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STATUS OF THE TAU-CHARM FACTORY PROJECT AND ASPECTS OF THE DETECTOR DESIGN[‡]

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Abstract

This paper reviews the status of the Tau-Charm Factory Project being proposed for construction in Spain. The paper also reviews characteristics of the detector design, and the issues surrounding the present choices of technologies.

THE TAU-CHARM PROJECT

The Tau-Charm Project is an e^+e^- storage ring of high luminosity presently being proposed for construction in Spain. The machine has a design luminosity of $10^{33} \text{ cm}^{-2} \text{ sec}^{-1}$ operating between 3 and 5 GeV cms energy. The physics program of the machine is broad, encompassing both τ and charm production and decay studies near threshold, as well as studies of bound state charmonium production and decay. At design luminosity, the collider will produce close to 10^8 τ - pairs/yr, 10^8 charmed mesons/yr., and 10^{10} J/ψ /yr.. As the tagging efficiency for charmed mesons and τ ~10-20%, samples exceeding 10^7 reconstructed decays will become available, far exceeding any presently existing or proposed facility. The machine is being designed with monochromator optics, allowing it to achieve cms energy spreads close to 100 KeV, and thereby allowing the production of J/ψ to be increased to 2×10^{11} /yr for CP studies in the hyperon system. Beam polarization may also be routinely available.

A high luminosity collider operated at threshold for the pair production of charm and τ , allows us to take advantage of unique kinematic constraints to enhance the analyzing power of the detector, and allowing the search for rare processes at the 10^{-7} to 10^{-8} level. The issues of detector design are critical, to fully exploit the high luminosity; they are described in the balance of this paper.

Tau-Charm Machine and Detector Interface

The characteristics of the machine strongly influence the detector design. The machine is a double ring design, with electrostatic separators to minimize parasitic crossings of the beams to the IP region alone. The rings of 360m circumference, each hold about 30 bunches, and have a crossing frequency of 25 MHz. The IR design has been evaluated by the CERN group. The vertical separators start at 3.5 meters from the IP (see Fig. 1), while the superconducting final focus quad. triplets penetrate the detector to $\pm 80\text{cm}$ from the IP.

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Aspects of The Tau -Charm Detector

The present design of the detector stresses a long barrel ($|\cos \theta| < 0.95$) containing a central drift chamber for tracking, followed by a 4m long two layer scintillating fiber time-of-flight system. The TOF itself is backed by a twofold longitudinally segmented CsI(Na) or CsI(Tl) crystal calorimeter, in a unique re-entrant geometry. The long barrel is chosen to avoid a significant endcap region. The hermeticity is completed by a forward silicon tracking system and calorimetry (luminosity monitor) comprised of BaF2 crystals. Unlike B-Factory detector designs, the use of TOF for particle identification is adequate at TCF; scintillator, being a fully active medium, doesn't compromise calorimetry performance, as most other particle ID systems would. The hermeticity is completed by a fine grained muon range system which also acts as a veto for the energy carried by neutral hadrons, and as a flux return yoke. In this way, even neutrinos can be tagged by missing energy, an essential ingredient to the effectiveness of the detector for the study of rare decays and the tagging of τ - decays.

Radiation Doses and Backgrounds

The radiation background from the large circulating currents (600ma/beam) of the TCF rings has been analyzed for the ORSAY-optics, and assuming an optimal masking scheme. The luminosity lifetime of 2 hrs is dominated by beam-beam bremsstrahlung assuming a 1nT vacuum can be maintained. Synchrotron radiation from the strong bending ($\rho \approx 12m$) of the main the rings does not strike the detector because of the vertical separation scheme, but it must be arranged to avoid illumination of the opposing electrostatic separator plates. Synch. radiation from the vertical bends themselves and the final focus quads is very soft, and can largely be absorbed in the beryllium beampipe. The major sources of detector background is from electroproduced hadrons from beam gas interactions, from lost beam particles that are coulomb or bremsstrahlung scattered, and particles lost from the injected beam itself. The results from injection losses are least understood,

and potentially the most serious as summarized in Table I. A model for injection frequency and loss rate is assumed

THE TAU-CHARM FACTORY DETECTOR

The basic features of the detector are described next. These requirements and potential solutions have largely been established through the work of the members of the US Tau-Charm Collaboration in R&D carried out since 1989. The design of the tracking system for the TCF detector is motivated by the desire to accurately measure the endpoint spectrum of the decay of the τ into multihadrons. This will ultimately lead to the best resolution possible for measurement of a limit on the ν_τ -mass to about 1 MeV. To do this, the tracker must achieve a momentum (p) resolution of $0.35\%p$ and a multiple coulomb scattering term of about 0.35%. At TCF energies, this is possible for a device having about 40 cylindrical planes of Al-guard/W-sense wires, over a radial extent of about 80 cm, immersed in a 1T axial magnetic field. We assume a low Z gas mixture such as He-CO₂-Isobutane is used, and knowledge of the IP position and sense and field wire multiple-scattering is included in the track fit. A single hit resolution of only about 150 μm is required. The TCF detector does not need a vertex chamber, so the beampipe aperture can be maintained at about $r=5cm$ through the IP, thereby reducing the broadband impedance of the machine. For the tracker, a first measurement in close proximity to the Be-pipe is important to eliminate the multiple scattering contribution to angular error in the long track extrapolation to the IP. The primary difficulty in TCF will be the 3.6m long wires required to meet the solid angleconstraints of the long barrel design. The electronics and endplate materials must also be minimized, or relocated, to avoid damage to the calorimetry performance for $0.95 > |\cos \theta| > 0.90$.

Following the tracking system, π, K, p particle ID is accomplished most directly by a system of about 96 azimuthally divided fiber TOF counters. Two radial layers, each 5cm thick of 1mm fiber, and readout on both ends with XP2020

Figure 1. Side View of the Tau - Charm Factory Detector With Proposed Technologies

equivalent photomultiplier tubes, are required to achieve a time resolution of 120ps over the 1 meter flight path. This is required primarily to give adequate background rejection ($>3\sigma$) in the measurement of $D^0\bar{D}^0$ mixing, where π and K momenta approach 1 GeV. Fiber counters are optimal because of their greater attenuation length - appropriate to the long barrel design. They can also be easily tailored to the re-entrant barrel design. In this present design, they match up to clear fiber highways, which then penetrate radially outward at the end-door breaks. The PMTs are then located outside the magnetic field. As indicated earlier, the TOF system will not degrade the calorimeter, as it itself is a fully active absorber medium.

In our present design, both CsI(Tl) and CsI(Na) crystals of 16 rl depth, read out with wavelength shifters and PIN photodiodes are being investigated (see Fig. 2). The primary advantage of CsI crystal is its large light output (about 50K

Section	Background Sources	Doses (Rad/Yr)
BGO	Beam-Gas	22
	Lost Particles	168
	0 ⁰ Rad. Bhabha	0.03
	Bhabhas	5
	Injection	3400
CDC	Non-Injection	1x10 ⁻² C/cm/yr
	Injection	0.2 C/cm/yr
TOF	Lost Particles	1.1
	0 ⁰ Rad. Bhabha	0.4
	Injection	22
CsI	Beam-Gas	10 ⁻⁴
	Lost Particles	0.23
	0 ⁰ Rad. Bhabha	0.1
	Injection	4.6

photons/MeV), however the drawback is its long time constant ($1\mu\text{sec}$, $0.5\mu\text{sec}$ for Tl and Na respectively). The consequence of this is the need to do coarse timing ($\sim 40\text{ ns}$) on the output, to reduce machine background pileup. Energy of showers must be uniquely assigned to beam crossings.

Figure 2. An Option For CsI Crystal Geometry

This will leave only the problem of albedo showers within the time window of 40 ns of a beam crossing.

To further address this remaining problem of in-time backgrounds we are investigating the use of increasing the longitudinal and transverse segmentation and readout of crystals, as illustrated in Fig. 2. These will provide additional information on shower shape, origin and direction. An extensive beam test program containing 22K cc of CsI, arranged with 132 readout channels was carried out in August 1992 at TRIUMF to study the strengths of these dual split crystal geometries and front/back wavelength shifter readout. Results will be reported in the near future.

The calorimetric solid angle is closed by a BaF₂ array in the very forward direction, that is mounted on the face of the final quads. This array doubles as the luminosity monitor, being supplemented by several planes of a large area silicon pad tracking system. This strip and pad system will be pre-processed by an onboard VLSI chip. Tests by members of the German TCF Collaboration have already been carried out at CERN on prototypes of the silicon tracking elements, and a detailed design for the crystal device is now under study by those groups.

Radially outward, following the CsI calorimeter, is the thin superconducting 1T coil. While it is outside the calorimetry, we still wish to do range measurements for muon identification, so the coil, its terminations and its cryostat must be minimized. A 0.25 absorption length device is acceptable. The muon system that follows is extremely fine grained, having steel plates of 1.5 to 2cm thickness. Typical range differences at 500 MeV/c between pions and muons are only 2 to 3 such gaps. Because of the low occupancy, we envision a large cell drift chamber between absorber layers. The requirement on the overall efficiency per layer of the device is extremely tight; it must achieve greater than 95% (including geometric occlusions) efficiency since the muon range system is also necessary as a veto for energy from neutral hadrons (K^0_L), and will impact rare decay studies if it is not sufficiently hermetic and efficient.

Electronics and Data Acquisition

Lastly, I will comment on the electronics and data acquisition. The drift chamber has about 5K wires, the calorimeter about 20K elements, forward calorimeters about 1K elements, the TOF about 1K ADC/TDC's, and the muon system about 20K wires. While this detector is then very comparable to present SLC and LEP detectors, perhaps unique to the TCF detector is the large dynamic range required in the trigger. The trigger must span from 100Hz to about 5KHz of *true* physics events. The purity of the trigger is expected to exceed 50% - namely a background rate under 100Hz from beam related sources is anticipated. A typical event should be processed in the front-end electronics to contain less than about 10 Kbytes of data. We envision extensive onboard detector signal processing and pattern recognition of raw data, directly into either digital or analog pipelines, thereby allowing a dead-timeless trigger to be formed. The final stages of the trigger would process the event more fully, applying geometric and time-varying constants, and allowing a very high level topological discrimination to take place, before writing events to a storage device. While such processing is not required for the running at energies suitable for τ and charm physics studies, this feature is essential if the machine is run with monochromator optics on the J/ψ , where the event rate will increase by 1 order of magnitude.

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