Standard-Model Higgs Boson Searches at CDF in Run II
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The SM Higgs boson has been searched in two channels: (1) \(q\bar{q} \rightarrow WH \rightarrow \ell\nu b\bar{b}\) for Higgs masses \(110 < M_H < 150\) GeV and (2) \(gg \rightarrow H \rightarrow WW \rightarrow \ell\nu\ell\nu\) for \(140 < M_H < 180\) GeV, where \(\ell \in \{e, \mu\}\), using about 200 \(\text{pb}^{-1}\) CDF Run II data. 95\% C.L. limits on each Higgs production at Tevatron at \(\sqrt{s} = 1.96\) TeV were set as a function of Higgs mass by fitting the distribution of dijet mass (1) and dilepton azimuthal angular separation (2). We have significantly advanced the sensitivity to the SM Higgs boson cf. Run I.

1. Introduction

In the Standard Model (SM) and its supersymmetric (SUSY) extensions, the Higgs boson is crucial to our understanding of electroweak symmetry breaking dynamics and the mass generation of electroweak gauge bosons and fermions. The Higgs boson mass is in addition an indicator of the scale of new physics. Electroweak precision measurements and LEP direct searches have constrained the Higgs mass to be within 114.4 and 251.1 GeV [1]. However, while almost all the high energy particle physics phenomena observed so far have been successfully explained by the Standard Model, the Higgs boson remains unobserved. It is the last missing piece and the most important one.

2. Higgs at Tevatron

Tevatron is currently the only machine capable of probing the Higgs sectors in the world, of which the luminosity delivery has surpassed the design goal for FY2004. Higgs reach of Tevatron had been evaluated in [2] and [3], Figure 1, and concluded hopeful with the Run II end goal (8 \(\text{fb}^{-1}\)) integrated luminosity and full strengths of both upgraded CDF and DØ experiments. CDF II has great improvement on energy measurement and dijet mass resolution over Run I [4]. The capabilities of CDF II on lepton ID and b-tagging can be found in [5] and [6]. Using data cumulated so far CDF and DØ can already probe the SM Higgs sector with precision measurement of electroweak parameters, such as top quark mass and \(W\) boson mass, Figure 2.

3. \(WH \rightarrow \ell\nu b\bar{b}\) Search at CDF II

CDF has searched for new particles decaying into \(b\bar{b}\) and produced in association with \(W\)
bosons decaying into $\ell \nu$, where $\ell \in \{e, \mu\}$, in $pp$ collisions at $\sqrt{s} = 1.96$ TeV [7]. Such new particles could be the SM Higgs bosons if with zero spin. The final state in search consists of one high-$p_T$ lepton, large missing transverse energy $E_T$, and two $b$ jets.

### 3.1. Data Samples and Event Selection

This analysis was based on 162 pb$^{-1}$ of TeVatron Run II data recorded by the CDF II detector [4] from March 2002 to September 2003, with an inclusive lepton trigger that requires an electron with $E_T > 18$ GeV or a muon with $p_T > 18$ GeV. From the inclusive lepton dataset, events reconstructed with

1. an isolated non-conversion electron with $E_T > 20$ GeV in the central region or non-cosmic muon with $p_T > 20$ GeV;
2. $E_T > 20$ GeV;
3. $E_T > 15$ GeV (cuts optimized for the best $S/\sqrt{B}$) and $|\eta| < 2$, with at least one jet $b$-tagged with CDF secondary vertex tagging algorithm (SECVTX) [6] to improve signal against $W +$ jets background;

were selected. To reduce top-contributed backgrounds, events with an isolated track having $p_T^{\text{seed}} > 20$ GeV and charge opposite to the primary lepton’s were removed. Events with additional jets (one with $E_T > 8$ GeV and $2 < |\eta| < 3$ or two with $8 < E_T < 15$ GeV and $|\eta| < 2$) were also removed.

### 3.2. Background

The dominant background was QCD production of $W +$ jets, which entered the signal sample when either one of the jets was $b$-tagged or a light-quark jet was mis-identified as a heavy one. Mistags were mostly due to symmetric resolution effects and were estimated by

1. parameterizing the negative tag rate in the generic QCD multijet sample as a function of jet $E_T$, track multiplicity, $\eta$, $\phi$ and the event sum of jet $E_T$s. The negative tag rate was enlarged by $1.2 \pm 0.1$ to account for the presence of long-lived particles and secondary interactions with detector material;
2. applying the predicted tag rates to the pretag sample.

Fraction of $W +$ heavy-flavour jets in the inclusive $W +$ jets sample was estimated using Alpgen [8] Monte Carlo. Background of $W +$ heavy-flavour jets was estimated by applying the estimated fraction and the tagging efficiency to the pretag sample.

Another substantial background came from mis-measured QCD events, in which one jet was mis-identified as a high-$p_T$ lepton and mis-measured energies made apparently large $E_T$. To estimate this background, the number of tagged events with an isolated lepton and low $E_T$ was extrapolated into the signal region.

Top-contributed backgrounds were estimated using PYTHIA [9] and assuming NLO cross-sections. Other backgrounds were estimated using suitable MC.

Summary of tagged data events and expected backgrounds as a function of jet multiplicity is in Table 1, showing good agreement between data and expectation.

### 3.3. Results

#### 3.3.1. Acceptance and Uncertainties

Signal acceptance was estimated using both data and PYTHIA MC, Figure 3. The total acceptance was calculated as a product of geometric and kinematic acceptance, lepton ID, trigger
and $b$-tagging efficiencies. Systematic uncertainty (11% in total) came from the modelling of initial (3%) and final (6%) state radiation, parton distribution function (2%), jet energy scale (3%) and resolution (1%), $b$-tagging (6%), lepton trigger (< 1%) and ID (5%) efficiencies.

![Figure 3. The total WH $\rightarrow \ell\nu b\bar{b}$ acceptance as a function of the Higgs mass.](image)

### 3.3.2. 95% C.L. Limits

A direct search for a resonant mass peak in the reconstructed dijet mass distribution did not yield any positive result, Figure 4. So 95% C.L. limits on $\sigma \cdot \text{BR}(p\bar{p} \rightarrow WH \rightarrow \ell\nu b\bar{b})$ were set as a function of the Higgs mass using a binned maximum likelihood technique, Equation (1), on the numbers from the $W + 2$ jets bin.

$$L \equiv \prod_i \frac{\mu_i N_i \cdot e^{-\mu_i}}{N_i!}$$

where $\mu_i$ was the expected signal plus background and $N_i$ was the number of observed events from the $W + 2$ jets sample in the $i$th mass bin. The limits, shown in Figure 7, were greatly improved over Run I [10] although the SM reach was still limited by statistics.

### 4. $H \rightarrow WW \rightarrow \ell\nu\ell\nu$ Search at CDF II

CDF has searched for the high-mass SM Higgs production and decay $p\bar{p} \rightarrow H \rightarrow W^+W^-$.
\[ \ell^+ \nu \ell^- \] where \( \ell = \{e, \mu\} \) at \( \sqrt{s} = 1.96 \text{ TeV} \) [11].

4.1. Data Samples and Event Selection

The analysis was based on \( 184 \text{ pb}^{-1} \) of CDF II data collected from March 2002 to September 2003, with one of the following inclusive high-\( p_T \) lepton triggers:

1. a central electromagnetic cluster with \( E_T > 18 \text{ GeV} \) and \( p_T > 9 \text{ GeV} \);
2. a central muon track with \( p_T > 18 \text{ GeV} \);
3. a forward electromagnetic cluster with \( E_T > 20 \text{ GeV} \) and \( \eta > 15 \text{ GeV} \).

The selection of a candidate event requires

1. two electron/muons with opposite charges;
2. \( E_T > 25 \text{ GeV} \);
3. no jets with \( E_T > 15 \text{ GeV} \) and \( |\eta| < 2.5 \);
4. large azimuthal angular separation between small \( E_t \) and any lepton or jet, \( \Delta \phi(E_t, \ell/j) > 20^\circ \) if \( E_t < 50 \text{ GeV} \), to lower \( Z \rightarrow \tau \tau \) background;
5. large \( E_t \) significance [12] if \( |M_{\ell\ell} - M_Z| < 15 \text{ GeV} \), to lower \( Z \rightarrow ee \) and \( Z \rightarrow \mu \mu \) background;
6. small \( M_{\ell\ell}^{-1} \), cut as shown in Table 3.

The \( H \rightarrow WW \) signals were discriminated from the \( WW \)

4.2. Acceptance

Signal acceptance, as shown in Table 3, was estimated using PYTHIA MC with CTEQ5L parton distribution functions.

4.3. Background

Drell-Yan \( Z/\gamma^* \rightarrow \ell^+ \ell^- \) where \( \ell = \{e, \mu, \tau\} \) and diboson (\( WW, WZ \) \& \( ZZ \)) backgrounds were estimated with PYTHIA MC, \( tt \) with HERWIG [13], normalized to their SM cross-sections.

\( W + \text{jet} \) entered the background when the jet was mis-identified as a lepton. The \( W + \text{jet} \) background was estimated by

1. calculating the probabilities (typically of order \( 10^{-3} - 10^{-4} \)) for a jet to fake a lepton using the inclusive jet data;

2. background using the spin-zero property of the Higgs boson, which would result in small invariant mass of the two leptons (\( M_{\ell\ell} \)). Figure 5, and azimuthal angular separation in between (\( \Delta \phi_{\ell\ell} \)).
2. counting the number of events containing a faking jet and a lepton and passing the $H \to WW$ dilepton event selection from the parent signal sample;

3. applying the fake probabilities to the number of events.

Summary of signal and background expectations for the 184⁻¹⁻¹ analyzed data is given in Table 2 and shown in Figure 6.

4.4. Results

The binned maximum likelihood method was used on the $\Delta \phi_{ll}$ distribution of data and MC signal, WW and the sum of other SM backgrounds to set 95% C.L. upper limits on $\sigma \cdot sBR_{SM}(p\bar{p} \to H \to WW)$ as a function of the Higgs mass for $140 < M_H < 180$ GeV. We compare the results to those obtained by using a simple counting method. Summary of results is in Table 3, showing good agreement between data and expectation, and Figure 7.

5. Conclusion

Higgs hunting Run II at CDF has started! With currently available data, we should/can set tight Higgs constraints with precision top and electroweak measurement besides getting exciting first Higgs direct search results. Eventually we will likely make sensitive tests of the SM Higgs sector with the Run II end goal luminosity of TeVatron and full strengths of both CDF and DØ experiments.

<table>
<thead>
<tr>
<th>$M_H$</th>
<th>140 GeV</th>
<th>150 GeV</th>
<th>160 GeV</th>
<th>170 GeV</th>
<th>180 GeV</th>
</tr>
</thead>
<tbody>
<tr>
<td>DY $ee + \mu\mu + \tau\tau$</td>
<td>0.18 ± 0.11</td>
<td>0.33 ± 0.19</td>
<td>0.65 ± 0.30</td>
<td>1.06 ± 0.42</td>
<td>1.33 ± 0.48</td>
</tr>
<tr>
<td>fake</td>
<td>0.40 ± 0.12</td>
<td>0.45 ± 0.14</td>
<td>0.53 ± 0.16</td>
<td>0.65 ± 0.19</td>
<td>0.81 ± 0.25</td>
</tr>
<tr>
<td>$t\bar{t}$</td>
<td>0.01 ± 0.01</td>
<td>0.01 ± 0.01</td>
<td>0.01 ± 0.01</td>
<td>0.01 ± 0.01</td>
<td>0.02 ± 0.01</td>
</tr>
<tr>
<td>$WZ + ZZ$</td>
<td>0.10 ± 0.01</td>
<td>0.13 ± 0.01</td>
<td>0.15 ± 0.01</td>
<td>0.19 ± 0.01</td>
<td>0.24 ± 0.02</td>
</tr>
<tr>
<td>WW</td>
<td>3.50 ± 0.41</td>
<td>3.82 ± 0.45</td>
<td>4.45 ± 0.52</td>
<td>5.38 ± 0.63</td>
<td>6.49 ± 0.76</td>
</tr>
<tr>
<td>total background</td>
<td>4.19 ± 0.45</td>
<td>4.74 ± 0.52</td>
<td>5.79 ± 0.64</td>
<td>7.29 ± 0.81</td>
<td>8.89 ± 0.98</td>
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</tbody>
</table>

Figure 6. The dilepton azimuthal angular separation $\Delta \phi_{ll}$ distribution of data and MC signal and background expectations.

REFERENCES

3. FERMILAB-PUB-03-320-E.
5. hep-ex/0406078, submitted to PRL.
7. CDF/PUB/EXOTIC/PUBLIC/7126.
Table 3: Inputs and 95% C.L. limits on $\sigma \cdot \text{BR}(gg \to H \to WW)$ at $\sqrt{s} = 1.96$ TeV for $140 < M_H < 180$ GeV. The total acceptance $\epsilon_{\text{tot}}$ includes $\text{BR}(W \to \ell \nu)^2$.

<table>
<thead>
<tr>
<th>Higgs mass $M_H$</th>
<th>140 GeV</th>
<th>150 GeV</th>
<th>160 GeV</th>
<th>170 GeV</th>
<th>180 GeV</th>
</tr>
</thead>
<tbody>
<tr>
<td>$\sigma_{\text{SM}}(gg \to H)$</td>
<td>0.45 pb</td>
<td>0.36 pb</td>
<td>0.30 pb</td>
<td>0.25 pb</td>
<td>0.21 pb</td>
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<tr>
<td>$\text{BR}_{\text{SM}}(H \to WW)$</td>
<td>0.48</td>
<td>0.68</td>
<td>0.90</td>
<td>0.97</td>
<td>0.94</td>
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<tr>
<td>$M_H$ cut</td>
<td>55.0 GeV</td>
<td>57.5 GeV</td>
<td>62.5 GeV</td>
<td>70.0 GeV</td>
<td>80.0 GeV</td>
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<tr>
<td>$\epsilon_{\text{tot}}$ (%)</td>
<td>0.124 ± 0.012</td>
<td>0.228 ± 0.023</td>
<td>0.402 ± 0.040</td>
<td>0.476 ± 0.048</td>
<td>0.449 ± 0.045</td>
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<tr>
<td>expected signal</td>
<td>0.10 ± 0.01</td>
<td>0.15 ± 0.02</td>
<td>0.22 ± 0.03</td>
<td>0.22 ± 0.03</td>
<td>0.17 ± 0.02</td>
</tr>
<tr>
<td>WW background</td>
<td>3.51 ± 0.41</td>
<td>3.82 ± 0.45</td>
<td>4.45 ± 0.52</td>
<td>5.38 ± 0.63</td>
<td>6.49 ± 0.76</td>
</tr>
<tr>
<td>other SM background</td>
<td>0.68 ± 0.16</td>
<td>0.92 ± 0.24</td>
<td>1.34 ± 0.35</td>
<td>1.91 ± 0.47</td>
<td>2.40 ± 0.55</td>
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<td>data candidates</td>
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<td>7</td>
<td>8</td>
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<td>limit – counting</td>
<td>18.4 pb</td>
<td>9.8 pb</td>
<td>6.2 pb</td>
<td>8.2 pb</td>
<td>8.8 pb</td>
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<td>expected limit – $\Delta \phi$</td>
<td>18.1 pb</td>
<td>9.8 pb</td>
<td>6.0 pb</td>
<td>7.4 pb</td>
<td>8.0 pb</td>
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<td>limit – $\Delta \phi_H$-binned</td>
<td>17.8 pb</td>
<td>9.4 pb</td>
<td>5.6 pb</td>
<td>5.6 pb</td>
<td>6.4 pb</td>
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</tbody>
</table>

![CDF Run II Preliminary](image.png)

Figure 7. CDF Run II 95% C.L. and expected upper limits for $WH$ and $H \to WW$ production at TeVatron, in comparison with the Standard-Model predictions.