RESEARCH MEMORANDUM

AN INVESTIGATION OF THE LOW-SPEED STABILITY AND CONTROL CHARACTERISTICS OF SWEPT-FORWARD AND SWEPT-BACK WINGS IN THE AMES 40- BY 80-FOOT WIND TUNNEL

By

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An investigation has been made at large scale of the characteristics of highly swept wings. Data were obtained at several angles of sideslip on wings having angles of sweep of $\pm 45^\circ$, $\pm 30^\circ$, and $0^\circ$. The airfoil sections of the wings varied from approximately NACA 0015 at the root to NACA 23009 at the tip. Each wing was investigated with flaps undeflected, partial-span split flaps deflected 60°, full-span split flaps deflected 60° and split-flap-type ailerons deflected $\pm 15^\circ$. Values of maximum lift were obtained at Reynolds numbers ranging from 5.7 to $9.2 \times 10^6$. In this report the summarized results are compared with the predictions made by use of the simplified theory for the effect of sweep and with existing small-scale data. The basic wind-tunnel results from which these summary data were taken are included in an appendix.

The primary problems accompanying the use of sweep as revealed by this investigation are the loss in maximum lift, the high effective dihedral, and the sharp reduction in lateral-control effectiveness. In general, simple theory enables good predictions to be made of the gross effects of sweep but further
refinements are necessary to obtain the accuracy required for design purposes. In cases where comparisons can be made, the indications are that, as sweep increases, scale effects diminish and large-scale results approach small-scale results.

INTRODUCTION

Theory indicates and experiment has shown that the prime aerodynamic effect of wing sweep is to reduce by the cosine of the angle of sweep the effective flight velocity experienced by the airfoil sections of the wing. This then enables increases in maximum flight speed to be attained before serious compressibility effects are encountered. Theory and experiment also show that wing sweep introduces a number of stability and control problems, the seriousness of which becomes accentuated at low flight speeds.

Small-scale tests have pointed out the general nature of these problems and indicated those which must be overcome if the high-speed benefits of sweep are to be realized. They have also suggested that boundary-layer flow and, hence, Reynolds number has a profound influence on measured characteristics and that the value of small-scale tests remain somewhat doubtful until the extent of this influence is understood.

Since no large-scale data were available for wings with large angles of sweep, an investigation of the effects of sweep was conducted in the Ames 40- by 80-foot wind tunnel and the results are reported herein. It is believed that
these data will go far towards establishing the datum required to estimate the effects of scale on highly swept wing plan forms. With this knowledge at hand it is evident that the value of future small-scale tests will be considerably increased.

This report discusses a summary of the basic results and compares them with simple swept-wing theories and, where possible, with existing small-scale data (references 1, 2, and 3). To make the basic data available for further analyses they are included as an appendix to this report.

DESCRIPTION OF MODELS

The five models tested were composed of wing panels from an available airplane which were given the desired plan form and sweep by individually fabricated tips and center sections. The resulting angles of sweep were 0°, 30°, and 45° sweep forward, and 30°, and 45° sweepback (measured with reference to the quarter-chord line of the airfoil sections). Aside from the angle of sweep, the prime plan-form variable was considered to be aspect ratio. The tips and center sections were constructed to give the smallest variation of this parameter possible without modification of the airplane wing panels. No special attempt was made to control the variation of taper ratio, area or span. Photographs of the wings and plan-form drawings with pertinent dimensions are shown in figures 1 and 2. The geometric characteristics of the five wings tested are listed in table I.
The airfoil sections of the swept wings were dictated by the sections of the airplane wing panels (an NACA 0015 at the inboard end and an NACA 23009 at the outboard end of the panel). The profiles of the center sections and tips were simply extensions of the wing-panel airfoil. To expedite construction, three tips only were fabricated: one for the swept-forward wings, one for the straight wing, and one for the swept-back wings. Thus for the swept-forward and swept-back wings compromise tips were used which were misaligned $7.5^\circ$ to the air stream. The twist in the chord plane of the wing panels was approximately $1/4^\circ$ of washout. The dihedral of the chord-plane leading edge was kept at $90^\circ$.

No attempt was made to improve the fairness of the wing panels beyond the original manufacturing condition. Thus, due to presence of various access plates, panel joints, etc., the wings were rough to a greater degree than that normally associated with latest construction requirements.

Partial-span and full-span split flaps were tested on all models. The flaps were 0.20 chord and were deflected $60^\circ$.\(^1\) The span of the partial-span flaps was 0.623 wing span for all models; the span of the full-span flaps varied slightly from full-span (in no case more than 0.064 wing span) as shown in table I.

Ailerons were simulated by attaching the outboard portion of one of the flaps to the right wing and deflecting it $\pm 15^\circ$ (up-deflection was obtained by attaching the flap to the upper surface of the wing). Thus the ailerons as tested were

\(^1\)Except where noted, all chords and spans used in this report were measured parallel and perpendicular to the plane of symmetry. Flap deflection angles were measured in a plane perpendicular to the flap hinge line.
0.20-chord split-flap-type ailerons.

The wings were mounted on a faired sting which in turn was attached to the three-strut support system. Photographs of the wing installations are shown in figure 1.

COEFFICIENTS AND SYMBOLS

The data are presented in the form of standard NACA coefficients and symbols as defined in figure 3 and the following tabulation. All forces and moments are presented about the stability axes with their origin located on the root chord, or root chord projected and at the same fore and aft location as the quarter L.A.C.

- $C_L$ lift coefficient (lift/$qS$)
- $C_D$ drag coefficient (drag/$qS$)
- $C_m$ pitching-moment coefficient ($\frac{\text{pitching moment}}{qSc}$)
- $C_r$ rolling-moment coefficient ($\frac{\text{rolling moment}}{qSb}$)
- $C_n$ yawing moment coefficient ($\frac{\text{yawing moment}}{qSb}$)
- $C_Y$ side-force coefficient ($\frac{\text{side force}}{qS}$)
- $C_{L_{\alpha}}$ rate of change of lift coefficient with angle of attack, per degree
- $\Delta C_L$ increment of lift coefficient due to deflecting flaps
- $C_{L_{\text{max}}}$ maximum lift coefficient
- $C_{r_{\beta}}$ rate of change of rolling-moment coefficient with sideslip, per degree
- $C_{r_{\psi}}$ rate of change of rolling-moment coefficient with wing-tip helix angle, per radian
\( C_{n_{\beta}} \) rate of change of yawing-moment coefficient with sideslip, per degree

\( C_{l_{\delta_a}} \) rate of change of rolling-moment coefficient with aileron angle, per degree

\( \frac{\partial C_{l_{c}}}{\partial C_{L}} \) rate of change of \( C_{l_{\beta}} \) with lift coefficient where lift is increased by changing angle of attack

\( \frac{\Delta C_{l_{c}}}{\Delta C_{L}} \) rate of change of \( C_{l_{\beta}} \) with lift coefficient where lift is increased by deflecting flaps

\( \frac{\partial C_{n_{\beta}}}{\partial C_{L}^2} \) rate of change of \( C_{n_{\beta}} \) with the lift coefficient squared

\( L/D \) ratio of lift to drag

\( q \) dynamic pressure, pounds per square foot

\( V \) velocity along flight path, feet per second

\( \alpha \) angle of attack, degrees

\( \beta \) angle of sideslip, degrees

\( \Lambda \) angle of sweep of quarter chord line of airfoil sections, degrees (Sweepback is positive and sweepforward is negative.)

\( \Gamma_{e} \) effective dihedral, degrees

\( \delta \) control surface deflection, degrees

\( A^s \) aspect ratio based on span \((\frac{b^s}{S})\)

\( A^l \) aspect ratio based on length of quarter chord line \((\frac{b^l}{S \cos^2 \Lambda})\)

\( E \) Jones' edge-velocity correction

\( \lambda \) taper ratio, ratio of tip chord to root chord \((\frac{c_t}{c_R})\)

\( S \) wing area, square feet
TESTS AND RESULTS

For each of the model configurations six-component force and moment data were obtained through an angle-of-attack range at each of several angles of sideslip. The data were obtained at dynamic pressures which range from 5 to 75 pounds per square foot (\( R = 2.8 \times 10^6 \) to \( R = 15.0 \times 10^6 \))^2; most of the data were obtained at dynamic pressures of 10 to 20 pounds per square foot (\( R = 4.0 \times 10^6 \) and \( R = 5.3 \times 10^6 \), respectively).

The basic data obtained from the wind-tunnel tests of the five swept wings are described in the appendix. Also included in the appendix is a description of the corrections and tares applied to the data.

DISCUSSION

In this discussion an evaluation is made of the effect of wing sweep on the more important aerodynamic parameters and of the consequent effect on airplane performance and stability.

^2These Reynolds numbers are based upon the N.A.C. as a reference length and are the minimum and maximum limits of the variation including the change in chord length with sweep.
Also, the accuracy with which the simplified sweep theory may be used to predict the characteristics of swept wings is evaluated by a comparison with the experimental data. Finally, an attempt is made to compare at least qualitatively the values of the various characteristics as obtained at small-scale ($R < 1.5 \times 10^6$) and full-scale Reynolds numbers. The summary data on which this discussion is based have been extracted (for a test dynamic pressure of 20 lb/sq ft) from the measured characteristics included in the appendix.

The concepts advanced by Botz in reference 4 form the groundwork for the theory of the aerodynamic effects of incorporating sweep in a wing plan form. These concepts are based on the assumption that for an infinite-span wing only the velocity component normal to the quarter-chord line influences the pressures over a wing; the spanwise component of velocity is neglected. Thus, if the velocity components are resolved perpendicular and parallel to the quarter-chord line of a wing, the effective dynamic pressure over the wing will decrease in proportion to the square of the cosine of the angle of sweep and the effective angle of attack will increase in proportion to the reciprocal of the cosine of the angle of sweep. These changes in effective dynamic pressure and angle of attack brought about by wing sweep form the basis for the existing simplified sweep theory.

In interpreting the comparisons to be made between the simplified theory and experimental results, the limitations
of the simplified theory must be born in mind. Over the root section of highly swept finite-span wings, particularly highly tapered low aspect ratio wings, the basic assumption that the wing reacts only to air velocities normal to the quarter-chord line probably does not hold. It should also be noted that simplified theory in its present form applies only to wings which generate an additional loading due to angle-of-attack change that is rectangular in form. Therefore appreciable deviations from rectangular loading such as produced by taper will result in discrepancies between the theoretical and experimental results.

Lift Characteristics

Lift-curve slope.-- The simplified theory indicates a decrease in lift-curve slope proportional only to $\cos \Lambda$. To account for induction effects, a correction must also be made for any variations of aspect ratio. Hence, the effect of sweep on lift-curve slope, when corrected for aspect ratio, will be in accordance with the relation:

\[
(C_{L,\alpha})_{\Lambda} = (C_{L,\alpha})_{\Lambda=0} \cos \Lambda \quad \frac{[A/(\Lambda+2)]_{\Lambda}}{[A/(\Lambda+2)]_{\Lambda=0}}
\]

In conformity with standard nomenclature, aspect ratio is based on the span of the wings; however, there is some contention that since only air flow perpendicular to the quarter-chord
line is considered to affect the aerodynamic characteristics, aspect ratio should be based on the length of the quarter-chord line. Such an assumption is used in the analysis included in reference 1. In figure 4, the experimental results (taken from the linear portion of the lift curve) are shown together with the predictions based on theory for both concepts of aspect ratio. For swept-back wings, basing the aspect ratio on the length of the quarter-chord line gives the better agreement; whereas for swept-forward wings, basing the aspect ratio on the conventional span gives the better agreement.

It is believed that neither of these aspect ratio concepts gives a correct picture of the induction effects of the vortex pattern on swept wings. It can be shown that if a wing is swept back, the induction influences of the trailing vortices on the wing should be reduced, and conversely, if a wing is swept forward, the induction influences on the wing should be increased. That is, the effective aspect ratio increases with sweepback and decreases with sweepforward for wings of constant geometric aspect ratio (b^2/S).

The lift-curve slopes for the wings of this report have been estimated using the method of Falkner (reference 5) which

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3 It is recognized that a further aspect ratio correction, namely, Jones' edge-velocity correction should be used. The effect is small, however, compared to the errors resulting from the use of simple sweep theory. It has been omitted, therefore, in an effort to indicate clearly the adequacy of simple sweep theory in indicating the lift-curve slope of highly swept wings.
takes into consideration the induction effects of the swept vortex system in a more precise manner. In applying this method the section lift-curve slope for these wings 
\(c_l\alpha = 0.103\) was used rather than the theoretical value 
\(c_l\alpha = \frac{2\pi}{57.3}\) used in reference 5. The results are shown in figure 4. The predicted values of \(c_l\alpha\) and their variation with sweep closely approximate the experimental results. This indicates that, when induction effects are properly accounted for, accurate predictions of lift-curve slope can be made. It can be inferred then that the failure of simple theory to accurately indicate the effect of sweep on lift-curve slope is a result of improper induction effects.

Taper appears to have a strong effect on lift-curve slope due to its inherent influence on induction effects. As previously mentioned, the simplified theory strictly applies only to rectangular loading and hence the taper of the wings of the subject investigation may account for some of the discrepancy between theoretical and experimental results. In an attempt to correlate the effect of taper on the lift-curve slope of swept wings, data from previous investigations of swept wings having different taper ratios (references 1, 2, 3, 6, and 7) are shown in figure 5. For most of the investigations the wing aspect ratio (defined as \(b^2/S\)) and taper ratio did not vary with sweep. For those cases where aspect ratio \((b^2/S)\) varied with sweep, the data were corrected to the aspect ratio \((b^2/S)\) for the unswept wing.
Examination of the data in figure 5 will reveal that, as taper ratio is decreased, the maximum value of lift-curve slope occurs at greater angles of sweepback. The relation between taper ratio and the angle of sweep at which the maximum value of lift-curve slope occurs is shown in figure 6. The figure discloses that in order to obtain maximum lift-curve slope the taper ratio should be reduced from 1.0 as the wing is swept back and, by inference, increased from 1.0 as the wing is swept forward.

A comparison of figures 4 and 5 shows that values of lift-curve slope determined from large-scale tests show no better or poorer agreement with simple theory than values from small-scale tests. It appears that the principal disagreement between theory and experiment lies in failure of the theory to properly account for the induction effects on swept wings and that in comparison the effects of scale are relatively small.

Examination of the nonlinear portion of the lift curves and comparison with small-scale data shows that no consistent effects exist which could be attributed directly to scale effect. Those differences which do exist are small and erratic in nature and probably result from differences in plan form, wing section, and local wing roughness.

Flap effectiveness. According to simple sweep theory, flap effectiveness decreases as $\cos^2\alpha$. An additional correction to account for induction effects must be applied
when comparing flap effectiveness on wings of different aspect ratio. The comments previously made regarding the effects of induction on lift-curve slope apply equally well to flap effectiveness. In fact when $\alpha', C_{L\alpha}'$, and $\delta$ are measured perpendicular to the quarter-chord line, the effectiveness parameter $\alpha' \delta$ is unaffected by sweep and the lift increment produced by flap deflection is directly proportional to $C_{L\alpha}'$ or to $C_{L\alpha} \cos \alpha$. Hence the theoretical effect of sweep on flap lift increment may be written either in terms of sweep and aspect ratio:

$$\left( \Delta C_L \right)_\alpha = \left( \Delta C_L \right)_\lambda = 0 \cos^2 \alpha \frac{[A/(A+2)]_\lambda}{[A/(A+2)]_\lambda=0}$$

or in terms of sweep and lift-curve slopes:

$$\left( \Delta C_L \right)_\alpha = \left( \Delta C_L \right)_\lambda = 0 \cos \alpha \frac{C_{L\alpha}'_\lambda}{C_{L\alpha}'_\lambda=0}$$

where $\alpha$ is the angle of attack of the root chord.

In figure 7 the experimental results are shown together with predictions made in accordance with both the foregoing relations. It can now be seen that predictions of flap lift increment made in terms of aspect ratio deviate from experiment the same as did the predictions of lift-curve slope. However, when predictions are made in terms of lift-curve slope:

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*Note that in correcting for aspect ratio, the aspect ratio was based on the span. As in the case of $C_{L\alpha}$, if aspect ratios were based on the length of the quarter-chord line better agreement in flap lift increment would have been obtained for swept-back wings and poorer agreement for swept-forward wings.*
slopes, the agreement with experiment is almost exact. Thus the control which $C_{L_{\alpha}}$ has on flap effectiveness emphasizes the importance of fully understanding the effect on $C_{L_{\alpha}}$ of the many factors involved; this is especially significant when flap effectiveness is considered in terms of airplane control and performance.

Since the flap lift increment is dependent upon lift-curve slope, the conclusions concerning the effect of scale on lift-curve slope apply equally well to flap lift increment. In general, it can be said that sweep introduces no new scale effect on flap lift increments measured at low angles of attack.

**Maximum lift.**—The effect of sweep on maximum lift of the wing without flaps, with 0.623 span flaps deflected 60°, and with full-span split flaps deflected 60° is shown in figure 8. Attention is called to the fact that the wings tested were composed largely of production wing panels with normal roughness and irregularities such as caused by access plates. As a result of the roughness, maximum lifts measured on these wings may be somewhat lower than those measured on smooth wings. However, since the measured values on the unswept wing appear to be reasonably high for the particular airfoil sections, it is believed that the roughness was not sufficient to seriously reduce the maximum lift measured.

As shown in figure 8, sweep in wing plan form produces serious losses in maximum lift. However, for all but one of
the wing configurations the measured maximum lift was equal to or greater than simple theory would indicate, that is, \( C_{L_{\text{max}}} \) did not in general decrease proportional to \( \cos^2 \alpha \). Furthermore, the geometric angle of attack for \( C_{L_{\text{max}}} \) (fig. 9) does not decrease as

\[
\cos \alpha \frac{\left[ A/(A+2) \right]_A=0}{\left[ A/(A+2) \right]_A}
\]

which would be predicted by simple theory. It is probable that spanwise boundary-layer flow prevents stall from spreading from tip to root on the swept-back wing and from root to tip on the swept-forward wing. It is also possible that this intense boundary-layer drain allows certain sections of the wing to reach abnormally high angles of attack prior to stall.

On the unswept wing the gain in maximum lift coefficient due to flap deflection is equal to the flap lift increment at low angles of attack; whereas on the swept wings the gains in maximum lift coefficient are somewhat less than the flap lift increments realized at low angles of attack. This is particularly true for the outboard portion of full-span flaps on the swept-back wings and the inboard flaps on the swept-forward wings. (For sweepback angles greater than 30°, full-span flaps produce no greater \( C_{L_{\text{max}}} \) than do partial-span flaps.) Such decreases in flap effectiveness with sweep are disappointing but not surprising, since near stall the air flow is separated on the outboard section of swept-back wings and the inboard section of swept-forward wings.
The measured loss in $C_{l_{\text{max}}}$ due to sweep seriously limits airplane performance. With either partial or full-span flaps, the loss in $C_{l_{\text{max}}}$ due to $45^\circ$ of sweep would require increases in landing speed of approximately 20 percent.

Because of differences in taper ratio and airfoil sections of models used for small- and large-scale tests to date, no quantitative conclusions can be drawn regarding the effects of scale on $C_{l_{\text{max}}}$ at various angles of sweep. In general, however, it can be inferred that as the angle of sweep is increased the effects of scale become smaller. For instance a comparison of figure 6 of this report with figure 7 of reference 8 shows that an increase in $C_{l_{\text{max}}}$ of 0.25 is obtained at $0^\circ$ of sweep when going from small- to large-scale tests. In contrast, increases of only 0.10 and 0.08 in $C_{l_{\text{max}}}$ are gained with $30^\circ$ and $45^\circ$ sweepback. The increase in $C_{l_{\text{max}}}$ at $0^\circ$ of sweep is in general accord with what past experience has shown to be a reasonable effect of scale; whereas the increases for the swept wings fall far short of what would be anticipated from experience on straight wings. These data indicate that large-scale tests show a much more rapid decrease of $C_{l_{\text{max}}}$ with sweep than do model tests. This seems true whether flaps are deflected or not. Since large-scale results tend to approach small-scale results at large angles of sweep, considerable care should be taken in trying to estimate large-scale airplane performance from swept-wing model tests. Expectations of improving $C_{l_{\text{max}}}$ commensurate with that experienced at zero sweep are not likely to be fulfilled. The importance of this problem would indicate a pressing need for swept-wing tests of a number of given models throughout the full Reynolds number range.
It should be noted that the above inferences have been drawn from results of tests of wings using conventional airfoil sections. It may well be that when sufficient data become available to make similar comparisons on wings using laminar-flow sections, the effects of roughness and Reynolds number may be markedly different.

L/D ratio. - The variation of L/D with lift coefficient is shown in figure 10 for each of the five wings with partial and with full-span flaps. That part of the drag attributed to induced drag has been corrected to the aspect ratio of the unswept wing, that is, an aspect ratio of 4.62. The L/D values for conditions where the drag coefficient was less than 0.1 are not shown because it is believed possible inaccuracies due to lack of precise tare values would invalidate any conclusion drawn from such results. This excluded study of the most important flight speed range for plain wings and hence the L/D values for plain wings are not shown. It is believed, however, that the results shown for the wings with flaps are sufficiently accurate to allow useful conclusions to be drawn as to the effect of sweep on L/D ratios for that region of the most interest, centering around gliding and landing.

The results show that at lift coefficients near a maximum lift coefficient of 0.3 the L/D for the swept wings

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*The use of aspect ratio based upon the span of wing rather than the length of the quarter-chord line is justified on the basis of results quoted in reference 6.*
approximates that for the unswept wing. As stall is approached the L/D ratios of the swept-back wings remain at least as high as those for the unswept wing; whereas the L/D ratios of the swept-forward wings show a rapid decrease.

Longitudinal Characteristics

The effects of sweep on the pitching-moment characteristics of the plain wing, wing with partial-span flaps, and with full-span flaps are shown in figures 11 to 13. The remarks which follow are based upon the data obtained on the plain wing (fig. 11) but in general apply also to the wings with partial- or full-span flaps (figs. 12 and 13).

For lift coefficients less than 0.5, the pitching-moment coefficients vary almost linearly with lift coefficient and indicate that forward sweep moves the aerodynamic center forward (4 percent M.A.C. at $-45^\circ$ sweep), while sweepback moves the aerodynamic center rearward (5 percent M.A.C. at $45^\circ$ sweep). At higher lift coefficients the $\pm 45^\circ$ swept wings, and to a lesser degree the $-30^\circ$ swept wing, exhibit an abnormal diving tendency. Similar diving tendencies of highly swept wings have been reported previously (reference 1). Such irregularities in moment characteristics do not appear serious if considered only in terms of the elevator power available with a conventional tail. However, the effect upon static stability and the abrupt variation of elevator position and of stick force with speed may prove objectionable to pilots. In
the case of a tailless design these irregularities would be more serious. For instance, if the 45° swept-back wing were considered a possible design with the 0.40-span ailerons used for longitudinal control (elevons) and with neutral stability at low lift coefficients, over 30° up-elevator travel would be required to maintain trim even if the elevator effectiveness at low lift coefficients were maintained. For the 45° swept-forward wing a similar but less extreme condition existed. Smaller control angles would be required but the data indicate an abruptness of control motion which, because of the low damping in pitch, might be serious in tailless designs. In considering the longitudinal stability it should be remembered that the effects of fuselage, tip shape, slots, etc., have been disregarded. It may be, and unpublished data so indicate, that minor configuration changes will remove the diving tendency and its associated problems.

For lift coefficients just less than $C_{L_{\text{max}}}$ the swept wings tested, with the exception of the 30° swept-forward wing, exhibited a strong climbing moment. This characteristic is obviously undesirable since it makes inadvertent stall quite likely. In reference 8 a chart was presented which defined, on the basis of small-scale data, the boundaries of aspect ratio and sweep angle which would give a wing either a climbing or a diving tendency near stall. This chart is reproduced herein as figure 14. Also shown on this figure are the data obtained in this investigation. Based upon these data, it appears that the chart as set forth in reference 8
applies as well to large-scale as to small-scale wings; furthermore, the chart applies to swept-forward as well as swept-back wings. Insofar as the over-all shape of the pitching-moment curve is concerned large-scale tests agree generally with small-scale results with the exception of minor differences. Again it should be noted that these comparisons have been made from examination of results of investigations on wings using conventional sections. The validity of the statements regarding these comparisons is as yet unsubstantiated in cases where laminar-flow sections are involved.

Lateral Characteristics

Dihedral effect. - The variation with lift coefficient of the rolling moment due to sideslip is shown in figure 15 for the plain wings and in figure 16 for the wings with flaps. The powerful influence of sweep on the dihedral effect is immediately apparent. (A scale of effective dihedral for the unswept wing has been shown on the figures to allow convenient comparisons.) Within limits, the dihedral effect due to sweep increases in proportion to lift coefficient.

Both the $30^\circ$ and $45^\circ$ plain swept-back wings reached a maximum value of $C_{l\beta}$ of $-0.0034$ ($17^\circ$ effective dihedral) at lift coefficients of 1.15 and 0.85, respectively. In the case of the swept-forward wing the maximum value of $C_{l\beta}$ increased with angle of sweep, being 0.0014 for the $-30^\circ$ swept wing and 0.0020 for the $-45^\circ$ swept wing. These maximum values for the swept-forward wings occurred in both cases near a lift coefficient of 0.9. It should be noted that while the swept-back wings show much greater dihedral effect than do the swept-forward wings, this is due largely to the dihedral effect of
the unswept wing. The incremental dihedral effect is roughly of the same order of magnitude for either direction of sweep.

The maximum dihedral effect of the wing with flaps deflected is considerably higher, about 32° for the 45° swept-back wing with full-span flaps. Such extreme dihedral would make maintenance of a wings-level attitude in the landing approach almost impossible because of extreme sensitivity in roll to slight angles of sideslip. Even with adequate lateral control it is felt that a pilot would have difficulty in reacting sufficiently fast to prevent reaching excessive angles of bank.

For the case where lift is changed by changing angle of attack (flap deflection constant), simple sweep theory gives the following relation for the parameter $\frac{\partial c_L}{\partial c_L}$:

$$\left(\frac{\partial c_L}{\partial c_L}\right)_0 = \left(\frac{\partial c_L}{\partial c_L}\right)_0 - \frac{1}{4} \frac{\tan \alpha}{57.3}$$

It is recognized that both of the terms on the right side of this equation should be modified further by a correction involving aspect ratio and edge velocity. Simple theory shows that, where asymmetrical lift exists, the corrections would be the form $A/(AE+4)$. Again the question arises as to what the value of aspect ratio should be. Obviously the choice is more complex than simply deciding whether the span should be based on conventional span or quarter-chord-line length. In attempting to correlate the subject data as well as other swept-wing data both these approaches were used. Since neither proved consistently superior to the other or to simple theory, it was decided to delete the correction entirely. It is possible that additional study of existing data together with future tests will reveal a means of determining an effective aspect ratio which when used in this connection will more accurately predict asymmetric loading conditions. It should be noted, then, that throughout the sections of this report dealing with asymmetric loading conditions (sideslip or ailerons deflected) no corrections for aspect ratio changes have been applied to the predictions for the effect of sweep.
This relation has been used to estimate the values of \( \frac{\partial C_{1\beta}}{\partial C_L} \) for the five wings tested and the results are compared with experimental values in figure 17. Reasonably good agreement is shown except in the case of the 45° swept-forward wing which had a somewhat lower value of \( \frac{\partial C_{1\beta}}{\partial C_L} \) than was predicted.

For the case where lift is changed by changing flap deflection (angle of attack constant), the theoretical effect of sweep on \( \frac{\partial C_{1\beta}}{\partial C_L} \) is twice that given by the foregoing expression, that is,

\[
\left( \frac{\partial C_{1\beta}}{\partial C_L} \right)_A = \left( \frac{\partial C_{1\beta}}{\partial C_L} \right)_{A=0} \cdot \frac{b_f}{b} \frac{\tan \Lambda}{57.3}
\]

where \( b_f/b \) is the ratio of flap span to wing span. The estimated and experimental results for this case are also shown in figure 17. The agreement between theory and experiment in this case is only fair. The discrepancy is probably due in great measure to failure of the theory to properly account for the spanwise center of load. Theory indicates rectangular loading — that is, that the center of additional load is applied at mid-semispan of wing or flap. Movement of the center of load inboard as much as 20 percent of the wing semispan would be required to make the discrepancy between theory and experiment vanish.

Thus, relations obtained by means of the simplified theory appear to estimate at least the gross effects of sweep on the parameter \( \frac{\partial C_{1\beta}}{\partial C_L} \). A notable exception is the case of the
45° swept-forward wing which exhibits a dihedral effect much less than theory would indicate but still greater than the 30° swept-forward wing.

Since the problem of determining the value of $\frac{\partial C_{\mu}}{\partial C_{L}}$ and the maximum value of $C_{\mu}$ is probably the most serious one faced by the designer of swept-wing airplanes, a considerable effort has been made to evaluate the effects of scale on swept-back wings from the data. Unfortunately, such an evaluation could not be obtained. Only the generalization can be made that the effects of scale appear much less important than the effects of wing geometry. Both large- and small-scale swept-back-wing tests show very similar characteristics. That is, the value of $\frac{\partial C_{\mu}}{\partial C_{L}}$ approximates that predicted by theory, with a maximum value of $C_{\mu}$ being reached prior to the stall, and followed by a reduction in $C_{\mu}$ as the stall is approached. Note that these and the following considerations regarding the effects of scale apply to plain swept-back wings only.

As previously noted, the value of $\frac{\partial C_{\mu}}{\partial C_{L}}$ indicated by simple theory is based upon the assumption that the additional load is concentrated at the mid-semispan. Therefore marked differences in this parameter would be expected where nonrectangular loading was known to exist. In comparing experimental results with the theory such was found to be the case. Reference 1 showed that for the rectangular swept-back wings the measured value was as much as 14 percent more than
the predicted value. The tests reported herein gave experimental results less (as much as 14 percent less) than the predicted value. Such differences might be anticipated since theory shows that sweepback tends to shift the load center towards the tips; and taper ratios less than 1.0 tend to shift the load center towards the root. A complete understanding of this action cannot be had until more thorough studies are made of the effects of sweep and taper on load center. A first approximation (probably an overcorrection) of the answer can be reached, however, by simply adjusting the load center to correspond to the area center. If this is done, theory would fall within 10 percent of the results shown in this report while, of course, the discrepancies of reference 1 would be unchanged. Such a procedure applied to the results of reference 3 would slightly overcorrect for the effects of taper that are shown. From this it can be concluded that the value of \( \frac{\delta C_{l_f}}{\delta C_L} \) can be approximated to within 15 percent by simple theory; that a closer approximation can be had—probably within 10 percent—if the centroid of load is assumed to lie on the centroid of area. It is believed that the effects of scale fall within this latter error and probably are of the same general magnitude as the effects of section or tip shape. No data could be found to aid in a quantitative evaluation of these effects.

With regard to the maximum values of \( C_{l_f} \) likely to be encountered with a highly swept wing, it appears impossible to
conclude more than the fact that a maximum value exists for every wing and that this maximum value tends to decrease with taper ratio. The data of this report and reference 1 show very nearly the same maximum value \( C_{l_p}^{\text{max}} = 0.0035 \) to \( 0.0038 \) for both swept-back wings. Reference 3 shows a very similar maximum for the untapered wings but shows the maximum decreasing with both taper and sweep for other wings. No relation seems to exist between the lift coefficient at which the maximum \( C_{l_p} \) occurs and the maximum lift coefficient of the wing. Since, however, the value of \( \frac{\partial C_{l_p}}{\partial C_L} \) in general, increases more rapidly with sweep than \( C_{l_p}^{\text{max}} \) decreases, the maximum value of \( C_{l_p} \) occurs at progressively lower percentages of \( C_{l_p}^{\text{max}} \) as sweep is increased. For instance, for a 45\(^\circ\) sweptback wing, \( C_{l_p}^{\text{max}} \) occurs at 0.55 \( C_{L}^{\text{max}} \) in reference 3, 0.61 \( C_{L}^{\text{max}} \) in reference 1, and 0.70 \( C_{L}^{\text{max}} \) in the data of this report; whereas \( C_{l_p}^{\text{max}} \) for a 30\(^\circ\) swept-back wing occurs at 0.60 \( C_{L}^{\text{max}} \), 0.82 \( C_{L}^{\text{max}} \), and 0.91 \( C_{L}^{\text{max}} \) for the same data, respectively. Since the phenomenon which causes the value of \( C_{l_p} \) to peak are not completely understood at the present time, an accurate prediction of its value is impossible. Examination of all available data leads, however, to the conclusion that if a maximum value of \(-0.0038\) is chosen the choice can be considered conservative, but for the present, wind-tunnel tests must be relied upon to give the exact answer. Certainly this problem is worthy of additional study. Until the governing factors are more clearly defined it remains
impossible to determine to what extent $C_{l}f_{\text{max}}$ is affected by scale.

In the case of swept-back wings with flaps deflected the amount of correlation possible between large- and small-scale data is extremely limited and the results far less amenable to interpretation. For most cases examined theory gave at least a slightly conservative value of $\delta C_{l}\beta/\delta C_{L}$ where the change in lift coefficient was due either to a flap deflection at constant angle of attack or a change of angle of attack with flaps deflected. All the data, large and small scale showed or indicated that a value of $C_{l}f_{\text{max}}$ existed and that it increased with sweep. While no systematic variation of $C_{l}f_{\text{max}}$ with wing geometry could be ascertained, none of the data showed a value greater than -0.007. For the present, therefore, if wind-tunnel tests are not available the best approach to predicting $C_{l}\beta$ characteristics of swept wings with flaps is to use simple theory to predict $\delta C_{l}\beta/\delta C_{L}$ and to consider -0.007 $C_{l}\beta$ as the maximum.

No correlation was attempted with the swept-forward-wing data because of the scarcity of low-scale tests.

Aileron effectiveness.—The variation of aileron effectiveness with lift coefficient is shown in figure 15. The values of aileron effectiveness shown in figure 15 were obtained as the $\Delta C_{l}$ produced by $-15^\circ$ or $15^\circ$ of aileron deflection and hence are $\Delta C_{l}/\Delta \delta_{a}$ rather than a true $C_{l}\delta_{a}$. It is immediately apparent that aileron effectiveness decreases with sweep,
decreasing as much as 50 percent for 45° of sweep. The effectiveness of the ailerons on the swept-back wings decreased with lift coefficient, rapidly at high lift coefficients. This is due to a loss in effectiveness of the upward deflected aileron which is in the wake of the separated flow and hence contributes little or nothing to the rolling moment. The ailerons on the swept-forward wings show a general increase in effectiveness with lift coefficient, probably due to a favorable effect of the spanwise boundary-layer drain.

According to the simplified theory, as a wing is swept, the aileron effectiveness will decrease, as for any flap, in proportion to the \( \cos^2 \Delta \). That is, when corrections for aspect ratio are ignored the value of \( C_l \delta_a \) is given by the following relation: (See footnote 6, p. 21)

\[
(C_l \delta_a)_\Lambda = (C_l \delta_a)_\Lambda = 0 \cos^2 \Delta
\]

This relation has been used to predict the variation of aileron effectiveness with sweep for the five wings tested, and the results are compared in figure 19 with the experimental data cross-plotted from figure 18 for zero lift.

---

7The ailerons on the wings of this investigation varied both in the relative amount of wing area affected and in the relative spanwise location of the center of pressure of the area affected. In comparing theory and experiment, these variations were accounted for by correcting the theoretical values of aileron effectiveness in proportion to the ratio of the relative area and spanwise center of pressure of the swept wing to the relative area and spanwise center of pressure of the unswept wing.
For either sweepforward or sweepback the experimental values of aileron effectiveness are as much as 20 percent lower than the theoretical values.

The foregoing results show that aileron effectiveness is reduced by wing sweep and that on swept-back wings the aileron effectiveness is further reduced at high lift coefficients. Insofar as rolling control at low lift coefficient is concerned, theory shows that $C_{Lp}$ is reduced in proportion to $\cos \Delta$; whereas $C_{L\delta a}$ is reduced in proportion to $\cos^2 \Delta$ and hence $pb/2V$ will be reduced in proportion to $\cos \Delta$ for a given size of aileron. In general, therefore, it appears that, to maintain a given value of $pb/2V$, aileron size must be increased as wings are swept. As higher lift coefficients are reached the lateral-control problem becomes particularly pronounced. Not only must powerful lateral control be provided to overcome great dihedral effects but the results reported herein show that available lateral control, at least for swept-back wings, decreases seriously with lift coefficients. For example, with the 45° swept-back wing equipped with full-span flaps and flying near $C_{l_{\text{max}}}$, 13° of total aileron deflection would be required to hold the wings level for only 1° of sideslip. The need for development of adequate aileron control or a means to reduce aileron control requirements is obvious.
Directional Characteristics

The variation with lift coefficient of the yawing moment due to sideslip is shown in figure 20 for the plain wings and in figure 21 for the wings with partial- and full-span flaps. Sweepback increased the directional stability and sweepforward decreased the stability; however, due to the initial positive stability of the unswept wing the stability of the swept-forward wings became negative only at higher lift coefficients and then only slightly so.

The theoretical effect of sweep on the directional stability is in accordance with the following relation:

\[
\left( \frac{\partial C_{n\beta}}{\partial C_{L}^2} \right)_{\Delta} = \left( \frac{\partial C_{n\beta}}{\partial C_{L}^2} \right)_{\Delta=0} + \frac{\tan \Delta}{2\pi \Delta \times 57.3}
\]

The directional stability estimated on the basis of this equation is compared with experimental results in figure 22. Although precise agreement is not obtained, the trend of the experimental data is indicated by theory.

The directional characteristics of the swept wings tested should not present any serious problems of a purely low-speed static directional stability and control nature since adequate stability and control should be obtainable by use of fins and rudders of normal proportions combined with normal tail lengths. However, a dynamic problem arises from the fact that \( C_{l\phi} \) increases with sweep more rapidly than \( C_{n\beta} \). This unbalance between \( C_{l\phi} \) and \( C_{n\beta} \) leads to the dutch-roll type
of instability which has been discussed in reference 8. The data obtained in the investigation reported herein substantiate previous tests conducted on small-scale models and consequently indicate that means must be found to balance $C_{l}$ and $C_{r}$.

CONCLUDING REMARKS

Large-scale tests indicate that the primary problems to be overcome before successful use can be made of high angles of sweep are (1) high dihedral effects accompanied by poor lateral control at high lift coefficients, (2) low maximum lift value together with low flap effectiveness, and (3) rapid shift in neutral point in the moderate to high lift-coefficient range coupled with a possibility of strong stalling moment at maximum lift resulting from poor plan-form choice.

In general, simple theory enables good predictions to be made of the gross effects of sweep on wing characteristics, but it is felt that the accuracy is inadequate for purposes of design. It appears that the majority of the inaccuracies result from an incomplete understanding of the effects of aspect ratio.

Where it has been found possible to compare large-scale data with small-scale data a comparison has shown that where scale effects exist at low angles of sweep, scale effects
tend to vanish at high angles of sweep with large-scale results approaching small-scale results:

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APPENDIX

Description of Basic Wind-Tunnel Test Results

For each swept wing, six-component force data were obtained at several angles of sideslip and several values of dynamic pressure. (See fig. 23 for variation of Reynolds number with dynamic pressure.) At each angle of sideslip, several model configurations were tested including plain wing, wing with partial-span split flaps, wing with full-span split flaps and wing with split-flap-type aileron. The data obtained are presented in figures 24 to 91 in terms of the variation of the measured characteristics with lift coefficient. Table II forms an index of these figures presenting the basic data.

All the data are referred to the stability axes whose origin is located at a point on the root chord or root chord projected and at the same fore and aft location as the quarter M.A.C. The test results are presented in the form of standard NACA coefficients as defined in the section Coefficients and Symbols.

All the basic wind-tunnel data have been corrected for air stream inclination and for wind-tunnel-wall effects. A brief analysis of the effect of sweep on tunnel-wall corrections indicated that the average correction either with or without sweep was approximately the same for the tunnel wing configurations considered. Hence the standard corrections for unswept wings were applied.
Force tests made with the sting support alone in the tunnel showed that its tare should be negligible except in the case of pitching moment, drag and yawing moment. Measured pitching-moment tares are believed reliable and were applied to all the data. While the drag tares are appreciable (approximately 0.02 in the case of the unswept wing where the area is small and decreasing for the swept wings where the area is larger), it is felt that they could not be determined with sufficient accuracy to warrant application. Hence no drag tares have been applied. Since the measured yawing-moment tares (fig. 92) were small, they were not applied to the basic data. However, in analysis of the data it was found that the tares were relatively large when compared to the effects of sweep. In order then to properly assess the effects of sweep, it was necessary to apply tares to the summary data which is, therefore, shown fully corrected in figures 20, 21, and 22.
REFERENCES


TABLE I.—GEOMETRIC CHARACTERISTICS OF THE FIVE SWEPT WINGS

<table>
<thead>
<tr>
<th>Angle of sweep (deg)</th>
<th>Span (ft)</th>
<th>Area, S (sq ft)</th>
<th>Aspect ratio, $b^2/S$</th>
<th>Taper ratio, $c_t/c_x$</th>
<th>1 Mean aerodynamic chord, c (ft)</th>
<th>2 Span of partial-span flaps (% b)</th>
<th>2 Span of full-span flaps (% b)</th>
<th>2 Span of ailerons (% b)</th>
</tr>
</thead>
<tbody>
<tr>
<td>-45</td>
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<td>335.5</td>
<td>3.12</td>
<td>0.38</td>
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<td>62.3</td>
<td>97.0</td>
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<td>10.00</td>
<td>62.3</td>
<td>96.2</td>
<td>33.9</td>
</tr>
</tbody>
</table>

1 All chords were measured parallel to the air stream.

2 Flaps and ailerons were 20 percent chord.
TABLE II.- INDEX TO THE BASIC DATA FIGURES

<table>
<thead>
<tr>
<th>Figure number</th>
<th>$\Lambda=-45$</th>
<th>$\Lambda=-30$</th>
<th>$\Lambda= 0$</th>
<th>$\Lambda=30$</th>
<th>$\Lambda=45$</th>
<th>Nominal angle of sideslip</th>
<th>Configuration</th>
<th>Data presented</th>
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<td>64</td>
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<td>37</td>
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<td>65</td>
<td>79</td>
<td>0</td>
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<td>$C_D, \alpha, C_m, C_l, C_n, C_Y$ vs $C_L$</td>
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<td>38</td>
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<td>66</td>
<td>80</td>
<td>0</td>
<td>Plain wing + full-span flaps</td>
<td>$C_D, \alpha, C_m, C_l, C_n, C_Y$ vs $C_l$</td>
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</tr>
<tr>
<td>27</td>
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<td>53</td>
<td>67</td>
<td>81</td>
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<td>73</td>
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<td>Plain wing + ailerons</td>
<td>$\alpha, C_m, C_l$ vs $C_L$</td>
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<td>34</td>
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<td>Plain wing + partial-span flaps</td>
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<td>$C_D, \alpha, C_m, C_l, C_n, C_Y$ vs $C_L$</td>
<td></td>
</tr>
</tbody>
</table>

$^1$Actual angle of sideslip is noted on each curve sheet.
(a) The $45^\circ$ swept-forward wing.

(b) The unswept wing.

Figure 1.— Photographs of three of the swept wings mounted in the Ames 40- by 30-foot wind tunnel.
Figure 1.- Concluded. (c) The 45° swept-back wing.
Figure 2a, b. - Geometric characteristics of the swept wings.
Fig. 2b - Concluded.

Sweep = 0°, Area = 301.6 sq ft, Aspect Ratio = 4.69, Taper Ratio = .55

Sweep = +30°, Area = 283.4 sq ft, Aspect Ratio = 4.84, Taper Ratio = .44

Sweep = +45°, Area = 308.6 sq ft, Aspect Ratio = 3.64, Taper Ratio = .43

(b) Figure 2b - Concluded.
Figure 3. - Sign convention for the standard NACA coefficients. All forces, moments, angles, and control surface deflections are shown as positive.
Figure 4.- Comparison between theoretical and experimental effects of sweep on lift-curve slope

Figure 5.- Effect of sweep on lift-curve slope
Figure 6. Chart summarizing the effect of angle on the angle of incidence back at which the maximum value of lift curve slope occurs.

Figure 7. Comparison between the theoretical and experimental effects of sweep on the increment of lift obtainable at 0° angle of attack by deflecting full-span and partial-span split flaps 60°.
Figure 8. Effect of sweep on the maximum lift coefficient.
Figure 10. Variation with lift coefficient of lift-drag ratio of the five swept wings.

Figure 11. Pitching-moment characteristics of the five swept wings. Plain wing configurations.
Figure 12.- Pitching-moment characteristics of the five swept wings, 65.9% span split flaps deflected 60°.

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Figure 14.- Chart summarizing the effect of aspect ratio on the pitching-moment curve of sweptback wings at the stall.

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Figure 17: Variation with sweep of the estimated and experimental values of $\frac{\Delta C_L_{16}}{\Delta C}$.

(a) $\frac{\Delta C_L}{\Delta C}$ at constant flap setting.

(b) $\frac{\Delta C_L}{\Delta C}$ at constant angle of attack.

Figure 16: The effect of dihedral on the lift coefficient for $20\%$ chord split flap.
Figure 18: Variation with lift coefficient of aileron effectiveness based on aileron deflection.

Figure 19: Effect of sweep on the aileron effectiveness for the plain wing at zero lift.

Note: Theoretical values of $C_{L}$ have been corrected for relative wing area affected by ailerons and for span-wise center of pressure of area affected.
Figure 20: The effect of sweep on the variation of directional stability with lift coefficient, plan wing.

Figure 21: The effect of sweep on the variation of directional stability with lift coefficient. 20% chord split flaps deflected 60°.
Figure 22. - Variation with sweep of the estimated and experimental values of wall boundary layer thickness.

Figure 23. - Variation of Reynolds number with dynamic pressure for wings tested.
Figure 24.- Aerodynamic characteristics of the 95° swept-forward wing at ±0.2° sideslip, plain wing.

(a) $C_D$, $C_m$ vs $C_L$
FIGURE 25.- AERODYNAMIC CHARACTERISTICS OF THE 95° SWEEPBACK WING AT 0.2° SIDESLIP,
40% CHORD, 62.3% SPAN SPLIT
FLAPS DEFLECTED 60°.
ROLLING-MOMENT COEFFICIENT, $C_m$

YAWING-MOMENT COEFFICIENT, $C_n$

SIDE-FORCE COEFFICIENT, $C_y$

(b) $C_f$, $C_n$, $C_y$ vs $C_L$

Figure 25 - Concluded.
Figure 25.-Aerodynamic characteristics of the 85° swept-forward wing at 150° dihedral 20% chord, full-span split flaps deflected 80°.
(b) $C_f, C_n, C_y$ vs $C_L$

Figure 26.- Concluded
Figure 28. Aerodynamic characteristics of the 75° swept-back wing at -3.5° incidence plain wing.
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ROLLING-MOMENT COEFFICIENTS $C_1$

YAWING-MOMENT COEFFICIENT $C_m$

SIDE-FORCE COEFFICIENT $C_y$

$C_l$, $C_d$, $C_y$ vs $C_l$

FIGURE 28.- CONCLUDED.
Figure 28 - Aerodynamic characteristics of the 95° swept-forward wing at -63° incidence.
20% chord, 62.3\% span split flaps deflected 60°.
Fig. 28b

(b) \( C_1, C_2, C_3, C_4 \) vs \( C_L \)

Figure 28.-Concluded.
Figure 30.- Aerodynamic characteristics of the 95° swept-forward wing at -3.5° incidence in a plain wing.
Figure 36.-Concluded.
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20% CHORD, 62.8% SPAN SPLIT FLAPS DEFLECTED 60°
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Figure 33.- Aerodynamic characteristics of the 95° sweep-forward wing at 15° angle of attack, plain wing.
Figure 33—Concluded.
Figure 34.- Aerodynamic characteristics of the 95° sweep-forward wing at 45° Mach number
20% chord, 62.5% span split flaps deflected 60°.
Figure 35.- Aerodynamic characteristics of the 95° swept forward wing at 4° sideslip 20% chord 94.4% span split flap type ailerons deflected 11°.
Figure 36. Aerodynamic characteristics of the 30° swept-forward wing at -0.2° sideslip. Plated wing.
Fig. 36b

ROLLING-MOMENT COEFFICIENT, $C_l$

YAWING-MOMENT COEFFICIENT, $C_m$

SIDE-FORCE COEFFICIENT, $C_y$

(a) $C_l$, $C_m$, $C_y$ vs $C_e$

FIGURE 36 - CONCLUDED.
Figure 37 - Aerodynamic Characteristics of the 30° Sweptback Wing at -12° Incidence
20% Chord, 52.3% Span Split Flaps
Deflected 60°
Figure 37: Continued.
Figure 38. - Aerodynamic characteristics of the 90° swept-forward wing at -0.2° sideslip.
20% chord, full span split flaps deflected 60°.
Fig. 38b

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ROLLING-MOMENT COEFFICIENT, $C_I$
YAWING-MOMENT COEFFICIENT, $C_n$
SIDE-FORCE COEFFICIENT, $C_y$

(b) $C_I$, $C_n$, $C_y$ vs $C_L$

Figure 38b.- Concluded.
Figure 40.- Aerodynamic characteristics of the 20° sweptforward wing at -5.3° sideslip. Plain wing.

\[ C_{D} \alpha, C_m \text{ vs } C_L \]
Figure 10. - Concluded.
Figure 41.- Aerodynamic characteristics of the 30° sweep-forward wing at -5° sideslip, 20% chord, 62.9% span split flaps deflected 60°.
Figure 41.- Concluded.
Figure 42. Aerodynamic characteristics of the 20° swept-forward wing at -18° sideslip. Plain wing.
Figure 42—Concluded.
FIG. 43 - AERODYNAMIC CHARACTERISTICS OF THE 30° SWEEP FORWARD WING AT -45° SIDESLIP
20% CHORD, 62.8% SPAN SPLIT FLAPS DEFLATED 60°
Fig. 45—Concluded.
Figure 45. - Aerodynamic characteristics of the 35° swept-forward wing at -3.8° incidence.
40% chord, 31% span, split flap type ailerons deflected 2.16°.
Figure 46. Aerodynamic characteristics of the 30° sweepforward wing at 450° incidence plain wing.
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ROLLING-MOMENT COEFFICIENTS, $C_n$

FORCE-MOMENT COEFFICIENTS, $C_m$

YAWING-MOMENT COEFFICIENTS, $C_y$

(b) $C_l, C_n, C_y$ vs $C_L$

FIGURE 45.—CONCLUDED.
Figure 47a - Aerodynamic characteristics of the 30° swept-forward wing at +50° sideslip.
20% chord, 62.3% span split flaps deflected 60°.
Figure 47 - Concluded.
Figure 48. Aerodynamic characteristics of the 30° swept-forward wing at a 50° incidence, 20% chord full span split flaps deflected 60°.
FIG. 48b

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(a) CL, CR, CY vs CL

(b) CL, CR, CY vs CL

FIGURE 48b.—Concluded.
Figure 48: Aerodynamic characteristics of the 30° sweptforward wing at +5° sideslip 20% chord 31.3% span split flap type ailerons deflected 15°.
Figure 50a: Aerodynamic characteristics of the unswept wing at -0.1° sideslip. Plain wing.
Figure 51 - Aerodynamic characteristics of the 0° swept wing at -6.1° dihedral, 50% chord, 62.8% span split flaps deflected 60°
Figure 52: Aerodynamic characteristics of the 0° swept wing at 0.1° sideslip, 20% chord, full span split flaps deflected 60°.
Figure 52b

Figure 62—Concluded.
Figure 53. Aerodynamic characteristics of the 0° swept wing at -0.1° sideslip, 80% chord, 35.2% span, split flap type aileron on right wing deflected 2.5°.
Figure 54. Aerodynamic characteristics of the 0° swept wing at -5.2° incidence. Plain wing.
Figure 56: Aerodynamic characteristics of the 0° swept wing at -6.5° sideslip at 20% chord, 62.5% span split flaps deflected 60°.
ROLLING-MOMENT COEFFICIENT, $C_I$

SIDE-FORCE COEFFICIENT, $C_Y$

YAWING-MOMENT COEFFICIENT, $C_N$

(b) $C_I$, $C_a$, $C_Y$ vs $C_L$

FIGURE 55b—Concluded.
Figure 56. - Aerodynamic characteristics of the 0° swept wing at -10° sideslip.
Plain wing.
$C_l, C_n, C_y \text{ vs } C_e$

Figure 56—Concluded.
Figure 5.1 - Aerodynamic Characteristics of the O° Swept Wing Fit - 10° Sidewash 20% Chord 62.5% Span Split Flaps Deflected 60°
Figure 57—Concluded.
Figure 58.—Concluded.
Figure 60. - Aerodynamic characteristics of the 0° swept wing at +4.4° sideslip. Plain wing.
Figure 60b—Concluded.

(b) $C_l$, $C_m$, $C_y$ vs $C_l$
Figure 61. Aerodynamic characteristics of the 0° swept wing at ±4.4° sideslip, 20% chord, 62.3% span split flaps deflected 60°.
Figure 61e—Concluded.
Figure 62.—Aerodynamic characteristics of the 0° swept wing at 14.4° sideslip 20% chord full span split flaps deflected 60°.
Figure 62.- Concluded.
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Figure 63. Aerodynamic characteristics of the 0° swept wing at 14.4° sideslip, 20% chord, 35.2% span, split flap type ailerons deflected 15°.
Figure 64. Aerodynamic characteristics of the 30° swept-back wing at -0.3° sideslip. Plain wing.
Figure 64.-Concluded.

(b) \( C_{l}, C_{m}, C_{y} \) vs \( C_{l} \)
Figure 66—Aerodynamic Characteristics of the 30° Sweepback Wing at -0.8° Side Slip, 20% Chord, 62.5% Span Split Flaps Deflected 60°.
Figure 65: Continued.

(b) $C_1, C_6, C_7$ vs $C_L$
Figure 66. — Aerodynamic characteristics of the 30° sweptback wing at -0.3° sideslip, 20% chord full span, split flaps deployed 60°.
Fig. 66b

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Rolling-moment coefficient, $C_I$

Yawing-moment coefficient, $C_n$
(b) $C_I, C_n, C_y$ vs $C_L$

Figure 66.- Concluded.
Figure 67: Aerodynamic characteristics of the 30° sweptback wing at -0.3° sideslip, 20% chord 36.5% span split-flap type aileron on right wing deflected 2 1/2°
Figure 68.- Aerodynamic characteristics of the 30° sweepback wing at -5.5° sideslip, plain wing.
Figure 68.—Concluded.
Figure 65.- Aerodynamic characteristics of the 30° sweepback wing at -6.3° sideslip, 80% chord, 62.5% span split flaps deflected 50°.
Figure 69. - Concluded.
FIG. 70a

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Drag coefficient, $C_d$  
Pitching-moment coefficient, $C_m$

Angle of attack, $\alpha$

$C_0$, $\alpha$, $C_m$ vs $C_L$

FIGURE 70. - Aerodynamic characteristics of the 30° sweptback wing at -9.5° sideslip, Plain wing.
Figure 70a — Concluded.
Figure 71. - Aerodynamic characteristics of the 30° sweepback wing at -9.5° sideslip 20% chord 62.5% span split flaps deflected 60°.
Figure 75. Aerodynamic characteristics of the 30° sweptback wing at -5.5° sideslip, 20% chord, full span split flaps deflected 60°.
Figure 73.- Aerodynamic characteristics of the 30° sweptback wing at -9.5° sideslip, 20% chord.
26.5% span split flap type ailerons deflected ±15°
Figure 74A - Aerodynamic characteristics of the 30° sweepback wing at +4.1° sideslip, plain wing.
(b) $C_l$, $C_n$, $C_y$ vs $C_L$

FIGURE 74b - CONCLUDED.
Figure 75.- Concluded.
Figure 76.- Aerodynamic characteristics of the 30° sweptback wing at +4.1° sideslip, 20% chord full span split flaps deflected 60°.
FIGURE 76.- Concluded.
FIGURE 77 - AERODYNAMIC CHARACTERISTICS OF THE 30° SWEPTBACK WING AT +1° SLIP, 25% CHORD 36.9% SPAN SPLIT FLAP TYPE, AILERONS DEFLECTED 15°
Figure 78. - Aerodynamic characteristics of the 45° sweptback wing at -0.1° sideslip, plain wing.
Figure 78.-Concluded.
Figure 79 - Aerodynamic characteristics of the 45° sweptback wing at -0.1° sideslip, 60% chord 62.5% span split flaps deflected 60°
(b) $C_L, C_m, C_y$ vs $C_L$

FIGURE 79. CONCLUDED.
Figure 60.—Aerodynamic characteristics of the 45° sweptback wing at -0.1° sideslip, 20% chord full span split flaps deflected 60°
Figure 80b - Concluded.
Figure 81: Aerodynamic characteristics of the 45° sweptback wing at -0.1° sideslip; 20% chord, 34.1% span split flap type ailerons deflected ±15°.
Figure 82: Aerodynamic characteristics of the 45° sweptback wing at -5.2° sideslip, Plain wing.
Figure 82—Concluded.

(b) $C_l$, $C_n$, $C_y$ vs $C_L$
Figure 83.- Aerodynamic characteristics of the 45° sweepback wing at -5.5° sideslip, 80° chord, 62.5% span, split flaps deflected 60°.
Figure 88.—Concluded.
Figure 84. Aerodynamic characteristics of the 45° sweptback wing at -0.9° sideslip, plain wing.
Fig. 84d

Figure 84.—Concluded.
Figure 86.-Aerodynamic Characteristics of the 45° Shearback Wing at -9.5° sideslip, 20% chord 62.5% span split flaps deflected 60°.
Figure 86.- Concluded.
Figure 86c. — Aerodynamic characteristics of the 45°
swepback wing at -9.9° sideslip, 20% chord
full span split flaps deflected 60°.
Figure 86. - Concluded.
Figure 87.- Aerodynamic characteristics of the 45° sweepback wing at -5.5° incidence, 50% chord, 34.1% span split flap type ailerons deflected ±15°.
Figure 88. - Aerodynamic characteristics of the 45° sweptback wing at +4.1° sideslip, plain wing.
Figure 88.- Concluded.
Figure 89. Aerodynamic characteristics of the 46° sweptback wing at 41° sideslip, 20% chord, 62.3% span spl. Flaps deflected 50°.
Figure 90: Aerodynamic characteristics of the 45° sweptback wing at +1.1° sideslip. 20% chord full span split flaps deflected 60°.
FIG. 90b

(b) $C_l$, $C_n$, $C_y$ vs $C_L$

Figure 90b, Concluded.
Figure 81. - Aerodynamic characteristics of the 45° sweptback wing at +1° sideslip, 20% chord 34.1% span split flap type ailerons deflected 5°.
Figure 92: Yawing-moment coefficient tare of the five swept wings. (Applied to basic data to obtain summary data.)