 1.575, 1.969, and 2.362 inches (4, 5, and 6 cm) for force tests and one with a chord of 3.15 inches (8 cm) for determining the pressure distribution along the profile.

Tests and results. - At Mach numbers of 0.4, 0.5, 0.6, 0.7, 0.8, and 0.9, the lift coefficient, the drag coefficient, and the pitching-moment coefficient about the quarter-chord point of the airfoil were determined for the three profiles having chords of 1.575, 1.969, and 2.362 inches. All the models were tested at two Reynolds numbers: approximately 250,000 and 480,000. The model with the 1.575-inch chord was also tested at a Reynolds number of 150,000. For the profile having a chord of 3.15 inches, pressure readings were made at angles of attack between -3.5° and 4.5° for Mach numbers
of approximately 0.7, 0.8, and 0.9. Values of lift and of pitching moment were obtained from the pressure distributions. Force-test results are shown in figures 3 to 8. In figures 9 to 11 the results of pressure measurements are presented. Figure 12 shows the results obtained from integration of the pressure diagrams compared with the results obtained by use of the balance.

Reynolds number effects.- The results of the preliminary tests of cylinders and spheres showed that for the range of Reynolds numbers covered in the tests the effect of Reynolds number decreased as the velocity increased. At Mach numbers close to 1.0 there was virtually no Reynolds number effect. In the airfoil tests the importance of Reynolds number was considerable at low Mach numbers and the effect of Reynolds number was noted up to the critical Mach numbers at which the phenomenon of shock began to appear (figs. 3 to 6). For supercritical Mach numbers, the effect of Reynolds number became less until it virtually disappeared for Mach numbers very near 1.0. In this range the formation of shock waves seems to control the aerodynamic phenomena and the development of the boundary layer. The boundary-layer thickness probably depends to a large extent on the angle of deviation of the air as it passes through the shock wave. The friction drag is a reduced part of the total drag and, therefore, the Reynolds number effect is small. The Reynolds number, however, could have an effect on the characteristics of the shock wave itself through its action on the boundary layer, but such an effect is not indicated. In general, these airfoil test results confirmed the results of the sphere tests. Large-scale tunnel tests made at the Deutsche Versuchsanstalt fur Luftfahrt (the DVL) in Germany and flight tests made at various times showed similar results.

Effect of air-stream boundaries.- The jet-boundary effects for the ratios of chord to jet height of 0.0755 to 0.113 covered in these tests appear to be negligible. Essentially equivalent results were obtained at a given Reynolds number for all values of the ratios employed in the tests. For a larger jet-boundary effect, a test was made of the model with a chord of 3.15 inches for which the ratio of the chord of the model to the height of the air stream (0.151) is twice that normally used in the tests. From the results of integration (fig. 12) the values obtained for $C_L$ and $C_{mC/4}$ are seen to
coincide at high Mach numbers with the values found by the force tests. This agreement indicates that the boundaries of the air stream probably did not interfere appreciably with the distribution of the pressures. For a Mach number of 0.94 the effect of the air-stream boundaries is important for the model of 3.15-inch chord but is not important for the models of 1.575- and 1.969-inch chord. For higher Mach numbers the boundaries also affected the results obtained with the two smaller models.

It is interesting to note that the phenomenon of choking of the air stream, which occurs in closed-throat wind tunnels at high speeds (reference 3), did not occur in the tunnel in which the present tests were made. For example, for model C-8, which had a chord of 3.15 inches, it is estimated that choking in a closed-throat tunnel would occur at a Mach number of 0.88 or lower. The choking Mach number for the 2.352-inch-chord model is estimated to be 0.90 or lower. These choking Mach numbers were calculated from one-dimensional theory for the zero-lift condition. They are therefore somewhat higher than the choking Mach numbers that would actually be obtained, especially for angles of attack other than that for zero lift. In the present tests it was possible to obtain data for these models at Mach numbers as high as 0.94, and the results of the jet-boundary-effect tests indicate that the data are essentially free from tunnel-wall effects at this Mach number.

EFFECT OF HUMIDITY AND CONDENSATION

The air becomes very cold in the expansion that occurs in the tunnel at high speeds. (The process is very nearly adiabatic.) Total condensation may occur in the whole jet at high speeds if the dew point is passed. Even if condensation does not occur in the jet, there is a possibility of its occurring in the low pressure regions over the test model where an additional expansion and temperature drop occur. Very low local temperatures, which are usually smaller than the local dew point, are found at high subsonic speeds; local condensation therefore could occur and could produce a "condensation shock" or a localized region in which condensation occurs.
Condensation complicates and modifies the flow over the body because it alters the values of the temperature, the pressure, and the speed in the air stream and, hence, modifies the values of the resultant aerodynamic forces. A complete examination of the effects of the phenomenon of condensation shock is very complicated. The variables involved include the value of the local humidity, the speed of the condensation, the possibility of the existence of supersaturated air, and the scale of the model.

The condensation process is not instantaneous but requires a finite time and its beginning may depend on such factors as the nuclei of condensation. (The supersaturated air may sometimes exist for a time at a temperature much lower than the critical.) If the tests are made at small scale, the air can pass through the low temperature region in so short a time that appreciable condensation does not occur. Condensation is therefore less likely to occur in small-scale tests than in large-scale tests. In flight, for example, when appreciable relative humidity is present, condensation normally occurs and is easily seen on propellers and wings in high-speed dives. Since the characteristics of the condensation vary with scale, it would appear to be practically impossible to simulate full-scale conditions in tests in which small models are employed. The problem is further complicated because the degrees of supersaturation existing in the tests in a wind tunnel may be different from in flight and the beginning of the condensation depends on certain variable conditions of the air. The condensation characteristics of different wind tunnels, even with the same setup, have in several instances been noted to be widely different. In the subsonic tunnel of the Aerodynamische Versuchsanstalt (the AVA) at Göttingen, for example, it is normally necessary to dry the air before it converges in the test section to prevent condensation; however, in the Langley 24-inch high-speed tunnel, which has a comparable entrance-cone shape and which operates under similar conditions, it is not necessary to dry the air, and complete condensation seldom occurs for relative humidities below 60 percent.

All the test data obtained up to the present time tend to indicate that even for large-scale models the effects of humidity are of secondary importance provided that the percentage of humidity is low. In the Guidonia high-speed tunnel previously described, it was very difficult to study humidity effects because of an automatic
drying up of the air which took place. A small quantity of water was removed from the tunnel air by the pump which was used to evacuate the tunnel to the low initial pressure. The condensation that occurred when the tunnel was started was believed to cause water to collect behind the test section and to adhere to the tunnel walls. As a result of this automatic water removal, fog did not occur in the test section even at supersonic velocities and no air-drying equipment was necessary. Because the humidity became less during the progress of a test in this tunnel, it was impossible to give precise results as to the effect of humidity, but the general indication of the data that have been obtained was that the humidity effects were not appreciable, at least not for the small-scale models tested.

Tests to study the effects of humidity have been conducted in the 8.86-foot high-speed tunnel of the DVL in Germany using an NACA 0015-64 airfoil section with a 1.6-foot chord. In this wind tunnel the amount of condensation existing in the test section can be controlled by varying the cooling of the tunnel and thus regulating the temperature of the air in the test section. For very high values of relative humidity, it is necessary to eliminate the cooling entirely in order to raise the temperature enough to avoid condensation. The results of the humidity-effect investigation in the DVL tunnel demonstrated that, even for the relatively large-scale model employed, the humidity effects were of secondary importance when the relative humidity was small.

In order to indicate the conditions under which condensation might occur in flight, figure 13 is presented showing the local Mach number as a function of the flight Mach number for which the conditions required for saturation are reached. (Adiabatic expansion of the air from its static condition to the conditions corresponding to local Mach number is assumed.) Also shown in figure 13 are the values of maximum local Mach number that are attained locally on two typical airfoils. The data calculated for the NACA 23015 airfoil (unpublished) were obtained from tests made in the Langley 24-inch high-speed tunnel. The data for the NACA 0015-64 airfoil were obtained from the DVL tests mentioned previously. Figure 13 indicates that, even for very low values of the relative humidity, local Mach numbers are obtained at which condensation is possible when the flight Mach number is 0.6 or greater.
The discussion in the preceding paragraphs has shown that humidity effects are likely to be most pronounced under large-scale conditions. Systematic tests to determine humidity effects could best be made in a large-scale wind tunnel in which the temperature of the circulating air could be varied by regulating the cooling. The tests in such a wind tunnel could be made at various periods in order to cover a wide range of relative humidities. Figure 14 has been prepared to indicate the conditions for saturation in the test section of a wind tunnel for three values of relative humidity and for various temperatures of the air in the entrance cone of the wind tunnel where the airspeed is low. Also shown in figure 14 is a comparison of the maximum local Mach numbers of the NACA 23015 and 0015-64 airfoils as functions of the stream Mach number to determine at what Mach number the conditions for saturation are locally reached. The figure shows that, for high relative humidity, it is necessary to have a high temperature of the tunnel air stream in order to eliminate condensation in the test section. It is also shown that, even if condensation is eliminated in the test section, the necessary conditions for the formation of local condensation over the test model will normally be attained.

COMPARISON OF TEST RESULTS FROM VARIOUS WIND TUNNELS AND FROM FLIGHT

Airfoil tests.- For a thorough examination of the accuracy and significance of the test results obtained in a given wind tunnel, it is essential that the results be compared with those obtained in other wind tunnels and in free flight on models of similar profile. As a step in this direction, tests were conducted on the NACA 0015-64 airfoil in both the Guidonia 1.31- by 1.74-foot rectangular high-speed tunnel and in the DVL 0.85-foot-diameter high-speed tunnel, which has closed circular walls. The model used had a rectangular plan form enclosed between two end plates. The chord of the model was 1.638 feet (50 cm), the span was 4.5 feet, and the end plates were 23.6 by 13.2 inches. The ratio of the model chord to the tunnel diameter was 0.185. With this setup, the choking Mach number was about 0.86, which is considerably higher than the choking Mach number that would have been obtained with the model completely
spanning the tunnel jet. The data obtained in these tests consisted of pressure distributions and wake surveys.

The test conditions were adjusted to produce an equivalent relative humidity of the air of 20 percent at sea level. The Reynolds number varied with the Mach number from about 5,800,000 to 6,400,000 in the high-speed range of the tests.

The model tested at Guidonia had the same profile but was of much smaller scale, the model chord being 1.575 inches (4 cm) and the ratio of model chord to tunnel depth being 0.0755. The relative humidity in the Guidonia tests was always very low. The Reynolds numbers were, of course, very much lower than those of the DVL tests and varied around a value of about 500,000. Force measurements of lift, drag, and moment were made in the Guidonia tests; pressure-distribution and wake-drag measurements were made in the DVL tests. The results obtained are compared in figures 15 to 17.

Figure 18 shows pressure-distribution measurements made at the DVL for one angle of attack, \( \alpha = -0.25^\circ \). It may be noted that the results from the two tunnels are at variance, especially at high speeds. This lack of agreement indicates that the testing technique and the proportions of the testing system are of great importance in high-speed wind-tunnel work.

The differences in the drag-coefficient values at low Mach numbers are probably due to the difference in Reynolds numbers. The largest differences between the results from the two tunnels are in the drag and pitching-moment coefficients at high Mach numbers. The abrupt changes in the coefficients from the DVL tests at Mach numbers in the vicinity of 0.8 are probably associated with the phenomenon of choking, and the results obtained in this range are therefore considered extremely questionable. Because of the much smaller relative size of the model in the Guidonia tests and also because of the fact that the jet was not restrained by top and bottom walls, similar effects did not occur. Further tests were made at the DVL tunnel of a smaller model of the same profile having a chord of 1.148 feet, the model-chord to tunnel-diameter ratio being 0.13. The results obtained with the smaller model are shown in figure 16. It will be noted that the rate of drag rise past the critical speed is appreciably less than with the larger model and
thus is in better agreement with the results of the Guidonia tests.

Free-flight tests were in the general research program at Guidonia, but they were interrupted by the war. The few flight tests made, however, indicated that the drag-coefficient curves had about the same slopes at supercritical speeds as were obtained in the Guidonia wind-tunnel tests.

Bomb tests.- Additional comparisons between high-speed wind-tunnel and flight data were obtained in tests of an airplane bomb of conventional shape. The approximate shape of the bomb is indicated in figures 19 and 20, which show the results of the tests. The bomb was launched in flight at an altitude of 39,300 feet, and its trajectory as a function of time was recorded with a phototheodolite. The speed, the Mach number, the acceleration, and the drag coefficient were obtained from the trajectory data. A one-third scale model of this bomb was tested in the DVL 8.86-foot-diameter high-speed tunnel (the ratio of bomb diameter to tunnel diameter was 0.0455, much lower than that normally used). A one-tenth scale model of the same bomb was tested in the Guidonia 1.31- by 1.74-foot rectangular high-speed tunnel using a ratio of model diameter to air-stream height of about 0.0714. Similar tests were made in a wind tunnel at the AVA in Göttingen, which has a partly free airstream similar to that at Guidonia but 47 inches high. The size of the model used in these tests is not known, but it is believed that the ratio of model diameter to tunnel air-stream height was considerably higher than that used in the tests in the other wind tunnels. The results shown in figure 19 indicate reasonably good agreement in the form of the drag curves obtained. As might be expected, however, the drag-coefficient values obtained at very high Mach numbers in the closed DVL tunnel are higher than those found at Guidonia in the relatively unrestricted jet. The results obtained in a subsequent launching of the bomb, with reinforcements to the tail structure, in flight tests at the DVL are shown in figure 20.
CONCLUSIONS

The following conclusions were drawn from the investigation of the effects of Reynolds number, air-stream boundaries, and humidity in tests of airfoils at high speeds:

1. It has been shown that the ratio of tunnel height to model size, the form of the test section, and the testing technique have a very great bearing on the results obtained at subsonic Mach numbers above 0.7.

2. Reynolds number effects were of secondary importance at very high Mach numbers for the range investigated.

3. In the absence of suitable correction factors, the only safe experimental technique consists in keeping the scale of the model small enough so that the corrections required are negligible.

4. In a closed air stream, the model must be small enough that the highest desired test Mach number is below the choking Mach number of the tunnel, at which the effects of the tunnel walls on the flow over the model become extremely large.

5. By use of a jet which is not restrained by top and bottom walls the maximum Mach number that can be used for a given value of the ratio of jet height to model chord is appreciably higher than the value that can be obtained in a closed jet.

6. The considerations of condensation phenomena that have been discussed have brought out the fact that the conditions under which condensation occurs depend on many variables and that only with great difficulty could flight conditions be simulated in wind-tunnel tests in which small-scale models are used. Wind-tunnel tests should be conducted with low values of the relative humidity, because under such conditions the effects of condensation are known to be negligibly small.
II. TEST RESULTS FOR 24 AIRFOILS IN THE MACH NUMBER RANGE OF 0.40 TO 0.94

APPARATUS AND METHODS

Twenty-four profiles were tested in the 1.51 by 1.71-foot high-speed tunnel at Caldonia with the partly free test section previously described. For every profile the lift, the drag, and the pitching moment at the quarter-chord point were measured by use of the three-component semiautomatic balance described in reference 4. The extremities of the models were fixed at the balance supports, and the models were checked during the tests to verify that the aerodynamic loads did not bend them appreciably. All the tests were repeated with the model inverted. For some models, the tests were repeated later when the static atmospheric conditions were completely different and with different humidities in the test section. (The values of the relative humidity were always low.) The differences in the results obtained were not appreciable.

All the models were made of well-polished steel and had chords of 1.575 inches for thickness ratios of 3 percent or greater. In order to prevent excessive bending, the models with thickness ratios of less than 3 percent had chords of 1.069 inches.

The profiles of small models seldom correspond exactly with the profile desired. For purposes of accuracy, therefore, an optical device was constructed that permitted photographing with extreme precision the true section of each model on a greatly increased scale. For each model two end sections were photographed and the true profile was projected on the photograph to provide the desired comparison. Because the airfoils were constructed by machine, the profile shape did not vary across the model span. This fact was confirmed by superimposing drawings of the two end sections. It was verified that the surface was adequately smooth by observing the tangential illuminated surface under great magnification.

Figure 21 shows the specified shapes of the profiles tested. In figure 22 the actual shapes of the profiles tested are compared with the specified shapes. In order
that the difference between the actual and the specified profiles may be clearly seen, the ordinate scale used in figure 22 has been enlarged. Table I shows the ordinates of the profiles tested.

All the tests were performed at an approximately constant Reynolds number varying in the range from $30,000$ to $420,000$. The density, and consequently the Reynolds number, had to be kept low for the thinner airfoils in order to prevent excessive loads.

AIRFOILS TESTED

The profiles listed in the following table were tested:

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<td>(2 percent thick)</td>
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<td>(10 percent thick)</td>
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<td>NACA M1</td>
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*Developed at Zurich University.*
RESULTS

In figures 23 to 46 the results of the tests of the 24 airfoils are shown in the form of the usual coefficients: $C_L$ and $C_D$ are plotted against the test Mach number at the same angles of attack, and $C_{m_c}/4$ is plotted against the Mach number at values of $C_L$ corresponding to the given angles of attack. Figures 47 to 70 show $\alpha$, $C_D$, and $C_{m_c}/4$ plotted against the corresponding $C_L$ for each airfoil at the same Mach numbers. In figure 71 the angle of zero lift is plotted against Mach number for representative airfoils of the group. Figure 72 gives the maximum lift-drag ratio $(L/D)_{max}$, $\alpha$ for $(L/D)_{max}$, and $C_L$ for $(L/D)_{max}$ as functions of Mach number for all the airfoils tested. Figures 73 and 74 present $C_{D_{min}}$ and $(L/D)_{max}$ as functions of the maximum percentage thickness for all the airfoils at various Mach numbers, and figures 75 and 76 show $C_{D_{min}}$ and $(L/D)_{max}$ plotted against Mach number for several groups of airfoils having the same maximum thickness.

It can be observed from the test results that:

The lift-coefficient curve as a function of Mach number presents a maximum and later a minimum value. The Mach numbers at these values can be defined as the first and the second critical Mach numbers for the lift. The Mach number at which the drag-coefficient curve abruptly bends upward is defined herein as the critical Mach number for the drag. It will be noted that the critical Mach numbers as defined herein are different for the lift and for the drag data. The critical Mach numbers used, furthermore, do not necessarily correspond to the stream Mach number at which local sonic velocity is reached.

The rate of drag rise past the critical Mach number increases as the lift coefficient, the angle of attack, and the thickness ratio are increased.

The first critical Mach number for $C_L$ and the critical Mach number for $C_D$ for each airfoil is lowered with the increase in angle of attack.
For each series of airfoils at the same angles of attack, these critical Mach numbers decrease as the thickness increases.

The critical Mach numbers at the same thickness and the same angle of attack are much lower for the cambered profiles than for the symmetrical profiles tested at the same angle of attack. At equal thickness and equal camber, the critical Mach numbers are higher where the maximum thickness was at the 40-percent-chord station than where it was at the 30-percent-chord station.

Above the critical Mach numbers, the drag increases and the lift decreases very rapidly; for a profile with a larger thickness and sharper curvature, the increase in drag and the decrease in lift is sharper.

These general phenomena agree with results of other laboratories. (See, for example, reference 4.)

Lift.- At subsonic Mach numbers the increase in lift coefficient with Mach number follows approximately the theoretical relation \( \frac{1}{\sqrt{1 - M^2}} \), especially for the low thickness ratios. After the first critical Mach number is reached, the lift coefficient decreases very rapidly until it reaches a minimum at the second critical Mach number when it again starts to increase. This second critical Mach number is lowered with the decrease in the first critical Mach number. Airfoils with larger camber had greater decreases in lift. For these airfoils, generally, the angle of attack for zero lift changed greatly and tended toward positive values (fig. 71). At the highest test Mach numbers all the wings functioned in a manner very similar to symmetrical profiles. This phenomenon agrees with the fact that the value of the angle of zero lift for an unsymmetrical profile changes sign and becomes considerably reduced in magnitude in passing from a subsonic to a supersonic velocity.

The lift-curve slope \( \frac{dC_L}{da} \) increases up to the first critical Mach number after which there is a considerable decrease up to the second critical Mach number. (See figs. 17 to 70.) The second critical Mach number is greatly affected by the value of the maximum percentage thickness. For the lower thicknesses tested,
the second critical Mach number was reached only at the maximum speed of the test.

Moment. - The curve of pitching-moment coefficient against Mach number has a fairly regular form (figs. 33 to 46). Generally, the value of $C_{m_c/4}$ remains constant up to the first critical Mach number and then tends to decrease. For the larger thickness ratios there is an increase in $C_{m_c/4}$ at the first critical Mach number, and it appears that the center of pressure moves forward. When the Mach number is increased beyond the critical value, $C_{m_c/4}$ decreases until it reaches a minimum and then tends to increase.

The center of pressure moves appreciably at low values of $C_L$ for profiles of large camber. For symmetrical profiles the variations of $C_{m_c/4}$ at low values of $C_L$ are small, especially if the maximum thickness is about 40 percent of the chord.

Drag and lift-drag ratio. - The value of $(L/D)_{max}$ decreased rapidly beyond the first critical Mach number for the lift and continued to decrease until the second critical Mach number was reached. It then varied very slowly with further increase in Mach number (fig. 72). The larger thicknesses suffer more pronounced relative changes in $(L/D)_{max}$. The angle of attack and the lift coefficient corresponding to the $(L/D)_{max}$ (fig. 72) decrease as far as the second critical point and then begin to increase rapidly. The variations are appreciably influenced by the value of maximum thickness ratio and by the mean camber.

In order to emphasize the importance of maximum percentage thickness on the values of $C_{D_{min}}$ and $(L/D)_{max}$, figures 73 and 74 were prepared to show the values of these factors as a function of maximum percentage thickness at constant Mach number for each series considered. These figures show that the effect of the maximum percentage thickness becomes greater as the Mach number increases for all profiles tested. For Mach numbers around 0.8, the effects of thickness ratio are very large. At low test speeds, for example, when the maximum thickness is
varied from 6 to 12 percent, $C_{D_{\text{min}}}$ increases about 30 percent; at Mach numbers around 0.8 or greater, the increase becomes 200 to 400 percent. The ratio of the values of $(L/D)_{\text{max}}$ for the NACA 0006-34 airfoil to those for the NACA 0012-34 airfoil changes from about 1.18 at a Mach number of 0.65 to 4.5 at a Mach number of 0.85 and to 2.2 at a Mach number of 0.94. It is also interesting to compare the aerodynamic characteristics of various profiles at equal maximum percentage thicknesses. (See figs. 75 and 76.)

The larger increments in $C_{D_{\text{min}}}$ occur for the profiles with larger camber, for which the critical Mach number is lower. With increase in the value of the maximum percentage thickness, the value of $dC_{D_{\text{min}}}/dM$ increases and even at very high Mach numbers this difference between various profiles is considerable.

The profile shape has considerable effect on $(L/D)_{\text{max}}$ (fig. 76); the unsymmetrical profiles have larger $(L/D)_{\text{max}}$ values at low Mach numbers. At higher speeds, the symmetrical profiles with the maximum thickness at about 40 percent of the chord had higher efficiencies than those with the maximum thickness at 30 percent of the chord. The difference between the various profile types is considerable for low Mach numbers; however, it decreases with increase in Mach number and is small for Mach numbers around 0.94.

CONCLUSIONS

The following conclusions may be drawn from the results obtained from tests of 24 small-scale airfoils in the Guidonia high-speed tunnel:

1. At subsonic Mach numbers both the profile shape and the thickness ratio had a large effect on the minimum drag coefficient.

2. Reducing the thickness ratio, moving the point of maximum thickness from 30 to 40 percent of the chord, and reducing the camber all tended to increase the critical Mach number.
3. Airfoils of large percentage thickness should not be used at high Mach numbers because of the radical adverse changes in their characteristics at supercritical speeds.

4. When the critical speed was exceeded, the drag coefficients increased rapidly. Abrupt decreases in lift and changes in moment occurred at somewhat higher critical Mach numbers.

5. The lift coefficient continued to decrease as the speed was advanced beyond the first critical Mach number until a second critical Mach number was reached, beyond which the lift coefficient increased in value.

6. At very high supercritical Mach numbers the thickness ratio is the dominating variable, the drag coefficient being almost directly proportional to the thickness at a Mach number of 0.94.

National Advisory Committee for Aeronautics
Langley Memorial Aeronautical Laboratory
Langley Field, Va.
REFERENCES


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Figure 2. Profile of C-8 airfoil section.

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Relative humidity in Tunnel entrance cone.

(Percent)

Mach number for dewpoint test section, $M_g$

Maximum local Mach number, $M$

Temperature, $T$, °F

Test Mach number, $M$
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- Actual airfoil
- Specified airfoil

(a) USNPS 1, 4 percent thick

(b) USNPS 2, 8 percent thick

(c) USNPS 3, 10 percent thick

(d) NACA 0006T

(e) NACA 2509

(f) Davis, 9 percent

(g) NACA M 1

(h) ETH 3609

Figure 22 - Concluded.
Figure 23. Effect of compressibility on the aerodynamic characteristics of the NACA 6006-64 airfoil.
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Figure 25. - Concluded.
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Figure 30. - Continued.
Fig. 30 Conc.

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Figure 30 — Concluded.
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Figure 32.—Continued.
Figure 32. Concluded.
Figure 33.—Effect of compressibility on the aerodynamic characteristics of the NACA 2309 airfoil.
Figure 33.—Continued.
Figure 33. - Concluded.
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Figure 34.—Concluded.
Figure 35. - Effect of compressibility on the aerodynamic characteristics of the NACA 2315 airfoil.
Figure 35—Continued.
Figure 35.—Concluded.
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Figure 36. Concluded.
Figure 37. Effect of compressibility on the aerodynamic characteristics of the NACA 2409 airfoil.
Figure 37 - Continued.
Figure 37. -- Concluded.
Figure 38: Effect of compressibility on the aerodynamic characteristics of the NACA 2412 airfoil.
Figure 38 - Continued.
Figure 38 - Concluded.
Figure 39. - Effect of compressibility on the aerodynamic characteristics of the USN.R5.1 airfoil, 4 percent thick.
Figure 39. - Continued.
Figure 39. - Concluded.
Figure 40. - Effect of compressibility on the aerodynamic characteristics of the U.S.N.F.S. 2 airfoil, 8 percent thick.
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Figure 42. - Continued.
Figure 42, - Concluded
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Figure 60.—Aerodynamic characteristics of the NACA 2406 airfoil.
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Figure 63.— Aerodynamic characteristics of the USNR-5.1 airfoil.
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Figure 65.—Aerodynamic characteristics of the USNPS. 3 airfoil.
Figure 66. - Aerodynamic characteristics of the M1 airfoil.
Figure 67. - Aerodynamic characteristics of the NACA 0006T airfoil.
Figure 68. Aerodynamic characteristics of the NACA 2509 airfoil.
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Figure 71: Variation of angle of zero lift with Mach number.
Figure 71e-g

(e) NACA 2406 airfoil.

(f) NACA 2409 airfoil.

(g) NACA 2312 airfoil.

Mach number, M
(h) USNPR. 1 airfoil, 4 percent thick.

(i) USNPR. 2 airfoil, 8 percent thick.

(j) USNPR. 3 airfoil, 10 percent thick.

Figure 71. - Continued.
Figure 71k-m

(k) NACA 2509 airfoil.

(l) Davis 8-percent-thick airfoil.

(m) ETH 3609 airfoil.

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Figure 71.- Concluded.
Figure 72a: Variation of $(L/D)_{max}$ and of the lift coefficient and angle of attack corresponding to $(L/D)_{max}$ with Mach number.
Figure 72b - Continued.

Airfoil
NACA 0006-34
NACA 0008-34
NACA 0012-34

Mach number, $\mathbf{M}$

(a) NACA 0000-34 series airfoils.

(b) NACA 0000-34 series airfoils.
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(c) NACA 0006-63 series and M and T airfoils.
Figure 72c - Continued.
Fig. 72d

(d) NACA 2300 series airfoils.

Figure 72. - Continued.
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Fig. 72e

\[ \alpha \text{ for } \left( \frac{b}{b_{\text{max}}} \right)_\text{max} \]

\[ \alpha \text{ for } \left( \frac{b}{b_{\text{max}}} \right)_\text{max} \]

Airfoil
- NACA 2406
- NACA 2409
- NACA 2412
- NACA 2509
- Davis 9 percent
- ETH 3609

\[ (\frac{b}{b_{\text{max}}}) \text{ max} \]

Mach number, M

Extrapolated NACA 2409

(e) NACA 24 and 25 series airfoils.

Figure 72 - Continued.
Figure 72f

Airfoil
USNPS. 1, 4 percent thick
USNPS. 2, 6 percent thick
USNPS. 3, 10 percent thick

\( C_r \) for \( (\frac{L}{D})_{\text{max}} \)

\( \alpha \) for \( (\frac{L}{D})_{\text{max}} \)

Mach number, \( M \)

\((\frac{L}{D})_{\text{max}}\)

USNPS. series airfoils.
Figure 73a - Variation of minimum drag coefficient with maximum thickness for various Mach numbers.
(b) NACA 0000-34 series airfoils.
Figure 73. - Continued.
Minimum drag coefficient, $C_{\text{min}}$

Maximum thickness, percent chord

(C) NACA 2300 series airfoils. Figure 73. – Continued.
(d) NACA 2400 series airfoils.
Figure 73. – Continued.
(E) U.S.N.P.S. series airfoils.
Figure 73.—Concluded.
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(b) NACA 0000-34 series airfoils. Figure 74b - Continued.
Figure 74c

(C) NACA 2300 series airfoils.
Figure 74c - Continued.
(d) NACA 2400 series airfoils. Figure 74. – Continued.
(e) U.S.N.P.S. series airfoils.
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(b) Maximum thickness, 9 percent.

Figure 75, - Continued.
Figure 75c

(c) Maximum thickness, 12 percent.

Figure 75. - Concluded.
Figure 76a - Variation of $(L/D)_{max}$ with Mach number for several airfoils having the same maximum thickness.
(b) Maximum thickness, 9 percent.

Figure 76.—Continued.
Figure 76c

Airfoil

NACA 2412
NACA 2312
NACA 0012-64
NACA 0012-34

(c) Maximum thickness, 12 percent.

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Figure 76c - Concluded.