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Surface Radiation Properties of ThO₂
and Other Reactor Materials

by

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and
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ABSTRACT

This technical report summarizes the results obtained on surface radiation properties of ThO_2 and other related reactor materials. The content herein is based on the M.S. dissertation submitted by H.H. Ma, Department of Mechanical Engineering, University of Wisconsin at Milwaukee, December 1979. It is a compliment to the previous report - CD0-4040-3, 1978, which summarizes the results of UO_2 and other pertaining materials.

Specifically, the present study provides surface radiation properties of three nuclear reactor materials; thorium dioxide, sodium and steel (99 Fe + 1 C). First, based on transmission spectra and Kramers-Kronig analysis, the complex refractive index, which is fundamental to surface radiation property calculations, of ThO_2 is presented over a wide spectral range. Then depending on the contacting media, three types of the surfaces are considered. They are the interfaces between ThO_2 and Sodium, ThO_2 and Steel (99 Fe + 1 C) as well as ThO_2 and a gas. For each interface, the spectral interface reflectance from one side of the medium to the other side is evaluated for all wavelengths. These spectral properties are further integrated to yield the total hemispherical properties over a temperature range, 1200°K to 6000°K, of interest to reactor safety analysis.

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NOMENCLATURE

- I - intensity of radiative flux
 k - extinction coefficient
 $K_{a\lambda}$ - absorption coefficient
 n - real index of refraction
 \hat{n} - complex index of refraction ($= n - ik$)
 T - temperature

Greek:

- α - total absorptance
 α_λ - spectral absorptance
 ε - total emittance
 ε_λ - spectral emittance
 ϕ - phase angle
 θ - incidence angle
 λ - wavelength
 ρ - total reflectance
 ρ_λ - spectral reflectance

Subscripts:

- 1 - medium 1
2 - medium 2
i - incident quantity
n - normal direction
s - source
t - transmitting quantity
 λ - wavelength

I. INTRODUCTION

In the safety analysis of nuclear reactors, the temperature of the reactor materials involved often is high enough (1200 to 6000°K) that thermal radiative transfer may become an important mode of heat transfer as compared to conduction and convection.¹⁻⁴ To perform radiative transfer calculations, the information on thermal radiation properties of reactor materials is needed. Chan and Tseng⁵ recently presented the surface radiation properties of uranium dioxide when in contact with various reactor media. The present study aims to extend their work to include thorium dioxide. Thorium dioxide has renewed interest as a potential reactor fuel, particularly, in view of recent concern about nuclear proliferation.

To use the well known Fresnel relations to calculate surface radiation properties, the values of the complex refractive index ($\hat{n} = n - ik$) of a substance must be provided. The real part of the index, n , is generally called as the index of refraction and the imaginary part, k , as the extinction coefficient. Since the optical constants of ThO_2 are only known at a few selected wavelengths, their values, as one of major purposes of the present study, will be made available over a wide range of the wavelength spectrum. The method to be employed to calculate these optical constants is Kramers-Kronig (K-K) theory, which is a powerful but relatively unknown to heat transfer engineers. A brief review of the theory seems to be in order.

The theory of Kramers-Kronig analysis of reflectance spectra measured with normal incidence on the plane sample was originally described by Robinson,⁶ in 1952. Then, in 1967, Berreman⁷ extended the applicability of the K-K analysis by presenting an algorithm for computing the phase shifts and n , k from a reflectance spectrum obtained from a radiant flux. The flux was linearly polarized either parallel or perpendicular to the plane of incidence and was specularly reflected at an oblique angle.

In a later work, Didrikil⁸ developed a method of K-K analysis for determining the optical constants from the reflectance spectra at an oblique angle of incidence with complete allowance for the degree of polarization, especially zero polarization. For any non-conductor, increasing the angle of incidence continuously increases the component of spectral reflectance. With nonpolarized light a considerable increase in the reflectance takes place for angles of incidence greater than the Brewster angle. In principle, increasing the incidence angle makes it possible to improve the accuracy of measurement of the reflectance in the case where the reflectance at normal incidence is low.

K-K analysis has the advantage of eliminating the need for a polarizer, and increasing the intensity of the reflected light. K-K analysis can be used for any wavelength range, and is suitable to any non-magnetic material. Thus, it simplifies the experimental equipment required and furnishes a better design for the reflection experiment.

This study will discuss two methods that K-K theory can be applied to determine the optical constants from spectral measurement: one from reflection spectra and the other from transmission spectra. After both methods are tested by the classical Lorentz dispersion model, the optical constants of ThO_2 will be determined, followed by evaluating the radiation properties of three types of interfaces:

- a. ThO_2 in contact with sodium.
- b. ThO_2 in contact with steel.
- c. ThO_2 in contact with water vapor.

For each interface, the spectral hemispherical reflectivity will be determined for all wavelengths of interest and the total hemispherical reflectivity from a black body source over a wide temperature range (1200 - 6000°K).

II. ELECTROMAGNETIC THEORY

To begin with, a brief review of reflection laws, primarily Fresnel's relation, from the electromagnetic theory will be made. They will be employed to calculate radiation properties in latter sections.

Fresnel's equations were derived by Fresnel in 1823. On the basis of the elastic theory of light, they give the ratios of the amplitudes of the incident, reflected, and transmitted waves.

Consider two semi-infinite media in direct contact as shown in Fig. 1. They can be either metals or dielectrics, with complex refractive indices $\hat{n}_{1\lambda}$ and $\hat{n}_{2\lambda}$ respectively. The interface between two media is assumed to be specular and optical smooth. The complex index of refraction is shown by $\hat{n}_\lambda = n_\lambda - ik_\lambda$. The subscript λ denotes that \hat{n}_λ , n_λ and k_λ are, in general, a function of radiation wavelength. According to Fresnel equations, the monochromatic directional reflectance of radiation from medium 1 back to the same medium is

$$P_{\lambda 11}(\theta) = \frac{1}{2} \{ P_{\lambda 11S}(\theta) + P_{\lambda 11P}(\theta) \} \quad (\text{II.1})$$

where

$$P_{\lambda 11S}(\theta_r) = \left| \frac{\hat{n}_{1\lambda} \cos \theta_i - \hat{n}_{2\lambda} \cos \theta_t}{\hat{n}_{1\lambda} \cos \theta_i + \hat{n}_{2\lambda} \cos \theta_t} \right|^2 \quad (\text{II.2})$$

$$P_{\lambda 11P}(\theta_r) = \left| \frac{\hat{n}_{2\lambda} \cos \theta_i - \hat{n}_{1\lambda} \cos \theta_t}{\hat{n}_{2\lambda} \cos \theta_i + \hat{n}_{1\lambda} \cos \theta_t} \right|^2 \quad (\text{II.3})$$

and $\theta_r = \theta_i$. θ_t can be calculated from Snell's law

$$\hat{n}_{2\lambda} \sin \theta_t = \hat{n}_{1\lambda} \sin \theta_i \quad (\text{II.4})$$

In Eq. (II.1), a circularly polarized light has been assumed, i.e. polarization of the light is assumed to be zero. In case the optical constants are known for all the wavelengths, the spectral directional reflectance over all wavelengths can be found. Then, the monochromatic hemispherical reflectance $\rho_{\lambda 11}$ can be determined by integration over the hemisphere,

$$\rho_{\lambda 11} = 2 \int_0^{\frac{\pi}{2}} \rho_{\lambda 11}(\theta) \sin\theta \cos\theta \, d\theta \quad (\text{II.5})$$

and the total hemispherical reflectance can be calculated by integrating over the wavelength,

$$\rho_{11}(T_s) = \frac{\int_0^{\infty} \rho_{\lambda 11} H_{\lambda}(T_s) \, d\lambda}{\int_0^{\infty} H_{\lambda}(T_s) \, d\lambda} \quad (\text{II.6})$$

In Eq. (II.6) H_{λ} is the incident radiation from a source at temperature T_s . Although the monochromatic reflectance ρ_{λ} is purely a function of the optical properties, the total reflectance ρ is not, as it depends also on the source temperature T_s . In the absence of a better radiating source, we take it as a black body source at $T_s = T_1$. Thus, the Planck's function serves as a weighting factor. Since it shifts to lower wavelengths as temperature T_s increases, ρ_{λ} 's at lower wavelengths are weighted more than those at higher wavelengths. Generally speaking, the temperature T_s has a significant effect on the reflectance. Care should be taken not to use the temperature of medium 2 on the other side of the interface. Strictly speaking, the optical constants are temperature dependent. However, due to the lack of data

at elevated temperature, such temperature dependence could not be accounted for at the present time.

The monochromatic directional transmittance from medium 1 to 2 is given by

$$t_{\lambda 12}(\theta) = 1 - \rho_{\lambda 11}(\theta) \quad (\text{II.7})$$

which is the same as the monochromatic directional absorptance of an incident radiation from medium 1 to medium 2.

$$\alpha_{\lambda 12}(\theta) = t_{\lambda 12}(\theta) \quad (\text{II.8})$$

In accordance with Kirchnoff's law, the directional monochromatic emittance from medium 2 to 1 is

$$\epsilon_{\lambda 21}(\theta) = \alpha_{\lambda 12}(\theta) \quad (\text{II.9})$$

if $T_1 = T_2$. Therefore, under the condition of a black body source and the condition that $T_2 = T_1$, the total hemispherical properties are inter-related by

$$\rho_{11} = 1 - \alpha_{12} = 1 - \epsilon_{21} = 1 - t_{12} \quad (\text{II.10})$$

Furthermore, the theory can be simplified under these two special cases. They are cited below.

First, if the medium 1 is a dielectric medium such as air or any other gases or vapors, then $k_1 \approx 0$, $n_1 \approx 1$, and $\hat{n}_1 \approx 1$. Upon the substitution of $\hat{n}_1 \approx 1$ into Eq. (II.2) and (II.3) yields expressions quoted by most textbooks.⁹⁻¹⁰ This is of particular interest to engineers because the measurements of all the surface properties are conducted by placing the specimen in air, vacuum, or some gaseous media.

The other one is the normal reflectance. When the incident radiation is normal to the interface,

$$\theta_i = \theta_r = \theta_t = 0$$

Then, from Eqs. (II.1) to (II.4),

$$\rho_{\lambda 11n} = \rho_{\lambda 11p} = \rho_{\lambda 11s} = \left| \frac{\hat{n}_{2\lambda} - \hat{n}_{1\lambda}}{\hat{n}_{2\lambda} + \hat{n}_{1\lambda}} \right|^2 \quad (\text{II.11})$$

which is indeed simple and useful.

In summary, the electromagnetic theory provides a prediction of the monochromatic specular reflectance, from which the spectral specular transmittance and emittance can be inferred, if necessary. Then, spectral and total hemispherical properties can be evaluated.

III. KRAMERS-KRONIG THEORY FOR THE DETERMINATION OF OPTICAL CONSTANTS

In accordance with K-K theory, there are two methods to extract the values of the fundamental spectroscopic parameters, namely, the optical constants n and k , from spectroscopic measurements. The first method is based on a measured reflection spectrum and the second one based on a measured transmission spectrum. The latter is applicable to gaseous media or semi-transparent solids such as ThO_2 or other dielectric materials while the former to solids only, but irrespectively whether the solid is semi-transparent or not, i.e. applicable to dielectrics as well as metals. Even though the second method will not be used in the present study, both methods for the sake of completeness will be described in detail and tested using the classical Lorentz oscillator model. Then, the first method will be employed to calculate the optical constants of ThO_2 .

III.1 Optical Constants from Reflection Measurement

III.1.1 Working Formula

According to Fresnel relations, the optical constants n and k can be determined from the spectrum of a perpendicular component R_s by means of the following equations:⁸

$$n = \left\{ \frac{((Q^2 - P^2 + \sin^2 \theta) - \{(Q^2 - P^2 + \sin^2 \theta)^2 + 4Q^2 P^2\}^{\frac{1}{2}})}{2} \right\}^{\frac{1}{2}} \quad (\text{III.1})$$

$$k = PQ/n \quad (\text{III.2})$$

where

$$P = \frac{2\sqrt{R_s} \sin\phi \cos\theta}{1 + R_s - 2\sqrt{R_s} \cos\phi} \quad (\text{III.3})$$

$$Q = \frac{(1 - R_s) \cos\theta}{1 + R_s - 2\sqrt{R_s} \cos\phi} \quad (\text{III.4})$$

and θ is the angle of incidence. Alternatively, P and Q can be expressed in terms of θ , n and k as follows

$$P = nk/Q \quad (\text{III.5})$$

$$Q = \left\{ \frac{(n^2 - k^2 - \sin^2\theta) + \left[(n^2 - k^2 - \sin^2\theta)^2 + 4n^2k^2 \right]^{\frac{1}{2}}}{2} \right\}^{\frac{1}{2}} \quad (\text{III.6})$$

In the above equations, the phase angle ϕ can be determined from the K-K relationship¹¹⁻¹³

$$\phi(a) = -\frac{2a}{\pi} \int_0^{\infty} \frac{\ln \left| \sqrt{R_s(\omega)} \right|}{\omega^2 - a^2} d\omega \quad (\text{III.7})$$

where ω is an angular frequency, a is the resonance frequency. It relates the phase angle ϕ to the perpendicular component of the specular reflectance R_s , and was obtained from a K-K analysis of an absolute, specular reflectance spectrum of a material. A brief derivation of the phase-shift dispersion relation is given in Appendix B.

The integral relation Eq. (III.7) shows that if the perpendicularly polarized reflectivity R_s is measured over the whole frequency spectrum, the phase angle and the indices n and k can be calculated for any arbitrary frequency at a given incident angle. Performing integration by parts, Eq. (III.7) becomes

$$\phi(a) = \pi^{-1} \int_0^{\infty} \ln \left| \frac{\omega-a}{\omega+a} \right| \frac{d \operatorname{In} R_s^{\frac{1}{2}}(\omega)}{d \omega} d \omega \quad (\text{III.8})$$

Based on the assumption that the slope of R_s be zero at $\omega=0$ and $\omega=\infty$, we can rewrite Eq. (III.8) into a more convenient form (see Appendix C) as,

$$\phi(a) = \pi^{-1} \sum_{k=2}^{n-1} p_k \left\{ (a-\omega_k) \operatorname{In} |a-\omega_k| + (a+\omega_k) \operatorname{In} |a+\omega_k| \right\} \quad (\text{III.9})$$

where

$$p_k = u_{k+1} - u_k$$

and

$$u_k = \frac{\operatorname{In} R_s^{\frac{1}{2}}(\omega_k) - \operatorname{In} R_s^{\frac{1}{2}}(\omega_{k-1})}{\omega_k - \omega_{k-1}}$$

The above equation can be easily programmed for numerical computation, a copy of FORTRAN program is given in Appendix E.

By examining Eq. (III.8), we find the values of $R_s(\omega)$ in the spectral regions where $\omega \gg a$ or $\omega \ll a$, do not contribute to the value of $\phi(a)$, because then the weighting function $\operatorname{In} \left\{ \left| \frac{\omega-a}{\omega+a} \right| \right\}$ approaches zero. Thus, we can neglect the variations of the reflectance spectrum, even if the variations are reasonably large, in those region. Consequently, we can assume the reflectance flatten out on both ends of spectrum. Moreover, in view that the spectral region in which the reflectance is constant do not contribute to the value of $\phi(a)$, only a finite spectral region of the reflectance spectrum needed to be measured, which reduces experimental efforts considerably.

The reflectance measured in the experiment is not R_s , but is R_o .

given by

$$R_o = \frac{1}{2} \left(R_s (1+\alpha) + R_p (1-\alpha) \right) \quad (\text{III.10})$$

where α is the degree of polarization

$$\alpha = \frac{I_s - I_p}{I_s + I_p}$$

I_s and I_p are the reflected intensities of light in the planes perpendicular and parallel to the plane of incidence respectively.

We can rewrite Eq. (III.10) in the form

$$R_s = \frac{2R_o}{(1+\alpha) + (R_p/R_s)(1-\alpha)} \quad (\text{III.10.a})$$

Because we are interested in using a nonpolarized light, $\alpha=0$. The ratio

R_p/R_s in turn is defined in the terms of P , Q and θ as

$$\frac{R_p}{R_s} = \frac{(Q - \sin\theta \tan\theta)^2 + P^2}{(Q + \sin\theta \tan\theta)^2 + P^2} \quad (\text{III.11})$$

Also,

$$R_s = \frac{(Q - \cos\theta)^2 + P^2}{(Q + \cos\theta)^2 + P^2} \quad (\text{III.11.a})$$

III.1.2 Direct Method

If we can determine the ratio R_p/R_s and knowing the value of polarization α through the experiment, then it is possible to calculate the R_s spectrum directly from the measured reflectance R_o using Eq. (III.10.a). The optical constants can then be calculated from Eqs. (III.1) and (III.2) with P and Q from Eqs. (III.3) and (III.4). However, it is not simple to determine the ratio R_p/R_s because R_s is unknown.

In order to determine the ratio R_p/R_s , experimental setup is somewhat complicated and expensive as the source light has to be polarized. Therefore, this method is avoided whenever possible. Instead the following method is more favorable.

III.1.3 Iterative Method

For the determination of the optical constants from the measured R_o spectrum, we use the method of successive approximations suggested by Ref. 8. We can substitute R_o in place of R_s in Eqs. (III.3), (III.4) and (III.7). The values of $\phi(a)$ obtained are used to calculate the zeroth approximation of P and Q from Eqs. (III.3) and (III.4), and R_p/R_s from Eq. (III.11). Based on the R_p/R_s value, the first approximation of R_s^* is obtained from Eq. (III.10.a) which is then used to redetermine ϕ , P, Q, R_p/R_s and the next approximation of R_s^* until convergence is reached. The final set of R_s^* is the desired R_s . Therefore, this process of successive approximations makes it possible to obtain the R_s spectrum from R_o spectrum with the required degree of accuracy and certain imposed conditions to be discussed later.

In certain wavelength range, sustained oscillations of the values of R_s^* may be produced, so the suggestions made by Ref. 8 can be followed. Since the difference between $R_s^* - R_s$ in the i th and $(i-1)$ th iteration have different signs, we introduce the additional condition

$$R_s^*(i) = \frac{1}{2}[R_s^*(i) + R_s^*(i-1)] \quad (\text{III.12})$$

to expedite the convergence. The above equation means that, in the process of iteration, instead of using the newly obtained R_s^* value from $(i-1)$ th iteration for i th iteration, a weighted average with the previous value should be used. Equation (III.12) is used only for the calculation of

those R_s^* which the corresponding R_o exceed a certain value B . The number of iterations for convergence depends on the choice of B , and increases with an increase in the incidence angle.

For the reflectance with a high maximum of R_o , R_{max} , and a relatively low edge, the value of B can be taken as $R_{max}/2$. For spectra with low R_{max} and a high long wave edge, the lower limit of B can be taken as the value of C , where $C = (R_{ok} + R_{max})/2$. R_{ok} is the value of the reflectivity of long wave edge at a sufficient distance from the reflection band. In the intermediate case, B has a value lying between $R_{max}/2$ and C . In practice, it is better not to use Eq. (III.12) in the first iteration. Instead it is applied to all points of the spectrum after 4-6 iterations. Or, simply setting $B=C$ in the first 4-6 cycles, then apply Eq. (III.12) to all the points of R_s^* spectrum.⁸

As to the selection of the angle of incidence θ , we use the value of the normal reflectivity at the maximum of the reflectance band, R_n , as a criterion. For high values of R_n , on the order of 0.9, the use of an oblique incidence angle may prove useful, though it is necessary to restrict the angle of θ not to exceed 60° . For R_n less than the order of 0.1, the angle θ may reach 80° . In the intermediate cases, it is convenient to select θ around 70° .

To test the validity of the iterative method for the determination of R_s and therefore the optical constants from Eqs. (III.1) to (III.5) from R_o spectrum, the classical Lorentz oscillator model is used. Details about the oscillator can be found in Appendix A.

Let us take the oscillator equations which best fit the reflectivity data of the material α -II-SiC for example. The four correlation parameters in the equations for this material have been found as¹⁴

$$\rho = 0.263$$

$$\tau = 0.006$$

$$a = 12.60$$

$$\epsilon_0 = 6.7$$

With these parameters, n and k of α -II-SiC can be calculated from Eqs. (A.17) to (A.21) in Appendix A and the results are listed in the last two columns of Table 1. Also can be calculated are P and Q from Eqs. (III.5) and (III.6), respectively. Finally by choosing a non-polarized light, $\alpha=0$, and an incidence angle of $\theta=60^\circ$ or $\theta=70^\circ$, R_o and R_s reflection spectra can be calculated through the use of Eqs. (III.11) and (III.11.a). They are displayed in Fig. 2 and listed in Table 1 as well.

Using the R_o spectrum and the iterative method as just described, the values of R_s^* , n^* and k^* at the end of the iteration are shown in Table 2 and Figs. 3 and 4. By comparing those in Table 2 with R_s , n and k values in Table 1 or by examining Figs. 3 and 4, we see that, for the case of $\theta=70^\circ$, they are in good agreement everywhere except near the peak region between 11 and 13. The sustained oscillations of R_s^* can not be avoided even if the additional condition, Eq. (III.12), is used. One possible reason is that the incidence angle is too high. If we check the reflectance spectra, there are many reflectance exceed the value of 0.95. In this instance, the incidence angle should be restricted to 60° or less. As we examine the case with $\theta=60^\circ$, we find the results in figs. 3

and 4 not only converge everywhere but also in good agreement through the whole spectrum. Thus the iterative K-K analysis can be successfully applied in determining the optical constants directly from the experimental reflection spectra.

As stated in Ref. 8, it should be mentioned in passing that the magnitude of error depends on the ratio of R_p/R_s . If both polarization α and R_p/R_s are zero, the magnitude of error reaches a maximum. When a nonpolarized light is used in an experiment, α is equal to zero, then the ratio R_p/R_s becomes critical to the accuracy of the iterative method using the K-K analysis.

III.2 Optical Constants from Transmission Measurement

In the thin-film theory, the absorption coefficient, K_a , which describes the fractional decrease in intensity I with distance s , is defined as

$$K_{a\lambda} = -\frac{1}{I} \frac{dI}{ds} \quad (\text{III.13})$$

Therefore,

$$K_{a\lambda} = \frac{\ln(I_0/I)}{s} = \frac{2.303 \times \text{O.D.}}{s} \quad (\text{III.14})$$

where: I_0 - is the intensity of the light beam before shining onto the sample.

s - sample thickness.

O.D. - optical density.

From the above, $K_{a\lambda}$ can be evaluated from transmission or absorption spectra which describe the optical density as a function of λ . Then one of the optical constants, namely, the extinction coefficient k ,

can be evaluated from $K_{a\lambda}$ by the relation

$$k = K_{a\lambda}\lambda/4\pi \quad (\text{III.15})$$

To evaluate the other optical constant, the K-K dispersion relations are introduced. They are integral relations between the real and imaginary parts of a function $f(\omega)$ which is frequency dependent, and analytic in the upper half of the complex frequency plane.

For a given frequency a , these relations can be written:¹⁵

$$\text{Re } f(a) = \frac{2}{\pi} \text{p.v.} \int_0^{\infty} \frac{\omega \text{Im } f(\omega)}{\omega^2 - a^2} d\omega \quad (\text{III.16})$$

$$\text{Im } f(a) = -\frac{2}{\pi} \text{p.v.} \int_0^{\infty} \frac{\text{Re } f(\omega)}{\omega^2 - a^2} d\omega \quad (\text{III.17})$$

where: p.v. signifies principle value of the integral for $\omega=a$.

These relations are valid for any physical quantity which expresses a linear relation between the response of system and the incident excitation, under the condition that the response of a system could never precede the incident excitation. These relations are useful if the complex function of $f(\omega)$ is replaced by the complex index, $\hat{n}(\omega) = n(\omega) - iK(\omega)$. We have,

$$n(a) = \frac{2}{\pi} \text{p.v.} \int_0^{\infty} \frac{\omega k(\omega)}{\omega^2 - a^2} d\omega \quad (\text{III.18})$$

$$k(a) = -\frac{2}{\pi} \text{p.v.} \int_0^{\infty} \frac{n(\omega)}{\omega^2 - a^2} d\omega \quad (\text{III.19})$$

If either $n(\omega)$ or $k(\omega)$ function is known, the other function can be determined directly through one of the integrals. Consequently, we don't really have to measure the reflectances for certain crystals if the transmission spectra are known, because from the transmission spectra the absorption coefficient $K_{a\lambda}$ can be calculated from Eq. (III.14) as a function of frequency. Therefore, the extinction coefficients

can be calculated from Eq. (III.15), and $n(\omega)$ from Eq. (III.18).

Similar to Eq. (III.7), Eqs. (III.16) and (III.17) can be expressed in the following conveniently numerical forms,¹⁵

$$n(a) - n(\infty) = \frac{1}{\pi} \sum_{k=2}^{n-1} q_k \left[(a-\omega_k) \ln|a-\omega_k| - (a+\omega_k) \ln(a+\omega_k) \right] \quad (\text{III.20})$$

$$k(a) = -\frac{1}{\pi} \sum_{k=2}^{n-1} p_k \left[(a-\omega_k) \ln|a-\omega_k| + (a+\omega_k) \ln(a+\omega_k) \right] \quad (\text{III.21})$$

where: ω_k - the frequency of a pole of a partial fraction.

q_k - the variation of the slope of $k(\omega)$ at these points.

p_k - the variation of the slope of $n(\omega)$ at these points. (see Appendix C)

It should be mentioned that any error in the determination of a measured function, even if it occurs only at a single frequency, will affect the whole assembly of the calculated values of the other function. In particular, the inherent assumption of zero slope at both integration limits in Eqs. (III.20) and (III.21) should be carefully observed by setting, for example, $k_1=k_2$ and $k_{n-1}=k_n$ in Eq. (III.20).

Noting that in the calculation of surface radiation properties, the refractive index n is dominant. Therefore, it is better off not to assume n to be constant. Instead a more realistic refractive index function, $n(\lambda)$, which varies with frequency should be used wherever possible. Thus Eq. (III.20) is of particular importance since it can be used to generate n values from k values.

To test Eq. (III.20), the values of the extinction coefficient k of Lorentz model listed in Table 1 are used in the equation to see if the reproduced index of refraction n^* agrees with the original n listed in Table 1. The results are compared in Table 3 and displayed in Fig. 5. It shows that the agreement is remarkably good.

IV. COMPUTATION & RESULT

Being successfully verified in the previous section, the K-K theory can now be employed to calculate the optical constants of thorium dioxide. It will then be followed by the calculation of surface radiation properties of thorium dioxide when it is in direct contact with sodium, steel or a gas.

IV.1 Optical Constants of Thorium Dioxide

Theoretically speaking, the optical constants of ThO_2 can be calculated by the equations from the Lorentz theory of dielectric media. However, in practical applications, we find the equations tend to fail for non-monatomic solids as it is difficult to fit data through the equations even if the criteria described in Appendix A.1 are followed. The calculated k values are found to deviate greatly from the experimental absorption spectrum at different temperatures.

a. Extinction Coefficient k :

According to Ref. 16, Bates reported the absorption spectra of a single crystal ThO_2 at room temperature from the wavelength range 2 to 13μ with 0.048 cm sample thickness. Also, he measured sample at high temperature with different annealed gases from 0.2 to 5μ . The extinction coefficient k , therefore, can be determined from the optical density by Eq. (III.14). The absorption spectra are reproduced in three figures in Appendix D.

b. Index of Refraction n :

Once the extinction coefficient is known, the K-K dispersion theory is applied to find the real part of the complex refractive index, namely, the index of refraction n . The difference between $n(\infty)$ and $n(\omega)$ is

first calculated from Eq. (III.20). In order to determine the values of $n(\omega)$, we have to determine the $n(\infty)$ next. As far as $n(\infty)$ is concerned, it is the square root of ϵ_0 (see Eq. (A.13) in Appendix A). Because the value of ϵ_0 for ThO_2 is not available, we have to interpolate the value from other means. Ellis & Lindstrom¹⁷ reported the index of refraction as following: 5893 Å, $n=2.105\pm 0.005$; 5641 Å, $n=2.110\pm 0.005$; 4358 Å, $n=2.135\pm 0.005$. Since the values of the integral term on the right hand side of Eq. (III.20) can be calculated at any of those three wavelengths, $n(\infty)$ can be evaluated from Eq. (III.20). A value of $n(\infty) = 2.072$ seems to fit well at all three wavelengths. With this $n(\infty)$ value, $n(\omega)$ at all wavelengths can be readily calculated as shown in Table 4. The results in the table show that the index of refraction varies only slightly with wavelengths.

IV.2 Optical Constants of Sodium & Steel

The Drude model for metals is obtained directly from the Lorentz model simply by setting a equal to zero in Eq. (A.13) and (A.14)

$$n^2 - k^2 = \epsilon_0 - \frac{4\pi N e^2}{m} \frac{1}{(\omega^2 + \gamma^2)}$$

$$2nk = \frac{4\pi N e^2}{m} \frac{\gamma}{(\omega^2 + \gamma^2)}$$

a. Sodium

According to Chan and Tseng⁵, $(4\pi N e^2/m)^{0.5}$ is $8.9 \times 10^{15} \text{ sec}^{-1}$, and the optimum value of γ appears to be $5.142 \times 10^{13} \text{ sec}^{-1}$. They computed the optical constants using the above equations and the results are drafted in Fig. 6 and Table 5 for convenient. The calculated value of the optical constant at $\lambda=0.2\mu$ is replaced by the same value as that at $\lambda=0.25\mu$

due to obvious error discussed in Ref. 5.

b. Steel

Again in Chan and Tseng's work, the optimal values $(4\pi N e^2/m)^{0.5}$ and γ of steel are found to be $2.005 \times 10^{16} \text{ sec}^{-1}$ and $8.26 \times 10^{15} \text{ sec}^{-1}$, respectively. The resulting optical constants are shown in Fig. 7 & Table 6.

IV.3 Spectral Hemispherical Reflectance

Knowing the optical constants of the two contacting materials, it is possible to calculate the spectral hemispherical properties from Eqs. (II.1) to (II.5). The 20-point Gaussian Quadrature is used to integrate Eq. (II.5) to yield the spectral hemispherical reflectance. A sample FORTRAN program is presented in Appendix F.

IV.4 Total Hemispherical Properties

The total hemispherical properties can be calculated from the spectral hemispherical properties by Eqs. (II.6) to (II.10). The trapezoidal method with 44 integration points between 0.2 and 10 is used to evaluate the spectral integral in the equations with a black source temperature ranging from 1200°K to 6000°K .

IV.5 Summary

In general, the radiation properties from the medium 1 side are different from the medium 2 side. It is necessary to indicate from which medium the radiation originates and to which medium it ends up.

In this study we always choose to originate from medium 1 to 2 (subscript "12"). By Kirchhoff's law the spectral absorptance $\alpha_{\lambda 12}$ equals the spectral hemispherical emittance from medium 2, $\epsilon_{\lambda 21}$. Since we take the radiation source as a black body at the same temperature as both media in contact, the total hemispherical absorption α_{12} is equal to the total hemispherical emittance ϵ_{21} .

IV.5.1 ThO₂-Sodium Interface (medium 1-ThO₂, medium 2-Sodium)

In reactor safety analysis, it often concerns with the interreaction between a molten fuel and a reactor coolant. When a molten ThO₂ drops into a liquid sodium, direct contact occurs and the radiation properties at the ThO₂-sodium interface will be needed for calculating radiative cooling from ThO₂. The calculated spectral hemispherical properties of this interface are given in Table 7 and Fig. 9 while the total hemispherical properties of the interface are given in Table 8 and Fig. 10.

IV.5.2 ThO₂-Steel Interface (medium 1-ThO₂, medium 2-Steel)

When a molten ThO₂ is in direct contact with a reactor material such as stainless steel, the properties of ThO₂-Steel interface would be of interest. Since the optical constants available are for the steel (99Fe + 1C) but not for the stainless steel, we can only calculate the properties of ThO₂-Steel (99Fe + 1C) interface. The spectral hemispherical properties of the contact interface are given in Table 9 and Fig. 11 while the total hemispherical properties of the interface are given in Table 10 and Fig. 12.

IV.5.3 Gas-ThO₂ Interface (medium 1-Gas, medium 2-ThO₂)

This interface exists when a vapor film (or any gas film) separates ThO₂ from sodium or water. The spectral hemispherical properties of such an interface are given in Table 11 and Fig. 13 while the total hemispherical properties of the interface given in Table 12 and Fig. 14.

V. COMPARISON & DISCUSSION

The optical constant of thorium dioxide and the radiation properties of three types of interfaces, namely, ThO_2 -Sodium, ThO_2 -Steel and Gas- ThO_2 , have been calculated in the last section. A number of issues concerning those properties will be discussed here, and a comparison of the results will be made with experimental data or other analytical results that are currently available.

V.1 ThO_2 -Sodium Radiation Properties

As for the total hemispherical reflectance of ThO_2 -Sodium, our calculation (see Fig. 10) shows that it lies in the range of 0.9766 to 0.9903 as the temperature increases from 1200°K to 6000°K . From Table 7, it is clear that the spectral hemispherical values may not be accurate at extremely short wavelengths, say $\lambda < 0.2\mu$ due to the inherent error in the optical constants of sodium (see Table 5). Fortunately, the possible error has practically no effect on total radiation property calculations because less than 0.5% of black body radiation energy lies in the region $\lambda < 0.2\mu$ even if the black body is at a temperature as high as 6000°K . No results are available for comparison, nor any experimental data are available since a direct measurement of, say, reflection from an interface is possible only when one side of the interface is vacuum or a gas.

V.2 ThO_2 -Steel Radiation Properties

From the analysis, we report the values of spectral and total hemispherical reflectance of ThO_2 -Steel interface in Table 9 & 10, respectively. Because the spectral hemispherical reflectance in the long wavelength region weights more in Eq. (II.6) as temperature increases, the total hemispherical reflectance increases with temperature. Again no comparison can be made in this case.

V.3 Gas- ThO_2 Radiation Properties

1) Spectral Hemispherical Properties

We report that the value of the reflectance is essentially constant, 0.174. Since the index of refraction is dominant in determining the radiation properties and the indices of refraction of both the interfaces change very little ($n_1 = 1$, $n_2 = 2.11$), a constant value is therefore expected. Some experimental values of normal spectral emissivity of thorium dioxide (thoria) were reported in connection with early works on thoria cathodes emission about thirty years ago. They vary over a wide range such as 0.35,²² 0.2 to 0.7²³ at $\mu=0.65$ while 0.35 was also assumed by Ref. 24 to determine thoria surface temperature from pyrometer measurements. Those values, nevertheless, can not be used for comparison with the present analytical calculations because they are the spectral emissivity from a thin coat of thoria on molybdenum²², tungsten²³ or tantalum.²⁴ Since the coating thickness is only 0.025 to 0.038 mm²², 0.05 mm²³ or 0.1 mm²⁴ and the extinction coefficient of thoria is about 26 cm^{-1} , the thoria coating is obviously in the optically thin region that

the radiation from the base metal penetrates through the coating. Therefore, the reported spectral emissivities do not represent that of a bulk thoria.

2) Total Hemispherical Properties

We integrate the spectral hemispherical reflectance of Gas-ThO₂ interface to obtain the total hemispherical reflectance from 1200°K to 6000°K, and find 0.1741 at 1200°K; 0.1750 at 6000°K (see Fig. 14). Again in connection with the investigations on thoria-coated cathodes, total normal emissivity was reported as 0.5²⁴ with a thoria thickness of 0.5 mm at 1900°K. Reference 25 claimed between 0.2 to 0.6 in the temperature range of 250 to 800°C for a sintered thoria with an unspecified thickness. Though comparison with the present calculation, which is valid for a pure bulk thoria with an optically smooth surface, may not meaningful for the same reason stated above, it may be appropriate to point out that the total normal reflectivity calculated from Eq. (II.11) is about 0.13 which would lead to a total normal emissivity of 0.87, higher than either reported values. Also should be pointed out is the inherent limitation imposed on the present calculation, namely, the optical constants used were derived from low temperature data (mainly the room temperature) as no higher temperature data are available. Thus the temperature dependence of optical constants can not be accounted for in the present investigation.

VI. CONCLUSION

The present study continues Chan and Tseng's work⁵ in providing the surface radiation properties of reactor materials that are useful in the safety analysis of nuclear reactors. Two methods of using K-K theory to calculate the optical constants from measured reflection and transmission spectra have been put into practical use. In particular, the optical constants of thorium dioxide have been calculated and presented here. Also presented are the spectral and total hemispherical properties of thorium dioxide in contacting with sodium steel or a gas. The spectral hemispherical properties are presented over a wide spectral range while the total hemispherical properties over a temperature range of 1200°K to 6000°K.

Table 1
 Reflectance Spectra of the Classical Dispersion for $\alpha=0$

Wavelength λ (μ)	Reflectance in Perpendicular Direction R_s		Experimental Reflectance R_o		Index of Refraction n	Extinction Coefficient k
	$\theta = 60^\circ$	$\theta = 70^\circ$	$\theta = 60^\circ$	$\theta = 70^\circ$		
2.0	0.4326	0.5625	0.2282	0.2816	2.572	0.0
3.0	0.4290	0.5592	0.2259	0.2801	2.550	0.0
4.0	0.4233	0.5541	0.2222	0.2777	2.516	0.0
5.0	0.4148	0.5463	0.2167	0.2742	2.466	0.00035
6.0	0.4019	0.5345	0.2087	0.2689	2.394	0.00073
7.0	0.3815	0.5155	0.1963	0.2607	2.286	0.00155
8.0	0.3460	0.4816	0.1757	0.2465	2.114	0.00337
9.0	0.2705	0.4054	0.1353	0.2156	1.805	0.00834
10.0	0.0082	0.0280	0.0049	0.0209	1.038	0.03485
10.5	0.9472	0.9650	0.9120	0.9377	0.068	0.90216
11.0	0.9781	0.9852	0.9556	0.9620	0.059	1.972
11.5	0.9847	0.9896	0.9658	0.9658	0.078	2.865
12.0	0.9856	0.9901	0.9659	0.9612	0.196	5.046
12.2	0.9842	0.9892	0.9623	0.9552	0.350	6.526
12.4	0.9800	0.9863	0.9517	0.9404	0.958	9.612
12.6	0.9412	0.9594	0.8631	0.8310	16.70	16.49
12.8	0.8273	0.8784	0.6471	0.6016	10.51	0.9323
13.0	0.7727	0.8382	0.5631	0.5242	7.801	0.3325
14.0	0.6600	0.7524	0.4214	0.4106	4.907	0.0503
15.0	0.6157	0.7174	0.3754	0.3777	4.234	0.0227

Table 2
Optical Constants by K-K Theory Using Iterative Method

Wavelength λ (μ)	R_s^*		n^*		k^*	
	$\theta = 60^\circ$; $\theta = 70^\circ$	$\theta = 60^\circ$; $\theta = 70^\circ$	$\theta = 60^\circ$; $\theta = 70^\circ$	$\theta = 60^\circ$; $\theta = 70^\circ$	$\theta = 60^\circ$; $\theta = 70^\circ$	$\theta = 60^\circ$; $\theta = 70^\circ$
2.0	0.4325	0.5624	2.571	2.571	0.0212	0.0188
3.0	0.4290	0.5592	2.550	2.550	0.0217	0.0183
4.0	0.4233	0.5540	2.516	2.515	0.0254	0.0216
5.0	0.4147	0.5463	2.465	2.466	0.0288	0.0233
6.0	0.4019	0.5344	2.394	2.393	0.0276	0.0204
7.0	0.3815	0.5155	2.285	2.285	0.0389	0.0309
8.0	0.3460	0.4816	2.113	2.113	0.0384	0.0297
9.0	0.2704	0.4053	1.804	1.804	0.0287	0.0204
10.0	0.0080	0.0272	1.012	1.017	0.0457	0.0430
10.5	0.9469	0.9648	0.068	0.068	0.8680	0.8794
11.0	0.9779	*	0.056	*	1.898	*
11.5	0.9845	0.9896	0.086	0.087	3.176	3.156
12.0	0.9854	*	0.169	*	4.649	*
12.5	0.9734	*	1.741	*	11.158	*
12.6	0.9411	*	8.597	*	14.438	*
12.7	0.8707	0.9078	12.195	7.269	5.250	7.031
12.8	0.8272	0.8771	9.860	8.234	2.659	4.256
13.0	0.7726	0.8372	7.569	7.120	1.342	2.018
13.5	0.6998	0.7826	5.602	5.509	0.660	0.8934
14.0	0.6599	0.7515	4.858	4.808	0.4713	0.6368
15.0	0.6155	0.7166	4.189	4.171	0.4044	0.4394

* divergence

Table 3

Index of Refraction from the Values of Extinction Coefficient
by Integral K-K Dispersion Relations

Wavelength λ (μ)	Index of Refraction from K-K Analysis n^*	Index of Refraction n	Extinction Coefficient k
5.0	2.516	2.466	0.00035
6.0	2.441	2.394	0.00073
7.0	2.331	2.286	0.00155
8.0	2.156	2.114	0.00337
9.0	1.842	1.805	0.00834
10.0	1.046	1.038	0.03485
11.0	0.0853	0.059	1.972
12.0	0.1760	0.196	5.046
12.1	0.2216	0.254	5.686
12.2	0.3010	0.350	6.526
12.3	0.4604	0.532	7.713
12.4	0.8756	0.958	9.612
12.5	3.419	2.592	13.412
12.6	12.323	16.70	16.495
12.7	15.590	14.04	2.498
12.8	11.001	10.51	0.9323
12.9	9.086	8.815	0.5115
13.0	7.979	7.801	0.3325
14.0	4.996	5.684	0.0503
15.0	4.307	4.234	0.0227
16.0	3.989	3.924	0.0138

Table 4
Optical Constants of ThO₂

Wavelength λ (μ)	Index of Refraction n	Extinction Coefficient k
0.20	2.118	0.0003053
0.25	2.117	0.0001717
0.30	2.116	0.0001694
0.35	2.116	0.0001376
0.40	2.115	0.0001374
0.45	2.114	0.0001168
0.50	2.114	0.0001202
0.55	2.114	0.0001322
0.60	2.113	0.0001488
0.65	2.113	0.0001637
0.70	2.113	0.0001736
0.75	2.112	0.0001717
0.80	2.112	0.0001526
0.85	2.112	0.0001297
0.90	2.112	0.0001030
0.95	2.111	0.0000906
1.00	2.111	0.0000763
1.05	2.111	0.0000641
1.10	2.111	0.0000545
1.15	2.111	0.0000482
1.20	2.111	0.0000458
1.25	2.111	0.0000452

Table 4 (Cont'd)

Wavelength λ (μ)	Index of Refraction n	Extinction Coefficient k
1.30	2.111	0.0000446
1.35	2.111	0.0000436
1.40	2.111	0.0000427
1.45	2.110	0.0000385
1.50	2.110	0.0000343
2.00	2.108	0.0000191
2.50	2.107	0.0000275
3.00	2.107	0.0000435
3.50	2.106	0.0000372
4.00	2.105	0.0000305
4.50	2.105	0.0000326
5.00	2.104	0.0000471
5.50	2.104	0.0000242
6.00	2.103	0.0000092
6.50	2.103	0.0
7.00	2.102	0.0
7.50	2.102	0.0
8.00	2.101	0.0000519
8.50	2.100	0.00026
9.00	2.099	0.000515
9.50	2.099	0.002624
10.00	2.099	0.003513

Table 5
Optical Constants of Sodium

Wavelength λ (μ)	Index of Refraction n		Extinction Coefficient k	
	Present	References	Present	References
0.20	0.007		0.008	
0.25	0.007	0.025*(18)	0.646	0.585*(18)
0.30	0.008	0.036*(18)	1.020	0.944*(18)
0.35	0.009	0.043*(18)	1.333	1.319*(18)
0.40	0.012	0.047*(18)	1.621	1.545*(18)
0.45	0.014	0.045*(18)	1.895	1.866*(18)
		0.058*(19)		1.864*(19)
0.50	0.017	0.036*(18)	2.177	2.102*(18)
		0.060*(19)		2.092*(19)
0.55	0.021	0.052*(19)	2.420	2.229*(19)
		0.047*(20)		2.238*(20)
0.60	0.024	0.045*(19)	2.676	2.476*(19)
0.65	0.029	0.049*(19)	2.929	2.733*(19)
0.70	0.033		3.180	
0.75	0.038		3.429	
0.80	0.043		3.676	
0.85	0.048		3.922	
0.90	0.054		4.167	
0.95	0.062		4.265	
1.00	0.066		4.656	
1.05	0.070		4.899	
1.10	0.080		5.141	
1.15	0.087		5.384	
1.20	0.094		5.626	
1.25	0.102		5.868	
1.30	0.111		6.110	

Table 5 (Cont'd)

Wavelength λ (μ)	Index of Refraction n		Extinction Coefficient k	
	Present	References	Present	References
1.35	0.119		6.351	
1.40	0.128		6.591	
1.45	0.138		6.833	
1.50	0.147		7.074	
2.00	0.261		9.474	
2.25		0.410(18)		11.60(18)
2.50	0.407	0.500(18)	11.855	12.50(18)
3.0	0.585	0.610(18)	14.261	14.70(18)
3.5	0.795	0.810(18)	16.655	17.20(18)
4.0	1.038	1.020(18)	19.049	19.40(18)
4.5	1.260		22.355	
4.75		1.71 (18)		22.30(18)
5.0	1.621		23.843	
5.5	1.961		26.251	
6.0	2.331		28.648	
6.5	2.734		31.056	
7.0	3.169		33.468	
7.2		3.66 (18)		32.9 (18)
7.5	3.636		35.881	
8.0	4.134		38.305	
8.5	4.633		40.731	
9.0	5.223		43.162	
9.5	5.815		45.599	
10.0	6.437		48.041	

Table 6
Optical Constants of Steel (99 Fe + 1 C)

Wavelength λ (μ)	Index of Refraction n		Extinction Coefficient k	
	Present	References	Present	References
0.20	0.806		1.626	
0.25	1.048	1.353*(18)	1.963	1.818*(18)
0.30	1.281	1.397*(18)	2.237	2.008*(18)
0.35	1.501	1.472*(18)	2.466	2.357*(18)
0.40	1.707	1.680*(18)	2.666	2.721*(18)
0.45	1.898	1.895*(18)	2.842	2.927*(18)
0.50	2.078	2.130*(18)	3.001	3.109*(18)
0.55	2.246	2.327*(18)	3.149	3.261*(18)
0.60	2.405	2.490*(18)	3.286	3.337*(18)
0.65	2.555	2.615*(18)	3.414	3.434*(18)
0.70	2.700	2.700(18)	3.537	3.537(18)
0.75	2.834		3.651	
0.80	2.964		3.762	
0.85	3.089		3.868	
0.90	3.210		3.971	
0.95	3.326		4.070	
1.00	3.438		4.167	
1.05	3.547		4.261	
1.10	3.653		4.352	
1.15	3.755		4.441	
1.20	3.855		4.528	
1.25	3.953		4.613	
1.30	4.047		4.696	

Table 6 (Cont'd)

Wavelength λ (μ)	Index of Refraction n		Extinction Coefficient k	
	Present	References	Present	References
1.35	4.140		4.778	
1.40	4.231		4.858	
1.45	4.320		4.937	
1.50	4.456		5.073	
2.00	5.200		5.723	
2.50	5.880		6.356	
3.00	6.494		6.929	
3.50	7.049		7.452	
4.00	7.566		7.943	
4.50	8.049		8.405	
5.00	8.505		8.843	
5.50	8.937		9.260	
6.00	9.350		9.659	
6.50	9.745		10.042	
7.00	10.125		10.410	
7.50	10.491		10.766	
8.00	10.844		11.111	
8.50	11.186		11.446	
9.00	11.518		11.771	
9.50	11.841		12.086	
10.00	12.155		12.395	

* Interpolated Values.

Table 7
 ThO₂-Sodium Interface Spectral Hemispherical
 Properties (medium 1-ThO₂ ; medium 2-Sodium)

Wavelength λ (μ)	$\alpha_{\lambda 12}$	$\rho_{\lambda 11}$
0.20	0.48338	0.51662
0.25	0.00559	0.99441
0.30	0.00731	0.99269
0.35	0.00816	0.99183
0.40	0.01031	0.98969
0.45	0.01116	0.98884
0.50	0.01237	0.98763
0.55	0.01403	0.98597
0.60	0.01464	0.98536
0.65	0.01613	0.98387
0.70	0.01675	0.98325
0.75	0.01763	0.98237
0.80	0.01827	0.98173
0.85	0.01872	0.98128
0.90	0.01937	0.98063
0.95	0.02149	0.97851
1.00	0.02020	0.97980
1.05	0.01988	0.98012
1.10	0.02112	0.97888
1.15	0.02140	0.97860
1.20	0.02159	0.97841

Table 7 (Cont'd)

Wavelength λ (μ)	$\alpha_{\lambda 12}$	$P_{\lambda 11}$
1.25	0.02193	0.97807
1.30	0.02237	0.97763
1.35	0.02253	0.97747
1.40	0.02282	0.97718
1.45	0.02317	0.97683
1.50	0.02330	0.97670
2.00	0.02504	0.97496
2.50	0.02616	0.97384
3.00	0.02683	0.97317
3.50	0.02731	0.97269
4.00	0.02768	0.97232
4.50	0.02488	0.97512
5.00	0.02819	0.97191
5.50	0.02835	0.97165
6.00	0.02847	0.97153
6.50	0.02853	0.97147
7.00	0.02857	0.97143
7.50	0.02861	0.97139
8.00	0.02861	0.97139
8.50	0.02843	0.97157
9.00	0.02859	0.97141
9.50	0.02877	0.97123
10.00	0.02878	0.97122

Table 8
 ThO₂-Sodium Interface Total Hemispherical
 Properties (medium 1-ThO₂ ; medium 2-Sodium)

Temperature (°K)	α_{12}	P_{11}
1200	0.02337	0.97663
1400	0.02239	0.97761
1600	0.02145	0.97855
1800	0.02054	0.97946
2000	0.01967	0.98033
2200	0.01884	0.98116
2400	0.01804	0.98196
2600	0.01729	0.98271
2800	0.01658	0.98342
3000	0.01591	0.98409
3200	0.01527	0.98473
3400	0.01468	0.98532
3600	0.01412	0.98588
3800	0.01360	0.98640
4000	0.01312	0.98688
4200	0.01266	0.98734
4400	0.01223	0.98777
4600	0.01183	0.98817
4800	0.01146	0.98854
5000	0.01111	0.98889
5400	0.01049	0.98951
5800	0.00994	0.99006
6000	0.00970	0.99030

Table 9
 ThO₂-Steel Interface Spectral Hemispherical
 Properties (medium 1-ThO₂ ; medium 2-Steel)

Wavelength λ (μ)	$\alpha_{\lambda 12}$	$\rho_{\lambda 11}$
0.20	0.48194	0.51806
0.25	0.53290	0.46710
0.30	0.56324	0.43676
0.35	0.58165	0.41835
0.40	0.59233	0.40767
0.45	0.59815	0.40185
0.50	0.60107	0.39893
0.55	0.60152	0.39848
0.60	0.60061	0.39939
0.65	0.59880	0.40120
0.70	0.59627	0.40373
0.75	0.59320	0.40680
0.80	0.58981	0.41019
0.85	0.58620	0.41380
0.90	0.58256	0.41744
0.95	0.57867	0.42133
1.00	0.57473	0.42527
1.05	0.57083	0.42917
1.10	0.56698	0.43302
1.15	0.56310	0.43690
1.20	0.55927	0.44073
1.25	0.55551	0.44449

Table 9 (Cont'd)

Wavelength λ (μ)	$\alpha_{\lambda 12}$	$P_{\lambda 11}$
1.30	0.55177	0.44823
1.35	0.54806	0.45194
1.40	0.54445	0.45555
1.45	0.54074	0.45926
1.50	0.53419	0.46581
2.00	0.50549	0.49451
2.50	0.47873	0.52127
3.00	0.45646	0.54354
3.50	0.43748	0.56252
4.00	0.42087	0.57913
4.50	0.40637	0.59363
5.00	0.39335	0.60665
5.50	0.38177	0.61823
6.00	0.37117	0.62883
6.50	0.36162	0.63838
7.00	0.35278	0.64722
7.50	0.34472	0.65528
8.00	0.33715	0.66285
8.50	0.33011	0.66989
9.00	0.32355	0.67645
9.50	0.31778	0.68222
10.00	0.31217	0.68783

Table 10
 ThO₂-Steel Interface Total Hemispherical
 Properties (medium 1-ThO₂ ; medium 2-Steel)

Temperature (°K)	α_{12}	ρ_{11}
1200	0.52343	0.47657
1400	0.54060	0.45940
1600	0.55386	0.44614
1800	0.56392	0.43608
2000	0.57138	0.42862
2200	0.57680	0.42320
2400	0.58059	0.41941
2600	0.58308	0.41692
2800	0.58452	0.41548
3000	0.58512	0.41488
3200	0.58502	0.41498
3400	0.58432	0.41568
3600	0.58321	0.41679
3800	0.58169	0.41831
4000	0.57985	0.42015
4200	0.57776	0.42224
4400	0.57546	0.42454
4600	0.57302	0.42698
4800	0.57047	0.42953
5000	0.56784	0.43216
5400	0.56248	0.43752
5800	0.55712	0.44288
6000	0.55449	0.44551

Table 11

Gas-ThO₂ Interface Spectral Hemispherical
Properties (medium 1-Gas ; medium 2-ThO₂)

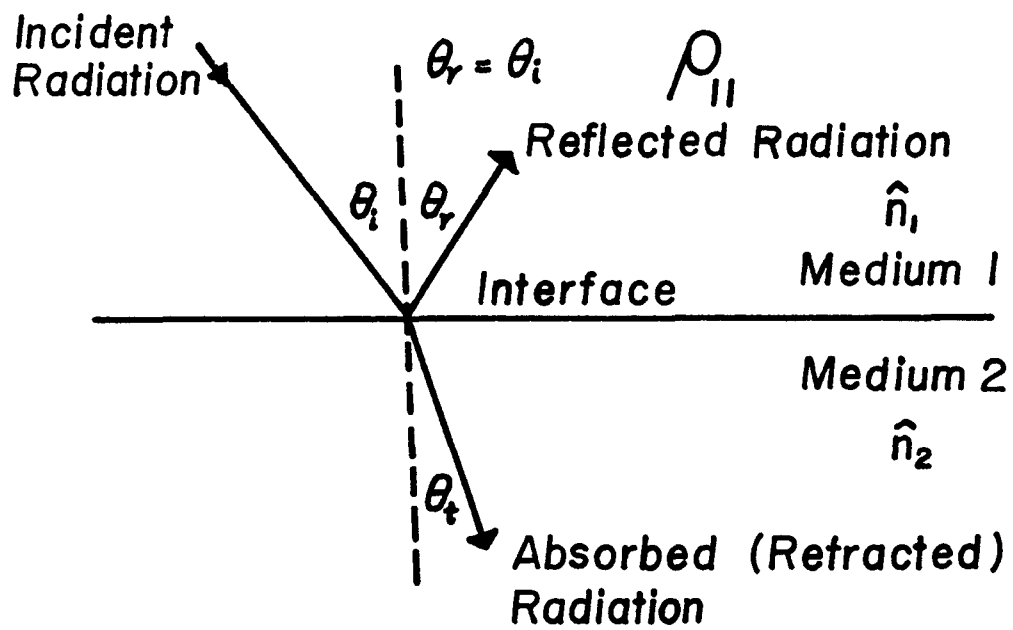
Wavelength λ (μ)	$\alpha_{\lambda 12}$	$\rho_{\lambda 11}$
0.20	0.82430	0.17570
0.25	0.82442	0.17558
0.30	0.82455	0.17545
0.35	0.82455	0.17545
0.40	0.82468	0.17532
0.45	0.82480	0.17520
0.50	0.82480	0.17520
0.55	0.82480	0.17520
0.60	0.82493	0.17507
0.65	0.82493	0.17503
0.70	0.82493	0.17503
0.75	0.82506	0.17494
0.80	0.82506	0.17404
0.85	0.82506	0.17404
0.90	0.82506	0.17404
0.95	0.82518	0.17482
1.00	0.82518	0.17482
1.05	0.82518	0.17482
1.10	0.82518	0.17482
1.15	0.82518	0.17482
1.20	0.82518	0.17482
1.25	0.82518	0.17482

Table 11 (Cont'd)

Wavelength λ (μ)	$\alpha_{\lambda 12}$	$\rho_{\lambda 11}$
1.30	0.82518	0.17482
1.35	0.82518	0.17482
1.40	0.82518	0.17482
1.45	0.82531	0.17469
1.50	0.82531	0.17469
2.00	0.82556	0.17444
2.50	0.82569	0.17431
3.00	0.82569	0.17431
3.50	0.82582	0.17418
4.00	0.82594	0.17406
4.50	0.82594	0.17406
5.00	0.82607	0.17393
5.50	0.82607	0.17393
6.00	0.82620	0.17380
6.50	0.82620	0.17380
7.00	0.82632	0.17368
7.50	0.82632	0.17368
8.00	0.82645	0.17355
8.50	0.82658	0.17342
9.00	0.82670	0.17330
9.50	0.82670	0.17330
10.00	0.82670	0.17330

Table 12
Gas-ThO₂ Interface Total Hemispherical
Properties (medium 1-Gas ; medium 2-ThO₂)

Temperature (°K)	α_{12}	ρ_{11}
1200	0.82588	0.17412
1400	0.82578	0.17422
1600	0.82570	0.17430
1800	0.82563	0.17437
2000	0.82556	0.17444
2200	0.82551	0.17449
2400	0.82546	0.17454
2600	0.82542	0.17458
2800	0.82538	0.17462
3000	0.82534	0.17466
3200	0.82531	0.17469
3400	0.82527	0.17473
3600	0.82524	0.17476
3800	0.82522	0.17478
4000	0.82519	0.17481
4200	0.82517	0.17483
4400	0.82514	0.17486
4600	0.82512	0.17488
4800	0.82510	0.17490
5000	0.82508	0.17492
5400	0.82504	0.17496
5800	0.82500	0.17500
6000	0.82499	0.17501



$$\alpha_{12} = 1 - \rho_{11}$$

Fig. 1 Incident, reflected and refracted radiation at Interface

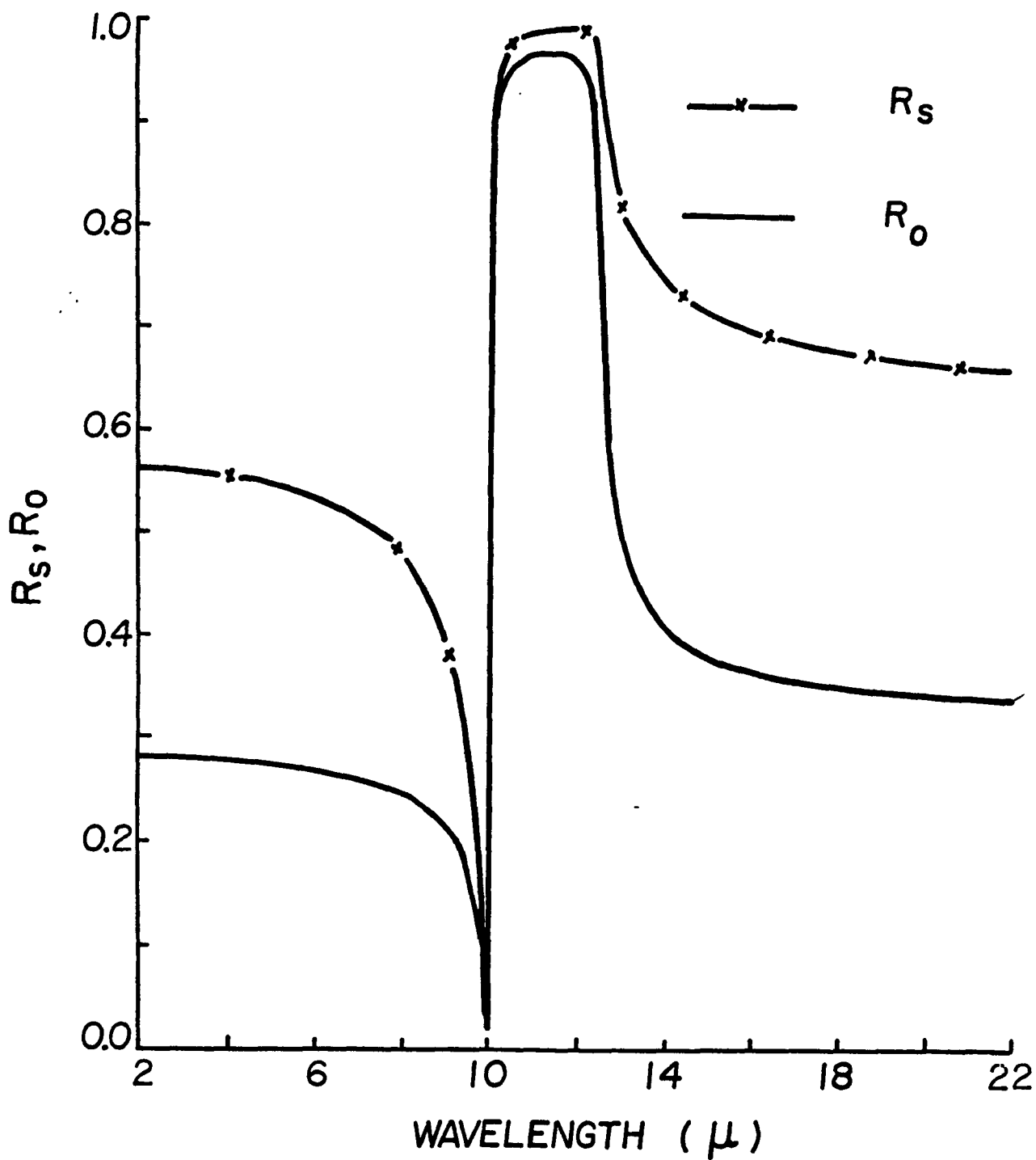


Fig. 2 Reflectance spectra of the Lorentz oscillator for the incidence angle $\theta = 70^\circ$; polarization $\alpha = 0$.

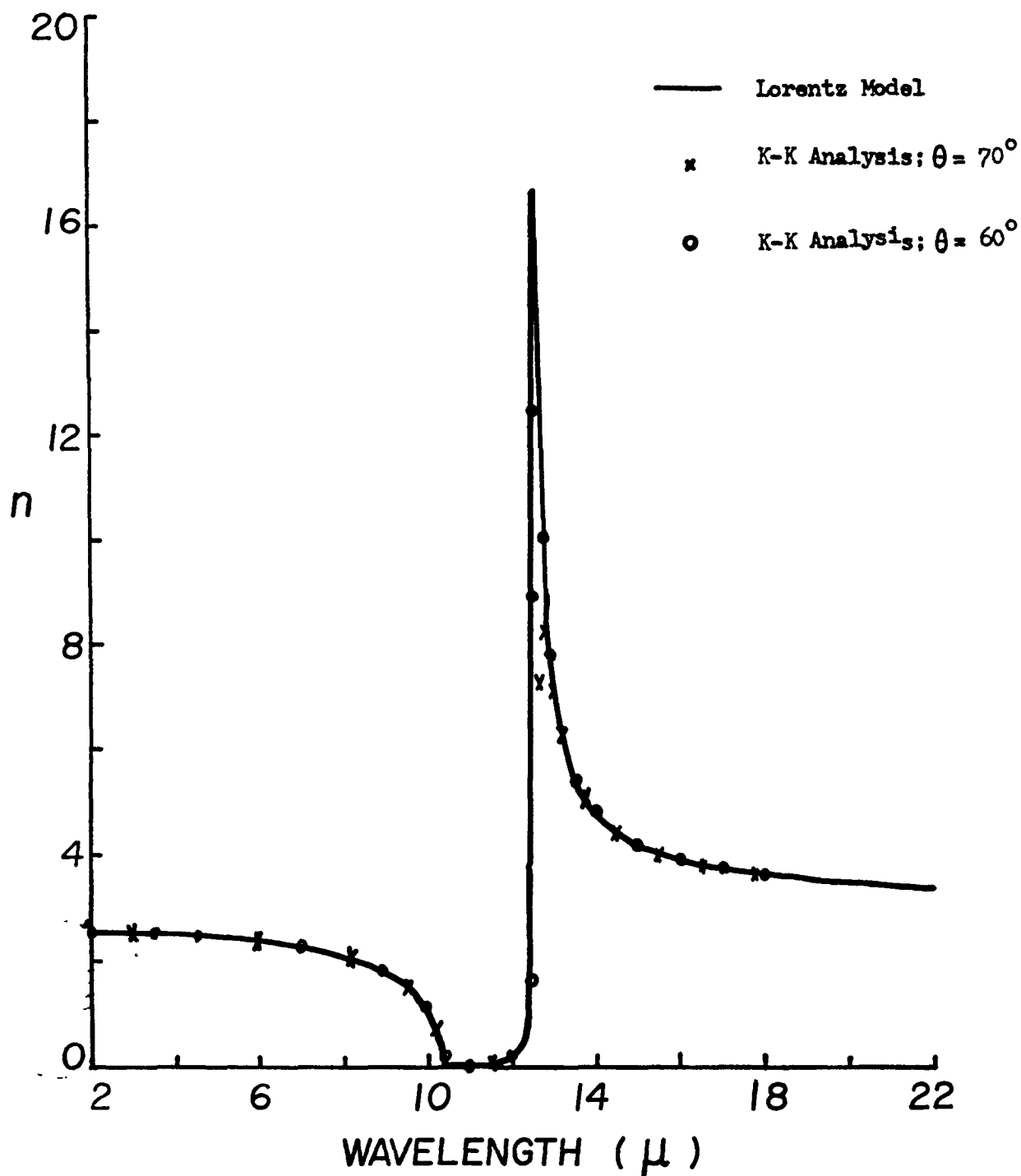


Fig. 3 Index of Refraction by Iterative K-K analysis compare with the values from Lorentz oscillator.

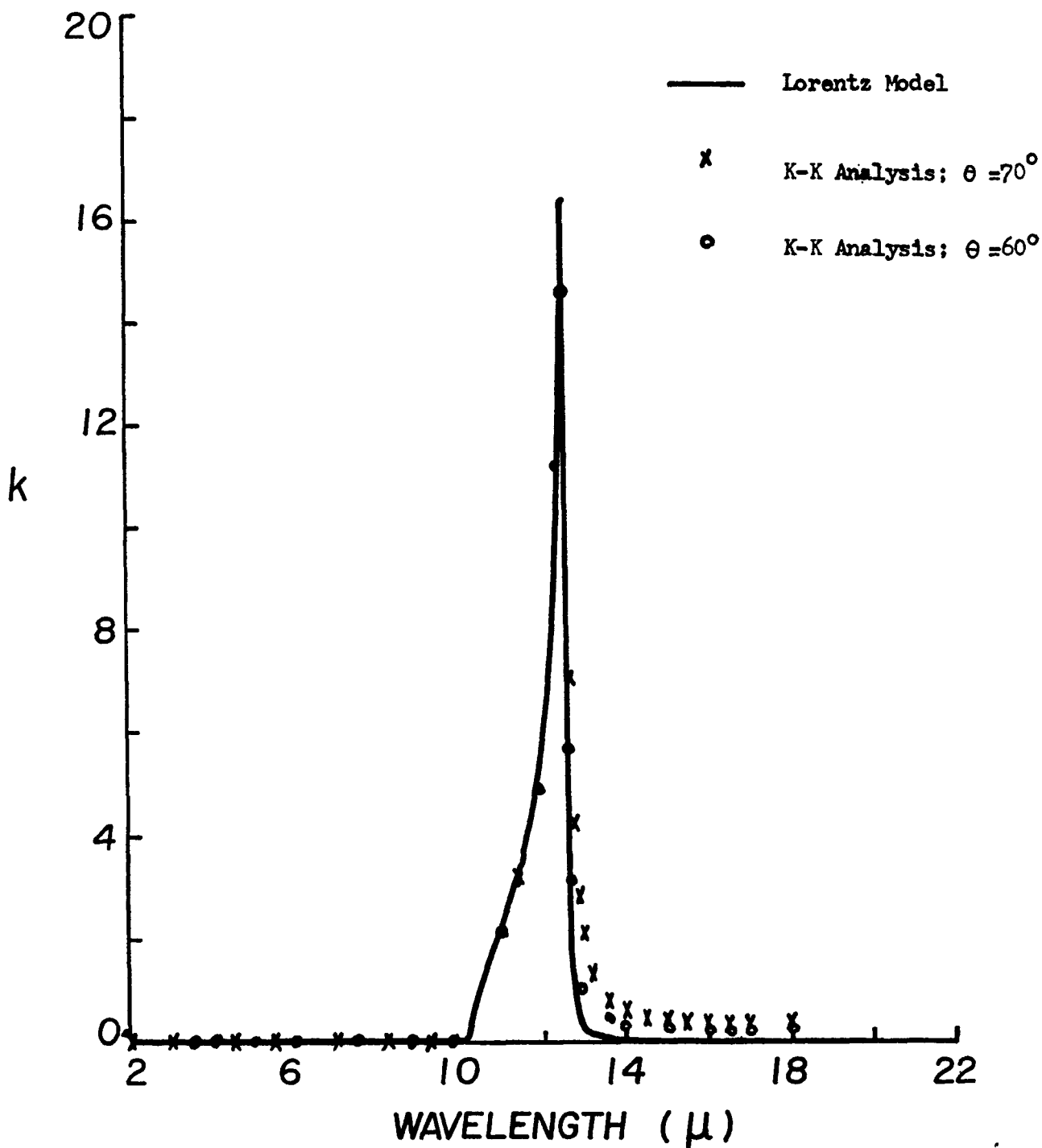


Fig. 4 Extinction Coefficient by Iterative K-K analysis compare with the values from Lorentz oscillator.

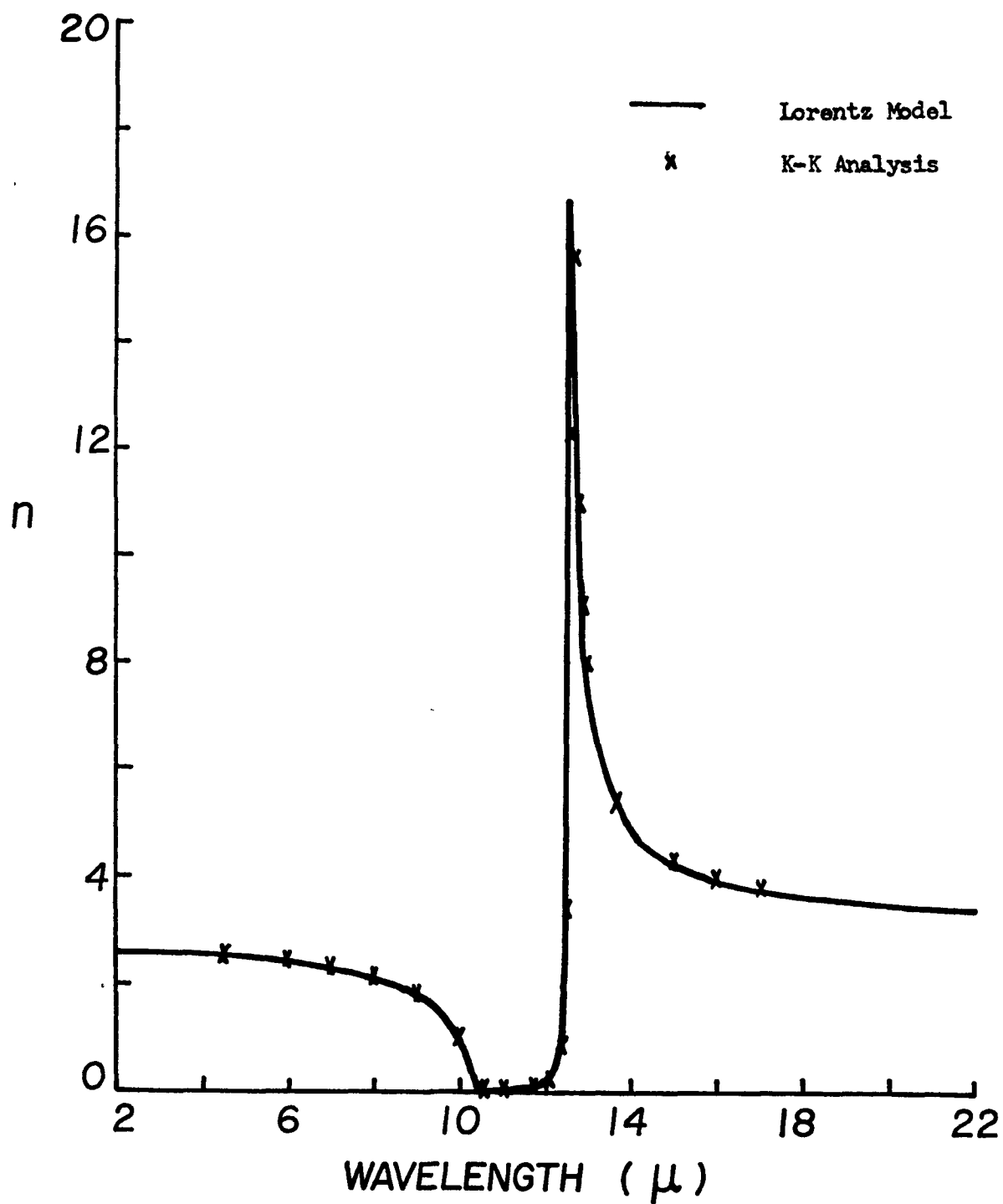


Fig. 5 Index of Refraction from the values of Extinction Coefficient with K-K analysis.

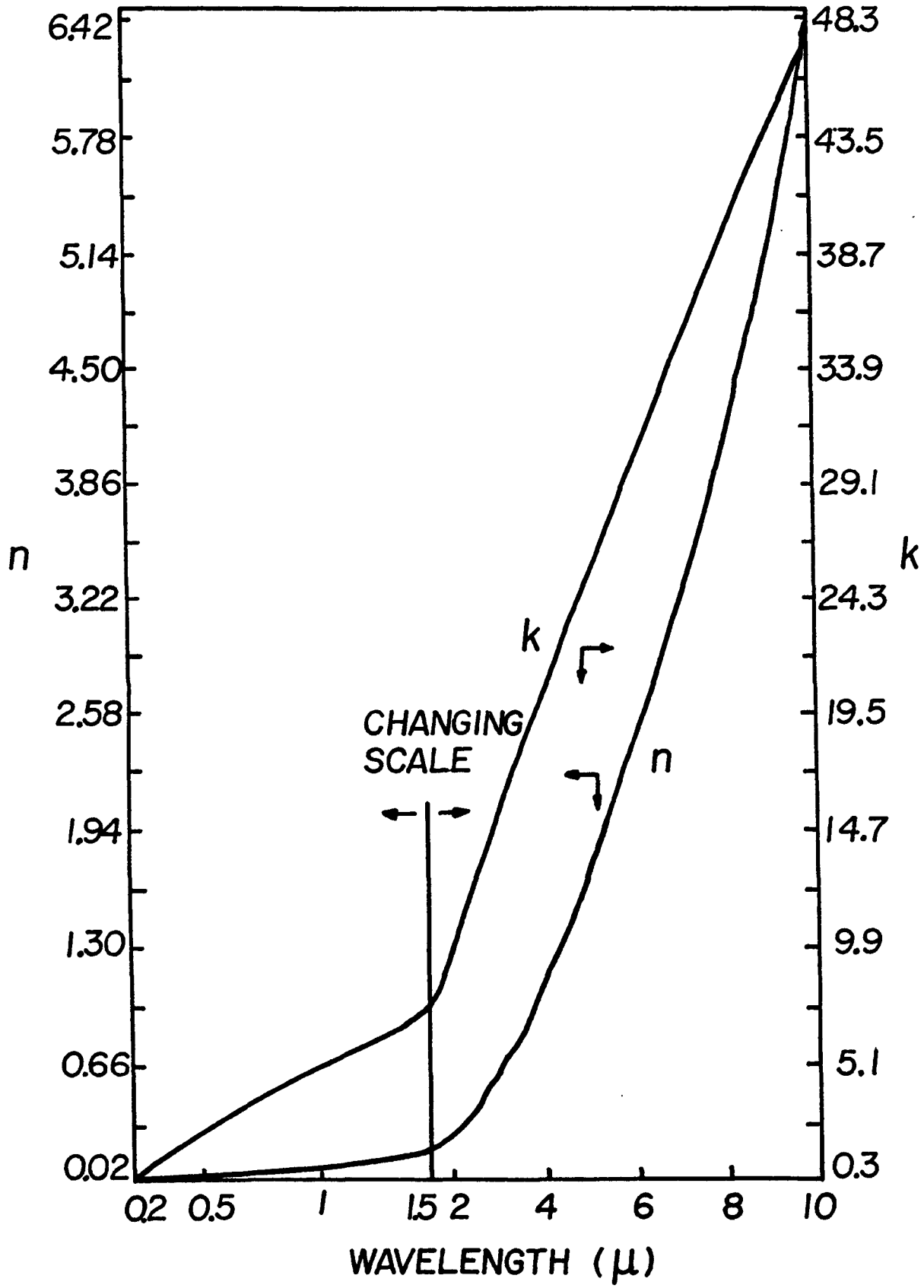


Fig. 6 Optical Constants of Sodium at Various Wavelength.

(From Ref. 5)

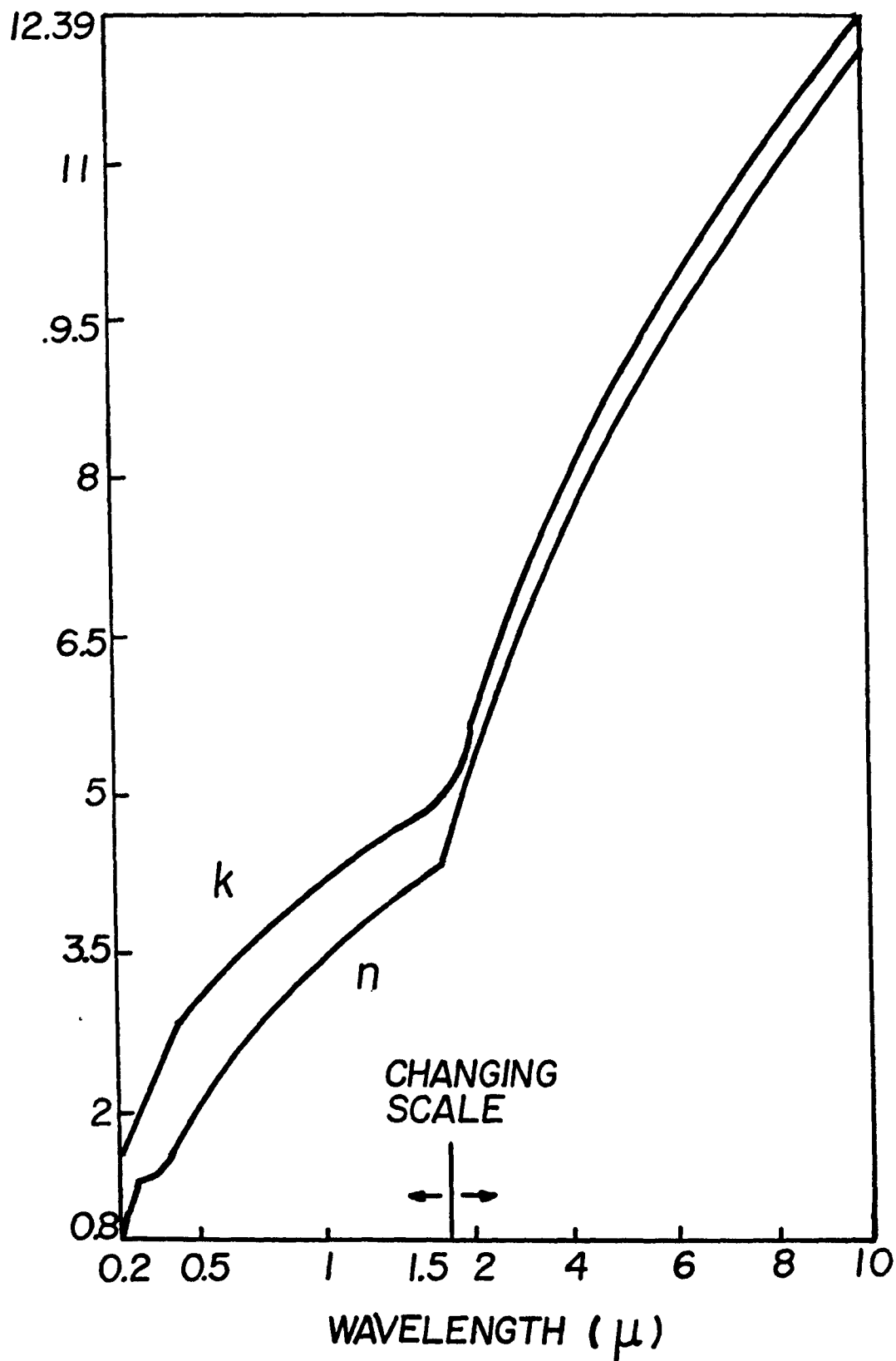


Fig. 7 Optical Constants of Steel at Various Wavelength
(From Ref. 5)

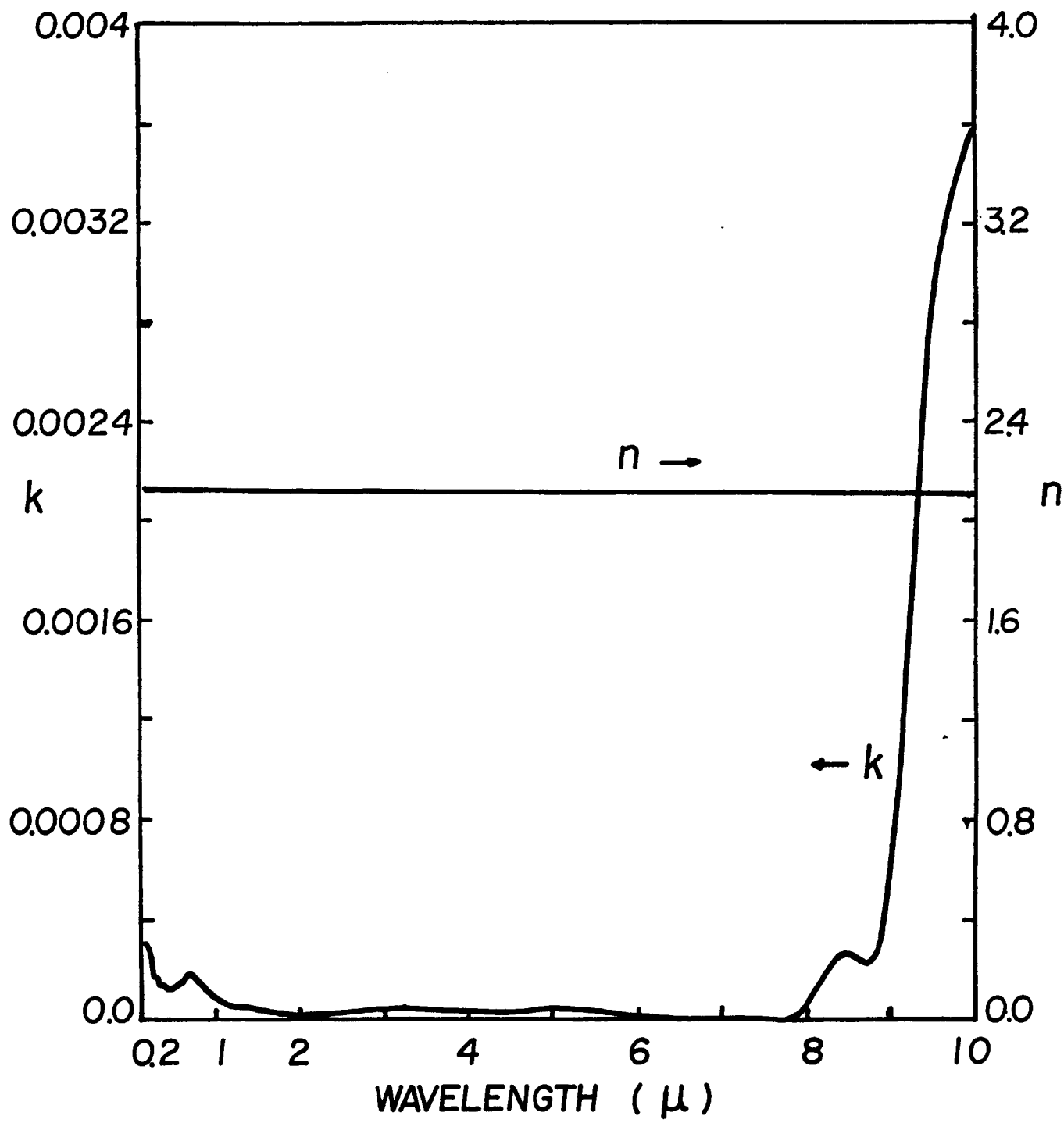


Fig. 8 Optical Constants of ThO₂ at Various Wavelengths

Based on the data of Table 4

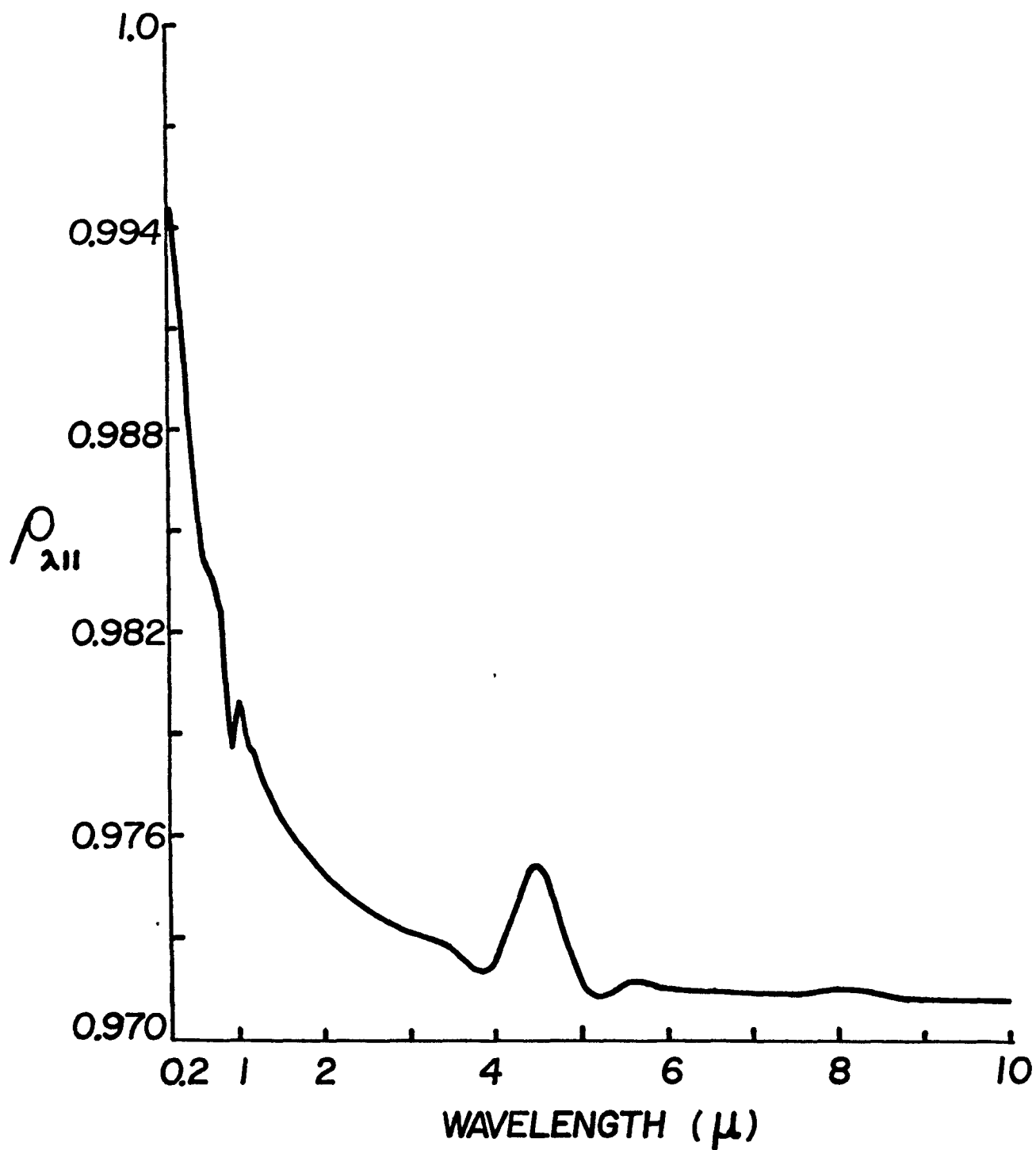


Fig. 9 ThO₂-Sodium Interface Spectral Hemispherical Reflectance

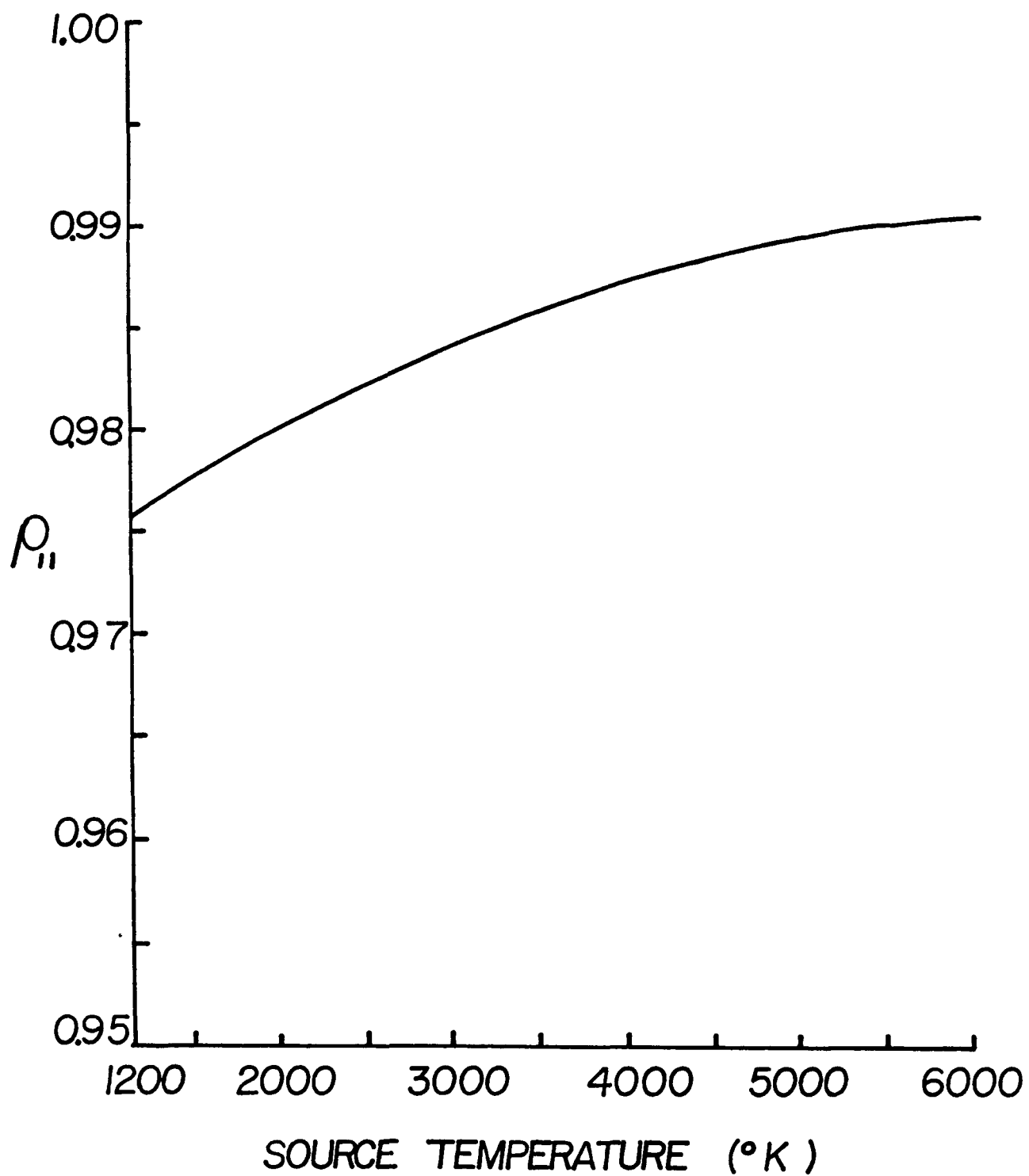


Fig. 10 ThO_2 -Sodium Interface Total Hemispherical Reflectance

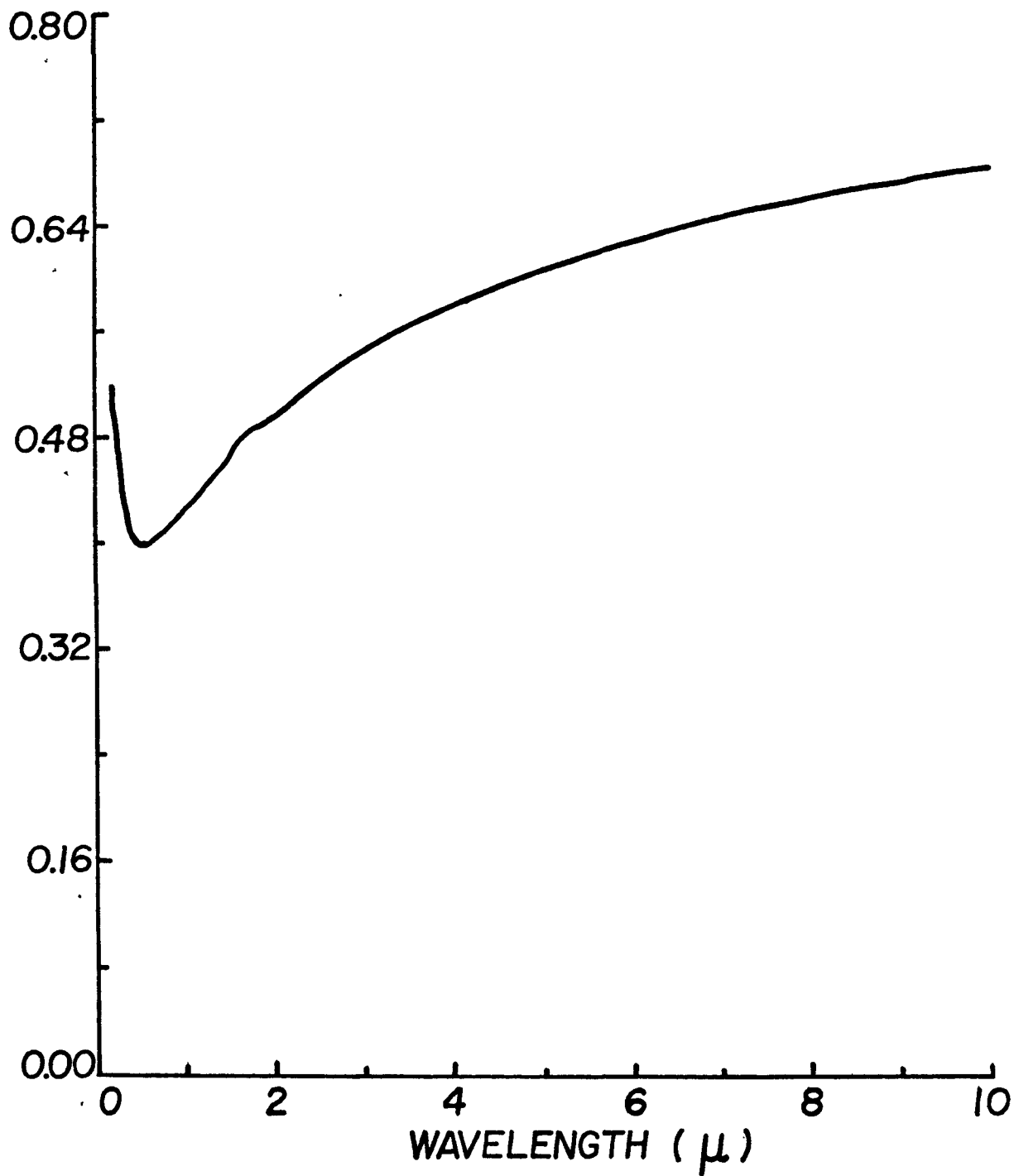


Fig. 11 ThO₂-Steel Interface Spectral Hemispherical Reflectance

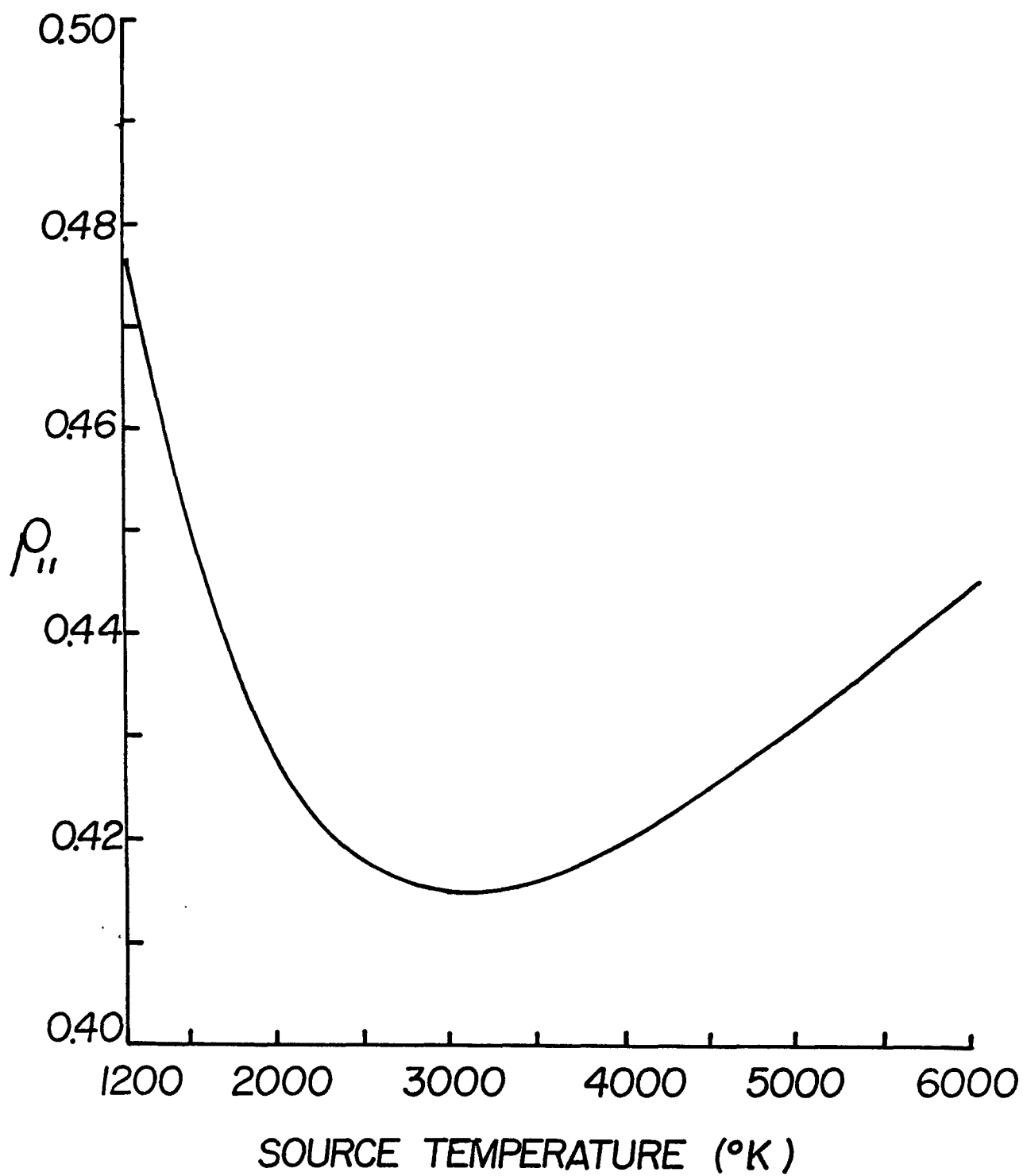


Fig. 12 ThO₂-Steel Interface Total Hemispherical Reflectance

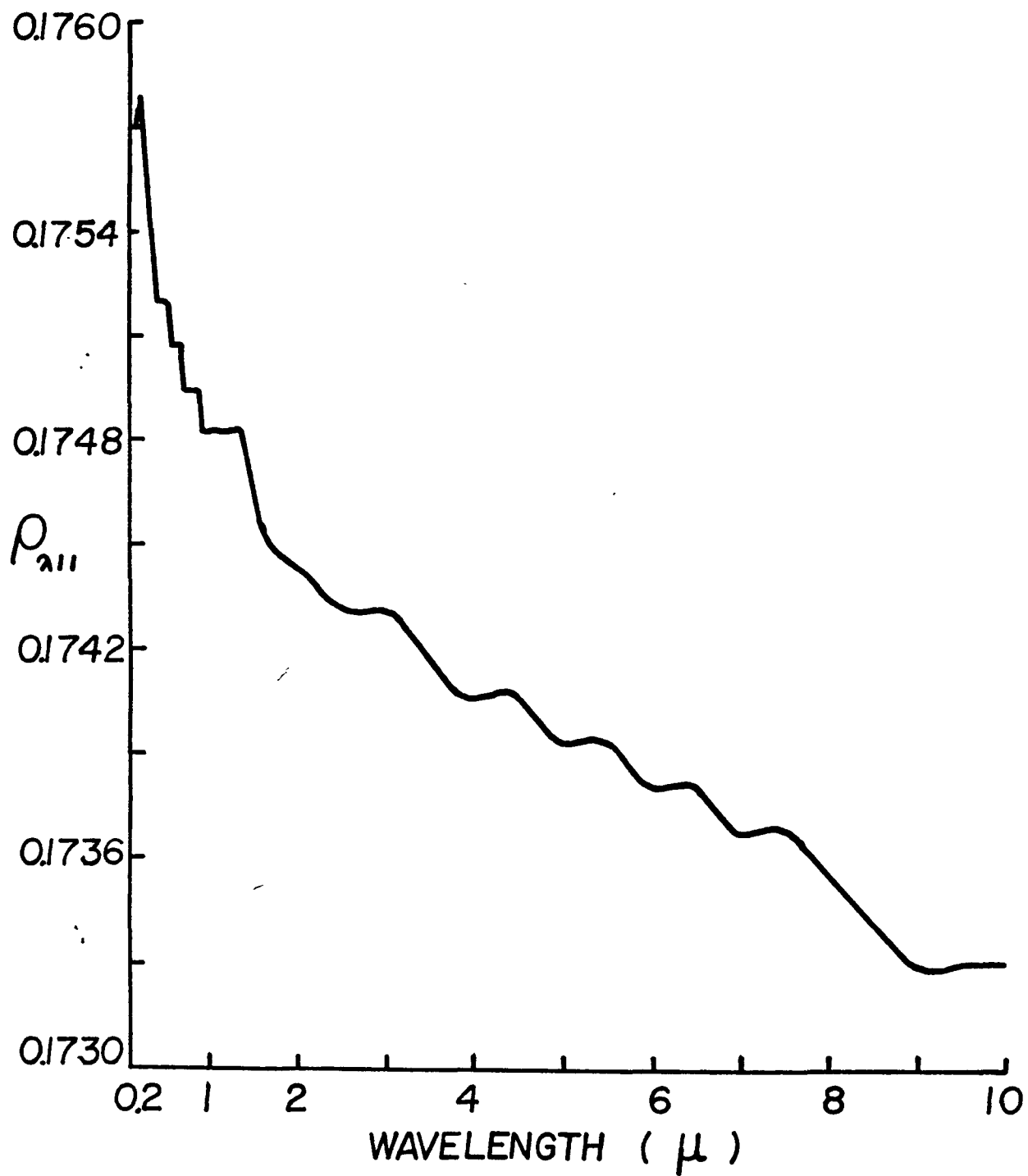


Fig. 13 Gas-ThO₂ Interface Spectral Hemispherical Reflectance

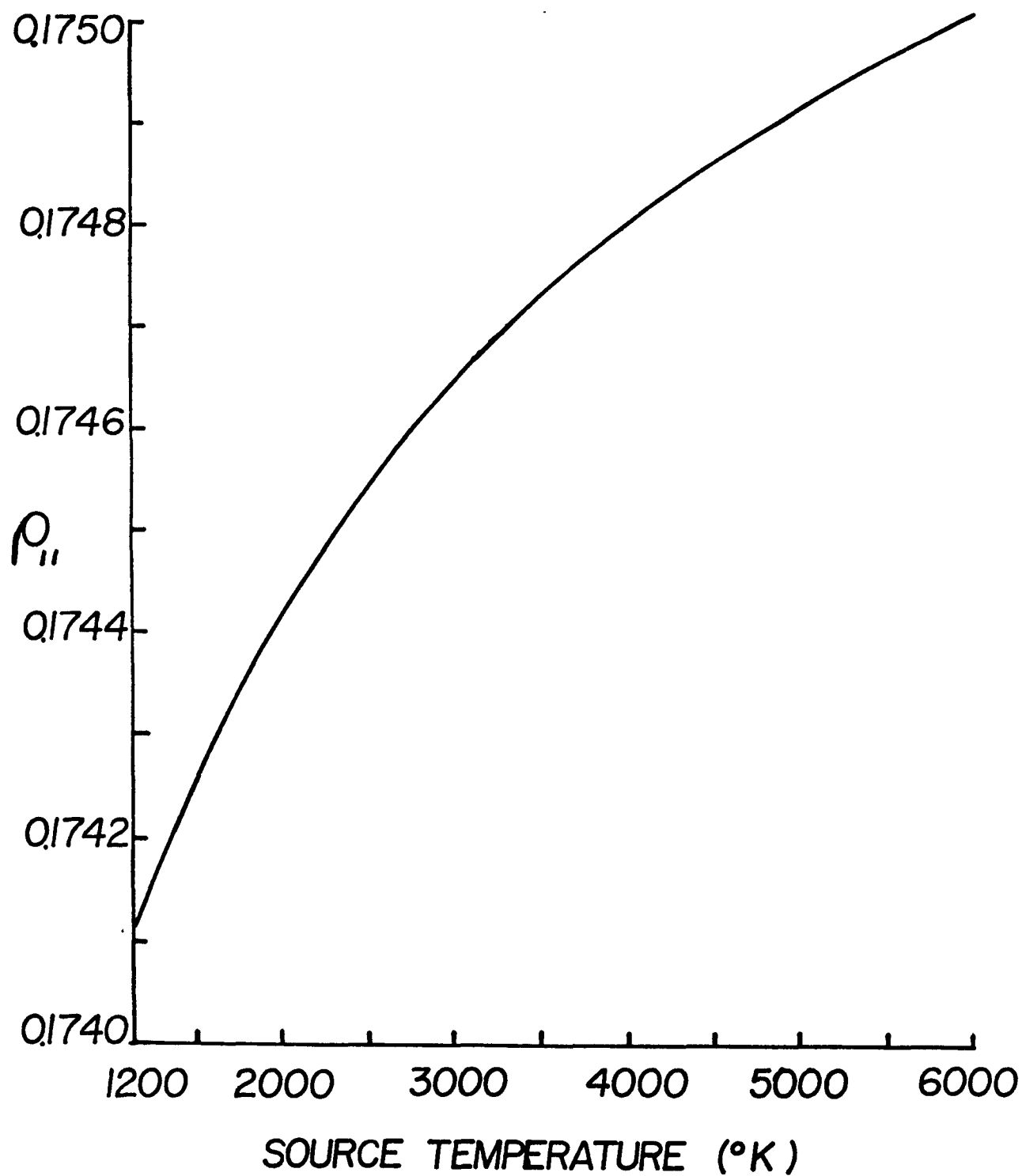


Fig. 14 Gas-ThO₂ Interface Total Hemispherical Reflectance

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APPENDIX A

LORENTZ OSCILLATOR (SEE REF. 13, pp. 42-45)

Consider an atom with electrons bound to the nucleus in much the same way as a small mass can be bound to a large mass by a spring.

The motion of an electron bound to the nucleus can be represented by,

$$m \frac{d^2 r}{dt^2} + m\gamma \frac{dr}{dt} + ma^2 r = -eE_{loc} \quad (A.1)$$

electronic acceleration force	viscous damping force	Hooke's restoring force	electric force
-------------------------------------	-----------------------------	-------------------------------	-------------------

where: m - electronic mass

e - magnitude of electronic charge

a - resonance frequency

γ - damping constant

r - displacement of electron

E_{loc} - local microscopic electric field acting on the electron

The local field can be taken to vary in time as $e^{i\omega t}$. Thus, the solution to Eq. (A.1) is

$$\hat{r} = \frac{-eE_{loc}/m}{(a^2 - \omega^2) - i\gamma\omega} \quad (A.2)$$

The induced dipole moment is

$$\hat{p} = \frac{e^2 E_{loc}/m}{(a^2 - \omega^2) - i\gamma\omega} \quad (A.3)$$

Note that it is important to be consistent in the form of the time variation in the electric field with the definition of the complex refractive index. The use of a time variation $e^{i\omega t}$ leads to a complex

refractive index $\hat{n} = n - ik$, and $e^{-i\omega t}$ leads to $\hat{n} = n + ik$.

Assume a linear relationship exists between \hat{p} and E_{loc} if the displacement r is sufficiently small,

$$\hat{p} = \hat{\alpha}(\omega) E_{loc} \quad (A.4)$$

where $\hat{\alpha}(\omega)$ is frequency-dependent atomic polarizability, and for a one-electron atom,

$$\hat{\alpha}(\omega) = (e^2/m) / [(a^2 - \omega^2) + i\gamma\omega] \quad (A.5)$$

For N one-electron atoms per unit volume, the macroscopic polarization is

$$P = N\hat{p} = N\hat{\alpha}(\omega) E_{loc} \quad (A.6)$$

Proceeding with the assumption the microscopic electric field E_{loc} can be approximated by the macroscopic field E , we have

$$P = N\hat{\alpha}(\omega)E = X_e E \quad (A.7)$$

X_e is macroscopic electric susceptibility.

Define a complex displacement \hat{D} such that

$$\hat{D} = \hat{\epsilon}E = \epsilon_0 E + 4\pi P \quad (A.8)$$

From Eqs. (A.7) and (A.8), the dielectric function $\hat{\epsilon} (= \epsilon_1 + i\epsilon_2)$ is

$$\hat{\epsilon} = \epsilon_0 + 4\pi N\hat{\alpha}(\omega) \quad (A.9)$$

Thus,

$$\hat{\epsilon} = \epsilon_0 + \frac{4\pi Ne^2}{m} \frac{1}{(a^2 - \omega^2) + i\gamma\omega} \quad (A.10)$$

From Maxwell's equations

$$\epsilon_1 = (n^2 - k^2) / \mu \quad (A.11)$$

$$\epsilon_2 = 2nk / \mu = 4\pi\sigma / \omega \quad (A.12)$$

μ : magnetic permeability. σ : conductivity. For non-magnetic materials, μ is equal to unity. We can find

$$\epsilon_1 = n^2 - k^2 = \epsilon_0 + \frac{4\pi N e^2}{m} \frac{(a^2 - \omega^2)}{(a^2 - \omega^2) + \gamma^2 \omega^2} \quad (\text{A.13})$$

$$\epsilon_2 = 2nk = - \frac{4 N e^2}{m} \frac{\gamma \omega}{(a^2 - \omega^2)^2 + \gamma^2 \omega^2} \quad (\text{A.14})$$

Again, from Eqs. (A.11) and (A.12), we find that for non-magnetic materials

$$n = \left\{ \frac{1}{2} \left[(\epsilon_1^2 + \epsilon_2^2)^{\frac{1}{2}} + \epsilon_1 \right] \right\}^{\frac{1}{2}} \quad (\text{A.15})$$

$$k = \left\{ \frac{1}{2} \left[(\epsilon_1^2 + \epsilon_2^2)^{\frac{1}{2}} - \epsilon_1 \right] \right\}^{\frac{1}{2}} \quad (\text{A.16})$$

Summary

Equations (A.13) and (A.14) can be cast into dimensionless form. Then the susceptibility χ_e and conductivity in the neighborhood of a resonance frequency a are given by

$$\chi_e = \rho' \frac{1 - \nu^2}{(1 - \beta^2)^2 + \gamma'^2 \nu^2} \quad (\text{A.17})$$

and

$$(\sigma'/\nu) = - 2\pi \rho' \frac{\gamma' \nu}{(1 - \beta^2)^2 + \gamma'^2 \nu^2} \quad (\text{A.18})$$

where: ν' - is the measured frequency divided by a .

σ' - is the conductivity divided by a .

γ' - is called the width of the resonance.

ρ' - is given by $N e^2 / 4\pi^2 m a^2$ which is strength of the resonance.

And n, k can be written alternatively as

$$n^2 = \frac{1}{2} \left\{ \left[\epsilon^2 + 4(\sigma/\nu)^2 \right]^{\frac{1}{2}} + \epsilon \right\} \quad (\text{A.19})$$

$$k^2 = \frac{1}{2} \left\{ \left[\epsilon^2 + 4(\sigma/\nu)^2 \right]^{\frac{1}{2}} - \epsilon \right\} \quad (\text{A.20})$$

respectively, where

$$\epsilon = \epsilon_0 + 4\pi X_e \quad (\text{A.21})$$

A.1 Criteria in Determining the Parameters of the Lorentz Model

In order to obtain a theoretical fit of Eqs. (A.17) and (A.18) to experimental data, the use of the following principles can be helpful:

a. The initial choice of a should correspond to the steepest point on the long wavelength side of the reflection peak. Alternatively one should use Havelock's²¹ formula

$$a = \nu_{\max} \left(1 + \frac{n_+^2 - n_-^2}{6n_-^2 - 4} \right)^{\frac{1}{2}}$$

where: ν_{\max} is frequency at the reflectivity maximum, n_+ and n_- are the limiting indices of refraction at long and short wavelength, respectively.

b. The initial choice of ϵ_0 should give the observed limiting reflectivity at short wavelength according to the formula

$$R_{\nu \rightarrow \infty} = \left[\frac{\epsilon_0^{\frac{1}{2}} - 1}{\epsilon_0^{\frac{1}{2}} + 1} \right]^2$$

c. The initial choice of ρ should give the correct reflectivity on the long wavelength side of the band. It may be chosen to give the static electric constant according to the formulas (A.17) and (A.21) by taking the limit $\nu \rightarrow 0$.

d. The initial choice of f is governed by the maximum reflectivity. The approximation

$$\max R \sim 1 - (\delta/\pi\rho)$$

is rough but very useful.

Repeat with the process a to d , we can determine the values of these parameters by trial and error until the best fit is reached. Therefore, a theoretical spectrum of R_0 can be calculated. In fact, with very limiting cases the theoretical curve can succeed in aligning with the experimental curve.

APPENDIX B

REFLECTANCE AND PHASE-SHIFT DISPERSION RELATION (REF. 13, pp. 244-248)

The derivation of the phase-shift dispersion relation requires that the complex reflectivity amplitude $\hat{r}(\omega)$ be analytically continued into the complex frequency plane. Suppose that $\hat{r}(\omega)$ and hence $\text{In}\hat{r}(\omega)$ can be analytically continued into the upper half of the complex frequency plane, which we denote as domain D. A contour integral as indicated in Fig. B.1 then yields the desired dispersion relation.

The linear causal relation between $E_r(\omega)$ and $E_i(\omega)$ require analyticity of $r(\omega)$ in the domain D as shown by the following:

$$E_r(t) = \int_{-\infty}^{\infty} E_r(\omega) e^{-i\omega t} d\omega \quad (\text{B.1})$$

$$E_r(\omega) = \hat{r}(\omega) E_i(\omega) \quad (\text{B.2})$$

$$E_i(\omega) = \frac{1}{2\pi} \int_{-\infty}^{\infty} E_i(t') \exp(i\omega t') dt' \quad (\text{B.3})$$

By substituting Eqs. (B.2) and (B.3) into Eq. (B.1), it yields

$$E_r(t) = \int_{-\infty}^{\infty} G(t-t') E_i(t') dt' \quad (\text{B.4})$$

where:

$$G(t-t') = \frac{1}{2\pi} \int_{-\infty}^{\infty} \hat{r}(\omega) \exp[-i\omega(t-t')] d\omega \quad (\text{B.5})$$

which is the Fourier transform of $\hat{r}(\omega)$. If this transform can be inverted, then,

$$\hat{r}(\omega) = \int_{-\infty}^{\infty} G(T) e^{i\omega T} dT \quad ; \quad T \equiv t-t' \quad (\text{B.6})$$

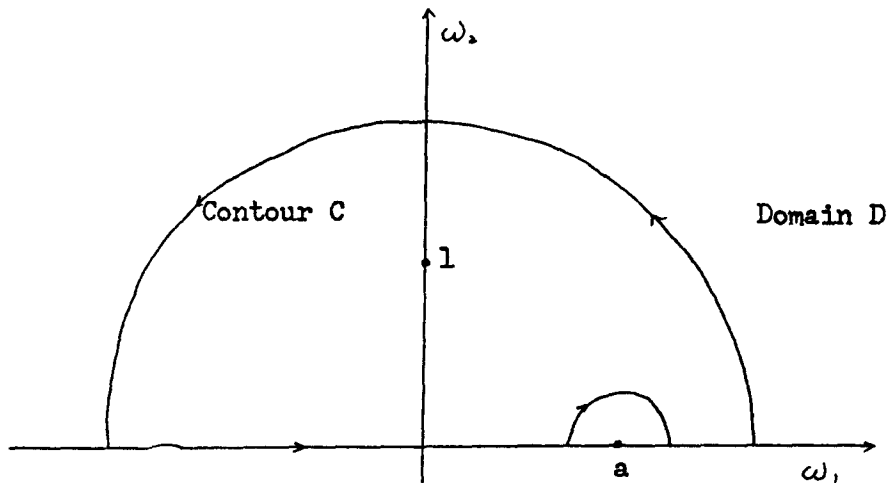


Fig. B.1 Integration Contour in Complex Frequency Plane
 $\hat{\omega} = \omega_1 + i\omega_2$

We want to derive a dispersion relation for $\text{In}\hat{r}(\hat{\omega})$, i.e., we want an integral relationship between $\text{In}\rho(\omega)$ and $\theta(\omega)$ so we can eliminate one of the unknowns in the equation

$$\text{In}\hat{r}(\omega) = \text{In}\rho(\omega) + i\theta(\omega) \quad (\text{B.7})$$

Therefore, a function other than $\text{In}\hat{r}(\hat{\omega})$ must be used for the integration over the contour shown in Fig. B.1. The function

$$f(\hat{\omega}) = (1+a\hat{\omega})\text{In}\hat{r}(\hat{\omega}) / [(1+\hat{\omega}^2)(a-\hat{\omega})] \quad (\text{B.8})$$

which can be integrated over the contour C satisfies the dispersion relation for $\text{In}\hat{r}(\hat{\omega})$.

Integrating Eq. (B.8), we have

$$\int_C \frac{(1+a\hat{\omega}) \text{In}[\hat{r}(\hat{\omega})]}{(1+\hat{\omega}^2)(a-\hat{\omega})} d\hat{\omega} = 2\pi i \sum_j \text{Res}(\hat{\omega}_j) \quad (\text{B.9})$$

This becomes

$$\text{p.v.} \int_{-\infty}^{\infty} \frac{(1+a\omega) \text{In}[\hat{r}(\hat{\omega})]}{(1+\omega^2)(a-\omega)} d\omega + i\pi \text{In}[\hat{r}(a)] = i\pi \text{In}[r(i)] \quad (\text{B.10})$$

since the only pole enclosed by the contour is $\hat{\omega}=i$, and the only pole on the contour is $\hat{\omega}=a$. From Eq. (B.6), we see $\hat{r}(i)$ is real, so is $\text{In}[\hat{r}(i)]$. Substituting Eq. (B.7) into (B.10) and equating the real parts of the resulting equation, we find

$$\phi(a) = \frac{1}{\pi} \text{p.v.} \int_{-\infty}^{\infty} \frac{1+a\omega}{1+\omega^2} \text{In}[\rho(\omega)] \frac{d\omega}{a-\omega} \quad (\text{B.11})$$

Using Eq. (B.2), we see that the complex conjugate of r is

$$\hat{r}^*(\omega) = E_r^*(\omega)/E_1^*(\omega) = E_r(-\omega)/E_1(-\omega) \quad (\text{B.12})$$

$$\text{so } \hat{r}^*(\omega) = \hat{r}(-\omega) \quad (\text{B.13})$$

$$\text{and } \text{In}[\rho(\omega)] = \text{In}[\rho(-\omega)] \quad (\text{B.14})$$

Considering integration of Eq. (B.11) over negative frequencies and using Eq. (B.14), we have

$$\frac{1}{\pi} \text{p.v.} \int_{-\infty}^0 \frac{(1+a\omega) \text{In}\rho(\omega)}{(1+\omega^2)(a-\omega)} d\omega = \frac{1}{\pi} \text{p.v.} \int_0^{\infty} \frac{(1-a\omega) \text{In}\rho(\omega)}{(1+\omega^2)(a+\omega)} d\omega \quad (\text{B.15})$$

Therefore, Eq. (B.11) becomes

$$\phi(a) = \frac{2a}{\pi} \text{p.v.} \int_0^{\infty} \frac{\text{In}\rho(\omega)}{a^2 - \omega^2} d\omega \quad (\text{B.16})$$

The singularity at $\omega=a$ is removed by subtracting

$$\frac{2a}{\pi} \text{p.v.} \int_0^{\infty} \frac{\text{In}\rho(a)}{a^2 - \omega^2} d\omega = 0 \quad (\text{B.17})$$

from Eq. (B.16). Then from the normal reflectivity relation

$$R(\omega) = \hat{r}(\omega)\hat{r}^*(\omega) = \rho^2 \quad (\text{B.18})$$

we obtain the desired dispersion relation

$$\phi(a) = \frac{a}{\pi} \int_0^{\infty} \frac{\text{In}[R(\omega)/R(a)]}{a^2 - \omega^2} d\omega \quad (\text{B.19})$$

There are several notable features of Eq. (B.19). A fractional error in reflectance, that is, a measurement which yields the same percentage error in reflectance at all frequencies, does not affect the phase. Such errors still yield the correct reflectance ratio $R(\omega)/R(a)$, and no effect to the phase at $\omega=a$.

APPENDIX C

NUMERICAL INTEGRATION OF THE PHASE-SHIFT EQUATION

The phase-shift equation,

$$\phi(a) = -\frac{2a}{\pi} \int_0^{\infty} \frac{\text{In } R_S^{\frac{1}{2}}(\omega)}{\omega^2 - a^2} d\omega \quad (\text{C.1})$$

can be written, after integrating by parts, as

$$\phi(a) = \frac{1}{\pi} \int_0^{\infty} \text{In} \left| \frac{\omega-a}{\omega+a} \right| \frac{d \text{In } R_S^{\frac{1}{2}}(\omega)}{d\omega} d\omega \quad (\text{C.2})$$

To integrate by part one more time, let us choose

$$u = d \text{In } R_S^{\frac{1}{2}}(\omega) / d\omega \quad ; \quad dv = \left\{ \text{In}|\omega-a| - \text{In}(\omega+a) \right\} d\omega$$

Then,

$$\begin{aligned} \phi(a) = \frac{1}{\pi} \left\{ \frac{d \text{In } R_S^{\frac{1}{2}}(\omega)}{d\omega} \left\{ (\omega-a) \text{In}|\omega-a| - (\omega+a) \text{In}(\omega+a) + 2a \right\} \right\}_0^{\infty} \\ - \int_0^{\infty} \left\{ (\omega-a) \text{In}|\omega-a| - (\omega+a) \text{In}(\omega+a) + 2a \right\} \frac{d^2 \text{In } R_S^{\frac{1}{2}}(\omega)}{d\omega^2} d\omega \quad (\text{C.3}) \end{aligned}$$

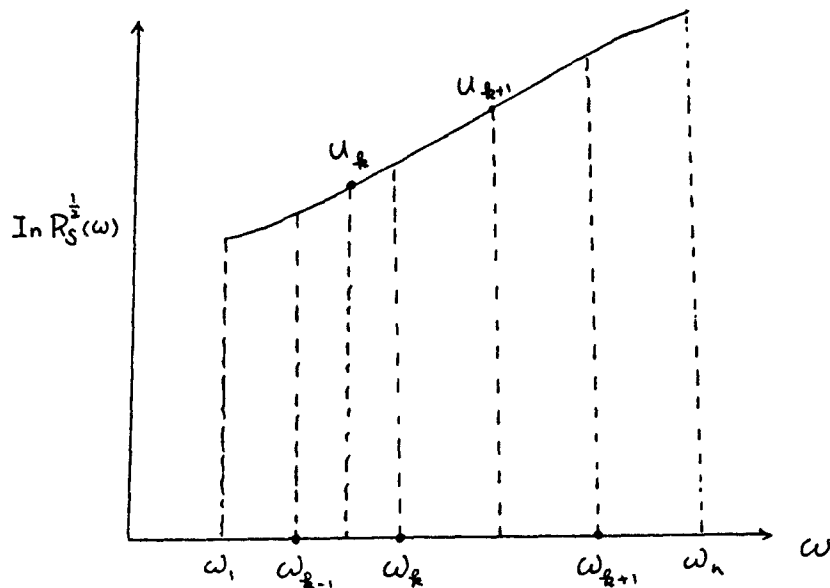
Because it is complex integral, we should use Cauchy principle value to simplify the integral to the desired form. Assume the slope of $R_S(\omega)$ is zero at $\omega=0$, and $\omega \rightarrow \infty$, the first term on the right hand side of Eq. (C.3) is then zero. Also for the last term in the integral,

$$\frac{2a}{\pi} \int_0^{\infty} \frac{d^2 \text{In } R_S^{\frac{1}{2}}(\omega)}{d\omega^2} d\omega = 0 \quad (\text{C.4})$$

Thus, it is only left with

$$\phi(a) = \frac{1}{\pi} \int_0^{\infty} \left\{ (\omega+a) \text{In}(\omega+a) + (a-\omega) \text{In}|\omega-a| \right\} \frac{d^2 \text{In } R_S^{\frac{1}{2}}(\omega)}{d\omega^2} d\omega \quad (\text{C.5})$$

In reference to the sketch, let



$$u_k = \frac{\ln R_S^{\frac{1}{2}}(\omega_k) - \ln R_S^{\frac{1}{2}}(\omega_{k-1})}{\omega_k - \omega_{k-1}} \quad \text{at } \frac{\omega_k + \omega_{k-1}}{2}$$

$$u_{k+1} = \frac{\ln R_S^{\frac{1}{2}}(\omega_{k+1}) - \ln R_S^{\frac{1}{2}}(\omega_k)}{\omega_{k+1} - \omega_k} \quad \text{at } \frac{\omega_k + \omega_{k+1}}{2}$$

$$\frac{d^2 \ln R_S^{\frac{1}{2}}(\omega)}{d\omega^2} d\omega = u_{k+1} - u_k = \gamma_k^{\circ}$$

Therefore,

$$\phi(a) = \frac{1}{\pi} \sum_{k=2}^{n-1} \gamma_k^{\circ} \left\{ (\omega+a) \ln(\omega+a) + (a-\omega) \ln|\omega-a| \right\} \quad (\text{C.6})$$

In using this formula it is essential to impose the first two values ($k=1, 2$) of $R_S(\omega)$ be the same. So are the last two values ($k=n-1, n$) because of the assumption imposed.

APPENDIX D

Optical Density of Single Crystal ThO_2
From Bates' Report (Ref. 16)

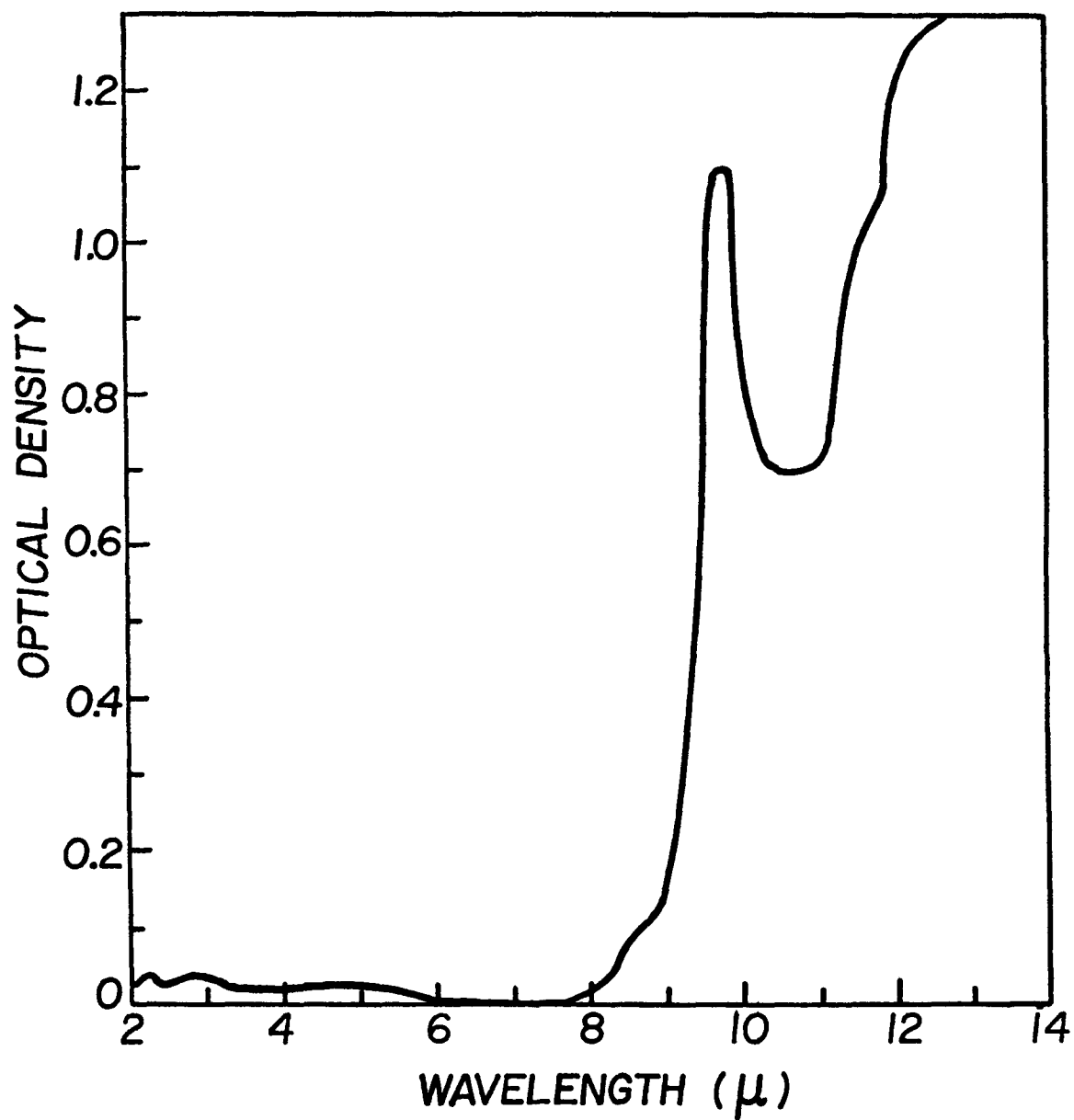


Fig. D.1 Infrared Absorption Spectrum of Single Crystal ThO₂
(Sample Thickness is 0.048 cm)

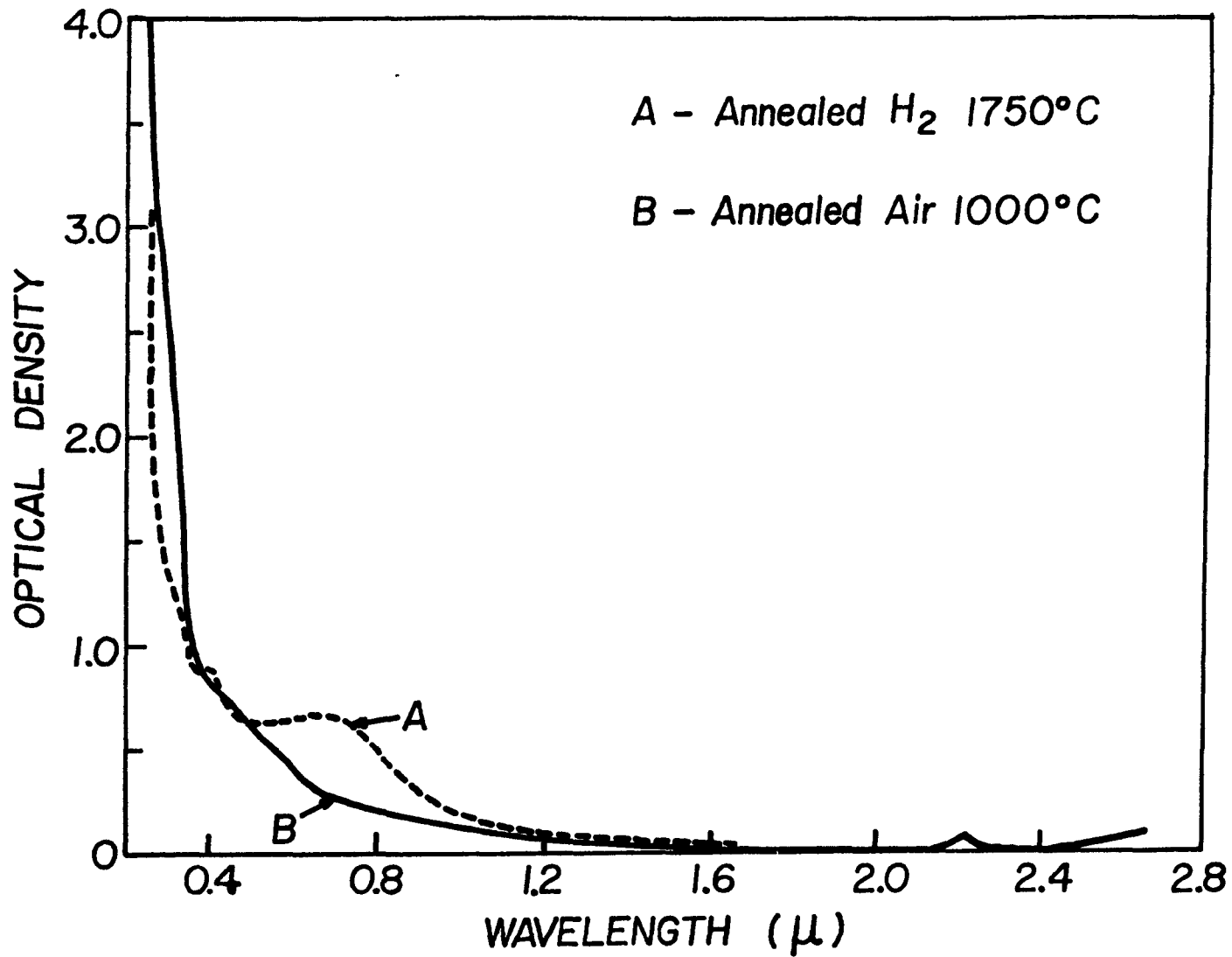


Fig. D.2 Ultraviolet, Visible, and Infrared Absorption Spectra of Single Crystal ThO₂

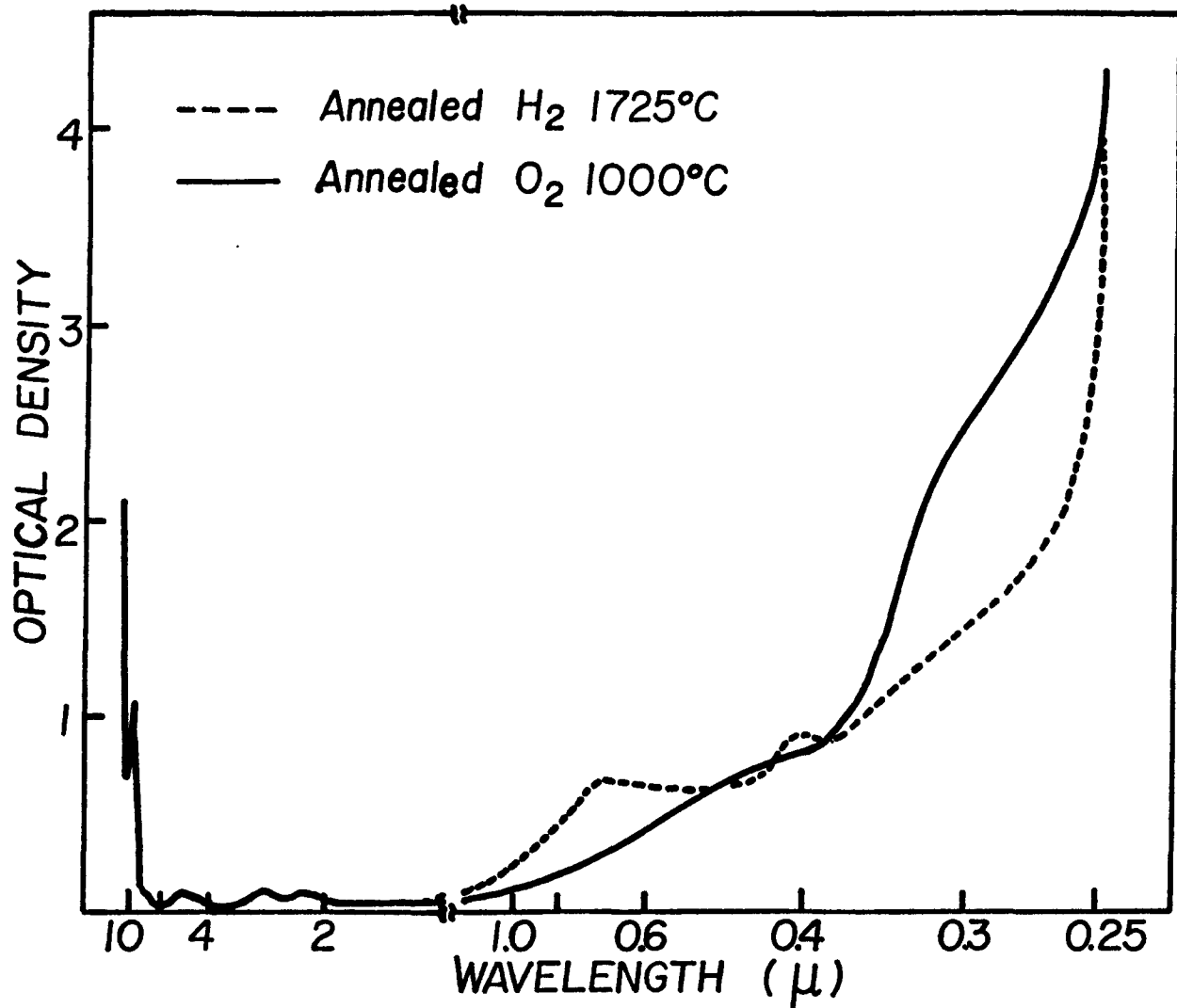


Fig. D.3 Ultraviolet, Visible, and Infrared Absorption Spectra of Single Crystal ThO₂

APPENDIX E

FORTRAN Program for the K-K Iterative Method

FORTRAN Implementation

Lorentz Oscillator:

Program Symbol	Definition
N	Number of segment
POL	Polarization
WI	Wavelength at short side (μ)
VO	Resonance wavelength (μ)
EO	Dielectric Constant at long wavelength
T	Strength of resonance
GARMA	Damping constant
THETA	Incidence angle (radian)
EPS	A positive number less than 10^{-5}
R	Reflectivity at normal incident angle
RO	Experimental reflectivity
ZN	Index of refraction
ZK	Extinction coefficient

Iterative K-K Analysis:

Program Symbol	Definition
M	Number of iteration
N	Number of Points
PHI	Incidence angle (radian)
POL	Polarization
RO(K)	Values of experimental reflectivity
THETA	Phase angle
X(K)	Wavelength (μ)

Index of Refraction from K-K Theory:

Program Symbol	Definition
EL	Dielectric constant at short wavelength (μ)
N	Number of Values
Y(K)	Values of wavelengths (μ)
ZN(K)	Values of index of refraction
ZK(K)	Values of extinction coefficient

APPENDIX F

FORTRAN Program for the Calculation of Surface Radiation
Properties of ThO₂ and the Related Reactor Materials

FORTRAN Implementation

Spectral Hemispherical Properties:

Program Symbol	Definition
WA	Resonance wavelength
Y1	Index of refraction at medium 1
Y2	Index of refraction at medium 2
Z1	Extinction coefficient at medium 1
Z2	Extinction coefficient at medium 2

Total Hemispherical Properties:

Program Symbol	Definition
EO(I)	Values of spectral hemispherical emittance
TE	Source temperature
YI(J)	Values of wavelength (μ)
VALUE	Total hemispherical emittance
VAL	Total hemispherical reflectance

BEFOR ASIZ
FOR 0077-0070/177-11112*03

MAIN PROGRAM

STORAGE UNITS: 007(1) 00724; DATA(0) 00711; BLANK COMMON(2) 00700

EXTERNAL REFERENCES (BLOCK, NAME)

0007 MINTP
0004 000
0005 000
0006 000
0007 000
0010 X00R
0011 N00J
0012 N00S
0013 N00PE

STORAGE ASSIGNMENT (BLOCK, TYPE, RELATIVE LOCATION, NAME)

0000	000044	100	0001	000004	1100	0011	000134	20L	0001	000155	70L	0001	000173	40L
0001	000174	70L	0000	R 000076	AA	0000	R 000025	AAA	0000	R 000022	E	0000	R 000034	03
0000	R 000000	70	0000	R 000073	PS	0000	R 000001	EPI	0000	R 000011	CARMA	0000	I 000017	I
0000	I 000000	N	0000	I 000012	NP1	0000	R 000037	P	0000	R 000007	FI	0000	R 000002	00L
0000	R 000035	0	0000	R 000075	050	0000	R 000033	R	0000	R 000043	R0	0000	R 000042	R2
0000	R 000041	000	0000	R 000071	STV	0000	R 000024	SIVS	0000	R 000010	T	0000	R 000004	THETA
0000	R 000010	V	0000	R 000016	VA	0000	R 000017	V0	0000	R 000005	V0	0000	R 000014	W
0000	R 000003	WI	0000	R 000000	X	0000	R 000040	Z	0000	R 000032	ZK	0000	R 000031	ZKK
0000	R 000030	ZN	0000	R 000027	ZNN									

00101	10	VEP00
00107	20	EPS=0.00000001
00104	30	POLE0
00105	40	WTE0
00106	50	THETA=1.0477
00107	60	V0=10.0
00110	70	0000.7
00111	80	PT=3.141592
00112	90	TE=0.267
00117	100	CARMA=0.000
00114	110	NP1=0.1
00115	120	DO=1.0*NP1
00120	130	WEWT=0.1*PLCAT(T-1)
00121	140	V=V0/W
00122	150	VA=1.-V**2
00123	160	V00CARMA*V
00124	170	Y=1+VA/(VA**2+V0**2)
00125	180	STV=0.*PI+T**P/(VA**2+V0**2)
00126	190	000004.*PI*X
00127	200	000000
00128	210	STV=STV**2

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000128

00131	02*	AA=DS+4.*ST	000137
00132	03*	AA=SQRT(AAA)	000137
00133	04*	ZN=1./T*(AA+)	000136
00134	05*	ZN=ABS(ZN)	000141
00135	06*	IF (ZN.LT.0.01) GO TO 10	000143
00137	07*	DN=1./T*(DN)	000146
00140	08*	GO TO 30	000150
00141	09*	10 ZN=0.	000154
00142	10*	20 ZK=0.5*(AA+)	000155
00143	11*	ZK=ABS(ZK)	000160
00144	12*	IF (ZK.LT.0.01) GO TO 40	000162
00145	13*	ZK=SQRT(ZK)	000165
00147	14*	GO TO 50	000171
00150	15*	40 ZK=0.	000177
00151	16*	50 R=((ZN-1.)**2+ZK**2)/(LN+1.)**2+ZK**2)	000174
00152	17*	BB=ZN**2-ZK**2-(SIN(THETA))**2	000212
00153	18*	DD=(BB*(BB**2+4.*(LN*ZK)**2))**.5/2.	000221
00154	19*	DD=DT(DD)	000235
00155	20*	P=(ZN*ZK)/D	000250
00156	21*	Z=((0-COS(THETA))**2+P**2)/(0+COS(THETA))**2+P**2)	000255
00157	22*	RFS=((0-SIN(THETA)*TAN(THETA))**2+P**2)/(0+SIN(THETA)*TAN(THETA)	000264
00157	23*)**2+P**2)	000264
00158	24*	DP=DD*7	000272
00159	25*	ZC=C.(Z*(1.+DCL)+D)*(1.+DCL)	000274
00162	26*	WRITE(6,10)W, LN, ZK, Z, ZC	000332
00172	27*	10 FORMAT(1H0.5X, 'W= ', F10.5, 'X, 'Z= ', F10.5, 'Y, 'ZK= ', F10.5, 'Y,	000317
00172	28*	'ZC= ', F10.5, 'X, 'ZE= ', F10.5, 'X, 'DC= ', F10.5)	000317
00173	29*	2 CONTINUE	000317
00175	30*	STOP	000317
00176	31*	END	000323

END OF COMPLETION: NO DIAGNOSTICS.

FOR, ASIZ
 FOR JF37-15/79-19:25:05

MAIN PROGRAM

STORAGE USED: CODE(1) 00000; DATA(0) 002703; BLANK COMMON(2) 00000

EXTERNAL REFERENCES (BLOCK, NAME)

0003 NINT91
 0004 NRDUS
 0005 COS
 0006 SIN
 0007 TAN
 0010 N1029
 0011 N1019
 0012 NRDUS
 0013 ALOG
 0014 SGPT
 0015 XPPR
 0016 NSTC03

STORAGE ASSIGNMENT (BLOCK, TYPE, RELATIVE LOCATION, NAME)

0007	002621	10F	0001	000060	114C	0001	000065	120G	0001	000077	127C	0001	000104	133G
0007	002612	141G	0001	000120	147G	0001	000125	155G	0001	000133	160G	0001	000174	171G
0001	000200	175G	0001	000247	2L	0000	002624	20F	0001	000526	25L	0001	000446	253G
0001	000470	26L	0001	000564	27L	0001	000504	276G	0001	000425	30L	0001	000542	321G
0001	000572	37L	0001	000254	4L	0000	002626	40F	0001	000455	42L	0001	000466	44L
0007	002670	57F	0001	000513	52L	0001	000524	54L	0001	000551	57L	0001	000562	59L
0001	000441	70L	0001	000477	71L	0001	000535	74L	0000	R 000000	A	0000	R 002607	AMW
0007	R 002610	APW	0000	R 000144	B	0000	R 000310	C	0000	R 000454	D	0000	I 002603	I
0007	I 002620	II	0000	I 002577	ITER	0000	I 002605	J	0000	I 002575	JJ	0000	I 002574	K
0007	I 002576	KK	0000	I 002606	L	0000	I 002671	LM1	0000	I 002570	M	0000	I 002571	N
0007	I 002672	MM	0000	R 000620	P	0000	R 002572	PHI	0000	R 002606	PI	0000	R 002614	PN
0007	R 002573	POL	0000	R 002617	PQ	0000	R 002610	W	0000	R 001604	R	0000	R 002612	RD
0007	R 002114	RNEW	0000	R 002260	RO	0000	R 002615	RPS	0000	R 002616	RS	0000	R 002611	THESUP
0007	R 001170	THETA	0000	R 002604	THETA	0000	R 002764	U	0000	R 001750	W	0000	R 002424	X
0007	R 001440	ZK	0000	R 001274	ZN									

00101 1* DIMENSION A(100),P(100),C(100),D(100),F(100),U(100),THETA(100),
 00101 2* 1Z(100),ZK(100),R(100),W(100),RNEW(100),RO(100),X(100)
 00103 3* READ(5,10)M,N,PHI,POL
 00111 4* 10 FORMAT(13,13,F10.4,F6.5)
 00112 5* READ(5,20)(X(K),K=1,N),(RO(K),K=1,N)
 00124 6* 20 FORMAT(8F10.4)
 00125 7* WRITE(6,40)(X(K),K=1,N),(FO(K),K=1,N)
 00137 8* 40 FORMAT(2LX,5F10.4)
 00140 9* DO 111 JJ=1,N
 00143 10* K(JJ)=1000./X(JJ)
 00144 11* 111 CONTINUE
 00145 12* DO 21 KK=1,N

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 000120

00151	13*	R(KK)=RC(KK)	000120
00152	14*	31 CONTINUE	000125
00154	15*	DO 33 I=L,N	000133
00157	16*	DO 80 L=L,N	000133
00162	17*	LM1=L-1	000137
00163	18*	A(L)=(1./2.)*(ALOG(R(L))-ALOG(R(LM1)))	000155
00164	19*	L(L)=w(L)-W(LM1)	000160
00165	20*	EU U(L)=A(L)/R(L)	000164
00167	21*	NM1=N-1	000166
00170	22*	DO 30 I=L,NM1	000174
00173	23*	THETA=0.	000200
00174	24*	DO 4 J=2,NM1	000203
00177	25*	P(J)=U(J+1)-U(J)	000205
00200	26*	FI=3.141592	000211
00201	27*	AMW=ABS(w(I)-w(J))	000215
00202	28*	APW=ABS(w(I)+w(J))	000224
00203	29*	D(J)=(w(I)+w(J))*(ALOG(APW))	000227
00204	30*	IF(I-J)3,2,3	000237
00207	31*	3 C(J)=(w(I)-w(J))*(ALOG(APW))	000245
00210	32*	THESUB=(1./PI)*P(J)+C(J)+D(J)	000247
00211	33*	GO TO 4	000254
00212	34*	2 THESUB=(1./PI)*P(J)+C(J)	000260
00213	35*	4 THETA=THETA+THESUB	000261
00215	36*	THETA(I)=THETA	000302
00216	37*	RD=1.+R(I)-2.*SQRT(R(I))*COS(THETA(I))	000307
00217	38*	G=(1.-R(I))*COS(PHI)/FD	000327
00220	39*	PN=2.*SQRT(R(I))*SIN(THETA(I))*COS(PHI)/RD	000327
00221	40*	RPS=((Q-SIN(PHI)*TAN(PHI))*2+PN**2)/((Q+SIN(PHI)*TAN(PHI))*2+	000343
00221	41*	1FN**2)	000351
00222	42*	RS=2.*RC(I)/((1.+POL)*RPS+(1.-POL))	000356
00223	43*	PG=Q**2-PN**2+(SIN(PHI))**2	000402
00224	44*	ZN(I)=((PG+(PG**2+4.*(FN*G)**2)**0.5)/2.)*0.5	000405
00225	45*	ZK(I)=PN*J/ZN(I)	000407
00226	46*	RNEW(I)=RS	000412
00227	47*	IF(ITER-1)30,61,61	000427
00232	48*	61 WRITE(6,50)ITER,X(I),RNEW(I),ZN(I),ZK(I)	000427
00241	49*	50 FORMAT(5X,'M=',13,10X,'X=',F10.5,10X,'RNEW=',F10.5,10X,'ZN=',	000427
00241	50*	\$F10.5,10X,'ZK=',F10.5)	000427
00242	51*	30 CONTINUE	000427
00244	52*	IF(ITER-7)70,71,72	000433
00247	53*	72 IF(ITER-13)71,74,74	000441
00252	54*	70 DO 20 II=2,NM1	000446
00255	55*	IF(X(II)-10.09)41,41,42	000451
00260	56*	41 R(II)=RNEW(II)	000453
00261	57*	GO TO 26	000455
00262	58*	42 IF(X(II)-15.01)43,44,44	000460
00265	59*	43 R(II)=(RNEW(II)+R(II))/2.	000464
00266	60*	GO TO 26	000466
00267	61*	44 R(II)=RNEW(II)	000471
00270	62*	20 CONTINUE	000473
00272	63*	R(1)=R(2)	000475
00273	64*	R(N)=P(NM1)	000477
00274	65*	GO TO 33	000504
00275	66*	71 DO 25 II=2,NM1	000507
00275	67*	IF(Y(II)-10.19)51,51,52	000511
00303	68*	51 F(II)=RNEW(II)	
00304	69*	GO TO 25	

00309	70*	52	IF(X(II)-13.61)53,54,54	000513
00310	71*	53	R(II)=(RNEW(II)+R(II))/2.	000516
00311	72*		GO TO 25	000522
00312	73*	54	R(II)=RNEW(II)	000524
00313	74*	25	CONTINUE	000527
00315	75*		R(1)=R(2)	000527
00316	76*		R(N)=R(NM1)	000531
00317	77*		GO TO 33	000533
00320	78*	74	DO 27 II=2,NM1	000535
00323	79*		IF(X(II)-10.29)56,56,57	000542
00326	80*	56	R(II)=RNEW(II)	000545
00327	81*		GO TO 27	000547
00330	82*	57	IF(X(II)-13.51)58,59,59	000551
00333	83*	58	R(II)=(RNEW(II)+R(II))/2.	000554
00334	84*		GO TO 27	000560
00335	85*	59	R(II)=RNEW(II)	000562
00336	86*	27	CONTINUE	000565
00340	87*		R(1)=R(2)	000565
00341	88*		R(N)=R(NM1)	000567
00342	89*	33	CONTINUE	000573
00344	90*		STOP	000573
00345	91*		END	000577

END OF COMPILATION: NO DIAGNOSTICS.

2.0000	2.2000	3.0000	3.5000	4.0000	4.5000	5.0000	5.5000
2.0000	7.0000	7.5000	8.0000	8.2000	8.4000	8.6000	8.8000
9.0000	9.2000	9.4000	9.6000	9.8000	9.9000	10.0000	10.1000
10.2000	10.3000	10.4000	10.5000	10.6000	10.7000	10.8000	11.0000
11.2000	11.4000	11.6000	12.0000	12.2000	12.4000	12.5000	12.6000
12.7000	12.8000	12.9000	13.0000	13.1000	13.2000	13.3000	13.4000
13.5000	13.6000	13.8000	14.0000	14.5000	15.0000	15.5000	16.0000
16.5000	17.0000	17.9000	18.0000	.2282	.2282	.2259	.2242
.2222	.2197	.2167	.2131	.2087	.1963	.1874	.1757
.1698	.1632	.1554	.1463	.1353	.1219	.1049	.0823
.0499	.0279	.0179	.0194	.7104	.8367	.8858	.9127
.9278	.9304	.9459	.9556	.9613	.9647	.9666	.9659
.9623	.9517	.9367	.9131	.7225	.6471	.5984	.5631
.5358	.5138	.4906	.4803	.4671	.4556	.4366	.4214
.3896	.3704	.3621	.3520	.3442	.3378	.3283	.3283

M= 1	X= 2.00000	RNEW= .45595	ZN= 1.65797	ZK= -.02534
M= 1	X= 3.00000	RNEW= .45124	ZN= 1.65001	ZK= -.03255
M= 1	X= 3.50000	RNEW= .44774	ZN= 1.64409	ZK= -.03778
M= 1	X= 4.00000	RNEW= .44361	ZN= 1.63710	ZK= -.04352
M= 1	X= 4.50000	RNEW= .43945	ZN= 1.62845	ZK= -.04926
M= 1	X= 5.00000	RNEW= .43225	ZN= 1.61910	ZK= -.05556
M= 1	X= 5.50000	RNEW= .42479	ZN= 1.60576	ZK= -.06211
M= 1	X= 6.00000	RNEW= .41566	ZN= 1.59796	ZK= -.06811
M= 1	X= 7.00000	RNEW= .38974	ZN= 1.54884	ZK= -.08710
M= 1	X= 7.50000	RNEW= .37111	ZN= 1.51989	ZK= -.09478
M= 1	X= 8.00000	RNEW= .34651	ZN= 1.48209	ZK= -.10473
M= 1	X= 8.20000	RNEW= .33412	ZN= 1.46379	ZK= -.10749
M= 1	X= 8.40000	RNEW= .32024	ZN= 1.44341	ZK= -.11061
M= 1	X= 8.60000	RNEW= .30325	ZN= 1.41979	ZK= -.11335
M= 1	X= 8.80000	RNEW= .28476	ZN= 1.39279	ZK= -.11556
M= 1	X= 9.00000	RNEW= .26175	ZN= 1.36102	ZK= -.11675
M= 1	X= 9.20000	RNEW= .23307	ZN= 1.32354	ZK= -.11627
M= 1	X= 9.40000	RNEW= .19885	ZN= 1.27809	ZK= -.11252
M= 1	X= 9.60000	RNEW= .15311	ZN= 1.22107	ZK= -.10266
M= 1	X= 9.80000	RNEW= .08950	ZN= 1.13987	ZK= -.08597
M= 1	X= 9.90000	RNEW= .04849	ZN= 1.08882	ZK= -.06382
M= 1	X= 10.00000	RNEW= .00781	ZN= .99576	ZK= -.03464
M= 1	X= 10.10000	RNEW= .22068	ZN= .88369	ZK= -.04380
M= 1	X= 10.20000	RNEW= .77035	ZN= .78213	ZK= -.03277
M= 1	X= 10.30000	RNEW= .88885	ZN= .62567	ZK= -.05321
M= 1	X= 10.40000	RNEW= .93366	ZN= .40985	ZK= -.09640
M= 1	X= 10.50000	RNEW= .95662	ZN= .13996	ZK= -.33410
M= 1	X= 10.60000	RNEW= .97029	ZN= .08877	ZK= -.61557
M= 1	X= 10.70000	RNEW= .97938	ZN= .07757	ZK= -.83087
M= 1	X= 10.80000	RNEW= .93581	ZN= .07453	ZK= -1.02300
M= 1	X= 11.00000	RNEW= .99430	ZN= .07844	ZK= -1.38934
M= 1	X= 11.20000	RNEW= .99974	ZN= .09138	ZK= -1.77736
M= 1	X= 11.40000	RNEW= 1.00367	ZN= .11525	ZK= -2.22944
M= 1	X= 11.60000	RNEW= 1.00621	ZN= .15818	ZK= -2.80717
M= 1	X= 12.00000	RNEW= 1.01257	ZN= .45779	ZK= -4.98073
M= 1	X= 12.20000	RNEW= 1.01616	ZN= 1.29569	ZK= -8.02623
M= 1	X= 12.40000	RNEW= 1.02213	ZN= 15.50962	ZK= -19.63038
M= 1	X= 12.50000	RNEW= 1.02947	ZN= 16.06409	ZK= 15.07863
M= 1	X= 12.60000	RNEW= 1.04752	ZN= 4.66971	ZK= 6.39406
M= 1	X= 12.70000	RNEW= 1.04300	ZN= 7.65978	ZK= 2.97209
M= 1	X= 12.80000	RNEW= 1.01577	ZN= 7.42097	ZK= 1.99679
M= 1	X= 12.90000	RNEW= .98719	ZN= 7.23372	ZK= 1.56092

M= 1	X= 13.00000	RNEW= .96062	ZN= 3.06707	ZK= 1.19783
M= 1	X= 13.10000	RNEW= .93652	ZN= 2.96100	ZK= .99258
M= 1	X= 13.20000	RNEW= .91480	ZN= 2.86767	ZK= .84403
M= 1	X= 13.30000	RNEW= .89526	ZN= 2.77707	ZK= .73147
M= 1	X= 13.40000	RNEW= .87772	ZN= 2.70669	ZK= .64368
M= 1	X= 13.50000	RNEW= .86176	ZN= 2.64527	ZK= .57301
M= 1	X= 13.60000	RNEW= .84726	ZN= 2.59335	ZK= .51428
M= 1	X= 13.70000	RNEW= .82159	ZN= 2.55077	ZK= .42667
M= 1	X= 14.00000	RNEW= .80065	ZN= 2.43612	ZK= .36319
M= 1	X= 14.50000	RNEW= .75295	ZN= 2.29469	ZK= .24103
M= 1	X= 15.00000	RNEW= .73056	ZN= 2.27730	ZK= .17175
M= 1	X= 15.50000	RNEW= .70858	ZN= 2.17876	ZK= .13130
M= 1	X= 16.00000	RNEW= .69147	ZN= 2.13551	ZK= .09787
M= 1	X= 16.50000	RNEW= .67207	ZN= 2.10245	ZK= .07094
M= 1	X= 17.00000	RNEW= .66667	ZN= 2.07541	ZK= .04795
M= 1	X= 17.90000	RNEW= .64927	ZN= 2.03521	ZK= .03423
M= 2	X= 2.00000	RNEW= .42654	ZN= 2.72767	ZK= -.01172
M= 2	X= 3.00000	RNEW= .42348	ZN= 2.68958	ZK= -.00627
M= 2	X= 3.50000	RNEW= .42119	ZN= 2.66686	ZK= -.00418
M= 2	X= 4.00000	RNEW= .41847	ZN= 2.64044	ZK= -.00305
M= 2	X= 4.50000	RNEW= .41501	ZN= 2.60794	ZK= -.00134
M= 2	X= 5.00000	RNEW= .41080	ZN= 2.56965	ZK= -.00047
M= 2	X= 5.50000	RNEW= .40564	ZN= 2.52472	ZK= .00119
M= 2	X= 6.00000	RNEW= .39919	ZN= 2.47133	ZK= .00653
M= 2	X= 7.00000	RNEW= .38319	ZN= 2.32962	ZK= -.00618
M= 2	X= 7.50000	RNEW= .36579	ZN= 2.23348	ZK= -.00430
M= 2	X= 8.00000	RNEW= .34594	ZN= 2.11622	ZK= -.01140
M= 2	X= 8.50000	RNEW= .33556	ZN= 2.06056	ZK= -.01116
M= 2	X= 8.60000	RNEW= .32366	ZN= 2.00062	ZK= -.01278
M= 2	X= 8.80000	RNEW= .30925	ZN= 1.92296	ZK= -.01513
M= 2	X= 8.90000	RNEW= .29198	ZN= 1.85806	ZK= -.01810
M= 2	X= 9.00000	RNEW= .27757	ZN= 1.77295	ZK= -.02181
M= 2	X= 9.20000	RNEW= .24366	ZN= 1.67464	ZK= -.02601
M= 2	X= 9.40000	RNEW= .20874	ZN= 1.56497	ZK= -.02972
M= 2	X= 9.60000	RNEW= .16158	ZN= 1.43233	ZK= -.03164
M= 2	X= 9.80000	RNEW= .09435	ZN= 1.26367	ZK= -.04439
M= 2	X= 9.90000	RNEW= .05965	ZN= 1.16212	ZK= -.03614
M= 2	X= 10.00000	RNEW= .00796	ZN= 1.01134	ZK= -.04532
M= 2	X= 10.10000	RNEW= .22479	ZN= .86704	ZK= -.07838
M= 2	X= 10.20000	RNEW= .77614	ZN= .70281	ZK= -.05701
M= 2	X= 10.30000	RNEW= .89191	ZN= .39927	ZK= -.12495
M= 2	X= 10.40000	RNEW= .93327	ZN= .11936	ZK= -.47189
M= 2	X= 10.50000	RNEW= .95312	ZN= .06121	ZK= -.76726
M= 2	X= 10.60000	RNEW= .96408	ZN= .06891	ZK= -.99384
M= 2	X= 10.70000	RNEW= .97075	ZN= .06279	ZK= -1.19145
M= 2	X= 10.80000	RNEW= .97503	ZN= .05934	ZK= -1.37461
M= 2	X= 11.00000	RNEW= .97977	ZN= .05627	ZK= -1.72078
M= 2	X= 11.20000	RNEW= .98167	ZN= .05624	ZK= -2.06676
M= 2	X= 11.40000	RNEW= .98260	ZN= .05857	ZK= -2.43292
M= 2	X= 11.60000	RNEW= .98250	ZN= .06753	ZK= -2.94240
M= 2	X= 12.00000	RNEW= .97989	ZN= .09461	ZK= -3.94066
M= 2	X= 12.20000	RNEW= .97674	ZN= .13365	ZK= -4.82477
M= 2	X= 12.40000	RNEW= .96965	ZN= .27975	ZK= -6.40517
M= 2	X= 12.50000	RNEW= .96045	ZN= .60290	ZK= -8.09271
M= 2	X= 12.60000	RNEW= .92141	ZN= 3.06891	ZK= -11.11707
M= 2	X= 12.70000	RNEW= .85550	ZN= 8.74957	ZK= -7.96124
M= 2	X= 12.80000	RNEW= .82199	ZN= 8.48655	ZK= -4.46860

M = 39	X = 4.50000	RNEW = .41941	ZN = 2.49268	ZK = -.02693
M = 39	X = 5.00000	RNEW = .41469	ZN = 2.46539	ZK = -.02878
M = 39	X = 5.50000	RNEW = .40898	ZN = 2.43291	ZK = -.02994
M = 39	X = 6.00000	RNEW = .40193	ZN = 2.39376	ZK = -.02764
M = 39	X = 7.00000	RNEW = .38152	ZN = 2.28507	ZK = -.03891
M = 39	X = 7.50000	RNEW = .36642	ZN = 2.20965	ZK = -.03600
M = 39	X = 8.00000	RNEW = .34597	ZN = 2.11289	ZK = -.03839
M = 39	X = 8.20000	RNEW = .33541	ZN = 2.06548	ZK = -.03574
M = 39	X = 8.40000	RNEW = .32339	ZN = 2.01330	ZK = -.03412
M = 39	X = 8.60000	RNEW = .30893	ZN = 1.95297	ZK = -.03246
M = 39	X = 8.80000	RNEW = .29169	ZN = 1.88427	ZK = -.03056
M = 39	X = 9.00000	RNEW = .27038	ZN = 1.80388	ZK = -.02872
M = 39	X = 9.20000	RNEW = .24376	ZN = 1.70988	ZK = -.02677
M = 39	X = 9.40000	RNEW = .20913	ZN = 1.59704	ZK = -.02447
M = 39	X = 9.60000	RNEW = .16216	ZN = 1.45829	ZK = -.02204
M = 39	X = 9.80000	RNEW = .09480	ZN = 1.27856	ZK = -.03650
M = 39	X = 9.90000	RNEW = .05083	ZN = 1.16944	ZK = -.03089
M = 39	X = 10.00000	RNEW = .00796	ZN = 1.01199	ZK = -.04573
M = 39	X = 10.10000	RNEW = .22203	ZN = .86037	ZK = -.07995
M = 39	X = 10.20000	RNEW = .77031	ZN = .67862	ZK = -.05773
M = 39	X = 10.30000	RNEW = .98476	ZN = .31118	ZK = -.15100
M = 39	X = 10.40000	RNEW = .92634	ZN = .09151	ZK = -.57525
M = 39	X = 10.50000	RNEW = .94608	ZN = .06754	ZK = -.86850
M = 39	X = 10.60000	RNEW = .95868	ZN = .05948	ZK = -1.10378
M = 39	X = 10.70000	RNEW = .96627	ZN = .05605	ZK = -1.31424
M = 39	X = 10.80000	RNEW = .97146	ZN = .05489	ZK = -1.51300
M = 39	X = 11.00000	RNEW = .97792	ZN = .05651	ZK = -1.89808
M = 39	X = 11.20000	RNEW = .98160	ZN = .06220	ZK = -2.29439
M = 39	X = 11.40000	RNEW = .98380	ZN = .07229	ZK = -2.72666
M = 39	X = 11.60000	RNEW = .98511	ZN = .08865	ZK = -3.22607
M = 39	X = 12.00000	RNEW = .98543	ZN = .16894	ZK = -4.64876
M = 39	X = 12.20000	RNEW = .98415	ZN = .28838	ZK = -5.86438
M = 39	X = 12.40000	RNEW = .97989	ZN = .71742	ZK = -8.29227
M = 39	X = 12.50000	RNEW = .97335	ZN = 1.74089	ZK = -11.15841
M = 39	X = 12.60000	RNEW = .94111	ZN = 8.59705	ZK = -14.43840
M = 39	X = 12.70000	RNEW = .87067	ZN = 12.19541	ZK = -5.24964
M = 39	X = 12.80000	RNEW = .82721	ZN = 9.86012	ZK = -2.65879
M = 39	X = 12.90000	RNEW = .79646	ZN = 8.46125	ZK = -1.77450
M = 39	X = 13.00000	RNEW = .77263	ZN = 7.56868	ZK = -1.34203
M = 39	X = 13.10000	RNEW = .75320	ZN = 6.94576	ZK = -1.08846
M = 39	X = 13.20000	RNEW = .73684	ZN = 6.48278	ZK = -.92332
M = 39	X = 13.30000	RNEW = .72279	ZN = 6.12374	ZK = -.80900
M = 39	X = 13.40000	RNEW = .71059	ZN = 5.83784	ZK = -.72490
M = 39	X = 13.50000	RNEW = .69977	ZN = 5.60214	ZK = -.65998
M = 39	X = 13.60000	RNEW = .69010	ZN = 5.40407	ZK = -.61061
M = 39	X = 13.80000	RNEW = .67362	ZN = 5.09353	ZK = -.53087
M = 39	X = 14.00000	RNEW = .65994	ZN = 4.85817	ZK = -.47129
M = 39	X = 14.50000	RNEW = .62977	ZN = 4.36937	ZK = -.40695
M = 39	X = 15.00000	RNEW = .61553	ZN = 4.18899	ZK = -.40439
M = 39	X = 15.50000	RNEW = .60178	ZN = 4.01561	ZK = -.37028
M = 39	X = 16.00000	RNEW = .59102	ZN = 3.88530	ZK = -.35887
M = 39	X = 16.50000	RNEW = .58251	ZN = 3.78589	ZK = -.35599
M = 39	X = 17.00000	RNEW = .57539	ZN = 3.70469	ZK = -.35878
M = 39	X = 17.90000	RNEW = .56455	ZN = 3.57793	ZK = -.39656
M = 40	X = 2.20000	RNEW = .43253	ZN = 2.57114	ZK = -.02108
M = 40	X = 3.00000	RNEW = .42901	ZN = 2.54977	ZK = -.02197
M = 40	X = 3.50000	RNEW = .42630	ZN = 2.53403	ZK = -.02352

M= 40	X= 4.00000	RNEW= .42330	ZN= 2.51558	ZK= -.02545
M= 40	X= 4.50000	RNEW= .41941	ZN= 2.49268	ZK= -.02697
M= 40	X= 5.00000	RNEW= .41469	ZN= 2.46538	ZK= -.02886
M= 40	X= 5.50000	RNEW= .40898	ZN= 2.43292	ZK= -.02982
M= 40	X= 6.00000	RNEW= .40193	ZN= 2.39376	ZK= -.02784
M= 40	X= 7.00000	RNEW= .38152	ZN= 2.28507	ZK= -.03987
M= 40	X= 7.50000	RNEW= .36642	ZN= 2.20965	ZK= -.03604
M= 40	X= 8.00000	RNEW= .34597	ZN= 2.11260	ZK= -.03840
M= 40	X= 8.20000	RNEW= .33541	ZN= 2.06548	ZK= -.03567
M= 40	X= 8.40000	RNEW= .32339	ZN= 2.01330	ZK= -.03411
M= 40	X= 8.60000	RNEW= .30893	ZN= 1.95293	ZK= -.03247
M= 40	X= 8.80000	RNEW= .29169	ZN= 1.88427	ZK= -.03059
M= 40	X= 9.00000	RNEW= .27038	ZN= 1.80368	ZK= -.02873
M= 40	X= 9.20000	RNEW= .24376	ZN= 1.70988	ZK= -.02676
M= 40	X= 9.40000	RNEW= .20913	ZN= 1.59704	ZK= -.02447
M= 40	X= 9.60000	RNEW= .16216	ZN= 1.45829	ZK= -.02203
M= 40	X= 9.80000	RNEW= .09480	ZN= 1.27858	ZK= -.03650
M= 40	X= 9.90000	RNEW= .05083	ZN= 1.16945	ZK= -.03089
M= 40	X= 10.00000	RNEW= .00796	ZN= 1.01199	ZK= -.04573
M= 40	X= 10.10000	RNEW= .22202	ZN= .86037	ZK= -.07994
M= 40	X= 10.20000	RNEW= .77031	ZN= .67862	ZK= -.05773
M= 40	X= 10.30000	RNEW= .88476	ZN= .31120	ZK= -.15099
M= 40	X= 10.40000	RNEW= .92634	ZN= .09151	ZK= -.57523
M= 40	X= 10.50000	RNEW= .94628	ZN= .06754	ZK= -.86848
M= 40	X= 10.60000	RNEW= .95868	ZN= .05048	ZK= -1.10350
M= 40	X= 10.70000	RNEW= .96627	ZN= .05605	ZK= -1.31426
M= 40	X= 10.80000	RNEW= .97145	ZN= .05489	ZK= -1.51295
M= 40	X= 11.00000	RNEW= .97792	ZN= .05651	ZK= -1.89807
M= 40	X= 11.20000	RNEW= .98160	ZN= .06220	ZK= -2.29434
M= 40	X= 11.40000	RNEW= .98380	ZN= .07229	ZK= -2.72672
M= 40	X= 11.60000	RNEW= .98511	ZN= .08865	ZK= -3.22610
M= 40	X= 12.00000	RNEW= .98543	ZN= .16893	ZK= -4.64866
M= 40	X= 12.20000	RNEW= .98415	ZN= .28935	ZK= -5.88399
M= 40	X= 12.40000	RNEW= .97989	ZN= .71735	ZK= -8.29175
M= 40	X= 12.50000	RNEW= .97335	ZN= 1.74071	ZK= -11.15813
M= 40	X= 12.60000	RNEW= .94111	ZN= 8.59324	ZK= -14.43625
M= 40	X= 12.70000	RNEW= .87067	ZN= 12.19351	ZK= -5.25138
M= 40	X= 12.80000	RNEW= .82721	ZN= 9.85938	ZK= -2.66004
M= 40	X= 12.90000	RNEW= .79646	ZN= 8.46114	ZK= -1.77478
M= 40	X= 13.00000	RNEW= .77263	ZN= 7.56850	ZK= -1.34254
M= 40	X= 13.10000	RNEW= .75320	ZN= 6.94561	ZK= -1.08887
M= 40	X= 13.20000	RNEW= .73684	ZN= 6.48271	ZK= -.92357
M= 40	X= 13.30000	RNEW= .72279	ZN= 6.12367	ZK= -.80922
M= 40	X= 13.40000	RNEW= .71059	ZN= 5.83784	ZK= -.72496
M= 40	X= 13.50000	RNEW= .69977	ZN= 5.60214	ZK= -.65999
M= 40	X= 13.60000	RNEW= .69010	ZN= 5.40403	ZK= -.61082
M= 40	X= 13.80000	RNEW= .67362	ZN= 5.09355	ZK= -.53079
M= 40	X= 14.00000	RNEW= .65994	ZN= 4.85819	ZK= -.47114
M= 40	X= 14.50000	RNEW= .62977	ZN= 4.38935	ZK= -.40701
M= 40	X= 15.00000	RNEW= .61553	ZN= 4.18900	ZK= -.40432
M= 40	X= 15.50000	RNEW= .60178	ZN= 4.01559	ZK= -.37035
M= 40	X= 16.00000	RNEW= .59102	ZN= 3.86528	ZK= -.35896
M= 40	X= 16.50000	RNEW= .58251	ZN= 3.78587	ZK= -.35607
M= 40	X= 17.00000	RNEW= .57539	ZN= 3.70466	ZK= -.35891
M= 40	X= 17.90000	RNEW= .56455	ZN= 3.57751	ZK= -.39659

8500, 8517
 FOR 8537-8544, 8557-8564

MAIN PROGRAM

STORAGE JOINTS (1) (START DATA) (4) (END DATA) (COMMON) (DIFF)

INTERNAL CHARACTERISTICS (BLOCK, NAME)

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EXPAND ASSIGNMENT (BLOCK, TYPE, RELATIVE LOCATION, NAME)

0000	014000	100	0001	010000	1110	0001	000044	1740	0001	000001	1700
0001	010000	1100	0002	010000	1100	0001	000111	1750	0001	000133	1750
0002	010000	1100	0003	010000	1100	0001	000225	1760	0000	000225	1760
0003	010000	1100	0004	014244	AMW	0000	000245	APL	0000	000210	1770
0004	010000	1100	0005	014247	PI	0000	000240	T	0000	000275	1780
0005	010000	1100	0006	014271	Y	0000	000275	L	0000	000237	1790
0006	014000	AMW	0007	014401	F	0000	000240	FT	0000	000240	1800
0007	010240	W	0008	014271	X	0000	000270	Y	0000	000300	1810
0008	014241	AMW	0009	014246	Z	74			0000	000270	1820

00101	1*	000100	000100	000100	000100	000100	000100	000100	000100	000100
00102	2*	000100	000100	000100	000100	000100	000100	000100	000100	000100
00103	3*	000100	000100	000100	000100	000100	000100	000100	000100	000100
00104	4*	000100	000100	000100	000100	000100	000100	000100	000100	000100
00105	5*	000100	000100	000100	000100	000100	000100	000100	000100	000100
00106	6*	000100	000100	000100	000100	000100	000100	000100	000100	000100
00107	7*	000100	000100	000100	000100	000100	000100	000100	000100	000100
00108	8*	000100	000100	000100	000100	000100	000100	000100	000100	000100
00109	9*	000100	000100	000100	000100	000100	000100	000100	000100	000100
00110	10*	000100	000100	000100	000100	000100	000100	000100	000100	000100
00111	11*	000100	000100	000100	000100	000100	000100	000100	000100	000100
00112	12*	000100	000100	000100	000100	000100	000100	000100	000100	000100
00113	13*	000100	000100	000100	000100	000100	000100	000100	000100	000100
00114	14*	000100	000100	000100	000100	000100	000100	000100	000100	000100
00115	15*	000100	000100	000100	000100	000100	000100	000100	000100	000100
00116	16*	000100	000100	000100	000100	000100	000100	000100	000100	000100
00117	17*	000100	000100	000100	000100	000100	000100	000100	000100	000100
00118	18*	000100	000100	000100	000100	000100	000100	000100	000100	000100
00119	19*	000100	000100	000100	000100	000100	000100	000100	000100	000100
00120	20*	000100	000100	000100	000100	000100	000100	000100	000100	000100
00121	21*	000100	000100	000100	000100	000100	000100	000100	000100	000100
00122	22*	000100	000100	000100	000100	000100	000100	000100	000100	000100

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00171  27*      DO 71 J=0,NM1
00172  28*      ZN=0.
00173  29*      DO 4 J=0,NM1
00174  30*      F(J)=U(J+1)-U(J)
00175  31*      P(J)=1.5*F(J)
00176  32*      AMW=ABS(W(I)-W(J))
00177  33*      ANW=ABS(W(I)+W(J))
00178  34*      Z(J)=(W(I)+W(J))*(ALUB(1P))
00179  35*      Z(I)=Z(J)*.7
00180  36*      Z(J)=(W(J)-W(I))*(ALUB(1MW))
00181  37*      ZN=ZN+(Z(J)+Z(I))*(C(J)+L(J))
00182  38*      GO TO 4
00183  39*      ZN=(1./PI)*P(J)*(Z(J))
00184  40*      ZN=ZN+ZCN
00185  41*      ZN(I)=ZN
00186  42*      ZI=ZCN(X)
00187  43*      ZN(I)=ZN(I)+ZI
00188  44*      W(I)=C(Y(I))+ZN(I)*K(I)
00189  45*      IF (C(I)*DX+.Y(I)*E+.F(10.7)*DX+.ZN(I)*E+.F(10.7)*DX+.ZK(I)*E+.F(10.7)
00190  46*      GO TO 71
00191  47*      STOP
00192  48*      END

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END OF COMPILE: NO DEPENDENTS.


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&MFOR*LIB.MFCR.IS MAIN
FORTRAN-MAC 1.14S-05/27/70-13:41:10 (00) MAIN
00101 1. DIMENSION R(20),W(20)
00102 2. EXTENSIONAL F
00103 3. COMMON Y1,Z1,Y2,Z2
00104 4. 7 READ(5,7,END=1001)WA
00107 5. 7 FORMAT(F5.2)
00110 6. READ(5,9,END=100)Y2,Z2,Y1,Z1
00116 7. 9 FORMAT(4F10.7)
00117 8. L=0
00120 9. CALL N*QUAD(0.0,1.570796,F,20,R,W,L,V,DUM)
00121 10. VALUE=1.-V
00122 11. PRINT 200,WA,V,VALUE
00127 12. 200 FORMAT(5X,'WA=',F5.2,'244 SPECTRAL HEMI. EMIT=',E17.7,'24H SPEC
00127 13. $TRAL HEMI. REF.=',E17.7)
00130 14. GO TO 3
00131 15. 100 STOP
00132 16. END

```

END OF COMPILATION: NO DIAGNOSTICS.

```

QMFOR*LIB.MFOR*IS F
FORTRAN-MACC 1.145-05/22/79-13241:13 (,C1 F
00101 1. FUNCTION F(X)
00102 2. COMMON Y1,Z1,Y2,Z2
00103 3. C=((Y1+Y2+Z1+Z2)/(Y2**2+Z2**2))*SIN(X)
00104 4. H=((Y1-Z2-Y2+Z1)/(Y2**2+Z2**2))*SIN(X)
00105 5. SS=1.-C**2+H**2
00106 6. TT=?.*C*H
00107 7. R=(SS**2+TT**2)**0.5
00110 8. Y=ATAN(-TT/SS)
00111 9. C=Y1*COS(X)
00112 10. D=(Y2*COS(Y/2.)+Z2*SIN(Y/2.))*R**0.5
00113 11. P=Z1*COS(X)
00114 12. Q=(Z2*COS(Y/2.)-Y2*SIN(Y/2.))*R**0.5
00115 13. S=(C+D)**2+(P+Q)**2
00116 14. T=C**2-D**2+P**2-Q**2
00117 15. U=?.*P*Q-2.*C*Q
00120 16. R01=(T/S)**2+(U/S)**2
00121 17. C1=Y2*COS(X)
00122 18. D1=(Y1*COS(Y/2.)+Z1*SIN(Y/2.))*R**0.5
00123 19. P1=Z2*COS(X)
00124 20. Q1=(Z1*COS(Y/2.)-Y1*SIN(Y/2.))*R**0.5
00125 21. S1=(D1+Q1)**2+(P1+Q1)**2
00126 22. T1=D1**2-C1**2+Q1**2-P1**2
00127 23. U1=2.*Q1*C1-2.*P1*Q1
00130 24. R02=(T1/S1)**2+(U1/S1)**2
00131 25. R0=0.5*(R02+R01)
00132 26. E=1.-R0
00133 27. F=2.*E*SIN(X)*COS(X)
00134 28. RETURN
00135 29. END

```

END OF COMPILATION: NO DIAGNOSTICS.

```

@MFOR=LI3.MFOR,IS MAIN
FORTRAN-MACC 1.14S-1.9/14/70-15:34:44 (0)      MAIN
00101  1.      DTYNSTON F0(44),Y1(44),E0(44),TC(44),F0(44),P(20),W(20),PLANG(44)
00102  2.      EXTERNAL F,FF
00103  3.      COMMON YY,YZ,TT
00104  4.      READ(5,10)(FC(I),I=1,44),(Y1(J),J=1,44)
00105  5.      10 FORMAT(1)
00106  6.      DO 5 I=1,40
00107  7.      TE(I)=1700.*100.*FLOAT(T-1)
00108  8.      TT=TE(T)
00109  9.      V1=C.
00110 10.      PLAN=G.
00111 11.      DO 7 J=1,27
00112 12.      H=(1.5-P.2)/20.
00113 13.      Y=P.2+4*FLOAT(J-1)
00114 14.      C1=(6.675E-13)*(3.F14)**2
00115 15.      C2=(6.675E-13)*(3.F14)/(1.78E-2)
00116 16.      C3=C2/(Y1(J)*Y*TT)
00117 17.      IF(C3.GT. 85.L)GO TO 2
00118 18.      EX= EXP(C3)-1.
00119 19.      EA=7.*3.141592*C1/(Y1(J)**2*Y**5*Y)
00120 20.      PLAND(J)=EA
00121 21.      PLAN=PLAN+PLAND(J)
00122 22.      EB(J)=EA
00123 23.      V1=V1+4*F0(J)*E0(J)
00124 24.      2 CONTINUE
00125 25.      PLAN=4*(PLAN-L.5*PLAND(1)-L.5*PLAND(27))
00126 26.      V1=V1-L.5*H*(E0(1)*EB(1)+E0(27)*E0(27))
00127 27.      V2=N.
00128 28.      PLA=N.
00129 29.      DO 3 J=27,44
00130 30.      S=(10.-1.5)/17.
00131 31.      X=1.5+C*FLOAT(J-27)
00132 32.      D1=(6.675E-13)*(3.F14)**2
00133 33.      D2=(6.675E-13)*(3.F14)/(1.38E-8)
00134 34.      D3=C2/(Y1(J)*X*TT)
00135 35.      IF(D3.GT. 85.L)GO TO 3
00136 36.      FX= EXP(D3)-1.
00137 37.      FA=7.*3.141592*D1/(Y1(J)**2*X**5*FX)
00138 38.      PLAND(J)=FA
00139 39.      PLA=PLA+PLAND(J)
00140 40.      F3(J)=FA
00141 41.      V2=V2+C*E0(J)*F3(J)
00142 42.      3 CONTINUE
00143 43.      PLA=C*(PLA-L.5*PLAND(27)-L.5*PLAND(44))
00144 44.      PLAN=PLAN+PLA
00145 45.      V2=V2-L.5*G*(E0(27)*F0(27)+E0(44)*F0(44))
00146 46.      YY=Y1(1)
00147 47.      YZ=Y1(44)
00148 48.      L=N
00149 49.      CALL NTGUAD(P.07,0.7,F,20,R,W,L,V,NUM)
00150 50.      V3=F0(1)*V
00151 51.      PLAN=PLAN+V
00152 52.      L=N
00153 53.      CALL NTGUAD(10.,150.,FF,20,R,W,L,V,NUM)
00154 54.      V4=F0(44)*V
00155 55.      PLAN=PLAN+V

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00213      56.      VALUE=(V1+V2+V3+V4)/PLON
00214      57.      VAL=1.-VALUE
00215      58.      WRITE(F,20)TT,VALUE,VAL
00222      59.      20 FORMAT(5X,'TLE ',F6.1,'X','TOTAL HEM. EMITTANCE= ',F15.7,'SX',
00223      60.      5'TOTAL HEM. REFLECTANCE=',F15.7)
00224      61.      Y1(1)=YY
00225      62.      Y1(44)=YZ
00226      63.      * CONTINUE
00227      64.      STOP
00230      65.      END

```

END OF COMPILATION: NO DIAGNOSTICS.

```

@MFOR*LIB.MFCP.1S F
FORTRAN-MACC 1.14S-09/14/70-15:34:50 (0.0) F
00101 1. FUNCTION F(X)
00102 2. COMMON YY,YZ,YY
00103 3. C1=(5.625E-10)*(3.514)**2
00104 4. C2=(6.625E-10)*(3.514)/(1.7E-9)
00105 5. IF(X.LT.0.00001)GO TO 6
00107 6. C3=C2/(YY*X*Y)
00110 7. IF(C3.GT. 99.0)GO TO 8
00112 8. GO TO 10
00113 9. 80 IF(C3.LT.1.0E-10)GO TO 9
00115 10. EX= EXP(C3)-1.
00116 11. EA=7.*7.141592*C1/(YY**2*X**5.*CY)
00117 12. F=EA
00120 13. GO TO 7
00121 14. 6 F=0.
00122 15. 7 RETURN
00123 16. END

```

END OF COMPILATION: NO DIAGNOSTICS.

6MFOR*LIB.MFOR,ASIZY FF
FORTRAN-MACC 1.145-19/14/79-15:35:05 (1.0) FF

FUNCTION FF ENTRY POINT 000064

STORAGE USED:

0 DATA 000022
1 CODE 000074

COMMON BLOCKS:

2 BLANKS 000003

EXTERNAL REFERENCES:

3 EXP
4 NERR33

STORAGE ASSIGNMENT:

1	000053 FL	1	000054 7L	1	000023 80L	0 R	000001 C1	0 R	000002 C2
	0 R 000003 C7		0 R 000005 EA		0 R 000004 EX		0 R 000000 FF		0 000015 INJP
	2 R 000002 TT		2	000000 YY		2 R	000001 YZ		

```
00101 1. FUNCTION FF(X)
00102 2. COMMON YY,YZ,TT
00103 3. C1=(F.675E-19)*(3.714)**2
00104 4. C2=(6.625E-19)*(3.714)/(1.37E-8)
00105 5. IF(X.LT.0.00001)GO TO 1
00107 6. C3=C2/(YZ*X+TT)
00110 7. IF(C3.GT. 95.0)GO TO 5
00112 8. GO TO 9
00113 9. 80 IF(C3.LT.1.0E-10)GO TO 6
00115 10. EX= EXP(C3)-1.
00116 11. EA=2.*3.141592*C1/(YZ**2*X**3*EX)
00117 12. FF=EA
00120 13. GO TO 7
00121 14. F FF=0.
00122 15. 7 RETURN
00123 16. END
```

END OF COMPILATION: NO DIAGNOSTICS.

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