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CALCULATIO OF GEOMETRICAL EPFICIEACY OF COUIFTRRS<br>Louis R. Henrich<br>Auguat 8, 1952

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## CALCULATION OF GEONETRICAL EFFICIEMCY OF COUNIERS

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August 8, 1952

Let $F$ be the flux of radiation going from a sample of radius $R_{2}$ tio a counter of radius $R_{1}$ then $\Phi$ the emitted intensity of radiation per unit solid angle. Then the flux between two elements of area $d \sigma_{1}$ and $d \sigma_{2}$ will be given by


$$
\begin{equation*}
d F=\varnothing \frac{\cos \phi}{r^{2}} d \sigma_{1} d \sigma_{2} \tag{1}
\end{equation*}
$$

Since

$$
\begin{equation*}
\cos \phi=\frac{h}{r}=\frac{h}{\sqrt{h^{2}+B^{2}}} \tag{2}
\end{equation*}
$$

and

$$
\begin{equation*}
d F=\bar{\Phi} \frac{h_{3}}{r} d \sigma_{1} d \sigma_{2} \tag{3}
\end{equation*}
$$

a. Point Source $\left(R_{1} \geq R_{2}\right)$

If there is only one element of area emitting $\left(\mathrm{dO}_{2}\right)$ then the total emitted $P l u x$ is $4 \pi \Phi d \sigma_{2}$ and the geometrical counting officiency will be given by

$$
\begin{equation*}
d a=\frac{d F}{4 \pi \Phi^{d \sigma_{2}}}=\frac{h}{r^{3}} d \sigma_{3} \tag{4}
\end{equation*}
$$

To integrate over $\sigma_{1}$ we set

$$
\begin{equation*}
d \sigma_{1}=\mathrm{Ed} d \boldsymbol{d} \tag{5}
\end{equation*}
$$

Then by symmetry wo can limit $0 \leqslant \alpha \leqslant \Pi$ and put in a factor 2 so that in this case

$$
\begin{equation*}
d G=\frac{h}{2 \pi^{\prime} \pi} d d \alpha \tag{6}
\end{equation*}
$$

Then integrate first over $\propto$.
For $R_{2}+R \leqslant R_{1}, 0 \leqslant * \leqslant \pi$;
and for $E_{2}+\mathrm{R}_{2} \geq \mathrm{R}_{1}$
the upper limit on of is determined by

$$
R_{1}^{2}=R^{2}+R_{2}^{2}-2 R_{2} R \cos \alpha,
$$


(7)
or

$$
\begin{equation*}
\psi=\cos ^{-1}\left[\frac{R^{2}+R_{2}^{2}-R_{1}^{2}}{2 R_{2} R}\right] \tag{8}
\end{equation*}
$$

$s$

$0=\frac{1}{2}\left[1-\frac{h}{\sqrt{\left(R_{1}-R_{2}\right)^{2}+h^{2}}}\right]+\frac{h}{2 \pi} \int_{R_{1}-R_{2}}^{R_{1}+R_{2}} \cos ^{-1}\left\{\frac{R^{2}+R_{2}^{2}-R_{1}^{2}}{2 R_{2} R}\right\} \frac{R d R}{r^{3}}$.
(10)
b. Point Source $\left(\mathbf{R}_{2} \geq \mathbb{R}_{1}\right)$

In this case of is determined
as in Eq. (8). The limits on Integraion are different so that here we have

$\theta=\frac{h}{2 i} \int_{R_{2}^{-R_{1}}}^{R_{1}+R_{1}} \cos -1\left\{\frac{R^{2}+B_{2}^{2}-R_{1}^{2}}{2 R_{2} R}\right\} \frac{R^{2} d R}{r^{3}}$

This agrees with the stated expression in Calvin, "Isotopic Carbon", App. IV, for $Z_{2}>B_{1}$.
b. Circular Source $\quad \mathbf{R}_{1} \geq \mathbf{R}_{2}$

In this case the total flux is $4 \pi \Phi\left(\pi R_{2}^{2}\right)$ so that
$d \sigma=\frac{d F}{4 \pi^{2} R_{2}^{2} \Phi}=\frac{h}{4 \pi^{2} R_{2}^{2}} \frac{d \sigma_{1} d \sigma_{2}}{r^{3}}$.
Let

$$
\begin{align*}
& d \sigma_{2}=r_{2} d r_{2} d \theta  \tag{13}\\
& d \sigma_{1}=R d R d \theta \tag{14}
\end{align*}
$$

We can integrate readily over $\theta$ and restrict $0 \leqslant \alpha \leqslant \pi$ so that we get
$d a\left(r_{2}, E, q\right)=\frac{h}{\pi R_{2}^{2}} \frac{r_{2} d r_{2} R d R d \in}{r^{3}}$.
Mow integrate over $\propto$

$$
\begin{equation*}
d G\left(r_{2}, E\right)=\left(\frac{h}{\pi R_{2}^{2}} \frac{R_{-} d R}{r^{3}}\right)\left(\alpha r_{2} d r_{2}\right) \tag{16}
\end{equation*}
$$

How either we shall have

$$
\alpha=\pi
$$

or

$$
\begin{equation*}
\phi=\cos ^{-1}\left\{\frac{\mathrm{R}^{2}+r_{2}^{2}-\mathrm{R}_{1}^{2}}{2 r_{2} R}\right\} \tag{17}
\end{equation*}
$$

depending on the elreumstances.

We now want to find out over what ranges the integration of $\mathbf{R}$ and $r_{2}$ must be carried for different values of $\mathcal{O}$. If we take a fixed $r_{2}$ we can draw as follows the ranges of integration for $R$.

$$
\text { Let } r_{2} \leqq R_{1} \text { then }
$$

for $r_{2}+R \leqslant R_{1} ; \quad \alpha=\mathbb{T}$ and $0 \leqslant R \leqslant R_{1}-r_{2}$

$$
\begin{array}{r}
r_{2}+R \geq R_{1} ; \quad \gamma=\cos ^{-1}\left\{\frac{R^{2}+r_{2}^{2}-R_{1}^{2}}{2 r_{2} R}\right\} ; \text { and } \\
R_{2}-r_{2} \leq R \leqslant R_{1}+r_{2} .
\end{array}
$$

while for $r_{2} \geq R_{1} ; \alpha=\cos ^{-2}\left\{\frac{R^{2}+r_{2}^{2}-R_{1}^{2}}{2 r_{2} R}\right\} ;$ and

$$
r_{2}-R_{1} \leqslant R \leqslant r_{2}+R_{1}
$$

These regions may be plotted in two dimensions as follows in the $\mathrm{R}, \mathrm{r}_{2}$ plane


If we change the order of integration to an integration over $\mathrm{r}_{2}$ first, we have various regions to consider.
b1 $\quad\left(R_{1}>R_{2}\right)$
Region

| $A_{1}$ | $0 \leqslant r_{2} \leqslant R_{2} ;$ | $0 \leqslant R \leqslant R_{1}-R_{2}$ |
| :--- | :--- | :--- |
| $A_{2}$ | $0 \leqslant r_{2} \leqslant R_{1}-R ; R-R_{2} \leqslant R \leqslant R_{1}$ |  |$\pi$

$B_{1}+B_{2}$
$R_{1}-R \leqslant r_{2} \leqslant R_{2} ; \quad R_{1}-R_{2} \leqslant R \leqslant R_{1}$
Eq. (17)
$B_{3}$
$R-R_{1} \leqslant r_{2} \leqslant R_{2} ; \quad R_{1} \leqslant R \leqslant R_{1}+R_{2}$
$b_{2} \quad\left(R_{1}<R_{2}\right)$
Region
Linits
$\alpha$

A
$0 \leqslant r_{2} \leqslant R_{1}-R \quad 0 \leqslant R \in R_{1}$
$\pi$
$B_{1}$
$B_{2}$
$R_{1}-R \leqslant r_{2} \leqslant R+R_{1} ; \quad 0 \leqslant R \leqslant R_{2}-R_{1}$
Eq. (17)
$R_{1}-R \leqslant r_{2} \leqslant R_{2} \quad ; \quad R_{2}-R_{1} \leqslant R \leqslant R_{1}$
Eq. (17)
$B_{3}$
$R-R_{1} \leqslant r_{2} \leqslant R_{2} \quad R_{1} \leqslant R \leqslant R_{1}+R_{2}$

Eq. (17)

Wow to evaluate the expressions for $G$ in the two cases we have to integrate expressions of the type

$$
\int d r_{2} d r_{2}
$$

where $\quad$ of $=\pi$ or $\gamma=\cos ^{-1}\left[\frac{R^{2}+r_{2}^{2}-R_{1}^{2}}{2 R r_{2}}\right]$
for $\quad \int_{a}^{b} \sigma r_{2} d r_{2}=\left.\frac{\pi}{2} r_{2}{ }^{2}\right|_{a} ^{b}$

For the other case

$$
\left.\begin{array}{rl}
\int_{\mathrm{a}}^{\mathrm{b}}\left\{\begin{array}{c}
\left.\cos ^{-1}\left[\frac{\mathrm{R}^{2}-\mathrm{R}_{1}^{2}+r_{2}^{2}}{2 R r_{2}}\right] r_{2}\right\} d r_{2} \\
\end{array}\right. & =\left\{\begin{array}{c}
\frac{r_{2}^{2}}{2} \cos ^{-1}\left[\frac{R^{2}-R_{1}^{2}+r_{2}^{2}}{2 R r_{2}}\right]-\frac{R_{1}^{2}}{2} \sin ^{-1} \frac{R^{2}+R_{1}^{2}-r_{2}^{2}}{2 R R_{1}}
\end{array}\right\} \\
-\frac{1}{4} \sqrt{4 r_{2}^{2} R^{2}-\left(R^{2}-R_{1}^{2}+r_{2}^{2}\right)^{2}}
\end{array}\right\}
$$

Evaluating this $\int$ at the various limits needed, we find that

$$
\begin{align*}
& \int^{R_{2}} \cos ^{-1}\left[\frac{R^{2}-R_{1}^{2}+r_{2}^{2}}{2 R r_{2}}\right] r_{2} d r_{2} \\
&= \frac{R_{2}^{2}}{2} \cos ^{-1}\left[\frac{R^{2}-R_{1}^{2}+R_{2}^{2}}{2 R R_{2}}\right]-\frac{R_{1}^{2}}{2} \sin ^{-1}\left[\frac{R^{2}+R_{1}^{2}-R_{2}^{2}}{2 R R_{1}}\right] \\
&-\frac{1}{4} \sqrt{4 R_{2}^{2} R-\left(R^{2}-R_{1}^{2}+R_{2}^{2}\right)^{2}}+c \\
&= \frac{\pi R_{2}^{2}}{4}-\frac{\pi R_{1}^{2}}{4}+r\left(R_{2}\right)+c \tag{20}
\end{align*}
$$

where
$r\left(R_{2}\right)=\frac{R_{1}^{2}}{2} \cos ^{-1}\left[\frac{R^{2}+R_{2}^{2}-R_{2}^{2}}{2 R_{1}}\right]-\frac{R_{2}^{2}}{2} \sin ^{-1}\left[\frac{R^{2}+R_{2}^{2}-R_{1}^{2}}{2 R R_{2}}\right]$

$$
\begin{equation*}
-\frac{1}{4} \sqrt{4 R^{2} R_{1}^{2}-\left(R^{2}+R_{1}^{2}-R_{2}^{2}\right)^{2}} \tag{21}
\end{equation*}
$$

$$
\begin{equation*}
\int^{R_{1}-R} \cos ^{-1}\left[\frac{R^{2}-R_{1}^{2}+r_{2}^{2}}{2 R r_{2}}\right] r_{2} d r_{2}=\frac{\pi}{2}\left(B_{1}-R\right)^{2}-\frac{\pi}{4} R_{1}^{2}+c \tag{204}
\end{equation*}
$$

$$
\begin{align*}
& \int^{R_{1}} \cos ^{-2}\left[\frac{R^{2}-R_{1}^{2}+r_{2}^{2}}{2 R r_{2}}\right] r_{2} d r_{2}=-\frac{\pi}{4} R_{1}^{2}+c \\
& \int_{(20 b)}^{R_{1}+R} \cos ^{-1}\left[\frac{R^{2}-R_{1}^{2}+r_{2}^{2}}{2 R \cdot 2}\right] r_{2} d r_{2}=\frac{\pi R_{1}^{2}}{4}+c \quad(20 \mathrm{c}) \\
& \int_{\left(R_{2}\right.}^{R_{2}} \cos ^{-1}\left[\frac{R^{2}-R_{1}^{2}+r_{2}^{2}}{2 R r_{2}}\right] r_{2} d r_{2}=\frac{T_{R_{2}}^{2}}{4}-\frac{\pi R_{1}^{2}}{4}+r\left(R_{2}\right)+c . \tag{20d}
\end{align*}
$$

We now substitute the limits in the expression for $G$.
Case $b_{1}$. $\quad\left(R_{1}>R_{2}\right)$

$$
\sigma=\frac{h}{\pi R_{2}^{2}}\left\{\begin{array}{l}
\int_{0}^{R_{1}-R_{2}} \frac{R d R}{r^{3}}\left[\frac{\pi}{2} R_{2}^{2}\right] \int_{R_{1}-R_{2}}^{R_{1}} \frac{R d R}{r^{3}} \frac{\pi}{2}\left(R_{1}-R\right)^{2} \\
+\int_{Z_{1}-R_{2}}^{R_{1}} \frac{R d R}{r^{3}}\left[\frac{\pi R_{2}^{2}}{4}-\frac{\pi R_{1}^{2}}{4}+r\left(R_{2}\right)-\frac{\pi}{2}\left(R_{1}-R\right)^{2}+\frac{\pi}{4} R_{1}^{2}\right] \\
+\int_{R_{1}}^{2} \frac{R_{2}}{r^{2}}\left[\frac{R d R}{4}\left[\frac{\pi R_{2}^{2}}{4}-\frac{\pi^{2} R_{1}^{2}}{4}+r\left(R_{2}\right)+\frac{\pi}{4} R_{1}^{2}\right]\right.
\end{array}\right.
$$

$$
\begin{align*}
& \text { - } 11 \text { - } \\
& =\frac{h}{2} \int_{0}^{R_{1}-R_{2}} \frac{R d R}{r^{3}}+\frac{h}{4} \int_{R_{1}-R_{2}}^{R_{1}+R_{2}} \frac{R d R}{r^{3}}+\frac{h}{R_{2}^{2}} \int_{R_{1}-R_{2}}^{R_{1}+R_{2}} r\left(R_{2}\right) \frac{R d R}{r^{3}} \\
& =\frac{1}{2}\left[1-\frac{h}{\sqrt{h^{2}+\left(R_{1}-R_{2}\right)^{2}}}\right]+\frac{1}{4}\left[\frac{h}{\sqrt{h^{2}+\left(R_{1}-R_{2}\right)^{2}}}-\frac{h}{\sqrt{h^{2}+\left(R_{1}+R_{2}\right)^{2}}}\right] \\
& +\frac{h}{/ / R_{2}{ }^{2}} \int_{R_{1}-R_{2}}^{R_{1}+R_{2}} f\left(R_{2}\right) \frac{R d R}{r^{3}} \\
& =\frac{1}{2}\left[1-\frac{1}{2} \frac{h}{\sqrt{h^{2}+\left(R_{1}-R_{2}\right)^{2}}}+\frac{h}{\sqrt{h^{2}+\left(R_{1}-R_{2}\right)^{2}}}\right] \\
& +\frac{h}{\pi R_{2}^{2}} \cdot \int_{R_{1}-R_{2}}^{R_{1}+R_{2}} f\left(R_{2}\right) \frac{R d R}{r^{3}} \tag{22}
\end{align*}
$$

which is formula (4) in Calvin, "Isotopic Carbon", App. Iv.

Case $b_{2}\left(R_{1}<R_{2}\right):$

$$
\begin{aligned}
G=\frac{h}{\pi R_{2}^{2}} & \left\{\int_{0}^{R_{1}} \frac{R d R}{r^{3}}\left[\frac{\pi\left(R_{1}-R\right)^{2}}{2}\right]+\int_{0}^{R_{2}} \frac{R R_{1}}{r^{3}}\left[\frac{\pi R_{1}{ }^{2}}{4}-\frac{\left.\pi\left(R_{1}-R\right)^{2}+\frac{\pi}{4} R_{1}{ }^{2}\right]}{}\right.\right.
\end{aligned}\left\{\begin{array}{l}
\int_{R_{2}-R_{1}}^{R_{1}} \frac{R d R}{r^{2}}\left[\frac{\pi R_{2}^{2}}{4}-\frac{\pi R_{1}^{2}}{4}+f\left(R_{2}\right)-\frac{\pi\left(R_{1}-R\right)^{2}}{2}+\frac{\pi R_{1}^{2}}{4}\right] \\
+\int_{R_{1}}^{R_{1}+R_{2}} \frac{R d R}{r^{3}}\left[\frac{\pi R_{2}^{2}}{4}-\frac{\pi R_{1}^{2}}{4}+f\left(R_{2}\right)+\frac{\pi R_{1}^{2}}{4}\right]
\end{array}\right\}
$$

$$
=\frac{h}{\pi R_{2}^{2}}\left\{\frac{\pi R_{1}^{2}}{2} \int_{0}^{R_{2}-R_{1}} \frac{R d R}{r^{3}}+\frac{\pi R_{2}^{2}}{4} \int_{R_{2}-R_{1}}^{R_{1}+R_{2}} \frac{R d R}{r^{3}}+\int_{R_{2}-R}^{R_{1}+R_{2}} f\left(R_{2}\right) \frac{R d R}{r^{3}}\right\}
$$

$$
=\frac{\frac{1}{2}}{2}\left\{\frac{R_{1}^{2}}{R_{2}^{2}}\left[1-\frac{h}{\sqrt{h^{2}+\left(R_{2}-R_{1}\right)^{2}}}\right]+\frac{1}{2}\left[\frac{h}{\sqrt{h^{2}+\left(R_{2}-R_{1}\right)^{2}}}-\frac{h}{\sqrt{h^{2}+\left(R_{1}+R_{2}\right)^{2}}}\right]\right\}
$$

$$
+\frac{h}{\pi R_{2}^{2}} \int_{R_{2}-R_{1}}^{R_{2}+R_{1}} f\left(R_{2}\right) \frac{R d R}{r^{3}}
$$

(23)
which is formula (5) in Calvin.


