COMPUTATIVE EXAMINATION OF BENDING STRENGTH OF GIRDER
ORIGINALLY CURVED AND SUBJECT TO LONGITUDINAL COMPRESSION.

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The 1918 military specifications (BLV) contained the following paragraphs:

"It must not only be demonstrated that the longitudinal stress generated in a girder is below the Euler compression force, but also that, with an initial deflection equal to 1/200 of the length of the girder, the stress produced in it by half the prescribed breaking load, does not exceed half the breaking strength.

"Short compression struts, whose cross-section may have a radius of gyration \( r \), must in so far as \( s/i<105 \) for cast iron and steel and \( s/i<110 \) for wood, not be computed according to the Euler, but according to the Tetmajer compression formula.

This stipulation is based on an article by Müller Breslau.** It is accordingly requested that at least double the actual load on the lever arm, 1/200 of the length of the girder, should at most create a stress near the limit of proportionality.

Since it is difficult to determine the actual load in airplane construction and since the introduction of the breaking load into the strength computation was customary, the Müller Breslau requirements, of double the actual load and the limit of proportionality, were converted into the stipulations of half the breaking load and half the breaking strength.

The following consideration is based on the designations in

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Fig. 1. The girder s(cm) long has an incipient parabolic curve indicated by the expression

\[ f = 4 \frac{r}{s} s \left[ \frac{x}{s} - \left( \frac{x}{s} \right)^2 \right] \]  

(1)

The curve is the greatest in the middle and has the value \( f_{\text{max}} = r \). Under the influence of the longitudinal forces \( S(\text{kg}) \), the girder is bent to the position of the dash line. At the point \( P \), with the coordinates \( x \) and \( y + f \), there is a bending moment

\[ M = S (y + f) \]  

(2)

This moment is held in equilibrium by the elastic forces of the girder.

\[ M = - E I \frac{d^2 y}{d x^2} \]  

(3)

in which \( E (\text{kg/cm}^2) \) is the modulus of elasticity and \( I (\text{cm}^4) \) the moment of inertia of the girder.

Expressions 1, 2 and 3, after the introduction of

\[ k = \sqrt{\frac{E I}{s}} \]  

(4)

are combined in the differential equation

\[ k^2 \frac{d^2 y}{d x^2} + y = 4 \frac{r}{s} s \left[ \left( \frac{x}{s} \right)^2 - \frac{x}{s} \right] \]  

(5)

This is integrated in the usual manner, as also, after introduction of the integration constants for \( x = 0 \) and \( x = s \) and the simplification
The solution:

\[ y = 4 r s \frac{2}{a^2} \left[ \cos \frac{x}{k} + \frac{1 - \cos \alpha}{\sin \alpha} \sin \frac{x}{k} - 1 \right] + \left[ \left( \frac{x}{s} \right)^2 - \frac{x}{s} \right] \]  

The bending moment from equation (3) is transformed to

\[ M = \frac{8 r s}{a^2} S \left[ \cos \frac{x}{k} + \frac{1 - \cos \alpha}{\sin \alpha} \sin \frac{x}{k} - 1 \right] \]  

and consequently increases proportionally with the maximum initial deflection \( r s \).

This bending moment attains, with \( x = 1/2 \), its maximum value

\[ M_{\text{max}} = 8 r s \left[ \frac{1}{\cos \frac{a}{2}} - 1 \right] \frac{E I}{s^2} \]  

On the other hand

\[ M_{\text{max}} = \sigma W \]  

in which \( \sigma \) (kg/cm\(^2\)) represents the bending stress in the middle of the girder and \( W \) (cm\(^3\)) the resistance moment of the girder.

According to the already cited stipulation of the 1918 military specifications, the bending stress under the action of the \( \text{th} \) part of the bending load as longitudinal load, must not exceed a certain \( \text{th} \) part of the breaking stress. From this requirement follows

\[ \sigma = \frac{K_p}{n} - \frac{S_k}{mF} \]
in which the designations $K_b$ (kg/cm²), for the breaking bending stress; $S_k$ (kg), for the Euler bending load and $F$ (cm³), for the cross-sectional area of the girder, are newly introduced.

By employing the Euler equation

$$\frac{S_k}{m} = \frac{\pi^2 E I}{m s^2}$$ (12)

in equation (6) we have

$$\alpha = \frac{\pi}{\sqrt{m}}$$ (13)

By combining equations (9) and (10) and introducing equations (12) and (13), we have

$$W \left[ \frac{K_b}{m} - \frac{\pi^2 E I}{m s^2} \right] = 5 \pi s \left[ \frac{1}{\cos \frac{\pi}{2 \sqrt{m}}} \right] - 1 \frac{E I}{s^2}$$ (14)

whence we obtain, for the length $s$ of the girder, the quadratic equation

$$s^2 - 3 \pi n \frac{E I}{K_b} \left[ \frac{1}{\cos \frac{\pi}{2 \sqrt{m}}} \right] s = \pi^2 \frac{n E I}{m K_b F}$$ (15)

If $D$ (cm) designates the distance between the outermost fibers of the cross-section of the bent girder under consideration, we may substitute

$$\frac{I}{W} = \frac{D^2}{2}$$ (16)

If we make the radius of gyration $i$, through the equation $i = \pi D$ dependent on $D$, we may also substitute

$$i = \frac{D}{\pi}$$ (17)
\[
\frac{1}{F} = 1^2 = e^2 D^2
\]

By employing the simplified equations (16) and (18), we finally obtain

\[
\frac{S}{D} = A + \sqrt{A^2 + B^2}
\]

in which are introduced

\[
A = 2 r n \frac{E}{K_b} \left[ \frac{1}{\cos \frac{\pi}{2 \sqrt{m}}} - 1 \right]
\]

and

\[
B = \pi e \sqrt{\frac{n E}{m K_b}}
\]

The quantity \( A \) is independent of the cross-section of the girder and increases in direct proportion with the quantities \( r \) and \( n \) and decreases with increasing values of \( m \). The quantity \( B \) increases with increasing \( n \) and decreasing \( m \).

This important result means that, for a given cross-section of a girder, designated by the factor \( e \), for a fixed ratio of \( E \) to \( K_b \), as well as for chosen quantities, \( r, m, \) and \( n \), there is a single ratio \( s/D \), which fulfills the requirement that, with an initial maximum deflection \( rs \) of the \( s \)-long girder for the \( n \)-part of the Euler bending load, also the \( n \)-th part of the breaking stress will be attained.

Since the value \( s/D \), thus computed for any cross-section, applies to all cross-sections possessing geometrically similar or like \( e \) values, it constitutes a great simplification, as will
subsequently be demonstrated in the example for cylindrical pipes.

The stress produced by the bending load $S_k$ is given by equation

$$\sigma_k = \frac{S_k}{F} = \left(\pi e \frac{D^2}{g}\right) E$$

(23)

Accordingly, for a fixed ratio $s/D$, also the stress produced by the bending load remains the same for all similar cross-sections.

Computation example. Cylindrical tubes. - For the cross-section of a cylindrical tube the equations

$$l = \frac{\pi}{64} (D^4 - d^4)$$

and

$$F = \frac{\pi}{4} (D^2 - d^2).$$

apply, when $D$ (cm) designates the outside and $d$ (cm) the inside diameter of the tube, hence:

$$e = \frac{1}{4} \sqrt{l + \left(\frac{d}{D}\right)^2}$$

In Fig. 2, $e$ and $d/D$ are plotted against each other.

After choosing the ratio $E/K_b = 550$, which corresponds approximately to the malleable, seamless steel tubes employed in airplane construction, the ratio $s/D$ is plotted against the quantity $r$ (Fig. 3) for the quantities $m = n = 1.8$, $m = n = 2.0$ and $m = n = 2.2$, as well for solid rods as for tubes with infinitely thin walls.

The curves of Fig. 3 indicate that the influence of the thic-
ness of the walls vanishes for high values of $r$, i.e. large initial deflections of the girder. For a very small initial deflection, hence small values of $r$, the influence of the thickness of the wall cannot be disregarded. If the initial deflection is taken too small, we obtain $s/D$ ratios, which lie beyond the region of application of the Euler bending formula for which the above consideration alone holds good. The Euler-Tetmajer-limits are introduced in Fig. 3, in accordance with the ratio which applies to cast iron and steel (See paragraph 2 of the quotation from the 1918 BLV at the beginning of this treatise).

$$\frac{s}{D} \geq \frac{110}{4} \sqrt{\left[1 + \left(\frac{d}{D}\right)^2\right]}$$

Solid cylindrical rod

$$\frac{s}{D} \geq 37.5$$

Tube with infinitely thin walls

$$\frac{s}{D} \geq 38.9$$

In the 1918 BLV, the stipulated quantity $r = 1/300 = 0.005$ gives values of $s/D$ which lie barely above the Tetmajer limits.

If we examine the struts of airplanes which have given satisfactory service, we find no simple illustration of the distribution of the $s/D$ values. Many of them lie in the Tetmajer and others far within the Euler region.

Since with slender wing struts, bucklings are not infrequent
as a consequence of some injury and since such struts have greater
\( \frac{a}{D} \) values than would correspond to an initial buckling of \( \frac{1}{200} \) or \( r = 0.005 \), a departure from the 1918 BLV is justified.

The ordinarily employed seamless steel tubes have a diameter
ratio of \( \frac{d}{D} = 0.9 \). For this value, as also for the quantities
\( r = 0.01 \) and \( m = n = 3 \), we have

\[
\frac{\sigma}{D} = 65.
\]

The stress \( \sigma_k \) for a cylindrical tube is found to be

\[
\sigma_k = \frac{n^2}{16} \left[ 1 + \left( \frac{d}{D} \right)^2 \right] \left( \frac{D}{s} \right)^2 E
\]

For \( E = 2150000 \text{ kg/cm}^2 \), for \( \frac{d}{D} = 0.9 \) and for \( \frac{s}{D} = 65 \), we have

\( \sigma_k = 570 \text{ kg/cm}^2 \)

If, as is usually the case, \( \frac{s}{D} \) is not computed for
\( r = 0.01 \) and \( m = n = 3 \), but is differently employed, then it
must be ascertained at what share of the bending load the desired
nth part of the breaking stress occurs. For this purpose, for
fixed values \( r \) and \( n \) with equations (19), (20), and (21), the
\( m \) value must be found at which equation (19) is fulfilled.

In Fig. 4 \( m \) and \( \frac{s}{D} \) are plotted against each other for the
fixed values \( E/K_b = 550 \), \( \frac{d}{D} = 0.9 \), \( r = 0.01 \), and \( n = 3 \):

For \( \frac{s}{D} = 65 \), corresponding to the previous computation,
\( m = 2 \). For smaller values of \( \frac{s}{D} \), a larger value of \( m \) must be
used, for instance, the value \( m = 2.9 \) corresponds to \( \frac{s}{D} = 40 \).
The application of the 1918 BLV results, therefore, in a lesser utilization of the resistance to axial compression in small girders than in large ones. In the computation, this circumstance is based on the bending strength.

The results explained in the example for cylindrical tubes may also be found in a similar way for other materials and other cross-sections. Here it should be remembered that the value \( n \) is specially adapted to the material. For girders with a high elasticity limit, \( n < 2 \). For wood, however, \( n = 2 \) is recommended.

**Conclusion.**—We examined the stipulation contained in the 1918 BLV, that a girder subjected to longitudinal compression under the influence of half of the specified breaking load, along with the Euler bending safety with an initial deflection of 1/300 of the length of the girder, can, at the most, be subjected to half the stipulated breaking stress. We found a generally applicable ratio for the relation of the length \( s \) (cm) of the girder to the distance \( D \) (cm) between the outermost fibers of the cross-section subjected to bending. This ratio gives, for all geometrically similar girder cross-sections (also \( e = i/D \) constant), a length limit fulfilling these conditions.

By means of an example, we showed that the stipulation of the initial deflection of 1/300 of the girder length is not fulfilled by compression struts in use, but that, on the contrary, the assumption of a greater initial deflection, e.g. 1/100 is advisable.
Since the MIL BLV stipulation is only fulfilled for a single value of $s/D$, we showed in what manner the necessary breaking portion of the bending load is found for the values of $s/D$ differing from said value.

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