Proceedings of RIKEN BNL Research Center Workshop

Volume 56

RBRC Scientific Review Committee Meeting

November 20-21, 2003

Organizers:

N. P. Samios and T. D. Lee

RIKEN BNL Research Center

Building 510A, Brookhaven National Laboratory, Upton, NY 11973-5000, USA

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Preface to the Series

The RIKEN BNL Research Center (RBRC) was established in April 1997 at Brookhaven National Laboratory. It is funded by the "Rikagaku Kenkyusho" (RIKEN, The Institute of Physical and Chemical Research) of Japan. The Center is dedicated to the study of strong interactions, including spin physics, lattice QCD, and RHIC physics through the nurturing of a new generation of young physicists.

During the first year, the Center had only a Theory Group. In the second year, an Experimental Group was also established at the Center. At present, there are three Fellows and seven Research Associates in these two groups. During the third year, we started a new Tenure Track Strong Interaction Theory RHIC Physics Fellow Program, with six positions in the first academic year, 1999-2000. This program had increased to include ten theorists and one experimentalist in academic year, 2001-2002. With five fellows having already graduated, the program presently has ten theorists and two experimentalists. Five of the ten RHIC Physics Fellows have been awarded/offered tenured positions, and this will be their final year in the program.

Beginning in 2001 a new RIKEN Spin Program (RSP) category was implemented at RBRC, presently comprising six RSP Researchers, and six RSP Research Associates. In addition, RBRC has four RIKEN Jr. Research Associates and one RBRC Young Researcher. A number of Visiting Scientists also contribute to the physics program at the Center.

RBRC also has an active workshop program on strong interaction physics with each workshop focused on a specific physics problem. Each workshop speaker is encouraged to select a few of the most important transparencies from his or her presentation, accompanied by a page of explanation. This material is collected at the end of the workshop by the organizer to form proceedings, which can therefore be available within a short time. To date there are fifty-five proceeding volumes available.

The construction of a 0.6 teraflops parallel processor, dedicated to lattice QCD, begun at the Center on February 19, 1998, was completed on August 28, 1998. A 10 teraflops QCDOC computer is under development and expected to be completed in JFY 2004.

**N. P. Samios, Director
November 2003**

CONTENTS

Preface to the Series	i
Introduction	
<i>N. P. Samios and T. D. Lee.....</i>	1
Meeting Agenda	3
RBRC Scientific Review Committee Membership	7
RBRC ADMINISTRATION PRESENTATIONS	
RIKEN BNL Research Center (RBRC)	
<i>N. P. Samios</i>	11
RBRC Theory Group	
<i>Larry McLerran</i>	39
Overview RBRC Experimental Group	
<i>Hideto En'yo</i>	49
EXPERIMENTAL PRESENTATIONS	
CCJ—The Computer Center in Japan	
<i>Yasushi Watanabe</i>	63
Polarimetry (AGS, RHIC, Jet)	
<i>Osamu Jinnouchi</i>	75
Forward Neutron Asymmetry and Longitudinal Polarization at PHENIX	
<i>Abhay Deshpande</i>	87
Central Arm Triggering for 2003 for PHENIX	
<i>Kensuke Okada</i>	101
Relative Luminosity for A_{LL} Measurements	
<i>Kiyoshi Tanida</i>	115
A_{LL} For π^0	
<i>Alexander Bazilevsky</i>	133
Belle Fragmentation Functions	
<i>Matthias Grosse Perdekamp</i>	143
J/Psi for pp and d-gold Collisions	
<i>Douglas Fields</i>	159

THEORY PRESENTATIONS (Continued)

Dark Matter and Pulsar Kicks from a Sterile Neutrino <i>Alexander Kusenko</i>	329
Charge Symmetry Breaking in Few-Nucleon Systems <i>Ubirajara van Kolck</i>	355
QCD and Dimensional Deconstruction <i>Mikhail Stephanov</i>	371
Kinetic Equation for ϕ^4 Theory with Generalized Kadanoff-Baym Ansatz <i>Takashi Ikeda</i>	381
QCDSF/QCDOC: Physics Results and Prospects/Project Status <i>Norman H. Christ</i>	391
<i>Robert Mawhinney</i>	411
RESEARCH SUMMARY - THEORY	
<i>Taku Izubuchi</i>	437
<i>Sangyong Jeon</i>	445
CURRICULA VITAE – SUMMARY	455
THEORY PUBLICATIONS	467
Additional RIKEN BNL Research Center Proceedings Volumes	515
Contact Information	

RBRC Scientific Review Committee Meeting

November 20-21, 2003

Brookhaven National Laboratory, Upton, NY 11973

The sixth evaluation of the RIKEN BNL Research Center (RBRC) took place on November 20-21, 2003, at Brookhaven National Laboratory. The present members of the Scientific Review Committee are Dr. Jean-Paul Blaizot, Professor Makoto Kobayashi, Dr. Akira Masaike, Professor Charles Young Prescott (absent), Professor Stephen Sharpe, and Professor Jack Sandweiss, Committee Chair. In order to illustrate the breadth and scope of the program, each member of the Center made a presentation on his research efforts. In addition, a special presentation was given jointly by our collaborators, Professors Norman Christ and Robert Mawhinney of Columbia University, on the progress and status of the RBRC QCDSF/QCDOC Supercomputer program. A demonstration of a 64-node (64 Gflops peak speed) QCDOC machine in action followed. Although the main purpose of this review is a report to RIKEN Management (Dr. Ryoji Noyori, RIKEN President) on the health, scientific value, management and future prospects of the Center, the RBRC management felt that a compendium of the scientific presentations are of sufficient quality and interest that they warrant a wider distribution. Therefore we have made this compilation and present it to the community for its information and enlightenment.

Thanks to Brookhaven National Laboratory and to the U. S. Department of Energy for providing the facilities to hold this meeting.

N. P. Samios & T. D. Lee

RIKEN BNL Research Center

Building 510A, Brookhaven National Laboratory, Upton, NY 11973 USA

11/12/03

**RBRC Scientific Review Committee (SRC) Meeting
Brookhaven National Laboratory, Upton, NY
Physics Department, Building 510, Open Sessions – Large Seminar Room
November 20-21, 2003
Draft Agenda**

Committee Members: Jean-Paul Blaizot, Makoto Kobayashi, Akira Masaike, Charles Prescott (Absent),
Stephen Sharpe, and Jack Sandweiss, Chair

Thursday, November 20, 2003

8:00 AM to 9:00 AM *Open Executive Session & Working Breakfast - (Room 2-160)*
(Presentations by RBRC Administration)

Large Seminar Room

9:00 AM to 10:15 AM EXPERIMENTAL PRESENTATIONS –GERRY BUNCE, CHAIR

9:00	CCJ—The Computer Center in Japan	Yasushi Watanabe
9:15	Polarimetry (AGS, RHIC, Jet)	Osamu Jinnouchi
9:30	Forward Neutron Asymmetry and Longitudinal Polarization at PHENIX	Abhay Deshpande
9:45	Central Arm Triggering for 2003 for PHENIX	Kensuke Okada
10:00	Relative Luminosity for A_{LL} Measurements	Kiyoshi Tanida
10:15 AM	Break	

10:45 AM to 12:15 PM EXPERIMENTAL PRESENTATIONS—HIDETO EN'YO, CHAIR

10:45	A_{LL} For π^0	Alexander Bazilevsky
11:00	Belle Fragmentation Functions	Matthias Grosse Perdekamp
11:15	J/Psi for pp and d-gold Collisions	Douglas Fields
11:30	π^0 Azimuthal Anisotropy in $\sqrt{s_{NN}} = 200$ GeV Au + Au Collisions	Masashi Kaneta
11:45	Electrons and Direct Photons	Yasuyuki Akiba
12:00	Vertex Detector for PHENIX	Yuji Goto
12:15 to 1:30 PM	SRC Executive Session (Committee Members) - Working Lunch (Room 2-160)	

Friday, November 21, 2003

8:00 AM to 8:30 AM

*SRC Executive Session (Committee Members) and Continental Breakfast
(Room 2-160)*

Large Seminar Room

8:30 AM to 9:30 AM

QCDSF/QCDOC: Physics Results and Prospects/Project Status
Norman H. Christ and Robert Mawhinney

9:30 AM to 10:00 AM

Demonstration of QCDOC

10:00 AM to 12:30 PM

Meetings with Individual RBRC Staff
Theorists (Room 2-160)
Experimentalists (Room 2-78)

(Possible Tours)

12:30 PM to 1:30 PM

*SRC Executive Session (Committee Members) and Working Lunch –
(Room 2-160)*

1:30 PM to 3:00 PM

Meetings with Individual RBRC Staff (Continued)
Theorists (Room 2-160)
Experimentalists (Room 2-78)

3:00 PM to 4:00 PM

SRC Executive Session (Committee Members) - (Room 2-160)

4:00 PM to 5:00 PM

Meeting with RBRC Administration - (Room 2-160)

5:00 PM

Adjourn

RBRC Scientific Review Committee Membership 2003

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RBRC ADMINISTRATION PRESENTATIONS

RIKEN BNL Research Center (RBRC)

Nicholas P. Samios

RBRC

1997 – 2003 In existence for 6 years

Director: T. D. Lee

Three Components: Theory, Experiment, Computing

RHIC

High Energy Density Nuclear Matter (RHIC)

Spin Component of Proton

Lattice QCD

Great Year (if not spectacular)

Previous Au x Au, p x p running

d x Au Run. Technical feat – asymmetric operation

Exciting results:

Multiplicities

Radial & Elliptic Flow

HBT size measurements

R_{AA} = Factor of 5 decrease

Jets

High y, small x effects

New Phenomena

**Hot, dense nuclear
matter**

Color glass condensate

Next run: Au x Au

Suppression of leading hadrons in RHIC data

$\sqrt{s_{nn}}=200$: Au+Au

$$R_{AA}(p_T) = \frac{d^2 N^{AA} / dp_T d\eta}{T_{AA} d^2 \sigma^{NN} / dp_T d\eta}$$

15

PHENIX:
Inclusive π^0 in central & peripheral collisions

STAR Charged hadrons:
Dependence on centrality

Anisotropy vs. p_T at 200 GeV

Significant anisotropy in Au+Au up to 7 GeV/c

A_{LL} : PHENIX vs theory

B.Jäger et al., PRD67, 054005 (2003)

RBRC

Director Emeritus:	T. D. Lee
Director:	N. P. Samios
Associate Director:	H. En'yo

Theory Group Leader (Designee):	L. McLerran
Deputy Group Leader:	A. Baltz

Experimental Group Leader:	H. En'yo
Deputy Group Leader:	G. Bunce

Theory Advisory Committee:

T. D. Lee
L. McLerran
A. Baltz
M. Creutz
M. Gyulassy
R. Pisarski (T. L. Trueman)

Experimental Advisory Committee:

A. Masaike
J. Sandweiss
S. Nagamiya

Workshops: 8
Publications: 95 Theory/33 Experiment
Seminars: 6/week
Other Articles: 2/8

- Volume 53 Theory Studies for Polarized pp Scattering (BNL-)*
August – September 2003
Organizers: S. Kretzer and W. Vogelsang
- Volume 54 RHIC Spin Collaboration Meetings XVII, XVIII, XIX (BNL-)
June 20, 2003, June 30, 2003, and August 18, 2003
Organizer: Brendan Fox
- Volume 55 Collective Flow and QGP Properties (In preparation)
November 17-19, 2003
Organizers: Nu Xu, E. Shuryak, P. Kolb, U. Heinz, S. Esumi, and S. Bass

*In press

THEORY PUBLICATION LIST

November 2002 to November 2003

RBRC-294 Konstantinos Orginos, (For the RBC Collaboration:
Y. Aoki, T. Blum, N. Christ, M. Creutz, C. Dawson,
T. Izubuchi, L. Levkova, X. Liao, G. Liu, R. Mawhinney,
Y. Nemoto, J. Noaki, S. Ohta, K. Orginos, S. Prelovsek,
S. Sasaki, and A. Soni, "Spin on the Lattice," to appear
in the *Proc. of the International Spin Physics
Symposium, SPIN 2002*, Upton, NY, USA,
September 9-14, 2002.

.....

RBRC-388 M. Hasenbusch, K. Jansen, D. Pleiter, H. Stüben,
P. Wegner, T. Wettig, and H. Wittig, "Benchmarking
Computer Platforms for Lattice QCD Applications,"
[hep-lat/0309149], to appear in Nucl. Phys. B
(Proc. Suppl.)

PROSPECTS OF THE GLUON POLARIZATION MEASUREMENT AT PHENIX, Y. Goto for the PHENIX Collaboration, SPIN 2002, AIP Conf. Proc. 675 (2003) 493-503.

FIRST POLARIZED PROTON COLLISIONS AT PHENIX, Y. Goto for the PHENIX Collaboration, DIS 2002, Acta Phys. Polon. B33 (2002) 3767-3771.

THE QUEST FOR NUCLEON SPIN, VERNON W. HUGHES; THE LAST DECADE: 1993-2003, A. Deshpande, Proceedings of "Symmetries and Spin 2003", Prague, to be published by World Scientific.

eRHIC: THE ELECTRON ION COLLIDER AT BNL, A. Deshpande, European Journal of Physics A, to be published in the Proceedings of International Meeting on Perspectives on Hadronic Physics, ICTP, Trieste, Italy, May 2003.

SPIN PHYSICS AT RHIC: PRESENT & FUTURE, A. Deshpande, to be published in PRAMANA Journal of Physics, Indian Academy of Sciences, Proceedings of QCD2002 Conference, IIT Kanpur, October 2002.

COMMISSIONING OF RHIC POLARIMETERS, A. Bravar, et al., BNL-Report-68976

TESTING OF QCD THROUGH SPIN OBSERVABLES IN NUCLEAR TARGETS, Workshop Summary of Experimental Talks, A. Deshpande, University of Virginia, April 20, 2003, World Scientific, Editors: D. Crabb, D. Day, J. Chen

MULTISTRANGE BARYON PRODUCTION IN AU-AU COLLISIONS AT $\sqrt{s_{NN}}=130$ GEV, Submitted to Phys. Rev. Lett, e-Print Archive: nucl-ex/0307024.

PHENIX DETECTOR OVERVIEW, PHENIX Collaboration (K. Adcox et al.), Nucl. Instrum. Meth. A499 (2003) 469-479.

PHENIX Magnet System, PHENIX Collaboration (S. H. Aronson et al.), Nucl. Instrum. Meth. A499 (2003) 480-488.

PHENIX Central Arm Particle ID Detectors, PHENIX Collaboration (M. Aizawa et al.), Nucl. Instrum. Meth. A499 (2003) 508-520.

PHENIX CALORIMETER, PHENIX Collaboration (L. Aphecetche et al.), Nucl. Instrum. Meth. A499 (2003) 521-536.

PHENIX Muon Arms, PHENIX Collaboration (H. Akikawa et al.), Nucl. Instrum. Meth. A499 (2003) 537-548.

PHENIX Inner Detectors, PHENIX Collaboration (M. Allen et al.), Nucl. Instrum. Meth. A499 (2003) 549-559.

Weekly Seminars

Spin Physics <i>(Theory & Exp)</i>	Tuesdays (10:00 a.m.)	Organized by Y. Goto W. Vogelsang A. Deshpande
Nuclear Physics	Tuesdays (11:00 a.m.)	Organized jointly by W. Vogelsang P. Petreczky with BNL Staff
High Energy-RIKEN Theory	Wednesdays (1:30 p.m.)	Organized jointly by S. Ohta with BNL Theorists
QCD and RHIC Physics <i>(Theory & Exp)</i>	Thursdays (12:30 p.m.)	Organized by W. Vogelsang
High Energy Theory Lunch Talks	Fridays (12:00 Noon)	Organized by S. Dawson
Nuclear Physics-RIKEN Theory	Fridays (2:00 p.m.)	Organized jointly by S. Kretzer, C. Dawson with BNL Staff

**The Outstanding Junior Investigator
Award in Nuclear Theory was
established by DOE in 2000. Since
Then, five RBRC members have
received the OJI Awards:**

Son, D. T.	(2000)
Stephanov, M.	(2001)
van Kolck, U.	(2001)
Schaefer, T.	(2002)
Bass, S.	(2003)

RHIC Physics Fellow Present Program	RBRC Dates	Position	New Position	Present Institution
van Kolck, Ubirajara	08/2000 - 08/2003 (Tenure)	RHIC Physics Fellow; U. of Arizona/Asst. Prof.	Tenured Associate Professor	RBRC/U. of Arizona
Vogelsang, Werner	4/1/2000 to 9/30/2003 10/1/2003 -	RIKEN BNL Fellow RHIC Physics Fellow BNL/Associate Physicist		RBRC/BNL
Wettig, Tilo	10/1999 -	RHIC Physics Fellow, Yale/Asst. Prof.	Associate Professor (2003)	RBRC/Yale
Experimentalists				
Fields, Douglas	09/2001	RHIC Physics Fellow, U. of New Mexico/Asst. Professor		RBRC/U. of New Mexico
Grosse Perdekamp, Matthias	01/1999 to 8/31/02 09/01/2002 -	RIKEN BNL Fellow RHIC Physics Fellow, U. of Illinois, Urbana Champaign/Asst. Prof.		RBRC/U. of Illinois Urbana Champaign

Name	Dates at RBRC	RBRC Position	New Position Title	New Position Institution
Saito, Naohito	[04/1995 to 03/1996] 04/1996 to 03/2001 04/2001 to 03/2002 03/2002 - present	[RIKEN Special Postdoc. Fellow] RIKEN/RBRC Res. RIKEN Spin Prog. Res RSP Visiting Scientist.	Associate Professor	Kyoto U., Japan
Sasaki, Shoichi	09/1998 to 09/2000	Research Associate	Assistant Professor	U. of Tokyo, Japan
Schaffner-Bielich, Jürgen	09/1998 to 10/2001	Research Associate	Research Associate	Columbia University University of Frankfurt
Son, Dam Thanh	10/1999 to 03/2002	RHIC Physics Fellow, Columbia U.	Associate Professor, Tenured	U. of Wash, Seattle, Natl. Inst. for Nucl. Th
Venugopalan, Rajagopal	10/1/2000 to 6/30/2003	RHIC Physics Fellow, BNL	Physicist with Tenure 6/1/02	BNL
Vogelsang, Werner	4/1/2000 to 9/30/2003 10/01/2003	RIKEN BNL Fellow	RHIC Physics Fellow/Associate Physicist	RBRC/BNL
Wingate, Matthew	09/1997 to 09/2000	Research Associate	Assistant Professor	Ohio State U., Columbus
Yasui, Yoshiaki	09/1997 to 09/2000	Research Associate	Postdoc Postdoctoral Fellow	Radiation Lab, RIKEN, KEK, Theory Group, Japan

N. P. Samios, RBRC Director
T. D. Lee, RBRC Director Emeritus
Hideto En'yo, RBRC Associate Director
* * * *

RBRC Research Scientists (2002 – 2003)
Experimental Group
Hideto En'yo, Group Leader
Gerry Bunce, Deputy Group Leader

Research Associates

Jinnouchi, O.

Kaneta, M.

Okada, K.

Fellows

Bazilevsky, A. (4/1/03)*

Deshpande, A.

Fox, B. (10/15/03)*

Messer, F. (4/1/03)*

Tenure Track/RHIC Fellow

Fields, D. (UNM)

Grosse-Perdekamp, M. (UIUC)

RIKEN Spin Program

(RSP) Researchers

Akiba, Y.

Goto, Y.

Ichihara, T.

Taketani, A.

Tanida, K.

Watanabe, Y.

RIKEN Spin Program

(RSP) Research Associates

Kobayashi, H. (9/30/03)*

Yokkaichi, S.

Tabaru, T.

Torii, H.

RIKEN Spin Program/

Visiting Scientists

Kurita, K.

Lange, J. S.

Li, Zheng

Ogawa, A.

Saito, N.

Advisory Committee

Experiment

Masaike, A.

Nagamiya, S.

Sandweiss, J.

Consultants

Jaffe, R.

Roser, T.

Makdisi, Y.

Tannenbaum, M.

RBRC Young

Researchers

Siegle, V. (3/31/03)

RIKEN Jr. Res. Assoc.

Fukao, Y.

Hachiya, T.

Horaguchi, T.

Visiting Res. Assoc.

Kiyomichi, A.

Sato, H.

Tojo, J.

RIKEN Researchers

Okamura, A.

RIKEN Young Res.

Kamihara, N.

Okada, H.

Togawa, M.

Technical Collaborator

Hiejima, H. (UIUC)

RBRC Theory Group

Larry McLerran

Theoretical Physics at RIKEN–BNL Center

Strong Interactions and QCD

**How are strongly interacting particles
made from fundamental constituents?**

**How do the fundamental interactions of
QCD produce quark confinement and
mass?**

**What is the behaviour of strongly
interacting matter in bulk?**

Nuclear Matter $\rightarrow$ Quark Gluon Plasma

**What is the physics beyond the standard
model?**

Tests of CKM matrix.

All issues intertwined!

Require understanding and computation

Current Interests in RIKEN–BNL Center

Lattice Gauge Theory

Masses and matrix elements of hadrons

CKM matrix

Properties of QGP and hadronic matter

Color Glass Condensate and Quark Gluon Plasma

RHIC Phenomenology

Everything for $x \leq 10^{-2}$

Structure of Hadrons

Origin of spin

Quark and gluon distribution functions

Perturbative QCD: RHIC and LHC

Matter at High Density

Energy densities high
Matter strongly interacts

High Density $\rightarrow$
Short Distances

Probes interesting
phase space

RIKEN-BNL Physics Fellows

Graduates

D. Bodeker Tenured C-4 Bielefeld
D. Kharzeev Tenured BNL
D. Rischke Tenured C-4 Frankfurt
D. Son Tenured U of Washington
R. Venugopalan Tenured BNL
T. Wettig Tenured C-3 Regensburg

Current Fellows

S. Bass Duke
T. Izubuchi Kanazawa
S. Jeon McGill
A. Kusenko UCLA (Tenure in process.)
T. Schaefer N. Carolina (Tenure in process)
M. Stephanov U of Ill. Chicago (Tenure in process)
U. van Kolck U of Arizona (Tenure in process)
T. Blum U of Connecticut
W. Vogelang BNL
P. Petreczky BNL

All at good research universities.

Sense of Community!

All succesful.

Overview RBRC Experimental Group

Hideto En'yo

Overview

RBRC Experimental Group

RBRC Review, 2003

Hideto Enyo

AGS warm snake, made in Japan

UPGRADE@CERN

Hiroaki Ohnishi

Johann Hauser

Hiroyuki Kano

Experiment G. Organization JFY03

RIKEN@RBRC

Yasuyuki Akiba

Yuji Goto

Hideyuki Kobayashi

Takahiro Kawabata

Tsuguchika Tabaru

RBRC

Hideto Enyo

Gerry Bunce

Matthias Grosse Perdekamp

Abhay L Deshpande

Douglas E. Field

Brendan Fox

Federica Messor

Alexander Bazilevsky

Kensuke Okada

Osamu Jin-nouchi

Masashi Kaneta

Naohito Saito (Visitor)

RBRC STUDENTS

Hiroki Sato

Viktor Siegel

Nobuyuki Kamihara

Takuma Horaguchi

Manabu Togawa

Yoshi Fukao

Hiromi Okada

Takashi Hachiya

BELLE

Soeren Lange (Visitor)

Akio Ogawa (Visitor)

Kazumi Hasuko

RIKEN/Wako

Masahiro Okamura

Kiyoshi Tanida

Atsushi Taketani

Rykov Vladimir

Junji. Tojo

Hisa Torii

Tetsuya Chujo

WAKO/CCJ

Takashi Ichihara

Yasushi Watanabe

Satoshi Yokkaichi

Akio Kiyomichi

**~40 active member
+~3 active visitors**

Students in 2003 (Students!)

- Longtime issue: RIKEN could support students only when they were Junior Research Associate or JSPS fellow.
- Hiroki Sato, now PhD, AGS-RHIC Thesis Award, JSPS Fellow
- Established the RBRC Young Researcher
- (Many Thanks to Review Committee)
 - Viktor Siegle (Heidelberg) =>Finishing his diploma at home
 - Yoshi Fukao (Kyoto-U, RBRC-YR =>JRA from April 03)
 - Takuma Horaguchi (TITech, RBRC-YR=>JRA from April 03)
 - Manabu Togawa (Kyoto-U JRA from April 02)
 - Hiromi Okada (Kyoto-U, became JRA2 from April 03)
 - Takashi Hachiya (Hiroshima-U, became JRA2 from April 03)
 - Kazuya Aoki (Kyoto-U, will be this year's candidate)
 - Theory JRA (Hatta (Kyoto) from 02, Yokoya(Hiroshima) 03)

Special Thanks to
Administrative Staffs
in BNL and Wako

Especially to M.Yokota

EXPERIMENTAL PRESENTATIONS

CCJ—The Computer Center in Japan

Yasushi Watanabe

Status

Working Center in Japan for spin physics at RHIC

Y. Watanabe, T. Ichihara, S. Yokkaichi, A. Kiyomichi,
O. Jinnouchi, H. En'yo, Y. Goto, H. Hamagaki⁽¹⁾
RIKEN, RBRC, CNS⁽¹⁾

Presented on 20th November 2003 at RBRC Review at BNL

CCJ growing history

	Feb-99	Dec-99	May-00	Nov-00	May-01	Oct-01	Feb-02	Mar-03
CPU (unit)	16	64	96	160	192	224	332	332
CPU (GHz)	8	34	56	136	168	200	374	378
Disk (0.1 TB)	3	19	35	92	152	212	252	332
HPSS (TB)	100	100	100	100	100	100	350	600

Current configuration of CCJ

StorageTek Tape Robot
Capacity: >1000TB
Installed 3000 tapes (600TB)

HPSS Server (IBM RS-6000/9

Analysis for the PHENIX experiment with CCJ

■ Official productions and simulations

- Study of the conversion and Albedo stuffs in the $pp \rightarrow \pi^0$ analysis
- Run2pp DST production
- Run3pp filtered data fast analysis
- run3pp data refiltering + spin-nanoDST + gamma3 process
- PbPb data analysis for the π^0 asymmetry fast analysis
- Run2-pp π^0 A_N analysis

■ Major analysis projects

- Analysis for silicon strip detector in PHENIX Upgrades - J. Tojo
- Searching for Non-hadronic sources of photons in Au+Au Collisions at $\sqrt{s_{NN}} = 200 \text{ GeV}$ at RHIC-PHENIX - T. Sakaguchi
- Identified Hadron Spectra and Yields at 200 GeV AuAu Collisions from the PHENIX experiment at RHIC - T. Chujo
- Reproduction of SpinLumi root files for analysis of BBC single spin asymmetry with ZDC coincidence - K. Tanida
- Pico DST production for π^0 anisotropy analysis from micro DST of 200 GeV Au+Au data - M. Kaneta
- Single particle MC simulation for hadron spectra(run2 pp) - S. Kaminaga
- 2nd scan of π^0 anisotropy analysis - M. Kaneta
- Monte-Carlo event generation for the study of dimuon production at CERN-SPS - H. Ohnishi
- Run#91720 MUIDS1D triggered events production - K. Aoki
- Direct photon measurement in RHIC-PHENIX - T. Horaguchi
- Simulation study for the Aerogel Cherenkov Counter - S. Takagi
- Belle simulation - K. Hasuko, A. Ogawa, M. G. Perdekamp

Plan and Current Status of RIKEN CCJ

Polarimetry (AGS, RHIC, Jet)

Osamu Jinnouchi

Polarimetry in RHIC, AGS and $\vec{H}$ Jet

Osamu Jinnouchi

(RIKEN BNL Research Center)

Outline

Brief Introduction

Outcomes from Run-03, RHIC & AGS

Towards the next run

Detector/Target and Event selection

79

non-relativistic kinematics

$$tof = \sqrt{\frac{M_C}{2T_{kin}}} L$$

RHIC SSDs = ■ + ■

6-Detectors : sensitive to X,Y
15cm radius

AGS SSDs = ■
2-detectors
25cm radius

AGS Polarization in Run-03

Main source of low polarization in RHIC Run-02 was due to large depolarization in AGS acceleration

Extremely important to know the polarization during the ramp to understand depolarization mechanism

Full spin flip at all imperfection resonances using a solenoidal partial Siberian snake

Full spin flip at strong intrinsic resonances using an rf AC dipole (spin flipper)

Remaining polarization losses are from coupling and weak intrinsic resonances

imperfection resonances every integer

RHIC : Polarization measurements in typical fill

- Measurements for every 2 hours (24GeV,100GeV)
- Long beam lifetime, stable polarization
- Spin rotators keep the polarization
- Polarization in store energy was up from run02
 - BLUE 29% YELLOW 25% in average

Summary & Outlook for Run04

First successful result from ramp measurement in AGS with using the CNI polarimeter as in RHIC

- Successful commissioning of Spin-rotator is confirmed
- Improved polarization in RHIC

- Add a 5% warm helical snake (run '04)

To avoid all depolarization build a strong (20%) superconducting helical Siberian snake (2005-2006)

Commissioning of H jet target in 2004 recoil detectors

Forward Neutron Asymmetry and Longitudinal Polarization at PHENIX

Abhay Deshpande

**Measurement of single transverse spin
asymmetry A_N in forward neutron
production at PHENIX**

**(Experimental confirmation of longitudinally
polarized proton beams at PHENIX)**

**Abhay Deshpande
RIKEN BNL Research Center**

Relativistic Heavy Ion Collider

16

RHIC accelerates heavy ions to 100 GeV/A
and polarized protons to 100 GeV (250 GeV in future)

Local Polarimeter at IP12 2001/2

- RUN-2: Neutron asymmetry observed in IP12 while testing a local polarimeter designed to look for π^0 , γ asymmetries:*

- “Left-Right” asymmetry measured for different slices in Φ :*

Abhay Deshpande
RBRC Review 2003

PHENIX Local Polarimeter 2003

Zero Degree Calorimeter (ZDC)

Tungsten-Fiber Sandwich

$5.1 \lambda_T, 149 X_0 \rightarrow$ 3 ZDCs each side

Shower Max Detector (SMD)

Plastic Scint. strips
 $\rightarrow (0.5\text{cm} \times 3)$

10 and 15 cm long

7 x 1.5cm columns

8 x 1.5cm rows

Read by
 Hamamatsu M16

Spin Rotator Commissioning (I)

Max at $\Phi = \pm \pi/2 \rightarrow$ Transverse

Max at $\Phi = 0 \rightarrow$ Transverse Radial ! Spin Rot PS wired wrong!

How longitudinal were the beams?

$$S_L = \sqrt{1 - S_T^2}, \quad S_T = \sqrt{S_{T\text{-vertical}}^2 + S_{T\text{-radial}}^2}$$

PHENIX Local
Polarimeter measures S_T

Left-Right vs. Fill No.

Up-Down vs. Fill No.

$$S_L(\text{blue}) = 99.3^{+0.5}_{-1.4} \quad {}^{+0.0}_{-0.9}$$

$$S_L(\text{yellow}) = 97.4^{+1.3}_{-3.2} \quad {}^{+0.1}_{-0.9}$$

Central Arm Triggering for 2003 for PHENIX

Kensuke Okada

PHENIX Central Arm Trigger in RUN3

Kensuke Okada
RBRC review
2003

Concept of EMCal-RICH trigger

PHENIX central arm
2 types of EMCal (PbSc, PbGl)

- $\gamma(\pi^0)$: EMCal
- Electron : EMCal and RICH
- $h^\pm$: EMCal and RICH (through hadronic interactions)

Run2 $\rightarrow$ Run3 (The second year of ERT)

Live area :

EMCal from 75% to 90~95%

We improved it with 2 steps

A) Replacement of bad FEM boards (Nov, Dec 2002)

B) Mask bad tiles instead of whole SM. (Apr 2003)

RICH from 15% to 70%

(In pair analysis for J/ψ , the efficiency is 90%.)

4x4 overlapping between super modules

4x4 turn-on width improvement

Trigger Circuit “variation”

The first assumption was it came from the deviation of the pedestal.

→ The pedestal scan was performed to check electrical zero level.
(It was independent to HV on/off)

Threshold setting was adjusted for each trigger tile.

($\sim \pm 1.5$ for 4x4, $\sim \pm 2$ for 2x2, FWHM)

	Trigger	Old RF	New RF	
Rejection power improvement	2x2 (800 MeV)	29.5	35	
	4x4a (2.8 GeV)	863	1989	← RF of 4x4 triggers
	4x4b (3.5 GeV)	3762	5197	← improved for more
	4x4c (2.1 GeV)	190	606	← than 2 times.
	e (2x2 $\otimes$ RICH)	977	1077	

Table 1: Rejection factors with old and new threshold settings.

Physics Results using ERT

Run2

π^0

"Midrapidity Neutral Pion Production in Proton-Proton Collisions at $\sqrt{s} = 200$ GeV"

S.S. Adler et al, accepted for publication in PRL on 19 September 2003

Preprint [http://arxiv.org/abs/hep-ex/0307001](#)

J/ψ

"J/psi production from proton-proton collisions at $\sqrt{s} = 200$ GeV"

S.S. Adler et al., submitted to Physical Review Letters July 8 2003

Preprint [http://arxiv.org/abs/hep-ex/0307002](#)

III

Run3

π^0

"Absence of Suppression in Particle Production at Large Transverse Momentum in $\sqrt{s_{NN}} = 200$ GeV d+Au Collisions"

S.S. Adler et al.,

Was

How to evaluate the threshold value?

If the time variation scale is a week or so.

Perform a calibration scan once a week.

Feed back to the new setting,

or

Record them for the offline-analysis.

Method	Time (hour) /tile	Comment
Pedestal scan	0.3 (Arcnet download)	Wide pitch, far from real threshold
Laser scan	~1 (rate, polarizer)	Only for PbSc 4x4
HV scan with LED for PbGl	~1 (HV setting time)	LED light length?
CALDAC scan	~1 (Arcnet download)	Good for PbGl, PbSc

It has to be multiplied by 36 to scan all.

Anyway it needs a speed-up for a regular calibration.

Relative Luminosity for A_{LL} Measurements

Kiyoshi Tanida

Relative luminosity for A_{LL} measurements – in PHENIX run3

K. Tanida (RIKEN/RBRC)

11/20/03 RBRC review

Outline

- What is relative luminosity? Why?
- How to get relative luminosity?
- Method of accuracy estimation
- Results and ZDC/BBC ΔA_{LL}
- Summary and outlook

Our goal in PHENIX run3

- $\delta A_{LL} \sim \delta R / (2P_B P_Y) < 3 \times 10^{-3}$
(c.f. expected A_{LL} for pions)
- Beam polarization $\sim 27\%$

- Relative luminosity accuracy:
 $\delta R < 4 \times 10^{-4}$

Method of accuracy estimation

- Cross-check by ZDCCLL1 trigger
(Zero Degree Calorimeter)
 - low backbround ($\sim 10^{-4}$)
 - statistics: moderate ($\sim 3\%$ of BBCLL1)
- $N(\text{ZDCCLL1})/N(\text{BBCCLL1})$ must be constant
- Crossing-by-crossing, fill-by-fill ratio calculation
 Blue: + - + - + - + -
 Yellow: + - - - + - - -
 → We have all spin combinations in each fill
- GL1p scaler counts both BBCCLL1 and ZDCCLL1 for each bunch
- Required accuracy: constant to 10^{-3} level for each fill.

Possible explanation – detector acceptance

- Extreme cases
 - Flat acceptance
→ $N \propto \text{area}$
 - Narrow acceptance
→ $N \propto \text{height}$

- ZDCLL1 has flat acceptance, while BBCLL1 has slightly center-peaked acceptance distribution

After correction

Relative luminosity accuracy $< 0.09\%$ (stat. limited)

Actual accuracy estimation

-- Bunch fitting

- Crossing-by-crossing analysis, taking care of possible A_{LL} effect
- For i-th crossing, calculate
$$r(i) = N_{ZDCLL1}(i)/N_{BBCLL1}(i)$$
corrected for vertex width
- Fit $r(i)$ by
$$r(i) = C[1 + A_{LL}P_B(i)P_Y(i)]$$
fit parameter: C, A_{LL}
- You will get C, A_{LL} , their errors, and χ^2
 - accuracy is obtained as $\delta A_{LL} (x P_B P_Y)$
 - χ^2 gives a good check for systematic errors.

Result & BBC/ZDC A_{LL}

- Systematic error seen $\rightarrow \chi^2$ correction required
- If χ^2/dof is larger than 1, then enlarge statistical error by $\sqrt{\chi^2/\text{dof}}$ – conservative correction
- Averaged result: $A_{LL} = (2.3 \pm 1.6) \times 10^{-3}$
- This corresponds to relative luminosity accuracy of $\delta R = 2.5 \times 10^{-4}$
 $\rightarrow$ Achieved the goal
- ZDC-BBC A_{LL} is 0 consistent
 $\rightarrow$ we can use both for reference of $A_{LL} = 0$

Summary and outlook

- Relative luminosity was calculated by using BBCLL1 trigger counts and were checked using ZDCCLL1 trigger.
- Accuracy: 2.5×10^{-4} achieved ($\delta A_{LL} = 1.6 \times 10^{-3}$) after vertex width correction.
 - error is dominated by ZDC statistics
- Outlook:
 - full accuracy estimation using real vertex distribution
 - different relative luminosity necessary for different reaction channels (acceptance correction).
 - study of multiple collision effect for future high-luminosity runs

A_{LL} for π^0

Alexander Bazilevsky

First π^0 A_{LL} results from PHENIX

RBRC Scientific Review Committee Meeting
BNL, November 20, 2003

A. Bazilevsky

PHENIX

A_{LL}

"Blue" beam

$$A_{LL} = \frac{\sigma_{++} - \sigma_{+-}}{\sigma_{++} + \sigma_{+-}} = \frac{1}{|P_B P_Y|} \frac{N_{++}/L_{++} - N_{+-}/L_{+-}}{N_{++}/L_{++} + N_{+-}/L_{+-}}$$

++ same helicity

+- opposite helicity

(P) Polarization

(L) Relative Luminosity

(N) Number of π^0 s

PHENIX

Relative Luminosity: Results

See talk by K. Tanida

Beam Beam Counter (BBC) used
Zero Degree Calorimeter (ZDC) as a cross check

Achieved relative luminosity precision $\delta R = 2.5 \cdot 10^{-4}$

Pessimistic estimation limited by ZDC statistics (30 times less than BBC statistics used in Rel. Lum. measurements)

A_{LL} of BBC relative to ZDC consistent with 0 ($< 0.2\%$)

Strong indication that both A_{LL} s are zero (very different kinematical regions, different physics signals)

π^0 Data Set

See talk by K. Okada

$\sqrt{s} = 200$ GeV data collected with high p_T photon trigger

Based on EMCal; Threshold ~ 1.4 GeV/c

Rejection factor ~ 110

Analyzed data sample:

42.7M events ($\sim 0.215 \text{ pb}^{-1}$)

$\text{sqrt}(\langle P_b P_y \rangle) \sim 26\%$

Minimum Bias data

To obtain "unbiased" π^0 cross section at low p_T

For high p_T photon trigger efficiency study

PHENIX

Pi0 counting for A_{LL}

N_{π^0} :
±25 MeV/c² around π^0 peak (and
also ±15 and ±35 MeV/c² for
cross checks)

N_{bck1} :
Two 50 MeV/c² wide areas
adjacent to π^0 peak

N_{bck2} :
250 MeV/c² wide area between
 π^0 and η peaks

N_{π^0} and N_{bck} accumulated statistics

pt GeV/c	N_{π^0} 15 MeV/c²	N_{π^0} 25 MeV/c²	N_{π^0} 35 MeV/c²	N_{bck1}	N_{bck2}
1-2	1278k	1777k	2129k	1470k	3478k
2-3	874k	1039k	1146k	335k	989k
3-4	176k	201k	208k	27k	85k
4-5	34k	38k	39k	3.9k	12k

PH ENIX

A_{LL} measurements

$$A_{LL} = \frac{\sigma_{++} - \sigma_{+-}}{\sigma_{++} + \sigma_{+-}} = \frac{1}{|P_R P_Y|} \frac{N_{++}/L_{++} - N_{+-}/L_{+-}}{N_{++}/L_{++} + N_{+-}/L_{+-}}, \quad \delta_{A_{LL}} \approx \frac{1}{|P_R P_Y|} \frac{1}{\sqrt{N_{++} + N_{+-}}}$$

++ same helicity

+- opposite helicity

Procedure

1. Collect N and L for ++ and +- configurations (sum over all crossings) and calculate A_{LL} for each fill (a few hours of running, constant beam polarization)
2. δA_{LL} is based only on N_{π^0} stat. error
3. Average A_{LL} over fills; use χ^2/NDF to control fit quality; use "bunch shuffling" to check syst. errors

PH ENIX

Cross Checks

- Different mass window cut for N_{jet}
15, 25 and 35 MeV/c² cuts
All results are consistent
- Different photon PID
All results are consistent
- Parity violating A_{LL} 's
++ vs -- and +- vs -+
All are consistent with 0 within 1.5 σ (as expected)
- Parity violating single spin asymmetries A_T
All are consistent with 0 within 1.5 σ (as expected)

PHENIX

$\pi^0 A_{LL}$ from pp at 200 GeV

p_T GeV/c	$A_{LL}^{\pi^0+bck}$ (r_{bck})	A_{LL}^{bck}	$A_{LL}^{\pi^0}$ (Background subtracted)
1-2	-0.028 ± 0.012 (45%)	-0.006 ± 0.014	-0.046 ± 0.025
2-3	-0.022 ± 0.015 (72%)	-0.035 ± 0.027	-0.019 ± 0.019
3-4	-0.003 ± 0.033 (77%)	0.094 ± 0.092	-0.009 ± 0.036
4-5	-0.023 ± 0.074 (5%)	0.38 ± 0.24	-0.045 ± 0.079

Polarization scaling error $\delta P \sim 30\%$: is not included

- ✓ Enters the A_{LL} quadratically
- ✓ Analyzing power $A_N(100 \text{ GeV}) \sim A_N(22 \text{ GeV})$ is assumed
- ✓ $\delta P \sim 30\%$: combined stat. and syst. error for $A_N(22 \text{ GeV})$ (AGS E950)

Relative luminosity contribution to $\pi^0 A_{LL}$ error is $< 0.2\%$

p_T smearing correction is not included

Belle Fragmentation Functions
Matthias Grosse Perdekamp

Prospects of Spin Dependent Fragmentation Function Measurements at Belle

RBRC Review

RIKEN BNL Research Center, Upton, New York, November 20-21, 2003

M. Grosse Perdekamp (UIUC and RBRC)

K. Hasuko (RIKEN/RBRC)

S. Lange (Frankfurt/RBRC)

A. Ogawa (BNL/RBRC) and V. Siegle (RBRC)
for the Belle Collaboration

Motivation: Transversity at RHIC I

Physics Channels

$$\int L dt = 320 \text{ pb}^{-1}, \sqrt{s} = 200 \text{ GeV}$$

- Collins Effect in Jets:

$$A_T(pp_{\perp} \rightarrow \pi + \text{Jet} + X)$$

J.C. Collins, Nucl. Phys. B396, 161(1993)

- π^+, π^- Interference Fragmentation:

$$A_T(p_{\perp} p \rightarrow (\pi^+, \pi^-) + X)$$

J. Collins, S. Heppelmann, G. Ladinsky,
Nucl.Phys. B420 (1994)565

R. Jaffe, X.Jin, J. Tang Phys. Rev. D57 (1999)5920

- Inclusive jet production

$$5 \cdot 10^{-4} \leq A_{TT} \leq 3 \cdot 10^{-3}$$

W. Vogelsang and M. Stratmann, RBRC
Workshop on Transversity (2000)

STAR

PHENIX

$$\text{Observables} \propto \delta q \cdot H$$

($O \propto \text{distribution - func.} \times \text{fragmentation func.}$)

Example: Interference Fragmentation in PHENIX

$$A_T(p_{\perp} \rightarrow \pi^+ \pi^- + X) \text{ vs } p_T^{\text{pair}}$$

Statistical sensitivity for 32 pb^{-1} (1 week of running)

KEKB and the Belle Experiment

- KEBB
 - Asymmetric collider
 - $8\text{GeV } e^- + 3.5\text{GeV } e^+$
 - $\sqrt{s} = 10.58\text{GeV (Y(4S))}$
 - Beam size: $100 \mu\text{m}(\sigma_x) \times 3 \mu\text{m}(\sigma_y)$
 - Crossing angle: 11mrad

Belle Detector Performance

Tracking part :

Central Drift Chamber(CDC)
(50 layers)

$$\sigma_{E/\eta} \sim 0.35\% \text{ @ } 1 \text{ GeV/e}$$

Silicon Vertex Detector(SVD)

(3 silicon layers double-sided)

$$\sigma \sim 55 \mu\text{m for } 1 \text{ GeV/e @ } 90^\circ$$

Particle identification (K^0 id up to 3.5 GeV/c)

Aerogel Cherenkov Counter(ACC) : $n = 1.01 \sim 1.03$

Time of Flight (TOF) $\sigma \sim 95 \text{ ps}$

CDC (dE/dx) : $\sigma_{dE/dx} \sim 7\%$

Electromagnetic Calorimeter(Csl(Tl) crystals) $\sigma_E/E \sim 1.8\% \text{ @ } 1 \text{ GeV}$

K_L and Muon detector (KLM = 14 layers) $\mu : \epsilon > 90\%$ and fake rate $\sim 2\%$

Collins Fragmentation Function

- $e^+e^- \rightarrow \pi_{\text{jet1}}^+ \pi_{\text{jet2}}^- X$
- Reaction plane defined with beam (z-axis) and jet axis
- Product (π) plane defined with π and jet axis
- ϕ : angle between the planes
- $A \propto H_1^\perp(z_1) H_1^\perp(z_2) \cos(\phi_1 + \phi_2)$
- Can analyze with/without using jet axis
 $A \propto H_1^\perp(z_1) H_1^\perp(z_2) \cos(2\phi)$

Analysis Method

- Reproduce spin-averaged FF from data
 - To estimate error sizes
 - Precision input to NLO pQCD FF-analysis
- Use off-resonance data (also can use on-resonance data)
- Select light-quark events
 - Event-shape cuts
 - charm tag: vertexing, mass, etc.
 - bottom tag: in on-resonance; easy
- Require two-jet events
 - Define fragmentation functions inclusively
- Study azimuthal distribution as functions of z , k_T
 - Ratio between different analyzing powers
Data/(MC w/o spin-dependent FF), unlike-sign/like-sign, etc.
 - Use τ productions for cross checks

Available Statistics (off-resonance)

- #events $\sim 27\text{M}$ (off-resonance data)
- Event selection
 - Thrust > 0.85
 - $\pi^+\pi^-$ in different hemispheres (for Collins function)
 - 8 bins in 2ϕ ($\phi_1+\phi_2$)
- $\delta A = \delta(H_T^\perp H_T^\perp) \cos(\phi_1+\phi_2)$
= 1.4% of statistical error
for high z-bin (> 0.5)
- More data available on-resonance

J/Psi for pp and d-gold Collisions

Douglas Fields

$J/\Psi \rightarrow \mu\mu$ in $\sqrt{s} = 200$ GeV p-p and d-Au Collisions

Douglas E. Fields

University of New Mexico/RIKEN-BNL Research Center

Detector Overview

Radiographs of active area for South arm in run-II. (Three stations with 3 or 2 gaps each)

Vast improvement for present (run-III) after extensive repairs to electronics and HV during shutdown

Newly installed North Arm is working well

Preliminary Run3 Results

Luminosities in this analysis (so far)

<u>Run</u>	<u>Arm</u>	<u>Trig</u>	<u>nDst</u> <u>Trigs</u>	<u>RC</u> <u>Trigs</u> (for runs used)	<u>BBCLL</u> <u>1</u> (x10 ⁹)	<u>BBCLL1</u> <u>Eff</u> (x10 ⁹)	<u>Lum</u> (nb ⁻¹)
dAu	N	2D	1088671 (96%)	1139145	1.816	1.736	0.873
		1D1S	3350008 (91%)	3696734	2.313	2.096	1.054
	S	2D	3462112 (94%)	3685783	1.786	1.678	0.844
		1D1S	0	0	0	0	Soon: ~1
pp	N	1D1S	2902091 (90%)	3210357	3.103	2.797	125
	S	1D1S	938597 (98%)	953267	3.418	3.365	151

Nominal Run3

Luminosities:

- dAu – 2.74 nb⁻¹
- pp – 352 nb⁻¹

~ 20 to 40%

More expected

67nb⁻¹ in Run2

Summary

- Both North and South Arms installed and performing well.
- Problems in previous runs with MuID trigger should be better this run with tunnel shielding and new LL1.
- Run3 analysis of J/Ψ approaching completion for both pp and d-Au runs.
- Challenge for Au-Au.
- More luminosity for pp - spin physics with muon arms!

π^0 Azimuthal Anisotropy in $\sqrt{s_{\text{NN}}} = 200 \text{ GeV Au} + \text{Au}$ Collisions

Masashi Kaneta

π^0 v_2 study
in $\sqrt{s_{NN}} = 200\text{GeV}$
Au+Au collisions

KANETA, Masashi

金田 雅司

for the PHENIX Collaboration

RIKEN-BNL Research Center

Event anisotropy in high p_T

- Hadron production in high p_T come from jet production
 - jet itself has NO event anisotropy for azimuthal direction
- Jet quench will make an event anisotropy due to different energy loss in different direction
 - for example M.Gyullasy, I.Vitev, X.N.Wong, PRL86 (2001) 2537

Method of π^0 v_2 Measurement

- Define reaction plane by charged multiplicity on Beam-Beam Counters
- π^0 reconstruction from gamma measured by Electro-Magnetic Calorimeter (EMC)
 - For each p_T , azimuthal angle, centrality
- Combine both information
 - Counting number of π^0 as a function of

$$E \frac{dN^3}{d^3 p} = \frac{1}{2\pi} \frac{d^2 N}{p_T dp_T dy} \left(1 + \sum_{n=1}^{\infty} 2 \underbrace{v_n^{\text{measured}}}_{\text{event anisotropy parameter measured}} \cos[n(\underbrace{\phi}_{\text{azimuthal angle of the particle}} - \underbrace{\Psi_r}_{\text{reaction plane angle}})] \right) \text{ where } n = 1, 2, 3, \dots$$

$$v_n^{\text{real}} = v_n^{\text{measured}} / (\text{reaction plane resolution})_n$$

Note: the detail of reaction plane definition will be found in **nucl-ex/0305013**

π^0 Identification

- Requirement for photon

- Dead and noisy EMC towers are removed for the analysis
- Energy of shower on EMC > 0.040 [GeV]
- PID cuts: $\chi^2 < 3$ for photon probability
- $|\text{TOF}| < 1.2$ [ns]
- Note: those cut are looser than single p_T distribution analysis

- For π^0

- Asymmetry cut: $|E_1 - E_2| / (E_1 + E_2) < 0.8$
- Combinatorial background is estimated by event mixing
 - Classes categorized for event mixing
 - centrality : every 10%
 - BBC Z Vertex : every 10cm in ± 30 cm
 - reaction plane direction in PHENIX detector : 24 bins in $\pm \pi$

One of difficulties (?) is..

- The analysis of π^0 v_2
 - Counting number of π^0 as a function of centrality, p_T and ϕ - Φ_R
 - The analyses (fitting of combinatorial background, its subtraction, counting π^0 and so on) are done by macro automatically
 - However, we need to think it is not perfect, that is, it may fail sometime
- Therefore,

V_2 vs. p_T vs. Centrality from 200GeV Au+Au

PHENIX PRELIMINARY

EXPERIMENTAL DATA

Statistical error is shown by error bar
Systematic error from π^0 count method and reaction plane determination is shown by gray box

v_2 VS. p_T (Minimum Bias) from 200GeV Au+Au

- Identified particle v_2 up to $p_T=10\text{GeV}/c$

36.3×10^6 [events] = $5.3^{+0.5}_{-0.4}$ [$(\mu\text{b})^{-1}$]

v_2 VS. p_T (Minimum Bias) from 200GeV Au+Au

PHENIX Preliminary Results, 2004-2005 Au+Au collisions at 200 GeV

- Identified particle v_2 up to $p_T=10\text{GeV}/c$

36.3×10^6 [events] = $5.3^{+0.5}_{-0.4}$ [$(\mu\text{b})^{-1}$]

Quark Coherence?

- Phys. Rev. Lett. 91 (2003) 092301, D.Molnar and S.A. Voloshin
- $q\bar{q} \rightarrow \text{meson}$, $qqq(\bar{q}\bar{q}\bar{q}) \rightarrow \text{Baryon}$

$$v_{2,M}(p_{\perp}) \approx 2v_{2,q}\left(\frac{p_{\perp}}{2}\right), \quad v_{2,B}(p_{\perp}) \approx 3v_{2,q}\left(\frac{p_{\perp}}{3}\right),$$

- How data looks like?
- Non-strange and strange meson and baryon seems to be merged around $p_T/n_{\text{quark}} \approx 1-3 \text{ GeV}/c$
- But we need more statistics to conclude it

Comparison with a model

Special thanks to
C. Nonaka (one of authors)
of nucl-th/0306027 for
data of model calculation

Comparison with a model which is described in nucl-th/0306027. Here we don't want to discuss which model can describe the data. To conclude which model can describe the data, we need much more statistics in high p_T region.

Outlook

• Feature plan of analysis

- Using high p_T gamma trigger in run2 Au+Au data
 - We will have about twice statistics in high p_T
 - need to study trigger bias
 - therefore, present analysis results are from minimum bias trigger events
- η v_2 is also available by same method
- PHENIX has photon v_2 also (STAR also, but not opened yet)
- photon v_2 , especially low p_T !

• RHIC run4 Au+Au, it will be

- Much more statistics
 - Detail study of v_2 shape around $p_T=2-4\text{GeV}/c$
- Much higher p_T
 - We want to know where is the end of finite v_2 in very high p_T

Electrons and Direct Photons

Yasuyuki Akiba

Single electron production and direct photon production at RHIC

Y. Akiba

Au+Au $\rightarrow$ e : motivation

- From single electron momentum distribution, charm production in Au+Au collision can be extracted.
- This method was used in RUN1 of RHIC.

$$\sigma_{c\bar{c}}^{0-10\%} = 380 \pm 60 \pm 200 \mu\text{b}$$

$$\sigma_{c\bar{c}}^{0-92\%} = 420 \pm 33 \pm 250 \mu\text{b}$$

- From charm production we can probe the earliest stage of Au+Au collision at RHIC
 - Gluon density
 - Energy loss of c-quark in dense matter
 - Reference to J/Ψ production

PHENIX PRL88 192303

Centrality Dependence of Non-Photonic Electrons

Working in Progress

- PYTHIA calculation is scaled with number of binary collisions.
- Well described for all centralities.

Presented in JPS

Direct photon in pp

- High Pt direct photon measurement in pp
 - Test of pQCD
 - Measurement of $G(x)$
 - Measurement of $\Delta G(x)$
- Backgrounds of direct photon are photons from light hadron decays
 - $\pi^0 \rightarrow \gamma\gamma$
 - $\eta \rightarrow \gamma\gamma$
- Expected S/B is small
 - $\gamma/\pi^0 \sim$ a few % at $p_T=4$ GeV/c

Quark-Gluon Compton Scattering

Present status

- T. Horaguchi (Tokyo Inst. Tech, JRA2) is working on the analysis.
- Analysis seems to indicate the presence of direct photon signal that is consistent with pQCD calculation.
- Other groups in PHENIX are also looking for direct photon signal in pp.
- Recently, pre-PPG has been formed.
- RUN3 pp data can be added to the analysis
- Present goal is to present “PHENIX preliminary” data of direct photon in QM2004 conference in January.

Plot shown in the fall JPS meeting
by T. Horaguchi

Vertex Detector for PHENIX

Yuji Goto

Vertex Detector for PHENIX

RBRC Review

November 20, 2003

Yuji Goto

Spin physics with VTX

- Gluon polarization measurement
 - covering x_{gluon} coverage
 - photon+jet and heavy-flavor production

Silicon pixel detector

- ALICE1 pixel hybrid
 - CERN test beam this month ...
 - ALICE production will start soon ???
 - It will include PHENIX chip (CERN-RIKEN contract)

Silicon pixel detector

- Pilot chip & bus
 - at CERN
 - ALICE pixel bus R&D
 - Cu (5 layers): tested
 - Cu+Al (7 layers): will be available, and used for test beam
 - Al (5 layers): will be available soon ...
 - » common power/ground lines for digital and analog parts
 - pilot chip
 - minor-modified version for PHENIX will be made: zero suppression study
 - parallelization study will be done next ...
 - » with bus study
 - In Japan
 - test production of pixel MCM board and bus will be done ...

Silicon strip detector

- Prototype v.1
 - single-sided
 - two-dimensional projection

Silicon strip detector

- Test results
 - β -source
 - KEK test beam
 - prototype v.1
 - position resolution
 - 23-26 μ m (23 μ m expected)
 - efficiency
 - ~97% for good sensors
 - charge division
 - need finer lines

THEORY PRESENTATIONS

Matrix Models and Lattice Data for the QCD Dirac Operator with Chemical Potential

Tilo Wettig

Matrix models and lattice data for the QCD Dirac operator with chemical potential

Tilo Wettig (Yale and RBRC)

work done with Gernot Akemann (Saclay), hep-th/0308003 → PRL

225

Outline:

1. Motivation
2. Matrix models and the domains of weak and strong non-hermiticity
3. Comparison with lattice data

Matrix models at $\mu > 0$

Stephanov (1996): matrix model for Dirac operator

$$Z(\mu) = \int_{\mathbb{C}^{(N+\nu) \times N}} dW e^{-N \operatorname{tr} W W^\dagger} \prod_{f=1}^{N_f} \det \begin{pmatrix} m_f & iW + \mu \\ iW^\dagger + \mu & m_f \end{pmatrix}$$

- explains failure of quenched approximation at $\mu > 0$
- difficult to compute eigenvalue correlations on the scale of the mean level spacing

Akemann (2002): complex eigenvalue model (with $1 - \tau^2 \hat{=} \mu^2$)

$$Z(\tau) = \int_{\mathbb{C}} \prod_{j=1}^N dz_j dz_j^* |z_j|^{2\nu+1} \prod_{f=1}^{N_f} (z_j^2 + m_f^2) e^{-\frac{N}{1-\tau^2} [|z_j|^2 - \frac{\tau}{2} (z_j^2 + z_j^{*2})]} \Delta^2(z^2)$$

- complex extension of the chiral Gaussian Unitary Ensemble
- spectral correlations computed for $N_f = 0$ and for $N_f > 0$
“phase-quenched” massless flavors

- lattice simulations with staggered Dirac operator (Wilson has complex eigenvalues even at $\mu = 0$, Neuberger and DW too expensive)
- need high statistics (20,000 configurations for each parameter set)
- $\beta = 5.0$ strong coupling (can be justified for this particular purpose)
- $\nu = 0$ because of zero-mode shift (Smit-Vink 1985)
- no free parameter: scale is set via mean level spacing

$$d_{\text{weak}} = \frac{\pi}{\sqrt{2}N} \quad d_{\text{strong}} = c \sqrt{\frac{1 - \tau^2}{2N}} \text{ with } c = 1.1429$$

- gauge fields generated in quenched approximation
 $N_f \rightarrow 0$ limit is subtle at $\mu > 0$
 - analytical results from $N_f \rightarrow 0$ limit of “phase-quenched” theory
 $\longrightarrow$ corresponds to theory with quarks and conjugate quarks
 $\longrightarrow$ should agree with quenched lattice data
 - results computed at $N_f > 0$ should describe unquenched lattice data (future work)

Comparison with lattice data: strong nonhermiticity

Dynamical Domain Wall Fermions

Christopher Dawson

Dynamical Domain Wall Fermions

Chris Dawson, RIKEN/BNL Research Center

[RBC Collaboration]

Details

- Dynamical simulations cost ≈ 100 times more than a quenched simulations.
- We are using several improvements to the HMC algorithm
 1. Improved fermion force term
 2. Chronological inverter
 3. Multiple gauge-step leapfrog.that give a factor of ≈ 3 speed-up.
- Total number of “trajectories” collected so far:
 - $am_f = 0.02$ - 4716 trajectories (7 months on 200GF)
 - $am_f = 0.03$ - 4785 trajectories
 - $am_f = 0.04$ - 3445 trajectoriesand we take one “independent” configuration every 50 trajectories.

Summary

- Modifications to the standard algorithms have provided approximately a factor of 2-3 in simulation speed.
- With the DBW2 action, dynamical DWF seem to be practical for lattice spacings around $a^{-1} \approx 2\text{GeV}$. Although the chiral symmetry breaking is larger than for a quenched calculation at a similar scale.

239

Future/Now

- In the progress of gathering more statistics while calculating
 - Spectrum (η')
 - Weak Matrix Elements ($B_K, \epsilon'/\epsilon$)
 - Structure Functions
 -

Pion Electromagnetic Form-factor with Domain Wall Fermions in Lattice QCD

Yukio Nemoto

Pion electromagnetic form-factor with domain wall fermions in Lattice QCD

Yukio Nemoto for RBC collaboration

RIKEN-BNL Research Center

Simulation parameters

- **Fermion action:**

Domain wall fermion ($L_S = 12, M_5 = 1.8$)

$$m_f a = 0.02, 0.04, 0.06, 0.08 \quad (m_\pi = 390 \sim 760 \text{ MeV})$$

- **Gauge action:**

DBW2 (quench) T.Takaishi, PRD54,1050(1996)

$$S_G = -\frac{\beta}{3} ((1 - 8c_1) \sum P[U] + c_1 \sum R[U])$$

P : plaquette, R : 1×2 rectangle, $c_1 = -0.14069$

$$\beta = 0.87, (a^{-1} = 1.31(4) \text{ GeV})$$

- **Lattice:**

Size: $16^3 \times 32$, **Volume** $\sim (2.4 \text{ fm})^3$

All the quark propagators are averaged over periodic and anti-periodic b.c. in time. $\rightarrow$ Effective temporal size is doubled, $32 \rightarrow 64$.

- **Computer:** QCDSP (RIKEN-BNL-Columbia)

$$m_\pi \sim 660 \text{ MeV} \ (m_f = 0.06)$$

$$F_\pi(Q^2 = 0.26 \text{ GeV}^2) \Rightarrow$$

Effective masses for PS mesons

$$F_\pi(Q^2 = 0.53 \text{ GeV}^2) \Rightarrow$$

Summary & Outlook

The pion EM form-factor in DWF + DBW2 gauge action

- $F_\pi(Q^2 = 0) = 1$ is well satisfied.
- small m_f and linear dependence?

Future:

- **More statistics**
- **Higher momenta**
functional form of $F_\pi(Q^2)$, comparison with VMD, charge radius...
Another way to measure the 3-point function is needed.
(on-going)
- **Related quantities**
transition form-factor of the vector meson

Application of DWF to Massive Quarks

Norikazu Yamada

On-shell improvement of Massive DWF

Norikazu Yamada

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RBRC Scientific Review Committee Meeting
November 20, 2003

Heavy quark on the lattice

Obstacles to 5% determination

- (quenched approximation)
- finite volume effect
- chiral extrapolation, . . .
- **heavy quark on the lattice**

Heavy quarks could cause a large discretization error.

255

e.g.) free quark propagator

$$S_q(p) \rightarrow \frac{i\not{p} + m}{p^2 + m^2} + O((am)^2, (a^2 m \vec{p}), (a\vec{p})^2)$$

(a : lattice spacing)

smaller $a \leftrightarrow$ larger cost

Presently $am_c \gtrsim 0.3$, $am_b \gtrsim 1.2$

↓

$O((am)^n)$ and $O((a\vec{p})(am)^n)$ errors $> 5\%$
must be removed!

On-shell improvement

Symanzik, Lüscher, Weisz

Improving the lattice action only on mass shell is enough for masses, WME,

e.g.) free quark propagator

$$S_q(p)|_{\text{on-shell}} \sim \frac{-\gamma_0 E + i\gamma_i p_i + m_1}{-E^2 + \vec{p}^2 + m_1^2} + O((a\vec{p})^2).$$

El-Khadra, Kronfeld, Mackenzie carried out this improvement to the Wilson heavy quark

- no $O((am)^n)$ and $O((a\vec{p})(am)^n)$
- only $a\vec{p} \ll 1$ is required. (possible)

What happen if it is applied DWF?

Results

We found that the improvement is realized by replacing $D_{\text{wilson}}(x, y)$ with

$$D_{\text{imp}}(x, y) = \left(\gamma_0 \Delta_0 + \nu \sum_i \gamma_i \Delta_i \right) - \left(\frac{r_t}{2} \Delta_0^2 + \frac{r_s}{2} \sum_i \Delta_i^2 \right) - \left(c_E \frac{ig}{2} \sum_i \sigma_{0i} F_{0i} + c_B \frac{ig}{4} \sum_{i,j} \sigma_{ij} F_{ij} \right),$$

and get the parameters proper values.

257

$$c_E = \nu \text{ and } c_B = r_s.$$

Vector and Chiral Gauge Theories on the Lattice
Takanori Sugihara

Vector and chiral gauge theories on the lattice

Takanori Sugihara

RIKEN BNL Research Center
Brookhaven National Laboratory

hep-lat/0310061

Vector and chiral gauge theories on the lattice – p.1/11

Chiral anomaly

Continuum theory

- $\text{tr}\gamma_5 \neq 0$ (Fujikawa's method)
- Infinite degrees of freedom

Lattice theory

- $\text{tr}\gamma_5 = 0$
- Finite degrees of freedom

Wilson terms reproduce chiral anomaly.

This work

- Derivative as a Fourier trans of momentum
- Finite lattice
- Antiperiodic boundary conditions
- Lanczos factor (TS, PRD 68, 034502)
- Modification of chiral transformation

Main results

- Chiral anomaly on the finite lattice
- A single Weyl fermion on the lattice

Vector and chiral gauge theories on the lattice – p.5/11

Doubler-free ultralocal derivative

Discrete Fourier transform with $\nabla_{n+N} = -\nabla_n$

$$\nabla_n = \frac{1}{N} \sum_{l=-N/2+1}^{N/2} ip_l e^{i2\pi \tilde{l}n/N}$$

Inverse transform

$$p_l = \sum_{n=1}^{N_c} F_n \nabla_n \left[(\delta_{n, \frac{N}{2}} - 1) 2 \sin \left(\frac{2\pi \tilde{l}n}{N} \right) + \delta_{n, \frac{N}{2}} (-1)^l \right]$$

where F_n is the Lanczos factor.

Vector gauge theory

Modified chiral transformation

$$\begin{aligned}\psi'_m &= \sum_n (1 + i\theta_n \hat{\gamma}_5)_{m,n} \psi_n \\ \bar{\psi}'_m &= \sum_n \bar{\psi}_n (1 + i\theta_n \hat{\gamma}_5)_{n,m}\end{aligned}$$

where

$$(\hat{\gamma}_5)_{m,n} \equiv \gamma_5 \left(1 - \frac{1}{2} a G \right)_{m,n}$$

G satisfies GW relation $\gamma_5 G + G \gamma_5 = a G \gamma_5 G$.

Correct chiral anomaly $\text{tr} \hat{\gamma}_5 \neq 0$ is reproduced.

Vector and chiral gauge theories on the lattice – p.9/11

Chiral gauge theory

Gauge transformation

$$\begin{aligned}\psi'_n &= (1 + i\theta_n^a T_a P_+) \psi_n \\ \bar{\psi}'_n &= \bar{\psi}_n (1 - i\theta_n^a T_a P_-)\end{aligned}$$

Fermion measure is trivially invariant because

$$\text{tr} \gamma_5 = 0$$

A single Weyl fermion can exist on the lattice without violating gauge symmetry.

Kaon Matrix Elements in Domain-Wall QCD with DBW2 Gauge Action

Jun-Ichi Noaki

Kaon Matrix Elements in Domain-Wall QCD with DBW2 Gauge Action

Jun Noaki for RBC Collaboration

RIKEN BNL Research Center

Numerical Simulation

● Quenched Simulation lat/0309026

- **size:** $24^3 \times 48$
- $\beta = 1.22$
- $L_s = 10, M_5 = 1.65$
- **5 light quark masses**
 $m_f a = 0.008 - 0.040, m_s a/2 \simeq 0.0155$
- **7 charm quark masses**
 $m_f a = 0.08 - 0.50, m_c a \simeq 0.45$

271

- #Config.	NPR	23
	B_K	89
	$m_\pi, m_\rho, K \rightarrow \pi$ ME	62

- Basic values

Lattice scale ($m_\rho = 770$ MeV)

$$a^{-1} = 2.88(9) \text{ GeV}$$

Phys. volume $\simeq 1.8$ fm

Chiral symm. breaking

$$m_{\text{res}} \simeq 0.28 \text{ MeV}$$

● Dynamical Simulation

- **size:** $16^3 \times 32$
- $\beta = 0.80$
- $L_s = 12, M_5 = 1.8$
- **3 sea quark masses** $m_{\text{sea}} a = 0.02, 0.03, 0.04$
- **5 valence quark masses**
 $m_{\text{valence}} a = 0.01 - 0.05$

- #Config.	$m_{\text{sea}} a = 0.02$	82
(all MEs	0.03	84
& spectrums)	0.04	58

See C. Dawson's talk for the details

$K \rightarrow \pi\pi$ matrix elements from $K \rightarrow \pi$ ($\Delta I = 1/2$ rule & ε'/ε)

subtraction for $\Delta I = 1/2$: $\langle \pi | Q_i - \alpha_i Q_{\text{subt}} | K \rangle / m_f a \propto \langle \pi\pi | Q_i | K \rangle$ Bernard *et al.*, 1986

Contribution of charmed quark to $\Delta I = 1/2$ rule ?

$$Q_2 = (\bar{s}u)_L(\bar{u}d)_L, \Delta I = 1/2$$

$$Q_{2c} = (\bar{s}c)_L(\bar{c}d)_L, m_c a = 0.40$$

$$A_0 \approx \frac{G_F}{\sqrt{2}} V_{us}^* V_{ud} z_2 \langle (\pi\pi)_{I=0} | (Q_2 - Q_{2c}) | K \rangle$$

Charmed quark tends to make minor contribution.

Calculation of ε'/ε : Large error ($\sim 200\%$) with current statistics.

We need a few times larger statistics to finalize this work.

The Hadronic Vacuum Polarization with 2+1 Flavors of Improved Kogut-Susskind Fermions

Thomas Blum

The Hadronic Vacuum Polarization in 2+1 flavor QCD with Improved Kogut-Susskind Fermions

Tom Blum, RIKEN BNL Research Center

RBRC Review, BNL, November 20, 2003

[Phys. Rev. Lett. 91 (2003) 052001; hep-lat/0310064]

Comparison with 3-loop continuum perturbation theory

Improved Kogut-Susskind: significant scaling violations.

Better as $m_v \rightarrow 0$ (?)

Summary

- Hadronic vacuum polarization is an interesting quantity for lattice calculations
- Important for muon $g-2$ (*c.f.* BNL-E821, Standard Model, ...)
- Investigate the structure of QCD, perturbative *vs.* non-pert, ...
- Important test for dynamical lattice calculations
- Encouraging lattice results for quenched and dynamical $\Pi(q^2)$, $g-2$

Outlook

- Need to understand lattice artifacts in Kogut-Susskind simulations
- Need to understand quark mass behavior (KS and DWF)
- Need better fit function and/or statistics
- Need dynamical DWF results (2+1 flavors: RBRC QCDOC)
- Bigger (and better) lattices for Kogut-Susskind becoming available (MILC)(SciDAC QCDOC(?))

High Density Effective Field Theory

Thomas Schaefer

High Density Effective Theory in QCD

Thomas Schaefer

North Carolina State University

and Riken BNL Research Center

Higher Order Terms

- power corrections: $1/\mu$ terms

$$\mathcal{L} = -\frac{1}{2\mu} \sum_v \psi_v^\dagger \left\{ D_\perp^2 - \frac{ag}{2} \sigma^{\mu\nu} G_{\mu\nu}^\perp - M M^\dagger \right\} \psi_v$$

higher order terms: $\frac{c_n}{\mu^n} \psi_v^\dagger D_\perp^{2n_1} (\bar{v} \cdot D)^{n_2} \psi_v, \dots$

four fermion operators: $\frac{1}{\mu^2} c^{\Gamma\Gamma'}(\vec{v}_i) (\psi_{v_1} \Gamma \psi_{v_2}) (\psi_{v_3}^\dagger \Gamma' \psi_{v_4}^\dagger) \delta_{\Sigma v_i}$

- power counting trivial? amplitude $\sim l^\delta$

$$\delta = 4 + \sum_i N_i (d_i - 4)$$

graphs with higher order vertices suppressed

- problem: “hard” loops

Hard Loops: Gluon Self Energy

- gluon self energy: sum over patches

$$\sum_{\vec{v}} \int \frac{d^4 k}{(2\pi)^4} \frac{v_\mu v_\nu}{(k_0 - l_k)(k_0 + p_0 - l_{k+p})}$$

- use

$$\frac{1}{2\pi} \sum_{\vec{v}} \int \frac{d^2 l_\perp}{(2\pi)^2} = \frac{\mu^2}{2\pi^2} \int \frac{d\Omega}{4\pi}.$$

- HDL polarization function

$$\Pi_{\mu\nu}^{ab}(p) = 2m^2 \delta^{ab} \int \frac{d\Omega}{4\pi} v_\mu v_\nu \left\{ 1 - \frac{p_0}{p_0 - l_p} \right\}$$

- need counter term $m^2 A_\perp^2$
- hard loops modify power counting
- have to resum for $E \sim g\mu$

leading order
hard loop diagrams

...

Superconductivity

- $E \sim \mu \exp(-1/g)$: BCS ladder diagrams IR divergent

introduce gap term

$$\mathcal{L} = \Delta R_l^\Gamma(\vec{v} \cdot \hat{\Delta})\psi_{-v}\sigma_2\Gamma\psi_v + h.c.$$

- gap equation: only two patches ($v, -v$)

$$\Delta(p_4) = \frac{g^2}{18\pi^2} \int_0^{\Lambda_{\parallel}} dq_4 \frac{\Delta(q_4)}{\sqrt{q_4^2 + \Delta(q_4)^2}} \log \left(\frac{\Lambda_{\perp}}{|p_4^2 - q_4^2|^{1/2}} \right).$$

- solution (can set $\Lambda_{\parallel} = \Lambda_{\perp}$ at this order)

$$\Delta_0 = 2\Lambda \exp(-3\pi^2/(\sqrt{2}g))$$

Summary: High Density Effective Theories

- Effective theories provide useful tools for systematic calculations in the regime $\Lambda < \Delta < \mu$.
- naive power counting modified by “hard loops”
- important questions: renormalization of four-fermion operators, counting and resumming logs

Hadronization at RHIC: Interplay of Recombination and Fragmentation

Steffen Bass

Hadronization @ RHIC: Interplay of Fragmentation and Recombination

Steffen A. Bass

Duke University
& RIKEN-BNL Research Center

- The proton puzzle at RHIC
- Recombination + Fragmentation Model
- Results and Predictions

- R.J. Fries, C. Nonaka, B. Mueller & S.A. Bass, PRL 90 202303 (2003)
- R.J. Fries, C. Nonaka, B. Mueller & S.A. Bass, PRC 68 044902 (2003)

Species dependent saturation of elliptic flow

299

Steffen A. Bass

- hyperon v_2 saturates later and higher than kaon v_2 .
 - same effect observed for protons and pions.
 - the phenomenology seems better described in $m_T - m_0$ than p_T ; why (kinetic energy)?
 - what drives the different p_T scales for K_S and Λ v_2 ?
- novel mechanism of baryon formation?

Hadronization @ RHIC #3

Recombination: new life for an old idea

High Energy Physics Phenomenology:

- K.P. Das & R.C. Hwa, Phys. Lett. B68, 459 (1977)

Quark-Antiquark Recombination in the Fragmentation Region

- description of leading particle effect

- T. Ochiai, Prog. Theo. Phys. 75, 1184 (1986)
- E. Braaten, Y. Jia & T. Mehen, Phys. Rev. Lett. 89, 122002 (2002)
- R. Rapp & E.V. Shuryak, Phys. Rev. D67, 074036 (2003)

Heavy-Ion Phenomenology:

- T. S. Biro, P. Levai & J. Zimanyi, Phys. Lett. B347, 6 (1995)

ALCOR: a dynamical model for hadronization

- yields and ratios via counting of constituent quarks

- R.C. Hwa & C.B. Yang, Phys. Rev. C66, 025205 (2002)
- R. Fries, B. Mueller, C. Nonaka & S.A. Bass, Phys. Rev. Lett. 90
- V. Greco, C.M. Ko and P. Levai, Phys. Rev. Lett. 90

Anisotropic flow:

- S. Voloshin, QM2002, nucl-ex/020014
- Z.W. Lin & C.M. Ko, Phys. Rev. Lett 89, 202302 (2002)
- D. Molnar & S. Voloshin, nucl-th/0302014

Input and Model Parameters

Input for the model is the momentum distributions of constituent quarks and anti-quarks at the time of hadronization

- the quark distribution is assumed to have a low p_t thermal component and a high p_t pQCD mini-jet component

- the thermal component is parameterized as:

$$w_a(\sigma, p) = g_a e^{-p \cdot v(\sigma)/T} e^{-\eta^2/2\Delta^2} f(\rho, \phi)$$

with a flavor dependent fugacity g_a , temperature T , rapidity width Δ and transverse distribution $f(\rho, \phi)$

- the pQCD component is parameterized as:

$$\left. \frac{dN_a}{p_t dp_t dy} \right|_{y=0} = K \frac{C}{(1 + p_t/B)^\beta}$$

with parameters C , B and β taken from a lo pQCD calculation

- parton energy loss is treated in the pQCD component
(currently utilizing the BDMPS approach)

Hadron Ratios vs. p_t

Elliptic Flow

- anisotropic or “elliptic” flow is sensitive to initial geometry

low p_t domain:

$$Force = -\nabla P$$

more flow in collision plane
than perpendicular to it

high p_t domain:

$$\Delta E \sim L$$

less absorption in collision
plane than perpendicular to it

- total elliptic flow is the sum of both contributions:

$$v_2(p_t) = r(p_t) v_2^{\text{recomb}}(p_t) + (1 - r(p_t)) v_2^{\text{frag}}(p_t)$$

$r(p_t)$: relative weight of the fragmentation contribution in spectra

Parton Number Scaling of v_2

• in leading order of v_2 ,
recombination predicts:

$$v_2^M(p_t) = 2v_2^p\left(\frac{p_t}{2}\right)$$

309

$$v_2^B(p_t) = 3v_2^p\left(\frac{p_t}{3}\right)$$

- smoking gun for recombination
- measurement of partonic v_2 !

Summary & Outlook

The Recombination + Fragmentation Model:

- provides a natural solution to the baryon puzzle at RHIC
- describes the intermediate and high p_t range of
 - hadron ratios & spectra
 - jet-quenching phenomena
 - elliptic flow
- provides a microscopic basis for the Statistical Model

issues to be addressed in the future:

- entropy production
- treatment of gluons
- realistic space-time dynamics of parton source
- need improved data of identified hadrons at high p_t

Interplay Between Soft and Hard Hadronic Components at RHIC

Tetsufumi Hirano

Interplay Between Soft and Hard Hadronic Components at RHIC

Tetsufumi Hirano

RIKEN BNL Research Center

Collaboration with Yasushi Nara (Arizona)

OUTLINE

- Introduction
- Hydro+jet model
- p_T spectra for π , K, p
- Jet quenching at $\eta \sim 2$
- Summary

References

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nucl-th/0307015;
nucl-th/0307087.

RBRC Scientific Review Committee Meeting, Nov. 20-21, 2003

p_T Spectra for Identified Hadrons

$p_{T,\text{cross}}$ depends on particle species!

$p_{T,\text{cross}} \sim$ 1.8 GeV/c for π
 2.7 GeV/c for K
 3.7 GeV/c for p

c.f.) $p_{T,\text{cross}} \sim$ inflection point for kaons and protons

Summary

QGP in Au+Au collisions at RHIC ?

→ Final state effects are consistent with intermediate ~ high p_T data (R_{AA} , b-to-b, ...).

→ Strongly encouraged by recent d+Au data.

- Dynamical model (hydro+jet) for heavy-ion physics

- Hadron species dependent transverse dynamics

→ Interplay between radial flow and jet quenching

- How large dense matter in longitudinal direction?

← Jet quenching in forward rapidity region ($\eta \sim 2$)

Calculations of Spin Asymmetries in pp Scattering

Werner Vogelsang

Calculations of spin asymmetries in pp scattering

Werner Vogelsang

(RBRC Scientific Review, Nov. 20/21 2003)

Today's example :

- High- p_T pion production

with S. Kretzer, Jäger, Stratmann

Other work :

- Soft-gluon resummation for heavy flavor production

with H. Yokoya

- Transverse-spin asymmetries at RHIC

with Mukherjee, Stratmann

- Higgs production at LHC

with Kulesza, Sterman

*Spin theory group is growing ! From Nov.'03 : Y. Koike, T. Kucs

Polarized pp collisions at RHIC :

- spin asymmetry

$$A_{LL} \equiv \frac{\sigma_{++} - \sigma_{+-}}{\sigma_{++} + \sigma_{+-}} \equiv \frac{\Delta\sigma}{\sigma}$$

- consider high- p_T reactions :

$$\vec{p}\vec{p} \rightarrow \text{jet}X, \vec{p}\vec{p} \rightarrow \gamma X, \vec{p}\vec{p} \rightarrow \pi X, \dots$$

$$p_T^3 \frac{d\Delta\sigma}{dp_T d\eta} = \left| \begin{array}{c} \text{Diagram: Two incoming protons (P) interact via a central vertex (hat{\sigma}). One vertex (a) emits a parton (a) which then interacts with a target (D_c) to produce a hadron (h). The other vertex (b) emits a parton (b) which interacts with the central vertex to produce X'. The central vertex also emits a parton (c) which interacts with the target (D_c) to produce a hadron (h).} \end{array} \right|^2 + \mathcal{O}\left(\frac{\lambda}{p_T}\right)^n$$

$$d\Delta\sigma^{pp \rightarrow \pi X} = \sum_{abc} \Delta f_a \otimes \Delta f_b \otimes d\Delta\hat{\sigma}^{ab \rightarrow cX'} \otimes D_c^\pi + \text{P.C.}$$

$$\Delta\hat{\sigma}^{(0)} + \alpha_s \Delta\hat{\sigma}^{(1)} + \dots \text{pert.}$$

Jäger, Schäfer, Stratmann, WV

- nucleon spin structure $\Delta f_{a,b}$
- fragmentation functions $D_c^\pi \rightsquigarrow e^+e^-$ annihilation

A lower bound on A_{LL}^π

Jäger, Kretzer, Stratmann, WV

- integrate LO cross section over all rapidities $|\eta| \leq \cosh^{-1}(x_T)$,
and take mellin moments in $x_T^2 \equiv 4p_\perp^2/S$:

$$\Delta\sigma^\pi(N) \equiv \int_0^1 dx_T^2 (x_T^2)^{N-1} \frac{p_\perp^3 d\Delta\sigma^\pi}{dp_\perp}$$

- obtain

$$\Delta\sigma^\pi(N) = \sum_{a,b,c} \Delta f_a^{N+1} \Delta f_b^{N+1} \Delta \hat{\sigma}_{ab}^{c,N} D_c^{\pi,2N+3}$$

- in terms of moments of Δg :

$$\Delta\sigma^\pi(N) = (\Delta g^{N+1})^2 \mathcal{A}^N + 2 \Delta g^{N+1} \mathcal{B}^N + \mathcal{C}^N$$

$\uparrow$
 $gg \text{ scatt.}$

$\uparrow$
 qg

$\uparrow$
 $qq, qq', q\bar{q}, \dots$

- a parabola – with a minimum given by

$$\mathcal{A}^N \Delta g^{N+1} = -\mathcal{B}^N$$

$$\Delta\sigma^\pi(N) \Big|_{\min} = -\frac{(\mathcal{B}^N)^2}{\mathcal{A}^N} + \mathcal{C}^N$$

- Mellin inverse :

$$\left. \frac{p_\perp^3 d\Delta\sigma^\pi}{dp_\perp} \right|_{\min} = \frac{1}{2\pi i} \int_\Gamma dN (x_T^2)^{-N} \Delta\sigma^\pi(N) \Big|_{\min}$$

- resulting asymmetry : negative, but tiny

$$p_\perp = 1.5 \text{ GeV} : \boxed{A_{LL}^\pi \approx -10^{-4}}$$

$$p_\perp = 4.5 \text{ GeV} : \boxed{A_{LL}^\pi \approx -10^{-3}}$$

Dark Matter and Pulsar Kicks from a Sterile Neutrino

Alexander Kusenko

Pulsar velocities

2

Pulsar kicks and dark matter from sterile neutrinos

- Pulsar kicks and possible explanations
- Sterile neutrinos: the missing dark matter particle?
- Pulsar kicks from active to sterile neutrino conversions

1

Pulsar kicks from neutrino emission?

Pulsar with $v \sim 500$ km/s has momentum

$$M_{\odot} v \sim 10^{11} \text{ g cm/s}$$

SN energy released: 10^{53} erg $\Rightarrow$ in neutrinos. Thus, the total neutrino momentum is

$$P_{\nu, \text{total}} \sim 10^{43} \text{ g cm/s}$$

a **1% asymmetry** in the distribution of **neutrinos**

is sufficient to explain the pulsar kick velocities

But what can cause the asymmetry??

6

Supernova neutrinos

Nuclear reactions in stars lead to a formation of a heavy iron core. When it reaches $M \approx 1.4M_{\odot}$, the pressure can no longer support gravity. $\Rightarrow$ collapse.

Energy released:

$$\Delta E \sim \frac{G_N M_{\text{Fe core}}^2}{R} \sim 10^{53} \text{ erg}$$

99% of this energy is emitted in neutrinos

5

Core collapse supernova

Thermal cooling: $t = 10 - 15$ s

Most of the neutrinos emitted during the cooling stage.

10

Core collapse supernova

Shock formation and "neutronization burst": $t = 1 - 10$ ms

Protoneutron star formed. Neutrinos are trapped. The shock wave breaks up nuclei, and the initial neutrinos come out (a few %).

9

Can the weak interactions asymmetry cause an anisotropy in the flux of neutrinos due to a large magnetic field?

No

Rescattering washes out the asymmetry [Vilenkin; AK, Segrè, Vilenkin].

In approximate thermal equilibrium the asymmetries in scattering amplitudes do not lead to an anisotropic emission. Only the outer regions, near *neutrinospheres*, contribute (a negligible amount).

However, if a weaker-interacting sterile neutrino was produced in these processes, the asymmetry would, indeed, result in a *pulsar kick*!

14

Can the weak interactions asymmetry cause an anisotropy in the flux of neutrinos due to a large magnetic field?

No

Neutrinos are trapped at high density.

13

A sterile neutrino in this range, consistent with dark matter, can also explain the observed velocities of pulsars through $\nu_e \rightarrow \nu_s$ oscillations in a supernova.

18

Sterile neutrinos in cosmology: dark matter

Sterile neutrinos are produced in primordial plasma through oscillations.

The resulting density of relic sterile neutrinos:

$$\Omega_{\nu_s} \sim 0.3 \left(\frac{\sin^2 2\theta}{10^{-8}} \right) \left(\frac{m_s}{\text{keV}} \right)^2$$

[Dodelson, Widrow; Dolgov, Hansen; Fuller, Shi; Abazajian, Fuller, Patel]

17

$$\begin{cases} |\nu_1\rangle = \cos \theta_m |\nu_e\rangle - \sin \theta_m |\nu_s\rangle \\ |\nu_2\rangle = \sin \theta_m |\nu_e\rangle + \cos \theta_m |\nu_s\rangle \end{cases} \quad (6)$$

The coupling of ν_2 to weak currents is also suppressed, and $\sigma \propto \sin^2 \theta$.

The probability of $\nu_e \rightarrow \nu_s$ conversion in presence of matter is

$$\langle P_m \rangle = \frac{1}{2} \left[1 + \left(\frac{\lambda_{\text{osc}}}{2\lambda_s} \right)^2 \right]^{-1} \sin^2 2\theta_m. \quad (7)$$

where λ_{osc} is the oscillation length, and λ_s is the scattering length.

22

Mixing angle in matter is different from vacuum:

$$\sin^2 2\theta_m = \frac{(\Delta m^2/2p)^2 \sin^2 2\theta}{(\Delta m^2/2p)^2 \sin^2 2\theta + (\Delta m^2/2p \cos 2\theta - V_m)^2}, \quad (3)$$

$$V_m = \frac{G_F \rho}{\sqrt{2} m_n} (3Y_e - 1 + 4Y_{\nu_e} + 2Y_{\nu_\mu} + 2Y_{\nu_\tau}) \quad (4)$$

$$\simeq (-0.2 \dots + 0.5) V_0, \quad (5)$$

where $V_0 = G_F \rho / \sqrt{2} m_n \simeq 3.8 \text{eV} (\rho / 10^{14} \text{gcm}^{-3})$

Mixing is suppressed when $V_m \gg (\Delta m^2/2k)$.

21

If the fraction of energy emitted in sterile neutrinos is

$$r_{\mathcal{E}} = \left(\frac{\mathcal{E}_s}{\mathcal{E}_{\text{tot}}} \right) \sim 0.05 - 0.7, \quad (9)$$

(as it can easily be), then the resulting momentum asymmetry is

$$\epsilon \sim 0.02 \left(\frac{k_0}{0.3} \right) \left(\frac{r_{\mathcal{E}}}{0.5} \right), \quad (10)$$

which is sufficient to explain the pulsar kick velocities.

26

Sterile neutrinos leave the star without scattering. Hence, they give the pulsar a kick.

25

Resonant active-sterile neutrino conversions in matter

Matter potential:

$$V(\nu_s) = 0$$

$$V(\nu_e) = -V(\bar{\nu}_e) = \frac{G_F \rho}{\sqrt{2} m_n} (3Y_e - 1 + 4Y_{\nu_e})$$

$$V(\nu_{\mu,\tau}) = -V(\bar{\nu}_{\mu,\tau}) = \frac{G_F \rho}{\sqrt{2} m_n} (Y_e - 1 + 2Y_{\nu_e}) + c_L^z \frac{\vec{k} \cdot \vec{B}}{k}$$

$$c_L^z = \frac{e G_F}{\sqrt{2}} \left(\frac{3N_c}{\pi^4} \right)^{1/3}$$

30

Allowed range of parameters (time scales, fraction of total energy emitted):

[Fuller, AK, Mocioiu, Pascoli, astro-ph/0307267]

29

A crude estimate of the kick

The mean energy of emitted sterile neutrinos is proportional to the temperature at the point of production. The point of resonant conversion depends on direction:

$$r(\phi) = r_0 + \delta \cos \phi, \quad (12)$$

where $\cos \phi = (\vec{k} \cdot \vec{B})/k$ and δ is determined by the equation:

$$2 \frac{dN_n(r)}{dr} \delta \approx e \left(\frac{3N_e}{\pi^4} \right)^{1/3} B. \quad (13)$$

This yields

34

The magnetic field shifts the position of the resonance because of the $\frac{\vec{k} \cdot \vec{B}}{k}$ term in the potential:

33

The range of parameters [AK, Segrè; Fuller, AK, Mocioiu, Pascoli]:

38

The core density $\rho \sim 10^{14} \text{ g/cm}^3$ determines the

$$\Delta m^2 \sim (10 \text{ keV})^2$$

Adiabaticity: the oscillation length

$$\lambda_{\text{osc}} \approx \left(\frac{1}{2\pi} \frac{\Delta m^2}{2k} \sin 2\theta \right)^{-1} \sim \frac{1 \text{ mm}}{\sin 2\theta}.$$

must be smaller than (1) the scale height of density (2) the mean free path of neutrinos. $\Rightarrow$

$$\sin^2 \theta \gtrsim 10^{-10}$$

37

Chandra non-observation of photons from $\nu_s \rightarrow \nu_e \gamma$

non-zero lepton asymmetry changes the dark matter range

42

Chandra non-observation of photons from $\nu_s \rightarrow \nu_e \gamma$

[Abazajian, Fuller, Tucker]

41

Resonant (1) & off-resonant (2) emissions combined:

[Fuller, **AK**, Mocoiu, Pascoli, Phys. Rev. D, in press; astro-ph/0307267]

46

Conclusions

- Sterile neutrinos in the 1-20 keV range can explain the observed pulsar kicks
- The same neutrino could be the dark matter
- One may be able to confirm this mechanism through observations of gravity waves in the event of a nearby supernova.

45

Charge Symmetry Breaking in Few-Nucleon Systems

Ubirajara van Kolck

Charge Symmetry Breaking in Few-Nucleon Systems

Bira van Kolck

University of Arizona

and

RIKEN-BNL Research Center

- ◇ What and why
- ◇ Chiral Symmetry
- ◇ $np \rightarrow d\pi^0$
- ◇ $dd \rightarrow \alpha\pi^0$
- ◇ $NN \rightarrow NN$
- ◇ Summary

Why

- A eventually to be calculated in lattice QCD

▷ extract $m_d - m_u$

- B estimated in models, *e.g.*

$$\bar{\delta}m_N = -(0.76 \pm 0.30) \text{ MeV} \quad \text{Cottingham sum rule} \quad \text{Gasser + Leutwyler '82}$$

$$\bar{\delta}m_N = -\frac{\alpha}{R_N} \simeq -1.5 \text{ MeV} \quad \text{quark models} \quad \text{Miller, Nefkens + Šlaus '90}$$

▷ test models

- BBN depends sensitively on Δm_N , *e.g.*

$$\text{if } \Delta m_N \lesssim \begin{cases} m_e : & \text{hydrogen} \\ -m_e : & \text{proton} \end{cases} \text{ unstable under } \begin{cases} \text{electron capture} \\ \beta \text{ decay} \end{cases}$$

→ BBN produces essentially ${}^4\text{He}$ + free neutrons

(→ no water ...)

▷ constrain variations of parameters with time

Kolb, Perry + Walker '86

...

$$np \rightarrow d\pi^0$$

Miller, Niskanen + v.K., '00

$$\begin{aligned}
 A_{fb} &= \frac{\int_0^{\pi/2} d\Omega (\sigma(\theta) - \sigma(\pi - \theta))}{\int_0^{\pi/2} d\Omega (\sigma(\theta) + \sigma(\pi - \theta))} \\
 &= 24.4 \cdot 10^{-4} \left\{ \frac{\delta m_N - \bar{\delta} m_N / 2}{\text{MeV}} - 1.15 \frac{g_{\eta NN}}{\sqrt{4\pi \cdot 3.68}} \frac{\langle \pi^0 | H_{QCD} | \eta \rangle}{(-0.0059 \text{ GeV}^2)} + \dots \right\}
 \end{aligned}$$

at $q_\pi/m_\pi = 0.17$.

For $g_{\eta NN}^2/4\pi = 3.68$ and $\langle \pi^0 | H_{QCD} | \eta \rangle = -0.0059 \text{ GeV}^2$,

$$\leadsto \begin{cases} \delta m_N = 1.66 \pm 0.27 \pm ?_{th} \text{ MeV} \\ \bar{\delta}m_N = -0.36 \pm 0.27 \pm ?_{th} \text{ MeV} \end{cases}$$

Observation of the Charge Symmetry Breaking $d + d \rightarrow {}^4\text{He} + \pi^0$ Reaction Near ThresholdE. J. Stephenson,¹ A. D. Bacher,^{1,2} C. E. Allgower,¹ A. Gärdestig,³ C. M. Lavelle,¹ G. A. Miller,⁴ H. Nann,^{1,2} J. Olmsted,¹ P. V. Pancella,⁵ M. A. Pickar,⁶ J. Rapaport,⁷ T. Rinckel,¹ A. Smith,⁸ H. M. Spinka,⁹ and U. van Kolck^{10,11}¹Indiana University Cyclotron Facility, Bloomington, Indiana 47408, USA²Department of Physics, Indiana University, Bloomington, Indiana 47405, USA³Indiana University Nuclear Theory Center, Bloomington, Indiana 47408, USA⁴Department of Physics, University of Washington, Seattle, Washington 98195, USA⁵Physics Department, Western Michigan University, Kalamazoo, Michigan 49008, USA⁶Department of Physics and Astronomy, Minnesota State University at Mankato, Mankato, Minnesota 56001, USA⁷Department of Physics and Astronomy, Ohio University, Athens, Ohio 45701, USA⁸Physics Department, Hillsdale College, Hillsdale, Michigan 49242, USA⁹Argonne National Laboratory, Argonne, Illinois 60439, USA¹⁰Department of Physics, University of Arizona, Tucson, Arizona 85721, USA¹¹RIKEN BNL Research Center, Brookhaven National Laboratory, Upton, New York 11973, USA

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We report the first observation of the charge symmetry breaking $d + d \rightarrow {}^4\text{He} + \pi^0$ reaction near threshold. Measurements using a magnetic channel (gated by two photons) of the ${}^4\text{He}$ scattering angle and momentum (from time of flight) permitted reconstruction of the π^0 "missing mass," the quantity used to separate ${}^4\text{He} + \pi^0$ events from the continuum of double radiative capture ${}^4\text{He} + \gamma + \gamma$ events. We measured total cross sections for neutral pion production of 12.7 ± 2.2 pb at 228.5 MeV and 15.1 ± 3.1 pb at 231.8 MeV. The uncertainty is dominated by statistical errors. These cross sections arise fundamentally from the down-up quark mass difference and quark electromagnetic effects that contribute in part through meson mixing (e.g., $\pi^0 - \eta$) mechanisms.

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PACS numbers: 24.80.+y, 24.85.+p, 25.10.+s, 25.45.-z

Charge symmetry is the symmetry of the quantum chromodynamic Lagrangian under the interchange of equal-mass up and down quarks [1,2]. This symmetry is broken by the different masses of the down and up quarks ($m_d > m_u$) and by their electromagnetic interactions. The combination of these two mechanisms leads, for example, to the neutron being heavier than the proton. Within the framework of chiral effective field theory [3,4], additional experimental information on the relative contributions of these two mechanisms to charge symmetry breaking (CSB) must, in leading order, come from pion-nucleon scattering. Direct experimental evidence is restricted to elastic scattering and charge exchange experiments with low-energy charged pions where the interpretation is complicated by corrections for the neutron-proton mass difference and electromagnetic interactions between the pions and nucleons [5,6]. Reactions in which a π^0 is emitted after being created by one nucleon and rescattered by a second are particularly clean. One example of such a CSB process is the measurement of a forward-backward asymmetry in the cross section for the $n + p \rightarrow d + \pi^0$ reaction [7].

The $d + d \rightarrow {}^4\text{He} + \pi^0$ reaction violates the conservation of isospin because the pion has isospin one while both the deuteron and ${}^4\text{He}$ have isospin zero. More specifically, it violates charge symmetry, which is a rotation by $\pi/2$ about the y axis in charge space that interchanges up and down quarks. (Isospin conservation is invariance under any rotation.) The π^0 , whose wave function is odd under

the interchange of down and up quarks, should not be produced from an initial state with charge symmetry even (in this case self-conjugate). The amplitude for CSB is weaker than a similar charge symmetry conserving amplitude by about 1/300, roughly the ratio of the quark mass difference to the nucleon mass. This suggests a $d + d \rightarrow {}^4\text{He} + \pi^0$ total cross section, which depends only on a CSB amplitude, that is as small as tens of picobarns.

Several searches for the $d + d \rightarrow {}^4\text{He} + \pi^0$ reaction have produced only upper limits [8]. A positive report at a deuteron energy of 1.1 GeV [9] has been questioned because the experiment did not clearly distinguish the photons from π^0 decay from photons that could have been produced by the double radiative capture $d + d \rightarrow {}^4\text{He} + \gamma + \gamma$ process. This isospin-allowed process was calculated to be of the same order as the reported cross section [10]. A further search that could clearly distinguish between these two reactions was therefore warranted.

We chose to look for the $d + d \rightarrow {}^4\text{He} + \pi^0$ reaction just above its threshold at 225.5 MeV to avoid other pion producing channels and to take advantage of the clean experimental conditions afforded by the Indiana University Cyclotron Facility's electron-cooled storage ring. A 6° bend located in one section of the ring provided a site where ${}^4\text{He}$ nuclei, produced in a narrow forward cone just above threshold, could be separated from the circulating deuteron beam. By placing a gas jet target sufficiently upstream of the 6° magnet, it became possible to

SUMMARY – a glimpse of what it might look like

Estimates

Niskanen '02

Friar, Payne, Coon + v.K. '03

Mechanism	$a_{pp} - a_{nn}$ (fm)	$B_3\text{H} - B_3\text{He}$ (keV)
short range	0.7	45
πNN coupling; OPE	0 ± 0.8	0 ± 50
kinetic energy	0.2	14
$\delta m_N, \bar{\delta} m_N$; TPE	< 0.3	< 20
...	...	...
"expt" (soft-EM subtracted)	1.5 ± 0.5	87

QCD and Dimensional Deconstruction

Mikhail Stephanov

QCD and Dimensional Deconstruction

M. Stephanov

U. of Illinois at Chicago and RIKEN-BNL

Reference: D.T. Son, M.S., hep-ph/0304182

QCD and Dimensional Deconstruction – p.1/11

Introduction/Motivation

- QCD is largely unsolved theory.
- Quark/gluon interactions are known and fixed by a principle of local gauge invariance. But ...
- Observable degrees of freedom are hadrons. What principles govern their interactions? Some can be derived from QCD, some not (e.g., VMD).
- Idea of dual description. Large N_c ('t Hooft, Witten) – dual theory weakly coupled?

QCD and Dimensional Deconstruction – p.2/11

a_1 , etc

- 2 hidden local symmetries: π, ρ, a_1 (Bando, *et al*)
- The diagram:

- More mesons (PDG): $\rho(770), a_1(1260), \rho(1450), a_1(1640), \rho(1700), \dots$
Large N_c – infinite tower of vector/axial mesons.
- Motivation to consider theories of arbitrary number K of hidden local symmetries.

$$\text{---} \quad K=0: \pi$$

$$\text{---} \bigcirc \text{---} \quad K=1: \pi, \rho$$

$$\text{---} \bigcirc \text{---} \bigcirc \text{---} \quad K=2: \pi, \rho, a_1$$

$$\text{SU}(2)_L \text{---} \bigcirc \text{---} \bigcirc \text{---} \dots \text{---} \bigcirc \text{---} \bigcirc \text{---} \text{SU}(2)_R$$

$(gA_\mu)^1 \quad (gA_\mu)^2 \quad \dots \quad (gA_\mu)^{K-1} \quad (gA_\mu)^K$
 $\Sigma^1 \quad \Sigma^2 \quad \Sigma^3 \quad \Sigma^{K-1} \quad \Sigma^K \quad \Sigma^{K+1}$

QCD and Dimensional Deconstruction – p.5/11

Open “moose”

$$\text{SU}(2)_L \text{---} \bigcirc \text{---} \bigcirc \text{---} \dots \text{---} \bigcirc \text{---} \bigcirc \text{---} \text{SU}(2)_R$$

$(gA_\mu)^1 \quad (gA_\mu)^2 \quad \dots \quad (gA_\mu)^{K-1} \quad (gA_\mu)^K$
 $\Sigma^1 \quad \Sigma^2 \quad \Sigma^3 \quad \Sigma^{K-1} \quad \Sigma^K \quad \Sigma^{K+1}$

$$\mathcal{L} = \sum_{k=1}^{K+1} f_k^2 \text{Tr} |D_\mu \Sigma^k|^2 - \sum_{k=1}^K \frac{1}{2} \text{Tr} (F_{\mu\nu}^k)^2.$$

Similar to the closed “moose” theory (Arkani-Hamed, Cohen, Georgi).
 f_k and g_k – parameters.

Physical observables:

$$\frac{4}{f_\pi^2} = \sum_{k=1}^{K+1} \frac{1}{f_k^2}.$$

Meson masses: solutions of normal mode eqs (“springs” f_k , “weights” $1/g_k$)

$$(M^2)_{kk'} b_n^{k'} = m_n^2 b_n^k.$$

Even/odd b_n alternate $\rightarrow \rho$ and a_1 alternate.

QCD and Dimensional Deconstruction – p.6/11

4 + 1d and dimensional deconstruction

$$\Sigma^k \approx 1 + iaA_5(u), \quad (5\text{th link variable})$$

$$S = -\text{Tr} \int du d^4x \left(-f^2(u) F_{5\mu}^2 + \frac{1}{2g^2(u)} F_{\mu\nu}^2 \right).$$

or

$$S = -\frac{1}{2g_0^2} \text{Tr} \int d^5x \sqrt{|g|} e^{-2\phi} F_{\hat{\mu}\hat{\nu}} F^{\hat{\mu}\hat{\nu}},$$

Background metric:

$$ds^2 = -du^2 + f^2 g^2 \eta_{\mu\nu} dx^\mu dx^\nu.$$

($f^2 g^2$ – warp factor), and dilaton field (varying gauge coupling)

$$g_0^2 e^{2\phi} = g^2(u).$$

GCD and Dimensional Deconstruction – p.9/11

AdS/CFT connection

In the continuum limit the procedure of calculating current-current correlators is the same as in AdS/CFT correspondence.

Introduce auxiliary vector fields A_μ^L and A_μ^R . In the continuum limit $\equiv$ to setting b.c. on the 5d (bulk) field A_μ on 2 boundaries:

$$A_\mu(-u_0) = A_\mu^L \quad \text{and} \quad A_\mu(+u_0) = A_\mu^R.$$

The classical action S_{cl} is a functional of the b.c. Calculate the 2nd derivative of the action with respect to the boundary fields, e.g.,

$$\langle \mathcal{L}_\mu(x) \mathcal{L}_\nu(y) \rangle = \frac{\delta^2 S_{cl}[A_\mu^i]}{\delta A_\mu^L(x) \delta A_\nu^L(y)}.$$

Unlike AdS/CFT there are 2 boundaries. Moose can be folded, if needed.

GCD and Dimensional Deconstruction – p.10/11

Current correlator and matching to QCD

$$g_{nV,A} = \sqrt{2n(n+1)(2n+1)} \frac{\Lambda^2}{g_5}.$$

$$\Pi_V(Q^2) = \sum_{n=1}^K \frac{g_{nV}^2}{m_n^2(Q^2 + m_n^2)} = \frac{2\Lambda^2}{g_5^2} \sum_{n \text{ odd}} \frac{2n+1}{Q^2 + n(n+1)\Lambda^2}.$$

For large Q^2 :

$$\Pi_V(Q^2) \rightarrow -\frac{1}{g_5^2} \ln(Q^2), \quad (Q^2 \gg \Lambda^2).$$

Matching with QCD: $\Pi_V(Q^2) \rightarrow -\frac{N_c}{24\pi^2} \ln(Q^2)$, one finds:

$$\frac{1}{g_5^2} = \frac{N_c}{24\pi^2}.$$

which gives, for example,

$$f_\pi^2 = \frac{N_c}{24\pi^2} m_\rho^2. \quad (\text{SVZ})$$

QCD and Dimensional Deconstruction – p.13/11

VMD in ‘Cosh’ background

Pion e.m. formfactor is saturated by a single ρ pole:

$$g_{n\pi\pi} = \frac{f_\pi^2}{4} \int_{-u_0}^{+u_0} \frac{du}{f^2(u)} g(u) b_n(u) = -\frac{c_n}{2} g_5 \int_{-1}^1 d\xi \sqrt{1-\xi^2} P_n^1(\xi),$$

vanishes unless $n = 1$.

More generally, contribution of higher resonances is suppressed due to smaller overlap of oscillating $b_n(u)$ with pion “wave-function” $1/f^2(u)$.

QCD and Dimensional Deconstruction – p.14/11

**Kinetic Equation for ϕ^4
Theory with Generalized Kadanoff-Baym Ansatz**

Takashi Ikeda

Kinetic equation for Φ^4 Theory with generalized Kadanoff-Baym ansatz

Takashi Ikeda

RIKEN BNL Research Center

RIKEN Review, Nov., 2003

Scalar Φ^4 theory & Gap equation

- Self-interacting scalar Φ^4 theory : $\mathcal{L} = \frac{1}{2} (\partial_\mu \Phi)^2 - \frac{1}{2} m^2 \Phi^2 - \frac{\lambda}{4!} \Phi^4$
(3+1 dimensions, spatially homogeneous)

- Focusing on decay of single particle excitation

(Taking selfenergies as in thermal equilibrium and

using quasiparticle approximation: $\rho(\mathbf{p}, p_0) = 2\pi \varepsilon(p_0) \delta(p_0^2 - \mathbf{p}^2 - M_T^2)$)

● Self-energies

- generating thermal mass M_T
by solving gap eq. selfconsistently

$$M_T^2 = \frac{\lambda}{2} \int \frac{d^3 \mathbf{k}}{(2\pi)^3} \frac{N_{\text{eq}}(E_k)}{E_k} + \frac{\lambda M_T^2}{32\pi^2} \left(\log \frac{M_T^2}{\mu^2} - 1 \right) \longrightarrow M_T/T$$

[$m=0$ and renormalization by modified MS scheme
Results agree with PRD63 065003 by J.-P. Blaizot et al.]

- Scattering processes leading to damping of single particle excitation

Numerical results for $p=0$ & Conclusions

- Weak coupling ($g = 0.1$) :
Results with KB and generalized KB ansatz
are in good agreement ; No memory effect.

Separation of scale : $1/\gamma(0, t=\infty) \gg t_{\text{osci}}$
(at which oscillation of $\gamma(t)$ ends.)

$$\delta N(t) \sim \exp(-2 \gamma(t=\infty) t) \delta N(0)$$

- Strong coupling ($g = 1.6$) :
($1/\gamma(0, t=\infty) \sim t_{\text{osci}}$)
 - KB ansatz — exponential damping
 - GKB ansatz — non-exponential damping, oscillation

Memory effect influences relaxation

**QCDSP/QCDOC: Physics Results
And Prospects/Project Status**

Norman Christ and Robert Mawhinney

**QCDOC Project
and
Physics Overview**

2003 RBRC REVIEW

November 21, 2003

**Norman H. Christ
Robert D. Mawhinney**

QCDOC PROJECT

- Physics: QCD is a complete, fundamental theory. For many important problems the only significant errors are numerical: Immediate scientific reward from increased computer capability.
- Architecture: Large gains possible from optimized design.
 - Space-time homogeneity supports easy parallelization and a mesh network.
 - System-on-a-chip technology permits a highly scalable and cost-effective design:
 - * Entire node (including interconnect logic) on a single chip.
 - * The only extra components:
 - Serial nearest-neighbor wires.
 - Commercial Ethernet tree for booting, diagnostics and I/O.
 - Low power, compact design.
- Goal: 10 Teraflops QCDOC machine offers $20\times$ boost to RBRC physics.

DESIGN

- IBM-fabricated, single-chip node.
[50 million transistors, 3-4 Watt, 1.3cm×1.3cm die]
- PowerPC 32-bit processor
 - 1 Gflops, 64-bit IEEE FPU.
 - Memory management.
 - GNU and XLC compilers.
- 4 Mbyte on-chip memory and up to 2.0 Gbyte/node on DIMM card.
- 6-dim communications network:
 - Efficient for small packet sizes, $\approx 200\text{ns}$ latency.
 - Global sum/broadcast functionality.
 - Minimal processor overhead.
 - Lower dimensional machine partitions.
- 100 Mbit/sec, Fast Ethernet
 - JTAG/Ethernet boot hardware.
 - Host-node OS communication.
 - Disk I/O.
 - RISCWatch debugger.
- ≈ 5 Watt, 10 in^3 per node.

MACHINE OVERVIEW

DAUGHTER BOARDS

MOTHER BOARDS

- Serial communications simulated and tested.

– Simulation:

– Real 6 meter cable:

PERFORMANCE[†]

Operation	Local Vol.	Mflops/node
Optimized assembly code :		
SU3-SU3	-	800
SU3-2spinor	-	780
DAXPY	-	190
ZAXPY	-	450
DAXPY-Norm	-	350
Wilson D_{eo}	2^4	470
Wilson D_{eo}	4^4	535
Clover D_{eo}	2^4	560
Clover D_{eo}	4^4	590
Staggered D_{eo}	2^4	370
Staggered D_{eo}	$2^2.4^2$	430
Asqtad D_{eo}	4^4	440
Comparison with compiled code :		
CloverTerm/asm	-	790
CloverTerm/gcc, no dcbt	-	150
CloverTerm/xlc, no dcbt	-	300
Generic (MILC) C code with QMP :		
Staggered D_{eo}	2^4	171
Staggered D_{eo}	4^4	207
Asqtad D_{eo}	4^4	157
Asqtad Force	2^4	140
Asqtad Force	4^4	200

From gate-level simulation of ASIC at “nominal” 500MHz speed. Communication, cache flush overhead and 20 ns pin-pin wire delay included.

[†] [Boyle, Jung, Wettig <http://arXiv.org/abs/hep-lat/0306023>]

OUTLOOK

- 2.0 Tflops (sustained) development machines February 2003.
- Production machines July 2004:
 - 5 Tflops (sustained) RBRC.
 - 5 Tflops (sustained) UKQCD.
 - 10 Tflops (sustained) US/SciDAC ?
- Price/performance of \$1/(sustained) Mflops for QCD.

$\pi - \pi$ Effective Mass — $8^2 \times 16 \times 32$ volume

$1/a = 0.978(14)\text{GeV}$, $\beta = 5.7$

$8^2 \times 16 \times 32$ and $\times 48$, $L_s = 10$

Wilson gauge action, domain wall fermions.

QCDSP/QCDOC: Physics Results and Prospects

Robert D. Mawhinney

QCDSF/QCDOC: Physics Results and Prospects

Robert D. Mawhinney
Columbia University
RIKEN-BNL Research Center

RBRC Review
Brookhaven National Laboratory
November 21, 2003

1. Actions, statistics, errors and warnings in lattice QCD
2. Features of domain wall fermions
3. Summary of QCDSF physics calculations
4. B_K : an example of actions, statistics, errors and warnings
5. 2+1 flavor dynamical domain wall fermion simulations with QCDOC

Basics of Lattice QCD

- Sample phase space for Feynman path integral for QCD using importance sampling

$$Z = \int [dU] [d\psi] [d\bar{\psi}] \exp \left\{ \sum_n [-\beta S_g + \bar{\psi}(\not{D} + m)\psi] \right\} = \int [dU] \det(\not{D} + m) \exp \left\{ \sum_n (-\beta S_g) \right\}$$

- Importance sampling: need solution to $(\not{D} + m)h = k$ for source k , not $\det(\not{D} + m)$,
- Numerically evaluate n -point functions for QCD. Example is 2-point pion correlator:

$$\begin{aligned} & \langle (\bar{\psi} \gamma_5 \tau^a \psi)_x (\bar{\psi} \gamma_5 \tau^b \psi)_y \rangle \\ &= \frac{1}{Z} \int [dU] \det(\not{D} + m) \exp \left\{ \sum_n (-\beta S_g) \right\} \text{Tr} [\gamma_5 \tau^a ((\not{D} + m)_{x,y}^{-1} \gamma_5 \tau^b (\not{D} + m)_{y,x}^{-1})] \\ & \quad \quad \quad \Updownarrow \quad \quad \quad \Updownarrow \quad \quad \quad \Updownarrow \\ & \quad \quad \quad m_{\text{dyn}} \quad \quad \quad m_{\text{val}} \quad \quad \quad m_{\text{val}} \end{aligned}$$

- For n -point functions, need h from $(\not{D} + m)h = k$ for point, box or volume source k
- Solving $(\not{D} + m)h = k$ dominates calculation, $\sim 10^8 \times 10^8$ matrix, ~ 1000 $\not{D}$ multiplies
- Leaving out $\det(\not{D} + m)$ is quenching, which removes

ASQTAD Fermion Action

- Improve discretization to reduce flavor breaking at finite lattice spacing
- Simple $\bar{\psi}(x + \hat{u}) U_u(x) \psi(x)$ changed to include other parallel transport paths
- Parallel transport well supported by grid-based computers.

417

Residual Chiral Symmetry Breaking for DWF

- Consider introducing in action a $SU(N_f)$ matrix Ω through term at $l \equiv L_s/2$

$$- \sum_x \{ \bar{\Psi}_{x,l-1} P_L (\Omega^\dagger - 1) \Psi_{x,l} + \bar{\Psi}_{x,l} P_R (\Omega - 1) \Psi_{x,l-1} \} \quad \Omega \rightarrow U_R \Omega U_L^\dagger$$

- QCD chiral Lagrangian, with $\Sigma \equiv \exp [2i\phi^a t^a / f]$ and quark mass matrix M is:

$$\begin{aligned} \mathcal{L}_{\text{QCD}}^{(2)} &= \frac{f^2}{8} \text{Tr} (\partial_\mu \Sigma \partial^\mu \Sigma^\dagger) + v \text{Tr} [M \Sigma + (M \Sigma)^\dagger] \\ &+ v' \text{Tr} [\Omega \Sigma + (\Omega \Sigma)^\dagger] + v'' \text{Tr} [\Omega M^\dagger + \Omega^\dagger M] \end{aligned}$$

- For modes bound to walls of fifth dimension, Ω enters Green's functions as

$$\Omega e^{-\alpha L_s}$$

- Chiral condensate found by differentiating w.r.t. mass

$$-\langle \bar{q} q \rangle(m_f, L_s) \sim v + v''(\Omega + \Omega^\dagger)$$

Operator Mixing and Chiral Symmetry

- Presence of lattice chiral symmetry markedly helps operator mixing
- Consider $Q^{(\Delta S=2)}$ as an example

$$\begin{aligned}\bar{s}^{\text{lat}}\gamma_\mu(1-\gamma_5)d^{\text{lat}}\bar{s}^{\text{lat}}\gamma_\mu(1-\gamma_5)d^{\text{lat}} &\equiv (\bar{s}^{\text{lat}}d^{\text{lat}})_{V-A}(\bar{s}^{\text{lat}}d^{\text{lat}})_{V-A} \\ &= Z_1(\mu a)(\bar{s}d)_{V-A}(\bar{s}d)_{V-A} \\ &\quad + Z_2(\mu a)(\bar{s}d)_{V+A}(\bar{s}d)_{V+A} \\ &\quad + Z_3(\mu a)(\bar{s}d)_{P-S}(\bar{s}d)_{P-S} \\ &\quad + Z_4(\mu a)(\bar{s}d)_{P+S}(\bar{s}d)_{P+S} \\ &\quad + Z_5(\mu a)(\bar{s}d)_T(\bar{s}d)_T\end{aligned}$$

- For DWF, Z_2, Z_3, Z_4, Z_5 are $\mathcal{O}(m_{\text{res}}^2)$, so small

The Kaon B Parameter, B_K

- RBC Collaboration: using NPR and DWF, $a^{-1} = 1.922(40)$ GeV, $16^3 \times 32 \times 16$ lattice, Wilson gauge action. $B_K(2 \text{ GeV}) = 0.532(11)$
- CP-PACS Collaboration: using pert. th. and DWF, $a^{-1} = 1.875(56)$ GeV, $16^3 \times 40 \times 16$ lattice, RG-improved gauge action. $B_K(2 \text{ GeV}) = 0.564(14)$
- CP-PACS Collaboration: for $a \rightarrow 0$, they quote $0.5746(61)(191)$, (stat) (scaling)

Results for Real A_2

- Lattice determination from $K^+ \rightarrow \pi^+$
- Uses known chiral logarithm in quenched theory

$$1 - \frac{6m_M^2}{(4\pi f)^2} \ln(m_M^2/\Lambda^2)$$

- Fit gives coefficients of m_M^4 terms

Major Physics Projects on QCDSF

- Both quenched and full QCD simulations with DBW2 action
- A wide variety of observables being measured.
- Full QCD lattices being analyzed on remaining machines
- Full QCD simulations have been sped up by $\approx 3\times$

427

Machines	a^{-1} (GeV)	L (fm)	N_f	m_f	L_s	m_{res}
300 Gflops RBRC $225 + 3 \times 25$	2.9	1.6	0		10	$\approx m_s/300$
200 Gflops CU	1.7	1.9	2	$\approx m_s/2$	12	$\approx m_s/10$
100 Gflops RBRC	1.7	1.9	2	$\approx 3m_s/4$	12	$\approx m_s/10$
100 Gflops RBRC	1.7	1.9	2	$\approx m_s$	12	$\approx m_s/10$
100 Gflops RBRC	1.3	2.5	2	$\approx m_s/2$	12	?

The Topological Charge

- Define by “smearing” gauge field
- Should preserve topological features beyond lattice scale

$m_{\text{dyn}}=0.02$ Topological Charge

Dynamical DWF at finite T

- $16^3 \times 6$ lattice with $L_s = 32$ and $m_f = 0.022$
- Large residual mass
- Pions still as light as other actions, except special Goldstone of staggered.
- See topology change in evolution

Plaquette distributions for dynamical DWF+DBW2

- Fermion determinant changes running of coupling
- Makes gauge coupling at lattice scale stronger Stronger coupling gives more lattice dislocations

RESEARCH SUMMARY - THEORY

Research Summary

Taku Izubuchi

We report preliminary results from an ongoing calculation of B_K for $N_f = 2$ dynamical QCD with domain wall fermions. Simulations have been done with three dynamical quark masses on $16^3 \times 32$ volumes with $L_s = 12$, where the lattice spacing is $a^{-1} = 1.81(6)$ GeV. Using measurements on ~ 70 lattices for each dynamical mass and extrapolating $m_{\text{dyn}} = m_{\text{val}}$ to the kaon point, we find $B_K^{\text{MS}}(\mu = 2\text{GeV}) = 0.503(20)$ and in the range of 0.50 to 0.55 from various way analysis.

I. INTRODUCTION

After two years of collaborative efforts of the Riken-BNL-Columbia collaboration addressing for the full QCD calculation, we are starting to produce important results for phenomenology this year. Here I report our preliminary results relevant for *Cabibbo-Kobayashi-Maskawa matrix*. Other interesting results including quark masses[5] and other Kaon matrix elements[6] would be reported by other members of RBC collaboration.

The uniqueness of our joint work is characterized by the use of the chirally improved fermion, domain wall fermions (DWF), as two flavor *dynamical* quarks on lattice. Using dynamical DWF we could investigate hadron physics on lattice under one of the most ideal conditions :

1. reduced quenching errors, $N_F = 0 \rightarrow N_F = 2$
2. small scaling violation, $O(a^2)$,
3. almost negligible unphysical (chirally disallowed) mixings between operators,
4. precise renormalization factors owing to the non-perturbative renormalization.

The kaon B parameter, B_K , is an essential bridge between $K_0 - \bar{K}_0$ mixing experiments and the CKM matrix. Using DWF, quenched calculations of B_K have been done [1–3], where operator mixing and $O(a)$ errors are small because of the good chiral properties of DWF. Including dynamical fermions is an obvious way to improve these calculations.

Using an improved HMC algorithm for DWF [4, 5], we have generated three ensembles of lattices of size $16^3 \times 32$ for $N_f = 2$ QCD with degenerate dynamical masses of $m_{\text{dyn}} = 0.02, 0.03, 0.04$. Our fermion action is the standard DWF action with $L_s = 12$ and $M_5 = 1.8$ and our gauge action is DBW2 with $\beta = 0.80$. Extrapolating the rho mass in m_{dyn} to the chiral limit, we find a lattice spacing of $a^{-1} = 1.81(6)$ and a residual mass $m_{\text{res}} = 0.00136(5)$. We measure on lattices separated by 50 trajectories and have 78, 84, 57 lattices for $m_{\text{dyn}} = 0.02, 0.03$ and 0.04 , respectively[10]. For more details see [5] and for results with $\Delta S = 1$ matrix elements see [6].

II. CALCULATION OF B_K

B_K is defined by

$$B_K \equiv \frac{\langle \bar{K}^0 | O_{LL} | K^0 \rangle}{\frac{8}{3} \langle \bar{K}^0 | A_\mu | 0 \rangle \langle 0 | A^\mu | K^0 \rangle} , \quad (1)$$

where $A_\mu = \bar{s} \gamma_\mu \gamma_5 d$ and O_{LL} is the $\Delta S = 2$ four-quark operator $(\bar{s}d)_{V-A} (\bar{s}d)_{V-A}$. For a general value of the mass of the pseudoscalars entering (1), we will use the symbol B_{PS} , reserving B_K for the case where the pseudoscalars have the kaon mass.

We first evaluate this by the conventional method, where the ratio of a three-point Green's function and two pseudoscalar-axial-vector correlators is taken:

$$R(t) = \frac{\langle J_5^a(t_0) O_{LL}(t) J_5^a(t_1) \rangle}{\frac{8}{3} \langle J_5^a(t_0) A_4^a(t) \rangle \langle A_4^a(t) J_5^a(t_1) \rangle} . \quad (2)$$

(no sum on a). We use Coulomb gauge fixed wall sources at $t_0 = 4$ and $t_1 = 28$ for valence masses, m_{val} , equal to 0.01, 0.02, 0.03, 0.04 and 0.05. Propagators are found with both periodic and anti-periodic temporal boundary conditions and averaged to double the temporal length as sketched in Figure 1.

FIG. 3: $B_{\text{PS}}^{\text{lat}}$ for different valence and dynamical masses. The filled symbols have $m_{\text{val}} = m_{\text{dyn}}$.

partially quenched formula [7]

$$B_{\text{PS}}(m) = b_0 \left(1 - \frac{6}{(4\pi f)^2} \left(M_{\text{PS}}^2 \log \frac{M_{\text{PS}}^2}{\Lambda_{\chi\text{PT}}^2} \right) \right) + b_1 M_{\text{PS}}^2, \quad (3)$$

fitting the three points with $m_{\text{val}} = m_{\text{dyn}}$ to determine b_0 and b_1 . We find $b_0 = 0.267(12)$, $b_1 = 1.256(79)$ and $B_K^{\text{lat}} = 0.541(9)$ (statistical error only), with a small $\chi^2/\text{d.o.f.}$ Note that we are using the known coefficient of the chiral logarithm, $M_{\text{PS}}^2 = 4.06 \times (m_f + m_{\text{res}})$, $f = 0.0786$ and $\Lambda_{\chi\text{PT}} = 1$ GeV. The statistical error in the lattice spacing gives an additional 2% systematic error in the value of B_K^{lat} , due to the uncertainty in the strange quark mass.

While our data makes extrapolation to $m_f = \bar{m}_s/2$ quite robust, the value of b_0 is much less solid. A small value in the chiral limit has also been seen in the quenched calculations, but more control over the extrapolation is needed.

We have also fit the dependence of B_{PS} on m_{val} to (3) for each fixed m_{dyn} , using the five values for m_{val} . The interpolated values for $B_K^{\text{lat}}(m_{\text{val}} = \bar{m}_s/2, m_{\text{dyn}})$ are 0.537(11), 0.557(9) and 0.568(10) for $m_{\text{dyn}} = 0.02$, 0.03 and 0.04 respectively. We reduce the fitting range to exclude the heaviest two points and see the results stay within error.

An alternative method of extracting B_K is to divide the three-point function in the numerator of (1) by a wall-wall pseudoscalar correlator [1]. This removes some zero mode effects in the quenched approximation, while giving larger statistical errors. However here we see no difference in the central values for results calculated using these two methods.

III. RENORMALIZATION AND MIXING

To get to $B_K^{\overline{\text{MS}}}$ we need to renormalize B_K^{lat} and control mixing with other dimension six, chirally disallowed operators. Using NPR [8] in the quenched theory we have seen that mixing coefficients with chirally disallowed operators are small. Preliminary NPR calculations on our ~ 40 dynamical DWF lattices also show that the mixing coefficients are small (less than a few percent).

CKM Unitarity Triangle from DynDWF

PRELIMINARY

FIG. 4: Preliminary results of CKM unitarity triangle. The allowed region is shown from this work ($N_F = 2$) and that of quenched calculations ($N_F = 0$).

quark masses. So the direct simulation at the light quark masses would be another interesting candidate for future projects on QCDOC.

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- [1] T. Blum *et. al.* (2001), hep-lat/0110075.
 - [2] A.Ali Khan *et. al.* (2001) Phys. Rev. D64 114506.
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 - [4] T. Izubuchi (RBC) NPB (PS) 119 (2002) 813.
 - [5] C. Dawson (RBC), to appear in NPB (Proc. Suppl.).
 - [6] R. Mawhinney (RBC), to appear in NPB (Proc. Suppl.).
 - [7] S. R. Sharpe, Phys. Rev. D46 (1992) 3146, M. F. L. Golterman and K.-C. Leung, Phys. Rev. D56 (1997) 2950.
 - [8] C. Dawson, NPB (PS) 94 (2001) 613.
 - [9] S. Aoki *et. al.* Phys. Rev. D67 (2003) 094502.
 - [10] This includes some data generated after the conference.
 - [11] The author thanks Dr. N. Yamada for providing the program to draw this figure.

Research Summary

Sangyong Jeon

Energy Loss of Leading Partons in QGP

Sangyong Jeon

McGill University

Riken-BNL Research Center

Work done with G. Moore (McGill)

Energy Loss Mechanism

- Medium induced gluon Bremsstrahlung
- Must take into account coherence effect
 - * Original study by Landau, Pomeranchuk and Migdal (LPM)
 - * QCD applications by Xin-Nian Wang and Miklos Gyulassy and also Baier, Dokshitzer, Mueller, Peign and Schiff (BDMPS)
 - * We use recent resummation method by Arnold, Moore and Yaffe (AMY) – Full leading order result in high T QCD
 - * Lesson: For asymptotically high energy partons

$$\Delta E \propto \sqrt{E}$$

instead of $\Delta E \propto E$ (Bethe-Heitler limit)

Unwarranted approximations? – I

- Approximating the whole passage by the average dE/dx is bad – The energy spectrum shifts as well as broadens

Vertical lines: Energy distribution shifted by $\int dt (dE/dt)$ at various times.

Curves: Actual energy distributions at various times.

Energy in unit of temperature.

Full calculation versus dE/dx approximation at various times.
 Energy in unit of temperature.

Experimental result from PHENIX (nucl-ex/0308006)

CURRICULA VITAE

CURRICULA VITAE - SUMMARY
RBRC FELLOWS/RESEARCH ASSOCIATES/RESEARCHERS

Yasuyuki Akiba **Birthplace:** Tokyo, Japan **DOB:** October 20, 1959

D.S. 1988, University of Tokyo

Experience Research Associate, Institute for Nuclear Study, University of Tokyo,
 March 1988 – March 1997;
 Research Associate, High Energy Accelerator Research Organization (KEK),
 April 1997 – March 2003;
 Senior Research Scientist, RIKEN, April 2003 – July 2003;
 Senior Researcher, RIKEN, August 2003 – present;
 RIKEN Spin Program Researcher, November 2003 – present

Yasumichi Aoki **Birthplace:** Gunma, Japan **DOB:** January 7, 1968

Ph.D. 1996, University of Tsukuba, Japan

Experience: COE Research Fellow at the Center for Computational Physics
 University of Tsukuba;
 Assistant (3 year research position) at the Center for Computational
 Physics, University of Tsukuba
 Research Associate, RIKEN BNL Research Center, May 1, 2000 – April 30, 2003.
 RBRC Visiting Scientist, with Columbia University, May 1 to July 31, 2003.

Steffen A. Bass **Birthplace:** Frankfurt, Germany **DOB:** May 17, 1968

Ph.D. 1997, J. W. Goethe Universität Frankfurt, Germany

Experience: Feodor Lynen Fellow and Research Associate, Duke University, 1998-99
 Visiting Assistant Professor, Michigan State University, 1999
 •RHIC Physics Fellow/Assistant Professor--RBRC/Duke,
 September 1, 2000 – present

Awards and Honors: W. E. -Heraeus-Award, W.E.-Heraeus Foundation, Germany, 1993
 Feodor-Lynen Fellow, A. v. Humboldt Foundation, Germany, 1997
 U.S. Department of Energy division of Nuclear Physics Outstanding
 Junior Investigator Award, 2003.

Alexander V. Bazilevsky **Birthplace:** Yaroslavl, Russia **DOB:** May 10, 1968

Ph.D. 1999, Institute for High Energy Physics, Protvino, Russia

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 Research Collaborator, BNL PHENIX
 Research Associate, BNL, RIKEN BNL Research Center, March 10, 1999 –
 March 9, 2002;
 RIKEN BNL Fellow, March 10, 2002 – April 1, 2003.

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Yuji Goto **Birthplace:** Shizuoka, Japan **DOB:** November 25, 1965
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 Postdoctoral Fellow, RIKEN, Japan, 1996-1999
 RIKEN BNL Fellow, November 1999 – March 31, 2002
 Scientist, RIKEN, April 2002 to March 2003;
 Senior Research Scientist, RIKEN, April 2003 – present.
 RIKEN Spin Program Researcher, RBRC, April 1, 2002 – present.

Matthias Grosse-Perdekamp **Birthplace:** Schwenningen, Germany **DOB:** December 1, 1963
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Experience: Associate Research Scientist, Yale University 1995-1998
 Research Scientist, Mainz University, 1998-1999
 RIKEN BNL Fellow, January 1999 – August 31, 2002.
 Member of the PHENIX and BELLE Collaborations;
 •RHIC Physics Fellow/Assistant Professor—RBRC-E; University of Illinois,
 Urbana Champaign, September 1, 2002 – present.
Awards and Honors: Foreign Scholar Award, UCLA, 1990; Gustav Mie Preis, Freiburg
 University, 1991.

Yoshitaka Hatta **Birthplace:** Kyoto, Japan **DOB:** August 5, 1976
Master Degree
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 Spring 2002 – 2003.
 RIKEN Junior Research Associate (Theory), April 1, 2003 – present.

Tetsufumi Hirano **Birthplace:** Kanagawa, Japan **DOB:** September 17, 1972
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 Postdoctoral Fellow at University of Tokyo, Japan, April 2001 to March
 2003.
 RIKEN Spin Program Research Associate, RBRC Theory Group,
 March 24, 2003 – present.
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 (2002).

Takuma Horaguchi **Birthplace:** Iwate, Japan **DOB:** September 29, 1977
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 Spring 2002 – 2003.
 RIKEN Junior Research Associate (Experiment), April 1, 2003 – present.

- Osamu Jinnouchi** **Birthplace:** Saga, Japan **DOB:** June 13, 1972
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Contract Researcher of RIKEN, April 2001-2003.
Research Associate, RBRC Experimental Group, April 1, 2003 – present.
- Masashi Kaneta** **Birthplace:** Hiroshima, Japan **DOB:** November 5, 1971
Ph.D. 1999, Hiroshima University, Japan
Experience: Member of NA44 Experiment at CERN, 1994-present
Member of Beam-Beam Counter (BBC) Group for PHENIX Experiment,
1994-1999; Member of TOF Detector Group in PHENIX Experiment at RHIC,
1999; Member of STAR Experiment at RHIC, 1999-present;
Postdoctoral Fellow at KEK, 1999; Postdoctoral Fellow Physicist at
Lawrence Berkeley National Laboratory, 1999-2002.
Research Associate, BNL, RBRC Experimental Group, Nov. 1, 2002-present.
- Hideyuki Kobayashi** **Birthplace:** Japan **DOB:** January 26, 1971
Ph.D. 2000, Tokyo Institute of Technology, Faculty of Science
Experience: Cooperate Researcher, RIKEN, CCJ, May 1-Sept. 30, 2000
Special Postdoctoral Researcher, RIKEN, October 1, 2000 – September 30, 2003;
RIKEN Spin Program Research Associate, April 1, 2001 – September 30, 2003.
Awards and Honors: Fellowship: Japan Society for Promotion of Science (JSPS), Research
Fellow, April 1, 1996 – March 31, 1999.
- Stefan A. Kretzer** **Birthplace:** Dortmund, Germany **DOB:** November 26, 1968
Ph.D. 1999, University of Dortmund
Experience: Visiting Research Associate at Michigan State University, 2000-2002
Research Associate, Brookhaven National Laboratory, October 2002 –present;
Research Associate, Joint RBRC/Nuclear Theory Group, BNL, April 2003 present.
Awards and Honors: Associate of the Graduiertenkolleg: Erzeugung und Zerfälle von
Elementarteilchen of the Deutsche Forschungsgemeinschaft (DFG)
Scholarship, 1996-2000.
- Kazuyoshi Kurita** **Birthplace:** Tokyo, Japan **DOB:** August 11, 1963
Ph.D. 1992, Columbia University, New York
Experience: June, 1991 Research Associate, Univ. of Tsukuba
April 1994 Assistant Prof., Univ. of Tsukuba
April 1997 Postdoctoral Researcher, RIKEN
Oct. 1997 Special Postdoctoral Researcher, RIKEN
RIKEN BNL Fellow, April 2000 – March 31, 2002;
Visiting Scientist, Experimental Group, RBRC, with Rikkyo University, Japan,
April 2002 - present
Awards and Honors: Educational Research Award, Tsukuba Gakuto Foundation, July 1992

Naohito Saito **Birthplace:** Aomori, Japan **DOB:** November 28, 1964
Ph.D. 1995, Kyoto University, Japan
Experience: 1993 July, Research Fellow, Kyoto University, Japan
 1995 April, RIKEN Special Post Doctoral Fellow
 RIKEN/RBRC Researcher, 1996 April to March 2001
 RIKEN Spin Program Researcher, April 1, 2001 – March 31, 2002;
 Visiting Scientist, Experimental Group, RBRC, with Kyoto University, Japan,
 Spring 2002 – present.

Awards and Honors: Fellowship: Japan Society for the Promotion of Science, 1993-1995

Thomas M. Schaefer **Birthplace:** Hanau, Germany **DOB:** May 18, 1965
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 Postdoctoral Research Associate Institute for Nuclear Theory, University of
 Washington; Member Institute for Advanced Study, Princeton;
 • RHIC Physics Fellow/Assistant Professor--RBRC/ SUNY, Stony Brook,
 January 2000 – December 31, 2002.
 •RHIC Physics Fellow/Associate Professor (with Tenure)—RBRC/North Carolina
 State University January 1, 2003 – present.

Awards and Honors: Member, Studienstiftung des deutschen Volkes;
 Fellowship, German Academic Exchange Service
 Feodor Lynen Fellowship, Alexander v. Humboldt Foundation
 U.S. Department of Energy Division of Nuclear Physics Outstanding Junior
 Investigator Award, 2002.

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Mikhail Stephanov **Birthplace:** Moscow, Russia **DOB:** April 19, 1966
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 Postdoctoral Research Associate, ITP, SUNY at Stony Brook
 •RHIC Physics Fellow/Assistant Professor --RBRC/
 University of Illinois at Chicago, October 1, 1999 to present
 (2003 Promoted to Associate Professor with tenure).
Awards and Honors: Soros Scholarship, Overseas Graduate Scholarship from Jesus College,
 Oxford; U.S. Department of Energy Division of Nuclear Physics Outstanding Junior
 Investigator Award, 2001; Alfred P. Sloan Fellowship, 2002.

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Experience: Research Associate, RCNP, Osaka University, May 1997 – March 1999;
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Ph.D. 1992, SUNY at Stony Brook

Experience: Research Associate, Theoretical Physics Institute, University of Minnesota, September 1992-September 1994
Research Associate, National Institute for Nuclear Theory, University of Washington, September 1994-December 1996
Research Assistant Professor and Danish Research Council Fellow, Niels Bohr Institute, Denmark, January 1997 -October 1998
Assistant Physicist, Physics Department, BNL - October 1998
Associate Physicist, Physics Department, BNL - Nov. 1999 – Oct. 2001
Physicist, Physics Department, BNL – Oct. 2001 – June 1, 2002
Physicist with Tenure, Physics Department, BNL – June 1, 2002 – present.
•RHIC Physics Fellow/ RBRC/BNL - October 1, 2000 – June 30, 2003.

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Ph.D. 1993, University of Dortmund, Germany

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Research Associate in Theory Group of the Rutherford Appleton Laboratory, Didcot, England, Oct. 1994-Dec. 1996
Fellow in Theory Division, CERN, Geneva, Switzerland, Jan. 1997-Dec. 1998
Post-doctoral Researcher at the Institute for Theoretical Physics at State University of New York, Stony Brook, January 1999 to March 2000
RIKEN BNL Fellow, April 1, 2000 – September 30, 2003.
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Awards and Honors: Scholar of the "Studienstiftung des Deutschen Volkes, 1984-1990

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Ph.D. 1993, University of Tokyo, Japan

Experience: Scientific Researcher, RIKEN, 1991
Research Collaborator, PHENIX/RBRC
Scientific Researcher, RIKEN, RBRC, 1998-2001
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Research Associate, Max-Planck-Institute for Nuclear Physics, Heidelberg, 1995-1996;
Research Associate, Institute for Theoretical Physics, Technical University of Munich; External Scientific Associate, Max-Planck-Institute for Nuclear Physics, Heidelberg, 1997-1999;
•RHIC Physics Fellow/Assistant Professor--RBRC/ Yale University, October 1999 - present; (2003 Promoted to Associate Professor at Yale)

Awards and Honors: 1st Prize, National Physics Olympiad, East Germany
Scholar of the "Studienstiftung des deutschen Volkes"
Heisenberg Fellowship of the German National Science Foundation

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RIKEN BNL Research Center
Theory Group
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RIKEN BNL RESEARCH CENTER

RBRC Scientific Review Committee Meeting

November 20-21, 2003

Li Keran

*Nuclei as heavy as bulls
Through collision
Generate new states of matter.
T.D. Lee*

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Speakers:

Y. Akiba	S. A. Bass	A. Bazilevsky	T. Blum	N. H. Christ
C. Dawson	A. Deshpande	H. En'yo	D. Fields	Y. Goto
M. Grosse Perdekamp	T. Hirano	T. Ikeda	O. Jinnouchi	M. Kaneta
A. Kusenko	R. D. Mawhinney	L. McLerran	Y. Nemoto	J. Noaki
K. Okada	N. P. Samios	T. Schaefer	M. Stephanov	T. Sugihara
K. Tanida	U. van Kolck	W. Vogelsang	Y. Watanabe	T. Wettig
N. Yamada				

Session Chairs: A. Baltz, G. Bunce, H. En'yo, S. Ohta
Organizers: N. P. Samios and T. D. Lee