

REVIEW OF PARTICLE PROPERTIES

Particle Data Group

LBL--100 Rev.

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This review of the properties of leptons, mesons, and baryons is an updating of Review of Particle Properties, Particle Data Group [Rev. Mod. Phys. 52 (1980) No. 2, Part II]. Data are evaluated, listed, averaged, and summarized in tables. Numerous tables, figures, and formulae of interest to particle physicists are also included. A data booklet is available.

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I. Introduction, credits, consultants

This review is an updating through December 1981 of our previous review of particle properties [Particle Data Group (1980)]. As in previous editions we have

attempted to make the text as complete and self-contained as possible.

As usual, the results of our compilation are presented in two sections, the Tables of Particle Properties and the Data Card Listings. The Tables summarize

the properties of only those particles whose existence is in our judgment experimentally well founded and which have a high probability of standing the test of time. This is a conservative judgment, and surely some genuine resonances are omitted, awaiting confirmation (see section V below).

The Data Card Listings give up-to-date information, with references, on all reported particles, whether considered well established or not. The Listings also contain mini-reviews on questions of interest.

A history of the Particle Data Group, with a discussion of procedures and problems, has been given by Rosenfeld (1975), and a short survey of the history of some of the constants we compile may be found in Appendix III.

As in previous editions, we include a section of miscellaneous tables, figures, and formulae. These are aimed at the practicing high energy physics experimentalist. We welcome all suggestions and comments regarding topics for deletion or inclusion, etc. This year we have revised many of these items, but no new ones have been added.

A pocket-sized Particle Properties Data Booklet, containing the Tables and a reprint of the figures and formulae from the first part of the review, is available on request. For North and South America, Australia, and the Far East, write to Technical Information Department, Lawrence Berkeley Laboratory, Berkeley, CA 94720, USA. For all other areas, write to CERN Scientific Information Service, CH-1211 Geneva 23, Switzerland.

As usual, we wish to emphasize that we compile the experimental results of others. It is inappropriate to give us the credit for their countless hours of effort. We urge that references be given directly to the original data, and we provide complete references in the Data Card Listings for that purpose.

The responsibilities for the various sections can be broken down as follows:

- (1) *Stable particles*: R. Frosch, T. Shimada, R.E. Shrock, T.G. Trippe, C.G. Wohl, and G.P. Yost.
- (2) *Meson resonances*: M. Aguilar-Benitez, M.J. Losty, L. Montanet, F.C. Porter, M. Roos, and Ch. Walck.
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- (4) *General, including text*: All authors.

Consultants: To overcome gaps in our coverage, both intellectual and geographical, we have solicited the help of consultants:

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The usefulness of this compilation depends in large part on the interaction between the users and the authors and consultants. We appreciate comments, criticisms, and suggestions for improvements of all stages of data retrieval, processing, evaluation, and presentation.

II. Selection of data

All particles are considered to fall into one of the three groups:

(1) Stable particles, defined to be those immune to decay via the strong interaction, including the photon, the leptons, the η , the D and F charmed mesons, the Λ_c charmed baryon, etc.

(2) Meson resonances, including the ψ , the χ , and the T particles.

(3) Baryon resonances.

These groups are maintained within the two main parts of the compilation:

(1) Tables of Particle Properties.

(2) Data Card Listings.

The Data Card Listings contain the original information (data, references, etc.), weighted averages, comments, and "mini-reviews". Immediately preceding the Data Card Listings is an illustrative key thereto. We attempt to give complete Data Card Listings up to our closing date (January 1, 1982) for all journals listed in the Illustrative Key. As a general rule, we do not include results from preprints or unpublished conference reports. Exceptions to this rule are made on a case-by-case basis.

Many of our encoded results, those set off in parentheses, are not used for averaging. The reasoning is then often given in a footnote below the data. If the reason is not given, it is one of the following:

The result was presented with no error stated.

The result comes from a preprint or conference report. It is our experience that such results (and particularly the errors) often change before final publication. Accordingly we keep these new results in parentheses until they are published (or explicitly verified to us by the authors).

It involves some assumptions that we do not wish to incorporate.

It is of poor quality, e.g. bad signal-to-noise ratio.

It is inconsistent with other results, e.g. because of

different methods employed, rendering averaging meaningless.

It is not independent of other results, e.g. it is a result from one of several partial-wave analyses all using the same data, again rendering averaging meaningless.

Upper limits are not averaged (except in rare cases which are re-expressible as numbers with gaussian errors).

When the data for a particle have received special treatment or present special problems, this is noted in a mini-review in the Data Card Listings.

As time goes by, some early results lose all their weight in the averages. We may then remove them from the Listings without further comment. We usually do not remove the corresponding reference cards, however, so that our reference sections preserve the historical record. In this edition the meson section has undergone an extensive "house-cleaning". As a result it appears more readable (or so we hope). The earlier data may be found in the 1980 edition).

The Tables of Particle Properties contain "best" values obtained from the data in the Data Card Listings by various methods. The statistical procedures of section VII are used to combine independent data which have gaussian errors. Upper limits in the Tables usually represent the strongest limit available from a single experiment. The extent to which these methods are tempered by critical judgment is explained in the footnotes to the Tables. In general, however, the footnotes are less complete than is the collection of notes and mini-reviews in the Data Card Listings. The reader is thus encouraged to become familiar with the Data Card Listings and, ultimately, with the original references.

III. Nomenclature

A. Quantum numbers

The symbols $I^G(J^P)C_n$ represent:

I = isospin,

G = G parity,

J = spin (also s),

P = space parity,

C_n = charge-conjugation parity for the neutral member of the isospin multiplet.

We also use:

B = baryon number,

S = strangeness,
 C = charm,
 l = orbital angular momentum.

1. Mesons. The charge-conjugation operator C turns particle into antiparticle and has eigenvalues ± 1 only for neutral states; so it is useful to define an operator G which has eigenvalues for charged states too. This is usually ^{†1} defined by

$$G = C \exp(i\pi I_y). \quad (1)$$

A neutral nonstrange, noncharmed state is an eigenstate of $\exp(i\pi I_y)$ with eigenvalue $(-1)^l$. Then we can write the eigenvalue equation for the whole multiplet as

$$G = C_n (-1)^l. \quad (2)$$

where C_n (n for neutral) is the eigenvalue C would have if applied to the neutral member of the multiplet. Thus, for a π^0 , C has the eigenvalue $+1$, and since $l = 1$, $G = -1$. For a charged π^+ , there are no eigenvalues corresponding to C and to the isospin rotation, but eqs. (1) and (2) still give $G = -1$.

Consider a meson as a bound state of fermion-antifermion, e.g. quark-antiquark $\bar{q}q$, with orbital angular momentum l , and with the two fermion spins coupling to give a spin s . Then one can show that the charge-conjugation eigenvalue [defined as in eq. (2)] is

$$C_n = (-1)^{l+s}. \quad (3)$$

Eqs. (2) and (3) combine to give

$$G = (-1)^{l+s+l}. \quad (4)$$

The parity is

$$P = -(-1)^l. \quad (5)$$

Eqs. (3) and (5) combine to give

$$C_n P = -(-1)^s, \quad (6)$$

so all singlets (1S_0 , 1P_1 , ...) have $C_n P = -1$, and all triplets (3S_1 , ...) have $C_n P = +1$. For proofs of the above, see our 1969 text [Particle Data Group (1969)]

^{†1} Most texts define it as in eq. (1); see e.g. Gasiorowicz (1966); however, sometimes the rotation is taken about I_x . The difference between the two conventions is mentioned in a footnote in Källén (1964).

and Appendix by C. Zemach.

If, instead of $\bar{q}q$, we consider the meson as a state of *boson-antiboson* (e.g. $A_2 \rightarrow \bar{K}K$), it turns out that some signs cancel, and eqs. (3) and (4) [not eq. (5)!] apply *unchanged*. Of course, the mesons are often spinless, so s is zero, but the equations are more general. Eqs. (3) and (4) can be considered as selection rules forbidding many decays.

We now use eqs. (3) and (4) to introduce the concept of "Abnormal- C_n " mesons, i.e. mesons that cannot be composed of $\bar{q}q$. For this, it is sufficient to consider the SU(3) subgroup of the full unitary group of flavors, containing the u, d, and s quarks in a {3} representation.

This triplet of quarks is of course defined to have isospin and hypercharge properties such that $\bar{q}q$ can combine (according to the SU(3) relations $\{3\} \otimes \{3\} = \{8\} \oplus \{1\}$) so as to form only octets and singlets. The non-observation of "exotic" mesons (i.e., mesons in larger SU(3) representations, or mesons requiring at least a $q\bar{q}q\bar{q}$ structure) is of course a direct consequence of the naive quark model. It is less obvious that even some *octets* are forbidden by the model, namely those with $(J^P)C_n = (0^+)_{-}, (1^-)_{+}, (2^+)_{-}, \dots$. Such states are not observed, and this is an additional success of the naive quark model classification scheme.

When the naive quark model is extended to QCD, one expects gg gluonium mesons also. Since the gluon g is a flavor singlet, all gluonium states must be flavor singlets which can be expected to mix with nearby $\bar{q}q$ singlets. No gluonium states have been definitely established yet.

In what follows, do not confuse "Abnormal- C_n " with "Normal" or "Abnormal" J^P , both of which are allowed by the quark model. The series $J^P = 0^+, 1^-, 2^+, \dots$ is called Normal because $P = (-1)^l$ as for normal spherical harmonics, and $J^P = 0^-, 1^+, \dots$ is called Abnormal.

The top part of table 1 shows all the low angular momentum states that can be formed from $\bar{q}q$. Note that half of the J^P states can be formed by both a triplet and a singlet $\bar{q}q$ state, e.g. $^3P_1, ^1P_1$, or $^3D_2, ^1D_2$, in spectroscopic notation $\{(2s+1)l_J\}$. Eq. (3) shows that 3P_1 and 1P_1 have opposite C_n , so the $\bar{q}q$ model allows both. But the states 3P_0 and 3P_2 have no 1P counterparts. According to eq. (6) they have $C_n P = +1$, and with the $\bar{q}q$ model there is no way to form a state with a J^P of $^3P_{0,2}$ (i.e. $J^P = \text{Normal}$) and

Table 1

Orbital excitations of the $\bar{q}q$ system, and corresponding mesons. For the distinction between Abnormal J^P and Abnormal C_n , see text following eq. (6) in section III. Strange and charmed mesons share the same values of J^P as the $I = 0$ and 1 states shown, but are not eigenstates of G . The second column, which gathers together $(J^P)_{N\text{ or }A} C_n P$, is a redundant intermediate step intended to make the table easier to read. The table repeats itself for each radial excitation.

$\bar{q}q$ State $(2s+1)_{LJ}$	$(J^P)_{N\text{ or }A} C_n P$ Normal or abnormal	Examples of ground-state mesons				
		$I^G(J^P)_{C_n}$	Non-strange, Non-charmed $S=C=0$	Strange $ S =1$ $(I=\frac{1}{2})$	Charmed $ C =1$ $(I=\frac{1}{2})$	Strange, Charmed $ S = C =1$ $(I=0)$
NORMAL- C_n STATES THAT CAN COME FROM $\bar{q}q$ MODEL						
Parity - $1S_0$	$(0^-)_{A^-}$	$\begin{pmatrix} 0^+(0^+) \\ 1^-(0^+) \end{pmatrix}$	$n, n', \eta_c(2980) ?$ π	\therefore	D(1870)	F(2030) ?
	Parity + $3S_1$	$(1^-)_{N^+}$	$\begin{pmatrix} 0^-(1^-) \\ 1^+(1^-) \end{pmatrix}$	$\omega, \phi, J/\psi(3100)$ ρ	$K^*(892)$	$D^*(2010)$
Parity - $1P_1$		$(1^+)_{A^-}$	$\begin{pmatrix} 0^-(1^+) \\ 1^+(1^+) \end{pmatrix}$	H B	Q_B	mix to give Q_1, Q_2
		Parity + $3P_0$	$(0^+)_{N^+}$	$\begin{pmatrix} 0^+(0^+) \\ 1^-(0^+) \end{pmatrix}$	$\kappa, S^*, \chi(3415)$ δ	
	Parity - $3P_1$		$(1^+)_{A^+}$	$\begin{pmatrix} 0^+(1^+) \\ 1^-(1^+) \end{pmatrix}$	D, E, $\chi(3510)$ A_1	
		Parity + $3P_2$	$(2^+)_{N^+}$	$\begin{pmatrix} 0^+(2^+) \\ 1^-(2^+) \end{pmatrix}$	F, F', $\chi(3555)$ A_2	
Parity - $1D_2$	$(2^-)_{A^-}$		$\begin{pmatrix} 0^-(2^-) \\ 1^+(2^-) \end{pmatrix}$	A_3	Regge recurrence of the Abnormal- C_n state $(J^P)_{C_n} = (0^-)$	
	Parity + $3D_1$	$(1^-)_{N^+}$	same as $3S_1$	$\psi(3770)$		
		Parity - $3D_2$	$(2^-)_{A^+}$	$\begin{pmatrix} 0^-(2^-) \\ 1^+(2^-) \end{pmatrix}$		$\omega(1670)$
	Parity + $3D_3$		$(3^-)_{N^+}$	$\begin{pmatrix} 0^-(3^-) \\ 1^+(3^-) \end{pmatrix}$		g
Parity + $1F_3$		$(3^+)_{A^-}$	$\begin{pmatrix} 0^-(3^+) \\ 1^+(3^+) \end{pmatrix}$	same as $3P_2$	h	
	Parity - $3F_2$	$(2^+)_{N^+}$	h			
		Parity + $3F_3$				$(3^+)_{A^+}$
	Parity - $3F_4$					$(4^+)_{N^+}$
ABNORMAL- C_n STATES THAT CANNOT COME FROM $\bar{q}q$ MODEL						
Abnormal C_n states Have no $\bar{q}q$ model	$(0^-)_{A^+}$	$\begin{pmatrix} 0^-(0^-) \\ 1^+(0^-) \end{pmatrix}$	All except $J^P = 0^-$ are $J^P = \text{normal},$ $C_n P = -1$			
	$(1^-)_{N^-}$	$\begin{pmatrix} 0^-(1^-) \\ 1^-(1^-) \end{pmatrix}$				
		$\begin{pmatrix} 0^-(0^+) \\ 1^-(0^+) \end{pmatrix}$				
	$(2^-)_{N^-}$	$\begin{pmatrix} 0^-(2^-) \\ 1^-(2^-) \end{pmatrix}$				
		$\begin{pmatrix} 0^-(3^-) \\ 1^-(3^-) \end{pmatrix}$				

with $C_n P = -1$. As mentioned, such octets have not been found. With the help of table 1 one can also see that the special state 1S_0 , $C_n P = +1$, cannot be formed, so has Abnormal C_n .

When, in addition to the l -excitation, there are radial excitations of the $\bar{q}q$ system, table 1 repeats itself, and we need a radial quantum number n for each repetition ($n = 1$ for the ground state). Examples of first radial excitations, $n = 2$, are $\rho'(1600)$, $\psi(3685)$, and $\Upsilon(10020)$. Examples of further possible radial excitations can be found in the ψ and Υ families.

2. General remarks. Well-established quantum numbers are underlined in the Tables of Particle Properties (except for stable particles, where most of the quantum numbers are established). We have used what evidence is available (sometimes flimsy) to guess many of the remaining ones, and we have indicated with "??" ones (in the Baryon Table) for which there is almost no evidence.

As is customary, we define antiparticles as the result of operating with CPT on particles, so both share the same spins, masses, and mean lives. Whenever there is a particularly interesting test of CPT invariance we include it in the Stable Particles Table.

B. Particle names

If a meson has a well-accepted colloquial name, we use it. If not, we name it by a single symbol which specifies its J^P -"Normality", its isospin I , its strangeness S and charm C , and, for a non-strange, non-charmed meson, its G parity.

The name conventions for mesons are given in the first part of table 2.

For some pairs of mesons with supposedly identical quantum numbers, we also use primes; e.g. η , η' ; f , f' ; ρ , ρ' . Note that primes and subscripts do not carry any further specific meaning.

For baryons no attempt has been made to attach a subscript about J and P . The name conventions are given in the second part of table 2. For stable baryons of each I and S we use the symbol standing alone; for resonances, the mass is in parentheses [i.e. $N(1680)$, $\Lambda(1405)$, $\Sigma(1775)$, etc.]. The J^P assignments are reported in the Baryon Table as $\frac{1}{2}+$, $\frac{3}{2}-$, $\frac{5}{2}+$, etc., and also by the symbols P_{11} , D_{13} , F_{15} , which refer to the $\pi\pi$ or $K\rho$ partial-wave amplitude in which the resonant state occurs (the first subscript refers to the isospin

Table 2
Particle name conventions.

Name	I	S	C	G
Mesons, J^P "Normal"				
$\omega, \phi, \psi, \Upsilon^a)$	0	0	0	-
ϵ	0	0	0	+
δ	1	0	0	-
ρ	1	0	0	+
K^*, κ	1/2	± 1	0	
D^*	1/2	0	± 1	
F^*	0	± 1	± 1	
Mesons, J^P "Abnormal"				
η	0	0	0	+
π	1	0	0	-
K, Q	1/2	± 1	0	
D	1/2	0	± 1	
F	0	± 1	± 1	
Baryons				
N	1/2	0	0	
Δ	3/2	0	0	
Z_I	0, 1	+1	0	
Λ	0	-1	0	
Σ	1	-1	0	
Ξ	1/2	-2	0	
Ω	0	-3	0	
Λ_c	0	0	1	
Σ_c	1	0	1	

a) We use the symbol ω for those $I^G = 0^-$ mesons which are mainly $u\bar{u}$ and $d\bar{d}$ quark states; ϕ for those which are mainly $s\bar{s}$ quark states, ψ for mainly $c\bar{c}$ states, and Υ for mainly $b\bar{b}$ states.

state: $2 \times I$ for N , Δ , and Ξ , and just I for Z , Λ , and Σ). When two or more baryons have identical quantum numbers we warn the reader by adding primes to the spectroscopic symbol as explained in footnote (a) of the Baryon Table.

IV. Conventions and parameters for strong interactions

A. Partial-wave amplitudes and resonance parameters

The vast majority of information concerning baryon resonances comes in the form of partial-wave analyses. In addition data concerning meson resonances ($\pi\pi$, $K\pi$, $\pi\pi\pi$) are, with increasing frequency, being subjected to partial-wave analyses. We thus find it natural to introduce the resonance parameters which we com-

pile in terms of a Breit-Wigner approximation for the partial-wave amplitude.

In general the elastic amplitude for a given angular momentum l may be written as

$$T_{11} = [\eta \exp(2i\delta) - 1]/2i, \tag{1}$$

where η is the absorption parameter ($0 \leq \eta \leq 1$) and δ is the phase shift. The subscripts 11 on T denote scattering from channel 1 to channel 1 (e.g. $\pi\pi \rightarrow \pi\pi$ or $\bar{K}K \rightarrow \bar{K}K$).

In fig. 1 we show an Argand plot of the elastic partial wave amplitude T_{11} . It illustrates geometrically how the real parameters η and δ are related to the real and imaginary parts of T_{11} . Many examples of such Argand plots may be found in the Baryon Data Card Listings.

Consider the so-called non-relativistic Breit-Wigner approximation for T_{11} :

$$T_{11} = \frac{1}{2} \Gamma_1 / (M - E - \frac{1}{2} i \Gamma), \tag{2}$$

where E is the c.m. energy or invariant mass, Γ_1 and Γ are the elastic and total widths, and M is the resonance mass. Eq. (2) is, of course, not the only possible description of a resonant amplitude; but it suffices to illustrate the properties of partial-wave amplitudes which we associate with resonance behavior in the absence of any background in the same partial wave (see, e.g., the $\pi N D_{15}$ and F_{15} waves in the Baryon Data Card Listings). Usually the widths contain barrier-penetration factors which can vary rapidly with energy. Near threshold, $\Gamma_1(E)$ should start up as q^{2l+1}

(also true for the inelastic width Γ_β). Various E dependences are then used for Γ_1 , mostly of the form

$$\Gamma_1(E) \propto (qR)^{2l+1} / [\text{const} + \dots + (qR)^{2l}]; \tag{3}$$

see Jackson (1964), Pišut and Roos (1968), and Barbaro-Galtieri (1968).

The BW approximation to the amplitude for an inelastic process leading from channel 1 to channel β ($\pi\pi \rightarrow \bar{K}K$ or $\bar{K}N \rightarrow \Sigma\pi$, for example) is

$$T_{1\beta} = \frac{1}{2} (\Gamma_1 \Gamma_\beta)^{1/2} / (M - E - \frac{1}{2} i \Gamma) \\ = (x_1 x_\beta)^{1/2} [\frac{1}{2} \Gamma / (M - E - \frac{1}{2} i \Gamma)], \tag{4}$$

where

$$\Gamma = \sum_1^N \Gamma_\beta, \quad x_\beta = \Gamma_\beta / \Gamma. \tag{5}$$

and x_1 (called the elasticity) is often written x_e . (Note that in the Data Card Listings we use the symbol P_β to denote x_β .) The channel cross section $\sigma_{1\beta}$ for the reaction $1 \rightarrow \beta$, for spin 0-spin 1/2 scattering, is

$$\sigma_{1\beta} = 4\pi \lambda^2 (J + \frac{1}{2}) |T_{1\beta}|^2, \tag{6}$$

where $J = l \pm \frac{1}{2}$.

The important features of eq. (4) which characterize resonant behavior in the Argand diagram ($\text{Im } T_{1\beta}$ versus $\text{Re } T_{1\beta}$) are:

energy variation given by circles with diameter $(x_1 x_\beta)^{1/2}$ and maximum amplitude at $E = M$ of

$$T_{1\beta}^{\text{max}} = i(x_1 x_\beta)^{1/2}; \tag{7}$$

a maximum in the speed near resonance, given approximately by

$$\text{"Speed"}(\text{res}) = |dT_{1\beta}/dE|_{E=M} = 2(x_1 x_\beta)^{1/2} / \Gamma(E), \tag{8}$$

for slowly varying $\Gamma(E)$. These features may be related to the η, δ representation of T_{11} . Thus when $E = M$, δ is either 90° ($x_1 > \frac{1}{2}$) or 0° ($x_1 < \frac{1}{2}$) and η dips to its minimum value.

These simple properties can be used to judge the presence or absence of resonance behavior in an Argand plot, but do not necessarily constitute the criteria we use (see section V). It must also be kept in mind that eqs. (2) and (4) are only approximations to the "true" amplitude. The simple picture given above can be distorted by various effects:

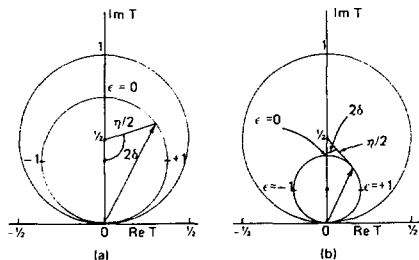


Fig. 1. Argand plots for the elastic partial wave amplitude T_{11} . The outer circles are the unitarity bound ($\eta = 1$). The inner circles correspond to the Breit-Wigner approximation of eq. (2) for (a) $x_1 = \Gamma_1/\Gamma = 0.75$ and (b) $x_1 = 0.4$. Note: $e = 2(M - E)/\Gamma$.

the presence of "background" in the same partial wave as the resonance,

two resonances in the same partial wave overlapping in energy,

the resonant energy M being close to an inelastic channel threshold, in which case a K -matrix-like parametrization is more appropriate,

the speed of the resonance being very slow so that the resonance is very broad, and the Breit-Wigner formula a bad approximation.

B. Sign conventions for resonance couplings

Consider the partial width Γ_β of a resonance decaying into the channel β . We can always define a coupling constant such that

$$\Gamma_\beta \propto G_\beta^2.$$

In this case the inelastic amplitude in the Breit-Wigner approximation, eq. (4), will go as

$$T_{1\beta} \propto G_1 G_\beta / (M - E - \frac{1}{2}i\Gamma),$$

where G_1 is the coupling constant for the elastic channel. In the context of exact SU(3) symmetry the relative signs of the product $G_1 G_\beta$ for different resonances are often useful as a consistency check on SU(3)

assignment of baryon resonances. See Appendix I for further details.

In the Data Card Listings for baryon resonances, we tabulate measured values for $(x_1 x_\beta)^{1/2} \propto G_1 G_\beta$. When the sign of the amplitude is determined, it is given; absence of an explicit sign indicates that it is undetermined (not that it is positive). For Λ and Σ resonances, the signs are chosen according to the convention advocated by Levi-Setti (1969) and used in the table of SU(3) Isoscalar Factors presented in this review. Thus the signs multiplying the Breit-Wigner amplitudes for $\bar{K}N \rightarrow \Sigma(1385) \rightarrow \Sigma\pi, \Lambda\pi$ and $\bar{K}N \rightarrow \Lambda(1405) \rightarrow \Sigma\pi$ are simply the product of the phases of the appropriate isoscalar factors. This convention is shown in fig. 2, adapted from Levi-Setti (1969).

C. Types of partial-wave analyses

Partial-wave analyses (PWA) are classified into three categories in the Data Card Listings: energy-independent partial-wave analyses (IPWA), energy-dependent partial-wave analyses (DPWA), and model-dependent partial-wave analyses (MPWA), in increasing order of the number of explicit supplementary hypotheses that are used to extract the amplitudes from experimental data.

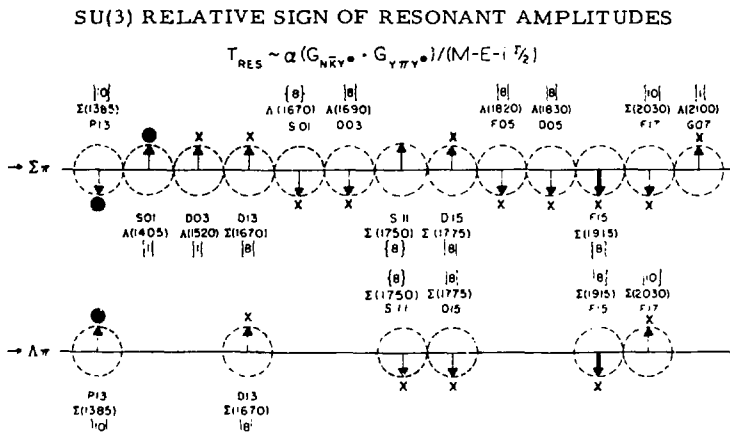


Fig. 2. Plot adapted from Levi-Setti (1969) showing the sign convention adopted here for the $\Sigma\pi$ and $\Lambda\pi$ amplitudes. Once the signs of one $l = 0$ and one $l = 1$ amplitude are fixed, the others can be measured relative to these two. Arrows here indicate signs predicted by SU(3); x marks indicate the observed phases; ● indicates phase chosen according to sign convention described in text. The $\Sigma(1915)$ predictions have been changed from Levi-Setti's original figure.

In an IPWA, data at different energies are analyzed separately. Usually each partial wave included in the fit is allowed to vary freely (subject to unitarity constraints) over some large region, and waves whose angular momenta are above some cutoff value are assumed to be negligible. The sharp cutoff in angular momentum resolves continuum ambiguities in the solution (such as the overall phase ambiguity), but there remains a finite number of indistinguishable "best" solutions (i.e., solutions corresponding to identical physical observables) which have been codified by Barrelet (1972). In addition, there are generally some nearby solutions (and their associated Barrelet ambiguities) which have chi-squared values close to the minimum one.

At the end of the analysis a choice is made among these many solutions, usually on the basis of energy continuity. A popular criterion for making this choice is the shortest path technique in which the total "length" of the preferred solution is chosen to be a minimum. The definition of "length" used here is not universal but is usually closely related to the total geometrical length of the lines representing the various partial-wave amplitudes in Argand plots (see the baryon section of the Data Card Listings for examples of Argand plots). Various other criteria which are also used in some analyses are, e.g., matching with known solutions at low energies, the presence of known resonances in the final results, and limited inelasticity in high partial waves.

In a DPWA, data at different energies are fit simultaneously by using an energy-dependent parametrization of the partial-wave amplitudes. The parametrization is usually chosen to include both resonances and nonresonant background of some sort and an attempt is made to keep it as "model independent" as possible. Often the data are grouped into several energy bins which are fit separately rather than trying to fit the whole energy range under consideration simultaneously. One of the main advantages of DPWA over IPWA is that sparse data spread over many different energies can be analyzed, e.g., nearly all $S = -1$ analyses are DPWA. In addition, the built-in energy continuity helps to resolve the ambiguities that plague IPWA and eases the problems associated with resonance parameter extraction. The price one pays for these advantages lies in the danger of systematic error in the amplitudes and poor fits to the data if the parametrization

is poorly chosen or insufficiently flexible.

An MPWA also uses an energy-dependent parametrization, but one based on explicit model-dependent theoretical assumptions such as Regge exchanges. This technique is usually applied to reactions where the data are incomplete. There is, of course, no sharp distinction between DPWA and MPWA, and a well chosen MPWA parametrization may actually be less biased than a model-independent but poorly chosen DPWA parametrization.

D. Production of resonances

Hereby, we mean the observation of statistically significant peaks in invariant mass plots or, loosely, in integrated cross sections. Many meson resonances are of this type. We expect most of these peaks to be associated with Breit-Wigner behavior in appropriate Argand plots; thus the ρ meson peak in $\pi\pi$ mass plots is firmly related to the $l = 1$, $l = 1$ $\pi\pi$ phase shift passing through 90° . The shape of the ρ meson is reasonably well described by the relativistic Breit-Wigner formula with the three parameters M , $\Gamma(M)$, and R of eq. (3).

From mass plots we can determine M , Γ , and the approximate branching ratios

$$x_\alpha/x_\beta = \Gamma_\alpha/\Gamma_\beta. \quad (9)$$

In the case of total cross sections, the peak above background gives us, using the optical theorem, the product $(J + \frac{1}{2})x_c$:

$$\sigma^{\text{tot}}(E = M) = 4\pi\chi^2(J + \frac{1}{2})x_c. \quad (10)$$

V. Criteria for resonances

An experimentalist who sees indications of a resonance in some energy (or mass) region will of course want to know what has been seen in that region in the past; hence, we strive to have the Data Card Listings serve as an archive for all substantial claims for resonances. (However, we do not intend to preserve a perennial record of claims proved to be wrong. Some such claims have been removed from the Meson Listings in 1980. We refer the interested reader to our 1980 edition for our complete archives.)

For the Tables of Particle Properties, on the other hand, we wish to be more conservative than for the

Listings and to include only those peaks or resonances which we feel have a large chance of survival. An arrow (\rightarrow) at the left of the Tables of Particle Properties indicates that a questionable candidate has been omitted from the Table, but that it can be found in the corresponding part of the Data Card Listings. One's betting odds for survival are of course subjective; therefore no precise criteria can be defined. Very slow speeds (ϵ and κ) make it quite difficult to decide what is a resonance and what is not. For more detailed discussions, see the mini-reviews in the Listings. In what follows we shall attempt to specify some guidelines.

(a) When energy-independent partial-wave analyses are available (mostly for N^* 's), approximate Breit-Wigner behavior of the amplitude appears to us to be the most satisfactory test for a resonance. We can check that the Argand plot follows roughly a left-hand circle, and that the "speed" of the amplitude also shows a maximum near the resonance energy; further, there should be data well above the resonance, showing that the speed again decreases. Indeed proper behavior of the partial-wave amplitude could accredit a resonance even if its elasticity is too small to make a noticeable peak in the cross section.

Of course even if Argand plots are available, it may still be a matter of opinion as to what behavior constitutes a resonance. Such an example is the $Z_0(1780)$ state seen in KN total cross-section experiments and in partial-wave analysis. The partial-wave analyses of Giacomelli (1974) and Martin (1975) find preferred solutions which exhibit a resonance-like loop in the P_{01} wave near 1740 MeV. However, Giacomelli et al. and Martin point out that, despite the resonantlike appearance of the loop, the evidence for resonant energy dependence is inconclusive. Thus we omit the $Z_0(1780)$ from the Baryon Table. A similar quandary has existed for some time concerning the $Z_1(1900)$, and it too has been omitted from the Tables.

(b) When there are insufficient data to perform energy-independent analyses, one often resorts to energy-dependent partial-wave analyses (mostly for Y^* 's). In this case Breit-Wigner behavior is an input. We therefore require that resonance solutions be found by several different analyses, preferably in different channels ($\bar{K}N \rightarrow \bar{K}N, \pi\Sigma$, etc.), before putting the claim in the table.

(c) Partial-wave analyses of three-body final states ($\pi N \rightarrow \pi\pi N$) are now available. While these analyses

are based on the isobar model ($\pi N \rightarrow \rho N, \pi\Delta$, etc.) and are subject to theoretical objections of varying importance, they provide increasingly reliable information on inelastic decay modes of otherwise established resonances.

(d) Most mesons, Ξ^* peaks, and high-mass N^* and Y^* peaks fall into a category for which no partial-wave analyses exist. In general we accept such peaks if they are experimentally reliable, of high statistical significance or observed in several different production processes.

Thus, we enter into the Tables of Particle Properties only states for which there is experimentally convincing evidence, and we expect that most of these will be confirmed as resonances.

VI. Conventions and parameters for weak and electromagnetic decays

A. Muon-decay parameters

The μ -decay parameters describe the momentum spectrum (ρ and η), the asymmetry (ξ and δ), and the helicity (h) of the electron in the process $\mu^- \rightarrow e^- + \nu_\mu + \bar{\nu}_e$. Assuming a local and lepton-conserving interaction, the matrix element in the charge-retention form may be written as

$$\sum_i \langle \bar{e} | \Gamma_i | \mu \rangle \langle \bar{\nu}_\mu | \Gamma_i (C_i + C'_i \gamma_5) | \nu_e \rangle,$$

where the summation is taken over $i = S, V, T, A, P$. Using the definitions and sign conventions of Sachs and Sirlin (1975) and Scheck (1978) for the Lorentz-covariant operators, we have for the momentum parameters:

$$\rho = (3g_A^2 + 3g_V^2 + 6g_T^2)/D,$$

$$\eta = (g_S^2 - g_P^2 + 2g_A^2 - 2g_V^2)/D;$$

for the asymmetry parameters:

$$\xi = \frac{-6g_S g_P \cos \phi_{SP} - 8g_A g_V \cos \phi_{AV} + 14g_T^2 \cos \phi_{TT}}{D},$$

$$\delta = (-6g_A g_V \cos \phi_{AV} + 6g_T^2 \cos \phi_{TT})/D\xi;$$

and for the parameter describing the helicity of the electron:

$$h = \frac{2g_S g_P \cos \phi_{SP} + 8g_A g_V \cos \phi_{AV} + 6g_T^2 \cos \phi_{TT}}{D}$$

Here

$$D = g_S^2 + g_P^2 + 4g_V^2 + 6g_T^2 + 4g_A^2,$$

$$g_i^2 = |C_i|^2 + |C_i'|^2,$$

and

$$\cos \phi_{ij} = \text{Re}(C_i^* C_j' + C_i' C_j^*) / g_i g_j.$$

The quantities g_i are defined to be real non-negative numbers, and the ϕ_{ij} are phase angles between the i -type and j -type interactions. Under the assumption of two-component neutrinos $C_V' = -C_V$ and $C_A' = -C_A$, the S, P, and T terms vanish, and ϕ_{AV} is the phase angle between C_A and C_V in the complex plane.

By using the above equations and the experimental determinations of ρ , η , ξ , δ , and h , limits can be placed on g_S/g_V , g_A/g_V , g_T/g_V , g_P/g_V , and ϕ_{AV} . The results, given in the Data Card Listings, assume neither two-component neutrinos nor time-reversal invariance. If, however, two-component neutrinos are assumed, then $\sin \phi_{AV}$ is the amplitude of time-reversal violation. Note that most experiments study only the upper end of the spectrum where ρ and η are highly correlated, so they can only report ρ for $\eta \equiv 0$ and η for $\rho = \frac{3}{4}$. The values for ρ and η we use here were obtained by combining measurements of both upper and lower ends of the spectrum and turn out to be nearly uncorrelated.

Note also that the radiative corrections are unambiguous only when $g_S = g_T = g_P = 0$. The same limits on g_A/g_V and ϕ_{AV} are obtained, however, as when g_S , g_T , and g_P are left free.

Current values for the asymmetry parameters as well as $|g_A/g_V|$ and ϕ_{AV} are given in the Addendum to the Stable Particle Table. In addition, upper limits on $|g_S/g_V|$, $|g_T/g_V|$ and $|g_P/g_V|$ are given in the μ section of the Stable Particle Data Card Listings.

B. K-decay parameters

1. *Dalitz plot for $K \rightarrow 3\pi$ decays.* The Dalitz plot distribution for the τ mode ($K^\pm \rightarrow \pi^\pm \pi^\pm \pi^\mp$), the τ' mode ($K^\pm \rightarrow \pi^0 \pi^0 \pi^\pm$), and the τ^0 mode ($K_L^0 \rightarrow \pi^+ \pi^- \pi^0$) of K decay can be parametrized by a series expansion such as that introduced by Weinberg (1960). We use the form

$$|M|^2 \propto 1 + g \frac{s_3 - s_0}{m_{\pi^+}^2} + h \left(\frac{s_3 - s_0}{m_{\pi^+}^2} \right)^2 + j \frac{s_2 - s_1}{m_{\pi^+}^2} + k \left(\frac{s_2 - s_1}{m_{\pi^+}^2} \right)^2 + \dots, \quad (1)$$

where $m_{\pi^+}^2$ has been introduced so as to make the coefficients g , h , j , and k dimensionless, and

$$s_i = (P_K - P_i)^2 = (m_K - m_i)^2 - 2m_K T_i, \quad i = 1, 2, 3,$$

$$s_0 = \frac{1}{3} \sum_i s_i = \frac{1}{3} (m_K^2 + m_1^2 + m_2^2 + m_3^2).$$

Here the P_i are four-vectors, m_i and T_i are the mass and kinetic energy of the i th pion, and the index 3 is used for the odd pion.

The coefficient g is a measure of the slope in the variable s_3 (or T_3) of the Dalitz plot, while h and k measure the quadratic dependence on s_3 and $(s_2 - s_1)$, respectively. The coefficient j is related to the asymmetry of the plot and must be zero if CP invariance holds (C stands for charge conjugation throughout the discussion in this section). Note also that if CP is good, g must be the same for τ^+ and τ^- , and similarly for h and k .

Since different experiments use different forms for $|M|^2$, in order to compare the experiments we have converted to g , h , j , and k whatever coefficients have been measured. See the mini-review in the K^\pm section of the Stable Particle Data Card Listings for details on this point. The results are given in the Addendum to the Stable Particle Table and in the K^\pm and K_L^0 sections of the Stable Particle Data Card Listings.

Relations among τ^\pm , τ'^\pm , and τ^0 are predicted by the $\Delta I = \frac{1}{2}$ rule.

2. *Form factors in $K_{\ell 3}$ leptonic decays.* Assuming that only the vector current contributes to these decays, we write the matrix element as

$$M \propto f_+(t) [(P_K + P_\pi)_\mu \bar{u}_\ell \gamma_\mu (1 + \gamma_5) u_\nu] + f_-(t) [m_\ell \bar{u}_\ell (1 + \gamma_5) u_\nu], \quad (2)$$

where P_K and P_π are the four-momenta of K and π mesons; m_ℓ is the lepton mass; f_+ and f_- are dimensionless form factors which can depend only on $t = (P_K - P_\pi)^2$, the square of the four-momentum

transfer to the leptons, f_+ and f_- are relatively real if time-reversal invariance holds for these decays. $K_{\mu 3}$ experiments measure f_+ and f_- , while $K_{e 3}$ experiments are sensitive only to f_+ because the presence of the lepton mass makes the f_- term negligible.

(a) $K_{\mu 3}$ experiments. Analyses of $K_{\mu 3}$ data frequently assume a linear dependence of f_+ and f_- on t , i.e.

$$f_{\pm}(t) = f_{\pm}(0)[1 + \lambda_{\pm}(t/m_{\pi}^2)] \quad (3)$$

Most $K_{\mu 3}$ data are adequately described by eq. (3) for f_+ and a constant f_- (i.e. $\lambda_- = 0$). There are two equivalent parametrizations commonly used in these analyses:

(1) λ_+ , $\xi(0)$ parametrization. Analyses of $K_{\mu 3}$ data often introduce the ratio of the two form factors

$$\xi(t) = f_-(t)/f_+(t).$$

The $K_{\mu 3}$ decay distribution is then described by the two parameters λ_+ and $\xi(0)$ (assuming time reversal invariance and $\lambda_- = 0$). These parameters can be determined by three different methods:

Method A. By studying the Dalitz plot or the pion spectrum of $K_{\mu 3}$ decay. The Dalitz plot density is [see, e.g. Chounet et al. (1972)]:

$$\rho(E_{\pi}, E_{\mu}) \propto f_+^2(t)[A + B\xi(t) + C\xi(t)^2],$$

where

$$A = m_K(2E_{\mu}E_{\nu} - m_K E_{\pi}^2) + m_{\mu}^2(\frac{1}{2}E_{\pi}' - E_{\nu}),$$

$$B = m_{\mu}^2(E_{\nu} - \frac{1}{2}E_{\pi}'),$$

$$C = \frac{1}{4}m_{\mu}^2 E_{\pi}'^2,$$

$$E_{\pi}' = E_{\pi}^{\text{max}} - E_{\pi} = (m_K^2 + m_{\pi}^2 - m_{\mu}^2)/2m_K - E_{\pi}.$$

Here E_{π} , E_{μ} , and E_{ν} are, respectively, the pion, muon, and neutrino energies in the kaon center of mass. The density ρ is fit to the data to determine the values of λ_+ , $\xi(0)$, and their correlation.

Method B. By measuring the $K_{\mu 3}/K_{e 3}$ branching ratio and comparing it with the theoretical ratio [see, e.g., Fearing et al. (1970)] as given in terms of λ_+ and $\xi(0)$, assuming μ -e universality:

$$\Gamma(K_{\mu 3}^{\pm})/\Gamma(K_{e 3}^{\pm}) = 0.6457 + 1.4115\lambda_+ + 0.1264\xi(0) + 0.0192\xi(0)^2 + 0.0080\lambda_+\xi(0),$$

$$\Gamma(K_{\mu 3}^0)/\Gamma(K_{e 3}^0) = 0.6452 + 1.3162\lambda_+ + 0.1246\xi(0) + 0.0186\xi(0)^2 + 0.0064\lambda_+\xi(0).$$

This cannot determine λ_+ and $\xi(0)$ simultaneously but simply fixes a relationship between them.

Method C. By measuring the muon polarization in $K_{\mu 3}$ decay. In the rest frame of the K, the μ is expected to be polarized in the direction A with $P = A/|A|$, where A is given [Cabibbo and Maksymowicz (1964)] by

$$A = a_1(\xi)p_{\mu} - a_2(\xi)\left[\frac{p_{\mu}}{m_{\mu}}\left(m_K - E_{\pi} + \frac{p_{\pi} \cdot p_{\mu}}{|p_{\mu}|^2}(E_{\mu} - m_{\mu})\right) + p_{\pi}\right] + m_K \text{Im} \xi(t)(p_{\pi} \times p_{\mu}).$$

If time-reversal invariance holds, ξ is real, and thus there is no polarization perpendicular to the K-decay plane. Polarization experiments measure the weighted average of $\xi(t)$ over the t range of the experiment, where the weighting accounts for the variation with t of the sensitivity to $\xi(t)$.

(2) λ_+ , λ_0 parametrization. Some of the more recent $K_{\mu 3}$ analyses have parametrized in terms of the form factors f_+ and f_0 which are associated with vector and scalar exchange, respectively, to the lepton pair. f_0 is related to f_+ and f_- by

$$f_0(t) = f_+(t) + [t/(m_K^2 - m_{\pi}^2)]f_-(t).$$

Here $f_0(0)$ must equal $f_+(0)$ unless $f_-(t)$ diverges at $t = 0$. The earlier assumption that f_+ is linear in t and f_- is constant leads to f_0 linear in t :

$$f_0(t) = f_0(0)[1 + \lambda_0(t/m_{\pi}^2)].$$

With the assumption that $f_0(0) = f_+(0)$, the two parametrizations, $(\lambda_+, \xi(0))$ and (λ_+, λ_0) are equivalent as long as correlation information is retained. (λ_+, λ_0) correlations tend to be less strong than $(\lambda_+, \xi(0))$ correlations.

The experimental results for $\xi(0)$ and its correlation with λ_+ are listed in the K^{\pm} and K_L^0 sections of the Stable Particle Data Card Listings in section XIA, XIB, or XIC depending on whether method A, B, or C discussed above was used. The corresponding values of λ_+ are listed in subsection L + M.

Because current experiments tend to use the (λ_+, λ_0) parametrization, we have added a subsection L0 for λ_0 results. Wherever possible we have converted $\xi(0)$ results into λ_0 results and vice versa.

(b) $K_{e 3}$ experiments. Analysis of $K_{e 3}$ data is simpler than that of $K_{\mu 3}$ because the second term of the

matrix element assuming a pure vector current [eq. (2) above] can be neglected. Here f_+ is usually assumed to be linear in t , and the linear coefficient λ_+ of eq. (3) is determined.

If we remove the assumption of a pure vector current, then the matrix element for the decay, in addition to the terms in eq. (2), would contain

$$+ 2m_K f_S \bar{u}_c (1 + \gamma_5) u_\nu + (2f_T/m_K)(P_K)_\lambda (P_\pi)_\mu \bar{u}_c \sigma_{\lambda\mu} (1 + \gamma_5) u_\nu,$$

where f_S is the scalar form factor, and f_T is the tensor form factor. In the case of the K_{e3} decays where the f_- term can be neglected, experiments have yielded limits on $|f_S/f_+|$ and $|f_T/f_+|$.

The K_{e3} results for λ_+ , $|f_S/f_+|$, and $|f_T/f_+|$ are listed in the subsections L + M, FS, and FT, respectively, of the K^\pm and K_L^0 sections of the Stable Particle Data Card Listings.

See also the Note on K_{e3}^\pm and K_{e3}^0 Form Factors in the K^\pm section of the Stable Particle Data Card Listings for additional discussion of the $K_{\mu 3}^0$ parameters, correlations, and conversion between parametrization and also for a comparison of the experimental results.

3. CP violation in K^0 decays. We list parameters for four different reactions in which CP can be tested [for details, see Okun and Rubbia (1967), Steinberger (1969), and Wolfenstein (1969)].

(a) $K_S \rightarrow \pi^+ \pi^- \pi^0$. The quantity measured here is the ratio of amplitudes

$$A_S(K_S \rightarrow \pi^+ \pi^- \pi^0)/A_L(K_L \rightarrow \pi^+ \pi^- \pi^0) \equiv x + iy. \quad (4)$$

If CPT invariance holds and there is no $I = 3$ state present, then x can be neglected and CP violation would be observed as a nonzero y . We give the result for eq. (4) in the K_L^0 section of the Stable Particle Table, and under Branching Ratio R4 in the K_S^0 section of the Stable Particle Data Card Listings. Our procedure is to assume that $x = 0$, and to list $(A_S/A_L)^2$ in the form of a branching ratio.

(b) *Charge asymmetry in $K_L \rightarrow 3\pi$ decays.* As mentioned above, the presence of a term in $(s_2 - s_1)$ in expression (1) describing the Dalitz plot distribution for τ^\pm, τ^0 decays of K mesons would be an indication of CP violation. Experimenters have used several forms for this CP-violation term. As described in the mini-review in the K^\pm section of the Stable Particle Data

Card Listings, we have converted all results to coefficient j in eq. (1) above. The latter is listed among the CP-violating parameters at the back of the K_L^0 section of the Stable Particle Data Card Listings. Note that only upper limits have been reported for this quantity.

(c) *Asymmetry in the $K_L \rightarrow \pi^+ \ell^+ \nu$ decays.* The quantity measured and compiled here is

$$\delta = \frac{\Gamma(K_L \rightarrow \pi^+ \ell^+ \nu) - \Gamma(K_L \rightarrow \pi^+ \ell^- \nu)}{\Gamma(K_L \rightarrow \pi^+ \ell^+ \nu) + \Gamma(K_L \rightarrow \pi^+ \ell^- \nu)}.$$

This asymmetry violates CP invariance. If CPT is good, for a pure K_L^0 beam, δ can be written as

$$\delta = 2[(1 - |x|^2)/(1 + |x|^2)] \operatorname{Re} \epsilon,$$

where x is the $\Delta S = \Delta Q$ -violating parameter defined in section B4, and ϵ is the parameter of the expansion

$$|K_L\rangle = [(1 + \epsilon)|K\rangle - (1 - \epsilon)|\bar{K}\rangle]/[2(1 + |\epsilon|^2)]^{1/2}, \quad (5a)$$

$$|K_S\rangle = [(1 + \epsilon)|K\rangle + (1 - \epsilon)|\bar{K}\rangle]/[2(1 + |\epsilon|^2)]^{1/2}. \quad (5b)$$

We give δ in the Addendum to the Stable Particle Table. In addition, in the K_L^0 CP-violation section of the Stable Particle Data Card Listings, we list δ separately for $K_L^0 \rightarrow \pi\mu\nu$ and $K_L^0 \rightarrow \pi e\nu$.

(d) $K_L \rightarrow 2\pi$ decay. The relevant parameters are

$$\eta_+ = A(K_L \rightarrow \pi^+ \pi^-)/A(K_S \rightarrow \pi^+ \pi^-)$$

$$= |\eta_+| \exp(i\phi_+).$$

$$\eta_{00} = A(K_L \rightarrow \pi^0 \pi^0)/A(K_S \rightarrow \pi^0 \pi^0)$$

$$= |\eta_{00}| \exp(i\phi_{00}).$$

ϵ , defined in eqs. (5) above, and

$$\epsilon' = \frac{1}{2} i\sqrt{2} \exp[i(\delta_2 - \delta_0)] \operatorname{Im}(A_2/A_0).$$

Here, A_j and δ_j are the amplitude and phase of $\pi\pi$ scattering at the K mass, defined by

$$\langle I = 0 | T | K \rangle = \exp(i\delta_0) A_0,$$

$$\langle I = 2 | T | K \rangle = \exp(i\delta_2) A_2.$$

Wu and Yang (1964) have derived the relationships

$$\eta_{+-} = \epsilon + \epsilon', \quad \eta_{00} = \epsilon - 2\epsilon'.$$

We give η_{+-} , η_{00} , ϕ_{+-} , and ϕ_{00} in the Addendum to the Stable Particle Table. The phases are measured directly, whereas the magnitudes η_{+-} and η_{00} are derived parameters. We use, as far as we can, the

directly measured quantities as input and calculate η_{+-} and η_{00} from the values given by our constrained fits. Therefore, if one looks at the Data Card Listings, most of the $|\eta|$ measurements appear in the form of branching ratios, with appropriate comments. We then give the values of η_{+-} and $|\eta_{00}|^2$ in a separate list at the end of the CP-violating parameters section of the K_L^0 section of the Stable Particle Data Card Listings.

4. $\Delta S = \Delta Q$ rule in K^0 decays. The relative amount of $\Delta S \neq \Delta Q$ component present is measured by the parameter x , defined as

$$x = A(\bar{K}^0 \rightarrow \pi^- \ell^+ \nu) / A(K^0 \rightarrow \pi^- \ell^+ \nu).$$

We list $\text{Re}\{x\}$ and $\text{Im}\{x\}$ for both K_{e3} and $K_{\mu 3}$ at the end of the Stable Particle Data Card Listings and give values in the Addendum to the Stable Particle Table.

C. η -decay parameters

1. C-violation in η decays. As a test of possible C-violation in electromagnetic interactions, a number of experiments have looked for possible charge asymmetries in the decays $\eta \rightarrow \pi^+ \pi^- \pi^0$ and $\eta \rightarrow \pi^+ \pi^- \gamma$. We list the following parameters:

(a) The left-right asymmetry

$$A = (N^+ - N^-) / (N^+ + N^-),$$

where $N^{(\pm)}$ means the number of events with the $\pi^{(\pm)}$ energy greater than the $\pi^{(\mp)}$ energy in the η rest frame.

(b) The sextant asymmetry

$$A_s = \frac{N_1 + N_3 + N_5 - N_2 - N_4 - N_6}{N_1 + N_2 + N_3 + N_4 + N_5 + N_6}$$

for the decay $\eta \rightarrow \pi^+ \pi^- \pi^0$. The numbers refer to the sextants of the Dalitz plot [see, for example, Layter (1972)]. A_s is sensitive to an $I = 0$ C-violating asymmetry.

(c) The quadrant asymmetry A_q , defined in a similar way as A_s , but with each sector of the Dalitz plot now containing $\pi/2$ rather than $\pi/3$ radians. A_q is sensitive to an $I = 2$ C-violating final state.

(d) The d-wave contribution to the C-violating amplitude in the decay $\eta \rightarrow \pi^+ \pi^- \gamma$. The upper limit for this contribution is measured by the parameter β , defined by

$$dN/d|\cos\theta| \propto \sin^2\theta(1 + \beta \cos^2\theta),$$

where θ is the angle between the π^+ and the γ in the di-pion center of mass. A term proportional to $\cos^2\theta$ could also be due to p- and f-wave interference.

We list A for the decay modes $\eta \rightarrow \pi^+ \pi^- \pi^0$ and $\eta \rightarrow \pi^+ \pi^- \gamma$, A_s and A_q for the decay $\eta \rightarrow \pi^+ \pi^- \pi^0$, and β for the decay $\eta \rightarrow \pi^+ \pi^- \gamma$ in the η section of the Stable Particle Data Card Listings.

2. Dalitz plot for $\eta \rightarrow \pi^+ \pi^- \pi^0$. The Dalitz plot for the decay $\eta \rightarrow \pi^+ \pi^- \pi^0$ may be fit by the distribution

$$|M(x, y)|^2 \propto 1 + ay + by^2 + cx + dx^2 + exy.$$

Here,

$$x = \sqrt{3}(T_+ - T_-)/Q, \quad y = (3T_0/Q) - 1,$$

T_+, T_-, T_0 are the kinetic energies of the $\pi^+, \pi^-,$ and π^0 in the η rest system, and $Q = m_\eta - m_{\pi^+} - m_{\pi^-} - m_{\pi^0}$. The coefficient of the term linear in x is sensitive to C-violation due to an $I = 0$ or $I = 2$ final state. We list papers presenting determinations of the parameters $a, b, c,$ and d in the η section of the Stable Particle Data Card Listings. However, we do not tabulate values of these parameters because the assumptions made by different authors are not compatible and do not allow comparison of the numerical values.

3. Dalitz plot for $\eta \rightarrow \pi^+ \pi^- \gamma$. The Dalitz plot for the decay $\eta \rightarrow \pi^+ \pi^- \gamma$ may be fit to the expression

$$|M|^2 \propto 1 + 2\alpha z,$$

where

$$z = \frac{2}{3} \sum_{i=1}^3 [3(m_\eta - 3m_\pi)^{-1}(E_i - \frac{1}{3}m_\eta)]^2 = \rho^2 / \rho_{\text{max}}^2.$$

Here E_i is the energy of the i th pion in the η rest frame, and ρ is the distance to the center of the Dalitz plot. We list the parameter α in the η section of the Stable Particle Data Card Listings.

D. Baryon-decay parameters

1. A/V ratio for baryon leptonic decays. Consider the decay

$$B_i \rightarrow B_f + \ell + \nu.$$

Assuming V, A theory, neglecting "induced" scalar, "induced" pseudoscalar, and axial weak-magnetism terms, and neglecting the q^2 dependence of the form factors, the baryon part of the matrix element for

these decays may be written [Goldberger and Treiman (1958)] as

$$\langle B_f | \gamma_\lambda (g_V - g_A \gamma_5) + (g_W/m_{B_i}) \sigma^{\lambda\nu} q_\nu | B_i \rangle,$$

where B_i and B_f represent initial and final baryons, g_A and g_V the axial and vector coupling constants, g_W the weak magnetism coupling constant, and q_ν the sum of the lepton momenta. Here the Pauli representation is used for the γ matrices. The ratio g_A/g_V may be written as

$$g_A/g_V = |g_A/g_V| \exp(i\phi),$$

where ϕ is $0 + n\pi$ if time reversal holds [see Jackson et al. (1957)].

Experiments on the leptonic decays of baryons other than the neutron have generally assumed ϕ to be either 0 or π , and have thus measured the magnitude and sign of g_A/g_V . In studying neutron beta decay, however, experiments have been sensitive enough to measure ϕ more precisely, and we include the phase angle in our Listings for this case. It is consistent with time-reversal invariance, and by using the above definition of the matrix element with the Pauli representations, the value of g_A/g_V in neutron beta decay is negative.

Due to statistical limitations the weak magnetism form factor g_W is usually assumed from CVC and SU(3), so that usually only g_A and g_V are determined experimentally. This determination is accomplished in a variety of ways.

(a) The lepton-neutrino angular correlation provides a measure of the absolute value of g_A/g_V [for relevant formulas, see, e.g., Albright (1959)].

(b) The up-down asymmetry of the lepton from polarized baryon decays provides a measure of g_A/g_V with its sign [for relevant formulas, see, e.g., Albright (1959)].

(c) The lepton spectrum, given enough statistics, provides a measure of g_A/g_V with its sign [for relevant formulas, see, e.g., Bender (1968)]. The lepton spectrum also provides a measure of g_W/g_A if the CVC-SU(3) assumption is relaxed.

(d) The polarization of the decay baryon, from polarized or unpolarized initial baryon, also provides g_A/g_V with its sign [for formulas, see, e.g., Willis and Thompson (1968)].

(e) The presence of a term proportional to

$$\sigma_{B_i} \cdot (\mathbf{p}_e \times \mathbf{p}_\nu),$$

xvi

where the initial baryon is polarized or

$$\sigma_{B_i} \cdot (\mathbf{p}_e \times \mathbf{p}_\nu),$$

where the polarization of the decay baryon is observed provides a measure of the deviation of ϕ from 0 or π , and is thus a test of time-reversal invariance [see, e.g., Willis and Thompson (1968)].

We compile the ratio g_A/g_V with its sign, for those decays for which it has been measured.

All the coupling constants and decay rates for baryon leptonic decays are related by Cabibbo's theory [Cabibbo (1963)], extended to six quarks (and three mixing angles) by Kobayashi and Maskawa (1973). A recent fit to this theory has been done by Shrock and Wang (1978).

2. *Asymmetry parameters in nonleptonic hyperon decays.* The transition matrix for the hyperon decay may be written as

$$M = s + p(\sigma \cdot \mathbf{q}), \tag{6}$$

where s and p are the parity-changing and the parity-conserving amplitudes, respectively; σ is the Pauli spin operator, and \mathbf{q} is a unit vector along the direction of the decay baryon in the hyperon rest frame.

The asymmetry parameters are defined by the relations

$$\alpha = 2 \operatorname{Re}(s^* p) / (|s|^2 + |p|^2),$$

$$\beta = 2 \operatorname{Im}(s^* p) / (|s|^2 + |p|^2),$$

$$\gamma = (|s|^2 - |p|^2) / (|s|^2 + |p|^2).$$

With the transition matrix (6), the angular distribution of the decay baryon, in the hyperon rest system, is of the form

$$I = 1 + \alpha P_Y \cdot \mathbf{q},$$

where $P_Y = \langle Y | \sigma | Y \rangle$ is the hyperon polarization.

In the notation of Lee and Yang (1957) the polarization P_B of the decay baryons is ^{1,2}

$$P_B = \frac{(\alpha + P_Y \cdot \mathbf{q})\mathbf{q} + \beta(P_Y \times \mathbf{q}) + \gamma\mathbf{q} \times (P_Y \times \mathbf{q})}{1 + \alpha P_Y \cdot \mathbf{q}},$$

^{1,2} Note that Lee and Yang (1957) contains a misprint. The minus sign in the definition of β should be replaced by a 2. In addition, our unit vector \mathbf{q} is the direction of the baryon, whereas their unit vector \mathbf{p} is the direction of the pion.

where P_B is defined in that rest system of the baryon obtained by a Lorentz transformation along q from the hyperon rest system in which q and P_Y are defined. Note that α is the helicity of the decay baryon for unpolarized hyperons.

The three parameters α , β , and γ satisfy the relation $\alpha^2 + \beta^2 + \gamma^2 = 1$.

It is then convenient to describe hyperon nonleptonic decays in terms of the two independent parameters α and the angle ϕ defined by

$$\beta = (1 - \alpha^2)^{1/2} \sin \phi,$$

$$\gamma = (1 - \alpha^2)^{1/2} \cos \phi,$$

which has a more nearly gaussian distribution than β or γ . Evidently

$$-\frac{1}{2}\pi \leq \phi \leq \frac{1}{2}\pi \quad \text{for } \gamma > 0,$$

$$+\frac{1}{2}\pi \leq \phi \leq \frac{3}{2}\pi \quad \text{for } \gamma < 0.$$

In discussing time-reversal invariance, the quantity of interest is Δ , defined by

$$\alpha = 2|s||p| \cos \Delta / (|s|^2 + |p|^2),$$

$$\beta = -2|s||p| \sin \Delta / (|s|^2 + |p|^2);$$

that is, Δ is the phase angle of s relative to p . Evidently

$$-\frac{1}{2}\pi \leq \Delta \leq \frac{1}{2}\pi \quad \text{for } \alpha > 0,$$

$$+\frac{1}{2}\pi \leq \Delta \leq \frac{3}{2}\pi \quad \text{for } \alpha < 0.$$

Under the assumption of time-reversal invariance, the angle Δ must satisfy the relation

$$\Delta = \delta_s - \delta_p,$$

modulo π , where δ_s and δ_p are the pion-baryon scattering phase shifts at the appropriate energy and for the appropriate isospin state. For Λ decay, assuming the validity of the $|\Delta I| = \frac{1}{2}$ rule,

$$\Delta = \delta_s - \delta_p = (7.0 \pm 1.0) \text{ deg}^{13}.$$

In the Stable Particle Data Card Listings we give α and

ϕ for each decay since they are the most closely related to the experiments and are essentially uncorrelated. Whenever necessary we have changed the signs of the reported values, so as to agree with our conventions. In the Stable Particle Table we give α , ϕ , and Δ with errors; and for convenience we also give the central value of γ , without an error.

VII. Statistical procedures

We divide this discussion on obtaining averages and errors into two sections:

A. The unconstrained case, or "simple averaging", and

B. The constrained case.

In what follows, the term "error" means one standard deviation (1σ); that is, for central value \bar{x} and error $\delta\bar{x}$, the range $\bar{x} \pm \delta\bar{x}$ constitutes a 68.3% confidence interval.

A. Unconstrained averaging

We are returning this year to the use of a standard gaussian procedure (with "scale factor") as our only method of averaging the data. The Student's distribution procedure, introduced in 1976 as a second method of averaging, has been discontinued. This results primarily from our observation that, although the data are better represented by a Student's distribution, the standard deviation (= the 68.3% confidence limit) of this Student's distribution turns out to be equal to the gaussian standard deviation. If one would choose to quote, e.g., 90% confidence limits, however, the gaussian procedure would give too small errors.

We begin by assuming that measurements of a given quantity obey a gaussian distribution, and thus we calculate a weighted average and error

$$\bar{x} \pm \delta\bar{x} = \left(\sum_i w_i x_i / \sum_i w_i \right) \pm \left(\sum_i w_i \right)^{-1/2},$$

$$w_i = [1/(\delta x_i)^2], \quad (1)$$

where x_i and δx_i are the value and error, respectively, reported by the i th experiment, and the sums run over N experiments. We also calculate χ^2 and compare it with its expectation value of $N - 1$.

If $\chi^2/(N - 1)$ is less than or equal to 1, and there are no known problems with the data, we accept the above results.

¹³ This value for $\delta_s - \delta_p$ is derived from the phase-shift analyses by Ayco (1976). The error is our estimation of the uncertainty allowing for possible correlations.

If $\chi^2/(N-1)$ is very large, or if there is prior knowledge of extremely large inconsistencies between experiments, we may choose not to average the data at all. Alternatively, we may quote the calculated average, but then give an educated guess as to the error; such a guess is generally a quite conservative estimate designed to take into account known problems with the data.

Finally, if $\chi^2/(N-1)$ is greater than 1, but not to such a large extent, we still average the data, but then try to make up for this fact in two ways:

(i) We plot an ideogram to guide the reader in deciding which data might be rejected before selected averages are made. An example of such an ideogram is given in fig. 3 below. Each experiment appearing in the plot is represented by a gaussian with central value x_i , error δx_i , and area proportional to $1/\delta x_i$. The choice of area is a somewhat arbitrary one: it is based on the assumption that an experimenter will work to reduce his (or her) systematic errors until they are slightly smaller (but seldom much smaller) than the statistical errors. Thus, as a bubble chamber physicist gets more events, he (or she) will use them both to reduce the statistical errors and to study the biases. Our

confidence that a significant systematic error has not been made in a given experiment, as compared with other contradictory experiments, then tends to go up as $1/\delta x_i$.

But why not assign a weight $1/\delta x_i^2$, as is done when computing a weighted average? We feel that this is equivalent to assuming that large systematic errors are as infrequent as large statistical fluctuations, and that this is unrealistic.

We emphasize the difference between least-squares averaging (where the weighting factor is the inverse square of the error) and the ideograms prepared for visual display. The former arithmetic is of course best if one has statistically distributed input, and yields a narrow gaussian distribution centered at the weighted mean. The ideogram (often multip peaked and certainly not gaussian) is based on the opposite hypothesis that some of the input is systematically in error. The idea behind least-squares averaging is that experiments 1, 2, 3, etc., are *all* valid (so we should multiply their probabilities). Our *ideograms* are based on the assumption that 1 or 2 or 3, etc., is valid, "hedged" with $1/\delta x_i$ betting odds; we then add their probabilities. Both approaches cannot simultaneously be right; we leave it to the reader to choose. A glance at the ideogram will show, however, that the discrepancy is often not severe for reasonably distributed input.

(ii) The second way in which we try to take account of $\chi^2/(N-1)$ being greater than 1 is to scale up our quoted error $\delta \bar{x}$ in eq. (1) by a factor

$$\text{SCALE} = [\chi^2/(N-1)]^{1/2}. \quad (2)$$

Our reasoning is as follows. Since we do not know which one or more of the experiments are wrong, we assume that all experimentalists underestimated their errors by the same scale factor (2). If we scale up all input errors by this factor, χ^2 returns to $N-1$, and of course the output error scales up by the same factor.

If all the experiments have errors of about the same size, the above (straightforward) procedure for calculating SCALE is carried out. If, however, we are to combine experiments with widely varying errors, we must modify the procedure slightly. This is because it is the more precise experiments that most influence not only the average value \bar{x} , but also the error $\delta \bar{x}$. Now, on the average, the low-precision experiments each contribute about unity to *both* the numerator and the denominator of SCALE, hence the χ^2 contri-

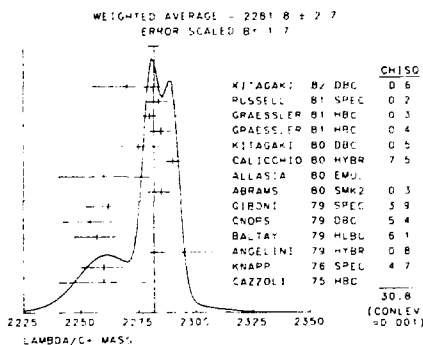


Fig. 3. Ideogram of measurements of the Λ_c^+ mass. The vertical line indicates the position of the weighted average, while the horizontal bar atop the line gives the error in the average after scaling by the SCALE factor. Only those experiments indicated by + error flags were precise enough to be accepted in the calculation of the SCALE factor; the column on the far right gives the χ^2 contribution of each of these experiments. The less precise experiments were included in the calculation of the weighted average, but not SCALE; they have \perp error flags.

bution of the sensitive experiments is diluted, i.e., reduced. Therefore, we evaluate SCALE by using *only* experiments for which the errors are not much greater than those of the more precise experiments. Explicitly, to calculate SCALE we use only the most sensitive experiments, i.e., those with errors less than δ_0 , where the ceiling δ_0 is (arbitrarily) chosen to be

$$\delta_0 = 3N^{1/2}\delta\bar{x}.$$

Here $\delta\bar{x}$ is the unscaled error of the mean of all the experiments. Note that if each experiment had the same error δx_i , then $\delta\bar{x}$ would be $\delta x_i/N^{1/2}$, so each individual experiment would be well under the ceiling on SCALE.

This scaling approach has the property that if there are two values with comparable errors separated by much more than their stated errors (with or without a number of other experiments of lower accuracy), the error on the mean value $\delta\bar{x}$ is increased so that it is approximately half the interval between the two discrepant values.

We wish to emphasize the fact that our scaling procedures for *errors* in no way affect central values. In addition, if one wishes to recover the unscaled error $\delta\bar{x}$, one need only divide the given error by the SCALE factor for that error.

B. Constrained fits

Except for trivial cases, all branching ratios and rate measurements are analyzed by the computer program AHR. This program makes a simultaneous least-squares fit to all the data, and outputs the partial-decay fractions \bar{P}_i , width Γ , partial widths Γ_i , and their error matrix.

The original version of AHR was written by J. Peter Berge. It is documented separately, and we wish here only to give the simplest nontrivial example that permits us to comment on the error matrix and the scale factor.

Assume that a state has only three partial-decay fractions, P_1 , P_2 , and P_3 ($\sum P_i = 1$), which have been measured in four different ratios, R_1, \dots, R_4 , where, e.g., $R_1 = P_1/P_2$, $R_2 = P_1/P_3$, etc.¹⁴ Further assume that *each* ratio has been measured by N experiments

¹⁴ We can handle any R of the form $R = \sum \alpha_i P_i / \sum \beta_j P_j$, where α_j and β_j are constants, usually 1 or 0. The forms $R = P_i \cdot P_j$ and $R = (P_i \cdot P_j)^{1/2}$ are also allowed.

(we designate each experiment with a subscript x , e.g., R_{1x}). Then AHR finds the best values of P_1, P_2 , and P_3 by minimizing χ^2 , namely

$$\chi^2 = \sum_{r=1}^4 \left[\sum_{x=1}^N \left(\frac{R_{rx} - R_r(P_1, P_2, P_3)}{\delta R_{rx}} \right)^2 \right]. \quad (3)$$

In addition to the fitted values \bar{P}_i , the program calculates an error matrix $\langle \delta \bar{P}_i \delta \bar{P}_j \rangle$. We tabulate the diagonal elements $\delta \bar{P}_i = \langle \delta \bar{P}_i \delta \bar{P}_i \rangle^{1/2}$ [except that some errors are scaled according to eq. (2) as discussed below]. In the listings we give the complete error matrix; we also calculate the fitted value of each ratio, for comparison with the input data, and list it below the relevant input, along with a simple unconstrained average of the same input.

Two further comments on the example above.

(1) There was no connection between measurements of the width and the branching ratios. But often we also have information on partial widths Γ_i as well as total width Γ . In this case AHR must introduce Γ as a parameter into the fit, along with the relations $\Gamma_i = \Gamma P_i$, $\sum \Gamma_i = \Gamma$. When appropriate, we tabulate the Γ_i along with the P_i , and give error matrices in the listings.

(2) Note that we do *not* allow for correlations between input data. We do try to pick those ratios and widths which are as independent and as close to the original data as possible.

In *asymmetric* errors, we use a continuous function of $\delta(P)^+$ and $\delta(P)^-$ in the fitting. When no errors are reported, we merely list the data for inspection.

Hyperon-decay parameters. The program AHR handles any type of input, α , ϕ , Δ , β , or γ , according to the definitions of section VI. If for a particular hyperon decay there are data for more than two of the decay parameters, they are analyzed by using the constraint

$$\alpha^2 + \beta^2 + \gamma^2 = 1.$$

Inconsistent constrained data. According to our simple example, which led to eq. (3), the double sum for χ^2 is summed over experiments $x = 1$ to N , leaving a single sum over ratios

$$\chi^2 = \sum_r \chi_r^2.$$

Even before fitting, some of the χ_r^2 may be too large. But if we scaled them before fitting, then the scaling would move the central value, contrary to our policy. So we do not scale until after the first fit; then, knowing the fitted χ_r^2 and its expectation value $\langle \chi_r^2 \rangle$ we form SCALE factors (just as before), i.e.,

$$(\text{SCALE})_r^2 = \chi_r^2 / \langle \chi_r^2 \rangle,$$

and if any $(\text{SCALE})_r$ is greater than 1, all N of the measurements of that particular ratio are equally penalized by having their errors increased by $(\text{SCALE})_r$. Program AHR then recycles on all the data, those with errors unchanged as well as those with errors increased. We then get new values, $\delta \bar{P}_i'$ for the errors in the partial-decay modes.

Because of the constraint ($\sum P_i = 1$) some SCALE factors may still be greater than 1 even after this second pass. If this is so, the whole procedure (i.e., increasing errors by the new SCALE factors and recycling through AHR) is repeated until AHR has converged.

At the end of AHR's final pass we have two measures of the errors for the \bar{P}_i . One is, of course, the $\delta \bar{P}_i'$, i.e., the errors in the final fitted values \bar{P}_i' which include the effects of scaling the input errors. The other measure of the errors is $(\bar{P}_i - \bar{P}_i')$, i.e., the *shift* in the central values of the i th mode between the first (unscaled) fit and the final (scaled) fit. In practice we find that on the average these two measures of the uncertainty are about equal. Rather than selecting just one or the other, our tabulated errors are given by the combination

$$(\delta \bar{P}_i)_{\text{tab}} = [\delta \bar{P}_i'^2 + (\bar{P}_i - \bar{P}_i')^2]^{1/2},$$

where \bar{P}_i is the fitted value of the i th partial-decay mode before scaling, \bar{P}_i' is its value after scaling, and $\delta \bar{P}_i'$ is the error in \bar{P}_i' . The SCALE factors we finally list in such cases are defined by

$$(\text{SCALE})_i = (\delta \bar{P}_i)_{\text{tab}} / \delta \bar{P}_i'.$$

However, in line with our policy of not letting SCALE affect the central values, we give the values of \bar{P}_i obtained from the original (unscaled) fits. [The differences between the \bar{P}_i calculated with either the scaled or the unscaled errors are, of course, always within the tabulated errors, $(\delta \bar{P}_i)_{\text{tab}}$.]

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TABLES OF PARTICLE PROPERTIES

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(Closing date for data: Jan. 1, 1982)

Stable Particle Table

For additional parameters, see Addendum to this table.

Quantities in italics have changed by more than one (old) standard deviation since Apr. 1 1980.

Particle	$I^G(J^P)C_n^a$	Mass ^o (MeV) Mass ² (GeV ²)	Mean life ^b (sec) $c\tau$ (cm)	Partial decay mode		ρ or P_{max}^c (MeV/c)
				Mode	Fraction ^b	
PHOTON						
γ	0,1(1 ⁻) ⁻	(< 6×10 ⁻²²)	-----	stable		
LEPTONS						
ν_e	$J=\frac{1}{2}$	(< 0.000046) ^d	stable (> 3×10 ⁸ m _{ν_e} (MeV))	stable		
e	$J=\frac{1}{2}$	0.5110034 ±0.0000014	stable (> ?×10 ²² y)	stable		
ν_μ	$J=\frac{1}{2}$	0 (< 0.52)	stable (> 1.1×10 ⁵ m _{ν_μ} (MeV))	stable		
μ	$J=\frac{1}{2}$	105.65943 ±0.00018 $m^2=0.01116392$	2.19714×10 ⁻⁶ ±0.00007 $c\tau=6.5868\times 10^4$	$\mu^- \rightarrow \mu^- \nu_\mu$ (or $\mu^+ \rightarrow CC$)		
				$e^- \bar{\nu}_e \nu_\mu$	(98.6 ± 0.4)%	53
				$e^- \bar{\nu}_e \nu_\mu \gamma$	(1.4 ± 0.4)%	53
				$\tau [e^- \nu_e \bar{\nu}_\mu]$	(< 9)%	53
				$e^- \gamma$	(< 1.9) × 10 ⁻¹⁰	53
$e^- e^+ e^-$	(< 1.9) × 10 ⁻⁹	53				
$e^- \gamma \gamma$	(< 5) × 10 ⁻⁸	53				
ν_τ	$J=\frac{1}{2}$	< 250				
τ	$J=\frac{1}{2}$	1784.2 ±3.2 $m^2=3.18$	(4.6±1.9)×10 ⁻¹³ $c\tau=0.014$	$\tau^- \rightarrow \tau^- \nu_\tau$ (or $\tau^+ \rightarrow CC$)		
				$\mu^- \bar{\nu}_\mu$	(18.5 ± 1.2)%	889
				$e^- \bar{\nu}_e$	(16.2 ± 1.0)%	892
				hadron ⁻ neutrals	(37.0 ± 3.2)%	
				3(hadron [±]) neutrals	(28.4 ± 3.0)%	
				5(hadron [±]) neutrals	(< 6)%	
				$\tau [3(\text{hadron}^\pm)\nu]$	(13 ± 8)%	
				3(hadron [±]) ν (≥1 γ)	(15 ± 7)%	
				$\dagger[\pi^- \nu]$	(10.7 ± 1.6)%	887
				$\rho^- \nu$	(21.6 ± 3.6)%	726
				K ⁻ neutrals	(small)	
				$\pi^- \pi^- \pi^+ \nu$	(7 ± 5)%	864
				$\pi^- \pi^- \pi^+ (\geq 0\pi^0)\nu$	(18 ± 7)%	864
$\dagger[K^{*-}(892)\nu]$	(1.7 ± 0.7)%	669				
$K^{*0}(1430)\nu$	(< 0.9)%	316				
$\pi^- \rho^0 \nu$	(5.4 ± 1.7)%	718				

(continued next page)

Stable Particle Table (cont'd)

Particle	$I^G(J^P)C_m^a$	Mass ^b (MeV) Mass ² (GeV ²)	Mean life ^b (sec) $c\tau$ (cm)	Partial decay mode		p or P_{max} (MeV/c)	
				Mode	Fraction ^b		
$\tau^- \rightarrow \mu^-$ (or $\tau^+ \rightarrow CC$)							
τ (continued)				e^- chgd.parts.			
				$+ \mu^-$ chgd.parts.	(<4)	(%)	
				$\mu^- \gamma$	(<5.5)	$\times 10^{-4}$	899
				$e^- \gamma$	(<6.4)	$\times 10^{-4}$	892
				$\mu^- \mu^+ \mu^-$	(<4.9)	$\times 10^{-4}$	876
				$e^- \mu^+ \mu^-$	(<3.3)	$\times 10^{-4}$	886
				$\mu^- e^+ e^-$	(<4.4)	$\times 10^{-4}$	889
				$e^- e^+ e^-$	(<4.0)	$\times 10^{-4}$	892
				$\mu^- \pi^0$	(<8.2)	$\times 10^{-4}$	884
				$e^- \pi^0$	(<2.1)	$\times 10^{-3}$	887
				$\mu^- K^0$	(<1.0)	$\times 10^{-3}$	819
				$e^- K^0$	(<1.3)	$\times 10^{-3}$	823
				$\mu^- \rho^0$	(<4.4)	$\times 10^{-4}$	722
				$e^- \rho^0$	(<3.7)	$\times 10^{-4}$	726
NONSTRANGE MESONS ^a							
$\pi^+ \rightarrow \mu^+$ (or $\pi^- \rightarrow CC$)							
π^\pm	$1^-(0^-)$	139.5673 ± 0.0007 $m^2=0.0194790$ $m_{\pi^\pm}-m_{\mu^\pm}=33.9079$ ± 0.0007	2.6030 $\times 10^{-8}$ ± 0.0023 $c\tau=780.4$ $(\tau^+ - \tau^-)/\bar{\tau} =$ $(0.05 \pm 0.07)\%$ (test of CPT)	$\mu^+ \nu$	100%		30
				$e^+ \nu$	$(1.267 \pm 0.023) \times 10^{-4}$		70
				$\mu^+ \nu \gamma$	$(1.24 \pm 0.25) \times 10^{-4}$		30
				$e^+ \nu \gamma$	$(5.6 \pm 0.7) \times 10^{-8}$		70
				$e^+ \nu \pi^0$	$(1.02 \pm 0.07) \times 10^{-8}$		5
				$e^+ \nu e^+ e^-$	< 5	$\times 10^{-9}$	70
π^0	$1^-(0^-)$	134.9630 ± 0.0038 $m^2=0.018715$ $m_{\pi^\pm}-m_{\pi^0}=4.6043$ ± 0.0037	0.83 $\times 10^{-16}$ ± 0.06 $S=1.8^*$ $c\tau=2.5 \times 10^{-6}$	$\gamma \gamma$	$(98.787 \pm 0.030)\%$		67
				$\gamma e^+ e^-$	(1.213)	(%)	67
				$\gamma \gamma \gamma$	< 3.8	$\times 10^{-7}$	67
				$e^+ e^- e^+ e^-$	(3.32)	$\times 10^{-5}$	67
				$\gamma \gamma \gamma \gamma$	< 4	$\times 10^{-6}$	67
				$e^+ e^-$	(2.2 ± 2.4)	$\times 10^{-7}$	67
				$\nu \nu$	< 2.4	$\times 10^{-5}$	67
η	$0^+(0^-)$	548.8 ± 0.6 $S=1.4^*$ $m^2=0.3012$	$\Gamma=(0.83 \pm 0.12) \text{keV}$ Neutral decays $(70.9 \pm 0.7)\%$	$\gamma \gamma$	$(39.1 \pm 0.8)\%$		274
				$3\pi^0$	$(31.8 \pm 0.8)\%$	$S=1.1^*$	180
				$\pi^0 \gamma \gamma$	< 0.3	(%)	258
				$\pi^+ \pi^- \pi^0$	$(23.7 \pm 0.5)\%$		175
				$\pi^+ \pi^- \gamma$	$(4.91 \pm 0.13)\%$		236
				$e^+ e^- \gamma$	$(0.50 \pm 0.12)\%$		274
				$\mu^+ \mu^- \gamma$	$(3.1 \pm 0.4) \times 10^{-4}$		253
				$e^+ e^-$	< 3	$\times 10^{-4}$	274
				$\mu^+ \mu^-$	$(6.5 \pm 2.1) \times 10^{-6}$		253
				$\pi^+ \pi^- e^+ e^-$	$(0.13 \pm 0.13)\%$		236
				$\pi^+ \pi^- \gamma \gamma$	< 0.21	(%)	236
				$\pi^+ \pi^- \pi^0 \gamma$	< 6	$\times 10^{-4}$	175
				$\pi^+ \pi^-$	< 0.15	(%)	236
				$\pi^0 e^+ e^-$	< 5	$\times 10^{-5}$	258
				$\pi^0 \mu^+ \mu^-$	< 5	$\times 10^{-6}$	211
				$\pi^0 \mu^+ \mu^- \gamma$	< 3	$\times 10^{-6}$	211

Stable Particle Table (cont'd)

Particle	$I^G(J^P)C_{\eta}$ ^a	Mass ^b (MeV) Mass ² (GeV ²)	Mean life ^b (sec) $c\tau$ (cm)	Partial decay mode		p or P _{max} ^c (MeV/c)
				Mode	Fraction ^b	
STRANGE MESONS^a						
$K^+ \bar{1}$ (or $K^- - \bar{1} - CC$)						
K^{\pm}	$\frac{1}{2}(0^-)$	493.667 ± 0.015 $m^2 = 0.2437$	1.2371 $\times 10^{-8}$ ± 0.0026 S=1.9 [*] $c\tau = 370.9$ $(\tau^+ - \tau^-)/\bar{\tau} =$ $(0.11 \pm 0.09)\%$ (test of CPT) S=1.2 [*]	$\mu^+ \nu$	(63.50 \pm 0.16)%	236
				$\pi^+ \pi^0$	(21.16 \pm 0.15)%	205
				$\pi^+ \pi^+ \pi^-$	(5.59 \pm 0.03)%	S=1.1 [*] 125
				$\pi^+ \pi^0 \pi^0$	(1.73 \pm 0.05)%	S=1.4 [*] 133
				$\pi^0 \mu^+ \nu$	(3.20 \pm 0.09)%	S=1.7 [*] 215
				$\pi^0 e^+ \nu$	(4.82 \pm 0.05)%	S=1.1 [*] 228
				$\mu^+ \nu \gamma$	ϵ (5.8 \pm 3.5) $\times 10^{-3}$	236
				$\pi^0 \pi^0 e^+ \nu$	(1.8 \pm $\frac{2.4}{0.6}$) $\times 10^{-5}$	207
				$\pi^+ \pi^- e^+ \nu$	(3.90 \pm 0.15) $\times 10^{-5}$	203
				$\pi^+ \pi^+ e^- \bar{\nu}$	(<1.2) $\times 10^{-8}$	203
				$\pi^+ \pi^- \mu^+ \nu$	(1.4 \pm 0.9) $\times 10^{-5}$	151
				$\pi^+ \pi^+ \mu^- \bar{\nu}$	(<3.0) $\times 10^{-6}$	151
				$e^+ \nu$	(1.54 \pm 0.07) $\times 10^{-5}$	247
				$e^+ \nu \gamma$ (SD+) ^h	(1.52 \pm 0.23) $\times 10^{-5}$	247
				$e^+ \nu \gamma$ (SD-) ^h	(<1.6) $\times 10^{-4}$	247
				$\pi^+ \pi^0 \gamma$	ϵ (2.75 \pm 0.16) $\times 10^{-4}$	205
				$\pi^+ \pi^+ \pi^- \gamma$	ϵ (1.0 \pm 0.4) $\times 10^{-4}$	125
				$\pi^0 \mu^+ \nu \gamma$	ϵ (<6) $\times 10^{-5}$	215
				$\pi^0 e^+ \nu \gamma$	ϵ (3.7 \pm 1.4) $\times 10^{-4}$	228
				$\pi^+ e^+ e^-$	(2.7 \pm 0.5) $\times 10^{-7}$	227
				$\pi^- e^+ e^+$	(<1) $\times 10^{-8}$	227
				$\pi^+ \mu^+ \mu^-$	(<2.4) $\times 10^{-6}$	172
				$\pi^+ \gamma \gamma$	ϵ (<3.5) $\times 10^{-5}$	227
				$\pi^+ \gamma \gamma \gamma$	ϵ (<3.0) $\times 10^{-4}$	227
				$\pi^+ \nu \bar{\nu}$	(<1.4) $\times 10^{-7}$	227
				$\pi^+ \gamma$	(<4) $\times 10^{-6}$	227
				$\pi^+ e^+ \mu^{\pm}$	(<7) $\times 10^{-9}$	214
				$\pi^+ e^- \mu^+$	(<5) $\times 10^{-9}$	214
				$e^+ \nu \nu \bar{\nu}$	(<6) $\times 10^{-5}$	247
				$\mu^+ \nu \nu \bar{\nu}$	(<6) $\times 10^{-6}$	236
				$\mu^+ \nu e^+ e^-$	(11 \pm 3) $\times 10^{-7}$	236
				$\mu^- \nu e^+ e^+$	(<2.0) $\times 10^{-8}$	236
$e^+ \nu e^+ e^-$	(2 \pm $\frac{2}{1}$) $\times 10^{-7}$	247				
$\mu^+ \nu_e$	(<4) $\times 10^{-3}$	236				
$\frac{K^0}{\bar{K}^0}$	$\frac{1}{2}(0^-)$	497.67 ± 0.13 S=1.1 [*] $m^2 = 0.24768$		50% K_{Short} , 50% K_{Long}		
K_S^0	$\frac{1}{2}(0^-)$		0.8923 $\times 10^{-10}$ ± 0.0022 $c\tau = 2.675$	$\pi^+ \pi^-$	(68.61 \pm 0.24)%	206
				$\pi^0 \pi^0$	(31.39 \pm 0.24)%	209
				$\pi^+ \pi^- \gamma$	ϵ (1.85 \pm 0.10) $\times 10^{-3}$	206
				$\mu^+ \mu^-$	(<3.2) $\times 10^{-7}$	225
				$e^+ e^-$	(<3.4) $\times 10^{-4}$	249
				$\gamma \gamma$	(<0.4) $\times 10^{-3}$	249

Stable Particle Table (cont'd)

Particle	$\Gamma^G(J^P)C_n^a$	Mass ^b (MeV) Mass ² (GeV ²)	Mean life ^b (sec) $c\tau$ (cm)	Partial decay mode		P or P _{max} ^c (MeV/c)				
				Mode	Fraction ^b					
K_L^0	$\frac{1}{2}(0^-)$	$m_{K_L^0} - m_{K_S^0} = 0.5349 \times 10^{10} \text{ } \hbar \text{ sec}^{-1} \pm 0.0022$	$5.183 \times 10^{-8} \pm 0.040$ $c\tau = 1554$	$\pi^0 \pi^0 \pi^0$	(21.5 ± 1.0)% S=1.7*	139				
				$\pi^+ \pi^- \pi^0$	(12.39 ± 0.20)% S=1.3*	133				
				$\pi^+ \pi^- \mu^+ \nu$	(27.1 ± 0.4)% S=1.4*	216				
				$\pi^+ \pi^- e^+ \nu$	(38.7 ± 0.5)% S=1.5*	229				
				$\dagger[\pi e \nu \gamma]$	$^e(1.3 \pm 0.8)\%$	229				
				$\pi^+ \pi^-$	$^j(0.203 \pm 0.005)\%$ S=1.1*	206				
				$\pi^0 \pi^0$	$^j(0.094 \pm 0.018)\%$ S=1.5*	209				
				$\pi^+ \pi^- \gamma$	$^e(4.41 \pm 0.32) \times 10^{-5}$	206				
				$\pi^0 \gamma \gamma$	(< 2.4) $\times 10^{-4}$	231				
				$\gamma \gamma$	(4.9 ± 0.4) $\times 10^{-4}$	249				
				$e \mu$	$^k(< 6) \times 10^{-6}$	238				
				$\mu^+ \mu^-$	(9.1 ± 1.9) $\times 10^{-9}$	225				
				$\mu^+ \mu^- \gamma$	(2.8 ± 2.8) $\times 10^{-7}$	225				
				$^j \mu^+ \mu^-$	(1.2) $\times 10^{-7}$	177				
				$e^+ e^-$	$^k(< 2.0) \times 10^{-7}$	249				
				$e^+ e^- \gamma$	(1.7 ± 0.9) $\times 10^{-5}$	249				
				$\pi^0 e^+ e^-$	(< 2.3) $\times 10^{-6}$	231				
				$\pi^+ \pi^- e^+ e^-$	(< 8.8) $\times 10^{-6}$	206				
				$\pi^0 \pi^\pm e^\mp \nu$	(6.2 ± 2.0) $\times 10^{-5}$	207				
				CHARMED NONSTRANGE MESONS^a						
D^\pm	$\frac{1}{2}(0^-)$	1869.4 ± 0.6 $m^2 = 3.495$ $m_{D^\pm} - m_{D^0} = 4.7 \pm 0.3$	$(9.1_{-1.3}^{+2.2}) \times 10^{-13}$ $c\tau = 0.027$	$D^{+,-} \rightarrow (\text{or } D^- \rightarrow \text{CC})$	$e^\pm \text{ anything}$ (19 \pm 4)%					
				$K^- \text{ anything}$ (16 ± 4)%						
				$\bar{K}^0 \text{ any} + K^0 \text{ any}$ (48 ± 15)%						
				$K^+ \text{ anything}$ (6.0 ± 3.3)%						
				$\eta \text{ anything}$ $^e(< 13)\%$						
				$\dagger[K^- \pi^+ \pi^+]$	(4.6 ± 1.1)%	845				
				$K^- \pi^+ \pi^+ \pi^+ \pi^-$	(< 4)%	772				
				$\bar{K}^0 \pi^+$	(1.8 ± 0.5)%	862				
				$\bar{K}^0 \pi^+ \pi^0$	(13 ± 8)%	845				
				$\bar{K}^0 \pi^+ \pi^+ \pi^-$	(8.4 ± 3.5)%	814				
				$\bar{K}^0 K^+$	(0.45 ± 0.30)%	792				
				$K^+ K^- \pi^+$	(< 0.6)%	744				
				$K^+ \pi^+ \pi^-$	(< 0.23)%	845				
				$\pi^+ \pi^0$	(< 0.5)%	925				
				$\pi^+ \pi^+ \pi^-$	(< 0.37)%	908				
				$\dagger[\bar{K}^0 \pi^+]$	(< 3.7)%	714				
				D^0	$\frac{1}{2}(0^-)$	1864.7 ± 0.6 $m^2 = 3.477$ $\frac{\Gamma(D^0 \rightarrow \bar{D}^0 \rightarrow K^+ \pi^-)}{\Gamma(D^0 \rightarrow K \pi)} < 0.16$	$(4.8_{-1.3}^{+2.4}) \times 10^{-13}$ $c\tau = 0.014$	$D^0 \rightarrow (\text{or } \bar{D}^0 \rightarrow \text{CC})$	$e^\pm \text{ anything}$ (< 6)%	
								$K^- \text{ anything}$ (44 ± 10)% S=1.3*		
								$\bar{K}^0 \text{ any} + K^0 \text{ any}$ (33 ± 10)%		
								$K^+ \text{ anything}$ (8 ± 3)%		
$\eta \text{ anything}$ $^e(< 13)\%$										
$\dagger[K^- \pi^+]$	(2.4 ± 0.4)%	861								
$K^- \pi^+ \pi^0$	(9.3 ± 2.8)%	844								
$K^- \pi^+ \pi^+ \pi^-$	(4.5 ± 1.3)% S=1.4*	812								
$K^- \pi^+ \pi^0 \pi^0$	(seen)	815								
$\bar{K}^0 \pi^0$	(2.2 ± 1.1)%	860								
$\bar{K}^0 \pi^+ \pi^-$	(4.2 ± 0.8)%	842								
$\pi^+ \pi^-$	(7.9 ± 3.8) $\times 10^{-4}$	922								
$\pi^+ \pi^+ \pi^- \pi^-$	(< 9) $\times 10^{-4}$	768								
$K^+ K^-$	(2.7 ± 0.8) $\times 10^{-3}$	791								
$\dagger[K^+ \pi^+]$	(3.4 ± 1.4)%	711								
$\bar{K}^0 \pi^0$	(1.4 \pm 2.3)%	711								
$K^- \rho^+$	(7.2 \pm 3.0)%	679								
$\bar{K}^0 \rho^0$	(0.1 \pm 0.6)%	677								

Stable Particle Table (cont'd)

Particle	$I^G(J^P)C_n^a$	Mass ^b (MeV) Mass ² (GeV ²)	Mean life ^b (sec) $c\tau$ (cm)	Partial decay mode		
				Mode	Fraction ^b	p or p_{max} (MeV/c) ^c
CHARMED STRANGE MESON^a						
F^\pm	$0(0^-)^m$	2021 ± 15	$(2.2^{+2.8}_{-1.1}) \times 10^{-13}$	$F^+ \rightarrow \eta\pi^+$ (or $F^- \rightarrow CC$)		
				$\eta\pi^+\pi^+\pi^-$	(seen)	930
				$\eta'\pi^+\pi^+\pi^-$	(seen)	885
				$\rho^+\phi$	(seen)	467
$\rightarrow B$						
NONSTRANGE BARYONS^a						
p	$\frac{1}{2}(\frac{1}{2}^+)$	938.2796 ± 0.0027 $m^2=0.880369$	stable ($\geq 8 \times 10^{30}y$)	stable		
				$ q_p - q_c < 10^{-21} q_c ^n$		
n	$\frac{1}{2}(\frac{1}{2}^+)$	939.5731 ± 0.0027 $m^2=0.882798$ $m_p - m_n = -1.29343$ ± 0.00004	925 \pm 11 $c\tau = 2.77 \times 10^{13}$	$pe^-\bar{\nu}$	100%	1
				$p\nu\bar{\nu}$ (chg.noncons.)	(<9) $\times 10^{-24}$	1
				$ q_n < 10^{-21} q_c ^n$		
STRANGENESS -1 BARYONS^a						
Λ	$0(\frac{1}{2}^+)$	1115.60 ± 0.05 $S=1.2^*$ $m^2=1.2446$ $m_\Lambda - m_{\Sigma^0} = -76.86$ ± 0.08	2.632×10^{-10} ± 0.020 $S=1.6^*$ $c\tau = 7.89$	$p\pi^-$	(64.2 \pm 0.5)%	100
				$n\pi^0$	(35.8 \pm 0.5)%	104
				$pe^-\nu$	(8.35 \pm 0.15) $\times 10^{-4}$	163
				$p\mu^-\nu$	(1.57 \pm 0.35) $\times 10^{-4}$	131
				$p\pi^-\gamma$	(8.5 \pm 1.4) $\times 10^{-4}$	100
Σ^+	$1(\frac{1}{2}^+)$	1189.36 ± 0.06 $S=1.8^*$ $m^2=1.4146$ $m_{\Sigma^+} - m_{\Sigma^-} = -7.97$ ± 0.07 $S=1.3^*$	0.800×10^{-10} ± 0.004 $c\tau = 2.40$	$p\pi^0$	(51.64 \pm 0.30)%	189
				$n\pi^+$	(48.36 \pm 0.30)%	185
				$p\gamma$	(1.20 \pm 0.13) $\times 10^{-3}$	$S=1.4^*$ 225
				$n\pi^+\gamma$	(4.5 \pm 0.5) $\times 10^{-4}$	185
				$\Delta c^+\nu$	(2.0 \pm 0.5) $\times 10^{-5}$	71
$n\mu^+\nu$	(<3.0) $\times 10^{-5}$	202				
$ne^+\nu$	(<5) $\times 10^{-6}$	224				
pe^+e^-	(<7) $\times 10^{-6}$	225				
				$\frac{\Gamma(\Sigma^+ \rightarrow \ell^+ n\nu)}{\Gamma(\Sigma^- \rightarrow \ell^- n\nu)} < .04$		
Σ^0	$1(\frac{1}{2}^+)^p$	1192.46 ± 0.08 $m^2=1.4220$	± 1.3 $c\tau = 1.7 \times 10^{-9}$	$\Lambda\gamma$	100%	74
				Δc^+e^-	(5.45) $\times 10^{-3}$	74
				$\Delta\gamma\gamma$	(<3) %	74
Σ^-	$1(\frac{1}{2}^+)$	1197.34 ± 0.05 $m^2=1.4336$ $m_{\Sigma^0} - m_{\Sigma^-} = -4.88$ ± 0.06	1.482×10^{-10} ± 0.011 $S=1.3^*$ $c\tau = 4.44$	$n\pi^-$	100%	193
				$ne^-\nu$	(1.08 \pm 0.04) $\times 10^{-3}$	230
				$n\mu^-\nu$	(0.45 \pm 0.04) $\times 10^{-3}$	210
				$\Delta c^-\nu$	(0.61 \pm 0.05) $\times 10^{-4}$	79
				$n\pi^-\gamma$	(4.6 \pm 0.6) $\times 10^{-4}$	193

Stable Particle Table (cont'd)

Particle	$I^G(J^P)C_u^a$	Mass ^b (MeV) Mass ² (GeV ²)	Mean life ^b (sec) cr (cm)	Partial decay mode		P or P _{max} ^c (MeV/c)		
				Mode	Fraction ^b			
STRANGENESS -2 BARYONS^a								
Σ^0	$\frac{1}{2}(\frac{1}{2}^+)^g$	1514.9	2.90×10^{-10}	$\Lambda\pi^0$	100%	135		
		± 0.6	± 0.10	$\Lambda\gamma$	(0.5 \pm 0.5)%	184		
		$m^2=1.729$	$cr=8.69$	$\Sigma^0\gamma$	(<7)%	117		
				$p\pi^-$	(<3.6) $\times 10^{-5}$	299		
				$pe^- \nu$	(<1.3) $\times 10^{-3}$	323		
		$m_{\Sigma^0}-m_{\Sigma^-}=-6.4$		$\Sigma^+ e^- \nu$	(<1.1) $\times 10^{-3}$	120		
		± 0.6		$\Sigma^- e^+ \nu$	(<0.9) $\times 10^{-3}$	112		
				$\Sigma^+ \mu^- \nu$	(<1.1) $\times 10^{-3}$	65		
				$\Sigma^- \mu^+ \nu$	(<0.9) $\times 10^{-3}$	49		
				$p\mu^- \nu$	(<1.3) $\times 10^{-3}$	309		
		Σ^-	$\frac{1}{2}(\frac{1}{2}^+)^g$	1321.32	1.641×10^{-10}	$\Lambda\pi^-$	100%	139
± 0.13	± 0.016			$\Lambda e^- \nu$	(2.9 \pm 1.1) $\times 10^{-4}$	190		
$m^2=1.7459$	$cr=4.92$			$\Sigma^0 e^- \nu$	(<1.4) $\times 10^{-4}$	123		
				$\Lambda\mu^- \nu$	(3.5 \pm 3.5) $\times 10^{-4}$	163		
				$\Sigma^0 \mu^- \nu$	(<8) $\times 10^{-4}$	70		
				$n\pi^-$	(<1.1) $\times 10^{-3}$	303		
				$ne^- \nu$	(<3.2) $\times 10^{-3}$	327		
				$n\mu^- \nu$	(<1.5)%	313		
				$\Sigma^- \gamma$	(<1.2) $\times 10^{-3}$	118		
				$p\pi^- \pi^-$	(<4) $\times 10^{-4}$	223		
				$p\pi^- e^- \nu$	(<4) $\times 10^{-4}$	304		
				$p\pi^- \mu^- \nu$	(<4) $\times 10^{-4}$	250		
				$\Xi^0 e^- \nu$	(<2.3) $\times 10^{-3}$	6		
STRANGENESS -3 BARYON^a								
Ω^-	$\alpha(\frac{3}{2}^+)^g$			1672.45	0.819×10^{-10}	ΛK^-	(68.6 \pm 1.3)%	211
		± 0.32	± 0.027	$\Xi^0 \pi^-$	(23.4 \pm 1.3)%	294		
		$m^2=2.7971$	$cr=2.46$	$\Xi^- \pi^0$	(8.0 \pm 0.8)%	290		
				$\Xi^0 e^- \nu$	(\sim 1)%	319		
				$\Xi^0(1530)\pi^-$	(\sim 2) $\times 10^{-3}$			
				$\Lambda\pi^-$	(<1.3) $\times 10^{-3}$	449		
				$\Xi^- \gamma$	(<3.1) $\times 10^{-3}$	314		
NONSTRANGE CHARMED BARYON^a								
Δ_c^+	$\alpha(\frac{1}{2}^+)^g$	2282.2	$(1.1^{+0.9}_{-0.4}) \times 10^{-13}$	$pK^- \pi^+$	(2.2 \pm 1.0)%	820		
		± 3.1	$cr=0.003$	$p\bar{K}^0$	(1.1 \pm 0.7)%	870		
		$S=1.8^*$		$p\bar{K}^0 \pi^+ \pi^-$	(<4)%	751		
		$m^2=5.21$		$\Delta\pi^+$	(0.6 \pm 0.5)%	861		
				$\Lambda\pi^+ \pi^+ \pi^-$	(<3.1, seen)%	804		
				$\Sigma^0 \pi^+$	(seen)	822		
				$\dagger[pK^0]$	(0.48 \pm 0.30)%	682		
				$\Delta^{++} K^-$	(0.45 \pm 0.27)%	707		
				$pK^+ \pi^+$	(seen)	576		
				e^+ anything	(4.5 \pm 1.7)%			
				$\dagger[pe^+]$ anything	(1.8 \pm 0.9)%			
		Λe^+ anything	(1.1 \pm 0.8)%					
<p>$\rightarrow \Delta_b^0$</p> <p>\rightarrow searches for massive neutrinos and lepton mixing</p> <p>$\rightarrow \nu$ bounds from astrophysics and cosmology</p> <p>\rightarrow heavy lepton searches</p> <p>\rightarrow weak gauge boson searches</p> <p>\rightarrow free quark searches</p> <p>\rightarrow magnetic monopole searches</p> <p>\rightarrow charm searches and evidence</p> <p>\rightarrow bottom hadron searches</p> <p>\rightarrow top hadron searches</p> <p>\rightarrow other stable particle searches</p>								

ADDENDUM TO

Stable Particle Table

Magnetic Moment		μ Decay parameters ^f				
e^w	1.001 159 652 209 $\pm .000\ 000\ 000\ 031$	$\frac{e\hbar}{2m_e c}$				
μ^w	1.001 165 924 $\pm .000\ 000\ 009$	$\frac{e\hbar}{2m_\mu c}$	$\rho = 0.752 \pm 0.003$ $\xi: P = 0.972 \pm 0.014$ $ g_A/g_V = 0.86^{+0.33}_{-0.11}$	$\eta = -0.12 \pm 0.21$ $\delta = 0.755 \pm 0.009$ $\phi = 180^\circ \pm 15^\circ$	$h = 1.01 \pm 0.06$	
η	Mode $\pi^+\pi^-\pi^0$ $\pi^+\pi^-\gamma$	Left-right asymmetry (0.12 \pm .17)% (0.88 \pm .40)%	Sextant asymmetry (0.19 \pm 0.16)%	Quadrant asymmetry (-0.17 \pm 0.17)%	$\beta = 0.047 \pm 0.062$	
K^\pm	Mode	Partial rate (sec ⁻¹)		Slope parameters for $K \rightarrow 3\pi^f$		
	$\mu\nu$ $\pi\pi^0$ $\pi\pi^+\pi^-$ $\pi\pi^0\pi^0$ $\mu\pi^0\nu$ $e\pi^0\nu$	(51.33 \pm 0.17) $\times 10^6$ (17.10 \pm 0.13) $\times 10^6$ (4.52 \pm 0.02) $\times 10^6$ (1.40 \pm 0.04) $\times 10^6$ (2.58 \pm 0.07) $\times 10^6$ (3.90 \pm 0.04) $\times 10^6$	S=1.2* S=1.1* S=1.1* S=1.4* S=1.7* S=1.1*	$K^+\pi^+\pi^+\pi^-$ $g = -0.215 \pm 0.004$ S=1.4* $K^-\pi^-\pi^-\pi^+$ $g = -0.217 \pm 0.007$ S=2.5* $K^\pm \rightarrow \pi^0\pi^0\pi^\pm$ $g = 0.607 \pm 0.030$ S=1.3* $K_{L,S}^0 \rightarrow \pi^+\pi^-\pi^0$ $g = 0.670 \pm 0.014$ S=1.6* See Data Card Listings for quadratic coefficients.		
K_S^0	$\pi^+\pi^-$ $\pi^0\pi^0$	J (0.7689 \pm .0033) $\times 10^{10}$ J (0.3517 \pm .0029) $\times 10^{10}$	S=1.1*	Form factors for $K_{L,S}^0$ decays		
	K_L^0	$\pi^0\pi^0\pi^0$ (4.14 \pm 0.20) $\times 10^6$ $\pi^+\pi^-\pi^0$ (2.39 \pm 0.04) $\times 10^6$ $\pi\mu\nu$ (5.23 \pm 0.09) $\times 10^6$ $\pi e\nu$ (7.47 \pm 0.11) $\times 10^6$ $\pi^+\pi^-$ J (3.91 \pm 0.10) $\times 10^4$ $\pi^0\pi^0$ J (1.81 \pm 0.36) $\times 10^4$	S=1.7* S=1.2* S=1.3* S=1.3* S=1.1* S=1.5*	$K^+ \begin{cases} \lambda_+^0 = 0.029 \pm 0.004 \\ \lambda_+^\pm = 0.026 \pm 0.008 \\ \lambda_+^0 = -0.003 \pm 0.007 \end{cases}$ S=1.5* $K_S^0 \begin{cases} \lambda_+^0 = 0.0300 \pm 0.0016 \\ \lambda_+^\pm = 0.034 \pm 0.006 \\ \lambda_+^0 = 0.020 \pm 0.007 \end{cases}$ S=1.2* S=2.5* See Data Card Listings for ξ , f_+ , and f_0 .		
CP violation parameters ^{u,j}						
$ \eta_{+-} = (2.27 \pm 0.022) \times 10^{-3}$ $ \eta_{00} = (2.33 \pm 0.08) \times 10^{-3}$ S=1.1* $\phi_{+-} = (44.6 \pm 1.2)^\circ$ $\phi_{00} = (54 \pm 5)^\circ$ $ \eta_{+-0} ^2 < 0.1$ $ \eta_{000} ^2 < 0.28$ $\delta = (0.330 \pm 0.012)^\circ$						
$\Delta S = -\Delta Q$ $\text{Re } x = 0.009 \pm 0.020$ S=1.4* $\text{Im } x = -0.004 \pm 0.026$ S=1.1*						
Magnetic moment ($e\hbar/2m_p c$)	Decay parameters ^v				Coupling Constant Ratios	
	Measured		Derived			
p	2.7928456 $\pm .0000011$	α	ϕ (degree)	γ	Δ (degree)	
n^w	-1.91304184 $\pm .00000088$	$pe^-\nu$				$g_A/g_V = -1.255 \pm 0.006$ $\phi_{AV} = (180.11 \pm 0.17)^\circ$
Λ^w	-0.613 $\pm .004$	$p\pi^-$ 0.642 \pm 0.013 $n\pi^0$ 0.646 \pm 0.044 $pe\nu$	(-6.5 \pm 3.5) $^\circ$	0.76	(7.7 \pm 4.1) $^\circ$	$g_A/g_V = -0.690 \pm 0.034$ S=1.4*
Σ^+	2.33 $\pm .13$	$p\pi^0$ -0.979 \pm 0.016 $n\pi^+$ +0.068 \pm 0.013 $p\gamma$ -0.72 \pm 0.29	(36 \pm 34) $^\circ$ (167 \pm 20) $^\circ$ S=1.1*	0.17 -0.97	(187 \pm 6) $^\circ$ (-73 $^{+13}_{-14}$) $^\circ$	
Σ^-	-1.41 $\pm .25$	$n\pi^-$ -0.068 \pm 0.008 $ne^-\nu$ $\Lambda e^-\nu$	(10 \pm 15) $^\circ$	0.98	(249 $^{+12}_{-18}$) $^\circ$	$g_A/g_V = \pm(0.385 \pm 0.070)$ S=2.3* $g_V/g_A = 0.14 \pm 0.24$ S=1.6* $g_{WM}/g_A = 2.4 \pm 1.7$
Ξ^0	-1.250 $\pm .014$	$\Lambda\pi^0$	-0.113 \pm 0.022 S=2.0*	(21 \pm 12) $^\circ$	0.85	(218 $^{+12}_{-18}$) $^\circ$
Ξ^-	-1.85 $\pm .75$	$\Lambda\pi^-$	-0.434 \pm 0.015 S=1.4* S=1.1*	(2 \pm 6) $^\circ$	0.90	(184 \pm 12) $^\circ$
Ω^-		ΛK^-	-0.10 \pm 0.38 S=1.2*			

Stable Particle Table (cont'd)

→ Indicates an entry in the Stable Particle Data Card Listings not entered in the Stable Particle Table.

* S = Scale factor = $\sqrt{\chi^2/(N-1)}$, where $N \approx$ number of experiments. S should be ≈ 1 . If $S > 1$, we have enlarged the error of the mean, $\delta\bar{x}$; i.e., $\delta\bar{x} \rightarrow S\delta\bar{x}$. This convention is still inadequate, since if $S \gg 1$ the experiments are probably inconsistent, and therefore the real uncertainty is probably even greater than $S\delta\bar{x}$. See text, and ideograms in Stable Particle Data Card Listings.

† Square brackets indicate subreactions of some previous unbracketed decay mode(s). Reactions in one set of brackets may overlap with reactions in another set of brackets.

a. The baryon number B, strangeness S, and charm C of the hadrons which appear in the tables are as follows:

Mesons (B=0)	S	C	Baryons (B=1)	S	C
π, η	0	0	p, n	0	0
K^+, K^0	+1	0	Λ, Σ	-1	0
K^-, \bar{K}^0	-1	0	Ξ	-2	0
D^+, D^0	0	+1	Ω^-	-3	0
D^-, \bar{D}^0	0	-1	Λ_c^+	0	+1
F^+	+1	+1			
F^-	-1	-1			

b. Quoted upper limits correspond to a 90% confidence level.

c. In decays with more than two bodies, p_{\max} is the maximum momentum that any particle can have.

d. 99% confidence level. Lower limit from same experiment, > 14 eV, not yet confirmed. See Stable Particle Data Card Listings.

e. See Stable Particle Data Card Listings for energy limits used in this measurement.

f. Theoretical value; see also Stable Particle Data Card Listings.

g. See note in Stable Particle Data Card Listings.

h. Structure-dependent part with positive (SD+) and negative (SD-) photon helicity.

i. The direct emission branching fraction is $(1.56 \pm .3) \times 10^{-5}$.

j. The $K_S^0 \rightarrow \pi\pi$ and $K_L^0 \rightarrow \pi\pi$ rates (and branching fractions) are from our branching fraction and rate fits and do not include results of $K_L^0 - K_S^0$ interference experiments. The $|\eta_{+-}|$ and $|\eta_{00}|$ values given in the addendum are these rates combined with the $|\eta_{+-}|$ and $|\eta_{00}|$ results from interference experiments.

k. The stronger limit $< 2 \times 10^{-9}$ of Clark et al., Phys. Rev. Lett. 26, 1667 (1971) is not listed because of possible (but unknown) systematic errors. See Stable Particle Data Card Listings.

l. This is a weighted average of D^{\pm} (44%) and D^0 (56%) branching fractions.

m. Quantum numbers shown are favored but not yet established. See Stable Particle Data Card Listings.

n. Limit from neutrality-of-matter experiments. Assumes $|q_n| = |q_p| - |q_e|$.

p. J^P not measured for Σ^0 . Assumed same as Σ^{\pm} to allow isotriplet association.

q. P for Ξ and J^P for Ω^- not yet measured. Values shown are SU(3) predictions.

r. J^P for Λ_c^+ not yet measured. Values shown are SU(4) predictions.

s. $|g_A/g_V|$ defined by $g_A^2 = |C_A|^2 + |C_A'|^2$, $g_V^2 = |C_V|^2 + |C_V'|^2$, and $\Sigma < \bar{\epsilon} | \Gamma_1 | \mu > < \bar{\nu} | \Gamma_1 (C_1 + C_1' \gamma_5) | \nu >$; ϕ defined by $\cos \phi = -\text{Re}(C_A^* C_V + C_A' C_V') / g_A g_V$, P_μ is muon longitudinal polarization [for more details, see text Section VI A].

t. The definition of the slope parameter of the Dalitz plot is as follows [see also text Section VI B.1]: $|M|^2 = 1 + g \left(\frac{s_3 - s_0}{m_\pi^2} \right)$

u. The definition for the CP violation parameters is as follows [see also text Section VI B.3]:

$$\eta_{+-} = |\eta_{+-}| e^{i\phi_{+-}} = \frac{A(K_L^0 \rightarrow \pi^+ \pi^-)}{A(K_S^0 \rightarrow \pi^+ \pi^-)} \quad \eta_{00} = |\eta_{00}| e^{i\phi_{00}} = \frac{A(K_L^0 \rightarrow \pi^0 \pi^0)}{A(K_S^0 \rightarrow \pi^0 \pi^0)}$$

$$\delta = \frac{\Gamma(K_L^0 \rightarrow \ell^+ \ell^-) - \Gamma(K_L^0 \rightarrow \ell^- \ell^-)}{\Gamma(K_L^0 \rightarrow \ell^+ \ell^-) + \Gamma(K_L^0 \rightarrow \ell^- \ell^-)}, \quad |\eta_{+-0}|^2 = \frac{\Gamma(K_S^0 \rightarrow \pi^+ \pi^- \pi^0) \text{CP viol.}}{\Gamma(K_L^0 \rightarrow \pi^+ \pi^- \pi^0)}, \quad |\eta_{000}|^2 = \frac{\Gamma(K_S^0 \rightarrow \pi^0 \pi^0 \pi^0) \text{CP viol.}}{\Gamma(K_L^0 \rightarrow \pi^0 \pi^0 \pi^0)}$$

v. The definition of these quantities is as follows [for more details on sign convention, see text Section VI B]:

$$\alpha = \frac{2|s| |p| \cos \Delta}{|s|^2 + |p|^2} \quad \beta = \sqrt{1 - \alpha^2} \sin \phi \quad g_A, g_V, g_{WM} \text{ defined by } \langle B_1 | \gamma_\lambda (g_V - g_A \gamma_5) + (g_{WM}/m_{B_1}) \sigma^{\lambda\nu} q_\nu | B_1 \rangle$$

$$\beta = \frac{-2|s| |p| \sin \Delta}{|s|^2 + |p|^2} \quad \gamma = \sqrt{1 - \alpha^2} \cos \phi \quad \phi_{AV} \text{ defined by } g_A/g_V = |g_A/g_V| e^{i\phi_{AV}}$$

w. For limits on electric dipole moment, see Data Card Listings. Forbidden by P and T invariance.

Meson Table

April 1982

In addition to the entries in the Meson Table, the Meson Data Card Listings contain all substantial claims for meson resonances. See contents of Meson Data Card Listings below.

Quantities in italics are new or have changed by more than one (old) standard deviation since April 1980.

J ^P	G	I	0	1	1/2	I ^G (J ^P)C _u	Mass M (MeV)	Full Width Γ (MeV)	M ² ±ΓM ^d (GeV ²)	Partial decay mode		
										Mode	Fraction(%) [Upper limits (%) are 90% CL]	p or b P _{max} (MeV/c)
NONSTRANGE MESONS												
π [±]						1 ⁻ (0 ⁻) ⁺	139.57	0.0	0.019479	See Stable Particle Table		
π ⁰							134.96	7.95 eV ±0.55 eV	0.018215			
η						0 ⁺ (0 ⁻) ⁺	548.8	0.83 keV	0.301	Neutral	70.9	See Stable Particle Table
							±0.6	±0.12 keV	±0.000	Charged	29.1	
ρ(770)						1 ⁺ (1 ⁻) ⁻	769 [±]	154 [±]	0.591	ππ	≈ 100	358
							±3 [§]	±5 [§]	±0.118	π [±] γ	0.044 ± 0.005	372
										μ ⁺ μ ⁻	0.0067 ± 0.0012 ^d	370
										e ⁺ e ⁻	0.0043 ± 0.0005 ^d	384
										ηγ	seen [†]	189
M and Γ from neutral mode.			For upper limits, see footnote e									
ω(783)						0 ⁻ (1 ⁻) ⁻	782.6	9.9	0.612	π ⁺ π ⁻ π ⁰	89.9 ± 0.5	327
							±0.2	±0.3	±0.008	π ⁰ γ	8.7 ± 0.5	380
							S=1.1*			π ⁺ π ⁻	1.4 ± 0.2	366
										π ⁰ μ ⁺ μ ⁻	0.010 ± 0.002	349
										e ⁺ e ⁻	0.0072 ± 0.0007 S=1.3*	391
										ηγ	seen [†]	199
For upper limits, see footnote f												
η'(958)						0 ⁺ (0 ⁻) ⁺ [‡]	957.57	0.28	0.917	ηππ	65.3 ± 1.6	231
							±0.25	±0.10	±0.0003	ρ ⁰ γ	30.0 ± 1.6	170
										ωγ	2.8 ± 0.5	159
										γγ	1.9 ± 0.2	479
										μ ⁺ μ ⁻ γ	0.009 ± 0.002	467
For upper limits, see footnote g												
S'(975)						0 ⁺ (0 ⁺) ⁺	975 ^c	33 ^c	0.951	ππ	78 ± 3	467
							±4	±6	±0.032	K \bar{K}	22 ± 3	
See note on ππ and K \bar{K} S wave. [‡]			S=1.4*									
δ(980) [‡]						1 ⁻ (0 ⁺) ⁺	983 ^h	54 ^h	0.966	ηπ	seen	320
							±2	±7	±0.053	K \bar{K}	seen	
φ(1020)						0 ⁻ (1 ⁻) ⁻	1019.61	4.21	1.040	K ⁺ K ⁻	49.1 ± 1.0	S=1.3*
							±0.07	±0.13	±0.004	K ₁ K _S	34.6 ± 1.0	S=1.3*
										π [±] π ⁰ (incl. ρπ)	14.8 ± 0.7	S=1.2*
										ηγ	1.5 ± 0.2	362
										π ⁰ γ	0.14 ± 0.05	501
										e ⁺ e ⁻	0.031 ± 0.001	510
										μ ⁺ μ ⁻	0.025 ± 0.003	499
									π ⁺ π ⁻	0.02 ± 0.01	490	
For upper limits, see footnote i												
H(1190)						0 ⁻ (1 ⁺) ⁻	1190	320	1.416	ρπ	seen	327
							±60	±50	±0.381			
Seen in one experiment only.												
B(1235)						1 ⁺ (1 ⁺) ⁻	1233	137	1.52	ωπ	only mode seen	349
							±10 [§]	±10 [§]	±0.17	[D/S amplitude ratio = 0.29 ± 0.05]		
For upper limits, see footnote j												

Meson Table (cont'd)

J ^P	G	I	C	N	ε	ρ	δ	K*	I ^G (J ^P)C _N	Mass M (MeV)	Full Width Γ (MeV)	M ² ±ΓM ^d (GeV ²)	Partial decay mode		
													Mode	Fraction(%) [Upper limits (%) are 90% CL]	S or P _{max} ^b (MeV/c)
f(1270)	0 ⁺ (2 ⁺) ⁺	1273 ±5 [§]	179 ±20 [§]	1.62 ±0.23	ππ	83.1 ± 1.9	S=1.4*	621							
					2π ⁺ 2π ⁻	2.8 ± 0.4	S=1.2*	558							
					K \bar{K}	2.9 ± 0.2	S=1.2*	397							
					γγ	0.0016 ± 0.0003		637							
					π ⁺ π ⁻ 2π ⁰	seen		561							
For upper limits, see footnote k															
A ₁ (1270)	1 ⁻ (1 ⁺) ⁺	1275 [‡] ±30	315 [‡] ±45	1.53 ±0.40	ρπ	dominant		389							
					π(ππ) _{S-wave}	seen		599							
D(1285)	0 ⁺ (1 ⁺) ⁺	1283 ±5 [§]	26 ±5 [§]	1.65 ±0.03	K \bar{K} π	11 ± 3		302							
					ηππ	49 ± 6		482							
					†[δπ	36 ± 7]		236							
					4π (prob. ρππ) [‡]	40 ± 7		564							
ε(1300)	0 ⁺ (0 ⁺) ⁺	~1300	200-600		ππ	~90		635							
					K \bar{K}	~10		418							
					ηη			348							
See note on ππ and K \bar{K} S wave. [‡]															
π(1300)	1 ⁻ (0 ⁻) ⁺	1300 [§] ±100 [§]	200-600		ρπ	seen		407							
					π(ππ) _{S-wave}	seen		612							
Not a well-established resonance.															
A ₂ (1320)	1 ⁻ (2 ⁺) ⁺	1318 ±5 [§]	110 ±5 [§]	1.74 ±0.14	ρπ	70.1 ± 2.2		419							
					ηπ	14.5 ± 1.2		534							
					ωππ	10.6 ± 2.5		361							
					K \bar{K}	4.8 ± 0.5		434							
					η'π	< 2 (CL=97%)		286							
					πγ	0.45 ± 0.11		652							
					γγ	0.0007 ± 0.0004		659							
E(1420) [‡]	0 ⁺ (1 ⁺) ⁺	1418 ±10 [§]	52 ±10 [§]	2.01 ±0.07	K \bar{K} π (prob. K* \bar{K} +K \bar{K} *)	seen		423							
					ηππ	possibly seen		565							
					†[δπ	possibly seen]		348							
f'(1515)	0 ⁺ (2 ⁺) ⁺	1520 ±10 [§]	75 ±10 [§]	2.31 ±0.11	K \bar{K}	dominant		574							
					ππ	possibly seen		747							
ρ'(1600)	1 ⁺ (1 ⁻) ⁻	1600 [‡] ±20 [§]	300 [‡] ±100 [§]	2.56 ±0.48	4π (incl. ρπ ⁺ π ⁻ , A ₁ π)	large		738							
					ππ	< 30 [‡]		788							
					K* \bar{K} + \bar{K} *K	~15		388							
					ηππ	~13		675							
					$\bar{K}\bar{K}$	~1		630							
					e ⁺ e ⁻	seen		800							
ω(1670)	0 ⁻ (3 ⁻) ⁻	1688 ±5	166 ±15 [§] S=1.1*	2.78 ±0.28	3π	seen		806							
					†[ρπ	seen]		648							
					5π	seen		740							
					†[ωππ (prob. Bπ)	seen]		616							
A ₃ (1680) [‡]	1 ⁻ (2 ⁻) ⁺	1680 [§] ±30 [§]	250 [§] ±50 [§]	2.82 ±0.42	fπ	55 ± 5		337							
					ρπ	36 ± 6		656							
					π(ππ) _{S-wave}	9 ± 5		813							
For upper limits, see footnote ℓ															
φ(1680)	0 ⁻ (1 ⁻) ⁻	1684 ±15 [§]	126 ±22	2.84 ±0.21	K* \bar{K} + \bar{K} *K	dominant		541							
					ωππ	seen		624							
					K \bar{K}	seen		682							
g(1690) [‡]	1 ⁺ (3 ⁻) ⁻	1691 ±5 [§]	200 [§] ±20 [§]	2.86 ±0.34	2π	23.8 ± 1.3		834							
					4π (incl. ππρ, ρρ, A ₂ π, ωπ)	70.9 ± 1.9		787							
					K \bar{K} π (incl. K* \bar{K})	3.8 ± 1.2		625							
					K \bar{K}	1.5 ± 0.3	S=1.3*	684							

† J^P, M, and Γ from the 2π and K \bar{K} modes.

‡

§

Meson Table (cont'd)

J^P	$G \times \eta$	0	1	$1/2$	$I^G(J^P)C_\eta$	Mass M (MeV)	Full Width Γ (MeV)	M^2 $\pm \Gamma M^a$ (GeV ²)	Partial decay mode		
									Mode	Fraction(%) [Upper limits (%) are 90% CL]	p or p_{max} (MeV/c)
$h(2040)$	$+$	ϵ	ρ	π	K^+	$2040^{\pm 8}$ $\pm 20^{\pm 8}$	$150^{\pm 8}$ $\pm 50^{\pm 8}$	4.16 ± 0.31	$\pi\pi$ $K\bar{K}$	seen seen	1010 890
$\eta_c(2980)$	$+$	η	π	K	$0^+(4^+) +$	2981 ± 6	< 20	8.89	$\eta\pi^+\pi^-$ $2(\pi^+\pi^-)$ $K^+K^-\pi^+\pi^-$ $p\bar{p}$	seen seen seen seen	1426 1458 1343 1158
$J/\psi(3100)$	$+$	$\omega/\phi/\delta$	π	K	$0^-(1^-) -$	3096.9 ± 0.1	0.063 ± 0.009	9.591 ± 0.000	e^+e^- $\mu^+\mu^-$ hadrons + radiative	7.4 ± 1.2 7.4 ± 1.2 85 ± 2	1548 1545
Decay modes into stable hadrons						Decay modes into hadronic resonances					
$\dagger[2(\pi^+\pi^-)\pi^0$						3.7 ± 0.5		1496	$\dagger[\rho\pi$	1.22 ± 0.12	1449
$3(\pi^+\pi^-)\pi^0$						2.9 ± 0.7		1433	$\omega 2\pi^+ 2\pi^-$	0.85 ± 0.34	1392
$\pi^+\pi^-\pi^0 K^+K^-$						1.2 ± 0.3		1368	ρA_2	0.84 ± 0.45	1126
$4(\pi^+\pi^-)\pi^0$						0.9 ± 0.3		1345	$\omega\pi\pi$	0.68 ± 0.19	1435
$\pi^+\pi^-K^+K^-$						0.72 ± 0.23		1407	$K^*(892)\bar{K}^*(1430) + c.c.$	0.67 ± 0.26	1007
$p\bar{p}\pi^+\pi^-$						0.53 ± 0.06		1107	$K^{\pm}K^{\mp}(892)$	0.34 ± 0.05	1373
$2(\pi^+\pi^-)$						0.4 ± 0.1		1517	$B^{\pm}(1235)\pi^{\pm}$	0.29 ± 0.07	1299
$3(\pi^+\pi^-)$						0.4 ± 0.2		1466	$K^0\bar{K}^*(892) + c.c.$	0.27 ± 0.06	1373
$n\bar{n}\pi^+\pi^-$						0.38 ± 0.36		1106	ωf	0.23 ± 0.08	$S=1.2^*$ 1144
$\Xi\bar{\Xi}$						0.32 ± 0.08		818	$\phi\pi^+\pi^-$	0.21 ± 0.09	1365
$2(\pi^+\pi^-)K^+K^-$						0.31 ± 0.13		1320	$\eta'p\bar{p}$	0.18 ± 0.06	596
$K_S^0 K_L^{\pm}\pi^{\mp}$						0.26 ± 0.07		1440	$\phi K K$	0.18 ± 0.08	1176
$\Sigma^+\Sigma^-$						0.24 ± 0.26		988	$\omega p\bar{p}$	0.16 ± 0.03	768
$p\bar{p}\eta$						0.23 ± 0.04		948	$\omega K\bar{K}$	0.16 ± 0.10	1265
$p\bar{p}$						0.22 ± 0.02		1232	$\phi\eta$	0.10 ± 0.06	1320
$p\bar{n}\pi^-$ or $\bar{p}n\pi^+$						0.21 ± 0.02		1174	$\phi f(1515)$	0.08 ± 0.05	874
$n\bar{n}$						0.18 ± 0.09		1231	$\pi^{\pm} A_2^{\mp}$	< 0.43	1263
$p\bar{p}\pi^+\pi^-$						$0.16 \pm 0.06^{**}$		1033	$K^*(430)\bar{K}^*(1430)$	< 0.29	584
$\Sigma^0\Sigma^0$						0.13 ± 0.04		988	$K^0\bar{K}^*(1430) + c.c.$	< 0.2	1184
$\Lambda\bar{\Lambda}$						0.11 ± 0.02		1074	$K^{\pm}K^{\mp}(1430)$	< 0.2	1154
$p\bar{p}\pi^0$						0.11 ± 0.01		1176	$\phi 2\pi^+ 2\pi^-$	< 0.15	1318
$2(K^+K^-)$						0.07 ± 0.03		1131	$\phi\eta'$	< 0.13	1192
K^+K^-						0.022 ± 0.008		1468	$K^*(892)\bar{K}^*(892)$	< 0.05	1266
$\pi^+\pi^-$						0.011 ± 0.005		1542	ϕf	< 0.037	1037
$\Delta\bar{\Sigma}$						< 0.015		1032	ωf	< 0.016	1006
$K_S^0 K_L^0$						< 0.009		1466	Radiative decay modes		
$\chi(3415)$	$0^+(0^+) +$					3415.0 ± 1.0		11.662	$\dagger[\gamma\pi(1440)$	$0.55 \pm 0.22^{\dagger}$	1224
									$\gamma\eta'$	0.36 ± 0.05	1400
									γf	0.15 ± 0.04	1287
									$\gamma\eta$	0.086 ± 0.009	1500
									$\gamma\pi^0$	0.007 ± 0.005	1546
									$\gamma D(1285)$	< 0.6	1283
									2γ	< 0.05	1548
									$\gamma f'(1515)$	< 0.03	1175
									$\gamma p\bar{p}$	< 0.01	1232
									3γ	< 0.006	1548
									$2(\pi^+\pi^-)$ (incl. $\pi\pi\rho$)	4.3 ± 0.9	1679
									$\pi^+\pi^-K^+K^-$ (incl. $\pi K\bar{K}^*$)	3.4 ± 0.9	1580
									$3(\pi^+\pi^-)$	1.7 ± 0.6	1633
									$\pi^+\pi^-$	0.9 ± 0.2	1702
									$\gamma J/\psi(3100)$	0.8 ± 0.2	303
									K^+K^-	0.8 ± 0.2	1635
									$p\bar{p}\pi^+\pi^-$	0.6 ± 0.2	1320
For upper limits, see footnote <i>o</i>											

Meson Table (cont'd)

J^P	G	I	Y	C	K^*	$G(J^P)_C$	Mass M (MeV)	Full Width Γ (MeV)	M^2 $\pm \Gamma M^2$ (GeV ²)	Partial decay mode		
										Mode	Fraction(%) [Upper limits (%) are 90% CL]	p or p_{max} (MeV/c)
A	+	0	1	1/2	-	estab.	10569 ± 10	14 ± 5	111.704 ± 0.15	e^+e^-	0.0019 \pm 0.0008	5285
$m_{T(10570)} - m_{T(9460)} = 1113 \pm 4$												
STRANGE MESONS												
							493.67		0.244	See Stable Particle Table		
							497.67		0.248			
							891.8 ± 0.4	50.8 ± 0.9	0.795 ± 0.045	$K\pi$ $K\gamma$ $K\pi\pi$	≈ 100 0.15 \pm 0.07 < 0.07 (CL=95%)	288 309 216
M and Γ from charged mode; $m^0 - m^\pm = 6.7 \pm 1.2$ MeV.												
							1270 ^S $\pm 10^S$	90 ^S $\pm 20^S$	1.61 ± 0.11	$K\rho$ $\kappa\pi$ $K^*\pi$ $K\omega$ $K\epsilon$	42 \pm 6 28 \pm 4 16 \pm 5 11 \pm 2 3 \pm 2	45 299
							~ 1350	~ 250	1.82 ± 0.34	$K\pi$	seen	574
See note on $K\pi$ S wave. †												
							1414 ± 13	180 ± 10	2.00 ± 0.25	$K^*\pi$ $K\rho$ $K\epsilon$ $K\omega$	94 \pm 6 3 \pm 3 2 \pm 2 1 \pm 1	410 308 294
							1434 ^S $\pm 5^S$	100 ^S $\pm 10^S$	2.06 ± 0.14	$K\pi$ $K^*\pi$ $K^*\pi\pi$ $K\rho$ $K\omega$ $K\eta$	44.8 \pm 2.3 24.6 \pm 2.0 13.0 \pm 2.6 8.8 \pm 1.1 4.2 \pm 1.5 5 \pm 5 ^S	S=2.7* S=1.1* S=1.1* S=1.3* 623 424 374 334 320 492
							$\sim 1770^S$	$\sim 200^S$	3.13 ± 0.35	$K^*(1430)\pi$ $K^*(892)\pi$ $K\Gamma$	dominant seen seen	278 652
See note on $L(1770)$. †												
							1775 ^S $\pm 10^S$	140 ^S $\pm 20^S$	3.15 ± 0.25	$K\pi\pi$ $\dagger\{K\rho$ $\dagger\{K^*\pi$ $K\pi$	large large] large] 17 \pm 5 ^S	793 616 654 812
See note on $K^*(1780)$. †												
CHARMED, NONSTRANGE MESONS												
							1869.4		3.495	See Stable Particle Table		
							1864.7		3.477			
							2010.1 ± 0.7	< 2.0	4.041	$D^0\pi^+$ $D^+\pi^0$ $D^+\gamma$	64 \pm 11 28 \pm 9 8 \pm 7	39 38 136
$m_{D^{*+}} - m_{D^0} = 145.4 \pm 0.2$ MeV												
							2007.2 ± 2.1	< 5	4.029	$D^0\pi^0$ $D^0\gamma$	55 \pm 15 45 \pm 15	44 137
CHARMED, STRANGE MESON												
							2021		4.084	See Stable Particle Table		

Meson Table (cont'd)

Contents of Meson Data Card Listings

Non-strange (S = 0; C, B = 0)						Strange (S = 1; C, B = 0)	
entry	$I^G(J^P)C_n$	entry	$I^G(J^P)C_n$	entry	$I^G(J^P)C_n$	entry	I (J ^P)
π	$1^-(0^-)+$	ρ' (1515)	$0^+(2^+)+$	$\rightarrow \delta$ (2450)	$1^-(6^+)+$	K	1/2(0 ⁻)
η	$0^+(0^-)+$	ρ' (1600)	$1^+(1^-)-$	$\rightarrow e^+e^-$ (1100-2200)		K^* (892)	1/2(1 ⁻)
ρ (770)	$1^+(1^-)-$	$\rightarrow \theta$ (1640)	$0^+(2^+)+$	$\rightarrow \bar{N}N$ (1400-3600)		Q_1 (1280)	1/2(1 ⁺)
ω (783)	$0^-(1^-)-$	ω (1670)	$0^-(3^-)-$	$\rightarrow X$ (1900-3600)		κ (1350)	1/2(0 ⁺)
η' (958)	$0^+(0^-)+$	A_3 (1680)	$1^-(2^-)+$	η_c (2980)	+	Q_2 (1400)	1/2(1 ⁺)
S^* (975)	$0^+(0^+)+$	ϕ (1680)	$0^-(1^-)-$	J/ψ (3100)	$0^-(1^-)-$	$\rightarrow K'$ (1400)	1/2(0 ⁻)
δ (980)	$1^-(0^+)+$	g (1690)	$1^+(3^-)-$	χ (3415)	$0^+(0^+)+$	K^* (1430)	1/2(2 ⁺)
ϕ (1020)	$0^-(1^-)-$	$\rightarrow \phi$ (1850)	0	P_c or χ (3510)	$0^+(1^+)+$	$\rightarrow L$ (1580)	1/2(2 ⁻)
H (1190)	$0^-(1^+)-$	$\rightarrow X$ (1850)	(2 ⁺)	χ (3555)	$0^+(2^+)+$	$\rightarrow K^*$ (1650)	1/2(1 ⁻)
B (1235)	$1^+(1^+)-$	$\rightarrow S$ (1935)		$\rightarrow \eta_c'$ (3590)	+	L (1770)	1/2(2 ⁻)
$\rightarrow \rho'$ (1250)	$1^+(1^-)-$	$\rightarrow \delta$ (2030)	$1^-(4^+)+$	ψ (3685)	$0^-(1^-)-$	K^* (1780)	1/2(3 ⁻)
f (1270)	$0^+(2^+)+$	h (2040)	$0^+(4^+)+$	ψ (3770)	(1 ⁻)	$\rightarrow K^*$ (2060)	1/2(4 ⁺)
A_1 (1270)	$1^-(1^+)+$	$\rightarrow \pi$ (2050)	$1^-(3^+)+$	ψ (4030)	(1 ⁻)	$\rightarrow K^*$ (2200)	
$\rightarrow \eta$ (1275)	$0^+(0^-)+$	$\rightarrow \pi$ (2100)	$1^-(2^-)+$	ψ (4160)	(1 ⁻)	Charmed (C = 1)	
D (1285)	$0^+(1^+)+$	$\rightarrow \rho$ (2150)	$1^+(1^-)-$	ψ (4415)	(1 ⁻)	D (1870)	1/2(0 ⁻)
ϵ (1300)	$0^+(0^+)+$	$\rightarrow \epsilon$ (2150)	$0^+(2^+)+$	T (9460)	(1 ⁻)	D^* (2010)	1/2(1 ⁻)
π (1300)	$1^-(0^-)+$	$\rightarrow \rho$ (2250)	$1^+(3^-)-$	T (10020)	(1 ⁻)	F (2020)	0 (0 ⁻)
A_2 (1320)	$1^-(2^+)+$	$\rightarrow \epsilon$ (2300)	$0^+(4^+)+$	T (10350)	(1 ⁻)	$\rightarrow F^*$ (2140)	
E (1420)	$0^+(1^+)+$	$\rightarrow \rho$ (2350)	$1^+(5^-)-$	T (10570)	(1 ⁻)	Bottom (Beauty) (B = 1)	
						$\rightarrow B$ (5200)	
						\rightarrow Exotics	

\rightarrow Indicates an entry in the Meson Data Card Listings not entered in the Meson Table. We do not regard these as established resonances. All the entries in the Listings can be found in the Table of Contents of the Meson Data Card Listings immediately preceding these footnotes.

‡ See Meson Data Card Listings.

* Quoted error includes scale factor $S = \sqrt{\chi^2/(N-1)}$. See footnote to Stable Particle Table.

† Square brackets indicate a subreaction of the previous (unbracketed) decay mode(s).

§ This is only an educated guess; the error given is larger than the error on the average of the published values. (See the Meson Data Card Listings for the latter.)

a. ΓM is approximately the half-width of the resonance when plotted against M^2 .

b. For decay modes into ≥ 3 particles, p_{\max} is the maximum momentum that any of the particles in the final state can have. The momenta have been calculated by using the averaged central mass values, without taking into account the widths of the resonances.

c. From pole position $(M - i\Gamma/2)$.

d. The e^+e^- branching fraction is from $e^+e^- \rightarrow \pi^+\pi^-$ experiments only. The $\omega\rho$ interference is then due to $\omega\rho$ mixing only, and is expected to be small. See note in the Meson Data Card Listings. The $\mu^+\mu^-$ branching fraction is compiled from 3 experiments, each possibly with substantial $\omega\rho$ interference. The error reflects this uncertainty; see notes in the Meson Data Card Listings. If μ universality holds, $\Gamma(\rho^0 \rightarrow \mu^+\mu^-) = \Gamma(\rho^0 \rightarrow e^+e^-) \times 0.99785$.

Meson Table (cont'd)

- e. Empirical limits on fractions for other decay modes of $\rho(770)$ are $\pi^{\pm}\eta < 0.8\%$ (CL=84%), $\pi^{\pm}\pi^{+}\pi^{-}\pi^{-} < 0.15\%$, $\pi^{\pm}\pi^{+}\pi^{-}\pi^{0} < 0.2\%$ (CL=84%).
- f. Empirical limits on fractions for other decay modes of $\omega(783)$ are $\pi^{+}\pi^{-}\gamma < 5\%$, $\pi^{0}\pi^{0}\gamma < 1\%$, $\eta + \text{neutral}(s) < 1.5\%$, $\mu^{+}\mu^{-} < 0.02\%$.
- g. Empirical limits on fractions for other decay modes of $\eta(958)$ are $\pi^{+}\pi^{-} < 2\%$ (CL=84%), $\pi^{+}\pi^{-}\pi^{0} < 5\%$ (CL=84%), $\pi^{+}\pi^{-}\pi^{-}\pi^{-} < 1\%$ (CL=95%), $\pi^{+}\pi^{+}\pi^{-}\pi^{-}\pi^{0} < 1\%$ (CL=84%), $6\pi < 1\%$, $\pi^{+}\pi^{-}e^{+}e^{-} < 0.6\%$, $\pi^{0}e^{+}e^{-} < 1.3\%$ (CL=84%), $\eta e^{+}e^{-} < 1.1\%$, $\pi^{0}\rho^{0} < 4\%$, $\eta\mu^{+}\mu^{-} < 1.5 \times 10^{-5}$, $\pi^{0}\mu^{+}\mu^{-} < 6 \times 10^{-5}$.
- h. The mass and width are from the $\eta\pi$ mode only. If the $K\bar{K}$ channel is strongly coupled, the width may be larger.
- i. Empirical limits on fractions for other decay modes of $\phi(1020)$ are $\pi^{+}\pi^{-}\gamma < 0.7\%$, $\omega\gamma < 5\%$ (CL=84%), $\rho\gamma < 2\%$ (CL=84%), $2\pi^{+}2\pi^{-}\pi^{0} < 1\%$ (CL=95%), $2\pi^{+}2\pi^{-} < 0.1\%$.
- j. Empirical limits on fractions for other decay modes of $B(1235)$ are $\pi\pi < 15\%$, $K\bar{K} < 2\%$ (CL=84%), $4\pi < 50\%$ (CL=84%), $\phi\pi < 1.5\%$ (CL=84%), $\eta\pi < 25\%$, $(\bar{K}K)^{\pm}\pi^{0} < 8\%$, $K_S K_S \pi^{\pm} < 2\%$, $K_S K_L \pi^{\pm} < 6\%$.
- k. Empirical limits (CL=95%) on fractions for other decay modes of $f(1270)$ are $\eta\pi\pi < 1\%$, $K^0 K^{-}\pi^{+} + c.c. < 0.5\%$, $\eta\eta < 2\%$.
- l. Empirical limits on fractions for other decay modes of $A_3(1680)$ are $\eta\pi < 10\%$, $5\pi < 10\%$.
- m. Includes $p\bar{p}\pi^{+}\pi^{-}\gamma$ and excludes $p\bar{p}\eta$, $p\bar{p}\omega$, $p\bar{p}\eta'$.
- n. The $\epsilon(1440)$ evidence is listed under $E(1420)$; see $E(1420)$ mini-review.
- o. Empirical limits on fractions for other decay modes of $\chi(3415)$ are $2\gamma < 0.17\%$, $p\bar{p} < 0.011\%$.
- p. Empirical limits on fractions for other decay modes of $\chi(3510)$ are $(\pi^{+}\pi^{-}$ and $K^{+}K^{-}) < 0.2\%$, $\gamma\gamma < 0.16\%$, $p\bar{p} < 0.13\%$.
- q. Empirical limits on fractions for other decay modes of $\chi(3555)$ are $2\gamma < 0.06\%$, $p\bar{p} < 0.10\%$, $J/\psi\pi^{+}\pi^{-}\pi^{0} < 1.5\%$.

Established Nonets, and octet-singlet mixing angles θ obtained from the Gell-Mann-Okubo mass formula [Appendix II, Eq. (3)]. Of the two isosinglets, the "mainly octet" one is written first, followed by the "mainly singlet" one; a semicolon separates the two. The angle $\delta = \theta - 35.3^{\circ}$ measures the deviation from ideal mixing.

$(J^P)C_n$	Nonet members	$\theta_{\text{lin.}}$	$\theta_{\text{quadr.}}$	$\delta_{\text{lin.}}$	$\delta_{\text{quadr.}}$
$(0^+) +$	$\pi, K, \eta; \eta'$	$-24.4 \pm 0.1^{\circ}$	$-11.1 \pm 0.2^{\circ}$	$-59.7 \pm 0.1^{\circ}$	$-46.4 \pm 0.2^{\circ}$
$(1^-) -$	$\rho, K^*, \phi; \omega$	$35.9 \pm 0.5^{\circ}$	$38.6 \pm 0.4^{\circ}$	$0.6 \pm 0.5^{\circ}$	$3.3 \pm 0.4^{\circ}$
$(2^+) +$	$A_2, K^*(1430), f'; \bar{f}$	$26 \pm 3^{\circ}$	$28 \pm 3^{\circ}$	$-9 \pm 3^{\circ}$	$-7 \pm 3^{\circ}$
$(1^+) +^{\dagger}$	$A_1, Q_A, E; D$	$52 \pm 13^{\circ}$	$51 \pm 12^{\circ}$	$16 \pm 13^{\circ}$	$15 \pm 12^{\circ}$

$^{\dagger} m(Q_A)$ is assumed to be the average of $m(Q_1)$ and $m(Q_2)$.

More generally, because of unitarity, the mixing angles are energy dependent and complex above the first threshold (see Appendix II C), which is important especially for the scalar and the axial mesons. Note also that the two axial strange mesons (Q_1 and Q_2) are mixtures of the exact SU(3) states: $Q_1 = \cos\phi Q_A + \sin\phi Q_B$, $Q_2 = -\sin\phi Q_A + \cos\phi Q_B$. Below we give the mixing angles δ and ϕ obtained in a unitary mixing scheme using both masses and widths as input data:

$(J^P)C_n$	Nonet members	Mixing angles
$(1^+) +$	$A_1, Q_A, E; D$	$\delta_{DE}(1283) = 14^{\circ} + i1^{\circ}$
		$\delta_{DE}(1418) = 25^{\circ} + i8^{\circ}$
$(1^+) -$	$B, Q_B, H;^{\dagger} H$	$\delta_{HH}(1190) = -6^{\circ} + i4^{\circ}$
		$\delta_{HH}(1400) = -15^{\circ} + i10^{\circ}$
		$\phi_{Q_1 Q_2}(1270) = 50^{\circ} + i3^{\circ}$
		$\phi_{Q_1 Q_2}(1414) = 61^{\circ} - i3^{\circ}$
$(0^+) +$	$\delta, \kappa, S;^{\dagger} \epsilon$	$\delta_{S\kappa}(975) = +4^{\circ} + i29^{\circ}$
		$\delta_{S\kappa}(1300) = -33^{\circ} + i7^{\circ}$

† as yet, not seen experimentally

Baryon Table

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The following short list gives the name, the nominal mass, the quantum numbers (where known), and the status of each of the Baryon States in the Data Card Listings. States with 3- or 4-star status are included in the Baryon Table below; the others are omitted because the evidence for the existence of the effect and/or for its interpretation as a resonance is open to question.

N(939) P ₁₁ ****	Δ(1232) P ₃₃ ****	Λ(1115) P ₀₁ ****	Σ(1193) P ₁₁ ****	Ξ(1317) P ₁₁ ****
N(1440) P ₁₁ ****	Δ(1550) P ₃₁ **	Λ(1405) S ₀₁ ****	Σ(1385) P ₁₃ ****	Ξ(1530) P ₁₃ ****
N(1520) D ₁₃ ****	Δ(1600) P ₃₃ ****	Λ(1520) D ₀₃ ****	Σ(1480) ****	Ξ(1630) **
N(1535) S ₁₁ ****	Δ(1620) S ₃₁ ****	Λ(1600) P ₀₁ **	Σ(1560) **	Ξ(1680) S ₁₁ **
N(1540) P ₁₃ *	Δ(1700) D ₃₃ ****	Λ(1670) S ₀₁ ****	Σ(1580) D ₁₃ **	Ξ(1820) 13 **
N(1650) S ₁₁ ****	Δ(1900) S ₃₁ **	Λ(1690) D ₀₃ ****	Σ(1620) S ₁₁ **	Ξ(1940) **
N(1675) D ₁₅ ****	Δ(1905) F ₃₅ ****	Λ(1800) S ₀₁ **	Σ(1660) P ₁₁ ****	Ξ(2030) 1 **
N(1680) F ₁₅ ****	Δ(1910) P ₃₁ ****	Λ(1800) P ₀₁ **	Σ(1670) D ₁₃ ****	Ξ(2120) *
N(1700) D ₁₃ ****	Δ(1920) P ₃₃ **	Λ(1800) G ₀₉ Dead	Σ(1670) **	Ξ(2250) *
N(1710) P ₁₁ ****	Δ(1930) D ₃₅ ****	Λ(1800) *	Σ(1690) **	Ξ(2370) 1 **
N(1720) P ₁₃ ****	Δ(1940) D ₃₃ *	Λ(1820) F ₀₅ ****	Σ(1750) S ₁₁ **	Ξ(2500) **
N(1990) F ₁₇ **	Δ(1950) F ₃₇ ****	Λ(1830) D ₀₅ ****	Σ(1770) P ₁₁ Dead	
N(2000) F ₁₅ **	Δ(2150) S ₃₁ *	Λ(1890) P ₀₃ ****	Σ(1775) D ₁₅ ****	Ω(1672) P ₀₃ ****
N(2080) D ₁₃ **	Δ(2160) *	Λ(2000) *	Σ(1840) P ₁₃ *	
N(2100) S ₁₁ *	Δ(2200) G ₃₇ **	Λ(2020) F ₀₇ *	Σ(1880) P ₁₁ **	Λ _c (2282) ****
N(2100) P ₁₁ *	Δ(2300) H ₃₉ **	Λ(2100) G ₀₇ ****	Σ(1915) F ₁₅ ****	
N(2190) G ₁₇ ****	Δ(2350) D ₃₅ *	Λ(2110) F ₀₅ **	Σ(1940) D ₁₃ **	Σ _c (2450) **
N(2200) D ₁₅ ****	Δ(2400) F ₃₇ *	Λ(2325) D ₀₃ *	Σ(2000) S ₁₁ *	
N(2220) H ₁₉ ****	Δ(2400) G ₃₉ *	Λ(2350) ****	Σ(2030) F ₁₇ ****	Λ _b (5500) *
N(2250) G ₁₉ ****	Δ(2420) H ₃₁₁ **	Λ(2585) **	Σ(2070) F ₁₅ *	
N(2600) I ₁₁₁ **	Δ(2500) *		Σ(2080) P ₁₃ **	Dibaryons
N(2700) K ₁₁₃ *	Δ(2750) 1313 *		Σ(2100) G ₁₇ *	NN(2170) 1F2 **
N(2800) G ₁₉ *	Δ(2850) **		Σ(2250) ****	NN(2250) 3F3 **
N(3000) *	Δ(2950) K ₃₁₅ *		Σ(2455) **	NN(?) *
N(3030) **	Δ(3230) **		Σ(2620) **	ΛN(2130) 3S1 **
N(3245) *			Σ(3000) **	ΞN(?) *
N(3690) *	Z ₀ (1780) P ₀₁ *		Σ(3170) *	
N(3755) *	Z ₀ (1865) D ₀₃ *			
	Z ₁ (1900) P ₁₃ *			
	Z ₁ (2150) *			
	Z ₁ (2500) *			

- **** Good, clear, and unmistakable.
 *** Good, but in need of clarification or not absolutely certain.
 ** Needs confirmation.
 * Weak.

Particle ^a	I(J ^P)L _{21-2J} ^b	P _{beam} ^c (GeV/c) σ = 4πχ ² (mb)	Mass ^d M (MeV)	Full ^e width Γ (MeV)	M ^{2 f} ±ΓM (GeV ²)	Partial decay mode		
						Mode ^g	Fraction ^h (%)	p ⁱ (MeV/c)
S=0 I=1/2 NUCLEON RESONANCES (N)								
p	1/2(1/2 ⁺)		938.3		0.880	See Stable Particle Table		
n			939.6		0.883			
N(1440)	1/2(1/2 ⁺)P ₁₁	p = 0.61 σ = 31.0	1400 to 1480	120 to 350 (200)	2.07 ±0.29	Nπ Nη Nππ	50-70 8-18 ~30	397 † 342
						Δπ Nρ Nε	12-28 ~7 ~5	143 † †
N(1520)	1/2(3/2 ⁻)D ₁₃	p = 0.74 σ = 23.5	1510 to 1530	100 to 140 (125)	2.31 ±0.19	Nπ Nη Nππ	50-60 <1 35-50	456 149 410
						Δπ Nρ Nε	15-25 15-25 <5	228 † †
N(1535)	1/2(1/2 ⁻)S ₁₁	p = 0.76 σ = 22.4	1520 to 1560	100 to 250 (150)	2.36 ±0.23	Nπ Nη Nππ	35-50 40-65 ~5	467 182 422
						Δπ Nρ Nε	~1 ~3 ~2	242 † †

Baryon Table (cont'd)

Particle ^a	$I(J^P)_L^b$ _{21, 2J}	P _{beam} ^c (GeV/c) $\sigma = 4\pi\chi^2$ (mb)	Mass ^d M (MeV)	Full ^e width Γ (MeV)	M^2 ^f $\pm \Gamma M$ (GeV ²)	Partial decay mode		
						Mode ^g	Fraction ^h (%)	p^i (MeV/c)
N(1650)	$1/2(1/2^-)S''_{11}$	$p = 0.96$ $\sigma = 16.4$	1620 to 1680	100 to 200 (150)	2.72 ± 0.25	N π N η ΔK ΣK N $\pi\pi$ [$\Delta\pi$ N ρ N ϵ]	55-65 ~ 1 5-10 3-10 ~ 30 4-15 ~ 20 < 5	547 346 161 † 511 344 † †
N(1675)	$1/2(5/2^-)D''_{15}$	$p = 1.01$ $\sigma = 15.4$	1660 to 1690	120 to 180 (155)	2.81 ± 0.26	N π N η ΔK N $\pi\pi$ [$\Delta\pi$ N ρ]	30-40 < 2 < 1 55-70 50-65 ~ 5	563 374 209 529 364 †
N(1680)	$1/2(5/2^+)F''_{15}$	$p = 1.01$ $\sigma = 15.2$	1670 to 1690	110 to 140 (125)	2.82 ± 0.21	N π N η N $\pi\pi$ [$\Delta\pi$ N ρ N ϵ]	55-65 < 1 ~ 40 ~ 12 ~ 10 ~ 20	567 379 532 369 † †
N(1700)	$1/2(3/2^-)D''_{13}$	$p = 1.05$ $\sigma = 14.5$	1670 to 1730	70 to 120 (100)	2.89 ± 0.17	N π N η ΔK N $\pi\pi$ [$\Delta\pi$ N ρ N ϵ]	8-12 ~ 4 ~ 1 ~ 85 15-40 ~ 5 < 40	580 400 250 547 385 † †
N(1710)	$1/2(1/2^+)P''_{11}$	$p = 1.07$ $\sigma = 14.2$	1680 to 1740	90 to 130 (110)	2.92 ± 0.19	N π N η ΔK ΣK N $\pi\pi$ [$\Delta\pi$ N ρ N ϵ]	10-20 5-35 5-15 2-10 > 50 10-25 25-65 15-40	587 410 264 138 554 393 48 †
N(1720)	$1/2(3/2^+)P''_{13}$	$p = 1.09$ $\sigma = 13.9$	1690 to 1800	125 to 250 (200)	2.96 ± 0.34	N π N η ΔK ΣK N $\pi\pi$ [$\Delta\pi$ N ρ N ϵ]	10-20 3-6 2-12 2-5 ~ 70 ~ 20 45-70 ~ 20	594 420 278 162 561 401 104 †
N(1990)	$1/2(7/2^+)F_{17}$	$p = 1.62$ $\sigma = 8.34$	1950 to 2050	120 to 400 (350)	3.96 ± 0.70	N π N η ΔK ΣK	~ 5 ~ 3 seen seen	766 648 554 497
N(2080)	$1/2(3/2^-)D''_{13}$	$p = 1.82$ $\sigma = 7.26$	2030 to 2100	115 to 300 (275)	4.33 ± 0.57	N π ΔK ΣK	~ 10 seen seen	821 627 576
N(2190)	$1/2(7/2^-)G_{17}$	$p = 2.07$ $\sigma = 6.21$	2120 to 2230	200 to 500 (350)	4.80 ± 0.77	N π N η ΔK	~ 14 ~ 2 < 1	888 790 712
N(2200)	$1/2(5/2^-)D''_{15}$	$p = 2.10$ $\sigma = 6.12$	1900 to 2230	150 to 400 (300)	4.84 ± 0.66	N π N η ΔK	~ ? seen seen	894 797 718
N(2220)	$1/2(9/2^-)H_{19}$	$p = 2.14$ $\sigma = 5.97$	2150 to 2300	300 to 500 (400)	4.93 ± 0.89	N π N η	~ 18 ~ 1	905 811

Baryon Table (cont'd)

Particle ^a	$I(J^P)L_{21-2J}$	P_{beam}^c (GeV/c) $\sigma = 4\pi\lambda^2$ (mb)	Mass ^d M (MeV)	Full ^e width Γ (MeV)	M^{2J} $\pm \Gamma M$ (GeV ²)	Partial decay mode		
						Mode ^f	Fraction ^h (%)	p^j (MeV/c)
N(2250)	$1/2(9/2^-)G_{19}$	$p = 2.21$ $\sigma = 5.74$	2130 to 2270	200 to 500 (300)	5.06 ± 0.68	N π	~ 10	923
						N η	~ 2	831
N(2600)	$1/2(11/2^-)I_{111}$	$p = 3.12$ $\sigma = 3.86$	2580 to 2700	>300 (400)	6.76 ± 1.04	N π	~ 1	1126
N(3030)	$1/2(?)$	$p = 4.41$ $\sigma = 2.62$	~ 3030	~ 400 (400)	9.18 ± 1.21	N π	$(J+1/2)_x$	1366
							$< 0.1^j$	
S=0 I=3/2 DELTA RESONANCES (Δ)								
$\Delta(1232)$	$3/2(3/2^+)P_{33}$	$p = 0.30$ $\sigma = 95.0$	1230 to 1234	110 to 120 (115)	1.52 ± 0.14	N π	99.4	227
						N γ	0.6	259
$\Delta(++)$ pole position: ^k $M-i\Gamma/2 = (1210.6 \pm 0.5) - i(49.7 \pm 0.3)$ ^l $\Delta(0)$ pole position: ^k $M-i\Gamma/2 = (1210.3 \pm 1.0) - i(53.0 \pm 1.0)$ ^l								
$\Delta(1600)$	$3/2(3/2^+)P_{33}$	$p = 0.87$ $\sigma = 18.7$	1500 to 1900	150 to 350 (250)	2.56 ± 0.40	N π	15-25	512
						N $\pi\pi$	~ 80	473
						$[\Delta\pi$ N ρ]	$20-65^p$ < 10	301 †
$\Delta(1620)$	$3/2(1/2^-)S_{31}$	$p = 0.91$ $\sigma = 17.7$	1600 to 1650	120 to 160 (140)	2.62 ± 0.23	N π	25-35	526
						N $\pi\pi$	~ 70	488
						$[\Delta\pi$ N ρ]	$35-50^p$ < 40	318 †
$\Delta(1700)$	$3/2(3/2^-)D_{33}$	$p = 1.05$ $\sigma = 14.5$	1630 to 1740	190 to 300 (250)	2.89 ± 0.43	N π	10-20	580
						N $\pi\pi$	~ 85	547
						$[\Delta\pi$ N ρ]	$< 50^p$ ~ 40	385 †
$\Delta(1900)$	$3/2(1/2^-)S_{31}$	$p = 1.44$ $\sigma = 9.71$	1850 to 2000	130 to 300 (150)	3.61 ± 0.29	N π	6-12	710
						ΣK	~ 10	410
$\Delta(1905)$	$3/2(5/2^+)F_{35}$	$p = 1.45$ $\sigma = 9.63$	1890 to 1920	250 to 400 (300)	3.63 ± 0.57	N π	8-15	713
						ΣK	< 3	415
						N $\pi\pi$	~ 80	687
						$[\Delta\pi$ N ρ]	$10-30^p$ ~ 60	542 421
$\Delta(1910)$	$3/2(1/2^+)P_{31}$	$p = 1.46$ $\sigma = 9.54$	1850 to 1950	200 to 330 (220)	3.65 ± 0.42	N π	20-25	716
						ΣK	2-20	421
						N $\pi\pi$	> 40	691
						$[\Delta\pi$ N ρ]	small ^p < 40	545 426
$\Delta(1920)$	$3/2(3/2^+)P_{33}$	$p = 1.48$ $\sigma = 9.39$	1860 to 2160	190 to 300 (250)	3.69 ± 0.48	N π	14-20	722
						ΣK	~ 1	431
$\Delta(1930)$	$3/2(5/2^-)D_{35}$	$p = 1.50$ $\sigma = 9.21$	1890 to 1960	150 to 350 (250)	3.72 ± 0.48	N π	4-14	729
						ΣK	< 10	441
$\Delta(1950)$	$3/2(7/2^+)F_{37}$	$p = 1.54$ $\sigma = 8.91$	1910 to 1960	200 to 340 (240)	3.80 ± 0.47	N π	35-45	741
						ΣK	< 1	460
						N $\pi\pi$	~ 50	716
						$[\Delta\pi$ N ρ]	$\sim 40^p$ ~ 20	574 469
$\Delta(2420)$	$3/2(11/2^+)H_{311}$	$p = 2.64$ $\sigma = 4.68$	2380 to 2450	300 to 500 (300)	5.86 ± 0.73	N π	5-15	1023
$\Delta(2850)$	$3/2(?)^+$	$p = 3.10$ $\sigma = 3.05$	2800 to 2900	~ 400 (400)	8.12 ± 1.14	N π	$(J+1/2)_x$	1266
							$\sim 0.25^j$	
$\Delta(3230)$	$3/2(?)$	$p = 5.08$ $\sigma = 2.25$	3200 to 3350	~ 440 (440)	10.43 ± 1.42	N π	$(J+1/2)_x$	1475
							$\sim 0.05^j$	

Baryon Table (cont'd)

Particle ^a	$I(J^P)L_{T,2J}$	P_{beam}^c (GeV/c) $\sigma = 4\pi\lambda^2$ (mb)	Mass ^d M (MeV)	Full ^e width Γ (MeV)	M_c^f $\pm \Gamma M$ (GeV ²)	Partial decay mode		
						Mode	Fraction ^h (%)	p^i (MeV/c)
S=-1 I=0 LAMBDA RESONANCES (Λ)								
Λ	$0(1/2^+)$		1115.6		1.245	See Stable Particle Table		
$\Lambda(1405)$	$0(1/2^-)S'_{01}$	Below K ⁻ p threshold	1405 $\pm 5^e$	40 ± 10^e	1.97 ± 0.06	$\Sigma\pi$	100	152
$\Lambda(1520)$	$0(3/2^-)D'_{03}$	$p = 0.395$ $\sigma = 82.2$	1519.4 $\pm 1.0^e$	15.6 ± 1.0^e	2.31 ± 0.02	$N\bar{K}$ $\Sigma\pi$ $\Lambda\pi\pi$ $\Sigma\pi\pi$	45 \pm 1 42 \pm 1 10 \pm 1 0.9 \pm 0.1	244 267 252 152
$\Lambda(1600)$	$0(1/2^+)P'_{01}$	$p = 0.58$ $\sigma = 41.6$	1560 to 1700	50 to 250 (150)	2.56 ± 0.24	$N\bar{K}$ $\Sigma\pi$	15-30 10-60	343 336
$\Lambda(1670)$	$0(1/2^-)S''_{01}$	$p = 0.74$ $\sigma = 28.5$	1660 to 1680	25 to 50 (35)	2.79 ± 0.06	$N\bar{K}$ $\Sigma\pi$ $\Lambda\eta$	15-25 20-60 15-35	414 393 64
$\Lambda(1690)$	$0(3/2^-)D''_{03}$	$p = 0.78$ $\sigma = 26.1$	1685 to 1695	50 to 70 (60)	2.86 ± 0.10	$N\bar{K}$ $\Sigma\pi$ $\Lambda\pi\pi$ $\Sigma\pi\pi$	20-30 20-40 ~25 ~20	433 409 415 350
$\Lambda(1800)$	$0(1/2^-)S''_{01}$	$p = 1.01$ $\sigma = 17.6$	1720 to 1850	200 to 400 (300)	3.24 ± 0.54	$N\bar{K}$ $\Sigma\pi$ $\Sigma(1385)\pi$ $N\bar{K}^*(892)$	25-40 seen seen seen	528 493 345 †
$\Lambda(1800)$	$0(1/2^+)P''_{01}$	$p = 1.01$ $\sigma = 17.6$	1750 to 1850	50 to 250 (150)	3.24 ± 0.27	$N\bar{K}$ $\Sigma\pi$ $\Sigma(1385)\pi$ $N\bar{K}^*(892)$	20-50 10-40 seen 30-60	528 493 345 †
$\Lambda(1820)$	$0(5/2^+)F'_{05}$	$p = 1.06$ $\sigma = 16.5$	1815 to 1825	70 to 90 (80)	3.29 ± 0.15	$N\bar{K}$ $\Sigma\pi$ $\Sigma(1385)\pi$	55-65 8-14 5-10	545 508 362
$\Lambda(1830)$	$0(5/2^-)D'_{05}$	$p = 1.08$ $\sigma = 16.0$	1810 to 1830	60 to 110 (95)	3.35 ± 0.17	$N\bar{K}$ $\Sigma\pi$ $\Sigma(1385)\pi$	3-10 35-75 >15	553 515 371
$\Lambda(1890)$	$0(3/2^+)F'_{03}$	$p = 1.21$ $\sigma = 13.6$	1850 to 1910	60 to 200 (100)	3.57 ± 0.19	$N\bar{K}$ $\Sigma\pi$ $\Sigma(1385)\pi$ $N\bar{K}^*(892)$	20-35 3-10 seen seen	599 559 420 233
$\Lambda(2100)$	$0(7/2^-)G'_{07}$	$p = 1.68$ $\sigma = 8.68$	2090 to 2110	100 to 250 (200)	4.41 ± 0.42	$N\bar{K}$ $\Sigma\pi$ $\Lambda\eta$ ΞK $\Delta\omega$ $N\bar{K}^*(892)$	25-35 ~ 5 < 3 < 3 < 8 10-20	751 704 617 483 445 514
$\Lambda(2110)$	$0(5/2^+)F'_{05}$	$p = 1.70$ $\sigma = 8.54$	2090 to 2140	150 to 250 (200)	4.45 ± 0.42	$N\bar{K}$ $\Sigma\pi$ $\Lambda\omega$ $\Sigma(1385)\pi$ $N\bar{K}^*(892)$	5-25 10-40 seen seen 10-60	757 711 455 589 524
$\Lambda(2350)$	$0(9/2^+)$	$p = 2.29$ $\sigma = 5.84$	2340 to 2370	100 to 250 (150)	5.52 ± 0.35	$N\bar{K}$ $\Sigma\pi$	~12 ~10	915 867
$\Lambda(2585)$	$0(?)$	$p = 2.92$ $\sigma = 4.35$	~2585	~300 (300)	6.68 ± 0.78	$N\bar{K}$	(J+1/2) _x ~1.0 _J	1060

Baryon Table (cont'd)

Particle ^a	$I(J^P)L_{\frac{1}{2}}^b$	P_{beam}^c (GeV/c) $\sigma = 4\pi\chi^2$ (mb)	Mass ^d M (MeV)	Full ^e width Γ (MeV)	M^2 ^f $\pm \Gamma M$ (GeV ²)	Partial decay mode		
						Mode	Fraction ^h (%)	p^i (MeV/c)
S=-1 I=1 SIGMA RESONANCES (Σ)								
Σ	$1(1/2^+)$		(+)1189.4 (0)1192.5 (-)1197.3		1.415 1.422 1.434	See Stable Particle Table		
$\Sigma(1385)$	$1(3/2^+)P'_{13}$	Below K^-p threshold	(+)1382.3 \pm 0.4 S=1.6 ^m (0)1382.0 \pm 2.5 S=1.6 ^m (-)1387.4 \pm 0.6 S=2.2 ^m	35 \pm 1 S=1.0 ^m ~35 40 \pm 2 S=1.9 ^m	1.92 \pm 0.05	$\Delta\pi$ $\Sigma\pi$	88 \pm 2 12 \pm 2	208 127
$\Sigma(1660)$	$1(1/2^+)P'_{11}$	$p = 0.72$ $\sigma = 29.8$	1630 to 1690	40 to 200 (100)	2.76 \pm 0.17	$N\bar{K}$ $\Delta\pi$ $\Sigma\pi$	10-30 seen seen	405 439 385
$\Sigma(1670)$	$1(3/2^-)D'_{13}$	$p = 0.74$ $\sigma = 28.5$	1665 to 1685	40 to 80 (60)	2.79 \pm 0.10	$N\bar{K}$ $\Delta\pi$ $\Sigma\pi$	7-13 5-15 30-60	414 447 393
$\Sigma(1750)$	$1(1/2^-)S'_{11}$	$p = 0.91$ $\sigma = 20.7$	1730 to 1800	60 to 160 (90)	3.06 \pm 0.16	$N\bar{K}$ $\Delta\pi$ $\Sigma\pi$ $\Sigma\eta$	10-40 seen < 8 15-55	486 507 455 81
$\Sigma(1775)$	$1(5/2^-)D'_{15}$	$p = 0.96$ $\sigma = 19.0$	1770 to 1780	105 to 135 (120)	3.15 \pm 0.21	$N\bar{K}$ $\Delta\pi$ $\Sigma\pi$ $\Sigma(1385)\pi$ $\Lambda(1520)\pi$	37-43 14-20 2-5 8-12 17-23	508 525 474 324 198
$\Sigma(1915)$	$1(5/2^+)F'_{15}$	$p = 1.26$ $\sigma = 12.8$	1900 to 1935	80 to 160 (120)	3.67 \pm 0.23	$N\bar{K}$ $\Delta\pi$ $\Sigma\pi$ $\Sigma(1385)\pi$	5-15 seen seen < 5	618 622 577 440
$\Sigma(1940)$	$1(3/2^-)D'_{13}$	$p = 1.32$ $\sigma = 12.1$	1900 to 1950	150 to 300 (220)	3.76 \pm 0.43	$N\bar{K}$ $\Delta\pi$ $\Sigma\pi$ $\Sigma(1385)\pi$ $\Lambda(1520)\pi$ $\Delta(1232)\bar{K}$ $N\bar{K}^*(892)$	<20 seen seen seen seen seen seen	637 639 594 460 354 410 320
$\Sigma(2030)$	$1(7/2^+)F'_{17}$	$p = 1.52$ $\sigma = 9.93$	2025 to 2040	150 to 200 (180)	4.12 \pm 0.37	$N\bar{K}$ $\Delta\pi$ $\Sigma\pi$ ΞK $\Sigma(1385)\pi$ $\Lambda(1520)\pi$ $\Delta(1232)\bar{K}$ $N\bar{K}^*(892)$	17-23 17-23 5-10 < 2 5-15 10-20 10-20 < 5	702 700 657 412 529 430 498 438
$\Sigma(2250)$	$1(?)$	$p = 2.04$ $\sigma = 6.76$	2210 to 2280	60 to 150 (100)	5.06 \pm 0.23	$N\bar{K}$ $\Delta\pi$ $\Sigma\pi$	<10 seen seen	851 842 803
$\Sigma(2455)$	$1(?)$	$p = 2.57$ $\sigma = 5.08$	~2455	~120 (120)	6.03 \pm 0.29	$N\bar{K}$	(J+1/2) _x ~0.2 _J	981
$\Sigma(2620)$	$1(?)$	$p = 3.02$ $\sigma = 4.19$	~2600	~200 (200)	6.86 \pm 0.52	$N\bar{K}$	(J+1/2) _x ~0.3 _J	1081

Baryon Table (cont'd)

Particle ^a	$1(J^P)L_{21}^b$	P ^c beam (GeV/c) $\sigma = 4\pi\lambda^2$ (mb)	Mass ^d M (MeV)	Full ^e width Γ (MeV)	M ^{2 f} $\pm \Gamma M$ (GeV ²)	Partial decay mode		
						Mode	Fraction (%)	p ⁱ (MeV/c)
S=-2 I=1/2 CASCADE RESONANCES (Ξ)								
Ξ	1/2(1/2 ⁺)		(0)1314.9 (-)1321.3		1.729 1.746	See Stable Particle Table		
$\Xi(1530)$	1/2(3/2 ⁺)P ₁₃		(0)1531.8±0.3 S = 1.3 ^m (-)1535.0±0.6	9.1±0.5 10.1±1.9	2.34 ±0.02	$\Xi\pi$	100	148
$\Xi(1820)$	1/2(3/2 ⁻)		1823 ±6 ^ℓ	20 ⁺¹⁵ ₋₁₀ ^ℓ	3.31 ±0.04	$\Lambda\bar{K}$ $\Sigma\bar{K}$ $\Xi\pi$ $\Xi(1530)\pi$	~45 ~10 small ~45	396 306 413 231
$\Xi(2030)$	1/2(?)		2024 ±6 ^ℓ	16 ⁺¹⁵ ₋₅ ^ℓ	4.12 ±0.03	$\Lambda\bar{K}$ $\Sigma\bar{K}$ $\Xi\pi$ $\Xi(1530)\pi$	~20 ~80 small small	587 524 573 418
Ω^-	0(3/2 ⁺)		1672.4±0.3		2.797	See Stable Particle Table		
Λ_c^+	0(1/2 ⁺)		2282±3		5.21	See Stable Particle Table		
S=0 DIBARYONS								
d	0(1 ⁺)		1875.6		3.518			
NN(2170)	1(2 ⁺) ¹ D ₂	p = 1.26 σ = 16.5	2140 to 2190	50 to 125 (90)	4.71 ±0.20	NN πd	10-20 seen	545 241
NN(2250)	1(3 ⁻) ³ F ₃	p = 1.49 σ = 12.7	2200 to 2400	75 to 225 (150)	5.06 ±0.34	NN πd	20-30 seen	621 318
S=-1 DIBARYONS								
$\Lambda N(2130)$	1/2(1 ⁺) ³ S ₁	p = 0.64 σ = 61.5	2100 to 2200	5 to 25 (15)	4.54 ±0.03	ΛN	seen	282

Baryon Table (*cont'd*)

- Each arrow in the left-hand margin indicates there is an entry in the Data Card Listings for a baryon that is not well enough established (status less than 3 stars) to be included here. There is a short list of *all* the baryons in the Listings, whatever their status, at the front of this Table.
- f.* This mode is energetically forbidden when the nominal mass of the decaying resonance (and of any resonance in the final state) is used, but is in fact allowed due to the finite width of the resonance(s).
- g.* The modes in brackets are sub-reactions of the first preceding unbracketed mode.
- a.* The nominal mass here (in MeV) is used for identification. See column 4 for the actual mass.
- b.* When there is more than one baryon with the same quantum numbers, one prime is attached to the spectroscopic symbol for the first of them (e.g., S'_{11}), two primes to the second, etc.
- c.* The quantities here are calculated using the nominal mass of column 1.
- d.* Usually a conservatively large range of masses rather than a statistical average of various determinations of the mass is given. In these cases, the mass determinations are nearly entirely from various phase-shift analyses of more or less the same data. It is thus not appropriate to treat the determinations as independent measurements or to average them together. The masses, widths, and branching fractions in this Table are Breit-Wigner parameters. The Data Card Listings also include pole parameters where they are available.
- e.* Usually a conservatively large range of widths rather than a statistical average of various determinations of the width is given (see note *d* for the reason). The nominal value in parentheses is then simply a best guess.
- f.* The quantities here are calculated using the nominal mass of column 1 and the nominal width of column 5.
- g.* For information on the $N\gamma$ decay modes of the N and Δ baryons, see the mini-review on these states in the Listings.
- h.* Most of the inelastic branching fractions come from partial-wave analyses, and these determine $\sqrt{x x'}$, where x and x' are the elastic and inelastic branching fractions, not x' directly. Thus any uncertainty (and it is often considerable) in x carries over into x' . When x' so determined is really poorly known, we here simply note that the mode is seen. The values of $\sqrt{x x'}$ are given in the Data Card Listings.
- i.* For a 2-body decay mode, this is the momentum of the decay products in the rest frame of the decaying particle. For a mode with more than two decay products, this is the maximum momentum any of the products can have in this frame. The nominal mass of column 1 is used, as is the nominal mass of any resonance in the final state.
- j.* The size of the bump in the total cross section gives $(J+1/2)x$, where x is the elastic branching fraction, but the value of J is not known.
- k.* These pole positions are from fits to phase shifts (without Coulomb corrections). The Data Card Listings now include pole positions and residues for most of the N and Δ resonances. See Sect. I of the N and Δ mini-review in the Listings for a brief discussion of the advantages of pole parameters over the usual Breit-Wigner parameters.
- l.* The error given here is only an educated guess. It is larger than the error on the weighted average of the published values (the error on this average is given in the Listings).
- m.* The error given here has been scaled up by the "S factor" (see the * footnote to the Stable Particle Table for how S is defined) because the various measurements disagree more seriously than one would expect from statistics.

PHYSICAL AND NUMERICAL CONSTANTS*

PHYSICAL CONSTANTS		Uncert. (ppm)
N_A	$= 6.022\ 045(31) \times 10^{23}$ mole ⁻¹	5.1
V_m	$= 22413.83(70)$ cm ³ mole ⁻¹ = molar volume of ideal gas at STP	31
c	$= 2.997\ 924\ 58(1.2) \times 10^{10}$ cm sec ⁻¹	0.004
e	$= 4.803\ 242(14) \times 10^{-10}$ esu = $1.602\ 189\ 2(46) \times 10^{-19}$ coulomb	2.9; 2.9
1 MeV	$= 1.602\ 189\ 2(46) \times 10^{-6}$ erg	2.9
$\hbar = h/2\pi$	$= 6.582\ 173(17) \times 10^{-22}$ MeV sec = $1.054\ 588\ 7(57) \times 10^{-27}$ erg sec	2.6; 5.4
$\hbar c$	$= 1.973\ 285\ 8(51) \times 10^{-11}$ MeV cm = $197.32858(51)$ MeV fermi	2.6; 2.6
$(\hbar c)^2$	$= 0.389\ 385\ 7(20)$ GeV ² mb	5.2
α	$= e^2/\hbar c = 1/137.03604(11)$	0.82
$k_{\text{Boltzmann}}$	$= 1.380\ 662(44) \times 10^{-16}$ erg °K ⁻¹	32
	$= 8.61735(28) \times 10^{-11}$ MeV °K ⁻¹ = $1\text{ eV}/11604.50(36)$ °K	32; 31
$\sigma_{\text{Stef. Boltz.}}$	$= 5.67032(71) \times 10^{-5}$ erg sec ⁻¹ cm ⁻² °K ⁻⁴	125
	$= 3.53911(44) \times 10^7$ eV sec ⁻¹ cm ⁻² °K ⁻⁴	125
m_e	$= 0.511\ 003\ 4(14)$ MeV = $9.109\ 534(47) \times 10^{-28}$ g	2.8; 5.1
m_p	$= 938.2796(27)$ MeV = $1836.15152(70)$ m_e = $6.722\ 775(39)$ m_{π^+}	2.8; 0.38; 5.8
	$= 1.007\ 276\ 470(11)$ amu	0.011
1 amu	$= 1/12\ m_{^{12}\text{C}}$ = $931.5016(26)$ MeV	2.8
m_d	$= 1875.6280(53)$ MeV	2.8
r_e	$= e^2/m_e c^2 = 2.817\ 938\ 0(70)$ fermi (1 fermi = 10^{-13} cm)	2.5
λ_e	$= \hbar/m_e c = r_e \alpha^{-1} = 3.861\ 590\ 5(64) \times 10^{-11}$ cm	1.6
$a_{\infty\text{Bohr}}$	$= \hbar^2/m_e e^2 = r_e \alpha^{-2} = 0.529\ 177\ 06(44)$ Å (1 Å = 10^{-8} cm)	0.82
σ_{Thomson}	$= (8/3)\pi r_e^2 = 0.665\ 244\ 8(33)$ barn (1 barn = 10^{-24} cm ²)	4.9
μ_{Bohr}	$= e\hbar/2m_e c = 0.578\ 837\ 85(95) \times 10^{-14}$ MeV gauss ⁻¹	1.6
μ_N	$= e\hbar/2m_p c = 3.152\ 451\ 5(53) \times 10^{-18}$ MeV gauss ⁻¹	1.7
μ_p/μ_{Bohr}	$= 0.001\ 521\ 032\ 209(16)$	0.011
$1/2\omega_{\text{cyclotron}}$	$= e/2m_e c = 8.794\ 024(25) \times 10^6$ rad sec ⁻¹ gauss ⁻¹	2.8
$1/2\omega_{\text{cyclotron}}$	$= e/2m_p c = 4.789\ 378(14) \times 10^3$ rad sec ⁻¹ gauss ⁻¹	2.8
Hydrogen-like atom (nonrelativistic, μ = reduced mass):		
	$\frac{v}{c}_{\text{rms}} = \frac{Z\alpha}{n}$; $E_n = \frac{\mu}{2} v^2 = \frac{\mu}{2} \left(\frac{cZ\alpha}{n}\right)^2$; $a_n = \frac{n^2 \hbar}{\mu Z c \alpha}$	
R_{∞}	$= m_e c^4/2\hbar^2 = m_e c^2 \alpha^2/2 = 13.605\ 804(36)$ eV (Rydberg)	2.6
	$= m_e c \alpha^2/2h = 109\ 737.3177(83)$ cm ⁻¹	0.075
ρ_c	$= 0.3$ H ρ (MeV, kilogauss, cm)	
1 year (sidereal)	$= 365.256$ days = 3.1558×10^7 sec ($\approx \pi \times 10^7$ sec)	
density of dry air	$= 1.204$ mg cm ⁻³ (at 20°C, 760 mm)	
acceleration by gravity	$= 980.62$ cm sec ⁻² (sea level, 45°)	
gravitational constant	$= 6.6720(41) \times 10^{-8}$ cm ⁵ g ⁻¹ sec ⁻²	615
1 calorie (thermochemical)	$= 4.184$ joules	
1 atmosphere	$= 1.01325$ bar (1 bar = 10^6 dynes cm ⁻²)	
1 eV per particle	$= 11604.50(36)$ °K (from $E = kT$)	31

NUMERICAL CONSTANTS

π	$= 3.141\ 592\ 7$	1 rad	$= 57.295\ 779\ 5$ deg	$\sqrt{\pi}$	$= 1.772\ 453\ 85$
e	$= 2.718\ 281\ 8$	1/e	$= 0.367\ 879\ 4$	$\sqrt{2}$	$= 1.414\ 213\ 6$
$\ln 2$	$= 0.693\ 147\ 2$	$\ln 10$	$= 2.302\ 585\ 1$	$\sqrt{3}$	$= 1.732\ 050\ 8$
$\log_{10} 2$	$= 0.301\ 030\ 0$	$\log_{10} e$	$= 0.434\ 294\ 5$	$\sqrt{10}$	$= 3.162\ 277\ 7$

* Revised 1982 by Barry N. Taylor. Originally prepared by Stanley J. Brodsky, based mainly on the "1973 Least-Squares Adjustment of the Fundamental Constants," by E. R. Cohen and B. N. Taylor, J. Phys. Chem. Ref. Data 2, 663 (1973). The figures in parentheses correspond to the one-standard-deviation uncertainty in the last digits of the main number. The equivalent uncertainty in parts per million (ppm) is given in the last column. Note that the uncertainties of the output values of a least-squares adjustment are in general correlated, and the general law of error propagation must be used in calculating additional quantities.

The set of constants resulting from the 1973 adjustment of Cohen and Taylor has been recommended for international use by CODATA (Committee on Data for Science and Technology), and is the most up-to-date, generally accepted set currently available. However, since the publication of the 1973 adjustment, a number of new experiments have been completed, yielding improved values for some of the constants: $N_A = 6.022\ 097\ 8(63) \times 10^{23}$ mole⁻¹ (1.04 ppm); $\alpha^{-1} = 137.035\ 963(15)$ (0.11 ppm); $m_p/m_e = 1836.15300(25)$ (0.14 ppm); and $R_{\infty} = 109\ 737.31521(11)$ cm⁻¹ (0.001 ppm). But it must be realized that, since the output values of a least-squares adjustment are related in a complex way and a change in the measured value of one constant usually leads to corresponding changes in the adjusted values of others, one must be cautious in carrying out calculations using both the output values from the 1973 adjustment and the results of more recent experiments. A new adjustment is planned for completion by mid 1982.

CLEBSCH-GORDAN COEFFICIENTS, SPHERICAL HARMONICS, AND d FUNCTIONS

Note: A $\sqrt{\quad}$ is to be understood over every coefficient; e. g., for $-8/15$ read $-\sqrt{8/15}$.

Notation:

J	J	\dots
M	M	\dots
m_1	m_2	Coefficients
\cdot	\cdot	
\cdot	\cdot	
\cdot	\cdot	

$$1/2 \times 1/2$$

$1/2$	$1/2$	1	0	0
$+1/2$	$+1/2$	1	0	0
$-1/2$	$-1/2$	$1/2$	$1/2$	1
$-1/2$	$+1/2$	$1/2$	$-1/2$	-1
$-1/2$	$-1/2$	1	0	0

$$Y_1^0 = \sqrt{\frac{3}{4\pi}} \cos \theta$$

$$Y_1^1 = -\sqrt{\frac{3}{8\pi}} \sin \theta e^{i\phi}$$

$$2 \times 1/2$$

$5/2$	$3/2$	2
$+5/2$	$+3/2$	2
$+2$	$+1/2$	1
$+1$	$+1/2$	1
$1/2$	$1/2$	1
0	$1/2$	1
$-1/2$	$1/2$	1
-1	$1/2$	1
$-3/2$	$1/2$	1
$-5/2$	$1/2$	1

$$Y_2^0 = \sqrt{\frac{5}{4\pi}} \left(\frac{3}{2} \cos^2 \theta - \frac{1}{2} \right)$$

$$Y_2^1 = -\sqrt{\frac{15}{8\pi}} \sin \theta \cos \theta e^{i\phi}$$

$$Y_2^2 = \frac{1}{4} \sqrt{\frac{15}{2\pi}} \sin^2 \theta e^{2i\phi}$$

$$3/2 \times 1/2$$

$3/2$	$1/2$	2
$+3/2$	$+1/2$	2
$+1$	$+1$	1
$1/2$	$1/2$	1
0	$1/2$	1
$-1/2$	$1/2$	1
-1	$1/2$	1
$-3/2$	$1/2$	1

$$1 \times 1/2$$

$3/2$	$1/2$
$+3/2$	$+1/2$
$+1$	$+1/2$
$1/2$	$1/2$
0	$1/2$
$-1/2$	$1/2$
-1	$1/2$
$-3/2$	$1/2$

$$2 \times 1$$

3	2	1
$+3$	$+2$	$+1$
$+2$	$+1$	0
$+1$	0	-1
0	-1	-2
-1	-2	-3

$$3/2 \times 1$$

$5/2$	$3/2$	2
$+5/2$	$+3/2$	$+2$
$+3$	$+1$	0
$+1/2$	$1/2$	1
0	$1/2$	1
$-1/2$	$1/2$	1
-3	$1/2$	1
$-5/2$	$1/2$	1

$$1 \times 1$$

2	1
$+2$	$+1$
$+1$	0
0	-1
-1	-2

$3/2$	$1/2$	2
$+3/2$	$+1/2$	$+2$
$+1$	$+1$	1
$1/2$	$1/2$	1
0	$1/2$	1
$-1/2$	$1/2$	1
-1	$1/2$	1
$-3/2$	$1/2$	1

2	1	0
$+2$	$+1$	0
$+1$	0	-1
0	-1	-2
-1	-2	-3

$3/2$	$1/2$	2
$+3/2$	$+1/2$	$+2$
$+1$	$+1$	1
$1/2$	$1/2$	1
0	$1/2$	1
$-1/2$	$1/2$	1
-1	$1/2$	1
$-3/2$	$1/2$	1

2	1	0
$+2$	$+1$	0
$+1$	0	-1
0	-1	-2
-1	-2	-3

$3/2$	$1/2$	2
$+3/2$	$+1/2$	$+2$
$+1$	$+1$	1
$1/2$	$1/2$	1
0	$1/2$	1
$-1/2$	$1/2$	1
-1	$1/2$	1
$-3/2$	$1/2$	1

2	1	0
$+2$	$+1$	0
$+1$	0	-1
0	-1	-2
-1	-2	-3

$3/2$	$1/2$	2
$+3/2$	$+1/2$	$+2$
$+1$	$+1$	1
$1/2$	$1/2$	1
0	$1/2$	1
$-1/2$	$1/2$	1
-1	$1/2$	1
$-3/2$	$1/2$	1

2	1	0
$+2$	$+1$	0
$+1$	0	-1
0	-1	-2
-1	-2	-3

$3/2$	$1/2$	2
$+3/2$	$+1/2$	$+2$
$+1$	$+1$	1
$1/2$	$1/2$	1
0	$1/2$	1
$-1/2$	$1/2$	1
-1	$1/2$	1
$-3/2$	$1/2$	1

2	1	0
$+2$	$+1$	0
$+1$	0	-1
0	-1	-2
-1	-2	-3

$3/2$	$1/2$	2
$+3/2$	$+1/2$	$+2$
$+1$	$+1$	1
$1/2$	$1/2$	1
0	$1/2$	1
$-1/2$	$1/2$	1
-1	$1/2$	1
$-3/2$	$1/2$	1

2	1	0
$+2$	$+1$	0
$+1$	0	-1
0	-1	-2
-1	-2	-3

$3/2$	$1/2$	2
$+3/2$	$+1/2$	$+2$
$+1$	$+1$	1
$1/2$	$1/2$	1
0	$1/2$	1
$-1/2$	$1/2$	1
-1	$1/2$	1
$-3/2$	$1/2$	1

$$Y_l^{-m} (-1)^m Y_l^m e^{-im\phi}$$

$$Y_{lm,0}^m \sqrt{\frac{4\pi}{2l+1}} Y_l^m e^{-im\phi}$$

$$\langle J_1 J_2 m_1 m_2 | J_1 J_2 J M \rangle = (-1)^{J_2 - J_1 - J} \langle J_1 J_2 m_1 m_2 | J_1 J_2 J M \rangle$$

$$d_{m',m}^{j,j} = (-1)^{m-m'} d_{m,m'}^{j,j} = d_{-m,-m'}^{j,j}$$

$$2 \times 3/2$$

$7/2$	$5/2$
$+7/2$	$+5/2$
$+3$	$+1$
1	1
0	1
-1	1
-3	1
$-5/2$	1
$-7/2$	1

$7/2$	$5/2$	4	3	2	1
$+7/2$	$+5/2$	$+4$	$+3$	$+2$	$+1$
$+2$	$+1$	0	-1	-2	-3
$+1$	0	-1	-2	-3	-4
0	-1	-2	-3	-4	-5
-1	-2	-3	-4	-5	-6
-2	-3	-4	-5	-6	-7
-3	-4	-5	-6	-7	-8
-4	-5	-6	-7	-8	-9
-5	-6	-7	-8	-9	-10

$7/2$	$5/2$	4	3	2	1
$+7/2$	$+5/2$	$+4$	$+3$	$+2$	$+1$
$+2$	$+1$	0	-1	-2	-3
$+1$	0	-1	-2	-3	-4
0	-1	-2	-3	-4	-5
-1	-2	-3	-4	-5	-6
-2	-3	-4	-5	-6	-7
-3	-4	-5	-6	-7	-8
-4	-5	-6	-7	-8	-9
-5	-6	-7	-8	-9	-10

$7/2$	$5/2$	4	3	2	1
$+7/2$	$+5/2$	$+4$	$+3$	$+2$	$+1$
$+2$	$+1$	0	-1	-2	-3
$+1$	0	-1	-2	-3	-4
0	-1	-2	-3	-4	-5
-1	-2	-3	-4	-5	-6
-2	-3	-4	-5	-6	-7
-3	-4	-5	-6	-7	-8
-4	-5	-6	-7	-8	-9
-5	-6	-7	-8	-9	-10

$7/2$	$5/2$	4	3	2	1
$+7/2$	$+5/2$	$+4$	$+3$	$+2$	$+1$
$+2$	$+1$	0	-1	-2	-3
$+1$	0	-1	-2	-3	-4
0	-1	-2	-3	-4	-5
-1	-2	-3	-4	-5	-6
-2	-3	-4	-5	-6	-7
-3	-4	-5	-6	-7	-8
-4	-5	-6	-7	-8	-9
-5	-6	-7	-8	-9	-10

$7/2$	$5/2$	4	3	2	1
$+7/2$	$+5/2$	$+4$	$+3$	$+2$	$+1$
$+2$	$+1$	0	-1	-2	-3
$+1$	0	-1	-2	-3	-4
0	-1	-2	-3	-4	-5
-1	-2	-3	-4	-5	-6
-2	-3	-4	-5	-6	-7
-3	-4	-5	-6	-7	-8
-4	-5	-6	-7	-8	-9
-5	-6	-7	-8	-9	-10

$7/2$	$5/2$	4	3
-------	-------	-----	-----

SU(3) ISOSCALAR FACTORS, SU(N) MULTIPLICITIES, AND PROPERTIES OF QUARKS

The most commonly used isoscalar factors, corresponding to the singlet, octet, and decuplet content of $8 \otimes 8$ and $10 \otimes 8$, are displayed at the right. The notation uses particle names to identify the coefficients, so that the pattern of relative couplings can be seen at a glance. We illustrate the use of the coefficients by example; see J. J. de Swart, *Rev. Mod. Phys.* **35**, 916 (1963) for detailed explanation and phase conventions.

A $\sqrt{\quad}$ is understood over every integer in the matrices; the exponent $\frac{1}{2}$ is a reminder of this. For example, in de Swart's notation the $\Xi \rightarrow \Sigma K$ element of our $10 \rightarrow 10 \otimes 8$ matrix reads

$$\begin{pmatrix} 10 & & & & & \\ 0 & -2 & & \frac{1}{2} & & \\ & & & & & \\ & & & & & \\ & & & & & \\ & & & & & \end{pmatrix} \begin{matrix} 8 \\ 1 \\ \frac{1}{2} \\ \frac{1}{2} \\ -1 \end{matrix} = \frac{-\sqrt{6}}{\sqrt{24}}$$

Intra-multiplet relative decay strengths can be read directly from our matrices. Thus, the partial widths for $\Delta^* \rightarrow (N\pi)_{I=3/2}$ and $\Sigma^* \rightarrow (\Sigma K)_{I=0}$ are in the ratio

$$\frac{\Gamma(\Delta^* \rightarrow (\Sigma K)_{I=0})}{\Gamma(\Delta^* \rightarrow (N\pi)_{I=3/2})} = \frac{12}{6} \times (\text{threshold factors})$$

Supplying isospin Clebsch-Gordan coefficients one obtains, e.g.,

$$\frac{\Gamma(\Delta^* \rightarrow \Sigma^0 K^+)_{I=0}}{\Gamma(\Delta^* \rightarrow p\pi^0)_{I=3/2}} = \frac{1/2 \times 1/2}{2/3 \times 6} \times \text{cf} = \frac{3}{2} \times \text{cf}$$

Partial widths for $8 \rightarrow 8 \otimes 8$ involve a linear superposition of B_1 (symmetric) and B_2 (anti-symmetric) couplings. For example,

$$\Gamma(\Xi^* \rightarrow \Xi\pi) \sim \left(-\sqrt{\frac{9}{20}} g_1 + \sqrt{\frac{3}{12}} g_2 \right)^2$$

The relation between g_1, g_2 (with de Swart's normalization) and the standard D,F couplings appearing in the interaction Lagrangian,

$$\mathcal{L} = -\sqrt{2} D \text{Tr}([\bar{B}, B]_+ M) + \sqrt{2} F \text{Tr}([\bar{B}, B]_- M)$$

is

$$D = \frac{\sqrt{30}}{40} g_1, \quad F = \frac{\sqrt{6}}{24} g_2$$

Thus,

$$\Gamma(\Xi^* \rightarrow \Xi\pi) \sim (1 - 2a)^2$$

where $a \equiv D/(D+F)$.

$1 \rightarrow 8 \otimes 8$

$$\begin{pmatrix} \Lambda \\ \Sigma \\ \Xi \end{pmatrix}_1 \rightarrow \begin{pmatrix} N\bar{K} & \Sigma\pi & \Lambda\eta & \Sigma K \\ \Sigma\bar{K} & \Lambda\bar{K} & \Xi\pi & \Xi\eta \end{pmatrix}_{8 \otimes 8} = \frac{1}{\sqrt{8}} \begin{pmatrix} 2 & 3 & -1 & -2 \end{pmatrix}^{\frac{1}{2}}$$

$B_1 \rightarrow 8 \otimes 8$

$$\begin{pmatrix} N \\ \Sigma \\ \Lambda \\ \Xi \end{pmatrix}_{B_1} \rightarrow \begin{pmatrix} N\pi & N\eta & \Sigma K & \Lambda K \\ N\bar{K} & \Sigma\pi & \Lambda\eta & \Sigma K \\ \Sigma\bar{K} & \Lambda\bar{K} & \Xi\pi & \Xi\eta \end{pmatrix}_{8 \otimes 8} = \frac{1}{\sqrt{24}} \begin{pmatrix} 9 & -1 & -9 & -1 \\ -6 & 0 & 4 & 4 & -6 \\ 2 & -12 & -4 & -2 \\ 9 & -1 & -9 & -1 \end{pmatrix}^{\frac{1}{2}}$$

$B_2 \rightarrow 8 \otimes 8$

$$\begin{pmatrix} N \\ \Sigma \\ \Lambda \\ \Xi \end{pmatrix}_{B_2} \rightarrow \begin{pmatrix} N\pi & N\eta & \Sigma K & \Lambda K \\ N\bar{K} & \Sigma\pi & \Lambda\eta & \Sigma K \\ \Sigma\bar{K} & \Lambda\bar{K} & \Xi\pi & \Xi\eta \end{pmatrix}_{8 \otimes 8} = \frac{1}{\sqrt{12}} \begin{pmatrix} 3 & 3 & 3 & -3 \\ 2 & 8 & 0 & 0 & -2 \\ 6 & 0 & 0 & 6 \\ 3 & 3 & 3 & -3 \end{pmatrix}^{\frac{1}{2}}$$

$10 \rightarrow 8 \otimes 8$

$$\begin{pmatrix} \Delta \\ \Sigma \\ \Xi \\ \Omega \end{pmatrix}_{10} \rightarrow \begin{pmatrix} N\pi & & \Sigma K \\ N\bar{K} & \Sigma\pi & \Lambda\eta & \Sigma\eta & \Xi\eta \\ \Sigma\bar{K} & \Lambda\bar{K} & \Xi\pi & \Xi\eta \\ \Xi\bar{K} & & & \Omega K \end{pmatrix}_{8 \otimes 8} = \frac{1}{\sqrt{12}} \begin{pmatrix} -6 & 6 \\ -2 & 2 & -3 & 3 & 2 \\ 3 & -3 & 3 & 3 \\ & & & & 12 \end{pmatrix}^{\frac{1}{2}}$$

$8 \rightarrow 10 \otimes 8$

$$\begin{pmatrix} N \\ \Sigma \\ \Lambda \\ \Xi \end{pmatrix}_8 \rightarrow \begin{pmatrix} \Delta\pi & \Sigma K \\ \Delta\bar{K} & \Sigma\pi & \Lambda\eta & \Sigma K \\ \Sigma\pi & \Sigma\eta & \Xi\eta \\ \Sigma\bar{K} & \Xi\pi & \Omega K \end{pmatrix}_{10 \otimes 8} = \frac{1}{\sqrt{15}} \begin{pmatrix} -12 & 3 \\ 8 & -2 & -3 & 2 \\ -9 & 6 \\ 3 & -3 & -3 & 6 \end{pmatrix}^{\frac{1}{2}}$$

$10 \rightarrow 10 \otimes 8$

$$\begin{pmatrix} \Delta \\ \Sigma \\ \Xi \\ \Omega \end{pmatrix}_{10} \rightarrow \begin{pmatrix} \Delta\pi & \Delta\eta & \Sigma K \\ \Delta\bar{K} & \Sigma\pi & \Lambda\eta & \Sigma K \\ \Sigma\bar{K} & \Xi\pi & \Omega K \\ \Xi\bar{K} & \Omega\pi & & \end{pmatrix}_{10 \otimes 8} = \frac{1}{\sqrt{24}} \begin{pmatrix} 15 & 3 & -6 \\ 8 & 8 & 0 & -8 \\ 12 & 3 & -3 & -6 \\ 12 & -12 \end{pmatrix}^{\frac{1}{2}}$$

SU(n) Multiplicities

The table below gives the multiplicities of the multiplets that occur in qq, q \bar{q} , and qqq systems in various SU(n). Normal mesons are q \bar{q} systems, and normal baryons are qqq systems. Also given are the multiplets that occur in meson-baryon scattering when the meson multiplet is the one to which the pion belongs and the

baryon multiplet is the one to which the proton belongs. Complex-conjugate representations are indicated by a bar. The two 20-dimensional representations of SU(4) are indicated as 20 (which contains the SU(3) decuplet) and 20' (which contains the SU(3) octet). The C(N,M)'s are the binomial coefficients N!/[M!(N-M)!].

qq	
SU(2):	2 ⊗ 2 = 3 ⊕ 1
SU(3):	3 ⊗ 3 = 6 ⊕ 3̄
SU(4):	4 ⊗ 4 = 10 ⊕ 6̄
SU(n):	n ⊗ n = n(n+1)/2 ⊕ n(n-1)/2
q \bar{q} (Mesons)	
SU(2):	2 ⊗ 2̄ = 3 ⊕ 1
SU(3):	3 ⊗ 3̄ = 8 ⊕ 1
SU(4):	4 ⊗ 4̄ = 15 ⊕ 1
SU(n):	n ⊗ n̄ = (n ² -1) ⊕ 1

qqq (Baryons)	
SU(2):	2 ⊗ 2 ⊗ 2 = 4 ⊕ 2 ⊕ 2
SU(3):	3 ⊗ 3 ⊗ 3 = 10 ⊕ 8 ⊕ 8 ⊕ 1
SU(4):	4 ⊗ 4 ⊗ 4 = 20 ⊕ 20' ⊕ 20' ⊕ 4̄
SU(n):	n ⊗ n ⊗ n = C(n+2,3) ⊕ 2C(n+1,3) ⊕ 2C(n+1,3) ⊕ C(n,3)
Meson-Baryon Scattering	
SU(2):	3 ⊗ 2 = 4 ⊕ 2
SU(3):	8 ⊗ 8 = 27 ⊕ 10 ⊕ 10̄ ⊕ 8 ⊕ 8 ⊕ 1
SU(4):	15 ⊗ 20' = 140 ⊕ 60 ⊕ 36 ⊕ 20 ⊕ 20' ⊕ 20' ⊕ 4̄

SU(3) ISOSCALAR FACTORS, SU(N) MULTIPLICITIES, AND PROPERTIES OF QUARKS (Cont'd)

Properties of Quarks

Quantum number	Quark					
	d	u	s	c	b	t
baryon number B	1/3	1/3	1/3	1/3	1/3	1/3
charge Q	-1/3	+2/3	-1/3	+2/3	-1/3	+2/3
isospin z component I _z	-1/2	+1/2	0	0	0	0
strangeness S	0	0	-1	0	0	0
charm C	0	0	0	+1	0	0
bottom (beauty) B	0	0	0	0	+1	0
top (truth) T	0	0	0	0	0	+1

WEAK INTERACTIONS OF QUARKS AND LEPTONS

The "standard" SU(2)_L ⊗ U(1) model^{1,2} is described here for six quarks and six leptons in left-handed doublets of SU(2)_{weak} and right-handed singlets of SU(2)_{weak} (T₃ = third component of weak isospin):

$$\begin{aligned}
 T_3 = +1/2 & \begin{pmatrix} \nu_e \\ e^- \end{pmatrix}_L, \begin{pmatrix} \nu_u \\ u^- \end{pmatrix}_L, \begin{pmatrix} \nu_c \\ c^- \end{pmatrix}_L, \begin{pmatrix} \nu_s \\ s^- \end{pmatrix}_L, \begin{pmatrix} \nu_b \\ b^- \end{pmatrix}_L, \begin{pmatrix} \nu_t \\ t^- \end{pmatrix}_L \\
 T_3 = -1/2 & \begin{pmatrix} \bar{\nu}_e \\ e^+ \end{pmatrix}_R, \begin{pmatrix} \bar{\nu}_u \\ u^+ \end{pmatrix}_R, \begin{pmatrix} \bar{\nu}_c \\ c^+ \end{pmatrix}_R, \begin{pmatrix} \bar{\nu}_s \\ s^+ \end{pmatrix}_R, \begin{pmatrix} \bar{\nu}_b \\ b^+ \end{pmatrix}_R, \begin{pmatrix} \bar{\nu}_t \\ t^+ \end{pmatrix}_R \\
 T = T_3 = 0 & e^-_R, u^-_R, \bar{\nu}_e, \bar{\nu}_u, \bar{\nu}_c, \bar{\nu}_s, \bar{\nu}_b, \bar{\nu}_t, d_R, c_R, s_R, t_R, b_R.
 \end{aligned}$$

Mixing occurs between quarks d, s, b of charge -1/3 (by convention the charge 2/3 quarks, u, c, t, are unmixed) and can be written

$$\begin{pmatrix} d' \\ s' \\ u' \end{pmatrix} = \begin{pmatrix} V_{ud} & V_{us} & V_{ub} \\ V_{cd} & V_{cs} & V_{cb} \\ V_{td} & V_{ts} & V_{tb} \end{pmatrix} \begin{pmatrix} d \\ s \\ b \end{pmatrix}$$

where, for example, the mixing matrix element V_{ud} modulates the strength of the udW vertex. In the Kobayashi-Maskawa parametrization² this matrix is expressed

$$\begin{pmatrix} d' \\ s' \\ b' \end{pmatrix} = \begin{pmatrix} c_1 & s_1 c_3 & s_1 s_3 \\ -s_1 c_2 & c_1 c_2 c_3 + s_2 s_3 e^{i\delta} & c_1 c_2 s_3 - s_2 c_3 e^{i\delta} \\ -s_1 s_2 & c_1 s_2 c_3 - c_2 s_3 e^{i\delta} & c_1 s_2 s_3 + c_2 c_3 e^{i\delta} \end{pmatrix} \begin{pmatrix} d \\ s \\ b \end{pmatrix}$$

where c_i = cos θ_i, s_i = sin θ_i, i = 1, 2, 3. In the limit θ₂ = θ₃ = δ = 0, this reduces to the usual Cabibbo mixing with θ₁ th. Cabibbo angle.

The interaction Lagrangian is

$$\mathcal{L}_{int} = e \left[A \bar{\psi} \gamma_\mu \psi + \frac{1}{\sin \theta_w \cos \theta_w} 2 \partial_\mu N + \frac{1}{\sqrt{2} \sin \theta_w} (W^+ \partial_\mu \psi^c + W^- \partial_\mu \psi^c) \right]$$

Here θ_w is the weak mixing angle in the relations

$$W^0 = Z \cos \theta_w + A \sin \theta_w$$

$$B = -Z \sin \theta_w + A \cos \theta_w$$

which relate the physical fields A (photon) and Z (neutral weak gauge boson) to W⁰ (SU(2)_{weak} partner of W[±] and W⁰) and B (U(1) gauge field). The charged current is written

$$\mathcal{J}_\mu^c = (\bar{\nu}_e \bar{\nu}_u \bar{\nu}_c) \left[\gamma_\mu \frac{(1-\gamma_5)}{2} \begin{pmatrix} e^- \\ \mu^- \\ \tau^- \end{pmatrix} \right] + (\bar{u} \bar{c} \bar{t}) \left[\gamma_\mu \frac{(1-\gamma_5)}{2} \begin{pmatrix} u \\ c \\ t \end{pmatrix} \right]$$

i.e., V-A structure. The neutral current is written

$$\begin{aligned}
 \mathcal{J}_\mu^N &= (\bar{\nu}_e \bar{\nu}_u \bar{\nu}_c) \left[\frac{1}{2} \gamma_\mu \frac{(1-\gamma_5)}{2} \begin{pmatrix} \nu_e \\ \nu_u \\ \nu_c \end{pmatrix} \right] \\
 &+ (\bar{e} \bar{\mu} \bar{\tau}) \left[-\frac{1}{2} \gamma_\mu \frac{(1-\gamma_5)}{2} + \sin^2 \theta_w \gamma_\mu \begin{pmatrix} e \\ \mu \\ \tau \end{pmatrix} \right] \\
 &+ (\bar{u} \bar{c} \bar{t}) \left[\frac{1}{2} \gamma_\mu \frac{(1-\gamma_5)}{2} - \frac{2}{3} \sin^2 \theta_w \gamma_\mu \begin{pmatrix} u \\ c \\ t \end{pmatrix} \right] \\
 &+ (\bar{d} \bar{s} \bar{b}) \left[-\frac{1}{2} \gamma_\mu \frac{(1-\gamma_5)}{2} + \frac{1}{3} \sin^2 \theta_w \gamma_\mu \begin{pmatrix} d \\ s \\ b \end{pmatrix} \right]
 \end{aligned}$$

where for fermion f the coupling [f^c]_μ has a V-A term depending on T₃ and a vector term depending on charge Q_f:

$$[f^c]_\mu = \left[\frac{1}{2} \gamma_\mu \frac{(1-\gamma_5)}{2} - Q_f \sin^2 \theta_w \gamma_\mu \right]$$

The effective Lagrangian for exchange of W[±] and Z between two currents reduces at low q² to

$$\mathcal{L}_{weak} = \frac{G}{\sqrt{2}} 4 \left(\bar{\psi} \gamma_\mu \psi^c + 2 \partial_\mu N \bar{\psi} \psi \right)$$

with G/√2 = g_w/(2M_w² sin² θ_w), α = e²/(4π), and ρ = M_w²/(M_Z² cos² θ_w). Assuming the simplest Higgs structure, ρ = 1, and the W and Z masses are related by M_w = M_Z cos θ_w. Currently reported values of the weak interaction parameters are

$$\left. \begin{aligned}
 |V_{ud}| &= |\cos \theta_1| = 0.9737 \pm 0.0025 \\
 |V_{us}| &= |\sin \theta_1 \cos \theta_3| = 0.219 \pm 0.011 \\
 |V_{ub}| &= |\sin \theta_1 \sin \theta_3| = 0.06 \pm 0.06 \\
 |V_{cb}|^2 &\gg |V_{ub}|^2 \\
 G = G_W &= (1.16632 \pm 0.00002) \times 10^{-5} \text{ GeV}^{-2} \\
 \sin^2 \theta_w &= 0.224 \pm 0.019, \quad \rho = 0.992 \pm 0.020 \\
 \sin^2 \theta_w &= 0.229 \pm 0.010, \quad \rho = 1 \text{ (fixed)}
 \end{aligned} \right\} \text{ Refs. 3, 4, 5, 8, 9}$$

The resulting mass estimates for W[±] and Z are M_w = 37.3 GeV/sin θ_w = 77.9 ± 1.7 GeV, and M_Z = 88.8 ± 1.4 GeV, where the numerical values are obtained using the simplest Higgs structure (ρ = 1). Electroweak radiative corrections to these estimates may be as large as 5%.¹⁰

Lepton-Nucleon Inclusive Scattering

For reactions ℓN + ℓ' + X, differential cross sections can be written using several choices of independent variables. These are related by

$$\begin{aligned}
 \frac{d^2\sigma}{dx dy} &= 2HE_L^2 \frac{d^2\sigma}{dxdy^2} = 2HE_L^2 \frac{d^2\sigma}{dx dQ^2} = \frac{2\pi ME_L^2 y}{|P_i| |P_f|} \frac{d^2\sigma}{dE_L dE_L'} \\
 &\approx \frac{2\pi ME_L^2 y}{E_L^2} \frac{d^2\sigma}{dE_L dE_L'}
 \end{aligned}$$

where ν, Q², x, and y are defined in the Relativistic Kinematics section IV(b), E_L, E_{L'} and E_p, E_{p'} are the incident and outgoing lepton lab energies and momenta, and M is the target nucleon mass.

WEAK INTERACTIONS OF QUARKS AND LEPTONS (Cont'd)

Structure Functions^{11,12}

For charged-current (C.C.) and neutral-current (N.C.) reactions, we have

$$\frac{d^2\sigma_{\nu}(\bar{\nu})}{dx dy} = \frac{G^2 M_E^2}{\pi} \left[(1 - \gamma - \frac{y}{2E}) xy F_2^{\nu}(\bar{\nu})(x, Q^2) - \frac{y^2}{2} 2x F_3^{\nu}(\bar{\nu})(x, Q^2) \pm (y - \frac{y^2}{2}) x F_3^{\nu}(\bar{\nu})(x, Q^2) \right],$$

where the upper and lower signs refer to ν and $\bar{\nu}$ scattering, respectively, and F_3 is defined as a positive quantity.¹³ The other common structure functions W_1 are related by $W_1 = F_1$, $W_2 = F_2$, and $W_3 = F_3$. For electron and muon scattering, $F_3 = 0$, and G^2 is replaced by $8\pi^2\alpha^2/(Q^2)^2$. The ratio of the longitudinally to transversely polarized photon absorption cross section is

$$R = \frac{\gamma}{\alpha} = \frac{1}{2x F_1(x, Q^2)} \left[F_2(x, Q^2) - 2x F_1(x, Q^2) - \frac{4M^2 x^2}{Q^2} F_2(x, Q^2) \right].$$

To compare with the parton-model predictions below, we write for $E_k \gg M$:

$$\frac{d^2\sigma_{\nu}(\bar{\nu})}{dx dy} = \frac{G^2 M_E^2}{\pi} \left[\frac{1}{2} \left(2x F_1^{\nu}(\bar{\nu}) + x F_3^{\nu}(\bar{\nu}) \right) + \frac{1}{2} (1-y)^2 \left(2x F_1^{\nu}(\bar{\nu}) + x F_3^{\nu}(\bar{\nu}) \right) + (1-y) \left(F_2^{\nu}(\bar{\nu}) - 2x F_1^{\nu}(\bar{\nu}) \right) \right],$$

$$\frac{d^2\sigma_{e, \mu}}{dx dy} = \frac{8\pi^2\alpha^2 M_E^2}{(Q^2)^2} \left[\frac{1 + (1-y)^2}{2} 2x F_1^{e, \mu} + (1-y) \left(F_2^{e, \mu} - 2x F_1^{e, \mu} \right) \right].$$

The Free-Quark-Parton-Model Predictions¹⁴

For this model in the Bjorken limit ($Q^2, \nu \rightarrow \infty$ with x fixed), $F_1(x, Q^2) = F_1(x)$,¹⁵ For spin- $\frac{1}{2}$ quark partons, we have $2x F_1(x) = F_2(x)$, the Callan-Gross relation. Thus, in this approximation, $R=0$ and there is no $(1-y)$ term in the cross section.

$$\bullet \text{ (C.C.) } \frac{d^2\sigma_{\nu N} - \mu^+ X}{dx dy} = \frac{G^2 M_E^2}{\pi} 2x \sum_q \left[f_q(x) + f_{\bar{q}}(x)(1-y)^2 \right].$$

For $\bar{\nu} N + \mu^+ X$, interchange $f_q(x)$ and $f_{\bar{q}}(x)$ in the formula. Here $f_q(x) dx$ is the number of quarks q in the target nucleus with momentum fraction x to $x+dx$. We include $f_q(x)$ and $f_{\bar{q}}(x)$ in the sum only for neutrons (positive) charged quarks and antiquarks in $\nu(\bar{\nu})$ reactions.

$$\bullet \text{ (N.C.) } \frac{d^2\sigma_{\nu N} - \nu X}{dx dy} = \frac{G^2 M_E^2}{\pi} 2p^2 x \sum_q \left\{ (e_q^2)^2 \left[f_q(x) + f_{\bar{q}}(x)(1-y)^2 \right] + (e_q^2)^2 \left[f_{\bar{q}}(x)(1-y)^2 + f_q(x) \right] \right\},$$

and the sum runs over all quarks. Here the neutral-current coupling is decomposed according to

$$f_q^N = c_L^q \gamma_5 \frac{(1-\gamma_5)}{2} + c_R^q \gamma_5 \frac{(1+\gamma_5)}{2}$$

with left- and right-handed coupling constants c_L^q and c_R^q .¹⁶ In the "standard" SU(2) \otimes U(1) model

$$c_L^q = T_3^q - Q_q \sin^2 \theta_W, \quad c_R^q = -Q_q \sin^2 \theta_W.$$

For $\bar{\nu} N + \bar{\nu} X$, interchange c_L^q and c_R^q in the cross-section formula.

$$\bullet \text{ (E.M.) } \frac{d^2\sigma_{e, \mu}^{\nu, \bar{\nu}}}{dx dy} = \frac{8\pi^2\alpha^2 M_E^2}{(Q^2)^2} x \sum_q \left[f_q(x) + f_{\bar{q}}(x) \right] \frac{1 + (1-y)^2}{2}.$$

Comparison with earlier structure function formulas gives:

$$\text{(C.C.) } F_2(x) = 2x \sum_q \left[f_q(x) + f_{\bar{q}}(x) \right],$$

$$x F_3(x) = 2x \sum_q \left[f_q(x) - f_{\bar{q}}(x) \right];$$

$$\text{(N.C.) } F_2(x) = 2e^2 x \sum_q \left[(e_q^L)^2 + (e_q^R)^2 \right] \left[f_q(x) + f_{\bar{q}}(x) \right],$$

$$x F_3(x) = 2e^2 x \sum_q \left[(e_q^L)^2 - (e_q^R)^2 \right] \left[f_q(x) - f_{\bar{q}}(x) \right],$$

$$F_3^{\nu}(x) = F_3^{\bar{\nu}}(x);$$

$$\text{(E.M.) } F_2(x) = x \sum_q Q_q^2 \left[f_q(x) + f_{\bar{q}}(x) \right].$$

In the examples below, $u(x)$, $\bar{u}(x)$, $d(x)$, $\bar{d}(x)$, etc., mean f_q ($f_{\bar{q}}$) for the individual quark (antiquark) in the proton (for neutron, interchange $u(x)$ and $d(x)$). Charm production is taken into account.

$$\bullet F_2^{\nu p} + \mu^+ X = 2x[d(x) + s(x) + \bar{u}(x) + \bar{c}(x)],$$

$$F_2^{\bar{\nu} p} + \mu^+ X = 2x[u(x) + c(x) + \bar{d}(x) + \bar{s}(x)],$$

$$x F_3^{\nu p} + \mu^+ X = 2x[d(x) + s(x) - \bar{u}(x) - \bar{c}(x)],$$

$$x F_3^{\bar{\nu} p} + \mu^+ X = 2x[u(x) + c(x) - \bar{d}(x) - \bar{s}(x)].$$

Hereafter we neglect small contributions of the s , \bar{s} , c , \bar{c} quarks in the sea.

$$\bullet \text{ For charge-symmetric nuclei with } q(x) = u(x) + d(x), \bar{q}(x) = \bar{u}(x) + \bar{d}(x),$$

$$F_2^{\nu N} + \mu^+ X = F_2^{\bar{\nu} N} + \mu^+ X = x[q(x) + \bar{q}(x)],$$

$$x F_3^{\nu N} + \mu^+ X = x F_3^{\bar{\nu} N} + \mu^+ X = x[q(x) - \bar{q}(x)].$$

$$\bullet F_2^{e p, \mu e}(x) = x \left[\frac{4}{9} (u(x) + \bar{u}(x)) + \frac{1}{9} (d(x) + \bar{d}(x)) \right]$$

$$F_2^{\nu d}(x) = \frac{5}{18} \frac{\nu(\bar{\nu})}{2 \text{ C.C.}}$$

$$\left(\frac{5}{18} = \text{average squared charge of } u, d \text{ quarks} \right).$$

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RELATIVISTIC KINEMATICS

I. BASICS

(a) Lorentz transformations -- Let E and \vec{p} be the energy and 3-momentum of a particle or system as seen from a certain inertial frame, and let E^* and \vec{p}^* be the same quantities as seen from a second inertial frame that moves with velocity $\vec{\beta}$ relative to the first. Then starred and unstarred quantities are related by

$$\begin{pmatrix} E^* \\ \vec{p}^* \end{pmatrix} = \begin{pmatrix} \gamma & -\gamma\vec{\beta} \\ -\gamma\vec{\beta} & \gamma \end{pmatrix} \begin{pmatrix} E \\ \vec{p} \end{pmatrix}, \quad p_1^* = p_1.$$

Here $\gamma = (1 - \beta^2)^{-1/2}$, and subscripts 1 and 1 indicate components of \vec{p} or \vec{p}^* that are parallel or perpendicular to $\vec{\beta}$ (often η is used for $\gamma\beta$). The inverse transformation is given by changing $\vec{\beta}$ to $-\vec{\beta}$. A particle of mass m at rest in the second frame, so that it is moving at velocity $\vec{\beta}$ relative to the first, has $E^* = m$ and $\vec{p}^* = 0$, so here

$$E = \gamma m, \quad \vec{p} = \gamma m \vec{\beta}.$$

In any frame, the energy, momentum, and mass are related by

$$E^2 = p^2 + m^2.$$

(b) Four momenta; scalar products -- The 4-momentum vector of a particle or system having energy E and 3-momentum \vec{p} is

$$q = (E, \vec{p}) = \{E, p_x, p_y, p_z\}.$$

Conservation of energy and the components of 3-momentum for any process $a + b + \dots \rightarrow 1 + 2 + \dots$ may then be written as

$$q_a + q_b + \dots = q_1 + q_2 + \dots$$

Although the components of a 4-momentum are different in different frames, the scalar product of any two 4-momenta q and q' , defined as

$$q \cdot q' = EE' - \vec{p} \cdot \vec{p}',$$

is an invariant; i.e., in numerical calculations the same result is obtained in any frame, and in algebraic calculations results obtained in different frames may be equated. For a particle of mass m , the scalar product $q \cdot q$ is

$$q \cdot q = q^2 = E^2 - p^2 = m^2.$$

The invariant mass M (or total c.m. energy) of an n -particle system is given by

$$M^2 = \left(\sum_{i=1}^n q_i \right)^2 = \left(\sum_i E_i \right)^2 - \left(\sum_i \vec{p}_i \right)^2,$$

where $q_i = (E_i, \vec{p}_i)$ is the 4-momentum of the i th particle.

(c) Electric and magnetic forces -- In Gaussian cgs units, the force on a particle with charge q moving with velocity \vec{v} in electric and magnetic fields \vec{E} and \vec{B} is

$$\vec{F} = q\vec{E} + q\vec{v} \times \vec{B},$$

where $\vec{\beta} = \vec{v}/c$. The units are \vec{F} in dynes, q in esu, \vec{E} in statvolts/cm, and \vec{B} in gauss. In mksa units, the force is

$$\vec{F} = q\vec{E} + q\vec{v} \times \vec{B},$$

where the units are \vec{F} in newtons (1 N = 10⁵ dynes), q in coulombs (1 C = 3 × 10⁹ esu; each 3 in this section is really 2.9979...), \vec{E} in volts/m (1 V = 1/300 statvolt), and \vec{B} in tesla (1 T = 10⁴ G). The force is zero if \vec{E} and \vec{B} are at right angles, $\vec{\beta}$ (or \vec{v}) is in the direction $\vec{E} \times \vec{B}$, and $\beta = E/B$ (cgs) or $v = E/B$ (mksa).

In a uniform, static magnetic field, the path of a charged particle is a helix of constant radius R and constant pitch angle λ , with the axis of the helix being along \vec{B} . The momentum is related to the other quantities by

$$p \cos \lambda = 3 \times 10^{-4} qBR,$$

where the units (very mixed!) are p in GeV/c, q in multiples of the electronic charge e , B in kG, and R in cm. The angular velocity about the axis of the helix is

$$\omega = 3 \times 10^{-4} qB/\gamma m,$$

where the units are ω in rad/sec, q in multiples of the electronic charge e , B in kG, and the energy γm in GeV.

II. DECAYS

(a) Survival probabilities -- Let a particle have mass m and proper mean life τ_0 . In a frame in which its 4-momentum is (E, \vec{p}) , the probability that it survives a time greater than t before decaying is

$$\text{Prob.}(> t) = \exp(-t/\gamma\tau_0) = \exp(-mt/E\tau_0).$$

The probability that it goes a distance greater than x before decaying is

$$\text{Prob.}(> x) = \exp(-x/\gamma\beta c\tau_0) = \exp(-mx/pct_0);$$

values of $c\tau_0$ (in cm) are given in the Stable Particle Table. If the particle has charge ze and is in a uniform magnetic field \vec{B} [see I(c)], then the probability that the projection of its helical path on the plane perpendicular to \vec{B} turns through an angle greater than θ before decaying is

$$\text{Prob.}(> \theta) = \exp(-Cm\theta/B\tau_0),$$

where, if m is in GeV, θ in deg, B in kG, and τ_0 in sec, then C is numerically 1.942×10^{-9} . This last distribution is independent of p or the helical pitch angle λ ; its only dependence is geometrical.

(b) Two-body decays -- A particle of mass m decays into two particles, masses m_1 and m_2 . In the rest frame of m , the energies of m_1 and m_2 are

$$e_1 = (m^2 + m_1^2 - m_2^2)/2m$$

$$e_2 = (m^2 + m_2^2 - m_1^2)/2m.$$

In this frame, the 3-momenta of m_1 and m_2 are equal and opposite and of magnitude

$$k = (e_1^2 - m_1^2)^{1/2} = (e_2^2 - m_2^2)^{1/2} \\ = \frac{1}{2} \{ [m^2 - (m_1 + m_2)^2] [m^2 - (m_1 - m_2)^2] \}^{1/2} / 2m.$$

See also the third paragraph of III(b).

(c) Three-body decays -- A particle of mass m decays into three particles, masses m_1 , m_2 , and m_3 . The invariant masses m_{ij} of the 2-particle systems, where $m_{ij}^2 = (q_i + q_j)^2$, satisfy the relation

$$m_{12}^2 + m_{13}^2 + m_{23}^2 = m^2 + m_1^2 + m_2^2 + m_3^2,$$

so that only two of the three m_{ij} 's are independent. In a rectangular Dalitz plot, m_{13}^2 (say) is plotted against m_{12}^2 . The kinematic boundaries may be calculated as follows: (i) The lower and upper limits on m_{12}^2 are $(m_1 + m_2)^2$ and $(m - m_3)^2$. (ii) For any m_{12}^2 between these limits, the lower and upper limits on m_{13}^2 are given by taking the + and - signs in

$$m_{13}^2 = (E_1 + E_3)^2 - (p_1 \pm p_3)^2,$$

where

$$E_1 = (m_{12}^2 + m_1^2 - m_3^2)/2m_{12}$$

$$E_3 = (m^2 - m_{12}^2 - m_3^2)/2m_{12}$$

$$p_1 = (E_1^2 - m_1^2)^{1/2}$$

$$p_3 = (E_3^2 - m_3^2)^{1/2}.$$

(These are the energies and momenta of particles 1 and 3 in the rest frame of m_{12} .) The phase-space density is uniform over the areas of both the above and the following form of the Dalitz plot.

In a triangular Dalitz plot, the kinetic energies T_1 , T_2 , and T_3 of the final-state particles in the rest frame of m are plotted as the distances inward from the sides of an equilateral triangle whose altitude is the energy Q released by the decay:

$$Q = T_1 + T_2 + T_3 = m - m_1 - m_2 - m_3.$$

The kinetic energies are related to the 2-particle invariant masses by

$$2mT_1 = (m - m_1)^2 - m_{23}^2 = (m_{23}^{\text{max}})^2 - m_{23}^2,$$

etc.

RELATIVISTIC KINEMATICS (Cont'd)

(d) Four-body decays -- A particle of mass m decays into four particles, masses $m_1, m_2, m_3,$ and m_4 . In a triangle (or Goldhaber) plot, the invariant mass of two of the particles is plotted against that of the other two, say m_{34} versus m_{12} , where $m_{ij}^2 = (q_i + q_j)^2$. The kinematic boundaries of this plot are the sides of the triangle whose vertices are at the points $(m_{12}, m_{34}) = (m_1 + m_2, m_3 + m_4), (m_1 + m_2, m - m_1 - m_2),$ and $(m - m_1 - m_4, m_3 + m_4)$. The phase-space density is not uniform over the enclosed area.

III. REACTIONS (MAINLY 2-BODY)

(a) Initial state -- Two particles, masses m_1 and m_2 , interact. In the lab frame, where particle 2 is at rest, the 4-momenta are (E_1, p_1) and $(m_2, 0)$. In the c.m. frame, where the 3-momenta are equal and opposite, the 4-momenta are (ϵ_1, \vec{k}) and $(\epsilon_2, -\vec{k})$. Then the total c.m. energy E is given by

$$E^2 = (\epsilon_1 + \epsilon_2)^2 = m_1^2 + m_2^2 + 2E_1m_2.$$

The c.m. energies of particles 1 and 2 are

$$\epsilon_1 = (m_1^2 + E_1m_2)/E = (E^2 + m_1^2 - m_2^2)/2E$$

$$\epsilon_2 = (m_2^2 + E_1m_2)/E = (E^2 + m_2^2 - m_1^2)/2E.$$

The c.m. momentum k is

$$k = p_1 m_2 / E.$$

See also the expression in II(b) for k , in which replace m with E .

The velocity of the c.m. relative to the lab is

$$\beta = p_1 / (E_1 + m_2).$$

The parameters for the Lorentz transformation between these frames [see II(a)] are

$$\gamma = (E_1 + m_2)/E$$

and

$$\gamma\beta = p_1/E.$$

(b) Two-body final states -- In the reaction $1 + 2 \rightarrow 3 + 4$, let the masses be m_i and the final-state c.m. 4-momenta be (ϵ_3, \vec{k}') and $(\epsilon_4, -\vec{k}')$. Then

$$\epsilon_3 = (E^2 + m_3^2 - m_4^2)/2E$$

$$\epsilon_4 = (E^2 + m_4^2 - m_3^2)/2E;$$

and

$$k' = (\epsilon_3^2 - m_3^2)^{1/2} = (\epsilon_4^2 - m_4^2)^{1/2}$$

$$= \{[E^2 - (m_3 + m_4)^2][E^2 - (m_3 - m_4)^2]\}^{1/2}/2E.$$

Let θ_3 be the lab production angle of particle 3 (the angle between p_3 and p_1), and let θ_3 be the c.m. production angle (the angle between \vec{k}' and \vec{k}). These angles are related by

$$\tan \theta_3 = \frac{p_{3\perp}}{p_{3\parallel}} = \frac{\sin \theta_3}{\gamma(\cos \theta_3 + \beta/\beta_3)},$$

where $p_{3\perp}$ and $p_{3\parallel}$ are the components of \vec{p}_3 perpendicular and parallel to p_1 , and $\beta_3 = k'/\epsilon_3$ is the c.m. velocity of particle 3. [See III(a) for γ and β .] If $\beta > \beta_3$, then particle 3 can only go forward in the lab, the maximum θ_3 being given by

$$\tan \theta_3^{\max} = \beta_3 \left(\frac{1 - \beta^2}{\beta^2 - \beta_3^2} \right)^{1/2}.$$

The components of \vec{p}_3 satisfy

$$\left(\frac{p_{3\parallel} - \gamma\beta\epsilon_3}{\gamma k'} \right)^2 + \left(\frac{p_{3\perp}}{k'} \right)^2 = 1,$$

which is the equation of an ellipse with semi-major axis $\gamma k'$ and semi-minor axis k' . Thus the possible lab momenta of particle 3 are the vectors to the ellipse from the point a distance $\gamma\beta\epsilon_3$ back along the major axis from the center of the ellipse.

The results of the preceding paragraph also apply to 2-body decay. Just set $m_2 = 0$, in which case $E = m_1$. [The decay-product masses are here m_1 and m_4 , not m_1 and m_2 as in II(b).]

The Mandelstam variables $s, t,$ and u are the Lorentz scalars defined in terms of the particle 4-momenta q_i as

$$s = (q_1 + q_2)^2 = (q_3 + q_4)^2$$

$$t = (q_1 - q_3)^2 = (q_2 - q_4)^2$$

$$u = (q_1 - q_4)^2 = (q_2 - q_3)^2.$$

They satisfy the relation

$$s + t + u = m_1^2 + m_2^2 + m_3^2 + m_4^2,$$

so that only two of the three are independent. Evaluating s in the c.m. frame gives

$$s = (\epsilon_1 + \epsilon_2)^2 = E^2,$$

and evaluating t and u , the 4-momentum-transfer-squared variables, in this frame gives

$$t = m_1^2 + m_3^2 - 2\epsilon_1\epsilon_3 + 2kk' \cos \theta_3$$

$$= t_0 - 4kk' \sin^2(\theta_3/2)$$

$$u = m_1^2 + m_4^2 - 2\epsilon_1\epsilon_4 + 2kk' \cos \theta_4$$

$$= u_0 - 4kk' \sin^2(\theta_4/2),$$

where θ_4 is the c.m. production angle of particle 4 ($\theta_3 + \theta_4 = \pi$), and

$$t_0 = t(\theta_3 = 0) = (\epsilon_1 - \epsilon_3)^2 - (k - k')^2$$

$$u_0 = u(\theta_4 = 0) = (\epsilon_1 - \epsilon_4)^2 - (k - k')^2.$$

The differences $\Delta t = t_0 - t_n$ and $\Delta u = u_0 - u_n$, where $t_n = t(\theta_3 = \pi)$ and $u_n = u(\theta_4 = \pi)$, are

$$\Delta t = \Delta u = 4kk'.$$

For elastic scattering, where $m_1 = m_3 = m$ and $m_2 = m_4 = M$, t_0 is zero and

$$t = -2k^2(1 - \cos \theta_3) = -4k^2 \sin^2(\theta_3/2).$$

And now

$$u_0 = (m^2 - M^2)^2/s.$$

Evaluating t in the lab frame gives

$$t = -2MT_4,$$

where $T_4 = E_4 - M$ is the lab kinetic energy of particle 4. For small-angle elastic scattering,

$$(-t)^{1/2} \approx k\theta_3 \approx p_1\theta_3 \approx p_4,$$

where $p_1, \theta_3,$ and p_4 are lab quantities.

IV. OTHER VARIABLES

(a) Rapidity -- For a system of energy E and momentum \vec{p} , the rapidity y is given by

$$y = \frac{1}{2} \ln \left(\frac{E + p_{\parallel}}{E - p_{\parallel}} \right) = \tanh^{-1} \left(\frac{p_{\parallel}}{E} \right) = \ln \left(\frac{E + p_{\parallel}}{m_1} \right),$$

where p_{\parallel} is the component of \vec{p} along a particular axis (the "rapidity axis", chosen, for example, parallel to the direction of an incoming beam), and $m_1 = (m^2 + p_{\perp}^2)^{1/2}$. Inverting these equations, we find

$$E = m_1 \cosh y$$

$$p_{\parallel} = m_1 \sinh y.$$

The shape of a rapidity distribution is invariant under a Lorentz transformation between inertial frames with relative motion parallel to the rapidity axis. Such a transformation is given by

$$y^* = y - \ln[\gamma(1 + \beta)] = y - \frac{1}{2} \ln \left(\frac{1 + \beta}{1 - \beta} \right),$$

where the sign of β is positive in the direction of increasing rapidity and p_{\parallel} .

RELATIVISTIC KINEMATICS (Cont'd)

(b) Scaling variable, hadron reactions -- In the inclusive reaction $h + 2 \rightarrow 3 + X$, with h any hadron, Feynman's x for particle 3 is defined as

$$x = k'_3/k'_{\max}$$

where k' is the c.m. momentum of particle 3. k'_{\max} is obtained [see Sec. III(b)] using the smallest mass m_X [called m_3 in III(b)] consistent with quantum conservation laws. At high energies, $k'_{\max} \approx \sqrt{s}/2$. Rapidity and x are related at large \sqrt{s} by

$$x \approx \frac{2m_3}{\sqrt{s}} \sinh y^*$$

where y^* is evaluated in the c.m.

(c) Scaling variables, lepton reactions -- For the inclusive reaction $l + 2 \rightarrow l' + X$, with particles l and l' leptons, we define the 4-velocity

$$q = (p_l - p_{l'})$$

so that

$$Q^2 \equiv -q^2 = 2E_l E_{l'} - 2|\vec{p}_l||\vec{p}_{l'}| \cos \theta - m_l^2 - m_{l'}^2 \geq 0$$

where θ is the $l \rightarrow l'$ scattering angle, and the preceding relation is valid in any frame. Also useful are

$$\nu = p_2 q / m_2 = [E_l - E_{l'}]_{\text{LAB}} = [E_X - m_2]_{\text{LAB}}$$

and

$$W = \sqrt{p_X^2} = (-Q^2 + 2m_2\nu + m_2^2)^{1/2} = m_X$$

Q^2 , ν , and W are Lorentz invariants, and the notation "LAB" refers to the reference frame with particle 2 at rest. (Note: ν is sometimes written $\nu = p_2 q$, leading to the replacement of $m_2\nu$ with ν throughout.)

Scaling variables in common use include

$$x \equiv \omega^{-1} = Q^2/2m_2\nu, 0 \leq x \leq 1$$

and

$$y = m_2\nu/p_l p_{l'} = [(E_l - E_{l'})/E_l]_{\text{LAB}}, 0 \leq y \leq 1$$

Both x and y are dimensionless.

Cross sections for inclusive reactions in the energy region where masses are negligible can be written in terms of E_l and certain pairs of these variables, usually Q^2 and ν , x and y , or Q^2 and x . If, in any frame, $|\vec{p}_l||\vec{p}_{l'}| \approx E_l E_{l'}$ and $E_l E_{l'} \sin^2(\theta/2) \gg m_l^2$ and $m_{l'}^2$ (i.e., $m_l, m_{l'}$ small), then

$$Q^2 \approx 4E_l E_{l'} \sin^2(\theta/2)$$

and

$$x \approx \frac{2E_l E_{l'} \sin^2(\theta/2)}{m_2\nu}$$

[†]Inequality sometimes violated unless $m_X \geq m_2$ and $m_l \geq m_{l'}$.

LORENTZ INVARIANT PHASE SPACE FORMULAE

For a system of n particles with overall four-momentum p and final four momenta p_1, \dots, p_n ($p_i = (E_i, \vec{p}_i)$), Lorentz Invariant Phase Space is given by

$$d \text{LIPS}(s; p_1, \dots, p_n) = (2\pi)^4 \delta^4(p - \sum_i p_i) \frac{1}{(2\pi)^{3n}} \prod_{i=1}^n \frac{d^3 \vec{p}_i}{2E_i} \quad (1)$$

$$\text{For 2-body: } d \text{LIPS}(s, p_1, p_2) = \frac{1}{(2\pi)^2} \delta^4(p - p_1 - p_2) d^4 p \frac{|\vec{p}^{\text{cm}}|}{4\sqrt{s}} d\Omega_1^{\text{cm}} \quad (2)$$

$$\text{For 3-body: } d \text{LIPS}(s, p_1, p_2, p_3) = \frac{1}{(2\pi)^5} \delta^4(p - p_1 - p_2 - p_3) d^4 p \frac{1}{32s} ds_{12} ds_{23} d\alpha d\beta d\gamma \quad (3)$$

where α, β , and γ are Euler angles.

For $a + b \rightarrow n$ particles or $X \rightarrow n$ particles, in general $i_f \rightarrow i_f'$,

$$\sigma_{if} = \frac{1}{4F} \int |\mathcal{M}_{if}|^2 d \text{LIPS}(s; p_1, \dots, p_n) \quad (4)$$

or

$$\Gamma_{if} = \frac{1}{2m_X} \int |\mathcal{M}_{if}|^2 d \text{LIPS}(m_X^2; p_1, \dots, p_n) \quad (5)$$

where \mathcal{M}_{if} is an invariant matrix element. F is Moller's invariant flux factor, $F^2 = (p_a \cdot p_b)^2 - m_a^2 m_b^2$. If a is beam, b , target ($\vec{p}_b^{\text{lab}} = 0$), then $F = |\vec{p}_a^{\text{lab}}| m_b = |\vec{p}_a^{\text{cm}}| \sqrt{s}$.

For elastic scattering in c.m., $\vec{p}_a^{\text{cm}} = -\vec{p}_b^{\text{cm}}$, and (2) and (4) yield

$$\frac{d\sigma}{d\Omega} = \frac{|\mathcal{M}|^2}{(8\pi)^2 s} \quad \text{or} \quad \frac{d\sigma}{dt} = \frac{|\mathcal{M}|^2}{64\pi |\vec{p}_a^{\text{cm}}|^2 s} \quad (6)$$

The normalization is such that the optical theorem reads

$$\text{Im } \mathcal{M}|_{t=0} = 2 |\vec{p}_a^{\text{cm}}| \sqrt{s} \sigma_{\text{tot}} \quad (7)$$

The choice of Eq. (4) implies a particular normalization of any spinors that may occur in \mathcal{M} . The advantage of this normalization is that it greatly simplifies the structure of \mathcal{M} by putting factors such as $\frac{1}{(2\pi)^3} \frac{1}{2E}$ into the phase space where they really belong. In addition, the labels, i, f , refer to specific spin (helicity) states, so that the usual "average and sum" rule is implicit.

C.M. ENERGY AND MOMENTUM VS. BEAM MOMENTUM

E_cm dE_cm = m_p dT_beam = m_p v_beam dP_beam ≈ m_p dP_beam

Table with 12 columns: PBEAM (GeV/c), C.M. ENERGY (GeV), MOMENTUM IN C.M. (GeV/c), PBEAM (GeV/c), C.M. ENERGY (GeV), MOMENTUM IN C.M. (GeV/c), PBEAM (GeV/c), C.M. ENERGY (GeV), MOMENTUM IN C.M. (GeV/c). Each column contains 10 sub-columns for different particle types (TP, VP, KP, PP).

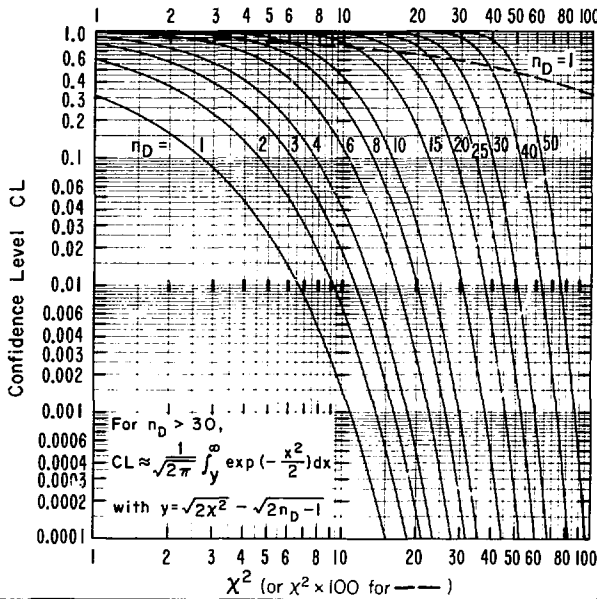
PROBABILITY AND STATISTICS

A. PROBABILITY DISTRIBUTIONS AND CONFIDENCE LEVELS

We give here properties of the three probability distributions most commonly used in high energy physics: Normal (or Gaussian), Chi-squared, and Poisson. We warn the reader that there is no universal convention for the term "confidence level"

as used by physicists; thus, explicit definitions are given for each distribution, and we have attempted to choose definitions that correspond to common usage. It is explained below how confidence levels for all three distributions can be extracted from the following figure.

χ^2 Confidence Level vs. χ^2 for Degrees of Freedom

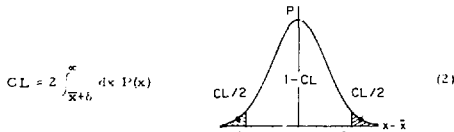


A.1. Normal Distribution

The normal distribution with mean \bar{x} and standard deviation σ (variance σ^2) is:

$$P(x)dx = \frac{1}{\sqrt{2\pi}\sigma} e^{-(x-\bar{x})^2/2\sigma^2} dx \quad (1)$$

The confidence level associated with an observed deviation from the mean, \hat{b} , is the probability that $|x - \bar{x}| > \hat{b}$, i.e.,



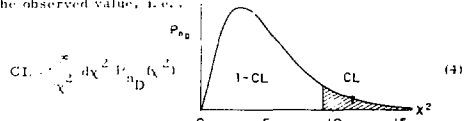
[The small figure in Eq. (2) is drawn with $\hat{b} = 2\sigma$.] CL is given by the ordinate of the $n_D = 1$ curve in the figure at $\chi^2 = (\hat{b}/\sigma)^2$. The confidence level for $\hat{b} = 1\sigma$ is 31.7%; 2σ , 4.6%; 3σ , 0.5%; 4σ , 99.7%. The odds against exceeding \hat{b} , $(1-CL)/CL$, for $\hat{b} = 1\sigma$ are 2.15:1; 2σ , 21:1; 3σ , 370:1; 4σ , 16,000:1; 5σ , 1,700,000:1. Relations between σ and other measures of the width: probable error (CL = 0.5 deviation) = 0.67σ ; mean absolute deviation = 0.80σ ; RMS deviation = σ ; half width at half maximum = 1.18σ .

A.2. Chi-squared Distribution

The chi-squared distribution for n_D degrees of freedom is:

$$P_{n_D}(\chi^2) d\chi^2 = \frac{1}{2^{n_D/2} \Gamma(n_D/2)} (\chi^2)^{n_D/2 - 1} e^{-\chi^2/2} d\chi^2 \quad (\chi^2 > 0) \quad (3)$$

where Γ ("half") $n_D/2$. The mean and variance are n_D and $2n_D$, respectively. In evaluating Eq. (3) one may use Stirling's approximation: $\Gamma(h) \approx (h-1)! \approx 2.507 e^{-h} h^{h-1/2} \times (1 + 0.0833/h)$ which is accurate to $\pm 0.1\%$ for all $h > 1/2$. The confidence level associated with a given value of n_D and an observed value of χ^2 is the probability of chi-squared exceeding the observed value, i.e.,



[The small figure in Eq. (4) is drawn with $n_D = 5$ and $CL = 10\%$.] CL is plotted as a function of χ^2 for several values of n_D in the above figure. For large n_D , χ^2 becomes normally distributed about n_D . Thus,

$$y_1 = (\chi^2 - n_D) / \sqrt{2n_D} \quad (5)$$

becomes normally distributed with unit standard deviation. A better approximation, due to Fisher,¹ is that y_2 , not χ^2 , becomes normally distributed, specifically

$$y_2 = \frac{\sqrt{2\chi^2} - \sqrt{2n_D - 1}}{\sqrt{2n_D - 1}} \quad (6)$$

approaches normality with unit standard deviation. For small CL's in particular, y_2 is much more accurate than y_1 . Thus, for $n_D = 50$ and $\chi^2 = 80$, the true CL is 0.45%, but y_1 is 3.0 corresponding to a CL of 0.13%, while y_2 is 2.7 corresponding to a CL of 0.35%.

PROBABILITY AND STATISTICS (Cont'd)

A.3. Poisson Distribution

The Poisson distribution with mean \bar{n} is:

$$P_n(n) = \frac{e^{-\bar{n}} (\bar{n})^n}{n!} \quad (n = 0, 1, 2, \dots) \quad (7)$$

The variance is equal to the mean. **Confidence levels** for Poisson distributions are usually defined in terms of quantities called "**upper limits**" as follows: The confidence level associated with a given upper limit N and an observed value n_0 of n is the probability that $n > n_0$ if $\bar{n} = N$, i. e.,

$$CL = \sum_{n=n_0+1}^{\infty} P_n(n) \quad (8)$$

$$= 1 - \sum_{n=0}^{n_0} P_n(n)$$

[The small figure in Eq. (8) is drawn with $n_0 = 2$ and $CL = 90\%$.] A useful relation between Poisson and chi-squared confidence levels allows one to look up this quantity on the above figure. Specifically, the quantity $1-CL$ is given by the ordinate of the $n_0 = 2(n_0+1)$ curve at $\chi^2 = 2N$. Thus, 90% confidence level upper limits for $n_0 = 0, 1, 2$ are given by half the χ^2 value corresponding to an ordinate of 0.1 on the $n_0 = 2, 4, 6$ curves, respectively; the values are $N = 2.3, 3.9, 5.3$.

Tables of confidence levels for all three of these distributions, the relation between Poisson and chi-squared confidence levels, and numerous other useful tables and relations may be found in Ref. 2.

B. STATISTICS

We consider here the situation in which one is presented with N independent data, $y_n \pm \sigma_n$, and it is desired to make some inference about the "true" value of the quantity represented by these data. For this purpose we interpret each datum y_n as a single sample point drawn randomly (and independently of the other data) from a distribution having mean \bar{y} (which we wish to estimate) and variance σ_y^2 . (Identification of the true σ_n with the σ_n datum is an approximation which may become seriously inaccurate when σ_n is an appreciable fraction of y_n .) Some methods of estimation commonly used in high energy physics are given below; see Ref. 3 for numerous applications. Section B.1. deals with the case in which all \bar{y}_n are the same, e. g., several different measurements of the same quantity; Sec. B.2. deals with the case in which $\bar{y}_n = \bar{y}(x_n)$, where x_n represents some set of independent variables, e. g., cross-section measurements at various values of energy and angle, $x_n = (E_n, \theta_n)$.

B.1. Single Mean and Variance Estimates

(1) If the y_n represent a set of values all supposedly drawn from a single distribution with mean \bar{y} and variance σ^2 (i. e., the σ_n are all the same, but the common value is unknown) then

$$\bar{y}_e = \frac{1}{N} \sum_{n=1}^N y_n \quad (9)$$

$$\sigma_e^2 = \frac{1}{N-1} \sum_{n=1}^N (y_n - \bar{y}_e)^2 = \frac{1}{N-1} \left[\left(\sum_{n=1}^N y_n^2 \right) - N \bar{y}_e^2 \right] \quad (10)$$

are unbiased estimates of \bar{y} and σ^2 . The variance of \bar{y}_e is σ^2/N . If the parent distribution is normal and N is large, the variance of σ_e^2 is $2\sigma^4/N$.

(2) If the \bar{y}_n all have the common value \bar{y} and the σ_n are known, then the weighted average

$$\bar{y}_w = \frac{1}{w} \sum_{n=1}^N w_n y_n \quad (11)$$

where $w = 1/\sigma_n^2$ and $w = \sum w_n$, is an appropriate unbiased estimate of \bar{y} . This choice of weighting factors in Eq. (11) minimizes the variance of the estimate; the variance is $1/w$.

B.2. Linear Least Squares Fit

A least squares fit of the function $y(x) = \sum_{j=1}^m a_j f_j(x)$ to independent data $y_n \pm \sigma_n$ at points x_n (e. g., a Legendre fit in which the f_j are Legendre polynomials and the a_j are Legendre coefficients) gives the following estimates of the parameters a_j :

$$\hat{a}_{e,i} = \sum_{j=1}^m V_{ij} f_j(x_n) y_n / \sigma_n^2 \quad (12)$$

Here V is the covariance matrix of the fitted parameters

$$V_{ij} = \frac{1}{\sigma_n^2} \left(\sum_{n=1}^N f_i(x_n) f_j(x_n) / \sigma_n^2 - \bar{f}_i \bar{f}_j \right) \quad (13)$$

which is given by

$$(V^{-1})_{ij} = \sum_{n=1}^N f_i(x_n) f_j(x_n) / \sigma_n^2 \quad (14)$$

The variance of an interpolated or extrapolated value of y at point x , $y_e = \sum a_{e,i} f_i(x)$, is:

$$\overline{(y_e - \bar{y}_e)^2} = \sum_{i,j=1}^m V_{ij} f_i(x) f_j(x) \quad (15)$$

For the case of a straight line fit, $y(x) = a + bx$, one obtains the following estimates of a and b ,

$$a_e = (S_y S_{xx} - S_x S_{xy}) / D, \quad (16)$$

$$b_e = (S_1 S_{xy} - S_x S_y) / D,$$

where

$$S_1, S_x, S_y, S_{xx}, S_{xy} = \sum (1, x_n, y_n, x_n^2, x_n y_n) / \sigma_n^2 \quad (17)$$

$$D = S_1 S_{xx} - S_x^2$$

The covariance matrix of the fitted parameters is:

$$\begin{pmatrix} V_{aa} & V_{ab} \\ V_{ab} & V_{bb} \end{pmatrix} = \frac{1}{D} \begin{pmatrix} S_{xx} & -S_x \\ -S_x & S_1 \end{pmatrix} \quad (18)$$

The variance of an interpolated or extrapolated value of y at point x is:

$$\overline{(y_e - \bar{y}_e)^2} = \frac{1}{S_1} \left(\frac{S_1}{D} x - \frac{S_x}{S_1} \right)^2 \quad (19)$$

C. ERROR PROPAGATION

We consider here the situation in which one wishes to calculate the value and error of a function of some other quantities with errors, e. g., in a Monte Carlo program. Let $\{y\}$ be a set of random variables with means $\{\bar{y}\}$ and covariance matrix V . Then the mean and variance of a function of these variables are approximately (to second order in $\{y-\bar{y}\}$):

$$\bar{f} \approx f(\{\bar{y}\}) + \frac{1}{2} \sum_{m,n} V_{mn} \frac{\partial^2 f}{\partial y_m \partial y_n} / \{y\} = \{y\} \quad (20)$$

$$\overline{(f - \bar{f})^2} = \sum_{m,n} V_{mn} \left(\frac{\partial f}{\partial y_m} / \{y\} - \frac{\partial f}{\partial y_n} / \{y\} \right) \{y\} = \{y\} \quad (21)$$

E. g., the mean and variance of a function of a single variable with mean \bar{y} and variance σ^2 are:

$$\bar{f} = f(\bar{y}) + \frac{1}{2} \sigma^2 f''(\bar{y}) \quad (22)$$

$$\overline{(f - \bar{f})^2} = \sigma^2 f'(\bar{y})^2 \quad (23)$$

Note that these equations will usually be applied by substituting some measured quantities, $\{y\}$ say, for the true means, $\{\bar{y}\}$. If, as is often the case, $\bar{y}_n - \bar{y}$ is of order $\sqrt{V_{nn}}$, then there is no point in keeping the second order terms in Eq. (20) or (22) since the substitution itself introduces first order errors.

1. R. A. Fisher, Statistical Methods for Research Workers (Oliver and Boyd, Edinburgh and London, 1958).
2. M. Abramowitz and I. Stegun, eds., Handbook of Mathematical Functions (National Bureau of Standards, Applied Mathematics Series, Vol. 55, Washington, 1964).
3. W. T. Eadie, D. Drijard, F. E. James, M. Roos, and B. Sadoulet, Statistical Methods in Experimental Physics (North-Holland, Amsterdam and London, 1971).

PARTICLE DETECTORS, ABSORBERS, AND RANGES*

A. DETECTOR PARAMETERS

In this section we give various parameters for common detectors. The quoted numbers represent at best an order of magnitude, and are useful only for preliminary design. A more detailed introduction to detectors can be found in "A Consumer's Guide to Particle Detectors," by D.J. Miller, Rutherford Lab Report RL-76-072, July 1976.

A.1 Scintillators: Photon yield $\approx 1\gamma/100$ eV in plastic scintillator¹ and $\approx 1\gamma/25$ eV in NaI.^{1,2}

A.2 Cerenkov:³ Half-angle θ_c of cone aperture in terms of velocity β and index of refraction n :

$$\theta_c = \arccos \left(\frac{1}{\beta n} \right) \sim \left[2 \left(1 - \frac{1}{\beta n} \right) \right]^{1/2}$$

Threshold velocity: $\beta_t = 1/n$; $\gamma_t = 1/\sqrt{1-\beta_t^2}$.

Therefore, $\beta_t \gamma_t = 1/\sqrt{2\delta + \delta^2}$, where $\delta = n-1$. Values of δ for various commonly used gases are given as a function of pressure and wavelength in Ref. 4; for values at atmospheric pressure, see the Table of Atomic and Nuclear Properties, following.

Number of photons, N , per cm:

$$N = \frac{\alpha}{c} \int \left(1 - \frac{1}{\beta^2 n^2} \right) 2\pi d\omega = \frac{\alpha}{c} \beta_t^2 \int \left(\frac{1}{\beta_t^2 \gamma_t^2} - \frac{1}{\beta^2 \gamma^2} \right) 2\pi d\omega$$

$$\approx 500 \sin^2 \theta_c / \text{cm (visible spectrum)}$$

A.3 Photon Collection: In addition to the photon yield, one should take into account the light collection efficiency ($\leq 10\%$ for typical 1-cm-thick scintillator), attenuation length (≈ 1 to 4 m for typical scintillators⁵), and quantum efficiency of the photomultiplier cathode ($\leq 25\%$).

A.4 Typical Detector Characteristics:

Detector Type	Accuracy (rms)	Resolution Time	Dead Time
Bubble chamber	$\approx \pm 10$ to $\approx \pm 150\mu$	≈ 1 ms	$\approx 1/20$ s ^a
Streamer chamber	$\pm 300\mu$	≈ 2 μ s	≈ 100 ms
Optical spark chamber	$\pm 200\mu$ ^b	≈ 2 μ s	≈ 10 ms
Magnetostrictive spark chamber	$\pm 500\mu$	≈ 2 μ s	≈ 10 ms
Proportional chamber	$\geq \pm 300\mu$ ^{c,d}	≈ 50 ns	≈ 200 ns
Drift chamber	± 50 to 300μ	≈ 2 ns ^e	≈ 100 ns
Scintillator	--	≈ 150 ps	≈ 10 ns
Emulsion	$\pm 1\mu$	--	--

^a Multiple pulsing time.

^b 60μ for high pressure.

^c 300μ is for 1 mm pitch.

^d Delay line cathode readout can give $\pm 150\mu$ parallel to anode wire.

^e For two chambers.

A.5 Shower Detectors: Typical energy resolutions (FWHM) for incident electron in the 1 GeV range, E in GeV. For a fixed number of radiation lengths, FWHM in the last three detectors would be expected to be proportional to \sqrt{t} for t (= plate thickness) ≥ 0.2 radiation lengths.⁶

NaI (20 rad. lengths):⁷ $\frac{2\%}{E^{1/4}}$

Lead Glass (14 rad. lengths):⁸ $\frac{10 - 12\%}{\sqrt{E}}$

Lead-Liquid Argon (15.75 rad. lengths):⁶ $\frac{16\%}{\sqrt{E}}$
(42 cells: 1.1 mm lead, 2 mm liquid argon, 2.3 mm lead-G10, 2 mm liquid argon)

Lead-Scintillator Sandwich (12.5 rad. lengths):⁹ $\frac{17\%}{\sqrt{E}}$
(66 cells: 1 mm lead, 5 mm scintillator)

Proportional Wire Shower Chamber (17 rad. lengths):¹⁰ $\frac{40\%}{\sqrt{E}}$
(36 cells: 0.474 rad. length type-metal + Al, 9.5 mm 80% Ar - 20% CH₄ gas)

A.6. Proportional Chamber Wire Instability: The limit on the voltage V for a wire tension T , due to mechanical effects when the electrostatic repulsion of adjacent wires exceeds the restoring force of wire tension, is given by¹¹

$$V \leq \frac{sT^{1/2}}{\epsilon C}$$

where s , ℓ , and C are the wire spacing, length, and capacitance per unit length. An approximation to C for chamber half-gap t and wire diameter d (good for $s \leq t$) gives¹²

$$V \leq 59T^{1/2} \left[\frac{t}{\ell} + \frac{s}{\pi \ell} \epsilon n \left(\frac{s}{\pi d} \right) \right],$$

where V is in kV, and T is in grams.

A.7 Proportional and Drift Chamber Potentials: Potential distributions and fields for an array of parallel line charges q (coul./m) along z and located at $y = 0, x = 0, \pm a, \pm 2a, \dots$, can usually be calculated with good accuracy from (MKS-A):

$$V(x,y) = -\frac{q}{4\pi\epsilon_0} \epsilon n \left\{ 4 \left[\sin^2 \left(\frac{\pi x}{a} \right) + \sinh^2 \left(\frac{\pi y}{a} \right) \right] \right\}$$

B. COSMIC RAY FLUXES

The fluxes of particles of different types depend on the latitude, their energy, and the conditions of measurement. Some typical sea-level values¹³ for charged particles are given below:

I_V	flux per unit solid angle above vertical direction crossing unit horizontal area
J_1	perpendicular component of total flux crossing unit horizontal area from above
J_2	total flux crossing unit horizontal area

	Total Intensity	Hard Component	Soft Component	
I_V	1.1×10^{-2}	0.8×10^{-2}	0.3×10^{-2}	$\text{cm}^{-2} \text{sec}^{-1} \text{sterad}^{-1}$
J_1	1.8×10^{-2}	1.3×10^{-2}	0.5×10^{-2}	$\text{cm}^{-2} \text{sec}^{-1}$
J_2	2.4×10^{-2}	1.7×10^{-2}	0.7×10^{-2}	$\text{cm}^{-2} \text{sec}^{-1}$

Very approximately, about 75% of all particles at sea-level are penetrating, and are muons. The absolute flux of protons at sea-level, in a momentum range 700-1100 MeV/c, is $1.5 \times 10^{-5} \text{cm}^{-2} \text{sec}^{-1} \text{sterad}^{-1}$, or $\sim 0.1\%$ of all particles.

The muon flux at sea-level has a mean energy of 2 GeV and a differential spectrum falling as E^{-2} , steepening smoothly to $E^{-3.6}$ above a few TeV. The angular distribution is $\cos^2 \theta$, changing to $\sec \theta$ at energies above a TeV, where θ is the zenith angle at production. The $+/-$ charge ratio is 1.25-1.30. The mean energy of muons originating in the atmosphere is roughly 300 GeV at slant depths \geq a few hundred meters. Beyond slant depths of ~ 10 km water-equivalent, the muons are due primarily to in-the-earth neutrino interactions (roughly $1/8$ interaction $\text{ton}^{-1} \text{year}^{-1}$ for $E_\nu > 300$ MeV, \sim constant throughout the earth).¹⁴ Muons from this source arrive with a mean energy of 20 GeV, and have a flux of $2 \times 10^{13} \text{cm}^{-2} \text{sec}^{-1} \text{sterad}^{-1}$ in the vertical direction and about twice that in the horizontal,¹⁵ down at least as far as the deepest mines.

C. PASSAGE OF PARTICLES THROUGH MATTER

C.1 Energy Loss Rates for Heavy Charged Projectiles: A heavy projectile (much more massive than an electron) of charge $Z_{\text{inc}} e$, incident at speed βc ($\beta \gg 1/137$) through a slowing medium, dissipates energy principally via interactions with the electrons of the medium. The mean rate of such energy loss per unit path length x may be written as:¹⁶

$$\left(\frac{dE}{dx} \right)_{\text{inc}} = \frac{D Z_{\text{med}} \rho_{\text{med}}}{A_{\text{med}}} \left(\frac{Z_{\text{inc}}}{\beta} \right)^2$$

$$\times \left[\epsilon n \left(\frac{2m_e \gamma^2 \beta^2 c^2}{I} \right) - \beta^2 - \frac{\delta}{2} - \frac{c}{Z_{\text{med}}} \right] \left\{ 1 + \nu \right\},$$

PARTICLE DETECTORS, ABSORBERS, AND RANGES (Cont'd)

where $D = 4\pi N_A r_e^2 m_e c^2 = 0.3070 \text{ MeV cm}^2/\text{g}$ (see Physical and Numerical Constants Table).

Here Z_{med} and A_{med} are the charge and mass numbers of the medium and ρ_{med} is the mass density of the medium; l , δ , C , and ν are phenomenological functions. Frequently, the values of δ , C , and ν are negligibly small; the parameter l characterizes the binding of the electrons of the medium. As a rule of thumb, we may estimate l for an idealized medium as $l \approx 16 (Z_{\text{med}})^{0.9} \text{ eV}$ when $Z_{\text{med}} > 1$. For realistic media the value of l will vary at the 10% level from this estimate; for H_2 , $l = 20.0 \text{ eV}$. We may approximately treat media which are chemical mixtures or compounds by computing

$$\frac{dE}{dx} \approx \sum_{n=1}^N \left(\frac{dE}{dx} \right)_n$$

with $(dE/dx)_n$ appropriate to the n^{th} chemical constituent (using $\rho_{\text{med}}^{(n)}$ as the partial density in the formula for dE/dx).¹⁷

The function δ represents the density effect upon the energy loss rate; it is non-negligible only for highly relativistic projectiles in denser media.¹⁸ For ultra-relativistic projectiles, δ approaches $2\epsilon n\gamma + \text{constant}$, where the value of the constant depends upon the density of the medium as well as its chemical composition.

The function C represents shell corrections to the energy loss rate.¹⁶ These effects are non-negligible only for projectiles with speeds not much faster than the speeds of the fastest electrons bound in the medium.

The function ν represents corrections due to higher-order electrodynamic effects.¹⁹ These effects become important when $|Z_{\text{inc}}/\beta|$ is comparable to 137. For relativistic unit-charge projectiles, $|\nu|$ is of the order of 1%; positively charged projectiles lose energy more rapidly than do their charge conjugates.^{19,20}

For non-relativistic projectiles, our formulae above are inapplicable. At the very slowest speeds, total energy loss rates are believed to be proportional to β , rising through a peak at projectile speeds comparable to atomic speeds, after having passed through a smaller peak (due to elastic Coulomb collisions with the nuclei of the slowing medium²¹) at intermediate speeds. In some cases, energy loss rates depend significantly upon the relation of the projectile trajectory to the crystalline structure of the slowing medium.²²

For relativistic projectiles, $(dE/dx)_{\text{inc}}$ falls rapidly with increasing β until reaching a minimum around $\beta = 0.96$ (almost independent of medium), followed by a slow rise. Because of the density effect, the quantity in square brackets approaches $\epsilon n\gamma + \text{constant}$ for large γ .

The quantity $(dE/dx)_{\text{inc}} \delta x$ is the mean total energy loss via interactions with electrons of the medium in a layer of thickness δx . For any finite δx , Poisson fluctuations can cause the actual energy loss to deviate from the mean. For thin layers, the distribution is broad and skewed, being peaked below $(dE/dx)\delta x$, and having a long tail toward large energy losses.²³ Only for a very thick layer $[(dE/dx)\delta x \gg 2m_e\beta^2\gamma^2c^2]$ will the distribution of energy losses become nearly Gaussian. The large fluctuations of the total energy loss rate from the mean are due to a small number of collisions involving large energy transfers. The fluctuations are greatly reduced for the so-called restricted energy loss rate, described in Section C.4.

C.2 Ionization Yields: Physicists frequently relate total energy loss to the number of ion pairs produced in the stopping medium. This relation becomes complicated for relativistic projectiles due to the wandering of energetic knock-on electrons whose ranges exceed the dimensions of the fiducial volume. For a qualitative appraisal of the non-locality of energy deposition by such modestly energetic knock-on electrons in various media, see Ref. 24. The mean energy loss per ion pair produced, W , is essentially constant for relativistic projectiles, but increases at slow projectile speeds.²⁵ The numerical value of W for gases can be surprisingly sensitive to trace amounts of various contaminants.²⁵ Of course, in addition to the preceding effects, practical ionization yields may be greatly influenced by subsequent recombinations, etc.²⁶

C.3 Energetic Knock-On Electrons: For a relativistic spinless point-charge projectile, the production of high energy (kinetic energy $T \gg I$) electrons is given by (neglecting the spin of the electron):

$$\frac{d^2N}{dTdx} = \frac{1}{2} D \left(\frac{Z_{\text{med}}}{A_{\text{med}}} \right) \left(\frac{Z_{\text{inc}}}{\beta} \right)^2 \rho_{\text{med}} \frac{1}{T^2}$$

for $l \ll T \leq T_{\text{max}}$, where

$$T_{\text{max}} = \frac{2m_e\beta^2\gamma^2c^2}{1 + 2\gamma \frac{m_e}{M_{\text{inc}}} + \left(\frac{m_e}{M_{\text{inc}}} \right)^2}$$

M_{inc} is the mass of the incident projectile, and all other quantities are as in Section C.1. This formula does not differ significantly from the precise result, incorporating spin effects, for any projectile (including e^{\pm}) in the restricted range $l \ll T \ll T_{\text{max}}$; more accurate formulae are available for various projectiles.^{27,28} Our formula is inaccurate for T close to l ; for $2l \leq T \leq 10l$, the $1/T^2$ dependence above becomes $\approx T^{-\eta}$ with $3 \leq \eta \leq 5$.²⁹

C.4 Rates of Restricted Energy Loss for Relativistic Charged Projectiles: The variability of energy loss for heavy projectiles is due primarily to the variability in the production of energetic knock-on electrons.

Bremsstrahlung and pair-production processes make this variability even greater for electrons than for heavy particles as projectiles (see, e.g., the figure "Fractional Energy Loss for e^+ and e^- in Lead," following). If an instrument, such as a bubble chamber, is capable of isolating these high-energy-loss interactions, then it is appropriate to consider the rate of energy loss excluding them, i.e., a restricted energy loss rate. The mean energy loss rate via all collisions which have energy transfer T such that $T \leq E_{\text{max}} \ll T_{\text{max}}$ is:¹⁶

$$\left(\frac{dE}{dx} \right)_{\leq E_{\text{max}}} = \frac{1}{2} D \frac{Z_{\text{med}}^2 \rho_{\text{med}}}{A_{\text{med}}} \left(\frac{Z_{\text{inc}}}{\beta} \right)^2 \times \left[\epsilon n \left(\frac{E_{\text{max}} T_{\text{max}}}{l^2} \right) - \beta^2 - \delta - \frac{2C}{Z_{\text{med}}} \right]$$

Notice the overall factor of $1/2$.

The density effect causes the restricted energy loss rate to approach a constant, the Fermi plateau value, for the fastest projectiles.

C.5 Multiple Scattering through Small Angles: As a charged particle traverses a medium it is deflected by many small-angle elastic scatterings. The bulk of this deflection is due to elastic Coulomb scattering from the nuclei within the medium, hence the usual identification as multiple Coulomb scattering (note, however, that strong interactions do contribute to the total multiple scattering for hadronic projectiles). For both Coulomb and strong interactions, the Central Limit Theorem provides little useful guidance in establishing the precise nature of the distribution of the total deflections resulting from multiple scattering. The true distribution is roughly Gaussian only for small deflection angles while it shows much greater probability for large-angle scatterings \geq a few θ_0 , see below, depending on absorber) than the Gaussian would suggest. These tails on the distribution (a few % of peak height in the region where the Gaussian part becomes negligible) are more pronounced for hadrons than for muons as projectiles. The large-angle behavior of these distributions is best estimated by computing the exact distribution for the vectorial sum of the largest deflections based upon the true elastic scattering cross section of the projectile against the medium,³⁰ or, when applicable, by interpolation from tabular data.³¹ An easier alternative which may suffice for non-critical applications would be to use a Gaussian approximation with the following width:³²

$$\theta_0 = \frac{14.1 \text{ MeV}/c}{\beta\bar{\rho}} Z_{\text{inc}} \sqrt{L/L_R} \left[1 + \frac{1}{9} \log_{10} \left(L/L_R \right) \right] \text{ (radians)}$$

PARTICLE DETECTORS, ABSORBERS, AND RANGES (Cont'd)

where p , β , and Z_{inc} are the momentum (in MeV/c), velocity, and charge number of the incident particle, and L/L_R is the thickness, in radiation lengths, of the scattering medium. L_R for certain materials is given in the Table of Atomic and Nuclear Properties of Materials (following). The angle, θ_0 , is a fit to Moliere³⁰ theory, accurate to about 5% for $10^{-3} < L/L_R < 10$ except for very light elements or low velocity where the error is about 10 to 20%. In this Gaussian approximation, θ_0 has the meaning

$$\theta_0 = \theta_{\text{plane}}^{\text{rms}} = \frac{1}{\sqrt{2}} \theta_{\text{space}}^{\text{rms}}$$

The non-projected (space) and projected (plane) angular distributions are given approximately³⁰ by the Gaussian forms:

$$\frac{1}{2\pi\theta_0^2} \exp\left[-\frac{\theta_{\text{space}}^2}{2\theta_0^2}\right] d\Omega,$$

$$\frac{1}{\sqrt{2\pi}\theta_0} \exp\left[-\frac{\theta_{\text{plane}}^2}{2\theta_0^2}\right] d\theta_{\text{plane}}$$

where θ is the deflection angle.

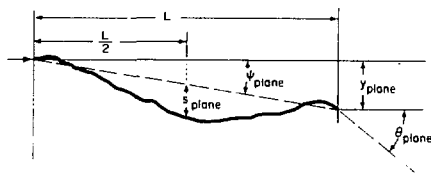
Other quantities are sometimes used to describe the amount of multiple Coulomb scattering: the auxiliary quantities $\psi_{\text{plane}}^{\text{rms}}$, $y_{\text{plane}}^{\text{rms}}$, and $s_{\text{plane}}^{\text{rms}}$ (see the figure) obey:

$$y_{\text{plane}}^{\text{rms}} = \frac{1}{\sqrt{3}} L \theta_{\text{plane}}^{\text{rms}} = \frac{1}{\sqrt{3}} L \theta_0,$$

$$\psi_{\text{plane}}^{\text{rms}} = \frac{1}{\sqrt{3}} \theta_{\text{plane}}^{\text{rms}} = \frac{1}{\sqrt{3}} \theta_0,$$

and

$$s_{\text{plane}}^{\text{rms}} = \frac{1}{4\sqrt{3}} L \theta_{\text{plane}}^{\text{rms}} = \frac{1}{4\sqrt{3}} L \theta_0.$$



All the quantitative estimates in this section apply only in the limit of small $\theta_{\text{plane}}^{\text{rms}}$ and in the absence of large-angle scatters.

C.6 Longitudinal Distribution of Electromagnetic Showers: A photon of energy E (GeV) ≥ 0.1 GeV converting in a semi-infinite medium produces an electromagnetic cascade whose intensity initially increases with depth and then falls off. The average number of e^\pm with kinetic energy above 1.5 MeV, crossing a plane at a depth t of radiation lengths from the beginning of the medium, in a material of atomic number Z , calculated using the Monte Carlo program EGS,³³ can be fit by the empirical formula³⁴

$$N = N_0 t^a e^{-bt},$$

where $N_0 = 5.51 E \sqrt{Z} b^{a+1} / \Gamma(a+1)$ and $b = 0.634 - 0.0021 Z$. For $Z \geq 26$, $a = 2.0 - Z/340 + (0.664 - Z/340) \ln E$. For $Z = 13$, $a = 1.77 - 0.52 \ln E$. The maximum intensity, N_{max} , occurs at the depth $t = a/b$. The maximum error of the fit occurs in the vicinity of this depth

and is less than $0.15 N_{\text{max}}$. The integral of the tail, $\int_{1.5 a/b}^{\infty} N dt$ is fit to

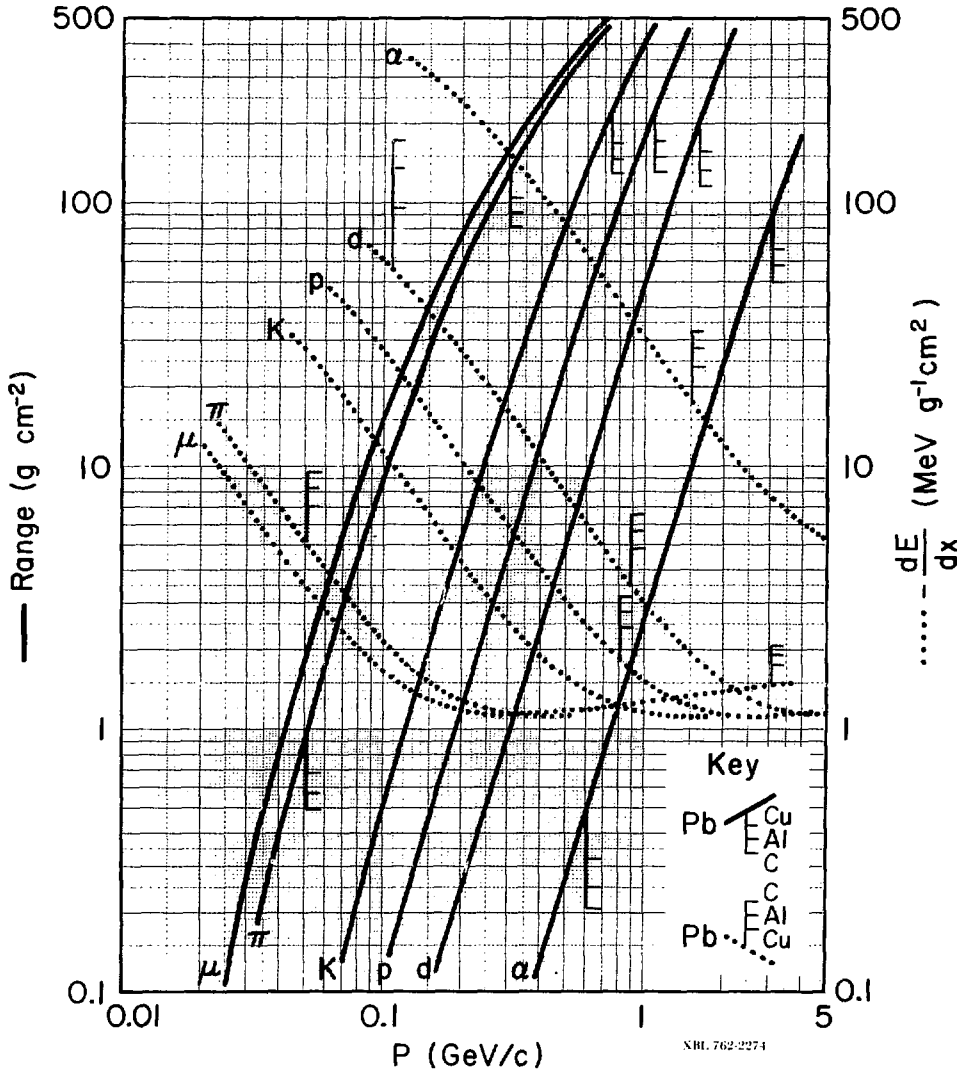
better than 2.5%. The total longitudinally-projected e^\pm path length, $\int_0^{\infty} N dt = 5.51 E \sqrt{Z}$, is less than the total e^\pm path length due primarily to multiple Coulomb scattering.

³Prepared April 1974 by Sherwood Parker and Bernard Sadoulet. Revised April 1982 by Sherwood Parker, Ray Hagstrom, Didier Besset, and John Learned.

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PARTICLE DETECTORS, ABSORBERS, AND RANGES (Cont'd)

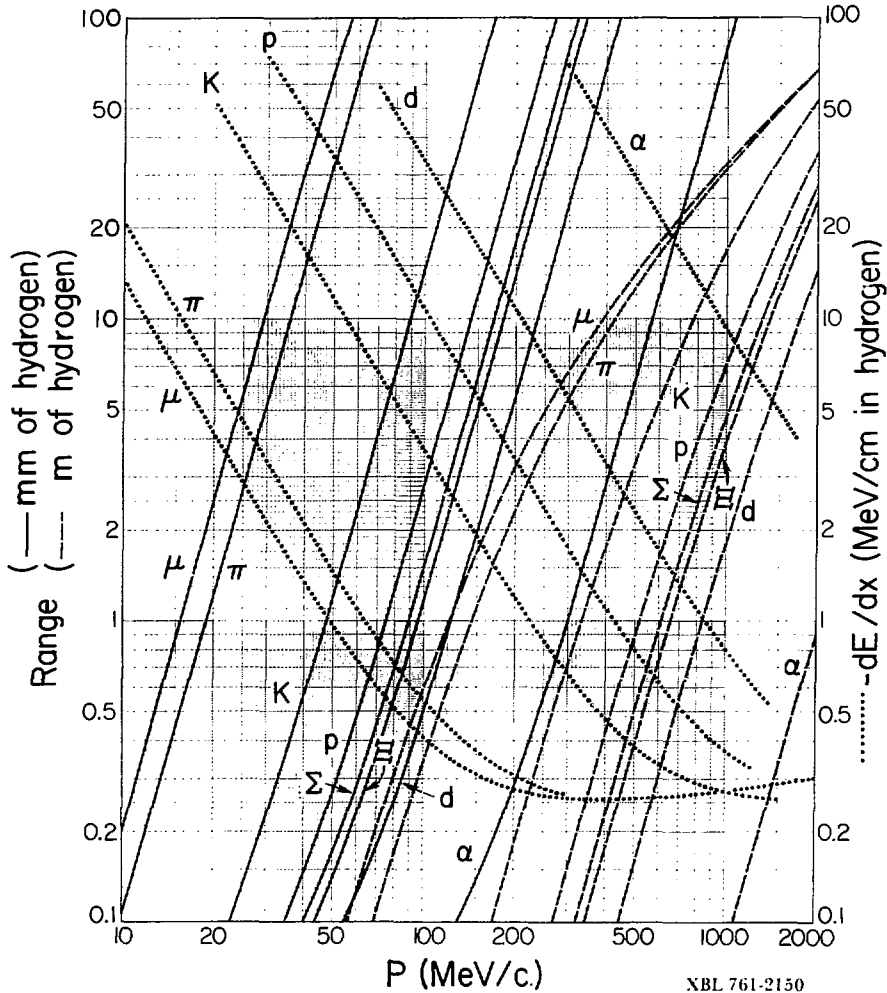
Mean Range and Energy Loss in Lead, Copper, Aluminum, and Carbon



Mean range and energy loss due to ionization for the indicated particles in Pb, with scaling to Cu, Al, and C indicated, using Bethe-Bloch equation (Section C.1 above) with corrections. Calculated using program of Hans Bichsel (UCRL-17538), with density correction added (Hans Bichsel, private communication). See also Joseph F. Janni [Air Force Weapons Laboratory Technical Report No. AFWL-TR-65-150 (1966)]. The average ionization potentials (I) assumed were: Pb (820 eV), Cu (320 eV), Al (166 eV), and C (77.5 eV). Figure indicates total path length; observed range may be smaller (by $\sim 1\% - 2\%$ in heavy elements) due to multiple scattering, primarily from small energy-loss collisions with nuclei. The functional forms have not been experimentally verified to better than roughly $\pm 1\%$. For higher energies refer to discussion by Cobb [“A Study of Some Electromagnetic Interactions of High Velocity Particles with Matter,” University of Oxford Report HEP/T/55 (1973)] and by Turner [“Penetration of Charged Particles in Matter: A Symposium”, National Academy of Sciences, Washington D. C. (1970), p. 48]. Scaling to other beam particles is, to a good approximation, described by the expression on the next page.

PARTICLE DETECTORS, ABSORBERS, AND RANGES (Cont'd)

Mean Range and Energy Loss in Liquid Hydrogen



Range and energy loss in liquid hydrogen bubble chamber, based on Bethe-Bloch equation (Section C.1 above), using an average ionization potential for H_2 of $I = 20.0$ eV, which is an approximate average of the experimental result of Garbincius and Hyman [Phys. Rev. A2, 1834 (1970)] and the theoretical result of Ford and Browne [Phys. Rev. A7, 418 (1973)]. Bubble chamber conditions are chosen to be those of Garbincius and Hyman: parahydrogen of density = 0.0625 g/cm³ (note: range \propto 1/density), with vapor-pressure 60.8 lb/in² (absolute) and temperature $26.2^\circ K$. The functional dependence of the Bethe-Bloch equation is not experimentally verified to better than about ± 14 over large momentum ranges. It should be noted that the number of bubbles per cm of a track in a bubble chamber is nearly proportional to $1/\delta^2$, not dE/dx . For the linear portions of the range curves, $R \propto p^{2.6}$. **Scaling law for particles of other mass or charge (except electrons):** for a given medium, the range R_b of any beam particle with mass M_b , charge z_b , and momentum p_b is given in terms of the range R_a of any other particle with mass M_a , charge z_a , and momentum $p_a = p_b M_a / M_b$ (i.e., having the same velocity) by the expression:

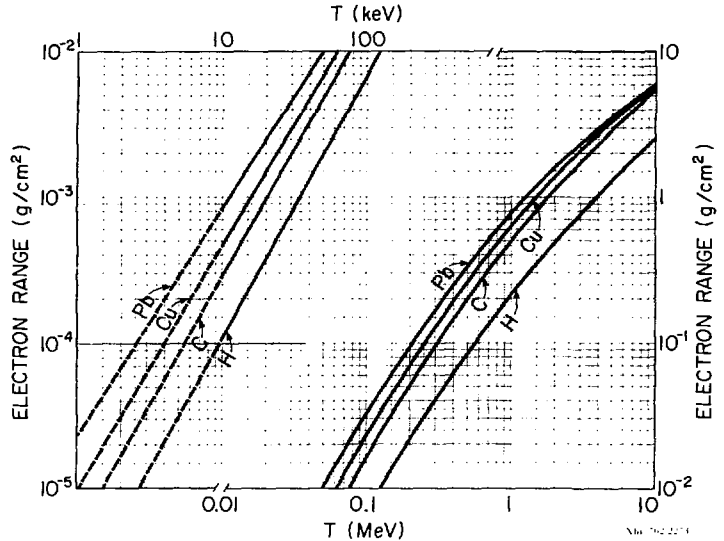
$$R_b(M_b, z_b, p_b) = \left[\frac{M_b/M_a}{z_b^2/z_a^2} \right] R_a(M_a, z_a, p_a = p_b M_a / M_b)$$

PARTICLE DETECTORS, ABSORBERS, AND RANGES (Cont'd)

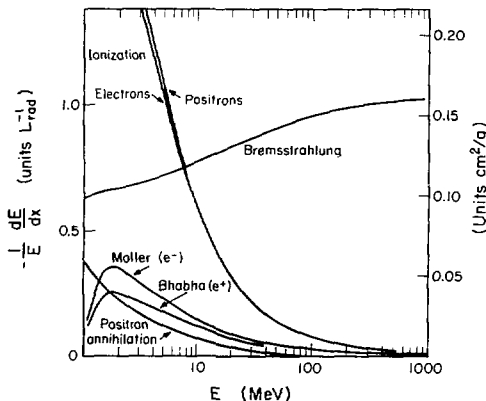
Mean Electron Range in Lead, Copper, Carbon, and Liquid Hydrogen

Mean range of electrons in the continuous-slowing-down approximation, taking into account energy loss by collisions with atomic electrons and by bremsstrahlung; strong fluctuations are to be expected for individual tracks. This range is the total path length; the "practical range" — a common measure of straight-line penetration distance — is shorter because of multiple Coulomb scattering, which becomes increasingly important as the electron slows down. E.g., for a fast electron the rms projected angle due to multiple Coulomb scattering reaches 1 radian by the time the electron has slowed to 0.4 MeV in hydrogen, 1.5 MeV in carbon, 9 MeV in copper, and 24 MeV (off scale) in lead. Electron energy deposition and penetration probability vs. range are discussed by L. V. Spencer, "Energy Dissipation by Fast Electrons," NBS Monograph #1, 1959, and S. M. Seltzer, "Transmission of Electrons through Foils," NBSIR 74, 457 (1974). Electrons which have energy less than 0.2 MeV in Ar, 1.5 MeV in Cu, 3.5 MeV in Sn, and 5 MeV in Pb are likely to deposit 10% of their energy behind their starting plane.

The practical range, R_p , is defined as that absorber thickness obtained by extrapolating to zero the linearly decreasing part of the curve of penetration probability vs. absorber thickness. Data for Al in the T range of the figure are available, and fit (to ~10%) $R_p = AT[1-B/(1+CT)]$ mg cm⁻² [a form suggested by K.-H. Weber, Nucl. Inst. Meth. 25, 261 (1964)], with $A=0.55$ mg cm⁻² keV⁻¹, $B=0.9841$, and $C=0.0030$ keV⁻¹. At this penetration depth, 90-95% of the incident electrons have stopped. Data for other elements are sketchy, but suggest that higher-Z (<50) elements have $1 \leq R_p/R_p(\text{Al}) \leq 1.4$ below ~10 keV, and $0.6 \leq R_p/R_p(\text{Al}) \leq 1$ above ~100 keV. The "critical energy" (above which the energy loss due to bremsstrahlung exceeds that due to ionization, and showering becomes important) is 400 MeV for hydrogen, 100 MeV for carbon, 25 MeV for copper, and 10 MeV for lead. The mean positron range may differ from the mean electron range by several percent. See Berger and Seltzer, NASA SP-3012 (1964) and SP-3036, and P. Trower, UCRL-2426, Vol. 111, Rev. (1966). 1-10 keV range was obtained by linear extrapolation; in this region the true range may actually lie above the curves.



Fractional Energy Loss for e⁺ and e⁻ in Lead

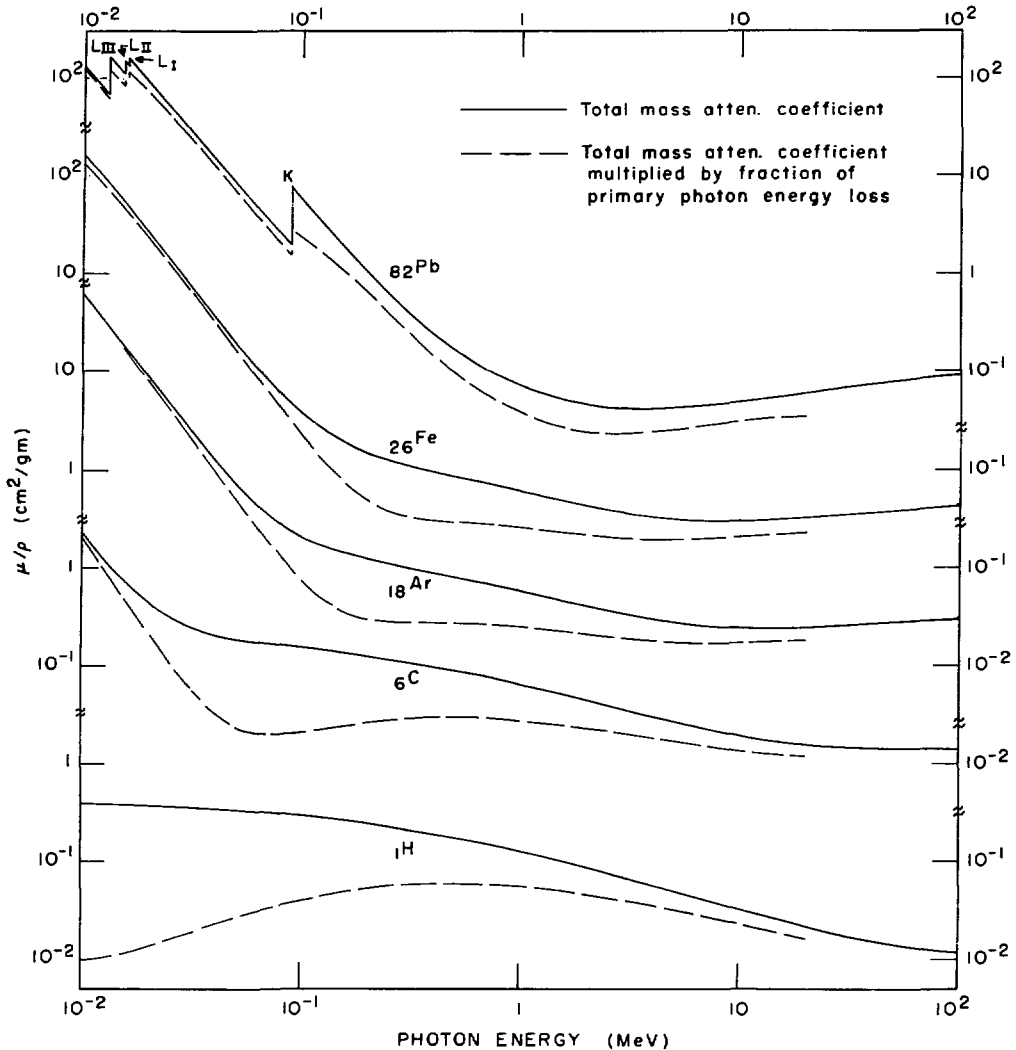


Fractional energy loss per radiation length in lead as a function of electron or positron energy. Electron (positron) scattering is considered as ionization when the energy loss per collision is below 0.255 MeV, and as Moller (Bhabha) scattering when it is above.

Adapted from Fig. 3.2 from Messel and Crawford, Electron-Photon Shower Distribution Function Tables for Lead, Copper and Air Absorbers, Pergamon Press, 1970. Messel and Crawford use $L_r(\text{Pb}) = 5.82$ g/cm², but we have modified the figures to reflect the value given in the Table of Atomic and Nuclear Properties of Materials (following), namely $L_r(\text{Pb}) = 6.4$ g/cm². The development of electron-photon cascades is approximately independent of absorber when the results are expressed in terms of inverse radiation lengths (i.e., scale on left of plot).

PARTICLE DETECTORS, ABSORBERS, AND RANGES (Cont'd)

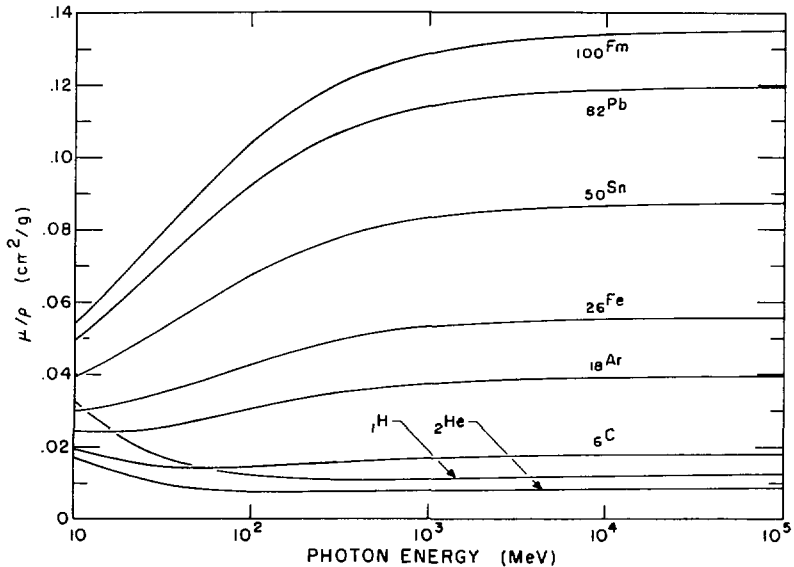
Photon Mass Attenuation Coefficients, Energy Deposition



The photon mass attenuation coefficient for various absorbers as a function of photon energy (solid curves). For a homogeneous medium of density ρ , the intensity I remaining after traversal of thickness t is given by $I = I_0 \exp(-\mu t)$. The accuracy is a few percent. Interpolation to other Z should be done in the cross section $\sigma = (\mu/\rho) M/N_A \text{ cm}^2/\text{atom}$, where M is the atomic weight of the absorber material and N_A is Avogadro's number. For a chemical compound or mixture, use $(\mu/\rho)_{\text{eff}} \approx \sum_i w_i (\mu/\rho)_i$, accurate to a few percent, where w_i is the proportion by weight of the i^{th} constituent. The dashed curve is the mass energy-absorption coefficient, giving μ/ρ multiplied by the fraction of photon energy deposited in a small volume (assumed large enough to contain the ranges of most secondary electrons) about the interaction. This fraction is smaller than 1.0 because such processes as Compton scattering and electron bremsstrahlung imply radiation of some of the energy away from the immediate area. The absorption coefficient is an approximation to the energy available for chemical, biological, and other effects associated with exposure to ionizing radiation. See next page for high energy range. From Hubbell, Gimm, and Øverbø, J. Phys. Chem. Ref. Data 9, 1023 (1980). Figures courtesy J. H. Hubbell.

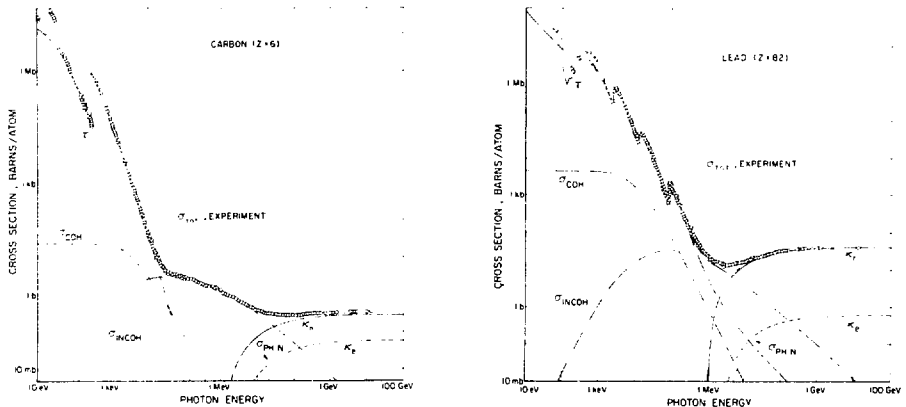
PARTICLE DETECTORS, ABSORBERS, AND RANGES (Cont'd)

Photon Mass Attenuation Coefficients (High Energy)



The photon mass attenuation coefficient, high energy range (note that ordinate is linear scale). See caption on previous page for details.

Contributions to Photon Cross Section in Carbon and Lead



Photon total cross sections as a function of energy in carbon and lead, showing the contributions of different processes.

- τ = Atomic photo-effect (electron ejection, photon absorption)
- σ_{COH} = Coherent scattering (Rayleigh scattering \rightarrow atom neither ionized nor excited)
- σ_{INCOH} = Incoherent scattering (Compton scattering off an electron)
- K_N = Pair production, nuclear field
- K_e = Pair production, electron field
- $\sigma_{\text{PH.N.}}$ = Photonuclear absorption (nuclear absorption, usually followed by emission of a neutron or other particle)

From Hubbell, Gimm, and Øverbø, J. Phys. Chem. Ref. Data 9, 1023 (1980). Figures courtesy J. H. Hubbell.

PARTICLE DETECTORS, ABSORBERS, AND RANGES (Cont'd)

Atomic and Nuclear Properties of Materials*

Material	Z	A	Nuclear total cross section σ_T [barn]	Nuclear ^b inelastic cross section σ_I [barn]	Nuclear ^c collision length λ_T [μm^2]	Nuclear ^c interaction length λ_I [μm^2]	dE/dx min ^d		Radiation length ^e L_{rad} [cm]	Density ^f [g/cm^3] (^g) is for gas [g/ℓ]	Refractive index n; ^f (ⁱ) is (n-1) $\times 10^6$ for gas	
							[MeV/c ²]	[MeV/cm]				
H ₂	1	1.01	0.0387	0.032	43.3	52.4	4.12	0.292	63.05	890	0.0708(0.090)	1.112(140)
D ₂	1	2.01	0.073	0.062	45.7	53.8	2.07	0.342	126.1	764	0.165	1.28
He	2	4.00	0.136	0.109	49.9	60.9	1.94	0.243	94.32	755	0.125(0.178)	1.024(35)
Li	3	6.94	0.216	0.158	53.8	72.9	1.58	0.843	82.76	155	0.534	--
Be	4	9.01	0.268	0.199	55.8	75.2	1.61	2.97	65.19	35.3	1.848	--
C	6	12.01	0.331	0.231	60.2	86.3	1.78	4.03	42.70	18.8	2.265 ^f	--
N ₂	7	14.01	0.379	0.262	61.4	88.8	1.82	1.47	37.99	47.0	0.808(1.25)	1.205(300)
O ₂	8	16.00	0.420	0.288	63.2	92.2	1.82	2.07	34.24	30.0	1.14(1.43)	1.22(266)
Ne	10	20.18	0.502	0.340	66.7	98.5	1.73	2.09	28.94	24.0	1.207(0.90)	1.092(67)
Al	13	26.98	0.634	0.421	70.6	106.4	1.62	4.37	24.01	8.9	2.70	--
Ar	18	39.95	0.850	0.554	78.0	119.7	1.51	2.11	19.55	14.0	1.40(1.78)	1.233(283)
Fe	26	55.85	1.113	0.703	83.3	131.9	1.48	11.6	13.84	1.76	7.87	--
Cu	29	63.54	1.232	0.782	85.6	134.9	1.44	12.9	12.86	1.43	8.96	--
Sn	50	118.69	1.967	1.191	100.2	165.5	1.26	9.2	8.82	1.21	7.31	--
W	74	183.85	2.767	1.649	110.3	185.1	1.16	22.4	6.76	0.35	19.3	--
Pb	82	207.19	2.960	1.776	116.2	193.7	1.13	12.8	6.37	0.56	11.35	--
U	92	238.03	3.378	1.983	117.0	199.3	1.09	≈ 20.7	6.00	≈ 0.32	≈ 18.95	--
Air (20°C)					62.0	90.0	1.82	0.0022	36.66	30423	0.001205(1.29)	1.000273(293)
H ₂ O					60.1	84.9	2.03	2.03	36.08	36.1	1.00	1.33
Shielding concrete ^h					67.4	99.9	1.72	4.25	26.7	10.7	2.5	--
SiO ₂ (quartz)					67.0	99.2	1.72	3.78	27.05	12.3	2.2	1.458
H ₂ (bubble chamber 26°K)					43.3	52.4	4.12	≈ 0.26	63.05	≈ 1000	$\approx 0.063^j$	1.112
D ₂ (bubble chamber 31°K)					45.7	53.8	2.07	≈ 0.29	126.1	≈ 900	$\approx 0.140^j$	1.110
H-Ne mixture (50 mole percent) ^j					65.0	94.5	1.84	0.75	29.70	73.0	0.407	1.092
Propane (C ₃ H ₈) ^k					56.5	77.2	2.25	0.92	45.38	111	0.41(2.0)	1.25(1005)
Freon 13B1 (CF ₃ Br) ^k					76.8	117	1.56	≈ 2.34	16.53	≈ 11.0	$\approx 1.50(8.71)$	1.238(750)
Iford emulsion G5					82.0	134	1.44	5.49	11.02	2.89	3.815	--
NaI					94.8	152	1.32	4.84	9.49	2.59	3.67	1.775
LiF					62	89.2	1.63	4.30	39.25	14.9	2.64	1.394
BGO (Bi ₄ Ge ₂ O ₁₂)					97.4	156	1.27	9.0	7.98	1.12	7.1	2.15
Polystyrene, scintillator (CH) ^f					58.4	82.0	1.95	2.01	43.8	42.4	1.032	1.581
Lucite, Plexiglas (C ₅ H ₈ O ₂)					59.2	83.6	1.95	≈ 2.30	40.55	≈ 34.4	1.16-1.20	≈ 1.49
Polyethylene (CH ₂)					56.9	78.8	2.09	≈ 1.95	44.8	≈ 47.9	0.92-0.95	--
Mylar (C ₃ H ₄ O ₂)					60.2	85.7	1.86	2.58	39.95	28.7	1.39	--
Borosilicate glass (Pyrex) ^m					66.2	97.6	1.72	3.84	28.3	12.7	2.23	1.474
CO ₂					62.4	90.5	1.82	0.0033	36.2	20220	(1.79)	(410)
Methane CH ₄					54.7	74.0	2.41	0.0017	46.5	64850	0.423(0.717)	(444)
Isobutane C ₄ H ₁₀					56.3	77.4	2.22	0.0059	45.2	16930	(2.67)	(1270)
Freon 12 (CCl ₂ F ₂) ⁿ					70.6	106	1.62	0.0080	23.7	4810	(4.93)	(1080)
Freon 13 (CClF ₃) ⁿ					68.1	101	1.64	0.0070	27.15	6370	(4.26)	(720)
Silica Aerogel ^o					65.5	95.7	1.83	≈ 0.36	29.85	≈ 150	0.1-0.3	1.0+0.25 _p
Spark or proportional chamber ^p					0.028%	0.020%	--	0.034	0.067%		0.019	

* Table revised March 1982 by J. Engler. For details, see Report KfK B000, Kernforschungszentrum D 75 Karlsruhe, P.O. Box 3640, Federal Republic of Germany.

- σ_{total} at 80-240 GeV for neutrons ($\approx \sigma$ for protons) from Murthy et al., Nucl. Phys. B92, 269 (1975).
- $\sigma_{\text{inelastic}} = \sigma_{\text{total}} - \sigma_{\text{elastic}} - \sigma_{\text{neutronic}}$; for neutrons at 60-375 GeV from Roberts et al., Nucl. Phys., B159, 56 (1979). For protons and other particles, see Carroll et al., Phys. Lett. 80B, 319 (1979); note that $\sigma_I(p) \approx \sigma_I(n)$.
- Mean free path between collision (AT) or inelastic interaction (AI), calculated $\lambda = A/(N \times \sigma)$.
- For minimum-ionizing protons and pions from Barkas and Berger, Tables of Energy Losses and Ranges of Heavy Charged Particles, NASA-SP-3013 (1964). For electrons see: Penetration of Charged Particles in Matter, NAS-NS39 (1964).
- From Y.S. Tsai, Rev. Mod. Phys. 46, 815 (1974).
- Values for solids, or the liquid phase at boiling point. Values in parentheses for gaseous phase STP (0°C, 1 atm.), except where noted. Refractive index given for sodium D line.
- For pure graphite, industrial graphite density may vary 2.1 - 2.3 g/cm³.
- Standard shielding blocks, typical composition O, 52%; Si 32.5%, Ca 6%, Na 1.5%, Fe 2%, Al 4% plus reinforcing iron bars. Attenuation length $\ell = 115 \pm 5$ g/cm², also valid for earth (typical $\rho = 2.15$) from CERN-LRL-RHEI Shielding exp., UCRL-17841 (1968).
- Density may vary about $\pm 3\%$, depending on operating conditions.
- Values for typical working conditions with H₂ target: 50 mole percent, 29°K, 7 atm.
- Values for typical chamber working conditions: Propane $\approx 57^\circ\text{C}$, 8-10 atm. Freon 13B1 $\approx 28^\circ\text{C}$, 8-10 atm.
- Typical scintillator; e.g., PILOT B and NE 102A have an atomic ratio H/C = 1.10.
- Main components: 80% SiO₂ + 12% B₂O₃ + 5% Na₂O.
- Used in Cerenkov counters. Values at 26°C and 1 atm. Indices of refraction from E.R. Hayes et al., ANL-6916 (1964).
- n(SiO₂) + 2n(H₂O) used in Cerenkov counters, $\rho = \text{density in g/cm}^3$. From M. Cantin et al., Nucl. Instr. Meth. 118, 177 (1974).
- Values for typical construction: 2 layers 50 μm Cu/Be wires, 8 mm gap, 60% argon, 40% isobutane or CO₂; 2 layers 50 μm Mylar/Aclar foils.

ELECTROMAGNETIC RELATIONS

Maxwell's Equations

Quantity	CGS (statvolt, statamp., sec/cm ²)	MKSA (coulomb, amp., ohm)
Potentials:	$V = \sum \frac{\text{charges}}{r}$ $\vec{A} = \frac{1}{c} \sum \frac{\text{currents}}{r}$ <p>$c = \text{speed of light in vacuum}$</p>	$V = \frac{1}{4\pi\epsilon_0} \sum \frac{\text{charges}}{r}$ $\vec{A} = \frac{\mu_0}{4\pi} \sum \frac{\text{currents}}{r}$ <p>$\epsilon_0 = \frac{1}{36\pi} 10^{-9} \text{ MKSA}$</p> <p>$\mu_0 = 4\pi 10^{-7} \text{ MKSA}$</p>
Fields:	$\vec{E} = -\vec{\nabla}V - \frac{\partial \vec{A}}{\partial t}$ $\vec{B} = \vec{\nabla} \cdot \vec{A}$	$\vec{E} = -\vec{\nabla}V - \frac{\partial \vec{A}}{\partial t}$ $\vec{B} = \vec{\nabla} \cdot \vec{A}$
Materials:	$\vec{D} = \epsilon \vec{E}, \quad \vec{B} = \mu \vec{H}$	$\vec{D} = \epsilon \vec{E}, \quad \vec{B} = \mu \vec{H}$
Force:	$\vec{F} = q \left(\vec{E} + \frac{\vec{v}}{c} \times \vec{B} \right)$	$\vec{F} = q \left(\vec{E} + \vec{v} \times \vec{B} \right)$
Maxwell:	$\vec{\nabla} \cdot \vec{D} = 4\pi \rho$ $\vec{\nabla} \cdot \vec{E} = -\frac{1}{c} \frac{\partial \rho}{\partial t}$ $\vec{\nabla} \times \vec{B} = 0$ $\vec{\nabla} \cdot \vec{H} = \frac{4\pi \vec{j}}{c} + \frac{1}{c} \frac{\partial \rho}{\partial t}$	$\vec{\nabla} \cdot \vec{D} = \rho$ $\vec{\nabla} \cdot \vec{E} = -\frac{\rho}{\epsilon_0}$ $\vec{\nabla} \times \vec{B} = 0$ $\vec{\nabla} \times \vec{H} = \vec{j} + \frac{\partial \rho}{\partial t}$
relativistic transformations:	$\vec{E}'_{\parallel} = \vec{E}_{\parallel}$ $\vec{E}'_{\perp} = \gamma \left(\vec{E}_{\perp} + \frac{1}{c} \vec{v} \times \vec{B} \right)$ $\vec{B}'_{\parallel} = \vec{B}_{\parallel}$ $\vec{B}'_{\perp} = \gamma \left(\vec{B}_{\perp} - \frac{1}{c} \vec{v} \times \vec{E} \right)$	$\vec{E}'_{\parallel} = \vec{E}_{\parallel}$ $\vec{E}'_{\perp} = \gamma \left(\vec{E}_{\perp} + \vec{v} \times \vec{B} \right)$ $\vec{B}'_{\parallel} = \vec{B}_{\parallel}$ $\vec{B}'_{\perp} = \gamma \left(\vec{B}_{\perp} - \frac{1}{c^2} \vec{v} \times \vec{E} \right)$

Impedances: Alternating Currents (MKSA)

Ohm's law: $V = ZI, \quad V = V_0 e^{i\omega t}$

- Impedance of self-inductance L : $Z = i\omega L$.
- Impedance of a capacitor of capacitance C : $Z = \frac{1}{i\omega C}$.
- Impedance of a flat conductor of width w at high frequency:

$$Z = \frac{(1 + i)}{\sqrt{2}} \rho \sqrt{\frac{2}{\omega \mu_0}}$$

$\rho = \text{resistivity in } 10^{-8} \text{ m}$

- ~ 1.7 for Cu
- ~ 2.4 for Au
- ~ 2.8 for Al
- (Al alloys may have up to double this value.)
- ~ 5.5 for W
- ~ 73 for SS 304
- ~ 100 for Nichrome

$\delta = \text{effective skin depth.}$

$$\delta = \sqrt{\frac{\rho}{\pi \omega \mu_0}} = \frac{6.6 \text{ cm}}{\sqrt{f(\text{sec}^{-1})}} \quad \text{for Cu.}$$

4. Impedance of free space: $Z = \sqrt{\mu_0/\epsilon_0} = 376.7 \Omega$

Capacitance C and Inductance L per Unit Length (MKSA)

1. For flat plates of width w , separated by d in vac:

$$C = \frac{\epsilon_0 w}{d}; \quad L = \frac{\mu_0 d}{w}$$

2. For coax cable of interior and exterior radii r_1 and r_2 :

$$C = \frac{2\pi\epsilon}{\ln(r_2/r_1)}; \quad L = \frac{\mu_0}{2\pi} \ln(r_2/r_1)$$

- $\epsilon = \text{dielectric constant}$ (2 to 6 for plastics)
- $\mu = \text{magnetic susceptibility}$ (4 to 6 for porcelain, glasses)

Transmission Lines (No Loss) (MKSA)

Velocity = $1/\sqrt{LC} = 1/\sqrt{\epsilon\mu}$

Impedance = $\sqrt{L/C}$

L and C are inductance and capacitance per unit length.

Synchrotron Radiation (CGS)

Energy loss/revolution = $\frac{4\pi}{3} \frac{r_e^2}{3} \beta^3 \gamma^4 \dots$ (omit radius)

For electrons ($\beta \approx 1$), $\frac{dE(\text{MeV})}{drev} = 0.0885 [E(\text{GeV})]^4$ (MeV/rev)

Critical frequency: $\omega_c = 3\gamma^3 \frac{v}{R}$

Frequency spectrum (for $\gamma \gg 1$):

$$I(\omega) \approx 3.3 \frac{e^2}{c} \left(\frac{\omega}{\omega_c} \right)^{1/3} e^{-2\omega/\omega_c}$$

$$I(\omega) \approx (1.0, 1.6, 1.4, 1.5, 0.88) \frac{e^2}{c} \left(\frac{\omega}{\omega_c} \right)^{1/3}$$

at $\frac{\omega}{\omega_c} = 0.01, 0.1, 0.2, 1.0, 2.0$, respectively.

$$I(\omega) \approx \sqrt{3\pi} \frac{e^2}{c} \left(\frac{\omega}{\omega_c} \right)^{1/2} e^{-2\omega/\omega_c}, \quad \omega \ll \omega_c$$

The radiation is confined to angles $\lesssim 1/\gamma$ relative to the instantaneous direction of motion.

See J. D. Jackson, *Classical Electrodynamics*, 2nd edition (John Wiley & Sons, New York, 1975) for more formulae and details. (Prepared April 1974; revised April 1980.)

RADIOACTIVITY AND RADIATION PROTECTION

Unit of activity = Curie:

$$1 \text{ Ci} = 3.7 \times 10^{10} \text{ disintegrations/sec}$$

Unit of exposure dose for x and γ radiation = Roentgen:

$$1 \text{ R} = 1 \text{ esu/cm}^2 = 87.8 \text{ erg/g} (5.49 \times 10^4 \text{ MeV/g}) \text{ of air}$$

Unit of absorbed dose = rad:

$$1 \text{ rad} = 100 \text{ erg/g} (6.25 \times 10^7 \text{ MeV/g}) \text{ in any material}$$

Unit of dose equivalent (for protection) = rem:

$$\text{rems (Roentgen equivalents for man)} = \text{rads} \times \text{QF}$$

where QF (quality factor) depends upon the type of radiation and other factors. For γ rays and HE protons, QF = 1; for thermal neutrons, QF = 3; for fast neutrons, QF ranges up to 10; and for α particles and heavy ions, QF ranges up to 20.

Maximum permissible occupational dose for the whole body:

$$5 \text{ rem/year (maximum 3 rem/calendar quarter)}$$

Fluxes (per cm²) to liberate 1 rad in carbon:

$$3.5 \times 10^{10} \text{ minimum ionizing singly charged particles}$$

$$1.0 \times 10^9 \text{ photons of 1 MeV energy}$$

(These fluxes are correct to within a factor of 2 for all materials.)

Natural background: 120 to 140 millirem/year

cosmic radiation (charged particles + neutrons) ~ 25 mrem/yr

cosmic radiation (γ rays) ~ 25 mrem/yr

radiation from rocks and air (γ rays) ~ 73 mrem/yr

Cosmic ray background in counters: ~ 1/min/cm²/ster

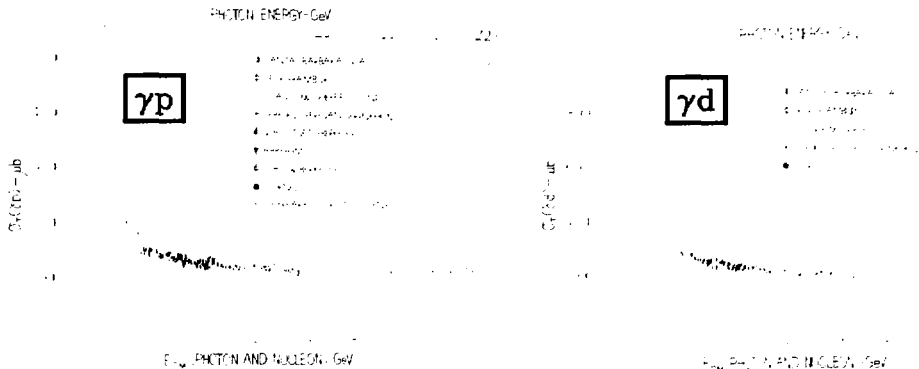
PERIODIC TABLE OF THE ELEMENTS

IA		IIA		H 1.0079										IIIA					IVA		VA		VIA		VIIA		2 He 4.00260																																																																										
3 Li 6.94	4 Be 9.01218			5 B 10.81	6 C 12.011	7 N 14.0067	8 O 15.9994	9 F 18.998403	10 Ne 20.17	11 Na 22.98977	12 Mg 24.3105	13 Al 26.98154	14 Si 28.0855	15 P 30.97376	16 S 32.06	17 Cl 35.453	18 Ar 39.948	19 K 39.0983	20 Ca 40.08	21 Sc 44.9559	22 Ti 47.90	23 V 50.9415	24 Cr 51.996	25 Mn 54.9380	26 Fe 55.847	27 Co 58.9332	28 Ni 58.71	29 Cu 63.546	30 Zn 65.38	31 Ga 69.723	32 Ge 72.59	33 As 74.9216	34 Se 78.96	35 Br 79.904	36 Kr 83.80	37 Rb 85.467	38 Sr 87.62	39 Y 88.9059	40 Zr 91.22	41 Nb 92.9064	42 Mo 95.94	43 Tc 98.9062	44 Ru 101.07	45 Rh 102.9055	46 Pd 106.4	47 Ag 107.868	48 Cd 112.41	49 In 114.82	50 Sn 118.69	51 Sb 121.75	52 Te 127.60	53 I 126.9045	54 Xe 131.30	55 Cs 132.9054	56 Ba 137.33	57-71 Rare Earths	72 Hf 178.49	73 Ta 180.947	74 W 183.85	75 Re 186.207	76 Os 190.2	77 Ir 192.22	78 Pt 195.09	79 Au 196.9665	80 Hg 200.59	81 Tl 204.37	82 Pb 207.2	83 Bi 208.9804	84 Po (209)	85 At (210)	86 Rn (222)	57 La 138.9055	58 Ce 140.12	59 Pr 140.9077	60 Nd 144.24	61 Pm (145)	62 Sm 150.4	63 Eu 151.96	64 Gd 157.25	65 Tb 158.9254	66 Dy 162.50	67 Ho 164.9304	68 Er 167.26	69 Tm 168.9342	70 Yb 173.04	71 Lu 174.967	89 Ac (227)	90 Th 232.0381	91 Pa 231.0389	92 U 238.029	93 Np 237.0482	94 Pu 244	95 Am (243)	96 Cm (247)	97 Bk (247)	98 Cf (251)	99 Es (254)	100 Fm (257)	101 Md (258)	102 No (259)	103 Lr (260)

Upper number is atomic number, expressing the positive charge of the nucleus in multiples of the electronic charge e . Lower number is atomic mass weighted by isotopic abundance in earth's surface, relative to the mass of the carbon 12 isotope, which has been arbitrarily assigned a mass of 12.00000 atomic mass units (amu). Numbers in parentheses are mass numbers (the whole number nearest the value of the atomic mass, in amu) of most stable isotope of that element. Adapted from the Handbook of Chemistry and Physics, 62nd Ed., 1981-1982. (Particle Data Group update, April 1982.)

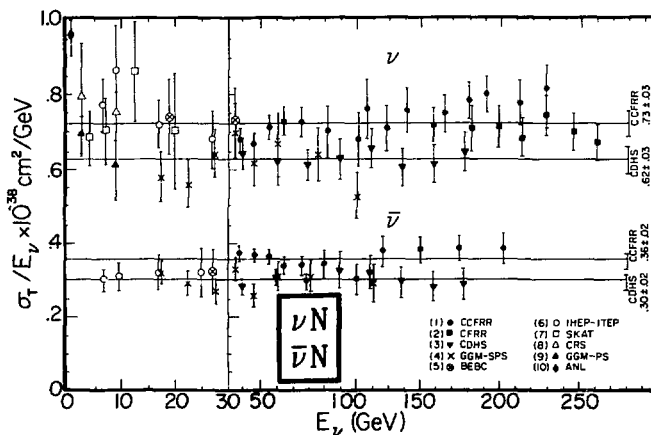
PLOTS OF CROSS SECTIONS AND RELATED QUANTITIES

NOTE: THE FIGURES IN THIS SECTION ARE INTENDED TO SHOW THE "BEST" OR "MOST REPRESENTATIVE" DATA IN THE OPINION OF THE COMPILER. THEY ARE NOT NECESSARILY COMPLETE COMPILATIONS OF ALL THE WORLD'S RELIABLE DATA.



σ_T total cross section versus photon energy (top scale) and photon-plus-nucleon total center-of-mass energy (lower scale). References: SANTA BARBARA-SLAC: D.O.Caldwell et al., Phys. Rev. **17**, 130 (1973); DESY-HAMBURG: H.Meyer et al., Phys. Lett. **33B**, 189 (1970); GLASGOW-SHEFFIELD-DNPL: T.A.Armstrong et al., Phys. Rev. **D5**, 1640 (1972); LEBEDEV-YEREVAN-SFRUKHOV: A.S.Belousov et al., Preprint 19, Moscow (1973); A.S.Belousov et al., Sov. Phys. Doklady **19**, 123 (1974), and A.S.Belousov et al., Sov. J. Nucl. Phys. **21**(3), 289 (1975); SLAC-BERKELEY-TITUS: J.Ballam et al., Phys. Rev. **D5**, 545 (1972); ABHHEN: H.L.Hilpert et al., Phys. Lett. **27B**, 474 (1968); SLAC and BERKELEY: J.Hallam et al., Phys. Rev. Lett. **21**, 1544 (1958), and H.H.Bingham et al., Phys. Rev. **DB**, 1277 (1973); CORNELL: S.Michalowski et al., Phys. Rev. Lett. **39**, 737 (1977); SANTA BARBARA-TORONTO-FNAL: D.O.Caldwell et al., Phys. Rev. Lett. **40**, 1222 (1978). See, also, the ep data of E.D.Bloom et al., SLAC-PP-054 (1969). Courtesy Gethin M. Lewis, Glasgow.

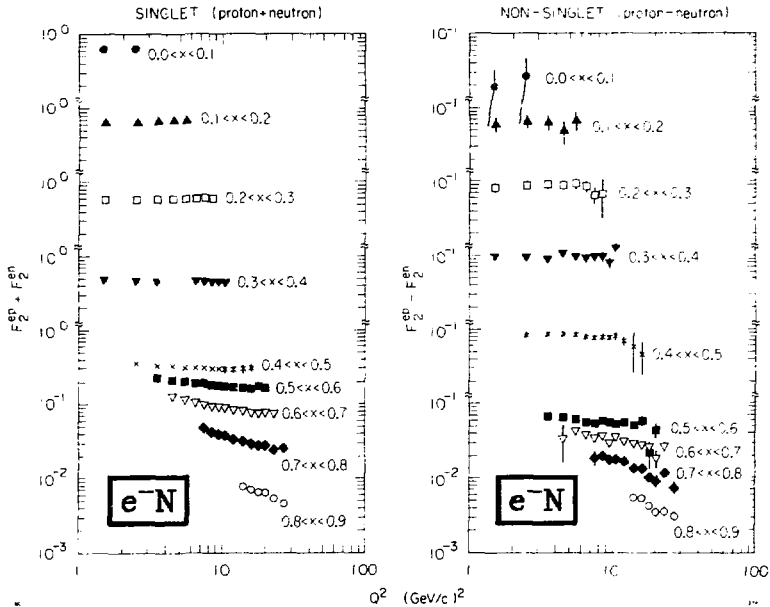
γd total cross section versus photon energy (top scale) and photon-plus-single-nucleon total center-of-mass energy (lower scale). References: SANTA BARBARA-SLAC: D.O.Caldwell et al., Phys. Rev. **D7**, 1362 (1973); DESY-HAMBURG: H.Meyer et al., Phys. Lett. **33B**, 189 (1970); GLASGOW-SHEFFIELD-DNPL: T.A.Armstrong et al., Nucl. Phys. **B41**, 445 (1972); LEBEDEV-YEREVAN-SFRUKHOV: A.S.Belousov et al., Sov. J. Nucl. Phys. **21**(3), 289 (1975); CORNELL: S.Michalowski et al., Phys. Rev. Lett. **39**, 737 (1977). Courtesy Gethin M. Lewis, Glasgow.



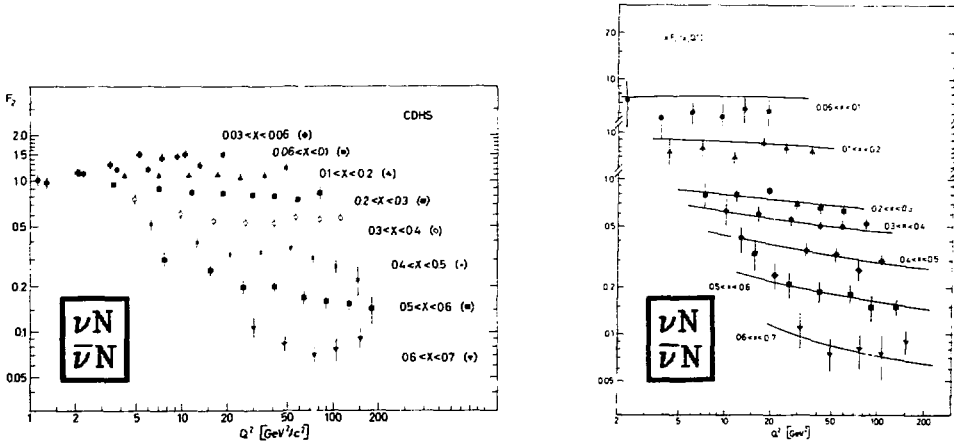
σ_T/E_ν for the muon neutrino and antineutrino charged-current total cross section as a function of neutrino energy. The error bars include both statistical and systematic errors. The straight lines are averages for the CCFRR and CDHS measurements. References: (1) R. Blair et al., in *Proc. of Neutrino '81*, Univ. of Hawaii (1981); (2) James Roy Lee, Ph.D. Thesis, Caltech (1981), "Measurements of νN Charged Current Cross Sections from $E_\nu=25$ GeV to $E_\nu=260$ GeV"; (3) J. G. H. de Groot et al., *Zeit. für Physik C - Particles and Fields* **1**, 143 (1979); (4) J. Morfin et al., Phys. Lett. **104B**, 275 (1981); (5) D. C. Colley et al., *Zeit. für Physik C - Particles and Fields* **2**, 187 (1979); (6) A. S. Vovenko et al., *Sov. J. Nucl. Phys.* **30**, 527 (1979); (7) D. S. Baranov et al., Phys. Lett. **81B**, 255 (1979); (8) C. Baltay et al., Phys. Rev. Lett. **44**, 916 (1980); (9) S. Ciampolillo et al., Phys. Lett. **84B**, 281 (1979); (10) S. J. Barish et al., Phys. Rev. **D19**, 2521 (1979). Courtesy M. Shaevitz, Nevis Laboratory.

PLOTS OF CROSS SECTIONS AND RELATED QUANTITIES (Cont'd)

Structure Functions

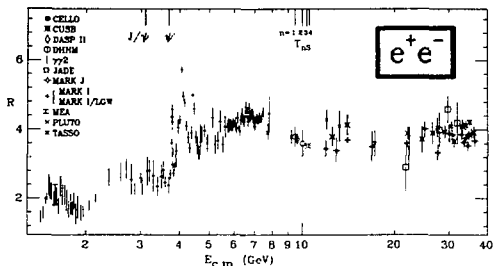


F_2 structure functions derived from inelastic electron-nucleon data taken at $s = Q^2 + 4$ with recoil mass > 2 GeV and four-momentum transfer squared $Q^2 > 1$ (GeV/c)² are shown. For definitions of F_2 , x , and Q^2 , see the "Relativistic Kinematics" section and the "Weak Interactions of Quarks and Leptons" section. $R = \sigma_L/\sigma_T = 0.21^3$ was assumed. Systematic errors are comparable in size to the data point symbols. Corrections for nucleon motion in deuterium have been made. These corrections are small except for $x > 0.7$. No error was included to account for uncertainties in this correction. References: 1) A. Bedek et al., Phys. Rev. D22, 1471 (1979); 2) W.B. Atwood, SLAC Report No. 185 (1975); 3) M.D. Mestayer, SLAC Report No. 214 (1978); 4) S. Stein et al., Phys. Rev. D12, 1884 (1975). Courtesy W. B. Atwood, SLAC.



Nucleon structure functions as measured by the CERN collaboration in high energy (30-200 GeV) charged-current neutrino- and anti-neutrino-nucleon scattering [J.G.H. de Groot et al., Z. Physik C - Particles and Fields 1, 143 (1979); reproduced by permission]. Definitions, and a discussion of the significance of these structure functions, may be found in the above reference, and also in the "Weak Interactions of Quarks and Leptons" section of the present work. See de Groot et al. for a discussion of experimental details, including corrections, etc. Curves are based on a QCD parametrization of Buras and Gaemers [Nucl. Phys. B132, 249 (1978)].

PLOTS OF CROSS SECTIONS AND RELATED QUANTITIES (Cont'd)



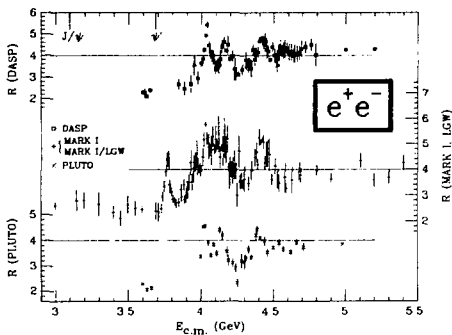
Measurements of $R = \sigma(e^+e^- \rightarrow \text{hadrons})/\sigma(e^+e^- \rightarrow \mu^+\mu^-)$, where the annihilation proceeds via one photon. The denominator is a calculated quantity:

$$\sigma(e^+e^- \rightarrow \mu^+\mu^-) = (hc)^2 \frac{4\pi^2}{3} \frac{e^4}{4E_{cm}^2} \int_0^1 d\epsilon \int_{-1}^1 d\cos\theta (1 - \epsilon^2 - \epsilon^2 \cos^2\theta) \quad (47)$$

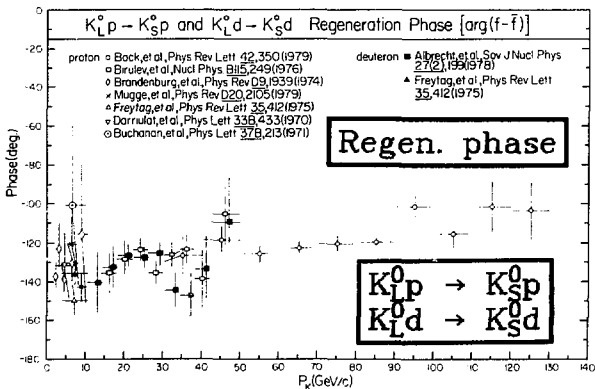
$$R = \frac{4\pi}{3} (hc)^2 \frac{e^4}{E_{cm}^2} = \frac{3\pi \alpha^2}{[E_{cm}(\text{GeV})]^2} \sum_k \beta_k^2 \quad \beta_k = \frac{p_k}{E_k}$$

for $e^+e^- \rightarrow \mu^+\mu^-$ for energies shown. Radiative corrections and, where pertinent, corrections for two-photon processes and τ production have been made. Note that the ADONE data (YY2 and MEA) is for τ hadrons. The points in the ψ (4770) region are from the MARK I Lead Glass Wall experiment. The DASP and PLUTO measurements have been omitted in the charm threshold region for clarity - they are shown in the expanded (lower) figure. Also for clarity, some points have been combined or shifted slightly (4%) in $E_{c.m.}$. Systematic normalization errors are not included; they range from 5-20%, depending on experiment. Note the suppressed zero. The horizontal extent of the plot symbols has no significance. The positions of the ψ , ψ' , and the four known Υ vector-meson resonances are indicated at the top of the figure. References:

- CELE - J. J. Behrns et al., Phys. Lett. **61**, to be published (preprint DESY 81-079); CUSB - L. Basso et al., submitted to Phys. Rev. Lett.; DASP - E. Brandelik et al., Phys. Lett. **76B**, 361 (1978); DASP II - S. Meseferi, thesis, DESY 80/AP25/41-3, to be published in Phys. Lett. B; DPHM - L. Bock et al. (DESY-Hamburg-Heidelberg-München collab.), Zeit. für Physik **CL**, 125 (1980); YY2 - C. Bacci et al., Phys. Lett. **86B**, 344 (1979); JADE - W. Bartel et al., Phys. Lett. **88B**, 121 (1979); MARK J - B. Newman, private communication; MARK I - J. E. Storaasli et al., SLAC-PUB-2831, 181-184, submitted to Phys. Rev. **2** (1982); MARK I - Lead Glass Wall - P. A. Kopialis et al., Phys. Rev. Lett. **39**, 526 (1977); P. A. Rapidis, thesis, SLAC-Report-29 (1979); MEA - B. Esposito et al., Lett. Nuovo Cimento **19**, 21 (1977); PLUTO - A. Böcker, thesis, Gesamthochschule Siegen, DESY 81-7/73 (1977); C. Gerke, thesis, Hamburg Univ. (1979); G. Berger et al., Phys. Lett. **81B**, 410 (1979); W. Lackner, thesis, RWTH Aachen, DESY 81/TC-81/11 (1981); TASSO - E. Brandelik et al., submitted to Phys. Lett. **8** (1982).

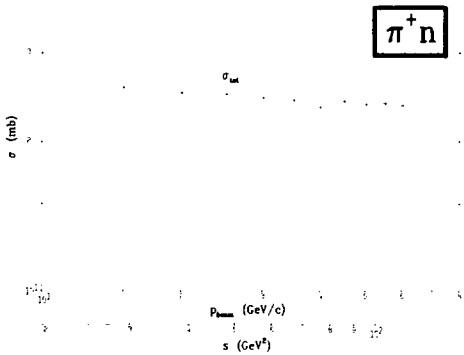
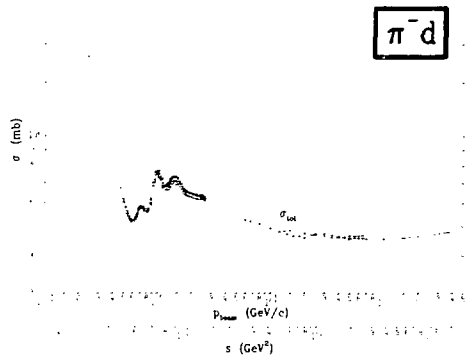
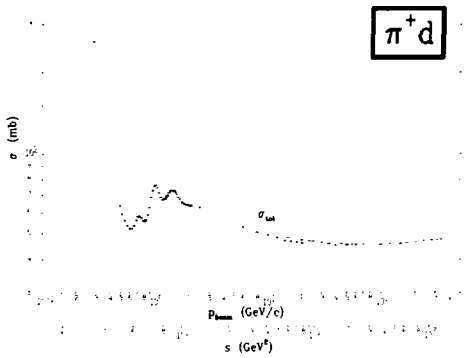
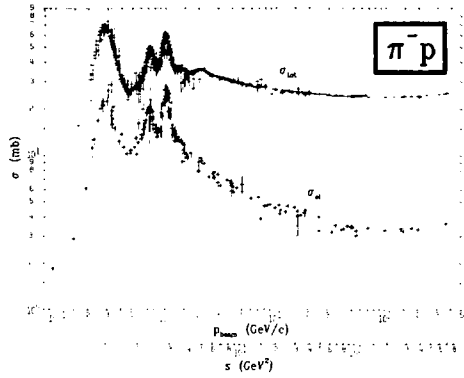
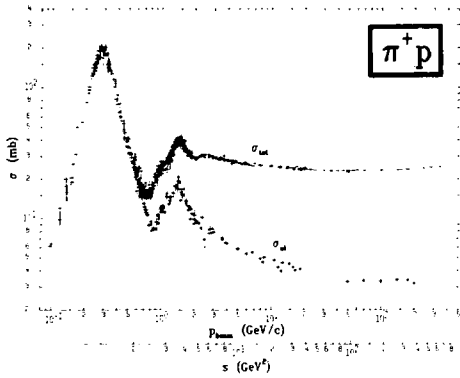


Lower figure: An expanded view of R measurements around charm threshold (the data points have been combined in this figure). We have arbitrarily added a horizontal line at $R=4$ as an aid to visual comparison of the three sets of data.



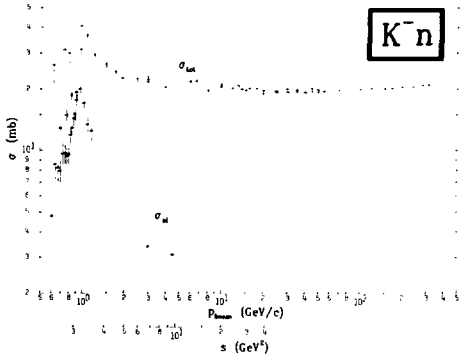
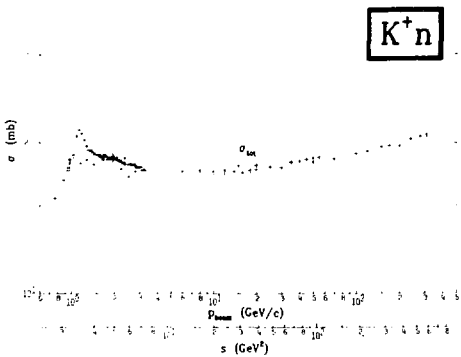
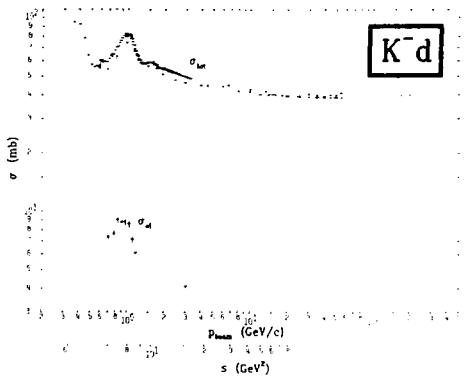
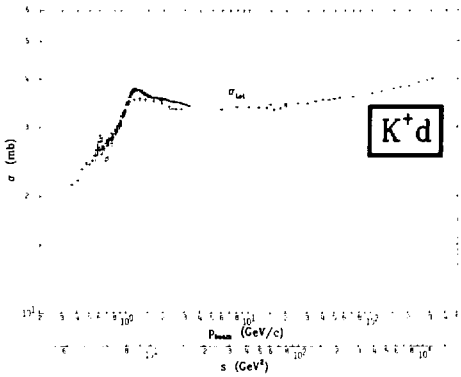
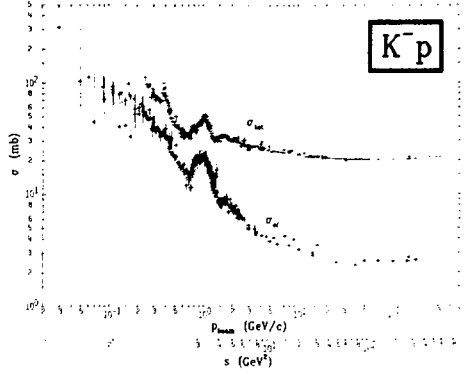
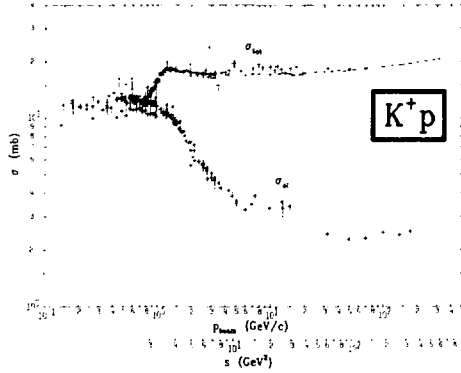
Phases of forward amplitudes for $K_L^0 p \rightarrow K_S^0 p$ (open symbols) and $K_L^0 d \rightarrow K_S^0 d$ (solid symbols). Courtesy S. Aronson, Brookhaven National Laboratory.

PLOTS OF CROSS SECTIONS AND RELATED QUANTITIES (Cont'd)



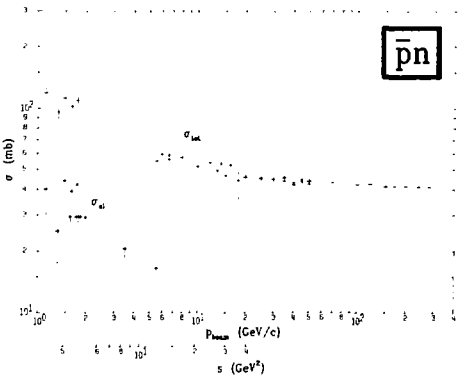
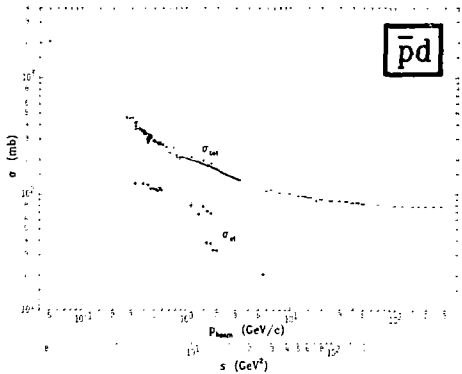
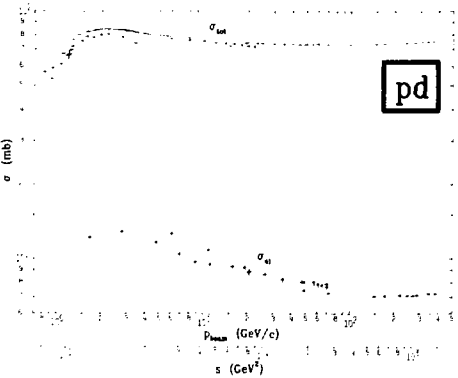
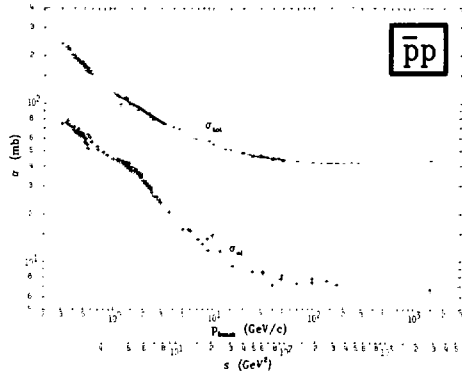
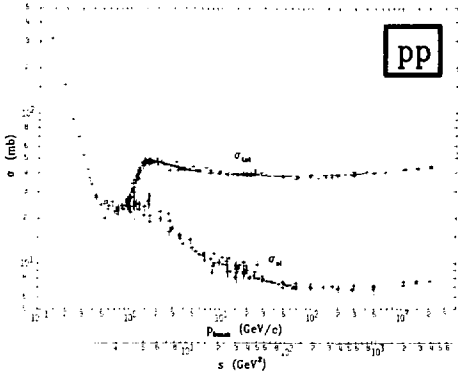
Hadronic total and elastic cross sections vs. laboratory beam momentum, P_{beam} and center-of-mass energy squared, s . Figures courtesy V. Flaminio, W. G. Moorhead, and D. R. O. Morrison, CERN.

PLOTS OF CROSS SECTIONS AND RELATED QUANTITIES (Cont'd)



Hadronic total and elastic cross sections vs. laboratory beam momentum, P_{beam} , and center-of-mass energy squared, s . Figures courtesy V. Flaminio, W. G. Moorhead, and D. R. O. Morrison, CERN.

PLOTS OF CROSS SECTIONS AND RELATED QUANTITIES (Cont'd)



Hadronic total and elastic cross sections vs. laboratory beam momentum, P_{beam} , and center-of-mass energy squared, s . Figures courtesy V. Flaminio, W. G. Moorhead, and D. R. O. Morrison, CERN.

DATA CARD LISTINGS

Illustrative Key

Name of particle as it appears in table. **XX(1200)** (74) KA MESON (1200, PPG. 2) 1-1

Arrow indicates this particle omitted from table. **ORIGINALLY CALLED XXX, OMITTED FROM TABLE**

Quantity tabulated below. **74 XX(1200) MASS (MEV)**

Code for quantity tabulated (M=mass, W=width, etc.).

M	116.1	11.	MERRILL	66 HBC	0 3.2 K-P	7/66
M	118.1	10.	LYNCH	67 HBC	2.7 P1-P	6/67
M	119.8	10.	PIERCE	68 ASPK	2.1 K-P	9/68
M	120.8	8.	FENNER	69 HBC	4.2 P1-P	9/69
M	121.0	8.	SMITH	81 MMS	3.5 P1-P	1/81*
M	AVG	1206.9	5.1	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)		

Symbol used to key together data card and related comments. **LYNCH DATA HAS QUESTIONABLE BACKGROUND SUBTRACTION**

Number of events above background. **AVG 1206.9 5.1**

Measured value (parentheses indicate value not used in average). **74 XX(1200) WIDTH (MEV)**

M	35.	5.	MERRILL	66 HBC	0 3.2 K-P	7/66
M	50.	10.	PIERCE	68 ASPK	2.1 P-P	9/68
M	70.	40.	FENNER	69 HBC	4.2 P1-P	9/69
M	(50.)	0P-3.5	SMITH	81 MMS	3.5 P1-P	1/81*
M	AVG	38.4	6.0	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.3)		

± error in measured value (— field blank if error symmetric; parentheses on error only indicate data not used in average due to problems with error estimation).

Average value (and error) of quantity tabulated. **WEIGHTED AVERAGE = 38.4 ± 6.0**
ERROR STA BY 1.3

Vertical bar indicates average; width of horizontal bar on top is (scaled) error in average.

Value and error for each experiment.

Particle name, and quantum numbers (if arrow). **74 XX(1200) PARTIAL DECAY MODES**

P1	XX(1200) INTO $\pi^0 \pi^0$	DELAY MASSES
P2	XX(1200) INTO $\pi^+ K^0$	139. 135. 139.
		493. 493

Representative masses of decay products (used for calculating last column of Particle Property Tables).

Branching ratio (labeled by R_i). **74 XX(1200) BRANCHING RATIOS**

R1	XX(1200) INTO $\pi^0 \pi^0$ / TOTAL	(P1)	
R1	.66	.02	MERRILL 66 HBC 0 3.2 K-P 7/66
R1	(.68)	(.03)	LYNCH 67 HBC 2.7 P1-P 6/67
R1	LYNCH DATA HAS QUESTIONABLE BACKGROUND SUBTRACTION		
R1	0.675	0.021	FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.3)
R2	XX(1200) INTO $\pi^+ K^0$ / TOTAL	(P2)	
R2	.35	.05	PIERCE 68 ASPK 2.1 K-P 9/68
R2	0.325	0.017	FROM F. (ERROR INCLUDES SCALE FACTOR OF 1.3)
R3	XX(1200) INTO $\pi^+ K^0$ / $\pi^0 \pi^0$	(P2)/(P1)	
R3	.50	.03	FENNER 69 HBC 4.2 P1-P 9/69
R3	.41	.04	SMITH 81 MMS 3.5 P1-P 1/81*
R3	0.468	0.043	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.8)
R3	0.480	0.027	FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.3)

Value (and error) of quantity tabulated, as determined from constrained fit (using all measured branching ratios for this particle).

References, listed by year, then author.

MERRILL	66	PRL	16	143
LYNCH	67	PR	155	610
PIERCE	68	PL	278	230
FENNER	69	NC	618	372
SMITH	81	PRL	86	121

Abbreviated reference form used on data cards above.

Journal, report, preprint, etc. (see abbreviations on next page).

REFERENCES FOR XX(1200)

A. MERRILL	(SACLAY-CERN) 0.0P
B. LYNCH	(BNL)
N. PIERCE	(LRL)
D. FENNER, B. BEANE	(NYS&AMEX)
J. SMITH	(SLAC)

Author(s).

Quantum number determinations in this reference.

Institution(s) of author(s) (see abbreviations on next page).

Stable Particles

 γ, ν

		D GAMMA(10, J=1)					
		D GAMMA MASS (IN UNITS OF 10**+21 MEV)					
M	F	16.1	ER LESS	PATEL	65	SATELLITE DATA	10/69
M		(5.1)	OR LESS	GINTSRUP	64	SATELLITE DATA	10/69
M		(2.3)	OR LESS	GOLDBABER	69	SATELLITE DATA	10/69
M	F	(10.06)	OR LESS	FRANKEN	71	LOW FREQ RES CIA	3/72
M		(10.1)	OR LESS	WILLIAMS	71	TESTS GAUSS LAW	3/71
M		(6.2-13.9)	OR LESS	LOWENTHAL	73	GEN. RELATIVITY	2/72
M		(10.73)	OR LESS	HOLLWEG	74	ALFVEN WAVES	7/74
M		(0.6)	OR LESS	CLAYTON	75	JUPITER MAGFIELD	7/76
M	F		VALIDITY QUESTIONABLE, SEE CRITICISM IN #97LL	71	AND GOLDBABER	71	3/76

		REFERENCES FOR GAMMA			
GINTSRUP	64	SOV. ASTR. BJ	536	M. A. GINTSRUP	(ACAD SCI. USSR)
PATEL	65	PL	14	V. L. PATEL	(DURHAM)
GOLDBABER	69	PHL	21	A. GOLDBABER, M. NIETO	(STONY BROOK)
FRANKEN	71	PHL	26	D. A. FRANKEN, G. W. AMPULSKI	(MICH.)
WILLIAMS	71	PHL	26	WILLIAMS	(MELBURN)
LOWENTHAL	73	PH	OR	D. D. LOWENTHAL	(UCLA)
HOLLWEG	74	PHL	32	J. V. HOLLWEG (NATL CENTER FOR ATMOS RESRCH)	(CITYSTON/CASLE)
DAVIS	75	PHL	35	GOLDBABER, NIETO	

		PAPERS NOT REFERRED TO IN DATA CARDS			
GOLDBABER	71	PH	43	A. S. GOLDBABER, M. M. NIETO	(STONY BROOK/UCSB)
KROLL	71	PHL	26	M. KROLL	(SLAC)
BYRNE	77	ASTR.	73	J.C. BYRNE	(UCLA)

Neutrinos

(by R. E. Shrock, State Univ. of New York, Stony Brook)

With this issue the section on neutrino properties has been expanded and reorganized. As before, there are listings which deal specifically with ν_e , ν_μ , and ν_τ . In addition, in the category of searches near the end of the Stable Particle Listings, we include sections which deal with correlated bounds on neutrino masses and lepton mixing but which do not pertain to any one weak eigenstate individually. Furthermore, we include constraints from cosmological and astrophysical data. (Since this Review is a compendium of particle properties, traditionally derived more or less directly from particle and nuclear physics, we treat astrophysical data on a different footing from particle physics data and have been somewhat less comprehensive in our coverage of the former.)

In contrast to the other particles in this Review, the neutrinos ν_e , ν_μ , and ν_τ are defined as weak eigenstates (that is, states which couple weakly with unit strength to e , μ , and τ) and are not, in general, states of definite mass. In the conventional case, where all neutrinos were assumed to be massless and hence degenerate, it was possible to define the weak eigenstates to be simultaneously mass eigenstates. However, in the

Data Card Listings

For notation, see key at front of Listings.

general case of massive (non-degenerate) neutrinos, the weak eigenstates have no well-defined masses, but instead are linear combinations of mass eigenstates. Thus, if one considers this general case, as, of course, one does in quoting mass limits, it is inconsistent to assume that the weak and mass eigenstates coincide. Let us denote the charged leptons as the set $\{\ell_a\}$, $a = 1, \dots, n$, where $n \geq 3$ is the number of generations, with $\ell_1 \equiv e$, $\ell_2 \equiv \mu$, and $\ell_3 \equiv \tau$. In the standard $SU(2)_L \times U(1)$ electroweak theory¹ the mixing of the left-handed components of the mass eigenstates $(\nu_j)_L$ to form the weak gauge-group eigenstates $(\nu_a)_L$ is specified by the transformation

$$(\nu_a)_L = \sum_{j=1}^n U_{aj} (\nu_j)_L$$

where $U^+ = U^{-1}$. (In the case of Dirac neutrinos there are right-handed components of the ν_j , but they are singlets under the gauge group; in the case of Majorana neutrinos in the standard theory there are no right-handed components.) The ordering of the mass eigenbasis is defined such that U is as nearly diagonal as possible, i.e. $|U_{jj}|$ (no sum on j) $\geq |U_{jk}|$, $k \neq j$. This does not imply that $m(\nu_j) > m(\nu_k)$ if $j > k$, although this ordering might be regarded as natural in view of the similar one that obtains in the quark sector. The virtue of this convention is that a mass limit on " $m(\nu_{\lambda_a})$ " can be used as a definite limit on ν_j , $j = a$, the dominantly coupled mass eigenstate in ν_{λ_a} .

Thus, in this general case of n massive (Dirac or Majorana) neutrinos, decays such as $H^3 \rightarrow He^3 + e^- + \bar{\nu}_e$ and $\pi^+ \rightarrow \mu^+ + \nu_\mu$, which have been used to set the best bounds on the respective neutrino masses, really consist of incoherent sums of the separate decay modes $H^3 \rightarrow He^3 + e^- + \bar{\nu}_j$ and $\pi^+ \rightarrow \mu^+ + \nu_k$, where the ν_j , ν_k are mass eigenstates, and the indices j and k range over the subset $\{1, \dots, n\}$ allowed by phase space in these two respective decays.² The coupling strengths for the j 'th modes are given for the two decays by the factors $|U_{1j}|^2$ and $|U_{2j}|^2$, respectively. There are, in addition, certain kinematic factors depending on the $m(\nu_j)$ which

Data Card Listings

For notation, see key at front of Listings.

Stable Particles

v

enter in determining the branching ratio for the j 'th decay mode. Assuming that the off-diagonal elements of the lepton mixing matrix U are small relative to the diagonal elements, the dominantly coupled decays are the ones with coupling strength $|U_{aj}|^2$, $a = j$, i.e. $H^3 \rightarrow He^3 + e^- + \bar{\nu}_1$ and $\tau^+ \rightarrow \mu^+ + \nu_2$.

It follows that the old neutrino mass limits quoted in the literature for " $m(\nu_e)$ ", " $m(\nu_\mu)$ ", and " $m(\nu_\tau)$ " are meaningful only insofar as they are reinterpreted as limits on the corresponding mass eigenstates. Specifically, a bound such as the Bergkvist limit,³ " $m(\nu_e)$ " < 60 eV (90% CL), really constitutes a weighted limit on each of the mass eigenstates ν_j in the weak eigenstate ν_e which are kinematically allowed to occur in tritium decay and which are coupled with strength $|U_{1j}|^2$ sufficiently large to make a significant contribution to the observed spectrum. It is thus certainly a limit on ν_1 . If leptonic mixing is hierarchical as quark mixing is known to be (as least for the first three generations), i.e.

$|U_{jj}|^2 \gg |U_{jk}|^2$, $k \neq j$, then ν_1 is the only mass eigenstate significantly constrained by a bound on " $m(\nu_e)$." Furthermore, a neutrino mass limit cannot be stated in isolation; it always contains some implicit dependence on the relevant lepton mixing angles. Fortunately, this dependence is relatively unimportant for the dominantly coupled decay modes, i.e. $e\bar{\nu}_1$, $\mu\bar{\nu}_2$, and $\tau\bar{\nu}_3$. Since these modes were the ones responsible for the mass limits given previously, the latter can be reinterpreted without significant complication as proper limits on $m(\nu_j)$, $j = 1, 2$, and 3 , respectively.

In addition to mass and lifetime limits, we have added data on neutrino magnetic dipole moments. These are of interest because a massless, purely chiral (empirically, left-handed) Dirac neutrino cannot have a magnetic (or electric) dipole moment. The same is true for a Majorana neutrino, whether massless or massive, because of its defining property of being self-conjugate.

If one considers the possibility of nonzero masses for neutrinos, for consistency one must also consider the leptonic mixing which would in general occur concomitantly. Accordingly we have devoted one category in the searches section to

correlated bounds on neutrino masses and lepton mixing angles. These can be divided into two types. First, there are those due to decays involving neutrinos in the final state, which must be recognized to have the general multi-mode structure pointed out above. In the two most sensitive cases suggested as tests for neutrino masses and mixing,² one obtains a limit on $m(\nu_j)$ and $|U_{aj}|^2$ individually for each j . Second, there are those due to processes involving the propagation and subsequent interaction of neutrinos. The latter are often called neutrino "oscillation"³ limits, although this term is correct only if the differences in neutrino masses are sufficiently small relative to their momenta that the propagation is effectively coherent in a quantum mechanical sense; otherwise, the individual ν_j from a given decay such as $\pi_{\nu 2}$ or $K_{\nu 2}$ propagate in a measurably incoherent manner and there is no "oscillation". Experimentalists usually present their results in terms of a simplifying model in which mixing is assumed to occur only between two neutrino species. Then the transformation equation becomes

$$\begin{pmatrix} \nu_{\ell a} \\ \nu_{\ell b} \end{pmatrix} = \begin{pmatrix} \cos \theta & \sin \theta \\ -\sin \theta & \cos \theta \end{pmatrix} \begin{pmatrix} \nu_i \\ \nu_j \end{pmatrix}.$$

Let the distance between the source of the neutrinos and their point of interaction be labeled as x , and their energy as E . Assume furthermore that the $m(\nu_j)$ are such that the coherence assumption is valid. Then, the probability of an initial $\nu_{\ell a}$ being equal to $\nu_{\ell b}$ at time t or equivalently (given the above assumption) at distance $x = t$, is

$$|\langle \nu_{\ell b}(0) | \nu_{\ell a}(t) \rangle|^2 = \sin^2 2\theta \sin^2 \left(\frac{\Delta m^2 x}{4E} \right),$$

where

$$\Delta m^2 = m(\nu_i)^2 - m(\nu_j)^2.$$

Thus, neutrino oscillation experiments cannot measure individual neutrino masses, but only differences of masses squared, and indeed these are generally weighted in a more complicated way by mixing-matrix coefficients than in the two-species

Stable Particles

 ν, ν_e

model. Experimental results are presented as allowed regions on a plot, the axes of which are $|\Delta m^2|$ and $\sin^2 2\theta$. These are often summarized in terms of the asymptotic limits $|\Delta m^2|_{\max}$ for $\sin^2 2\theta = 1$, and $\sin^2 2\theta$ for "large" $|\Delta m^2|$, i.e., sufficiently large $|\Delta m^2|$ that the detector averages over many cycles of oscillation (or there ceases to be any coherence). We refer the reader to the original papers for the two-dimensional plots; for the purpose of these Listings we shall give only the asymptotic limits.

An important question has to do with whether neutrinos are Dirac or Majorana (self-conjugate) particles. In the former case neutrinoless double beta decay, $(Z, A) \rightarrow (Z+2, A) + e^- + e^-$, is forbidden from occurring.⁴ In the Majorana case it may occur, if (a) neutrinos are massive and/or (b) there are right-handed leptonic currents. In the light-neutrino case an upper limit on neutrinoless double beta decay yields a correlated upper bound on the quantity

$$\bar{m} \equiv \left| \sum_{j=1}^n U_{1j}^2 m(\nu_j) \right|$$

and η , the fractional admixture of right-handed leptonic current.

Further explanatory notes are included in the Listings.

References

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Data Card Listings

For notation, see key at front of Listings.

- For recent studies of neutrinoless double beta decay see H. Primakoff and S. P. Rosen, Ann. Rev. Nucl. Sci. **31**, 145 (1981); S. P. Rosen, Proc. of 1981 Intl. Conf. on Neutrino Physics and Astrophysics (Maui, Hawaii), eds. R. J. Sens et al., v.2, p. 76; W. C. Haxton, G. L. Stephenson, Jr., and D. Strottman, Phys. Rev. Lett. **47**, 153 (1981); M. Doi, T. Kotani, H. Nishiura, K. Okuda, and E. Takasugi, Phys. Lett. **103B**, 219 (1981), and Prog. Theor. Phys. **66**, 1739 and 1765 (1981).

U_e

1 NU-E(1,1/2)

NOT IN GENERAL A MASS EIGENSTATE

1 NU-E **MASS** (EV)

APPLIES TO NU-1, THE PRIMARY MASS EIGENSTATE IN NU-E. WOULD ALSO APPLY TO ANY OTHER NU-J WHICH MIXES STRONGLY IN NU-E AND WAS SUFFICIENTLY SMALL MASS THAT IT CAN OCCUR IN THE RESPECTIVE DECAYS.

NOTE -- THE ABBREVIATION NUJ IS USED BELOW FOR ANTIMU

M	1250.1	OR LESS	LANGER	57 CNTR	ANU-E, TRITIUM	
M	1590.1	OR LESS	HAMILTON	53 CNTR	ANU-E, TRITIUM	11/73
M	1550.1	1200.1	FRIEDMAN	58 CNTR	ANU-E, TRITIUM	11/73
M	14100.1	9R LESS	CL-67	BECK	68 CNTR	NU, SODIUM 22
M	1500.1	OR LESS	CL-90	DARIS	69 CNTR	ANU-E, TRITIUM
M	1320.1	OR LESS	CL-90	SALGO	69 CNTR	ANU-E, TRITIUM
M	160.1	OR LESS	CL-90	BERGVIST	72 CNTR	ANU-E, TRITIUM
M	180.1	OR LESS	CL-90	ROOF	72 CNTR	ANU-E, TRITIUM
M	1100.1	PR LESS	PIEL	73 CNTR	ANU-E, TRITIUM	1/81*
M	14.7551R	OR LESS	CL-90	CLARK	74 ASPK	KE3 DECAY
M	135.1	OR LESS	CL-90	TRETYAKOV	76 SPEC	ANU-E, TRITIUM
M	116.1	TO 46.	CL-99	LURIMOV	80 SPEC	ANU-E, TRITIUM
M	105.1	OR LESS	CL-95	SIMPSON	81 CNTR	ANU-C, TRITIUM
M	D	DARIS 69	WELCH(1.57)			DISAGREES WITH THEIR FIG.6. WE USE
M	D	FIG.6.				4/82*
M	L	TRETYAKOV 76	DATA INCLUDED. AT LEAST IN PART, IN LURIMOV 80.			4/82*
M	L	SEE THE DISCUSSION OF THE LURIMOV 80 RESULT BY BERGVIST 80.				1/82*
M	L	WE USE UPPER LIMIT FROM LURIMOV 80 IN THE STABLE PARTICLE TABLE.				4/82*
M	L	THEIR LOWER LIMIT NEEDS CONFIRMATION.				4/82*

1 (MU-1) - (LANU-1) MASS DIFF. (EV)

TEST OF CPT FOR A DIRAC NEUTRINO

DM	14.5551R	OR LESS	CL-90	CLARK	74 ASPK	KE3 DECAY	11/73
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1 NU-1 MEAN LIFE/MASS (UNITS SECEV)

I	R	3. E 2	CR MORE	REINES	74 CNTR	ANTI-NEUTRINO	3/76
I	R	REINES 74	LOCKED FOR NU-E OF NU-1 ZERO MASS DECAYING TO A NEUTRAL				3/78
I	R	OF LESSER MASS	GAMMA. USED LIQUID SCINT. DETECTOR NEAR 10% SEC				3/78
I	R	REACTOR.	FINDS LAB LIFETIME 0.67 SEC OR MORE. ABOVE VALUE OF				3/78
I	R	MEAN LIFE/MASS ASSUMES AVG. EFFECTIVE NEUTRINO ENERGY OF 3.2MEV.					3/78

1 NU-1 MAGNETIC MOMENT (UNITS EV/GAUSS)

MUST VANISH FOR MAJORANA NEUTRINO OR PURELY CHIRAL MASSLESS DIRAC NEUTRINO

MM	B	11.16-1730R	LESS	BERNSTEIN	63		1/82*
MM	B	BERNSTEIN 63	IS A THEORETICAL ANALYSIS OF REACTOR ANTI-NU-E				1/82*
MM	B	SCATTERING DATA.					1/82*

REFERENCES FOR NU-E

LANGER	57	PR	88	689	L. M. LANGER, P. J. D. MOFFAT	(INDIANA)
HAMILTON	53	PR	92	1521	D. HAMILTON, W. P. REYNOLDS, G. DODS	(PRINCETON)
FRIEDMAN	58	PR	109	2214	L. E. FRIEDMAN, L. Y. CHEN, G. SMITH	(MICHIGAN)
BECK	68	PR	204	216 229	E. BECK, M. DANIEL	(MPSH)
BERNSTEIN	63	PR	132	1227	BERNSTEIN, RUBENMAN, FEINBERG	(YUKON COL)
DARIS	69	NO	A138	545	D. DARIS, C. ST-PIERRE	(LAVAL-QUEBEC)
SALGO	69	NO	A138	417	R. C. SALGO, H. STAUB	(ZURICH)
BERGVIST	72	NO	B39	317	KARL-ERIK BERGVIST	(UNIV STOKHOLM)
ROOF	72	LHC	5	134	B. ROOF, M. DANIEL	(MICHIGAN STATE)
PIEL	73	NO	A203	360	WILLIAM F. PIEL, JR.	(IITP)
CLARK	74	PR	D0	533	H. TIEFF, F. RAICH, J. JOHNSON, W. HERTH, SHEN	(LBL)
REINES	74	PHL	32	100	S. REINES, G. RUPP	(EDU)
					ALSO 78 PRIVATE COMM.	(PUBD)
TRETYAKOV	76	RUSP	40	10-1	TRETYAKOV, ISLJL, ACAD. SEC. USSR, DMV. I	(ITEP)
					ALSO 76 NU CONF. BACHEN	TRETYAKOV, MYASOUDOV, APALIKOV, KORYAEV, IITEP
BERGVIST	80	NEUTRINO	NO	VERTICE	K. E. BERGVIST	(ISTO)
LURIMOV	80	PL	94B	294	ANDRIYKOV, NOZIK, TRETYAKOV, KOSIK	(ITEP)
					ALSO 80 SJNP 32 154	K37 IN LURIMOV, NOVIKOV, NOZIK, TRETYAKOV IITEP
					J. J. SIMPSON	(IUGL)

Data Card Listings

For notation, see key at front of Listings.

Stable Particles

e, ν_e

e 3 ELECTRON, μ_e 1/271

1 ELECTRON MASS (MEV)
M 10.011000 0.000021 COHEN 65 PVUE 1/770
M 10.5110041 0.000011 TAYLOR USING NEW E/M 3/776
M 0.5110034 0.000014 COHEN 73 PVUE 3/776

3 ELECTRON MEAN LIFE τ BRANCHING FRACTION UNITS YSI
TEST OF CHARGE CONSERVATION
T M 12. 6211 CO MORE MDE 65 CNTR SEE NOTE 5 BELOW 6/66
T M 14. 6221 CO MORE MDE 65 CNTR E- -> NEU GAMMA 6/66
T M 15. 3E23 CO MORE STEINBERG 75 CNTR SEE NOTE 5 BELOW 2/76
T M 2. 822 CO MORE CL-68 KVALCHUK 70 CNTR SEE NOTE 5 BELOW 1/81
T M 15. 3E23 CO MORE CL-68 KVALCHUK 70 CNTR E- -> NEU GAMMA 1/81
T M SEE MOE 65 FOR DISCUSSION OF EARLIER EXPERIMENTS. 1/81
T M MOE AS LIMIT ESTABLISHED BY STEINBERG 75 TO BE 11. 8201. 1/81
T S THESE LIMITS ARE FOR ALL MODES IN WHICH DEBY PARTICLES ESCAPE 1/81
T S FROM THE DETECTOR WITHOUT DEPOSITING ENERGY. 1/81

3 ELECTRON MAGNETIC MOMENT (MEV)
M 10.011000 0.000021 COHEN 65 PVUE 1/770
M 10.5110041 0.000011 TAYLOR USING NEW E/M 3/776
M 0.5110034 0.000014 COHEN 73 PVUE 3/776
M 1.00118001 +-1241E-7 SCHUPP 61 CNTR - 8/56
M 1.1001196271 +-1271E-4 WILKINSON 63 CNTR - POSITION 8/60
M 1.10011801 +-1271E-4 RICH 60 CNTR A 8/60
M R 1.10011945471 +-1301E-4 RICH 68 CNTR - 2/71
M 1.10011963003 +-1311E-10 TAYLOR 69 PVUE 2/71
M 1.100119644 +-171E-9 WESLEY 70 CNTR 6/70
M 1.10011965711 +-1351E-10 WESLEY 71 CNTR - 2/72
M 1.1001196603 +-1121E-7 GILLILAND 72 CNTR - 2/72
M 1.100119665711 +-1351E-10 COHEN 73 PVUE 3/76
M 1.1001196679 +-1249E-6 WALLS 73 CNTR - BOLOMETRIC TECHN 1/77
M 1.100119669311 +-1201E-11 WANDYCK 77 CNTR - PRL BY WANDYCK 12/77
M V 1.00119652200 +-140E-12 WANDYCK 79 CNTR - PENNING TRAP 1/82
M V 1.00119652222 +-150E-12 WANDYCK 79 CNTR - PENNING TRAP 1/82
M WITH AN IS RECALCULATION OF WILKINSON 63 1/82
M V WANDYCK 79 CONFIRMED FINAL BY W. DEHMETL, PRIV. COMM. 1/82
M AVG 1.00119652209 +-131E-12 AVERAGE ERROR INCL. SCALE FACTOR 1.01
M AVERAGE ASSUMING EQUAL Q-FACTOR VALUES FOR E+ AND E- BY C1.

POSITION TO ELECTRON Q-FACTOR RATIO MINUS ONE. (Q+V)/-1
TEST OF CPT
M 1.10E-11 OR LESS CL-95 SERENYAN 77 CNTR M+M- ASSUMED 4/82
M 2.2E-11 6.4E-11 SCHMIDNER 81 ELEC PENNING TRAP 4/82

3 ELECTRON ELECTRIC DIPOLE MOMENT UNITS 10⁻²³ D-CM
FORBIDDEN BY BOTH T INVARIANCE AND P INVARIANCE

E/M 0.3 OR LESS CL-95 WEISSKOPF 69 MFS CESIUM 12/79
E/M 10.071 10.721 CL-90 PLYNER 70 MFS XENON 4/82
E/M 10.19 10.743 CL-90 SANDERS 75 MFS THALLIUM 4/82
E/M 18.11 11.6 VASILEV 78 12/79

REFERENCES FOR ELECTRON
SCHUPP 61 PR 121 1 A T SCHUPP, R W PIDD, M R GRAVE (MICH)
WILKINSON 63 PR 139 812 D A WILKINSON, H P GRAVE (MICH)
COHEN 65 PR 37 537 CHEN, GOUND, CHEN, AVERLON, SELIGER, CHEN (MICH)
MDE 65 PR 140 9 982 CHEN, MDE (MICH)

RICH 66 PRL 17 271 A RICH, R R CRAFF (MICH)
RICH 66 PRL 20 967 A RICH (MICH)
WEISSKOPF 68 PRL 21 1645 WEISSKOPF, CARBICO, GOULD, LEPPORETTA (BRAN)
TAYLOR 69 PR 140 9 982 PARKER, LANGERBERG (MICH)

PLYNER 70 PR 63 1620 W.A. PLYNER, G.M. SANDARS (MICH)
WESLEY 70 PRL 24 1320 J.C. WESLEY, A.RICH (MICH)
WESLEY 71 PR 44 313 J.C. WESLEY, A.RICH (MICH)
GILLILAND 72 PR 44 313 J.GILLILAND, RICH (MICH)
LAUTRUP 72 PRL 3 193 LAUTRUP, BRETERMAN, DE RAFALICEN, BUSCEL (MICH)
RICH 72 PR 44 250 A RICH, J.C. WESLEY (MICH)

COHEN 73 J. PHYS. CHEM. REF. DATA 3, P. 663, E.R. COHEN, B.N. TAYLOR
WALLS 73 PRL 31 975 F.L. WALLS, T.S. STEIN (MASH)
SANDERS 75 PR 113 475 P.C. SANDERS, W.A. STEINBERG (COFF) BR 73
STEINBERG 75 PR 112 2582 R. I. STEINBERG, K.M. FATHORHANI, M. ENHAUT (UMD)
SERENYAN 77 PR 166 102 SERENYAN, SIDOROV, SKRINSKIY (MVD)
WANDYCK 77 PRL 38 1045 SCHMIDNER, VAN DYCK, DEHMETL (MASH)
WANDYCK 78 (YND) REP. P. 371 T. KINOSHITA (CORNI)

VASILEV 78 JETP 47 243 +KOLYCHEVA (JINR)
KVALCHUK 79 JETP 27 145 +SCHMIDNER, DEHMETL (INRM)
WANDYCK 79 BULL. APS 24 575 +SCHMIDNER, DEHMETL (COFF) BR 73
ALSO IN AT. PHYS., P. P. 337 M. DEHMETL, EDS. KLEPPER, W. PLENY, NY. RIJFMASS
KINOSHITA 81 PRL 47 1473 T. KINOSHITA, W. BLINDENSTEIN (CORNI)
SCHMIDNER 81 PRL 47 1470 SCHMIDNER, VAN DYCK, DEHMETL (MASH)

v 2 ν_e MUON-1/271

NOTE IN GENERAL A MASS IS SHOWN. SEE NOTE ON NEUTRONS IN THE ELECTRON NEUTRINO SECTION ABOVE.

2 ν_e MUON MASS (MEV)
APPLIES TO ν_e 2. THE PRIMARY MASS MEASUREMENT IN ν_e MUON WOULD ALSO APPLY TO ν_e OTHER ν_e WHICH MIXES SIGNIFICANTLY IN ν_e MUON AND HAS SUFFICIENTLY SMALL MASS THAT IT CAN DECAY TO THE RESPECTIVE DECAYS. (THIS WOULD BE NEUTRALIZED ONLY FOR ν_e 3. 5. GIVEN THE ν_e MASS MASS LIMIT ABOVE.)

M 13.81 OR LESS BARNAT 56 E/M 5
M 14.01 OR LESS DUBOZIK 59 E/M 7
M 13.01 OR LESS FLEISCHER 63 PVUE 7/60
M 13.01 OR LESS ALLCOCK 65 PVUE 7/60
M 12.51 OR LESS BARDON 65 ASPR 5/71
M 12.81 OR LESS CL-90 SHAFER 65 CNTR 3/66
M 11.61 OR LESS CL-90 BODTH 67 CNTR 11/70
M 12.21 OR LESS CL-90 HYMAN 67 HECB 0. 5. 4E 11/70
B M 11.71 OR LESS CL-90 BACKENSTOSS 71 CNTR $\mu\mu\pi\pi$ -1.234-1.24 10/71
S 11.151 OR LESS CL-90 SHUM 71 CNIP $\mu\mu\pi$ -1.354-1.14 12/71
B M 11.151 OR LESS CL-90 BACKENSTOSS 73 CNIP $\mu\mu\pi\pi$ -0.294-0.90 1/73
M 10.61 OR LESS CL-90 CLARK 74 ASPR $\mu\mu\pi$ DECAY 1/74
M 10.571 OR LESS CL-90 DAUM 70 SPC $\mu\mu\pi$ -0.134-0.16 10/81
M L 10.52 OR LESS CL-90 LU 80 CUP $\mu\mu\pi\pi$ -1.024-1.10 1/76
M WE CALCULATE UPPER LIMIT AT CL-90 FROM $\mu\mu\pi$. 1/76
M BACKENSTOSS 73 REPLACES BACKENSTOSS 71 AND USES THE NEW PV. MASS. 1/73
S SHUM 71 USES UPPER OF PV. MASS VALUE AND GIVES THE PV. MASS VALUE. 1/73
M LU 80 COMBINES DATA TO PV. MASS. (MUMI) MEASUREMENT WITH NEW LU 80 1/82
M PV. MASS AND REPLACES DAUM 70. 1/82

2 ν_e 21 147U-23 MASS DIFF. (MEV)
TEST OF CPT AND A DIRAC ν_e 21
M 10.451 OR LESS CL-90 CLARK 74 ASPR $\mu\mu\pi$ DECAY 11/75

2 ν_e 21 MEAN LIFE/MASS (YEITS SEC/EV)
T B 0.13, E-31 OR MORE CL-90 BELLOTTI 76 HECB $\mu\mu\pi$ CEPN 1/78
T B 1.13, 3E-21 OR MORE CL-90 BELLOTTI 76 HECB ANTI- ν_e CEPN 1/78
T B 0.12, 2E-31 OR MORE CL-90 BARNAT 77 BR. NU. ALL 1/78
T B 0.13, 2E-31 OR MORE CL-90 BARNAT 77 BR. NU. ALL 1/78
T B 0.13, 2E-31 OR MORE CL-90 BELLOTTI 76 HECB $\mu\mu\pi$ CEPN 1/78
T B 0.13, 2E-31 OR MORE CL-90 BELLOTTI 76 HECB $\mu\mu\pi$ CEPN 1/78
T B 0.13, 2E-31 OR MORE CL-90 BARNAT 77 BR. NU. ALL 1/78
T B THE SEVERAL VALUES OF PV. MASS. (MUMI) MEASUREMENT WITH NEW LU 80 1/82
T B ANTI-ELECTRON. 1/82

2 ν_e 21 212Bi DECAY: BRANCHING FRACTION UNITS 10⁻⁴
EXPECTED TO BE ZERO FOR MASSLESS NEUTRINO
V 77 12.01 OR LESS CL-90 ALPERT 76 SPC 5500V M 1/78
V 76 14.01 OR LESS CL-90 ALPERT 76 SPC 5500V M 1/78
V 9800 10.41 OR LESS CL-90 BELLOTTI 76 SPC 1/79

2 ν_e 21 MAGNETIC MOMENT UNITS (GAUSS)
MUST VERIFY FOR MAJORITY NEUTRINO OF PURELY CHIRAL MASSLESS DIRAC NEUTRINO

M K 14.7 1710E LEM KIM 74 1/82
M K 14.7 1710E LEM KIM 74 1/82

REFERENCES FOR ν_e MUON
BARNAT 56 PR 101 378 A BARNAT, W. DUBOZIK, F. M. SMITH (LRL)
DUBOZIK 58 PR 114 376 W. F. DUBOZIK, S. GAGNE, J. VEDDER (LRL)
FLEISCHER 63 PR 139 812 G. FLEISCHER, L. M. LEIDERMAN (COLUMBIA)
ALLCOCK 65 PR 37 537 G. ALLCOCK (LIVERPOOL)
BARDON 65 PRL 14 449 R. BARDON, NORTON, PEOLIS + (EDLINGTON BRONX)

SHAFER 65 PRL 14 923 F. SHAFER, CONDE, JENNINS (LRL)
BODTH 67 PL 268 30 BODTH, JOHNSON, WILLIAMS, WARDHAL (LIVERPOOL)
HYMAN 67 PL 258 375 HYMAN, FLEISCHER, MCKENZIE + (LIVERPOOL)
BACKENSTOSS 71 PRL 38 1045 BACKENSTOSS, SCHMIDNER, KIM, GERN, KARL, MEIDI
SHUM 71 PL 37 114 E. V. SHUM, O. H. ZIEGLER (UNIV OF VIRGINIA)
BACKENSTOSS 73 PL 438 530 BACKENSTOSS, DEHMETL, WANDYCK (LIVERPOOL)

CLARK 74 PR 10 533 BELLOTTI, FERSON, JOHNSON, KERR, SHEN + (LRL)
KIM 74 PR 10 3550 J.C. KIM, S. WATSON, C. F. BROWN (LRL)
ALPERT 76 PRL 36 837 ALPERT + (LIVERPOOL)
BELLOTTI 76 PRL 16 1551 CAVALLOTTI, FERSON, HOLLIER (LRL)
GARNER 77 PRL 38 1045 GARNER, DAUVE, FERNANDEZ + (BRONX)
BLITSCH 78 PR 8133 205 BLITSCH, SHUM, KARL, M. BUCHER, E. POLAK, M. LARSON +

KALBFLEISCH 79 PRL 43 1361 KALBFLEISCH, BAGGETT, FOWLER + (FALCON) (LRL)
DAUM 79 PR 270 260 DAUM, DUBAL, EATON, FRIEDRICH, MCCULLOUGH + (LRL)
ALSD 76 PL 820 380 ALSD, TAYLOR, FRIEDRICH, MITSCHMAN, + (LRL)
DAUM, EATON, FRIEDRICH, MITSCHMAN, + (LRL)
ROD 80 PL 45 1066 +DELER, ZUBAN, MULLER, FAPPEN + (STATE COLLEGE)
FRANK 81 PR 24 2091 +BURN, GILKES, WALLS, M. B. GILKES, M. B. GILKES (MASH)

Stable Particles

T*

Data Card Listings
For notation, see key at front of Listings.

35 TAU+- (1705 J/psi)21 HEAVY LEPTON
E+- -> TAU+TAU- CROSS SECTION THRESHOLD BEHAVIOR
AND MAGNITUDE CONSISTENT WITH POINTLIKE SPIN 1/2
DIPAC PARTICLES. BRANDELK 78 RULES OUT POINTLIKE
SPIN 0 OR SPIN 1 PARTICLE. FELDMAN 78 RULES OUT
J=3/2-. KIRBY 79 ALSO RULES OUT J=INTEGER, J=3/2.

35 TAU MASS (MEV)
M P 64(1090.) (2001.) PERL 75 SMAG INCL. 14 REPL 77 2/78
B 220(1930.) (331.) BURNETT 77 PLUT ASSUMES V-A DECAY 12/77
P 144(1900.) (1001.) PERL 77 SMAG E+E= 3.0-7.0GEV ECM 12/77
A 492 1793. 30. BACIND 78 DLCD E+E= 3.1-7.4GEV ECM 2/82*
A 299 1787. 18. BARTEL 78 SPCE E+E= 3.0-6.4GEV ECM 7/79
M 1807. 20. BRANDELK 78 DASP E+E= 3.1-6.2GEV ECM 3/78
M 1787. 10. BLOCKER 80 SMAG E+E= 3.5-6.7GEV ECM 2/82*
M 1138(1890.) (1101.) BLOCKER 82 SMAG INCL. 14 BLOCKER 80 2/82*
B 0 BURNESTER 77 MASS VALUE ARE FROM ELEMENTS CONTAINING MU+ PLUS ONE
B 0 OTHER PONG, ORIGINATING FROM E+ E- -> TAU+ TAU- THE MASS 12/77
B 0 VALUES COME FROM A FIT TO THE SHAPE AND ECM DEPENDENCE OF THE 12/77
B 0 MU+- SPECTRA. ASSUMING THAT THE TAU SPIN IS 1/2 AND ITS ASSOC. 12/77
B 0 NEUTRINO MASS 46.0. 12/77
P P REAL 77 VALUE COMES FROM E+ E- -> MU+ MU- AND NO OTHER OBTAINED 12/77
P PARTICLES. ASSUMES V-A COUPLING AND ZERO MASS FOR ASSOC. NEUTRINO. 12/77
A BACIND 78 VALUE COMES FROM E+ E- -> THRESHOLD. PUBLISHED MASS 1782 2/82*
A MEV INCREASED BY 1 MEV USING THE HIGH PRECISION SPI-PRIME MASS 2/82*
A MEASUREMENT TECHNIQUE 90 TO ELIMINATE THE ABSOLUTE SPEAK ENERGY 2/82*
A CALIBRATION UNCERTAINTY. 2/82*
R BARTEL 78 FITS ENERGY DEPENDENCE OF CS FOR E+- AND MU+- EVENTS. 7/79
R MASS VALUE NOT DEPENDENT ON WHETHER V-A OR V+A DECAY ASSUMED. 7/79
M AVG 1786.2 (1101.) 3.2 AVERAGE ERROR INCLUDES SCALE FACTOR OF 1.01

35 TAU MEAN LIFE (UNITS 10^-13 SECS)
I 77 100.1 CR LESS CL-95 ALEXANDER 70 PLUT E+E= 3.0-5 GEV ECP 7/79
I 594 (23.) CR LESS CL-95 BRANIND 70 DLCD E+E= ECM=3.5-7.4GEV 7/79
I 78 (14.) CR LESS CL-95 BRANDELK 78 TASS E+E= ECM=3.0GEV 8/81M
I 102 4.6 1.9 FELDMAN 82 SMAG E+E= ECM=20 GEV 1/82*

35 TAU PARTIAL DECAY MODES
P1 TAU+- INTO MU+- NU(M) NUTAU(S) DECAY MASSES
P2 TAU+- INTO E+- NU(E) NUTAU(S) 105 0 0
P7 TAU+- INTO HADRONS+ NEUTRAL(S) 105 0 0
T8 TAU+- INTO 3 HADRONS+ NEUTRAL(S) 105 0 0
P9 TAU+- INTO NUTAU PHO(P) P(+)- 0+ 76% 13%
P10 TAU+- INTO NUTAU PI(0)PI(+)- 0+127%
P11 TAU+- INTO E+- NEUTRAL(S) 105 0 0
P12 TAU+- INTO NUTAU PI(+)- 0+ 13%
P13 TAU+- INTO NUTAU PI(0) P(+)- (INCL. PD. 923) 0+ 13% 13% 13%
P14 TAU+- INTO NUTAU PI(0) P(+)- (PD(S) (INCL. P13) 0+ 13% 13% 13%
P15 TAU+- INTO NUTAU GE. SHADRON(S)+ NEUTRAL(S) 5 26%
P17 TAU+- INTO 3 HADRONS+ NUTAU 105 0 0
P18 TAU+- INTO 3 HADRONS+ NUTAU GAMMA(S) 105 0 0
P19 TAU+- INTO HADRONS+ ENDF BHC(+)- P(+)- NEUTRAL(S) 105 0 0
P20 TAU+- INTO 3 HADRONS+ NEUTRAL(S) 105 0 0
P21 TAU+- INTO NUTAU K(K*0)P(+)- 0+ 6%
P22 TAU+- INTO NUTAU K(K*0) 0+14%
LEPTON NUMBER VIOLATING MODES.

P31 TAU+- INTO MU+- GAMMA 105+ 0
P32 TAU+- INTO E+- GAMMA 105+ 0
P33 TAU+- INTO MU+- CHARGED PARTICLES
P34 TAU+- INTO E+- CHARGED PARTICLES
P35 TAU+- INTO MU+- MU+ MU- 105+ 105+ 105
P36 TAU+- INTO E+- MU+ MU- 105+ 105+ 105
P37 TAU+- INTO MU+- E+ E- 105+ 105+ 105
P38 TAU+- INTO E+- E+ E- 105+ 105+ 105
P39 TAU+- INTO MU+- PI(0) 105+ 13%
P40 TAU+- INTO E+- PI(0) 105+ 13%
P41 TAU+- INTO MU+- K(0) 105+ 6%
P42 TAU+- INTO E+- K(0) 105+ 6%
P43 TAU+- INTO MU+- RHO(0) 105+ 6%
P44 TAU+- INTO E+- RHO(0) 105+ 6%

FITTED PARTIAL DECAY MODE BRANCHING FRACTIONS

The matrix below is derived from the error matrix for the fitted partial decay mode branching fractions, Pij, as follows. The diagonal elements are Pij^2, where Pij = sqrt(Bij) * 100, and the off-diagonal elements are the normalized correlation coefficient c_ij = Pij * c_ij / (P_i * P_j). For the definitions of the individual Pij, see the listings above, only those Pij appearing in the matrix are assumed in the fit to be nonzero and those constrained to add to 1.

Table with 9 columns (P1 to P9) and 9 rows of numerical values representing branching fractions and correlations.

35 TAU BRANCHING RATIOS
R1 TAU+- INTO MU+- NU(M) NUTAU(S)/TOTAL (P13)
R1 220 0.15 0.03 BURNETT 77 PLUT ASSUMES V-A DECAY 12/77
R1 0.15 0.040 PERL 77 SMAG E+E= 3.0-7.0GEV ECM 12/77
R1 0.12 0.027 0.37 CAVALLIUS 77 SPCE E+E= TO MU+- P+ 1/78
R1 11 0.22 0.07 BRANDELK 78 DASP E+E= 3.1-6.2GEV ECM 3/78
R1 0.15 0.050 FELDMAN 82 SMAG E+E= 3.5-6.7GEV ECM 2/82*
R1 0.178 0.027 BERGER 81 PLUT E+E= 3.0-6.2GEV ECM 1/82*
R1 AVG 0.175 0.017 AVERAGE ERROR INCLUDES SCALE FACTOR OF 1.01
R1 FIT 0.185 0.012 FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.01)
R2 TAU+- INTO E+- NU(E) NUTAU(S)/TOTAL (P21)
R2 B 459 0.180 0.013 BRANDELK 78 DLCD E+E= ECM=3.1-7.4GEV ECM 2/82*
R2 0.19 0.09 BRANDELK 80 TASS E+E= ECM=3.0GEV ECM 8/81M
R2 B BACIND 78 VALUES COME FROM FIT TO EVENTS WITH A 4X1 OTHER 1/75
R2 B NONELECTRON CHARGED PROM. 1/79
R2 AVG 0.161 0.013 AVERAGE ERROR INCLUDES SCALE FACTOR OF 1.01
R2 FIT 0.167 0.009 FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.01)
R3 TAU+- INTO E+- NU(E) NUTAU(S)/TOTAL (P12)(P13)
R3 WHERE L MEANS E OR NU. EQUALITY OF E AND NU MODES IS ASSUMED.
R3 P 105 0.19 0.09 PERL 76 SMAG 3/77
R3 B 144 0.186 0.030 PERL 77 SMAG 1/77
R3 B 21 0.224 0.055 BRANDELK 77 SMAG 1/77
R3 B 13 0.182 0.051 BRANDELK 77 DASP ASSUMES V-A DECAY 3/78
R3 B WE HAVE COMBINED STATISTICAL AND SYSTEMATIC ERRORS INDIVIDUALLY. 1/78
R3 P ASSUMES V-A COUPLING TAU+ MU+ S+ 1/2 GEV. TAU NEUTRINO MASS 50. 1/78
R3 ASSUMES V-A COUPLING TAU+ MU+ S+ 1/2 GEV. TAU NEUTRINO MASS 50. 1/78
R3 AVG 0.175 0.017 AVERAGE ERROR INCLUDES SCALE FACTOR OF 1.01
R3 FIT 0.1728 0.0074 FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.01)
R4 TAU+- INTO E+- NU(E) NUTAU(S)/MU+- NU(M) NUTAU(S) (P12)(P13)
R4 PREDICED TO BE 1 FOR SEQUENTIAL LEPTON 2 FOR PARALELIZATION.
R4 1/2 FOR PARALELIZATION. PARALELIZATION ALSO RULED OUT BY HELIX 78 1/77
R4 21 0.92 0.37 BURNESTER 77 PLUT ASSUMES V-A DECAY 1/77
R4 B 18 1.00 0.38 BRANDELK 78 DASP E+E= 3.1-6.2GEV ECM 3/78
R4 B 154 0.75 0.29 BRANDELK 82 SMAG E+E= ECM=3.5-7.4GEV ECM 2/82*
R4 B BRANDELK 78 QUOTES THE INVERSE OF THIS RATIO AS 0.28-0.32 1/78
R4 L BLOCKER 82 GIVES THE INVERSE OF THIS RATIO AS 1.33-1.10-0.26 2/82*
R4 AVG 0.86 0.17 AVERAGE ERROR INCLUDES SCALE FACTOR OF 1.01
R4 FIT 0.875 0.094 FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.01)
R5 TAU+- INTO MU+- NU(M) NUTAU(S)/E+- NU(E) NUTAU(S) (P13)(P21)
R5 B 20 0.034 0.009 BRANDELK 78 DLCD E+E= ECM=3.1-7.4GEV ECM 2/82*
R5 257 0.030 0.005 BLOCKER 82 SMAG E+E= ECM=3.5-7.4GEV ECM 2/82*
R5 B BACIND 78 VALUES COME FROM FIT TO EVENTS WITH A 4X1 OTHER 1/75
R5 B QUADRATURE ASSUMING BR(I)=0.16. WE MAY, BY 0.16 TO GET ABOVE VAL. 1/79
R5 AVG 0.034 0.004 AVERAGE ERROR INCLUDES SCALE FACTOR OF 1.01
R5 FIT 0.0249 0.0025 FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.01)
R6 TAU+- INTO SHADRON+ NEUTRAL(S)/TOTAL (P12)(P13)(P14)
R6 14 0.45 0.19 BRANDELK 77 SMAG 1/77
R6 B 20 0.45 0.19 BRANDELK 77 DASP ASSUMES V-A DECAY 3/78
R6 B (0.22) (0.14) BRANDELK 80 TASS E+E= ECM=3.0GEV ECM 8/81M
R6 B NOT INDEPENDENT OF BRANDELK 80 R1, R2 AND P21 VALUES. 8/81M
R6 AVG 0.330 0.095 AVERAGE ERROR INCLUDES SCALE FACTOR OF 1.01
R6 FIT 0.370 0.032 FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.01)
R9 TAU+- INTO E+- NEUTRAL(S)/TOTAL (P13)
R9 B BRANDELK 77 FINDS 0.074-0.06 K+ PER LVT IN E+- -> E+- BRANDELK 1/78
R9 B BRANDELK 77 FINDS 0.074-0.06 K+ PER LVT IN E+- -> E+- BRANDELK 1/78
R10 TAU+- INTO 3 HADRONS+ NEUTRAL(S)/TOTAL (P17)(P18)
R10 0.35 0.11 BRANDELK 78 DASP ASSUMES V-A DECAY 3/78
R10 35 0.24 0.06 BRANDELK 80 TASS E+E= ECM=3.0GEV ECM 8/81M
R10 AVG 0.265 0.053 AVERAGE ERROR INCLUDES SCALE FACTOR OF 1.01
R10 FIT 0.284 0.036 FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.01)
R11 TAU+- INTO ENLTAU PHO(P) PI(+)-/TOTAL (P9)
R11 78 (10-045) (0.0213) ALEXANDER 70 PLUT 64% BY WACHER 40 3/78
R11 78 (10-045) (0.0213) WACHER 80 PLUT 64% BY WACHER 40 3/78
R11 78 (10-045) (0.0213) WACHER 80 PLUT 64% BY WACHER 40 3/78
R12 TAU+- INTO ENLTAU PI(0)PI(+)-/TOTAL (P10)
R12 21 (0.101) (0.031) ALEXANDER 70 PLUT REPL. BY WACHER 40 3/78
R12 23 (0.101) (0.031) WACHER 80 PLUT 64% BY WACHER 40 3/78
R12 M ALEXANDER 70 QUOTE BR(I)=0.000-0.015. WE MAY, BY 0.16 TO GET ABOVE VAL. 1/79
R12 M ENU PHO(P) PI(+)- EVENTS ARE INU PI(+)- AND ARE INU NUTAU(S) 2/82*
R12 M -174-0.13 - 1/82*
R13 TAU+- INTO ENLTAU PHO(P) PI(+)-/TOTAL (P10)
R13 33 0.18 0.06 JAROS 78 SMAG E+E= 3.0-6.4GEV ECM 7/79
R14 TAU+- INTO ENLTAU PHO(P) PI(+)-/TOTAL (P10)
R14 J 13 0.07 0.009 JAROS 78 SMAG E+E= 3.0-6.4GEV ECM 7/79
R14 J JAROS 78 EVENTS CONSISTING WITH BEING HO PI(+)- 1/79
R15 TAU+- INTO ENLTAU GE. SHADRON(S)+ NEUTRAL(S) (P17)(P18)
R15 492 0.32 0.05 BACIND 78 DLCD E+E= ECM=3.1-7.4GEV ECM 2/82*
R15 10 0.32 0.05 BRANDELK 78 DASP ASSUMES V-A DECAY 3/78
R15 (SHADRON(S)+ NEUTRAL(S), I.E. THAT THE BRANCHING FRACTION FOR 5 OR MORE CHARGED HADRONS IS ZERO.
R15 FIT 0.32 0.030 FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.01)
R16 TAU+- INTO ENLTAU PI(+)-/TOTAL (P10)
R16 A 23 0.015 0.006 ALEXANDER 70 PLUT REPL. BY WACHER 40 3/78
R16 B 18 0.015 0.006 BRANDELK 78 DLCD E+E= ECM=3.1-7.4GEV ECM 2/82*
R16 A ALEXANDER 70 QUOTE BR(I)=0.000-0.015. WE MAY, BY 0.16 TO GET ABOVE VAL. 1/79
R16 A QUADRATURE ASSUMING BR(I)=0.16. WE MAY, BY 0.16 TO GET ABOVE VAL. 1/79
R16 B QUADRATURE ASSUMING BR(I)=0.16. WE MAY, BY 0.16 TO GET ABOVE VAL. 1/79
R16 AVG 0.015 0.002 AVERAGE ERROR INCLUDES SCALE FACTOR OF 1.01
R16 FIT 0.0172 0.0026 FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.01)
R17 TAU+- INTO ENLTAU PHO(P) PI(+)-/TOTAL (P10)
R17 (0.024) (0.005) WACHER 80 SMAG REPL. BY BLOCKER 82 12/79
R17 L3R 0.034 0.008 BLOCKER 80 SMAG E+E= ECM=3.5-7.4GEV ECM 2/82*
R17 FIT 0.0350 0.0059 FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.01)
R18 TAU+- INTO ENLTAU PHO(P) PI(+)-/MU+- NU(M) NUTAU(S) (P13)(P21)
R18 (0.024) (0.005) WACHER 80 SMAG REPL. BY BLOCKER 82 12/79
R18 L3R 0.034 0.008 BLOCKER 80 SMAG E+E= ECM=3.5-7.4GEV ECM 2/82*
R18 FIT 0.0350 0.0059 FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.01)

Stable Particles

π^+ , π^0

#2 CHAR. PION INTO E NEU (UNITS 10¹⁰+-) (P21/E11)
 R2 D 11.2471 10.0289 D1 CAPUA 84 CNTR 13/75
 R2 D 11.276 BRYMAN 75 BUVE 9/75
 R2 D BRYMAN 75 15 A RECALC. OF OCAPUA 84 EXPY USING LATEST P1 LIFETIME. 9/75
 #2 AVG 1.267 ± 0.023 AVERAGE ERROR INCLUDES SCALE FACTOR OF 1.0

#3 CHAR. PION INTO P10 E NEU (UNITS 10¹⁰+-) (P41/E11)
 R3 D 52 (11.15) 10.221 DEPOMNII 63 CNTR + 2/72
 R3 D 36 0.07 0.20 BARLETT 84 DSPK +
 R3 D 38 1.07 0.21 BACASTOW 65 CNTR +
 R3 D 1.10 0.26 BERTRAM 65 DSPK +
 R3 D 43 1.1 0.2 DUMITRESCU 65 CNTR +
 R3 332 1.00 0.08 0.10 DEPOMNII 68 CNTR + 2/68

#3 AVG 1.267 ± 0.089 AVERAGE ERROR INCLUDES SCALE FACTOR OF 1.0
 D DEPOMNII 68 STATES THAT THE RESULT OF DEPOMNII 63 IS AT LEAST 2/72
 D 10 PERCENT TOO LARGE BECAUSE OF A SYSTEMATIC ERROR IN THE P10
 D DETECTION EFFICIENCY. THIS MAY BE TRUE OF OTHER THE PREVIOUS 2/72
 D MEASUREMENTS ACCORDING TO DEPOMNII 68 AND V. SOERGEL, PRIVATE
 COMMUNICATION. 1972.

#4 CHAR. PION INTO E NEU GAMMA (UNITS 10¹⁰+-) (P53/E11)
 R4 143 13.0 0.7 DEPOMNII 63 CNTR + E+ GAMMA 6 MEV 2/82
 R4 5 226 5.0 0.7 STETZ 78 SPEC + E NOM 556 MEV/C. 12/79
 R4 5 STETZ 78 15 FOR E-GAMMA OPENING ANGLE >132 DEG. OBTAINS 5.7 WHEN 12/79
 5 USING SAME CUTOFFS AS DEPOMNII.

#5 CHAR. PION INTO E NEU (E+ E- UNITS 10¹⁰+-) (P63/E11)
 R5 13.43 CR LESS CL=90 KORENCHEN 71 DSPK + 10/71
 R5 0.48 CR LESS CL=90 KORENCHEN 76 SPEC + 1/78

REFERENCES FOR CHARGED PION

CRUKE 54 PR 96 470 K W CRUKE P N PHILLIPS (LRL)
 SHAFER 56 PR 117 1714 W M BARKAN + RICHARD W SMITH (LRL)
 CHWEE 57 NC 5 541 M W CRUKE (STANFORD HELP)
 CASTAGNO 58 PR 112 1774 C CASTAGNO L M MUCHNIK (LRL)

ANDERSON 60 PR 119 2050 M L ANDERSON I FUJII R W MILLER + (LRL)
 DEPOMNII 63 PL 7 285 ASHKEWITZ, FINEBO, LEPEN + (LRL)
 DEPOMNII 63 PL 7 480 DEPOMNII 63 INTZ, RUBIN, SOERGEL (CERN)
 DEPOMNII 63 PL 7 480 P. DEPOMNII + HEINTZ, RUBIN, SOERGEL (CERN)
 DI CAPUA 64 PP 1338 1333 BARLETT 84 DSPK +
 DI CAPUA 64 PP 1338 1333 DI CAPUA, GARLAND, DOMON, STRELFITZ (CERN)

BACASTOW 65 PR 139 8407 WGHESSEWER, WITLAND, LARSEN (LRL)
 BERTRAM 65 PR 139 B 617 BERTRAM, MEYER, CARRIGAN (LRL)
 DUMITRESCU 65 PR 20 56 DUMITRESCU, PETRICHEN, PROKOSHIN + (LRL)
 ECKHAUSEN 65 PR 19 344 ECKHAUSEN, HARRIS, SPODER (WILLIAM AND MARY)

BARDON 66 PL 28 1776 BARDON, DUFFE, COBAN, KRUEGER + (COLUMBIA)
 KJUTVIN, PROKOSHIN, RAJUSOVA, SIMONOV (UBNAI)
 INSEY 66 PR 144 1113 INSEY, LOEWENICZ, HODDORF (CORNELL)
 LOEWENICZ 66 PR 17 442 LOEWENICZ, HODDORF, HODDORF (CORNELL)

MYNEN 67 PL 259 376 ALDRICH, WHITE, DEFFNER + (LANCASHIRE UNIV)
 HODDORF 67 PL 249 596 HODDORF, LOEWENICZ, RICHMAN (CORNELL)
 SHAFER 67 PR 163 1451 ROBERT E. SHAFER (LRL)
 ALSO: 65 PRL 14 323 SHAFER, COHEN, JENKINS (LRL)

DEPOMNII 68 PR 84 1295 DEPOMNII 68 DUELS, HEINTZ, RAJUSOVA (CERN)
 PETRICHEN, SPYKALIN, WHEELER, GISEK (UBNAI)
 RUDIN 70 PL 328 1273 JIMSON + WILLIAMS, NORRALL (LIV)
 WRES 71 PR 30 1151 WGHESSEWER, GRENBERG, KENNEDY (LRL)
 ALSO 67 PR 12 1286 WRES 71 PR 30 1151 WRES 71 PR 30 1151 WRES 71 PR 30 1151 WRES 71 PR 30 1151
 ALSO 68 PRL 21 261 WRES 71 PR 30 1151 WRES 71 PR 30 1151 WRES 71 PR 30 1151
 ALSO 69 PRL 23 1267 WRES 71 PR 30 1151 WRES 71 PR 30 1151 WRES 71 PR 30 1151

BACASTOW 71 PL 36A 453 BACASTOW, DANIEL, KOCHE + (CERN, KARL HEISE)
 ALSO TO 71 515 C. VON DER MALSBURG (HEITHELBURG)
 KOCHE, KOCHE +, ZISTEN, WITELMACHNER (LIV)
 SHAFER 72 PR 143 1022 ROBERT E. SHAFER (LRL)
 BACASTOW 71 PL 418 1339 BACASTOW, DANIEL, KOCHE + (CERN, KARL HEISE)
 ALSO 75 SUBMITTED TO NP DUMITRESCU 73 SUPR 18 292 DUMITRESCU, PROKOSHIN, RAJUSOVA + (SERP)

ARMSTRONG 76 PR D11 1337 P. ARICCIOTTO (UNIV OF VICTORIA)
 BRADYDAD 76 INT 318 1150 RYANOD O'DRIVERA, DANIEL, VON EGLOY + (LIV)
 CARPER 76 PR 17 1190 R. CARPER, SUNDBERGAN (LRL)
 KOCHE 76 JEP 44 35 KOCHE, KOCHE +, ZISTEN, WITELMACHNER (LIV)
 MARUSHEV 76 JEP 23 72 MARUSHEV, W. ZEITSKY, DEBRUIN + (LIV)
 ALSO 76 PRIVATE COMM. P. SHAFER (LRL)
 A. I. SMIRNOV (LIV)

YITZ 78 PR A138 285 K. GAMBOLL, O'FLOIN, SPENCER + (MILWAUKEE)
 DAIN 78 PR 229 2602 J. TON, FROSCHE, HERSCHELMAN, MCELROY + (LIV)
 ALS 78 PR PAF 174 JON, FATHON, FROSCHE, HERSCHELMAN + (LIV)
 R4 80 PRL 45 1760 W. BELLETIN, DUBAN, W. CARPENE + (LIV)

PAPERS NOT REFERRED TO IN DATA CARDS
 CARLBERG 53 PR 91 677 CARLBERG, PETERMAN, WHITEHEAD, WILCOX (LRL)
 WERSON 62 ADP 11 1 WERSON (LIVERPOOL)
 SHAFER 62 PR 125 1022 R. SHAFER (LIV)
 CLEGG 63 PR 133 341 J. CLEGG + I. EDWARDS (LIV)
 CROOK 64 PR 133 341 J. CROOK + J. EDWARDS (LIV)
 DEPOMNII 60 NP 835 47 P. DEPOMNII (LIV)
 WILKIN 80 JPG 6 15 WILKIN (CERN)

NEUTRAL PION MASS DIFFERENCE (MEV)

D 15.371 13.91 PANDOLF 51 CNTR - 2/72
 D 4.40 0.31 CHITANSKY 54 CNTR - 2/72
 D 4.62 0.50 HADDUCK 59 CNTR - 2/72
 D 4.06 0.08 HELLMAN 59 CNTR - 2/72
 G 4.25 0.07 COFFEE 59 CNTR - 2/72
 G 4.06 0.07 SAUNDERS 60 CNTR - 2/72
 G 4.056 0.005 L'FRIO 63 CNTR - 2/72
 G 4.50 0.009 PETERMAN 63 CNTR - 2/72
 G 4.074 0.0052 WATLINS 64 CNTR - 2/72
 D 4.0643 0.0037 AVE J. (PANDOLF INCLUDES SCALE FACTOR OF 1.0)



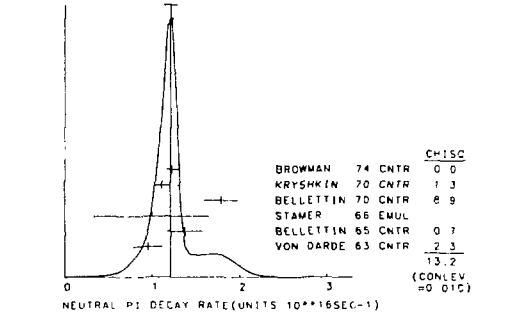
Data Card Listings

For notation, see key at front of Listings.

#9 NEUTRAL PION MEAN LIFE (UNITS 10⁻¹⁶ SEC)

T N 76 (1.91) (0.5) (0.5) GLASSER 61 EMUL 8/70
 T N 45 (2.31) (1.1) (1.0) TIETGE 62 EMUL 8/70
 T N 88 (2.81) (0.9) (0.9) KOLLER 63 EMUL 8/70
 T N 1.05 0.18 VON DARDE 63 CNTR 8/70
 T N 75 (1.71) (0.5) (0.5) SHNE 64 EMUL 8/70
 T N 61 (2.79) (0.10) BELLETIN 65 CNTR 8/70
 T N 67 (2.81) (0.10) EVANS 65 EMUL 8/70
 T N 232 1.0 0.5 STRYER 66 EMUL 8/70
 T N 0.54 0.06 BELLETIN 70 CNTR 8/70
 T N 0.9 0.08 KRYSHKIN 70 CNTR 8/70
 T N 0.82 0.04 BROWMAN 74 CNTR 8/70
 T N OLD EMULSION MEASUREMENTS NOT USED BECAUSE OF POSSIBLE SYSTEMATIC
 T N SHIFT TO LARGER MEAN LIFE VALUES.
 T N INCLUDES EVENTS OF KOLLER 63.
 T B BROWMAN GIVES P10 MEAN LIFE 0.95 ± 0.02EV. MEAN LIFE IS 0.94/0.02M.
 T AVG 0.828 ± 0.057 AVERAGE ERROR INCL. SCALE FACTOR OF 1.0
 ISEE IDEOGRAM BELOW 11/75

WEIGHTED AVERAGE = 1.207 ± 0.080
ERROR SCALE BY 1.8



NEUTRAL PION PARTIAL DECAY MODES

P1 P10 INTO 2 GAMMA 2+ 0
 P2 P10 INTO E+ E- GAMMA -4+ -4 0
 P3 P10 INTO SELECTIONS -5+ -5+ -5+ -5
 P4 P10 INTO 3 GAMMA 2+ 0 0
 P5 P10 INTO GAMMA 0+ 0+ 0+ 0
 P6 P10 INTO E+ E- -5+ -5+ 0
 P7 P10 INTO 2 NEUTRINOS 2+ 0

NEUTRAL PION BRANCHING RATIOS

R1 P10 INTO GAMMA E+ E- (1/2) GAMMA (PERCENT) (P21/E11)
 R1 1.1(16) THEORET. CALC. JOSEPH 60 QUANTUM ELECT. 9/56
 R1 207 1.176 0.15 BUDAGOV 60 HBC PI - P → P10 N 3/68
 R1 0 0.38 CR LESS CL=90 AUERBACH 60 CNTR 8/69
 R1 1.25 0.04 SCHAOT 61 SPEC PI - P → P10 N 1/82
 R1 S GAMMA VALUE USES PANOFKY RATIO = 1.62
 R1 AVG 1.213 ± 0.030 AVERAGE ERROR INCLUDES SCALE FACTOR OF 1.0

NEUTRAL PION TOTAL (UNITS 10¹⁰+-) (P4)

R2 D 0 (4.9) CR LESS CL=90 DUCLOS 64 CNTR 6/66
 R2 D (4.9) CR LESS CL=90 HUTIN 65 CNTR 3/68
 R2 D THESE EXPTS. GIVE BRIGAMMA/ZGAMMA = 0.10 ± 0.06
 R2 D (1.5) CR LESS CL=90 AUERBACH 64 CNTR 1/70
 R2 D 0.38 CR LESS CL=90 AUERBACH 60 CNTR 8/69

R3 P10 INTO (E+ E-)/2 GAMMA (UNITS 10¹⁰+-) (P17/E11)
 R3 N 144 13.10 (0.30) SAUNDERS 62 HBC SEE NOTE A P12 6/66
 R3 N 3.28 THEORET. CALC. MIYAZAKI 73 QUANTUM ELECT. 2/76
 R3 N ABOVE VALUE USES PANOFKY RATIO = 1.62

NEUTRAL PION (UNITS 10¹⁰+-) (P5)

R4 A 0 (6.4) CR LESS CL=90 ABRAMS 73 SPEC 8/73
 R4 A ABRAMS 73 GIVES BRIGAMMA/ZGAMMA = 0.10 ± 0.05
 R4 D (3.8) CR LESS CL=90 AUERBACH 70 CNTR 2/70
 R4 D 0.38 CR LESS CL=90 AUERBACH 60 CNTR 8/69

NEUTRAL PION ELECTROMAGNETIC FORM FACTOR

THE AMPLITUDE FOR THE PROCESS P10 → E+ E- GAMMA CONTAINS A FORM FACTOR (GAMMA → P2) AT THE (P10 GAMMA GAMMA) VERTEX WHERE GAMMA IS MASSIVE/MASSLESS THE PARAMETER B IN THE LINEAR EXPANSION (GAMMA → P2) = 1 + A + B Q^2 IS LISTED BELOW.

A LINEAR COEFFICIENT OF P10 ELECTROMAGNETIC FORM FACTOR
 A (-0.15) (0.10) COBAN 61 HBC NO RAD. CORR. 2/60
 A 2071 (1.0-24) (0.16) SAUNDERS 61 HBC NO RAD. CORR. 2/60
 A 300 (1.0-11) (0.11) DEVONS 60 DSPK NO RAD. CORR. 2/60
 A F 300 ± 0.10 (0.03) FISCHER 78 SPEC RAD. CORR. 2/80
 A F < 0.00 STATISTICAL ONLY. RESULT WITHOUT RAD. CORR. = 0.05 ± 0.03. 2/80

Data Card Listings

For notation, see key at front of Listings.

Stable Particles

π^0, η

REFERENCES FOR NEUTRAL PION

PANDOSKY 51 PR 81 365
 CHINDUSKY 54 PR 93 586
 CASSELS 50 PPS 74 92
 HADDOCK 50 PRL 3 478
 HELLMAN 50 NC 14 807

BUDAGOV 60 JETP 11 755
 JOSEPH 60 MC 10 987
 SARIOS 60 MC 18 154
 GLASSER 61 PR 123 1314
 KOBRAK 61 MC 20 1115
 SARIOS 61 PR 121 275
 SARIOS 62 PR 126 1844
 TIETGE 62 PR 127 1324

CZJARR 63 PR 130 341
 KOLLER 63 NC 27 1405

ALSO SEE STAMER
 PETERKUMH 63 SIENA CONF 70
 VON DARD 63 PL 4 201

SHNE 64 PR 1309 1839
 BELLETTI 65 MC 40 1139
 DUKLOS 65 PL 19 253
 EVANS 65 PR 136 982
 KUTIN 65 JETP 127 243

STAMER 66 PR 151 1108
 HASTLEYS 66 PL 23 281
 DEBONS 65 P 184 1356
 BELLETTI 70 NC 644 243
 KRISHNEN 70 JETP 30 1037

ADAMS 73 PL 658 66
 MIFRAJKE 73 PR 08 2091
 BUDMAN 74 PRL 33 1400
 DAVIES 74 NC 244 324

AUERBACH 78 PRL 41 275
 AUERBACH 78 PRL 788 355
 FISCHER 78 PL 738 359
 FISCHER 78 PL 738 364

AUERBACH 80 PL 908 217
 HIGHLAND 80 PRL 64 628
 HERZIG 81 PL 1008 147
 SCHARDT 81 PRD 23 639

W. C. H. PANDOSKY, R. L. AARNDT, J. MADLEY (LRL)
 W. CHINDUSKY, J. STEINBERGER (COLUMBIA)
 CASSELS, J. JONES, MURPHY, O. NEILL (ILVERPOOL)
 HADDOCK, A. BASHIAN, C. HOWE, CZJARR (LRL)
 HELLMAN, M. MODELKOPF, Y. KAGATA, Z. AVATY (CERN)

BUDAGOV, VIKTOR, DEMELEPOV, ERNDLOV + (JINR)
 D. M. JOSEPH (EPJ)
 N. P. SARIOS (COLUMBIA)
 W. G. GLASSER, M. SEEMAN, B. STILLER (LNL)
 H. KOBRAK (EPJ)
 N. P. SARIOS (COLUMBIA+BNL)
 SARIOS, P. LANDAU, PRODELL + (COLUMBIA+BNL)
 J. TIETGE, M. PUESCHEL (MAX PLANCK INST)

JOHN R. CZJARR (LRL)
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 W. VON DARD, O. DEKKERS, M. ERDOP, VAN PUTTEN (CERN)

M. SHNE, F. M. SMITH, M. H. BARRAS (LRL)
 BELLETTI, B. BERNARDI, BRACCINI + PISA (PISA)
 DUKLOS, F. RYKATZ, HEINZEL + (CERN+MIDELBERG)
 O. A. EVANS (OXFORD)
 KUTIN, V. PETERKUMH, PRODKOSHIN (JINR)

STAMER, TAYLOR, KOLLER, MUEYTER + (STEVENS)
 VASILEVSKY, V. SHAWARD, J. DIMA, J. SEV + (DUBNA)
 M. MENEMEN, N. SIEGHE-SABAT, D. CARINA (COLUMBIA)
 BELLETTI, B. BERNARDI, LUBELSEV + (PISA+BOVNI)
 KRYSHNEN, I. G. USOV (TMSX POLYTECH. INST.)

G. CARROLL, M. CYCILLI, MICHAEL, MCKEYTT + (BNL)
 T. MIFRAJKE, E. TERASUD (TOSK)
 O. WEINER, G. TITELMAN, HANSDON (CONN+BNL)
 O. GUY, ZIR (BIRMINGHAM+SNAP)

AUERBACH, HIGHLAND, JOHNSON, + (TEMP+LANS)
 AUERBACH, HIGHLAND, JOHNSON, + (TEMP+LANS)
 FISTERMANN, GU. SAN, MERRAD, MOORE (GEVA+SNL)
 FISTERMANN, GU. SAN, MERRAD, MOORE (GEVA+SNL)

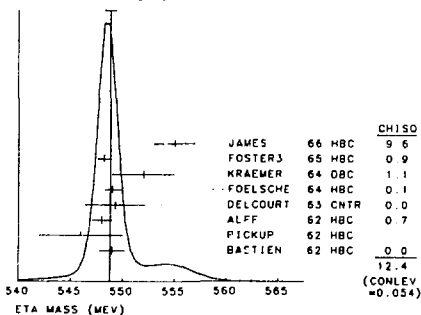
M. A. J. HIGHLAND, MCFARLANE, MACKER (TEMP+LANS)
 AUERBACH, M. A. J. HIGHLAND, MCFARLANE, MACKER (TEMP+LANS)
 P. HERZIG, G. H. HOFFMANN (LNL)
 F. FRANK, HOFFMANN, HESCHKE, MUDR + (ARZS+LANL)

14 ETA MASS (MEV)

53 549.0 1.2 BASTIEN 62 HBC
 35 546.0 4.0 PICKUP 62 HBC
 91 548.0 1.0 ALFF 62 HBC
 54 547.0 2.4 DELCOURT 63 CNTR
 148 549.0 0.7 FOELSCHKE 64 HBC
 325 552.0 3.0 KRAEMER 64 OBC
 548.0 2.0 FOSTERJ 65 HBC
 250 555.0 2.0 JAMES 66 CNTR

AVERAGE ERROR INCLUDES SCALE FACTOR OF 1.44
 (SEE TABLE BELOW)

WEIGHTED AVERAGE = 548.82 ± 0.56
 ERROR SCALED BY 1.4



14 ETA WIDTH

ETA WIDTH DETERMINED FROM MASS SPECTRUM (UNITS KEV)

91 (10.0) OR LESS ALFF 62 HBC
 140 (10.0) OR LESS FOELSCHKE 64 HBC
 31 (12.0) OR LESS JAMES 66 HBC
 14.0 (3) OR LESS DELCOURT 63 CNTR
 10.9 (1) OR LESS CL-95 JAMES 66 CNTR

ETA WIDTH DETERMINED FROM DECAY RATE (UNITS KEV)

THIS IS THE PARTIAL DECAY RATE (W) FOR THE MODE (ETA INTO ZGAMMA)
 DIVIDED BY THE FITTED BRANCHING FRACTION (P) FOR THIS MODE.

W FIT 0.83 0.12 FROM FIT

14 ETA PARTIAL DECAY MODES

Mode	Decay Modes	Decay Rates
P1	ETA INTO ZGAMMA	0+ 0
P2	ETA INTO 3PI0	134+ 134+ 134
P3	ETA INTO PI+ PI- PI0	139+ 139+ 134
PA	ETA INTO PI+ PI- GAMMA	139+ 139+ 0
P5	ETA INTO E+ E- PI0	134+ 134+ 0
P6	ETA INTO E+ E- PI+ PI-	139+ 139+ 134+ 0
P7	ETA INTO PI0 ZGAMMA	134+ 0+ 0
PA	ETA INTO E+ E- GAMMA	134+ 134+ 0
P10	ETA INTO 2PI0 GAMMA (VIOLATES C)	139+ 139+ 134+ 0
P11	ETA INTO PI+ PI- ZGAMMA	139+ 139+ 0+ 0
P12	ETA INTO MU+ MU- GAMMA	105+ 105+
P13	ETA INTO MU+ MU- PI0	105+ 105+ 0
P15	ETA INTO PI+ PI-	139+ 139+
P16	ETA INTO E+ E-	134+ 134+
P17	ETA INTO MU+ MU- PI0 GAMMA	105+ 105+ 134+ 0

FITTED PARTIAL DECAY MODE BRANCHING FRACTIONS

The matrix below is derived from the error matrix for the fitted partial decay mode branching fractions, P_i , as follows: The diagonal elements are $P_i \delta P_i$, where $\delta P_i = \sqrt{(\delta P_i)^2}$, while the off-diagonal elements are the normalized correlation coefficients $(\delta P_i \delta P_j) / (P_i P_j)$. For the definitions of the individual P_i , see the listings above; only those P_i appearing in the matrix are assumed in the fit to be nonzero and are thus constrained to add to 1.

P 1	P 2	P 3	P 4	P 5	
P 1	-.3509+-.0074				
P 2	-.6894	.3184+-.0080			
P 3	-.4922	-.3265	-.7369+-.0054		
P 4	-.3760	-.2549	-.8016	-.0491+-.0013	
P 5	-.0520	-.0461	-.0547	-.0540	-.0050+-.0012

FITTED PARTIAL DECAY MODE RATES

The matrix below is the branching fraction matrix above, transformed into rate space, i.e., $G_i = \Gamma_i / \Gamma_{total}$, in appropriate units. In analogy to the matrix above, the diagonal elements are $G_i \delta G_i$, where $\delta G_i = \sqrt{(\delta G_i)^2}$, while the off-diagonal elements are the normalized correlation coefficients $(\delta G_i \delta G_j) / (G_i G_j)$. Note that, because of the error in Γ_{total} , the errors and correlations here are not directly derivative from those above.

G 1	G 2	G 3	G 4	G 5	
G 1	.3240+-.0460				
G 2	-.9583	-.2647+-.0393			
G 3	-.9690	-.9669	-.1965+-.0288		
G 4	-.9694	-.9507	-.9988	-.0408+-.0003	
G 5	-.5151	-.5119	-.5151	-.5119	-.0094+-.0011

14 ETA DECAY RATES

W1 ETA INTO ZGAMMA (UNITS KEV) (G1)
 W1 B 11.001 10.222 BERNARDI 67 CNTR PRINAKOFF EFFECT 11/75
 W1 C 0.324 0.046 BROWN 74 CNTR PRINAKOFF EFFECT 7/74
 W1 D BERNARDI 67 GIVES W1(1.21)+.26 KEV ASSUMING THAT W1/TOTAL=0.314. 11/75
 W1 E BERNARDI PRIVATE COMMUNICATION GIVES MORE GENERAL RESULT AS 11/75
 W1 F W1/W1(TOTAL)=.380+-.083. HE EVALUATE THIS USING W1/TOTAL=.38+-.01. 11/75
 W1 G NOT INCLUDED IN AVERAGE BECAUSE THE UNCERTAINTY RESULTING FROM THE 2/76
 W1 H SEPARATION OF THE COULOMB AND NUCLEAR AMPLITUDES WAS APPARENTLY 2/76
 W1 I BEEN UNDERESTIMATED. 2/76
 W1 J FIT 0.324 0.046 FROM FIT ERROR INCLUDES SCALE FACTOR OF 1.01

14 ETA BRANCHING RATIOS

R1 ETA INTO NEUTRAL/CHARGED (P1+P2)/(P3+P4+P5)
 R1 N 1.0 (2.5) 1.0 PICKUP 62 HBC
 R1 M 53 (3.20) (1.26) BASTIEN 62 HBC
 R1 N JAMES 66 HBC
 R1 N THESE EXPERIMENTS HAVE NOT BEEN USED IN COMPUTING THE AVERAGES
 R1 M AS THEY WERE UNABLE TO SEPARATE CERTAIN PARTIAL MODES (31 AND 64)
 R1 N FROM EACH OTHER. THE REPORTED VALUES THIS PROBABLY CONTAIN
 R1 M SOME (UNKNOWN) FRACTION OF MODE (43).
 R1 J 2.64 0.23 BALTAY 57 DBC
 R1 FIT 2.436 0.076 FROM FIT ERROR INCLUDES SCALE FACTOR OF 1.01

R2 ETA INTO ZGAMMA/CHARGED (P11)/(P3+P4+P5)
 R2 0.99 0.48 CRAMPDORF 63 HBC
 R2 75 1.51 0.98 BENDALL 74 OBP
 R2 AVG 1.10 0.43 AVERAGE ERROR INCLUDES SCALE FACTOR OF 1.01
 R2 FIT 1.362 0.090 FROM FIT ERROR INCLUDES SCALE FACTOR OF 1.01

Note on $\eta \rightarrow \pi^0 \gamma \gamma$

(by A. V. Barnes, Lawrence Berkeley Laboratory)

For several years the measurements of the branching ratio $(\eta \rightarrow \pi^0 \gamma \gamma) / (\eta \rightarrow \text{neutrals})$ have been in disagreement. The recent upper limit measurement reported in DAVIDOV 81 is far below the previous positive results. The earlier

Stable Particles

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Data Card Listings

For notation, see key at front of Listings.

positive results are each suspect. The earliest report (DIGIUGNO 66) was a photon spectrum measurement. It depends on a Monte Carlo to fit the observed spectrum with that expected from a combination of photons from the decays $\eta \rightarrow \gamma\gamma$, $\eta \rightarrow \pi^0\gamma\gamma$, and $\eta \rightarrow \pi^0\pi^0\pi^0$. Their result is of course sensitive to the assumptions of this Monte Carlo. FELDMAN 67 is an optical spark chamber experiment. The scanning efficiency for the 5- and 6-photon events is hard to measure and critical to their result. STRUGALSK 71 does not address the problem of contamination from the $\eta \rightarrow \pi^0\pi^0\pi^0$ decay mode. Assuming that the $\pi^0\gamma\gamma$ mode is absent and that the decay ratio $(\eta \rightarrow \gamma\gamma)/(\eta \rightarrow \pi^0\pi^0\pi^0)$ is approximately 0.9 implies there are ≈ 700 $3\pi^0$ events in the $\leq 4\gamma$ sample of their data. The 4 ± 10 $\eta \rightarrow \pi^0\gamma\gamma$ decays observed could easily be misidentified $\eta \rightarrow \pi^0\pi^0\pi^0$ decays. DAVYDOV 81 accounts for the $\eta \rightarrow \pi^0\pi^0\pi^0$ decays properly and is much more sensitive than previous measurements.

The $\pi^0\gamma\gamma$ branching fraction is now assumed to be zero in our branching ratio fit. As a result, the fitted $\gamma\gamma$ and $3\pi^0$ branching fractions have increased by 1.1% and 1.9%, respectively.

R3	ETA INTO EPID ZGAMMA/NEUTRALS	(P7)(P1+P2)		
R3	OTHER RESULTS ARE IN SECTIONS P14, P22/P3/DAVYDOV 81, AND R26.			
R3	10.351 10.027	DIGIUGNO 66 CTR INVERSE BR REPORTED	7/66	
R3	10.071 10.101	GRUNHAUS 66 DSXP	8/67	
R3	10.221 10.004	MINISTROT 67 DSXP	11/67	
R3	10.284 10.051	FELDMAN 67 DSXP	8/67	
R3	10.021 10.019	BUTTRAM 70 DSXP	12/70	
R3	10.221 10.002	10.241/DSXP	7/66	
R3	10.071 10.155	DAVYDOV 70 DSXP	5/70	
R3	10.011 10.047	SCHWITT 70 DSXP	12/70	
R3	10.111 10.031	STRUGALSK 71 HRC	5/71	
R3	10.041 10.155	AROSIMOV 80 HRC	10/81	
R3	AVG	0.2074 0.2053	AVERAGE ERROR INCLUDES SCALE FACTOR OF 1.11	
R3	FIT	0.2074 0.2053	FROM FIT ERROR INCLUDES SCALE FACTOR OF 1.11	
R6	ETA INTO 3PI0/ZGAMMA	(P7)(P1)		
R6	10.303 10.036	CHRISTIAN 62 HRC	7/66	
R6	11.251 10.391	BAGCI 65 CTR INVERSE BR REPORTED	7/66	
R6	0.88 0.16	BALTAY 67 HRC	11/67	
R6	1.1 0.21	TINCI 67 CSXP	1/68	
R6	0.91 0.14	COY 70 HPE	6/70	
R6	3.75 0.70	JEVONS 70 ISXP	12/70	
R6	AVG	0.842 0.065	AVERAGE ERROR INCLUDES SCALE FACTOR OF 1.01	
R6	FIT	0.816 0.035	FROM FIT ERROR INCLUDES SCALE FACTOR OF 1.01	
R7	ETA INTO ZGAMMA/PI- PI- PI0	(P1)(P2)		
R7	1.1 0.20	FOSTER 65 HRC	7/66	
R7	4.01 1.72	BAGLIN 69 HRC	7/69	
R7	AVG	1.609 0.21	AVERAGE ERROR INCLUDES SCALE FACTOR OF 1.01	
R7	FIT	1.649 0.061	FROM FIT ERROR INCLUDES SCALE FACTOR OF 1.01	
R8	ETA INTO NEUTRAL/PI+ PI- PI0	(P1+P2)(P3)		
R8	3.6 0.8	KRAEHR 64 DBC	7/66	
R8	3.6 0.8	PARLI 66 HRC	8/66	
R8	2.89 0.56	ALFF-STFI 66 HRC	9/66	
R8	244 3.6	0.6	FLATTEZ 67 HRC	1/68
R8	29 2.4	1.1	AGUIE 68 72 HRC	11/72
R8	2 8.3	0.89	BLUMENWORT 72 HRC	11/72
R8	7A 2.56	1.89	KEWALL 76 CSXP	12/75
R8	BRP	INCREASED FROM PUBLISHED VALUE 0.4 BY BLODOWORTH, PRIV. COMM.	1/73	
R8	AVG	3.26 0.78	AVERAGE ERROR INCLUDES SCALE FACTOR OF 1.01	
R8	FIT	2.993 0.095	FROM FIT ERROR INCLUDES SCALE FACTOR OF 1.01	

R9	ETA INTO EPID ZGAMMA/NEUTRALS	(P7)(P1+P2)		
R9	SINGLE PHOTON PROCESS FORBIDDEN BY C-PARITY			
R9	(110.) CR LESS	PRICE 65 HRC		
R9	0.177 CR LESS	FOSTER 65 HRC		
R9	0.423 CR LESS CL-90	BAGLIN 67 HRC	8/67	
R9	0.116 CR LESS CL-90	BILLING 67 HRC	11/67	
R9	1.4 CR LESS CL-90	JANEI 76 DSXP	12/75	
R10	ETA INTO EPID ZGAMMA/NEUTRALS	(P7)		
R10	CR LESS	KITTSER 65 HRC		
R11	ETA INTO EPID ZGAMMA/NEUTRALS	(P7)(P4)		
R11	1 0.026	0.026	GROSSMAN 66 HRC	6/66
R12	ETA INTO ZGAMMA/NEUTRALS	(P1)(P1+P2)		
R12	10.416 10.044	DIGIUGNO 66 CTR	ERRP DOUBLED	6/66
R12	10.441 10.071	GRUNHAUS 66 DSXP		8/67
R12	10.579 10.052	FELDMAN 67 DSXP		8/67
R12	10.393 10.061	JONES 66 CTR		7/67
R12	THIS RESULT FROM COMBINING CROSS SECTIONS FROM TWO DIFFERENT EXPER.			
R12	0.95	0.033	MINISTROT 67 DSXP	11/67
R12	0.535	0.018	BUTTRAM 70 DSXP	12/70
R12	10.571 10.091	STRUGALSK 71 HRC		5/71
R12	0.60	0.16	MEDINA 76 SPK	12/75
R12	88 0.52	0.09	AROSIMOV 80 HRC	10/81
R12	AVG	0.547 0.016	AVERAGE ERROR INCLUDES SCALE FACTOR OF 1.11	
R12	FIT	0.551 0.011	FROM FIT ERROR INCLUDES SCALE FACTOR OF 1.11	
R13	ETA INTO 3PI0/NEUTRALS	(P7)(P1+P2)		
R13	10.291 10.054	DIGIUGNO 66 CTR	ERRR DOUBLED	6/66
R13	10.293 10.101	GRUNHAUS 66 CSXP		8/67
R13	10.171 10.031	FELDMAN 67 CSXP		8/67
R13	10.411 10.033	MINISTROT 67 CSXP	NOT INDEP. OF P12	11/67
R13	10.294 10.029	BUTTRAM 70 CSXP		12/70
R13	10.321 10.091	STRUGALSK 71 HRC		7/71
R13	75 3.44	0.88	AROSIMOV 80 HRC	10/81
R13	AVG	3.479 0.209	AVERAGE ERROR INCLUDES SCALE FACTOR OF 1.01	
R13	FIT	3.440 0.011	FROM FIT ERROR INCLUDES SCALE FACTOR OF 1.01	
R14	ETA INTO EPID ZGAMMA/ZGAMMA	(P7)(P1)		
R14	10.53 CR LESS CL-90	HAMEL 66 SPK		7/66
R14	10.28 CR LESS CL-90	BALTAY 67 DBC		11/67
R15	ETA INTO EPID ZGAMMA/NEUTRALS	(P7)		
R15	SINGLE PHOTON PROCESS FORBIDDEN BY C-PARITY			
R15	13.71 CR LESS	KITTSER 65 HRC		6/66
R15	13.094 CR LESS CL-90	BAGLIN 68 DBC		6/68
R15	13.016 CR LESS CL-90	MARTINOV 76 HRC		6/77
R17	ETA INTO EPID ZGAMMA/NEUTRALS	(P7)(P1+P2)		
R17	17.01 CR LESS	FLATTEZ 67 HRC		8/67
R17	13.91 CR LESS	PRICE 67 HRC		8/67
R17	11.24 CR LESS CL-90	BALTAY 67 DBC		11/67
R17	11.71 CR LESS CL-90	AROSIMOV 80 HRC		10/81
R17	0 2.28 CR LESS CL-90	INABE 73 ASP		6/73
R18	ETA INTO EPID ZGAMMA/PI+ PI- PI0	(P1)(P2)		
R18	0.009 CR LESS	PRICE 67 HRC		8/67
R18	13.036 CR LESS CL-90	BALTAY 67 DBC		11/67
R19	ETA INTO 3PI0/PI+ PI- PI0	(P7)(P1)		
R19	2.0 1.0	CRAWFORD 67 HRC		7/66
R19	3.0 1.0	FELSCH 66 HRC		7/66
R19	2.0 0.74	FOSTER 65 HRC		7/66
R19	1.4 0.4	4.6/12	67 HRC	8/67
R19	1.47 0.20	3.17	MULLER 68 HRC	7/68
R19	1.1 0.1	7.74	BAGLIN 69 HRC	7/69
R19	AVG	1.24 0.14	AVERAGE ERROR INCLUDES SCALE FACTOR OF 1.11	
R19	FIT	1.344 0.011	FROM FIT ERROR INCLUDES SCALE FACTOR OF 1.11	
WEIGHTED AVERAGE = 1.28 ± 0.14 ERRR SCALED BY 1.3				
R21	ETA INTO NEUTRAL/OTAL	(P1+P2)		
R21	3.79 0.08	MINISTROT 67 DSXP		11/67
R21	16 0.75 0.08	BAGLIN 71 CTR HP SPECTROMETER		8/71
R21	AVG	3.79 0.08	AVERAGE ERROR INCLUDES SCALE FACTOR OF 1.01	
R21	FIT	3.709 0.066	FROM FIT ERROR INCLUDES SCALE FACTOR OF 1.01	
R22	ETA INTO EPID ZGAMMA/TOTAL	(P7)		
R22	13.17 CR LESS CL-90	JACQUET 69 HRC		6/70
R22	0 2.003 CR LESS CL-90	DAVYDOV 81 CTR PI- PI- PI0		12/82
R23	ETA INTO NEUTRAL/TOTAL	(P7)		
R23	0 12.1 CR LESS CL-90	BERMAN 68 HRC		6/68
R23	7 2.6 0.21	3/MEYER 80 SPK PI- PI- PI0		9/81

Stable Particles

K⁺

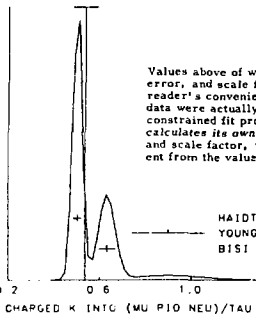
Data Card Listings

For notation, see key at front of Listings.

R18 CHAR. K INTO IPI 2P10/TAU (P41/1P3)
 R18 2027 0.303 0.009 BISI 65 HMLC + 8/66
 R18 17 0.303 0.009 YOUNG 65 EMUL + 8/66
 R18
 R18 AVG 0.3037 0.0090 AVERAGE ERROR INCLUDES SCALE FACTOR OF 1.03
 R18 FIT 0.3100 0.2080 FROM FIT ERROR INCLUDES SCALE FACTOR OF 1.03

R19 CHAR. K INTO IPI 2P10/TAU (P51/1P3)
 R19 2845 0.632 0.235 BISI 1 65 HMLC + 8/66
 R19 38 0.90 0.16 YOUNG 65 EMUL + 8/66
 R19 H 1505 10.5101 (0.037) EICHEN 68 HMLC + 11/68
 R19 H 1505 7.503 0.019 HAIDT 71 HMLC + 12/70
 R19 H HAIDT 71 IS A REANALYSIS OF EICHTEN 68.
 R19
 R19 AVG 0.536 0.054 AVERAGE ERROR INCLUDES SCALE FACTOR OF 3.21
 R19 FIT 0.572 0.016 FROM FIT ERROR INCLUDES SCALE FACTOR OF 1.03
 (SEE IDEOGRAM BELOW)

WEIGHTED AVERAGE = 0.536 ± 0.054
 ERROR SCALED BY 3.2



Values above of weighted average, error, and scale factor are for the reader's convenience only. The data were actually processed by a constrained fit program, which calculates its own values of \bar{X} , S , and scale factor, which are different from the values shown here.

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3.1

7.5

10.5

(CONLEV = 0.001)

R20 CHAR. K INTO IPI 2P10/TAU (P61/1P3)
 R20 219 0.90 0.06 BODRINI 64 HRC + 8/66
 R20 17 3.90 0.16 YOUNG 65 EMUL + 8/66
 R20 R 1505 10.5101 (0.037) EICHEN 68 HMLC + 11/68
 R20 H 4385 (12.844) (0.021) EICHEN 71 HMLC + 11/68
 R20 H 4385 0.850 0.019 HAIDT 71 HMLC + 12/70
 R20 2427 1.858 0.240 BRAUN 75 HMLC + 12/75
 R20 H HAIDT 71 IS A REANALYSIS OF EICHTEN 68.
 R20
 R20 AVG 0.858 0.018 AVERAGE ERROR INCLUDES SCALE FACTOR OF 1.03
 R20 FIT 0.8633 0.7098 FROM FIT ERROR INCLUDES SCALE FACTOR OF 1.13

R21 K+ INTO IPI P1-1+ NEU/TAU UNITS 100M-1 (P71/1P3)
 R21 44 5.7 1.5 BTJGE 65 FRC + 8/66
 R21 240 5.93 0.63 EL 69 HMLC + 11/68
 R21 500 7.36 0.98 ROUDOUVIN 71 ASPK + 12/71
 R21 106 7.0 0.9 SCHWITZE 71 HRC + 9/77
 R21 308 7.21 0.32 ROSSLEIT 71 SPEC + 12/71
 R21
 R21 AVG 6.59 0.26 AVERAGE ERROR INCLUDES SCALE FACTOR OF 1.03

R22 K+ INTO IPI P1-1+ MU NEU/TAU UNITS 100M-1 (P91/1P3)
 R22 1 (2.51) APPROD GHEINER 64 EMUL + 8/66
 R22 7 2.57 1.55 BISE 67 DBC + 11/67

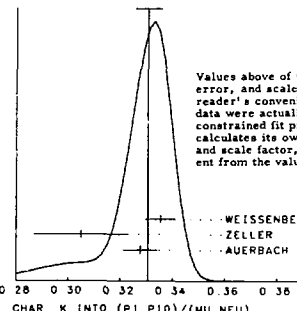
R23 CHAR. K INTO IPI 2P10/TAU (P11/1P3)
 R23 1670 5.85 0.21 CFSTER 66 DSPK + 8/67
 R23 5110 6.16 0.22 ESCHMUTZ 68 DSPK + 3/68
 R23 9 9.92 0.65 WEISSBERG 76 SPEC + 1/78
 R23 W VALUE CALCULATED FROM WEISSBERG 76 #53. KAUZ-APRIL VALUES
 R23 W TEST ESTIMATE DEPENDENCE ON OUR 1974 YAU AND TAU-PIE FRACTIONS.
 R23
 R23 AV 6.01 0.15 AVERAGE ERROR INCLUDES SCALE FACTOR OF 1.03
 R23 FIT 5.649 0.067 FROM FIT ERROR INCLUDES SCALE FACTOR OF 1.13

R24 CHAR. K INTO IPI 2P10/TAU (P12/1P1)
 R24 4517 0.126 0.003 AUERBACH 67 DSPK + 10/74
 R24 1020 0.305 0.018 ZELLER 69 ASPK + 10/76
 R24 W 25K (0.328) (0.005) WEISSBERG 74 SPEC + 7/77
 R24 C 424 0.048 0.003 GARLAND 68 DSPK + 1/78
 R24 A AUERBACH 67 CHANGED FROM .3275+0.005. SEE COMMENT WITH RATIO R26.
 R24 W WEISSBERG 74 WEISSBERG 74.
 R24
 R24 AVG 0.3307 0.0051 AVERAGE ERROR INCLUDES SCALE FACTOR OF 1.21
 R24 FIT 0.3377 0.0017 FROM FIT ERROR INCLUDES SCALE FACTOR OF 1.01
 (SEE IDEOGRAM BELOW)

R25 CHAR. K INTO IPI 2P10/TAU (P61/1P1)
 R25 A 295 0.0791 0.0054 AUERBACH 67 DSPK + 1/77
 R25 460 0.0775 0.0053 BOTTIPI 68 ASPK + 5/68
 R25 461 0.068 0.0037 GARLAND 68 DSPK + 1/78
 R25 350 0.069 0.004 ZELLER 69 ASPK + 10/76
 R25 A AUERBACH 67 CHANGED FROM .0797+0.006. SEE COMMENT WITH RATIO R26.
 R25 B VALUE -1.2784+0.0025 GIVEN IN AUERBACH 67 IS AN AVERAGE OF
 R25 A AUERBACH 67 225 AND CESTER 66 223.
 R25
 R25 AVG 0.0752 0.0024 AVERAGE ERROR INCLUDES SCALE FACTOR OF 1.01
 R25 FIT 0.07597 0.00091 FROM FIT ERROR INCLUDES SCALE FACTOR OF 1.11

R26 CHAR. K INTO IPI 2P10/TAU (P51/1P1)
 R26 A 307 0.0486 0.0040 AUERBACH 67 DSPK + 1/77
 R26 C 424 0.048 0.0037 GARLAND 68 DSPK + 1/78
 R26 240 0.046 0.004 ZELLER 69 ASPK + 10/76
 R26 A AUERBACH 67 CHANGED FROM .0602+0.006. BY ERROR WHICH BRINGS THE
 R26 A MU-SPECTRUM CALCULATION INTO AGREEMENT WITH GARLAND TO APPENDIX S.
 R26 C GARLAND 68 CHANGED FROM .055+0.004 IN AGREEMENT WITH MU-SPECTRUM
 R26 D CALCULATED BY GARLAND TO APPENDIX B. L.C. PONDROM, PRIC, COM-1731
 R26
 R26 AVG 0.0468 0.0027 AVERAGE ERROR INCLUDES SCALE FACTOR OF 1.01
 R26 FIT 0.0501 0.0014 FROM FIT ERROR INCLUDES SCALE FACTOR OF 1.71

WEIGHTED AVERAGE = 0.3307 ± 0.0051
 ERROR SCALED BY 1.2



Values above of weighted average, error, and scale factor are for the reader's convenience only. The data were actually processed by a constrained fit program, which calculates its own values of \bar{X} , S , and scale factor, which are different from the values shown here.

CHISO

0.7

2.0

0.2

3.0

(CONLEV = 0.228)

R27 CHAR. K INTO IPI 2P10/TAU (P11/1P1)
 R27 R 427 (10.38) (0.82) YOUNG 65 EMUL + 9/66
 R27 R DELETED FROM OVERALL FIT BECAUSE YOUNG 65 CONSTRAINS HIS RESULTS.
 R27 R TO ADD UP TO 1. ONLY YOUNG 65 DIRECTLY.
 R27
 R27 FIT 11.364 0.072 FROM FIT

R28 CHAR. K INTO IPI 2P10/TAU UNITS 100M-5 (P11/1P1)
 R28 10 1.9 0.7 0.5 BOTTIPI 67 ASPK + 11/67
 R28 B 1.8 0.8 0.6 WACEK 69 ASPK + 4/69
 R28 112 2.42 0.42 CLARK 72 DSPK + 1/73
 R28 534 2.37 0.17 WEARDE 75 SPEC + 11/75
 R28 404 2.51 0.15 HEINTZE 76 SPEC + 2/76
 R28
 R28 AVG 2.4 0.1 AVERAGE ERROR INCLUDES SCALE FACTOR OF 1.01

R29 CHAR. K INTO IPI 2P10/TAU IPI 2P10/TAU (P51/1P6)
 R29 C1505 0.703 0.056 CALLANA 66 HMLC + 6/68
 R29 5601 0.667 0.017 BOTTIPI 68 ASPK + 6/68
 R29 M 1598 (0.608) (0.021) EICHEN 68 HMLC + 10/68
 R29 M (0.596) (0.025) HAIDT 71 HMLC + 12/70
 R29 03480 0.698 0.025 CHANG 72 DSPK + 1.84 GEVVC + 9/72
 R29 B 1595 0.705 0.063 LUCASEZ 73 HRC + 11/73
 R29 B 1595 (0.608) (0.014) BRAUN 75 HMLC + 1/76
 R29 067 0.12 WEISSBERG 76 SPEC + 1/78
 R29 E (0.673) (0.014) GELVITZ 77 SPEC + 12/77

R29 C FROM CALLANA 68 WE USE ONLY THE MU/SZ RATIO AND DO NOT
 R29 C INCLUDE IN THE FIT THE RATIOS MU/TAU AND E3/TAU, SINCE THEY SHOW
 R29 C LARGE DISAGREEMENTS WITH THE REST OF THE DATA.
 R29 H HAIDT 71 IS A REANALYSIS OF EICHTEN 68.
 R29 H ONLY INDIVIDUAL RATIOS INCLUDED IN FIT (SEE R10 AND R20).
 R29 D CHANG 72 R20 IS STATISTICALLY INDEPENDENT OF CHANG 72 R5 AND P6.
 R29 D LUCAS 73 GIVES MU/SZ RATIO. HIS 73A-3-1-1. WE OBTAIN
 R29 D BRAUN 75 VALUE IS FROM FORM FACTOR FIT. ASSUMES MU-E UNIVERSALITY.
 R29 E HEINTZE 77 VALUE FROM FIT TO LANGRAB. ASSUMES MU-E UNIVERSALITY.
 R29
 R29 AVG 0.679 0.013 AVERAGE ERROR INCLUDES SCALE FACTOR OF 1.03
 R29 FIT 0.683 0.018 FROM FIT ERROR INCLUDES SCALE FACTOR OF 1.71

R30 CHAR. K INTO IPI 2P10/TAU MU GAMMA/IPI 67 HRC UNITS 100M-23
 R30 11.21 (0.8) BELLOTTI 67 HMLC + EGAM GT 30MEV 11/67
 R30 R 13 0.74 0.20 ROMANO 71 HMLC + EGAM GT 10MEV 10/71
 R30 R 100 10.53 (0.22) ROMANO 71 HMLC + EGAM GT 30 MEV 9/73
 R30 L 10 0.83 0.20 LUNG 73 HMLC + EGAM GT 30 MEV 9/73
 R30 L (0.22) (0.15) (0.10) LUNG 73 HMLC + EGAM GT 30 MEV 9/73
 R30 L FIRST LUNG VALUE IS FOR COS(ELECT-GAMMA). 0.9. SECOND VALUE IS
 R30 L FOR COS(ELECT-GAMMA) BETH 0.6 AND 0.9 FOR COMPARISON WITH ROMANO.
 R30 R BOTH ROMANO VALUES ARE FOR COS(ELECT-GAMMA) BETH 0.6 AND 0.9.
 R30 R SECOND VALUE IS FOR COMPARISON WITH SECOND LUNG VALUE.
 R30 R WE USE LOWEST EGAM CUT FOR TABLE VALUE. SEE ROMANO FOR EGAM DEPEND.
 R30

R31 K- INTO IPI 1+ E- 3P10/TAU UNITS 100M-5 (P1)
 R31 TEST OF LEPTON NUMBER CONSERVATION.
 R31 11.51 OR LESS CHANG 68 HRC + 3/68

R32 CHAR. K INTO IPI 2P10/TAU UNITS 100M-4 (P20)
 R32 C (1.4) CR LESS CL+90 KLEMS 71 DSPK + TPI1 67-127MEV 3/74
 R32 C (0.94) CR LESS CL+90 KLEMS 73 CNTR + TPI1 60-105 MEV 2/74
 R32 C (0.54) CR LESS CL+90 KLEMS 73 CNTR + TPI1 60-127 MEV 2/74
 R32 L 0.157 0.11 CR LESS CL+90 LUNG 73 HMLC + 1/73
 R32 C 0.14 CR LESS CL+90 ASANO 87 CNTR + TPI1 116-127MEV 1/67
 R32 C KLEMS 73 AND CALB 72 ASSUME PI SPECTRUM SAME AS HES DECAV
 R32 C SECOND CHARGE LIMIT COMBINES CALB AND KLEMS DATA FOR VECTOR INT.
 R32 L LUNG 73 ASSUMES VECTOR INTERACTION.

R33 CHAR. K INTO IPI 2P10/TAU UNITS 100M-51 (P21)
 R33 M 73 MACK 70 DSPK + PEEI 234 TO 247 12/70
 R33 M ABOVE IS MEASUREMENT OF STRUCTURE-DEPENDENT DECAY ONLY.

R34 CHAR. K INTO IPI 2P10/TAU UNITS 100M-6 (P22)
 R34 VIOLATES ANGULAR MOMENTUM CONSERVATION. NOT LISTED IN TABLES.
 R34 K (4.0) CR LESS CL+90 KLEMS 71 DSPK + 1/73
 R34 K TEST OF MODEL OF SELLEN. NC. 2911909.

R35 CHAR. K INTO IPI 2P10/TAU UNITS 100M-4 (P23)
 R35 3.0 CR LESS CL+90 KLEMS 71 DSPK + TPI1 GT 117MEV 6/71

R36 K+ INTO IPI 1+ E- NEU/IPI 2P10/TAU UNITS 100M-4 (P81/1P7)
 R36 C (130.) CR LESS CL+90 BOURQUIN 71 ASPK 8/76
 R36 B 3.0 CR LESS CL+90 BLOCH 76 SPEC 8/76
 R36 B CORRESPONDS TO SELLEN AT CL+90.

R37 CHAR. K INTO IPI 2P10/TAU IPI 2P10/TAU UNITS 100M-4 (P24/1P6)
 R37 0 (37.0) CR LESS CL+90 ROMANO 71 HMLC + 12/71
 R37 2 3.8 5.0 1.2 LUNG 73 HMLC + 9/73

R39 K+ INTO IPI 1+ E- MU-3/TOTAL UNITS 100M-8 (P25)
 R39 K- INTO IPI 1+ E- MU-3/TOTAL IS ALSO INCLUDED HERE
 R39 12.81 CR LESS CL+90 BIECK 72 DSPK + 9/72

Data Card Listings

For notation, see key at front of Listings.

Stable Particles

K⁺

840	K+ INTO EP1 E4 MU-1/TOTAL UNITS 1000+81	EP263
842	K+ INTO EP1 E4 MU-1/TOTAL IS ALSO INCLUDED HERE	
842	11.44 CP LESS CL+09 BETER	72 DSPK -- 9/72
841	CHAR. * INTO [MU NEUTRAL/TOTAL UNITS 1000+0]	EP271
841	P 0 2.0 CP LESS CL+09 PANG	73 CND + 11/73
841	P PANG 73 ASSUMES MU SPECTRUM FROM NEU-NEU INTERACTION OF RAMON 70.	3/74
847	CHAR. * INTO EP10 MU NEU 2441/TOTAL UNITS 1000-51(EP281)	9/73
847	0 0.1 CP LESS CL+09 LUNG	73 HBC + 1 JAN 01 30 NEV 9/73
843	CHAR. * INTO EP12 MU NEU 2103/EP1 PID1	EP31(EP21)
843	L 780 0.221 0.012 LUCAS2 73 HBC - DALLIEZ PAS ONLY	11/73
843	L LUCAS 73 GIVES NEU3-780+3.1PCL, NEP12+3564+3.1PCL, W DIVIDED	11/73
843	*** 0.2290 0.0231 FROM FIT ERROR INCLUDES SCALE FACTOR OF 1.13	
846	CHAR. * INTO EP12 MU NEU 2103/EP1 PID1	EP41(EP21)
846	L 574 0.091 0.009 LUCAS2 73 HBC - DALLIEZ PAS ONLY	11/73
846	L LUCAS 73 GIVES NEU1 2103+574+5.0 PCL, NEP12+3564+3.1 PCL.	11/73
846	L WE CORRECTIONED 2103/NEP12 WHERE PLS IS PEACH ONLY DALLIEZ	11/73
846	L PANG 73 ASSUMES MU SPECTRUM FROM NEU-NEU INTERACTION OF RAMON 70.	11/73
846	*** 0.0819 0.0222 FROM FIT ERROR INCLUDES SCALE FACTOR OF 1.31	
845	CHAR. * INTO EP12 MU NEU 2103/EP1 PID1	EP41(EP21)
845	L 12 2.4 0.5 MESSNER 74 SPEC + G-GAMMA DT 0 MEV	7/74
845	*** 0.0819 0.0222 FROM FIT ERROR INCLUDES SCALE FACTOR OF 1.31	
846	CHAR. * INTO EP12 MU NEU 2103/EP1 PID1	EP41(EP21)
846	P 41 2.5 0.5 MESSNER 74 SPEC + G-GAMMA DT 0 MEV	7/74
846	P MULTIP BY QUOTES THIS RESULT MULTIPLIED BY OUR 1974 KFA HP RATIO.	11/74
846	P MULTIP BY QUOTES THIS RESULT MULTIPLIED BY OUR 1974 KFA HP RATIO.	11/74
847	CHAR. * INTO EP12 MU NEU 2103/EP1 PID1	EP21(EP11)
847	STRUCTURE DEPENDENT PART WITH + GAMMA HELICITY.	
847	H 46 11.051 10.291 10.101 HEMER1 75 SPEC + PEI 730 TO 247	11/75
847	H THIS VALUE IS INCLUDED IN THE SECOND HEINZIE 79 VALUE IN SEC.054	11/75
847	H HEINZ.	11/75
848	K+ INTO EP1+MU+PI+EP1+PI-E NEU3 UNITS 1000+41	EP25+EP26(EP21)
848	D 0 2.0 CP LESS CL+09 DIAMANTRE 76 SPEC +	11/76
848	D DIAMANTRE 76 QUOTES THIS RESULT TIMES OUR 1975 KFA HP RATIO.	11/76
849	K+ INTO EP1+MU+PI+EP1+PI-E NEU3 UNITS 1000+41	EP25(EP21)
849	D 0 1.3 CP LESS CL+09 DIAMANTRE 76 SPEC +	11/76
849	D DIAMANTRE 76 QUOTES THIS RESULT TIMES OUR 1975 KFA HP RATIO.	11/76
850	CHAR. * INTO [MU NEU 611-1/EP1-NEU3 UNITS 1000+3]	EP31(EP21)
850	J 14 13.1 10.9 DIAMANTRE 76 SPEC + MIEI 51 140	11/76
850	J 14 27.0 10.9 DIAMANTRE 76 SPEC + EXTRAPOLATED CP	11/76
850	J DIAMANTRE 76 QUOTES THESE RESULTS TIMES OUR 1975 KFA HP RATIO.	11/76
850	J THE SECOND DIAMANTRE 76 VALUE IS THE FIRST VALUE EXTRAPOLATED TO 0	11/76
850	J TO INCLUDE CP MASS EFFECTS.	11/76
851	CHAR. * INTO EP1+MU+PI+EP1+PI-E NEU3 UNITS 1000+41	EP31(EP21)
851	H 15 15.1 10.9 DIAMANTRE 76 SPEC +	11/76
851	H DIAMANTRE 76 QUOTES THIS RESULT TIMES OUR 1975 KFA HP RATIO.	11/76
852	CHAR. * INTO EP1+MU+PI+EP1+PI-E NEU3 UNITS 1000+41	EP31(EP21)
852	D 0 2.5 CP LESS CL+09 DIAMANTRE 76 SPEC +	11/76
852	D DIAMANTRE 76 QUOTES THIS RESULT TIMES OUR 1975 KFA HP RATIO.	11/76
853	CHAR. * INTO EP1+MU+PI+EP1+PI-E NEU3 UNITS 1000+41	EP32(EP21)
853	H 4 2.4 0.56 0.27 DIAMANTRE 76 SPEC +	11/76
853	H DIAMANTRE 76 QUOTES THIS RESULT TIMES OUR 1975 KFA HP RATIO.	11/76
854	CHAR. * INTO EP1+MU+PI+EP1+PI-E NEU3 UNITS 1000+41	EP11(EP11)
854	STRUCTURE DEPENDENT PART WITH + GAMMA HELICITY	
854	H 51 12.141 10.421 HEINZIE 79 SPEC +	7/79
854	H 107 2.49 7.36 HEINZIE 79 SPEC +	7/79
854	H SECOND HEINZIE 79 VALUE IS FIRST COMBINED WITH HEINZIE 79 RESULT	7/79
854	H FROM SECTION 847 ABOVE.	7/79
855	CHAR. * INTO EP1+MU+PI+EP1+PI-E NEU3 UNITS 1000+41	EP211
855	STRUCTURE DEPENDENT PART WITH - GAMMA HELICITY	
855	H 14 1.4 CP LESS CL+09 HEINZIE 79 SPEC +	7/80
855	H IMPLIES LARVAL VEC. VECTOR) AMPL. RATIO OUTSIDE RANGE -1.8 TO -0.6.	7/80
856	CHAR. * INTO EP1+MU+PI+EP1+PI-E NEU3 UNITS 1000+41	EP31(EP11)
856	H 14 1.4 CP LESS CL+09 HEINZIE 79 SPEC +	7/79
857	K+ INTO EP12 MU NEU 2103/EP1 PID1	EP41
857	D 0 3.06 CP LESS CL+09 LYONS 81 HBC 2009 + + N.B. BLM	2/82

Note on Slope Parameter for K⁺ → 3π⁺ Decays

As was discussed in Section VI B.1 of the text, for the 3π⁺ decays of the K mesons we list the slope parameter "g" which is defined, as in that section, by

$$|N|^2 \propto 1 + g \frac{(s_3 - s_0)}{m_{\pi^+}^2} + h \frac{(s_3 - s_0)}{m_{\pi^+}^2}^2 + j \frac{(s_2 - s_1)}{m_{\pi^+}^2} + k \frac{(s_2 - s_1)}{m_{\pi^+}^2}^2 + \dots \quad (1)$$

where

$$s_i = (P_K - P_{\pi^+})^2 = (m_K - m_{\pi^+})^2 - 2m_K^2 \tau_i \quad (2)$$

$$s_0 = \frac{1}{3} \int s_i = \frac{1}{3} (m_K^2 + m_{\pi^+}^2 + m_{\pi^+}^2) \quad (3)$$

P_K , P_{π^+} are the four-vectors for the K and the i^{th} pion, and the index 3 refers to the odd pion, i.e., the third pion

in the decays listed below.

We refer to the three possible charged decays as

τ, τ', and τ'':

$$\tau^+ \rightarrow K^+ + \pi^+ \pi^+ \pi^+$$

$$\tau^{+\prime} \rightarrow K^+ + \pi^0 \pi^+ \pi^+$$

$$\tau^{+\prime\prime} \rightarrow K_L^0 + \pi^+ \pi^+ \pi^0$$

The measurements of g vary considerably beyond the authors' quoted errors as can be seen in the ideograms associated with the GT⁺, GT⁻, and GTP subsections of the K⁺ Data Card Listings and the GTO subsection of the K_L⁰ Listings.

There is no indication of a CP-violating asymmetry in K_L⁰ decay as measured by the coefficient j given in subsection JTO of the K_L⁰ Listings.

The high-statistics τ⁺-decay experiment of MESSNER 74 finds significant non-zero quadratic coefficients h and k. CHO 77, a lower-statistics τ⁰ experiment, obtains results in agreement with MESSNER 74 but can also obtain good fits with a linear term (g) only. The correlation between the linear and quadratic coefficients changes the CHO 77 g_± from 0.629 ± 0.017 (linear fit) to 0.681 ± 0.024 (quadratic fit). Another experiment, PEACH 77, does not observe this correlation and is in agreement only with the linear fit of CHO 77.

There is some evidence for a non-zero k coefficient from τ⁺ experiments. FORD 72 (1.5M events) have studied K⁺ → π⁺ π⁺ π⁺ and find that the χ²/DF goes from 1.38 to 1.20 for DF ≈ 150 when the second order and CP-violation terms are added.

Stable Particles K[±]

However, the authors state that since their Coulomb correction is larger than the experimental errors and is not well known, it is difficult to interpret these results. DEVAUX 77 also finds a non-zero k.

Because of the above evidence for quadratic terms, and for consistency in our treatment of τ^0 and τ^\pm decay, we now include in our averages only those τ^0 and τ^\pm experiments for which we have information on the three coefficients g, h, and k. Correlations prevent us from comparing fits which do not include these three parameters. For τ^\pm decays we compile g and h only since no experiments measure k.

Parametrizations

In the literature *other definitions* of slope parameters have appeared. We have converted to the definitions of g, h, j and k in Eq. (1) from whatever experimental quantity has been reported. We give the conversion to the definition (1) for the most widely used parametrizations and tabulate the conversion factors for the reader's convenience.

a) For analysis of charged K's and some K⁰ experiments, the expression often used is:

$$|M|^2 = 1 + a_Y Y + b_Y Y^2 + d_Y X + e_Y X^2$$

with

$$Y = \frac{3T_3 - Q}{Q},$$

$$X = \frac{\sqrt{3} (T_1 - T_2)}{Q},$$

$$Q = m_K - \sum m_i.$$

The relevant formulae are:

$$Y = -\frac{3}{2} \frac{s_3 - s_0}{m_K Q} + \Delta, \quad X = \frac{\sqrt{3}}{2} \frac{s_2 - s_1}{m_K Q}$$

with

$$\Delta = \frac{m_- - m_3}{Q} \left(2 - \frac{m_3 + m_1}{m_K} \right)$$

and

Data Card Listings

For notation, see key at front of Listings.

$$g = \frac{-c_Y (a_Y + 2b_Y \Delta)}{1 + a_Y \Delta + b_Y \Delta^2},$$

$$h = \frac{c_Y^2 b_Y}{1 + a_Y \Delta + b_Y \Delta^2},$$

$$j = \frac{c_Y d_Y}{\sqrt{3} (1 + a_Y \Delta + b_Y \Delta^2)},$$

$$k = \frac{c_Y^2 e_Y}{3(1 + a_Y \Delta + b_Y \Delta^2)},$$

with

$$c_Y = \frac{3}{2} \frac{m_\pi^2}{m_K Q}.$$

b) For the analysis of some K⁰ experiments the expression used is

$$|M|^2 = 1 + 2a_t \frac{m_K}{m_\pi^2} (2T_3 - T_{3\max}) + b_t \left(\frac{m_K}{m_\pi^2} \right)^2 (2T_3 - T_{3\max})^2,$$

with

$$T_{3\max} = \frac{(m_K - m_3)^2 - (m_1 + m_2)^2}{2m_K}.$$

The relevant transformations are

$$T_3 = -\frac{s_3 - s_0}{2m_K} + \frac{Q}{3} (1 + \Delta)$$

and

$$g = \frac{-2a_t - b_t c_t}{1 + a_t c_t + \frac{b_t c_t^2}{4}},$$

$$h = \frac{b_t}{1 + a_t c_t + \frac{b_t c_t^2}{4}},$$

with

$$c_t = \frac{2m_K}{m_\pi^2} \left[\frac{2}{3} Q(1 + \Delta) - T_{3\max} \right].$$

Data Card Listings

For notation, see key at front of Listings.

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K⁺

c) Other K⁰ authors use the same form of matrix element as given in b) above with a linear term only, but define

$$T_{\max} = \frac{2}{3} Q$$

The relevant transformation is then

$$g = \frac{-2a_u}{1 + a_u c_u}$$

with

$$c_u = \frac{4m_K}{3m_{\pi^+}} Q \Delta$$

d) Older K⁰ analyses were done using

$$|M|^2 = 1 + a_v \frac{T_3}{m_K}$$

The relevant transformation is then

$$g = \frac{-c_v a_v}{1 + d_v a_v}$$

with

$$c_v = \frac{m_{\pi^+}}{2m_K^2}$$

and

$$d_v = \frac{Q}{3m_K} (1 + \Delta)$$

e) The CP-violating term in $|M|^2$ for K_L⁰ → π⁺π⁻π⁰ experiments has been parametrized in several ways. BLANPIED 68 and SCRIBANO 70 use the parametrization given in (b) above with no quadratic term and with an additional CP violating term. BLANPIED 68 parametrizes the CP-violating term as

$$2\alpha_B \frac{m_K}{m_{\pi^+}} (T_1 - T_2)$$

The relevant transformation is then

$$j = \frac{\alpha_B}{1 + c_t a_t}$$

with c_t as defined in (b) above. SCRIBANO 70 parametrizes the CP-violating term as

$$\frac{2}{\sqrt{3}} \alpha_S \frac{T_1 - T_2}{m_{12\max}}$$

where T_{12max} is the maximum kinetic energy of particle 1 or 2, the charged π's, given by

$$T_{12\max} = \frac{(m_K - m_1)^2 - (m_2 + m_3)^2}{2m_K}$$

The resulting transformation is then

$$j = \frac{m_{\pi^+}^2}{\sqrt{3} m_K T_{12\max}} \frac{\alpha_S}{(1 + c_t a_t)}$$

SMITH 70 gives the asymmetry

$$a = \frac{N_+ - N_-}{N_+ + N_-}$$

where N₊ is the number of events with T₁ > T₂ and N₋ is the converse. BLANPIED 68 gives the relation α_B = α/1.16 which allows us to use the transformation to j given above for BLANPIED 68.

For the reader's convenience we give a table of numerical values for Q, T_{3max}, T_{12max}, Δ, c_γ, c_t, c_u, c_v, and d_v, obtained using the masses from the current edition.

	T [±]	T [±]	T ⁰
Q	74.97	84.17	83.57
T _{3max}	48.08	53.20	53.89
T _{12max}	48.08	53.99	53.12
Δ	0.0000	-0.0790	0.0798
c _γ	0.7895	0.7032	0.7025
c _t	0.0962	-0.0769	0.3204
c _u	0.0030	-0.2246	0.2272
c _v	0.0400	0.0400	0.0393
d _v	0.0506	0.0523	0.0604

References

See the reference sections of the K[±] and K_L⁰ Data Card Listings.

See also the review of T. J. Devlin and J. O. Dickey, Rev. Mod. Phys. 51, 237 (1979), which contains an analysis of K → 2π and K → 3π data in terms of transition amplitudes with appropriate energy dependence.

Stable Particles

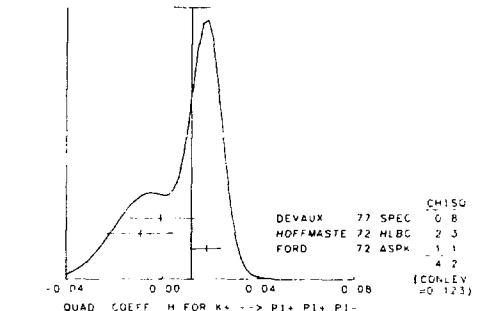
K⁺

Data Card Listings

For notation, see key at front of Listings.

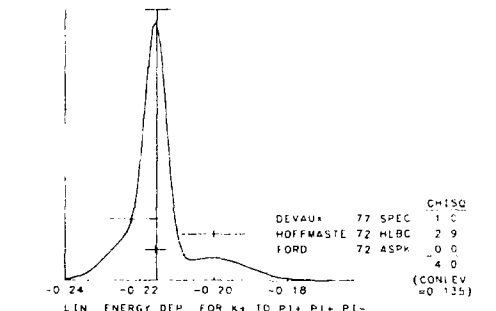
10 CHARGED * ENERGY DEPENDENCE OF DELTA PLOT
 RELATED TEXT SECTION VI B.1 AND MINI-REVIEW ABOVE
 MATRIX ELEMENT SQUARED = 1 * QW + HXJW2 + *WWW2 1/774
 WHERE UNLS=SQRT(HXJW2) AND HXCG=SQRT(WWWW2) 1/774
 G+ LINEAR COEFFICIENT G FOR TAU DECAYS K+ -> PI+ PI- PI+ 1/774
 G- LINEAR COEFFICIENT G FOR TAU DECAYS K+ -> PI+ PI- PI+ 1/774
 G+ TERM AT RIGHT, SEE MINI-REVIEW ABOVE. 1/774
 STAZL 5428 (-0.221) (0.024) ZEMCHENKO 67 HBC + AY=0.274+-03 10/69
 G+ L 1904 (-0.218) (0.016) BUTLER 68 HBC + AY=0.217+-02 10/69
 G+ 13198 (-0.196) (0.012) GRAMMAN 70 HLCB + AY=0.229+-03 8/70
 G+ 750K (-0.2157) (0.008) FORD 72 ASPK + AY=0.7734+-0025 1/79
 ST+H 39819 -0.200 0.009 HOFFMASTE 72 HLCB + 1/79
 G+ 225K -0.2221 (0.0085) DEVAUX 77 SPEC + AY=0.2814+-0062 1/79
 G+ L EXPERIMENTS WITH LARGE ERRORS NOT INCLUDED IN AVERAGE. 1/79
 G+ F ALSO INCLUDES DBC EVENTS 3/78
 G+ Q EMULS. DATA ADDED - ALL EVENTS INCLUDED BY HOFFMASTE 72 1/79
 G+ H HOFFMASTE 72 INCLUDES ORGANA. TO DATA. 1/78
 G+ * * * * * 1/79
 G+ AVG -0.2154 ± 0.0035 AVERAGE ERROR INCLUDES SCALE FACTOR OF 1.43
 (SEE IDENTIFAN BELOW)

WEIGHTED AVERAGE = 0.0122 ± 0.0016
 ERROR SCALED BY 1.4



HT+ QUADRATIC COEFF. H FOR K+ -> PI+ PI- PI- 1/79
 HT+ 750K -0.0187 (0.0067) FORD 72 ASPK + 1/79
 HT+ 39819 -0.0109 (0.014) HOFFMASTE 72 HLCB + 1/79
 HT+ 225K -0.0006 (0.014) DEVAUX 77 SPEC + 1/79
 HT+ * * * * * 1/79
 HT+ AVG -0.0127 ± 0.0035 AVERAGE ERROR INCLUDES SCALE FACTOR OF 1.43
 (SEE IDENTIFAN BELOW)

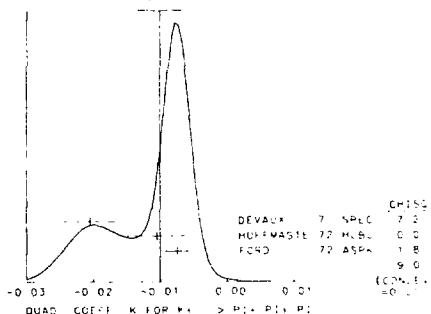
WEIGHTED AVERAGE = -0.2154 ± 0.0035
 ERROR SCALED BY 1.4



K+ QUADRATIC COEFF. K FOR K+ -> PI+ PI- PI- 1/79
 K+ 750K -0.0075 (0.0019) FORD 72 ASPK + 1/79
 K+ 39819 -0.0105 (0.004) HOFFMASTE 72 HLCB + 1/79
 K+ 225K -0.0205 (0.0039) DEVAUX 77 SPEC + 1/79
 K+ * * * * * 1/79
 K+ AVG -0.0101 ± 0.0035 AVERAGE ERROR INCLUDES SCALE FACTOR OF 2.11
 (SEE IDENTIFAN BELOW)

G+ LINEAR COEFFICIENT G FOR TAU DECAYS K+ -> PI+ PI- PI+ 1/79
 G- FOR DEFINITION OF AY SEE NOTE IN SECTION G1 ABOVE.
 G+ F 1347 (-0.220) (0.035) FERRO-LIJ 61 HBC + AY=0.284+-045 10/69
 ST+H 5378 (-0.193) (0.023) MSL 050 68 HBC + AY=0.242+-029 10/69
 G+ 50919 -0.193 (0.010) MSL 69 HBC + AY=0.244+-013 1/79
 G+ 750K -0.2158 (0.008) FORD 72 ASPK + AY=0.211+-0035 1/79
 G+ Q 81K (-0.198) (0.0089) LICHT 75 HBC + AY=0.252+-011 10/62
 G+ F NO RADITIVE CORRECTIONS ENCLOSED. 3/78
 G+ L EXPERIMENTS WITH LARGE ERRORS NOT INCLUDED IN AVERAGE.
 G+ M ALSO INCLUDES DBC EVENTS.
 G+ Q QUADRATIC DEPENDENCE IS REQUIRED BY KL EXPTS. FOR COMPARISON W/ 1/79
 G+ O AVERAGE ONLY THOSE K+ EXPERIMENTS WHICH QUOTE QUADRATIC FIT VALUES. 1/79
 G+ * * * * * 1/79
 G+ AVG -0.2167 ± 0.0066 AVERAGE ERROR INCLUDES SCALE FACTOR OF 2.51

WEIGHTED AVERAGE = -0.0107 ± 0.0034
 ERROR SCALED BY 2.1



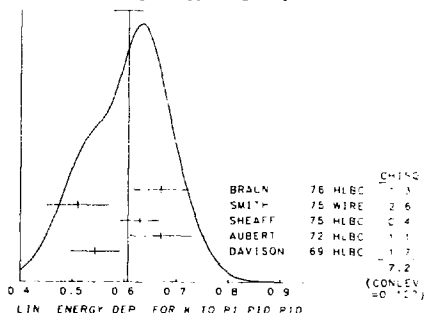
HT- QUADRATIC COEFF. H FOR K+ -> PI+ PI- PI+ 1/79
 HT- 50919 -0.001 (0.012) MSL 69 HBC + 1/79
 HT- 750K -0.0125 (0.0062) FORD 72 ASPK + 1/79
 HT- * * * * * 1/79
 HT- AVG -0.0007 ± 0.0055 AVERAGE ERROR INCLUDES SCALE FACTOR OF 1.33

K- QUADRATIC COEFF. K FOR K+ -> PI+ PI- PI+ 1/79
 K- 50919 -0.014 (0.012) MSL 69 HBC + 1/79
 K- 750K -0.0081 (0.0019) FORD 72 ASPK + 1/79
 K- * * * * * 1/79
 K- AVG -0.0084 ± 0.0019 AVERAGE ERROR INCLUDES SCALE FACTOR OF 1.33

JG (G10) -0.1375 (0.147) IN PERCENT
 DG A NON-ZERO VALUE FOR THIS QUANTITY INDICATES CP VIOLATION
 JG 3.24 -0.70 0.43 FORD 72 ASPK 11/77

G+ LINEAR COEFFICIENT G FOR TAU DECAYS CHARG. K+ -> PI+ PI- PI+ 1/79
 G+ UNLESS OTHERWISE STATED, ALL EXPTS. INCLUDE TERMS QUADRATIC IN 1/79
 STP (E3-5917/14/1972), SEE MINI-REVIEW ABOVE.
 STP K 1792 (0.481) (0.261) KALNIK 64 HLCB + 1/79
 STP K 1874 (0.508) (0.094) MSL 65 HBC + ALSO HBC 1/79
 STP 4048 0.544 0.088 DAVIDSON 69 HLCB + ALSO EMUL 1/79
 STP L 198 (0.527) (0.107) RANDOLPH 70 EMUL + 1/79
 STP 1365 0.481 0.036 AUBERT 72 HLCB + 1/79
 STP K 474 (0.484) (0.084) LICHT 73 HBC + DALITZ PHS 7/67 1/79
 STP 5835 0.630 (0.036) SHEAR 75 HLCB + 1/79
 STP 27K 0.510 (0.050) SMITH 75 HLCB + 1/79
 STP 1 4430 (0.809) (0.223) BERTRAND 76 EMUL + 1/79
 STP 3263 0.670 (0.054) BRAUN 76 HLCB + 1/79
 G+ A AUTHOR. GIVE LINEAR FIT ONLY.
 G+ L EXPERIMENTS WITH LARGE ERRORS NOT INCLUDED IN AVERAGE.
 G+ * * * * * 1/79
 G+ AVG -0.627 ± 0.277 AVERAGE ERROR INCLUDES SCALE FACTOR OF 1.33
 (SEE IDENTIFAN BELOW)

WEIGHTED AVERAGE = 0.607 ± 0.230
 ERROR SCALED BY 1.3



HTP QUADRATIC COEFF. H FOR CHRG. K+ -> PI+ PI- PI+ 1/79
 HTP 4048 0.026 (0.350) DAVIDSON 69 HLCB + ALSO EMUL 1/79
 HTP L 198 (0.018) (0.124) RANDOLPH 70 EMUL + 1/79
 HTP 1365 -0.01 (0.7) AUBERT 72 HLCB + 1/79
 HTP 5835 0.041 (0.030) SMITH 75 HLCB + 1/79
 HTP 27K 0.009 (0.040) SHEAR 75 WIRE + 1/79
 HTP L 4039 (0.094) (0.223) BERTRAND 76 EMUL + 1/79
 HTP 3263 0.152 (0.082) BRAUN 76 HLCB + 1/79
 HTP L EXPERIMENTS WITH LARGE ERRORS NOT INCLUDED IN AVERAGE. 1/79
 HTP * * * * * 1/79
 HTP AVG -0.036 ± 0.020 AVERAGE ERROR INCLUDES SCALE FACTOR OF 1.01

Data Card Listings

For notation, see key at front of Listings.

Stable Particles

K*

K_L → π⁰ on K_L⁺ and K_L⁰ Form Factors

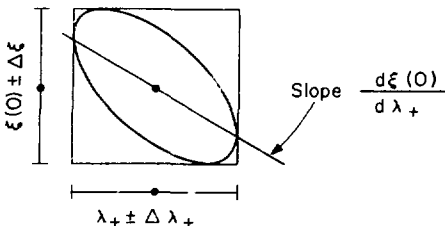
Definitions of the parameters λ_+ , $f_+(0)$, λ_0 , f_S/f_+ , and f_T/f_+ and a general discussion of the methods of analysis are given in Section VI B.2 of the text.

This note describes the contents of the Data Card Listings for the two K_L⁺ parametrizations, $(\lambda_+, f_+(0))$ and (λ_+, λ_0) , which were discussed in the text. Problems related to our data entries for individual experiments are discussed and a comparison of results is given.

K_L⁺ Experiments

The matrix element for K_L⁺ decay, assuming a pure vector current, is given by Eq. (2) in Section VI B.2 of the text. Most experiments appear to be compatible with the assumption that f_+ depends linearly on q^2 and that f_- is constant. Only DALLY 72 (K_L⁺) appears to require $f_- \neq 0$ (by about three standard deviations). A single data bin at low q^2 seems to be responsible. The effect is not observed in the high-statistics experiment of DONALDSON 74 (also K_L⁺).

$(\lambda_+, f_+(0))$ Parametrization: λ_+ data from K_L⁺ decay are entered in the K_L⁺ and K_L⁰ sections of the Data Card Listings in subsection L+N. The corresponding $f_+(0)$ values are entered in subsection XIA, XIB, or XIC, depending on whether Method A, B, or C, discussed below and in the text, was used. The data cards contain the values, one-standard-deviation errors $d\lambda_+$ and $d^2f_+(0)/d\lambda_+^2$, as well as the correlation $d^2f_+(0)/d\lambda_+^2$, all indicated on the $e^{-1/2}$ likelihood contour below. The correlations are given on the right side of the $f_+(0)$ data cards.



XBL 743-2602

(λ_+, λ_0) Parametrization: This parametrization is used in recent K_L⁺ analyses. To facilitate comparison between experiments, we convert earlier experiments from the $(\lambda_+, f_+(0))$ parametrization to (λ_+, λ_0) whenever possible (i.e., when λ_+ and $f_+(0)$ values, errors, and correlations are given). The transformation between these parametrizations is:

$$\lambda_0 = \lambda_+ + a f_+(0),$$

$$d\lambda_0^2 = (1 + 2a \frac{d^2 f_+(0)}{d\lambda_+^2}) d\lambda_+^2 + a^2 d^2 f_+^2,$$

$$\frac{d\lambda_0}{d\lambda_+} = 1 + a \frac{d^2 f_+(0)}{d\lambda_+^2},$$

where $a = m_\pi^2 / (m_K^2 - m_\pi^2)$. The λ_0 value, the one-standard-deviation error $d\lambda_0$, and the correlation $d\lambda_0/d\lambda_+$ are given in subsection L0 of the data cards.

We also convert (λ_+, λ_0) results into the $(\lambda_+, f_+(0))$ parametrization whenever possible so that subsection L0 is essentially equivalent to the three subsections XIA, XIB, and XIC.

Individual analyses have used a variety of parametrizations. Problems arise when trying to express their results in terms of the parametrizations used here. The discussion of these problems is divided into three sections corresponding to the three methods of analyses discussed in the text.

Method A: Dalitz plot analyses and pion spectrum analyses usually determine λ_+ and $f_+(0)$ (or λ_0) values, errors, and correlation. Such measurements are entered in the L+N, XIA, and L0 subsections. They give rise to the error ellipses shown in Figs. 1 and 2. These are approximations to likelihood contours.

Some analyses of this type fix λ_+ and determine $f_+(0)$, e.g., CARPENTER 66 and PEACH 73 (both K_L⁺). We enter $f_+(0)$ and $d^2 f_+(0)/d\lambda_+^2$ in the XIA section and give the fixed λ_+ value in the data card footnote. The $f_+(0)$ error is parenthesized because it does not include the uncertainty in the value of λ_+ . These results, transformed to λ_0 measurements, give rise to bands in Fig. 2. These bands are also approximations to the likelihood contours. The actual likelihood bands will not be straight.

Stable Particles

K^+

In some cases, we alter an error from its published value in order to obtain an error ellipse with a width which matches the error in $\xi(0)$ for fixed λ_+ . These adjustments are noted in the $\xi(0)$ data card footnotes, e.g., for CALLAHAN 66 and HAIDT 71 (K^+ subsection XIA), where the published errors and correlation violate the constraint $|C_{\lambda\xi}| < 1$ on the normalized correlation coefficient $C_{\lambda\xi}$ given by

$$C_{\lambda\xi} = \frac{d\lambda_+}{d\xi} \frac{d\xi(0)}{d\lambda_+}.$$

BIRULEV 81 ($K_{\mu 3}^+$) gave a correlation $d\lambda_+/d\xi$ which, while not unphysical, is so large that the resulting error ellipse is unreasonably narrow. As a result, in the fit to be discussed, BIRULEV 81 would dominate even DONALDSON 74 which has 10 times the BIRULEV 81 statistics. To prevent this, we make the assumption $d\lambda_+/d\xi = 0$, giving the broad error ellipse shown in Fig. 2.

In some cases, e.g., BRAUN 73, the parametri-

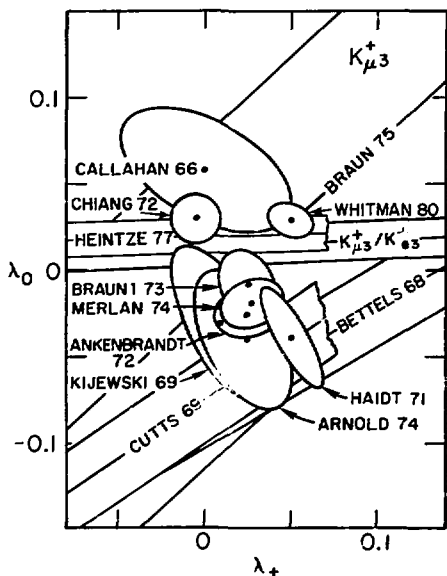


Fig. 1. One-standard-deviation ($e^{-1/2}$) likelihood contours in the (λ_+, λ_0) plane for $K_{\mu 3}^+$.

Data Card Listings

For notation, see key at front of Listings.

zation used is λ_+ , $\xi(0)$, $\xi(t^*)$, where t^* is the weighted average of t with weighting according to the sensitivity to ξ . In this case we do not use $\xi(0)$. It is a badly determined parameter comparable to λ_- or the slope of $\xi(t)$. Instead, we use

$$\xi(0) = \xi(t^*)(1 + \lambda_+ t^*),$$

$$\frac{d\xi(0)}{d\lambda_+} = \frac{d\xi(t^*)}{d\lambda_+} (1 + \lambda_+ t^*) + \xi(t^*) t^*.$$

With the BRAUN 73 values, $\lambda_+ = 0.027$, $\xi(6.6) = -0.34 \pm 0.20$, and $d\xi(6.6)/d\lambda_+ = -14$, we obtain

$$\xi(0) = (-0.40 \pm 0.24) - 19(\lambda_+ - 0.027);$$

or for their fitted $\lambda_+ = 0.025 \pm 0.017$, we get $\xi(0) = -0.36 \pm 0.40$.

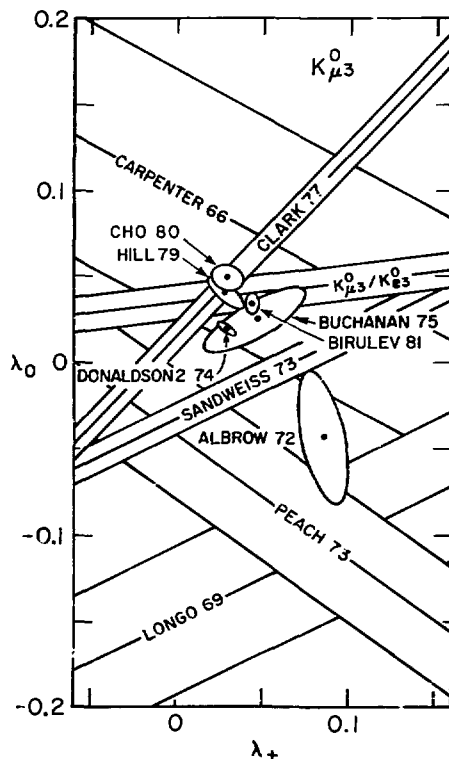


Fig. 2. One-standard-deviation ($e^{-1/2}$) likelihood contours in the (λ_+, λ_0) plane for $K_{\mu 3}^0$.

Data Card Listings

For notation, see key at front of Listings.

Stable Particles

 K^+

Method B: Branching ratio experiments cannot determine λ_+ and $\xi(0)$ simultaneously, but simply fix a relationship between them, given in Section VI B.2 of the text. Results are usually quoted as values of $\xi(0)$ at fixed λ_+ . We list these results in subsection XI B, but we do not average them because the λ_+ values differ. Instead, we compute a combined result by using the relations in the text and our fitted values of $\Gamma(K_{\mu 3}^+)/\Gamma(K_{e 3}^+)$ and $\Gamma(K_{\mu 3}^0)/\Gamma(K_{e 3}^0)$, which include the branching ratios from these experiments. The branching ratios from our 1980 edition and the results for $\xi(0)$ and λ_0 evaluated at $\lambda_+ = 0.030$ are

	K^+	K_L^0
$\Gamma(K_{\mu 3}^+)/\Gamma(K_{e 3}^+)$	$0.663 \pm .018 (S=1.7)$	$0.695 \pm .017$
$\xi(0)$	$-0.20 \pm .15 (S=1.7)$	$+0.08 \pm .13$
$d\xi(0)/d\lambda_+$	-11.9	-10.3
λ_0	$+0.014 \pm .012 (S=1.7)$	$+0.037 \pm .011$
$d\lambda_0/d\lambda_+$	$+0.04$	$+0.12$

The current (1982 edition) values differ only in the addition of the CHO 80 K_L^0 branching ratio measurement. We include their Dalitz-plot-analysis results in our overall form factor fit, and since their branching ratio measurement is not an independent result, we do not include it here. The scale factor S in the above table is the amount by which the error has been multiplied in order to compensate for discrepancies in the branching ratios. These λ_0 results give rise to the $K_{\mu 3}^+/K_{e 3}^+$ bands in Figs. 1 and 2.

Method C: Polarization experiments measure $\langle \xi(t) \rangle$, the weighted average of $\xi(t)$ over the t range of the experiment, where the weighting accounts for the variation with t of the sensitivity to $\xi(t)$. Such measurements are entered in subsection XI C.

To reinterpret these results in the $(\lambda_+, \xi(0))$ parametrization, we recognize that $\lambda_+ = 0$ corresponds to $\xi(t)$ constant (always assuming $\lambda_- = 0$) so that

$$\xi(0) \Big|_{\lambda_+=0} \equiv \langle \xi(t) \rangle .$$

The correlation with λ_+ is given by the following relations (valid for small λ_+):

$$\xi(0) \approx \langle \xi(t) \rangle (1 + \lambda_+ \langle \frac{t}{m_\pi^2} \rangle) ,$$

$$\frac{d\xi(0)}{d\lambda_+} \approx \langle \xi(t) \rangle \langle \frac{t}{m_\pi^2} \rangle ,$$

where $\langle t/m_\pi^2 \rangle$ is the average value of t weighted by the sensitivity to $\xi(t)$. These results, transformed to λ_0 and $d\lambda_0/d\lambda_+$ values, are entered in subsection I 0 and give rise to bands in Figs. 1 and 2.

In Figs. 1 and 2, we do not include those polarization measurements for which $d\xi(0)/d\lambda_+$ is not obtainable. Also we do not include the MERLAN 74 ($K_{\mu 3}^+$) polarization result because the signs of $\xi(0)$ and $d\xi(0)/d\lambda_+$ are opposite, whereas the above equation requires them to be the same (since $t > 0$).

Comparison of $K_{\mu 3}^+$ Experiments: Figures 1 and 2 show the likelihood contours in the (λ_+, λ_0) plane for $K_{\mu 3}^+$ and $K_{\mu 3}^0$ respectively.

Most $K_{\mu 3}^+$ Dalitz plot results (ellipses) shown are fairly consistent and appear to cluster between the $K_{\mu 3}^+/K_{e 3}^+$ result and the polarization results of BETTELS 68 and CUTTS 69. However, the latest experiment, WHITMAN 80, finds larger values of λ_0 and λ_+ than most others.

The $K_{\mu 3}^0$ results are even less consistent. The latest results, HILL 79, CHO 80, and BIRULEV 81 are not in very good agreement with the earlier high-precision experiment of DONALDSON 74.

χ^2 fits to the results shown in Fig. 1 and Fig. 2 yield the following values for λ_+ and λ_0 . The corresponding values of $\xi(0)$ are also given.

Stable Particles

 K^+

	$K_{\mu 3}^+$	$K_{\mu 3}^0$
λ_+	$+0.032 \pm 0.008^*$	$+0.034 \pm 0.005^*$
λ_0	$+0.004 \pm 0.007^*$	$+0.025 \pm 0.006^*$
$d\lambda_0/d\lambda_+$	-0.16	-0.16
χ^2/DF	56/21	88/16
s^*	1.6	2.3
$\xi(0)$	$-0.35 \pm 0.15^*$	$-0.11 \pm 0.09^*$
$d\xi(0)/d\lambda_+$	-14.	-14.

*All errors have been increased by the scale factor $S = (\chi^2/DF)^{1/2}$ to take into account the discrepancies between measurements.

In view of the large χ^2/DF , the fit results should be taken with a grain of salt. The largest contributors to χ^2 in the $K_{\mu 3}^+$ case are WHITMAN 80 with 12, CHIANG 72 with 8.3, the polarization results, BETTELS 68 with 7.4 and CUTTS 69 with 5.9, and the π spectrum result of ANKENBRANDT 72 with 5.3. In the $K_{\mu 3}^0$ case the largest contributors are the polarization results of SANDWEISS 73 with 21, LONGO 69 with 14, and CLARK 77 with 8.4, and the Dalitz plot results of ALBROW 72 with 12, CHO 80 with 10, DONALDSON2 74 with 6.0, BIRULEV 81 with 5.8, and PEACH 73 with 5.8. All other χ^2 values were less than 5.

The DONALDSON2 74 result

$$\lambda_+ = 0.030 \pm 0.003$$

$$\lambda_0 = 0.019 \pm 0.004$$

clearly dominates the statistics in the $K_{\mu 3}^0$ case. The λ_+ value is consistent with the K_{e3} value of λ_+ , and with the pole approximation

$$f_+(t) = f_+(0) \frac{m_{K^*}^2}{m_{K^*}^2 - t}$$

Their $f_0(t)$ extrapolates linearly to the Callan-Treiman point. It is less than two standard deviations from the $K_{\mu 3}/K_{e3}$ result.

 K_{e3} Experiments

The f_- term of the matrix element [Eq. (2) text

Data Card Listings

For notation, see key at front of Listings.

Section VI B.2] can be neglected for K_{e3} because it is proportional to the lepton mass. The f_+ term is usually assumed to be linear in $t = q^2 = (P_K - P_\pi)^2$, the square of the four-momentum transfer, i.e., the effective mass of the lepton pair. We quote the linear coefficient λ_+ (L+E on the data cards). There has been some suggestion of departure from linearity [CHIEN 71 (K_{e3}^0) and Chounet, Gaillard, and Gaillard¹ - Review] but no compelling evidence. The λ_+ results are fairly consistent and the average values are

$$K_{e3}^+: \lambda_+ = 0.0285 \pm 0.006$$

$$K_{e3}^0: \lambda_+ = 0.0300 \pm 0.0016 \quad (S=1.2)$$

where the K_{e3}^0 error has been multiplied by the scale factor 1.2 to compensate for inconsistencies (see ideogram in K_L^0 section L+E).

See also the excellent reviews of Gaillard and Chounet,¹ Chounet, Gaillard, and Gaillard,² and Pondrom.³

References

1. M. K. Gaillard and L. M. Chounet, "K₂₃ Form Factors," CERN 70-14 (May 1970), and Phys. Letters 32B, 505 (1970).
2. L. M. Chounet, J. M. Gaillard, and M. K. Gaillard, Physics Reports 4C, 199 (1972).
3. L. G. Pondrom, "Weak Decay Processes," Proc. Particles and Fields 1976, BNL, Oct. 6-8, 1976.

10 CHARGED K FORM FACTORS

RELATED TEXT SECTION VI B.2 AND MINI-REVIEW ABOVE.

IN THE FORM FACTOR COMMENTS, THE FOLLOWING ABBREVIATIONS ARE USED. F₊ AND F₋ ARE FORM FACTORS FOR THE VECTOR MATRIX ELEMENT. F_S AND F_T REFER TO THE SCALAR AND TENSOR TERM.

FD = (F₊ - F₋) / (F₊ + F₋) / (M_K + M_π)²

L₊, L₋ AND L₀ ARE THE LINEAR EXPANSION COEFFS. OF F₊, F₋ AND F₀.

L₊ REFERS TO THE K_{μ3} VALUE EXCEPT IN THE K_{e3} SECTIONS.

DX/DL IS THE CORRELATION BETWEEN X/DL AND L₊ IN K_{μ3}.

DL/DL₊ IS THE CORRELATION BETWEEN DL AND L₊ IN K_{μ3}.

M = MOMENTUM TRANSFER TO THE π IN UNITS OF M_π².

DP = DALITZ PLOT ANALYSIS

PI = π SPECTRUM ANALYSIS

MU = μ SPECTRUM ANALYSIS

POL = POLARIZATION ANALYSIS

BR = BRANCHING RATIO ANALYSIS

E = POSITRON OR ELECTRON SPECTRUM ANALYSIS

RC = RADIATIVE CORRECTIONS

XIA	XIA	F/F ₊ (DETERMINED FROM SPECTRA)						
XIA	76	1±1.81	(0.61)	BROWN	62 XERC	DP+BR, L=0	1/74	
XIA	87	1±0.71	(0.53)	GLACONELL	64 ERU	MU+BR (MU+K _{e3}), L=0	1/74	
XIA	J	1±0.03	(0.71)	JENSEN	64 XERC	DP+BR (MU+K _{e3})	1/74	
XIA	2448	1.0±0.1	(1.13)	(0.91)	CALLAN	66 FRUC	MU, L=0, T UNKN	1/74
XIA	C 444	+0.72	0.93		CALLAN	68 FRUC	PI, DX/DL=-12	1/74
XIA	78	-0.51	(0.91)	EISLER	68 HLBC	PI, L=0, NO DX/DL	1/74	
XIA	K2041	-0.5	0.8	KIJEWSKI	69 DSPA	PI, DX/DL=-26	1/74	
XIA	K3240	-1.1	0.56	HOT	71 HLBC	DP, DX/DL=-28	1/74	
XIA	A625	-0.62	0.28	ANKENBRANDT	72 ASPA	PI, DX/DL=-12	1/74	
XIA	B3460	+0.45	0.28	CHIANG	72 DSPA	DP, DX/DL=-15	1/74	
XIA	D1927	-0.2	0.40	REAU	73 HLBC	DP, DX/DL=-10	1/74	
XIA	N 490	-0.8	0.8	ARMOLD	74 HLBC	DP, DX/DL=-20	1/75	
XIA	M6527	-0.27	0.24	HELAND	75 ASPA	DP, DX/DL=-0	3/74	
XIA	3973	-0.27	0.25	WHITMAN	80 SPEC	DP, DX/DL=-17	4/82	
XIA	ANG	-0.32	0.15			AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.3)		
XIA	FIT	-0.35	0.15			FROM FIT (ERROR INCLUDES SCALE FACTOR OF 2.3)	5/82	
XIA	J	FIT DISCUSSED IN NOTE ON K _{L3} FORM FACTORS ABOVE.						
XIA	J	SHARPLE 64 PHYSICS 11					1/74	
XIA	C	CALLAN 66 TABLE 1 (PI AMU) GIVES DX/DL=1.72-0.017(C-0.04)-12.					1/74	
XIA	C	ERROR RAISED FROM .80 TO AGREE WITH DX/DL=37 FOR FIXED L ₊ .					1/74	

Data Card Listings

Stable Particles
K₀

For notation, see key at front of Listings.



12 SHORT-LIVED NEUTRAL (K₀NS), JPD=1 1=12

12 KOS MEAN LIFE (UNITS 10⁻¹⁰=10 SEC)

Table with columns for T, C, M, F, B, M, H, T, AVG and various data points for experiments like CREWFIELD 59 HBC, KAUFMAN 61 HBC, etc.

COMMENTS
M HILL AB HAS BEEN CHANGED BY THE AUTHORS FROM THE PUBLISHED VALUE 11/72
M IQ, 0.60+-0.001 GIVEN AS A CORRECTION IN THE SHEET DUE TO EXTRA...

12 KOS PARTIAL DECAY MODES

Table with columns for P1, P2, P3, P4, P5, P6, P7 and decay modes like KOS INTO PI+ PI-, KOS INTO PI0 PI0, etc.

12 KOS BRANCHING RATIOS

Table with columns for R1, R2, R3, R4, R5, R6, R7 and branching ratios for various decay channels like KOS INTO PI+ PI- / TOTAL, etc.

WEIGHTED AVERAGE = 0.316 +/- 0.014
ERROR SCALE BY 1.5

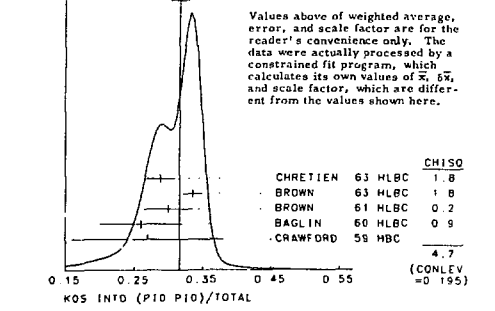


Table with columns for P3, P4, P5, P6, P7, P8, P9, P10, P11, P12 and various particle decays like KOS INTO PI+ PI- (PI0 PI0), KOS INTO PI+ PI- (K₀NS), etc.

Table with columns for K5, K6, K7, K8 and decays like KOS INTO PI+ PI- (K₀NS), KOS INTO PI+ PI- (K₀NS), etc.

Table with columns for K9, K10, K11, K12 and decays like KOS INTO PI+ PI- (K₀NS), KOS INTO PI+ PI- (K₀NS), etc.

Table with columns for K13, K14, K15, K16 and decays like KOS INTO PI+ PI- (K₀NS), KOS INTO PI+ PI- (K₀NS), etc.

REFERENCES FOR K₀

Table of references for K₀ decays, listing authors and publications like BAIRD 58 PRL 1 150, BOLDT 58 PRL 2 296, BAGLIN 60 NC 18 1043, etc.

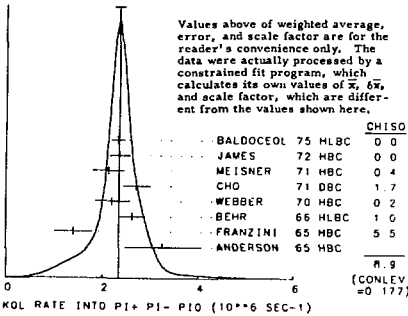
Data Card Listings

For notation, see key at front of Listings.

Stable Particles

K_L

WEIGHTED AVERAGE = 2.34 ± 0.11
ERROR SCALED BY 1.2



#3 KOL INTO PI E NEUTRINO I UNITS (10**6 SEC-1) (G4)
#3 7.52 0.89 0.72 AUBERT 65 HBC DS=00:CP ASSUMED 8/67
#3 620 7.81 0.56 CHU 71 HBC 2/72

#3 AVG 7.71 0.40 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)
#3 FIT 7.47 0.11 FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.3)

#4 KOL INTO CHARGED (3-BODY) UNITS (10**6 SEC-1) (G2+G3+G4)
#4 98 15.1 1.9 AUERBACH 66 CSPK 8/67

#4 FIT 15.10 0.21 FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.4)

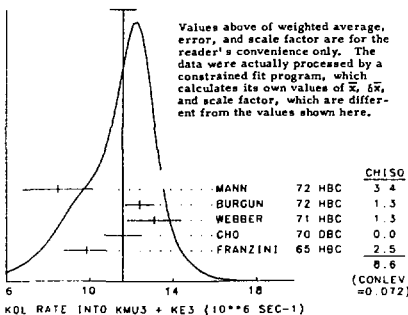
#5 KOL INTO LEPTONIC (KMU+K) UNITS (10**6 SEC-1) (G3+G4)
#5 0 109 9.85 1.15 1.05 FRANZINI 65 HBC 2/72

#5 C 335 (10.33) (0.8) HILL 67 DBC K4 TO K0 P 8/67
#5 D 393 11.0 0.9 CHO 70 DBC K4 TO K0P 10/70
#5 E 257 13.1 1.3 WEBBER 71 HBC K- P TO K0BAR N 2/72
#5 F 410 12.4 0.7 BURGUN 72 HBC K- P TO K0P1 1/73
#5 G 126 8.47 1.60 HANN 72 HBC K- P TO K0BAR N 1/72

#5 C CHO TO INCLUDES EVENTS OF HILL 67
#5 O ASSUMES DS=00 RULE

#5 AVG 11.00 0.65 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.5)
#5 FIT 12.70 0.18 FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.4)
(SEE IDEOGRAM BELOW)

WEIGHTED AVERAGE = 11.60 ± 0.65
ERROR SCALED BY 1.5



#6 KOL INTO PI MU NEUTRINO UNITS (10**6 SEC-1) (G21)
#6 19 4.54 1.24 1.08 LEWIS 67 HLCB 8/67

#6 FIT 5.232 0.086 FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.3)

13 KOL BRANCHING RATIOS

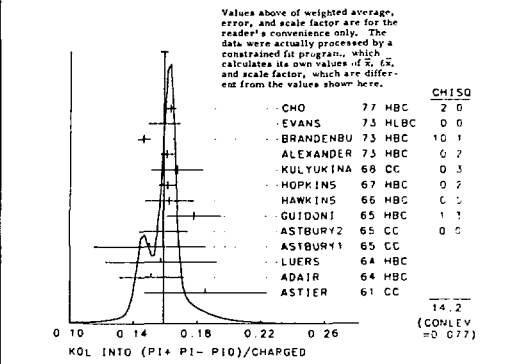
R1 KOL INTO (PI0 P10)/CHARGED (P11/(P2+P3+P4)) 6/66
R1 26 0.24 0.09 ANIKINA 68 CC 6/66
R1 349 0.251 0.014 BURDQD 68 HLCB ORSY MEASUR. 10/68
R1 444 0.277 0.021 BURDQD 68 HLCB EC. PDY TC. MEAS 10/68
R1 29 0.31 0.07 0.06 KULYUKINA 68 CC 2/71

R1 AVG 0.260 0.011 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.3)
R1 FIT 0.274 0.016 FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.7)

R2 KOL INTO (PI+ PI- P10)/CHARGED (P21/(P2+P3+P4)) 8/66
R2 59 0.185 0.038 ASTIER 41 CC 8/66
R2 79 0.151 0.020 ADAP 64 HBC 8/66
R2 75 0.157 0.03 LUERS 64 HBC 8/66
R2 86 0.15 0.03 0.04 ASTURBY 65 CC 8/66
R2 326 0.159 0.015 ASTURBY2 65 CC 8/66
R2 568 0.178 0.017 HOPKINS 65 HBC 8/66
R2 1729 (10.44) (10.04) HOPKINS 65 HBC SEE HOPKINS 67 8/66
R2 126 0.162 0.015 HAWKINS 66 HBC 8/66
R2 3200 0.161 0.006 HOPKINS 67 HBC 8/66
R2 1402 0.167 0.016 KULYUKINA 68 CC 2/71
R2 1500 0.1605 0.0088 ALEXANDER 73 HBC 1/74
R2 558 0.159 0.010 EVANS 73 HLCB 1/73
R2 6499 0.163 0.003 CHU 77 HBC 11/77

R2 AVG 0.1587 0.0024 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.3)
R2 FIT 0.1584 0.0020 FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.2)
(SEE IDEOGRAM BELOW)

WEIGHTED AVERAGE = 0.1587 ± 0.0024
ERROR SCALED BY 1.3



R3 KOL INTO (PI MU NEUTRINO)/CHARGED (P31/(P2+P3+P4)) 7/66
R3 C 291 (0.356) (0.7) LUERS 64 HBC 7/66
R3 C 172 (0.39) (0.08) 10.101 ASTURBY 65 CC 2/71
R3 C 330 (0.351) (0.055) KULYUKINA 68 CC 2/71
R3 C THIS VALUE MEASURED INDEPENDENTLY FROM R2 AND R4

R3 FIT 0.3466 0.0028 FROM FIT

R4 KOL INTO (PI MU NEUTRINO)/CHARGED (P41/(P2+P3+P4)) 2/76
R4 24 0.46 0.11 HEAGU 41 CC 2/76
R4 153 0.487 0.05 LUERS 64 HBC 7/66
R4 202 0.46 0.03 0.10 ASTURBY 65 CC 2/71
R4 300 0.498 0.052 KULYUKINA 68 CC 2/71

R4 AVG 0.485 0.032 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.3)
R4 FIT 0.4951 0.0029 FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.0)

R5 KOL INTO (PI E NEU)/(PI E NEU)/(PI MU NEU) (P43/(P3+P4))
R5 320 0.415 0.120 ASTIER 61 CC
R5 FIT 0.5882 0.0032 FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.0)

R6 KOL INTO (PI+ PI- P10)/TOTAL (P21/(P2+P3+P4))
R6
R6 FIT 0.1239 0.0020 FROM FIT

R7 KOL INTO (PI NEUTRINO)/TOTAL (P3+P4)
R7
R7 FIT 0.6505 0.0087 FROM FIT

R8 KOL INTO (E GAMMA)/TOTAL (10**6 SEC-1) (P51) 8/66
R8 C 11 (0.6) CRIGEE 66 CSPK REFL. CRIGEE66 11/68
R8 32 6.7 2.2 TODOROFF 67 CSPK REFL. CRIGEE66 11/68
R8 K 35 (7.4) (1.6) CROWIN 67 CSPK MDM-TU 3P1E+1 2/71
R8 90 5.5 1.1 68 DISK MDM-TU 3P1E+1 2/71
R8 23 4.5 1.0 ENSTRON 71 CSPK KOL 1.5-9 C/VJ 11/71
R8 B 5.0 0.11 REPELLEN 71 CSPK 11/71
R8 B 4.54 0.84 BANNER 72 HSK 8/72
R8 B THIS VALUE USES (EQUATION)=100-1.05-0.14 IN GENERAL. 5E38 = 8/72
R8 B (4-32-0.551110M+1111ED)/E IN GENERAL. 5E38 = 8/72
R8 R ASSUMES REGEN AMPL IN COPPER AT 2E5 IS 22 MU TO EVALUATE 11/71
R8 R FOR A GIVEN REGEN AMPL AND ERROR, MULTIPLY BY REGEN AMPL/2248 ** 11/71
R8 C CRIGEE 66 REPLACED BY TODOROFF 67 11/68
R8 K CROWIN 67 REPLACED BY KUMI 68. 2/71

R8 AVG 0.89 0.54 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)
R8 FIT 0.821 0.0024 FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.0)

R9 KOL INTO (PI+ PI-)/CHARGED UNIT (10**3) (P51/(P2+P3+P4))
R9 45 (2.0) (0.8) CHRISTENS 66 CSPK ERA = 1.95-0.23 2/76
R9 54 (2.05) (0.55) GILBERT 65 CSPK ERA = 1.90-0.16 2/76
R9 0 (1.93) (1.0) BASILE 66 CSPK ERA = 1.92-0.13 2/76
R9 0 (1.95) (1.0) BOTT-BODD 66 CSPK ERA = 1.93-0.24 2/76
R9 M4200 (2.63) (0.07) HESKNER 73 CSPK ERA = 2.23-0.05 8/73

R9 D 000 EXPERIMENTS EXCLUDED FROM FIT. SEE SUBSECTION E-4 BELOW FOR
R9 M FROM SAME DATA AS R27 MEASUR. BUT WITH DIFFERENT NORMALIZATION. 6/73

R9 FIT 2.569 0.000 FROM FIT

Stable Particles

K_L⁰

Data Card Listings

For notation, see key at front of Listings.

R10 KOL INTO IPI MU NEU/(PI E NEU) (P131/P14)
 R10 0.81 0.19 ADRI 64 HBC 6/66
 R10 0.87 0.10 DEBOUARD 67 DSPK 11/67
 R10 273 0.7 0.2 HAWKINS 67 HBC 8/67
 R10 0.81 0.08 MURPINS 67 HBC 8/67
 R10 770 0.7 0.05 BUDAGOV 68 HLCB 10/68
 R10 K (10.571) (0.133) KULYUMINA 68 CC 3/74
 R10 569 10.71 0.041 BELLEFLEUR 69 HLCB 10/69
 R10 1309 (10.448) (0.303) EVANS 69 HLCB REPL. BY EVANS 73 1/73
 R10 3548 0.68 0.08 BASILE 70 DSPK 10/70
 R10 1309 0.74 0.046 BRANDEBU 73 HBC 1/74
 R10 1309 0.682 0.030 EVANS 73 HLCB 1/73
 R10 1309 0.682 0.037 WILLIAMS 73 HBC 10/74
 R10 39K 0.702 0.011 HOBBS 80 HBC 2/82
 R10 K KULYUMINA 68 IS NOT MEASURED INDEPENDENTLY FROM R2 AND R4.
 R10 B BELLISSEAU 69 IS A SEAMING EXPT USING SAME EXPOSURE AS BUDAGOV 68
 R10 * * * * *
 R10 AVG 0.7001 0.0093 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)
 R10 FIT 0.0001 0.0002 FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.0)

R11 KOL INTO IEMUW-1/CHARGED (UNITS 10⁰⁰-6) (P61/P2P3P4)
 R11 (100.0) CR LESS ANIKINA 65 CC 6/66
 R11 (250.0) CR LESS CL-90 BLFF-STEL 66 DSPK 9/66
 R11 (25.0) CR LESS CL-90 BOTT-RODE 67 CSPK 6/67
 R11 (35.0) CR LESS CL-90 FITCH 67 DSPK 3/68
 R12 KOL INTO IPI MU NEU/(PI E NEU) (P131/P14)
 R12 (15.0) CR LESS ANIKINA 65 CC 6/66
 R12 0 (15.0) CR LESS BFLDITZ 66 HLCB GAM KE 40-130 MV 8/67
 R12 1 (3.0) CR LESS NEUMENS 66 DSPK GAM KE 120 MEV 6/66
 R12 (10.4) CR LESS CL-90 WATCHER 68 DSPK GAM KE 20-170 MV 2/71
 R12 (1.2) CR LESS CL-90 BOBBIOT 74 HLCB GAM KE 40-100 MEV 12/75
 R12 D 24 (10.021) (0.021) DONALDS 74 SPEC 10/74
 R12 (10.463) CR LESS CL-90 WOO 74 SPEC 12/75
 R12 M 518 (10.252) (0.036) CARROLL 72 SPC +-GDAM KE GT 20 MEV 12/80
 R12 F 548 (10.028) (0.028) CARROLL 72 SPC +-GDAM KE GT 20 MEV 12/80
 R12 K1062 0.0441 0.0032 CARROLL 80 SPC +-GDAM KE GT 20 MEV 12/80
 R12 D USES KOL TO P1P1-P10/ALL KOL DECAYS = 0.1239
 R12 M INTERNAL BREMSSTRAHLUNG COMPONENT ONLY. 12/80
 R12 F CORRECT FOR GAMMA EMISSION COMPONENT ONLY. 12/80
 R12 K BOTH COMPONENTS. USES KOL TO P1P1-P10/ALL KOL DECAYS = 0.1239 12/80

R13 KOL INTO IPI E-1/CHARGED (UNITS 10⁰⁰-6) (P71/P2P3P4)
 R13 (100.0) CR LESS ANIKINA 65 CC 6/66
 R13 (250.0) CR LESS CL-90 BLFF-STEL 66 DSPK 9/66
 R13 (25.0) CR LESS CL-90 BOTT-RODE 67 CSPK 6/67
 R14 KOL INTO IPI MU NEU/(PI E NEU) (P131/P14)
 R14 (15.0) CR LESS ANIKINA 65 CC 6/66
 R14 (1.0) CR LESS CL-90 CARPENTER 66 DSPK 6/67
 R14 (12.0) CR LESS CL-90 BOBBIOT 74 HLCB 10/74
 R14 0.08 0.08 0.08 FITCH 67 DSPK 3/68

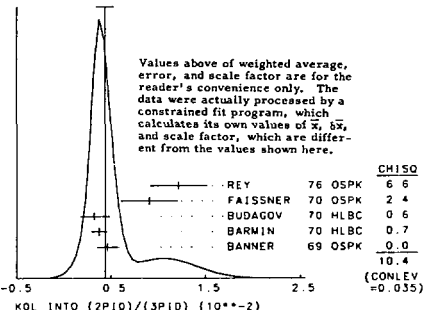
R15 KOL INTO IPI E-1/CHARGED (UNITS 10⁰⁰-6) (P71/P14)
 R15 0 97 (10.90) (0.18) LEUNG 61 CC 6/66
 R15 0 94 (11.08) (0.16) NEU 64 HBC 8/66
 R15 0 84 (12.94) (0.23) KULYUMINA 68 CC 3/74
 R15 0 1539 (11.06) (0.05) WEHNEY 66 DSPK 6/67
 R15 0 LOW PRECISION EXPTS NOT AVERAGED. FOR MORE PRECISE VALUE.
 R15 0 SEE S132 IN THE CP VIOLATION SECTION BELOW.
 R16 KOL INTO IPI MU NEU/(PI E NEU) (P131/P14)
 R16 1M 1.0081 0.0027 DORFAN 67 DSPK 11/67
 R16 SEE ALSO S132 AND S134 IN THE CP VIOLATION SECTION BELOW. 2/71

R17 KOL INTO IPI0 IPI0/TOTAL (UNITS 10⁰⁰-3) (P11)
 R17 C 7 (1.21) (1.5) (1.2) CREAGER 66 DSPK 7/66
 R17 C CRIDGE EXPT NOT DESIGNED TO MEASURE D P10 DECAY MODE 5/69
 R17 C 149 (2.5) (0.8) GALLARD 69 DSPK 5003.64-0.6 10/70
 R17 G LATEST RESULT OF THIS EXPERIMENT GIVEN BY FAISSNER TO R19 1/71
 R17 FIT 0.46 0.19 FROM FIT
 R18 KOL INTO IPI0 IPI0/(P1P1-P10) (P11/P12)
 R18 186 2.0 0.6 ALEKSANYA 64 FBC 9/66
 R18 1910 1.80 0.13 BUDAGOV 68 HLCB 10/68
 R18 883 (1.051) (0.07) BARNER 72 HLCB ERROR STAT. ONLY 3/74
 R18 * * * * *
 R18 AVG 1.88 0.13 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)
 R18 FIT 1.73 0.10 FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.0)

R19 KOL INTO IPI0 IPI0/(P1P1-P10) (UNITS 10⁰⁰-2) (P11/P11)
 R19 C 109 (1.89) (0.31) CROMIN 1 67 CSPK ET400.49+-0.5 6/67
 R19 C (1.36) (0.18) CROMIN 2 67 CSPK ET400.392+-0.3 11/67
 R19 C CROMIN IS FROM ANALYSES CROMIN 1 AND BOTH WITHORHAM 12/80
 R19 C NO EVENTS SEEN. BARTLEY 68 CSPK SEE EDD BELOW 11/68
 R19 R 57 0.46 0.31 BANNER 69 DSPK ET400.22+-0.3 2/72
 R19 R 133 (1.31) (0.31) CEMSE 69 DSPK ET400.37+-0.5 10/74
 R19 29 0.37 0.08 BARNIN 70 HLCB ET400.20+-0.23 12/70
 R19 19 0.32 0.15 CEMSE ET400.19+-0.5 10/70
 R19 F 172 0.20 0.30 FAISSNER TO DSPK ET400.32+-0.5 12/70
 R19 R 150 1.21 0.30 REY 76 CSPK ET400.38+-0.5 8/76
 R19 F FAISSNER TO CONTAINS SAME 2P10 EVENTS AS GALLARD 69 RIT 1/77
 R19 R CENCE 99 EVENTS ARE INCLUDED IN REY 76.
 R19 AVG 0.43 0.09 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)
 R19 FIT 0.437 0.085 FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.0)
 R19 FIT 1.73 0.10 (SEE IDEOGRAM BELOW)

R20 KOL INTO IPI P1-EKES + K⁰NI (UNITS 10⁰⁰-3) (P51/P3P4)
 R20 0 309 (2.5) (0.21) DEBOUARD 67 DSPK ETR=2.00+-0.09 2/76
 R20 0 525 (2.35) (0.19) FITCH 67 DSPK ETR=1.94+-0.08 2/76
 R20 0 7703 3.04 0.14 JEVD 77 SPC ETR=2.25+-0.05 11/77
 R20 0 OLD EXPERIMENT EXCLUDED FROM FIT. SEE SUBSECTION E+ BELOW FOR 2/76
 R20 D AVERAGE ETR-- IF THESE EXPERIMENTS AND FOR NOTE ON DISCREPANCY. 2/76
 R20 FIT 3.076 0.075 FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.0)

WEIGHTED AVERAGE = 0.437 ± 0.092
 ERROR SCALED BY 1.6



Values above of weighted average, error, and scale factor are for the reader's convenience only. The data were actually processed by a constrained fit program, which calculates raw own values of Σ , Σ^2 , and scale factor, which are different from the values shown here.

R21 KOL INTO I2GAMMA1/3 P101 (UNITS 10⁰⁰-3) (P1P1/P11)
 R21 10 2.5 0.7 ARNOLD 68 HLCB NEUMON DECAF 11/68
 R21 BANNER 69 IS NEW EXPT. NOT TO BE CONF. WITH REY OR CROMIN 67 2/72
 R21 115 2.13 0.43 BANNER 69 DSPK 11/68
 R21 26 2.13 0.43 BARNIN 71 HLCB 8/71
 R21 AVG 2.24 0.22 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)

R22 KOL INTO IEMUW-1/PIEPI1 (UNITS 10⁰⁰-1) (P2/P1P5)
 R22 0 (14.0) CR LESS CL-90 FUEH 69 SPC 5/70
 R22 0 (18.7) CR LESS CL-90 DARRILLAT 70 SPC 11/70
 R22 A 0 (1.53) CR LESS CL-90 CLARK 71 SPC 2/76
 R22 C 9 5.8 2.3 1.5 CARITHERS 73 SPC 2/76
 R22 F 3 4.2 5.1 2.6 FUKUSHIMA 76 SPC 2/76
 R22 15 4.0 1.4 0.9 SHOCHET 79 SPC 7/79
 R22 A CLARK 71 LIMIT PASSED FROM L.2 E-06 BY FIELD 74 REANALYSIS. 2/76
 R22 A NOT IN AGREEMENT WITH CARITHERS. NOT AVERAGED. 2/76
 R22 C CARITHERS 73 ERRORS ARE AT CL=0.68, W. CARITHERS, PRIV. COMM. 1979. 2/76
 R22 F FUKUSHIMA 76 ERRORS ARE AT CL=90 PERCENT. 2/76
 R22 AVG 4.47 0.95 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)

R23 KOL INTO IPI E-1/PIEPI1 (UNITS 10⁰⁰-5) (P71/P14)
 R23 0 (10.0) CR LESS CL-90 FUEH 69 ASPK 5/70
 R23 0 (1.0) CR LESS CL-90 CLARK 71 SPC 6/71
 R24 KOL INTO IEMUW1/PIEPI1 (UNITS 10⁰⁰-5) (P8/P14)
 R24 A (10.0) CR LESS CL-90 CLARK 71 ASPK 6/71
 R24 A POSSIBLE (BUT UNKNOWN) SYSTEMATIC ERRORS. SEE NOTE A IN R22 ABOVE. 4/82

R25 KOL INTO IPI E-1/NEU GAMI/IKL ES (UNITS 10⁰⁰-2) (P21/P14)
 R25 10 3.3 2.0 PEACH 71 HLCB GAM KE GT 15 MEV 6/71
 R26 KOL INTO IPI 0 TWO GAMMA51/3P101 (UNITS 10⁰⁰-3) (P131/P14)
 R26 0 (1.1) CR LESS CL-90 BANNER 69 DSPK 2/72
 R27 KOL INTO IPI P1-JEAL (UNITS 10⁰⁰-2) (P51/P12)
 R27 400 1.64 0.06 MESSNER 73 ASPK ETA = -2.23 6/73
 R27 FIT 1.635 0.055 FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.0)

R28 KOL INTO IPI E-1/GAMMA/TOTAL (UNITS 10⁰⁰-5) (P14)
 R28 B 0 (2.73) CR LESS CL-90 BARNIN 72 HLCB 12/80
 R28 C 4 1.74 0.87 CARROLL 80 SPC +- 12/80
 R28 B USES KOL TO P1P1-P10/ALL KOL DECAYS = 0.1239 12/80
 R28 C USES KOL TO P1P1-P10/ALL KOL DECAYS = 0.1239 12/80
 R29 KOL INTO IPI MU NEU/(PI E NEU) (UNITS 10⁰⁰-6) (P14)
 R29 D 17.811 CR LESS CL-90 DONALDS 74 SPC 6/77
 R29 C 1 0.28 0.28 CARROLL 80 SPC +- 12/80
 R29 D USES KOL TO P1P1-P10/ALL KOL DECAYS = 0.126 6/77
 R29 C USES KOL TO P1P1-P10/ALL KOL DECAYS = 0.126 12/80
 R30 KOL INTO IEMUW-1/PI0/TOTAL (UNITS 10⁰⁰-5) (P16)
 R30 D (15.66) CR LESS CL-90 DONALDS 74 SPC 6/77
 R30 C 0 0.12 0.12 CR LESS CL-90 CARROLL 80 SPC 12/80
 R30 D USES KOL TO P1P1-P10/ALL KOL DECAYS = 0.126 6/77
 R30 C USES KOL TO P1P1-P10/ALL KOL DECAYS = 0.1239 12/80

R31 KOL INTO IPI P1-E-1/TOTAL (UNITS 10⁰⁰-6) (P17)
 R31 (30.0) CR LESS ANIKINA 73 SPC 3/77
 R31 0 8.41 CR LESS CL-90 BLOWM 74 SPC 12/80
 R31 D USES KOL TO P1P1-P10/ALL KOL DECAYS = 0.126 6/77
 R32 KOL INTO IPI0 P1+- E-1 NEU/TOTAL (UNITS 10⁰⁰-3) (P18)
 R32 16 (2.2) CR LESS CL-90 DONALDS 74 SPC 6/77
 R32 16 0.082 0.020 CARROLL 80 SPC 12/80
 R32 D DONALDS 74 USES KOL TO P1P1-P10/ALL KOL DECAYS = 0.126 6/77

R33 KOL INTO IPI MU ATOM NEU/TOTAL (UNITS 10⁰⁰-7) (P19)
 R33 18 SEEN (DOBBS 76 WIRE) 6/77
 R34 KOL INTO IPI E-1/PI0/TOTAL (UNITS 10⁰⁰-6) (P20)
 R34 C 0 2.3 CR LESS CL-90 CARROLL 80 SPC 12/80
 R34 C USES KOL TO P1P1-P10/ALL KOL DECAYS = 0.1239 12/80

Stable Particles

Data Card Listings

For notation, see key at front of Listings.

LO LAMBDA 0 LINEAR ENERGY DEPENDENCE OF F0 IN KMUS DECAYS
WHEREVER POSSIBLE, WE HAVE CONVERTED THE ABOVE VALUES OF X(10) INTO
VALUES OF LO USING THE ASSOCIATED LWR AND DX/DLO.

LO FIT DISCUSSED IN NOTE ON KL3 FOM FACTORS IN K+- SEC. OF DATA CDS.
LO L VALUE IS FOM L=0.05 CALCULATED BY US FROM X(10) AND DX/DLO.

LME LAMBDA 0 LINEAR ENERGY DEPENDENCE OF F0 IN K0 DECAYS
FOR RAD. E.C.C. OF KE3 DP SEE GINSBURG 67 AND DECKMANN 70.

WEIGHTED AVERAGE = 0.0300 +/- 0.0016
ERROR SCALED BY 1.2

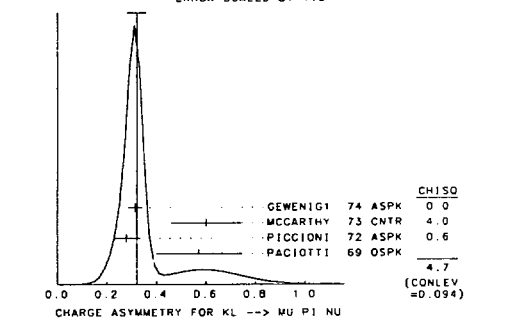
LAMBDA-0 FOR KE3 DECAY OF KOL
RATIO OF SCALAR F+ F- COUPLINGS FOR KE3 DECAYS (ABS. VALUES)
RATIO OF TENSOR TO F+ COUPLINGS FOR KE3 DECAYS (ABS. VALUES)

13 CP VIOLATION PARAMETERS IN KOL DECAYS
RELATED TEXT SECTION VI B.3 AND MINI-REVIEW BELOW

JTD COEFF OF TERM (S1-S2I) (M2I) IN MATRIX ELEMENT DEFINED AT BEGINNING
OF SECTION 6.0 ABOVE. SEE ALSO MINI-REVIEW ON SLOPE PARAMETERS IN
CP CHG K SECTION AND TEXT SEC. VI B.4. THIS SECTION REPLACES CHANGE

SUCH ASYMMETRY VIOLATES CP - IT IS RELATED TO REAL/REPSILON.
KOL INTO (M+P1-N1) (E+P1-N1) (M0+P1-N1) (E0+P1-N1) (PERCENT)

WEIGHTED AVERAGE = 0.319 +/- 0.038
ERROR SCALED BY 1.5



A2 KOL INTO (E+P1-N1) (E-P1-N1) (E+P1-N1) (E-P1-N1) (PERCENT)
A2 B 10M 10.224 10.363 BENNETT 67 CNTR KE3 11/67
A2 B 10M 0.246 0.059 SARL 69 CNTR KE3 10/70

AL B KOL INTO ((L-)(L-)) ((L+)(L+)) ((COMBINED A1 AND A2) (PERCENT)
AL B 10M 0.246 0.059 SARL 69 CNTR KE3 2/71
AL D 1M -5.7 0.17 PACIOTTI 69 OSKP KMUS 1/73

13 PARAMETERS FOR KOL INTO 2PI DECAY
TEXT SECTION VI B.3 C

ETA+ = A(KL TO P1) (P1) / (A(KL TO P1) (P1))
ETA0 = A(KL TO P1) (P1) / (A(KL TO P1) (P1))
THE FITTED VALUES OF ETA+ AND ETA0 GIVEN BELOW ARE THE RESULTS
OF THE FIT TO ETA+ AND ETA0 AND ETAD/ETA+ RESULTS. THE VALUES LISTED

FTM FT/F+ RATIO OF TENSOR TO F+ COUPLINGS FOR K0 DECAYS (ABS. VALUES)
FTM (0.121) (0.121) BIRULEV 81 SPEC 2/824

Data Card Listings

For notation, see key at front of Listings.

Stable Particles

K_L^0

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EOS (ETAD0)*2 = (CALC TO ZP01)A(KS TO ZP01)*2 (UNITS 10**+3) ---
EOS K 45 (14.95) (0.20) CHRISTENSEN 60 CSPK 10769
EOS K 57 (14.91) (1.2) BANNER 60 CSPK 2712
EOS KR 133 (14.11) (3.4) GELBER 60 CSPK 10769
EOS KF 180 (13.1) (4.1) GAILLARD 60 CSPK 10769
EOS X 29 (14.081) (0.9) BARNIN 70 HEGC 12770
EOS K 30 (13.911) (1.4) BOGARD 70 HEGC 10769
EOS C 8.7 3.7 CHOLLET 70 CSPK CU REG..4 GAMMAS 2712
EOS KF 172 (9.9) (3.4) FAISSNER 70 OSPK 12770
EOS C 26 (7.4) 2.0 WOLF 71 OSCM CU REG..4 GAMMAS 12771
EOS KR 150 (14.11) (3.4) REY 70 CSPK CU REG..4 GAMMAS 8776
EOS C 54.3 0.96 CHARLES 70 CSPK 12770
EOS K 1.3 0.3 CHARLES 72 WAVE OR SCALE FACTOR 1.3 47829
EOS X SEE NOTE ABOVE REGARDING FITTED VALUES OF ETA- AND ETAD0.
EOS B CENCE 60 EVENTS ARE INCLUDED IN THE 10769
EOS F FAISSNER TO CERTAINS SAME ZP01 EVENTS AS GAILLARD 60 1777
EOS C CHOLLET TO GIVES ETAD0=(1.23+0.24)*MGEN AMPL..ZGVCU CU31240D0M 2772
EOS C WOLF 71 GIVES ETAD0=(1.35+0.12)*MGEN AMPL..ZGVCU CU31240D0M 2772
EOS C WJ COMPUTE BOTH ETAD0*2 VALUES FOR (REGEN AMPL..ZGVCU CU31241*2M) 1772
EOS C THIS WAVEF AMPL. RESULTS FROM AVERAGING OVER FAISSNER 59. 1772
EOS C ESTABLISHED USING OPTICAL MODEL CALCULATIONS OF MURPHY ET AL. 1772
EOS C PL 270 594 (1968) AND THE DATA OF BALAIS 71 (FROM W. FAISSNER). 1772
EOS C PRIVATE COMMUNICATIONS. 1772
EOS F FAISSNER TO CERTAINS SAME ZP01 EVENTS AS GAILLARD 60 1777
EOS
EOS AVG 54.5 0.96 AVERAGE ERROR INCLUDES SCALE FACTOR OF 1.01 47829
EOS FIT 54.1 0.98 FROM FIT ERROR INCLUDES SCALE FACTOR OF 1.11 47829
EOS THIS FIT VALUE CORRESPONDS TO ETAD0=2.325+0.092
E--
E-- ETAD- = 1.31M TO P1+P1-JA(KS TO P1P1-J) UNITS 10**+3 -----
E-- F 56 (11.99) (0.16) GALBRAITH 65 LPM 2776
E-- F (11.92) (0.15) BASILE 65 LPM 2776
E-- F (11.95) (0.14) NOT-BURN 60 CSPK 2776
E-- F (7.00) (0.09) GEDRUM 57 LPM 2776
E-- F (11.94) (0.08) FITCH 67 CSPK 2776
E-- F (10.95) (0.09) POLJAKS 71 WAVE EXP.15, P1(KS 71 2776
E-- A AVERAGE OF ABOVE EXPERIMENT. THESE ARE EXCLUDED FROM THE ANALYSIS. 1780
E-- A AVERAGE AND FIT IS GIVEN BELOW IN THE COMMENTS SECTION WITH WAVE 1780
E-- A REFER PRECISE AND IN PRINCIPLE SUPERIOR EXPERIMENTAL. 1780
E-- Y 4290 (2.23) (0.05) MESSNER 73 ASPK 11775
E-- X 2.10 0.20 73 ASPK 11775
E-- X 2703 (2.24) (0.05) VEVE 73 SPEC 11775
E-- X 2.21 0.12 CHRISTEY 79 ASPK 12775
E-- X 2.255 0.029 CALVERTS 73 WAVE OR SP. AFTER 71 62644
E-- X SEE NOTE ABOVE REGARDING FITTED VALUES OF ETA- AND ETAD0.
E--
E-- AVG 2.272 0.022 AVERAGE ERROR INCLUDES SCALE FACTOR OF 1.01 47829
E-- FIT 2.274 0.022 FROM FIT ERROR INCLUDES SCALE FACTOR OF 1.11 47829
E#
E# RATIO OF F1AD0 EVEF LTA--
E# 124 1.03 0.07 BANNING 72 LPM 6776
E# 107 1.00 0.26 HOLDER 72 6776
E# C 11.000 (0.091) CHRISTIE 79 ASPK 2780
E# C NOT INDEPENDENT OF F1- AND F0S VALUES WHICH ARE INCLUDED IN FIT. 2780
E#
E#
E# AVG 1.013 0.045 AVERAGE ERROR INCLUDES SCALE FACTOR OF 1.01 47829
E# FIT 1.017 0.036 FROM FIT ERROR INCLUDES SCALE FACTOR OF 1.11 47829
E--
E-- PHASE OF F1A -- (DEGREES)
E-- THE DEPENDENCE OF THE PHASE OF THE KOL-K0S MASS DIFFERENCE
E-- IS GIVEN FOR EACH EXPERIMENT IN THE COMMENTS BELOW, WHERE THE IS
E-- (MASS DIFF. IN UNITS FORMED SEC.-1) WE HAVE EVALUATED THE
E-- MASS DEFERENCE USING OUR APRIL 1982 VALUE, DM=0.5269+0.0022
E-- TO OBTAIN THE VALUES AND AVERAGE GIVEN BELOW. WE ALSO GIVE THE
E-- REGENERATOR PHASE IN THE COMMENTS.
E--
E-- U (145.0) (50.0) FITCH 65 USPK RE REGEN 1277
E-- C 133.0 (45.0) FRENCH 66 HEGC 11769
E-- D 170.0 (21.0) NOT-BURN 60 CSPK C REGEN 1776
E-- D 125.0 (35.0) MISCHE 67 CSPK CU REGEN 1298
E-- D "TO EXPERIMENTALS WHOSE LARGE ERRORS NOT INCLUDE IN THE ABOVE 1298
E-- H (111.0) (11.0) BENNETT 68 LTMV OR REG. USES 8768
E-- C 36.2 10.0 BENNETT 68 CNTR CU REGEN 2774
E-- F 42.8 12.0 JENSON 60 CSPK WAVEFORM REGEN 2774
E-- F 45.2 1.4 FAISSNER 60 ASPK CU REGEN 2771
E-- J 42.8 6.2 JENSON 70 ASPK WAVEFORM REGEN 2771
E-- D 37.2 12.0 HALAIS 71 CSPK LC REGEN 5771
E-- P 36.2 6.1 CARNEGIE 72 ASPK CU REGEN 1273
E-- G 46.9 1.6 CARNEGIE 76 ASPK SACRIM REGEN 1276
E-- H 45.5 2.4 CARITHER 75 SPEC E REGEN 1175
E-- H 41.7 2.9 CHRISTEY 79 ASPK 12775
E--
E-- AVG 44.6 1.2 AVERAGE ERROR INCLUDES SCALE FACTOR OF 1.01 47829
E-- FIT 44.6 1.2 FROM FIT ERROR INCLUDES SCALE FACTOR OF 1.11 47829
E--
E-- COMMENTS
E-- Y BENNETT 68 IS A REEVALUATION OF BENNETT 68. 11769
E-- C BENNETT 68 WAS MEASUREMENT OF F1A-COMES IN ALL 11 ENDEAVOR DEC. 11769
E-- C BENNETT 68 IS (34.98+10.01+89710M-3451) DEG. FR=49.98+9.4 DEG. 2771
E-- B BURN 69 FR=441.00+12.91+479470M-5261 DEG. 2771
E-- F FAISSNER 60 IS ADDED TO INCLUDE WAVEFORM. WAVEFORM PHASE. 2771
E-- F FAISSNER 60 FR=149.3+7.61+205910M-5451 DEG. FR=42.74+9.0 DEG. 2771
E-- J JENSON 60 FR=142.44+0.01576+910M-5301 DEG. FR=42.74+9.0 DEG. 2771
E-- D HALAIS 71 FR=106.0+12.01+910M-5441 DEG. FR=41.98+8.0 DEG. 2771
E-- P CARNEGIE 72 IS IN UNITS IN 3M. FR=46.2+8.2 DEG. 1273
E-- G CARITHER 75 FR=145.4+2.81+22610M-5341 DEG. FR=41.98+9.4 DEG. 1175
E-- H CARITHER 75 FR=145.4+2.81+22610M-5341 DEG. FR=41.98+9.4 DEG. 1175
E#
E# PHASE OF F1A 00 (IN DEGREES)
E# FIRST QUADRANT REFERRED COBRI 60 CSPK 11669
E# 51 0.0 0.0 CHOLLET 70 CSPK CU REG..4 GAMMAS 12771
E# M 54 38.9 25.0 WOLF 71 OSCM CU REG..4 GAMMAS 12771
E#
E# AVG 55.7 4.9 AVERAGE ERROR INCLUDES SCALE FACTOR OF 1.01 12770
E# FIT 55.7 4.9 CHRISTEY 79 ASPK 12770
E# C CHOLLET 70 LIVES WAVEFORM PHASE FR=46.2+8.4 DEG. 1273
E# M WOLF 71 LIVES WAVEFORM PHASE FR=46.2+8.4 DEG. 1273
E#
E#
E# FIT 54.5 0.96 FROM FIT ERROR INCLUDES SCALE FACTOR OF 1.01 47829
E--
E-- PHASE DIFFERENCE FOR F1- = (DEGREES)
E-- F 1.6 (8.9) CHRISTIE 79 ASPK 11775
E-- C (12.0) (6.2) CHRISTIE 79 ASPK 2780
E-- B INDEPENDENT OF BENNETT MEASUREMENT AND EFFECTS. 1773
E-- C NOT INDEPENDENT OF F1- AND F0S VALUES WHICH ARE INCLUDED IN FIT. 2780
E--
E-- FIT 1.8 0.2 FROM FIT ERROR INCLUDES SCALE FACTOR OF 1.11 47829

```

Sec. VI B(d), the superweak model¹ predicts that²

$$|\eta_{00}/\eta_{+-}| = 1,$$

$$\phi_{+-} = \phi_{00} = \tan^{-1} \left(\frac{2\text{Im}T}{\text{Re}T} \right)$$

$$\text{Re}T = |\eta_{+-}| \left[1 + \left(\frac{2\text{Im}T}{\text{Re}T} \right)^2 \right]^{-1/2}$$

The latter two expressions and the values of the $K_L^0 - K_S^0$ mass difference $\Delta m = (0.5349 \pm 0.0022) \times 10^{10}$ sec⁻¹, the K_S^0 mean life $\tau_S = (0.8921 \pm 0.0027) \times 10^{-10}$ sec, and the magnitude of the $K_{L^0}^0 \rightarrow \pi^+ \pi^0 / K_{L^0}^0 \rightarrow \pi^+ \pi^-$ amplitude ratio $|\eta_{+-}| = (2.274 \pm 0.022) \times 10^{-3}$, all from the current edition, result in the predictions that

$$\phi_{+-} = \phi_{00} = (41.67 \pm 0.14)^\circ$$

$$\text{Re}T = (1.645 \pm 0.016) \times 10^{-3}$$

The above predictions can be compared with the experimental values

$$|\eta_{00}/\eta_{+-}| = 1.023 \pm 0.036,$$

$$\phi_{+-} = (44.6 \pm 1.2)^\circ,$$

$$\phi_{00} = (54.5 \pm 5.3)^\circ,$$

$$\text{Re}T = (1.621 \pm 0.088) \times 10^{-3},$$

where $\text{Re}T$ has been computed using the relation

$$\text{Re}T = \frac{i}{2} \left(\frac{1 - x^2}{1 + |x|^2} \right),$$

and our current values of the charge asymmetry parameter for leptonic K_L^0 decay $\delta = (0.330 \pm 0.012) \pm 0.004 \pm 0.026$.

The superweak predictions are in agreement with the data except for the measured value of ϕ_{00} , which is two standard deviations above the prediction.

This results primarily from the CHRISTENSON 79 measurement $\phi_{00} = (55.7 \pm 5.8)^\circ$.

References

1. I. Wolfenstein, Phys. Lett. **13**, 562 (1964).
2. T. D. Lee and L. Wolfenstein, Phys. Rev. **138B**, 1490 (1965).

Superweak Model Predictions for $|\eta_{00}/\eta_{+-}|$, ϕ_{+-} and $\text{Re}T$

In terms of the parameters defined in the text,

Stable Particles

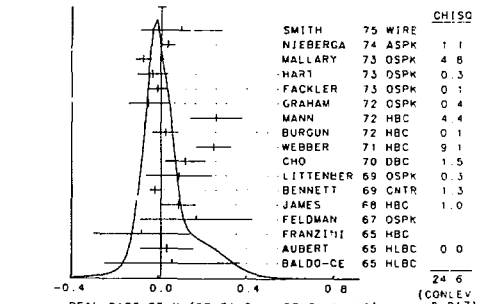
Data Card Listings

For notation, see key at front of Listings.

K0
RELATD TEXT SECTION VI 8.4
13 X = IDS+DD AMPLITUDE/IDS+DD AMPLITUDE

RELATD TEXT SECTION VI 8.4
REX REK PART OF X
REX C 152 0.06 0.18 0.44 BALDO-CE 65 HLCB
REX 196 0.035 0.11 0.13 AUBERT 65 HLCB
REX F 209 -0.08 0.16 0.24 FRANZINI 65 HLCB
REX 116 0.17 0.16 0.35 FELDMAN 67 DSPK
REX N 335 (0.17) (0.10) HILL 67 DBC
REX B (0.03) (0.03) BENNETT 68 CNTR
REX 121 0.08 0.07 0.09 JAMES 67 DBC
REX B -0.020 0.025 BENNETT 69 CNTR
REX 686 0.00 0.14 0.16 WEBBER 72 HBC
REX N 215 0.12 0.09 0.16 CHD 70 HBC
REX U 222 (0.04) (0.07) (0.08) BURGUN 71 HBC
REX 552 0.25 0.10 0.07 WEBBER 71 HBC
REX U 410 0.07 0.06 0.06 BURGUN 72 HBC
REX 126 0.24 0.10 0.14 MANN 72 HBC
REX G 362 (0.13) (0.11) GRAHAM 72 DSPK
REX G 100 (0.04) (0.10) (0.13) GRAHAM 72 DSPK
REX G 42 -0.05 0.09 0.10 WEBBER 72 DSPK
REX 1757 -0.08 0.044 FACKLER 73 DSPK
REX 1367 -0.03 0.07 HART 73 DSPK
REX 1076 -0.070 0.036 MALLERY 73 DSPK
REX 4724 0.04 0.03 NIEBERGA 74 ASPK
REX 79 0.10 0.18 SMITH 75 HBC
REX C BALDO-CE 65 GIVES X AND META-CONVERTED BY US TO REK AND IMX. 11/67
REX F FRANZINI 65 GIVES X AND META-CONVERTED BY US TO REK AND IMX. 11/67
REX N CHO TO 15 ANALYSIS OF UNBIASED EVENTS IN NEW DATA AND HILL 67. 11/73
REX U BURGUN 72 IS A FINAL RESULT WHICH INCLUDES BURGUN 71. 11/73
REX B BENNETT 67 IS A FINAL RESULT WHICH INCLUDES BENNETT 68. 10/69
REX G SECOND GRAHAM 72 VALUE IS FIRST GRAHAM 72 VALUE COMBINED WITH 2/72
REX G WANTSCH 72. 2/72
REK AVE -0.009 0.020 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.4)
(SEE IDEOGRAM REFLN.)

WEIGHTED AVERAGE = 0.009 ± 0.020
ERROR SCALED BY 1.4



SMITH 75 WIRE
NIEBERGA 74 ASPK
MALLERY 73 DSPK
HART 73 DSPK
FACKLER 73 DSPK
GRAHAM 72 DSPK
MANN 72 HBC
BURGUN 72 HBC
WEBBER 71 HBC
CHO 70 HBC
LITTENER 69 DSPK
BENNETT 67 CNTR
JAMES 67 HBC
FELDMAN 67 DSPK
FRANZINI 65 HBC
AUBERT 65 HLCB
BALDO-CE 65 HLCB
REK AVE -0.009 0.020 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.4)
(SEE IDEOGRAM REFLN.)

BAROON 58 ANP 5 156
CRAWFORD 59 PRL 2 361
ASTER 61 AIZ 001 1 227
FITCH 61 NC 2 110
GODD 61 PR 124 1223
CHRISTEN 64 PRL 13 138
ALSO 61 JETP 13 1138
CAMERINI 62 PR 128 357
DARMON 62 PL 3 2
ADAIR 64 PL 2 107
ALEKSANY 64 DDNA 2 167
ALSO 64 JETP 10 1010
ANIKINA 64 JETP 19 42
CHRISTEN 64 PRL 13 138
FUJII 64 DDNA 2 146
LUERS 64 PR 133 B 1276
M BARDON-K LONDE-L LEDERMAN (ICOLUBA+BNL)
CRAWFORD-CRESTE-R, DOUGLASS-GODD (ILL)
ASTER-BALASKOV-ARWET-SILAU D (EPOL)
V FITCH-PRIKRE-R PERKINS (PRINCETON)
GODD-MATSEN-MULLER-WICCIANO-POMELA (ILL)
MEAGOW-ORNDORF-PETROW-RIZENOV-BAKOV-LITNER
ITAGI, OKUNOY, PETROW, RIZENOV, NUSKOV, LIJNER
CAMERINI-FRY-GATDOS-ERICE-ELY (CNIS-CERN)
J DARMON-A ROUSSET-T (EPOL)
R K ADAIR-L B LETPNER (VAL-EBNL)
ALEKSANYAN, ALIKHANYAN, VARTANYAN (EREVAN)
ALEKSANYAN, LEBEDYANSKY EVD (MOSCOW)
ANIKINA, ZHURAVLVA-IGORING ACAD SGT (DUSNA)
CHRISTEN, CHRISTEN, FITCH, TUCKER (EPOL)
FUJII, JOVANOVIĆ, TURKOTI (BNL-MARYLAND, MIT)
LUERS, WITTHA, WILLIS, VANAMOTO (BNL)

ANIKINA 65 JYR 2 P48B
ANDERSON 65 PRL 14 475
ASTBURY 65 PRL 18 178
ALSO 65 HELV, PH, AC, 39 523
ASTBURY2 65 PL 18 178
ASTBURY3 65 PL 18 178
AUERT 65 PL 17 59
ALSO 65 LOMVS
BALDO-CE 65 NC 38 884
CRAWFORD 65 PRL 2 361
FISHER 65 ANL 7130 63
FITCH 65 PL 15 73
FRANZINI 65 PR 140 B 127
GALBRAITH 65 PL 28 383
LITTENER 65 PR 140 B 127
HOPKINS 65 ARGONNE CONF 67
VISHNEVSKY 65 PL 18 339
ALFF-STE 66 PL 21 957
ANIKINA 66 SMP 2 339
AUERBACH 66 PRL 17 980
AUERBACH 66 PR 149 1052
ALSO 66 ANL 14 192
BALDO-CE 66 NC 454 73
BASILE 66 BALATON CONF.
BERN 66 PL 22 540
BELLOTTI 66 NC 454 737
BOIT+BOEN 66 PL 23 217
CAMERINI 66 PRL 17 944
CANTER 66 PRL 17 944
CARPENTER 66 PR 142 871
CHAND 66 PL 23 102
CHICOEE 66 PRL 17 150
FIRESTONE 66 PRL 17 116
FIRESTONE 66 PRL 17 116
FUJII 66 PRL 17 116
FUJII 66 IS THE CORRECTED
HANKINS 66 PL 21 238
ALSO 67 PR 156 1444
JOVANOVIĆ 66 PRL 17 1075
KUKUKUNIN 66 PR 18 1611
MEISNER 66 PR 16 278
MEISNER 66 PRL 17 492
NEFKENS 66 PL 19 100
VERHEY 66 PRL 17 660
BENNETT 67 PRL 19 993
BOIT+BOEN 67 PL 248 194
BOIT+BOEN 67 PL 248 194
ALSO 66 PL 20 212
BOIT+BOEN 67 PL 248 194
CANTER 67 THESIS 15
CRONIN 67 PRL 18 253
CRONIN 67 PRINC CONF 11/67
DEBOUARD 67 NC 524 662
ALSO 67 PL 15 1644
DEVILIN 67 PRL 18 54
ALSO 68 PR 169 1045
DORFAN 67 PRL 19 987
KUKUKUNIN 67 PR 18 1611
FIRESTONE 67 PRL 18 170
FITCH 67 PR 164 371
HANKINS 67 PR 15 1944
HILL 67 PRL 19 660
HOPKINS 67 PRL 19 185
KADVA 67 PRL 19 347
KUKUKUNIN 67 PR 18 1611
LOMVS 67 PL 248 75
MISCHKE 67 PRL 18 133
NEFKENS 67 PR 157 1238
TODOROFF 67 THESIS
ABRAMS 68 PR 176 1003
ARNOLD 68 PL 286 56
ARNDSON 68 PRL 20 787
ALSO 68 PR 175 1708
BALZAT 68 PL 268 320
BARTELE-CHEVREY-VISHNEVSKY-GALINAVA (ITEP)
BASILE 68 PR 268 320
BASILE 68 PR 268 320
BENNETT 68 PR 178 244
BENNETT 68 PR 178 244
BLANPFLIED 68 PR 21 1650
BUDAGOV 68 NC 174 162
ALSO 68 PR 156 1114
CARMICHAEL 68 PRINC 1044 THESIS
F JAMES 68 PR 365
ALSO 68 PRL 21 257
KUKUKUNIN 68 PR 18 1611
HANKINS 68 THESIS (PU 66)
MELCHOR 68 PR 172 1013
WITTCHE 68 PR 174 1674
BANNER 68 PR 188 2033
BANNER 68 PR 21 1103
ALSO 68 PR 21 1107
ALSO 68 PR 208 202
BENNETT 68 PL 256 317
BOHM 68 PR 89 605
ALSO 68 PR 278 319
BOIT+BOEN 69 CERN 660 329
69 CERN 660 329
EVANS 69 PL 23 427
FALKNER 69 PL 23 427
FOSCH 69 PL 208 262
GALLMAN 69 NC 594 453
ALSO 67 PRL 18 20
LONDE 69 PL 22 885
LITTENER 69 PRL 22 885
GODD 69 PR 181 1808
MELCHOR 69 PR 172 1013
ALSO 69 THESIS 15
GALBRAITH, HUSSRI, JARVIS (CERN, BNL, JALHSH)
KRITEN, GALBRAITH, HUSSRI (CERN, BNL, AACH)
GORT+KAKEL-MOFFETT-ROSEN-GODD (ICOLUBA)
LITTENER-FIELD-PI-CHEVREY (EPOL)
M J LONGO-K YOUNG, J H HELL AND HICH, WUBLO
MELCHOR 69 PR 172 1013
ALSO 69 THESIS 15
J SALL (ICOLUBA)
*ABRISON, BALDO-CEOLIN, AUERT+ (PADDO, EPOL)
BELLOTTI, PHILIP, BALDO-CEOLIN (MILAN, EPOL)
BOIT+BOEN-HAUSSEN, DEBOUARD, CASSEL (CERN)
CAMERINI, CHLOE, ENGLISH, DEBOUARD, CASSEL (CERN)
CANTER, CHLOE, ENGLISH, DEBOUARD, CASSEL (CERN)
CARPENTER, ABRAMS, ABRAMS, FISHER (ILL INDI)
CHANG, BASTAND, FRAUCHI, OGDON (YERKOVSKY)
*OD, FRAUENFELDER, HANSON, MOSCATI (ILL INDI)
FIRESTONE, KIM, LACH, SANDWESS (VAL, BNL)
FUJII, JOVANOVIĆ, TURKOTI, ZORN (BNL+MARYLAND)
VALUE GIVEN BY JOVANOVIĆ 66 (VAL)
C J B HAWKINS
C J B HAWKINS
JOVANOVIĆ, FUJII, TURKOTI, ZORN (BNL+MANNHEIM)
KUKUKUNIN, NEFKENS, MILLER, LINGG, PEIER (BNL)
O W MEISNER, R B CRAWFORD, P CRAWFORD (ILL)
G MEISNER, B CRAWFORD, P CRAWFORD (ILL)
NEFKENS, ABRAMS, ABRAMS, CARPENTER, FISHER (ILL)
VERHEY, NEFKENS, ABRAMS (ILL)
BENNETT, NYGREN, SALL, STEINBERGER (ICOLUBA)
BOIT+BOEN, DEBOUARD, CASSEL (CERN)
BOIT+BOEN, HAUSSEN, DEBOUARD, CASSEL (CERN)
BOIT+BOEN, HAUSSEN, DEBOUARD, CASSEL (CERN)
BOIT+BOEN, HAUSSEN, DEBOUARD, CASSEL (CERN)
*HUNZ, RISK, WHEELER (PRINCETON)
*HUNZ, RISK, WHEELER (PRINCETON)
DEBOUARD, NEFKENS, JORDAN, MERMO (CERN)
DEBOUARD, NEFKENS, JORDAN, MERMO (CERN)
DEVILIN, SLODAN, SHEPARD, BEALL (PRINCIPAL)
*YAYER, BEALL, DEVILIN, SHEPARD (UM+PRINCIPAL)
DORFAN, ENSTRÖM, RAYMOND, SCHWARTZ (VAL+BNL)
FELDMAN, FRANKEL, MESTVIR, SHILL, INEAGU (JIF)
FIRESTONE, KIM, LACH, SANDWESS (VAL+BNL)
FITCH, ROTH, RUSS, VERHOY (PRINCETON)
HANKINS
HILL, LUERS, ROBINSON, LANTER (BNL, CERNIEGIE)
HOPKINS, BACON, EISLER (BNL)
KADVA, CHAND, ORNDORF, DREN, SHELTON (ILL)
KUKUKUNIN, MESTVIR, SHILL, INEAGU (JIF)
LOMVS, AUERT, CHODNET, PASCAUD (EPOL, ORSA)
MISCHKE, ABRAMS, ABRAMS (ILL INDI)
*ABRAMS, ABRAMS, CARPENTER, FISHER (ILL)
JOHN A TODOROFF (ILL INDI)
ABRAMS 68 PR 176 1003
ARNOLD 68 PL 286 56
ARNDSON 68 PRL 20 787
ALSO 68 PR 175 1708
BALZAT 68 PL 268 320
BARTELE-CHEVREY-VISHNEVSKY-GALINAVA (ITEP)
BASILE 68 PR 268 320
BASILE 68 PR 268 320
BENNETT 68 PR 178 244
BENNETT 68 PR 178 244
BLANPFLIED 68 PR 21 1650
BUDAGOV 68 NC 174 162
ALSO 68 PR 156 1114
CARMICHAEL 68 PRINC 1044 THESIS
F JAMES 68 PR 365
ALSO 68 PRL 21 257
KUKUKUNIN 68 PR 18 1611
HANKINS 68 THESIS (PU 66)
MELCHOR 68 PR 172 1013
WITTCHE 68 PR 174 1674
BANNER 68 PR 188 2033
BANNER 68 PR 21 1103
ALSO 68 PR 21 1107
ALSO 68 PR 208 202
BENNETT 68 PL 256 317
BOHM 68 PR 89 605
ALSO 68 PR 278 319
BOIT+BOEN 69 CERN 660 329
69 CERN 660 329
EVANS 69 PL 23 427
FALKNER 69 PL 23 427
FOSCH 69 PL 208 262
GALLMAN 69 NC 594 453
ALSO 67 PRL 18 20
LONDE 69 PL 22 885
LITTENER 69 PRL 22 885
GODD 69 PR 181 1808
MELCHOR 69 PR 172 1013
ALSO 69 THESIS 15
GALBRAITH, HUSSRI, JARVIS (CERN, BNL, JALHSH)
KRITEN, GALBRAITH, HUSSRI (CERN, BNL, AACH)
GORT+KAKEL-MOFFETT-ROSEN-GODD (ICOLUBA)
LITTENER-FIELD-PI-CHEVREY (EPOL)
M J LONGO-K YOUNG, J H HELL AND HICH, WUBLO
MELCHOR 69 PR 172 1013
ALSO 69 THESIS 15
J SALL (ICOLUBA)

Stable Particles

D⁰, F[±], B, p

Data Card Listings

For notation, see key at front of Listings.

REFERENCES FOR NEUTRAL D

GOLDHABER 76 PRL 37 255
 FELDMAN 77 PRL 38 1313
 GOLDHABER 77 PRL 38 1030
 PERUZZI 77 PRL 39 1149
 PICCOLI 77 PRL 70 260

BAITAY 78 PRL 41 77
 SCHARRE 78 PRL 40 74
 VUILLEMIN 78 PRL 41 1149
 ABRAMS 78 PRL 43 481
 AMPHIESE 78 PRL 80B 115
 AITTA 78 PRL 43 414

ADAMOVICH 80 PL 898 427
 ALLASTIA 80 CP B176 13
 ASTON 80 P. 948 113
 AVERY 80 PRL 45 1709
 BACINO 80 PRL 45 329
 USHIDA 80 PRL 45 1049
 ZHOLENTZ 80 PL 908 214
 ALSO 81 YAD. PHYS. 34 1471

ADAMOVICH 81 PL 908 271
 ADAMOVICH 81 DESY-L-TRANS-260
 ALSO 80 ADMETECH
 ALSO 81 FIORINO

ADEVA 81 PL 1028 289
 BALLON 81 PR 024 7
 ALSO 80 PL 898 473
 FIORINO 81 LNC 30 164
 FUCHI 81 LNC 31 159
 SCHINDLER 81 PRO 24 78

ABE 82 PRL 46 THE PUBL. L'ESLAC HYBRID FACILITY PROTON COLLARADYTON
 REUCROFT 82 CERNAEP/0903R/5
 USHIDA 82 PRL 48 864

QUANTUM NUMBER DETERMINATIONS NOT REFERRED TO IN THE DATA CARDS

NGUYEN 77 PRL 39 262

REVIEWS

A. BARBERO-GALIERI (SLAC) (1978) (LJL)
 S. KODICKI (SLAC SUMMER INST.) (1978) (SLAC)
 J. KIRBY (FLETPUN CONF. BATAVIA, 1979) (SLAC)
 G.H. TRELLING (FLETPUN) (1979)

F[±]

74 F#-(1202J, JP)

24 F#-(1202J) MASS (MEV)

M B 412030.1 (60.) BRANDELIK 77 DASP E+ E- ECH=4.2GEV 12/77
 M B 0 2030.1 00. BRANDELIK 78 PASP E+ E- ECH=4.2GEV 1/80
 M 1.017. 25. AMNAR 80 HYPER + NEU MIDERAND 7/82
 M 2026. 50. USHIDA 80 EMUL - FINAL MU WIDERAND 7/82
 M 1.7080. 121. USHIDA 80 EMUL + FINAL MU WIDERAND 2/82
 M A 400 202C. ASTON 81 OEGE + GAMMA P->F+ 1/82
 M B BRANDELIK 77 EVENTS INCLUDED IN BRANDELIK 79 VALUES 1/80
 M A ERROR QUOTED BY ASTON 81 IS 10 MEV STAT AND <20 MEV SYST. 1/82
 M AVG 15.2 AVERAGE ERROR INCLUDES SCALE FACTOR OF 1.03

34 F#-(1202I) MEAN LIFE (UNITS 10⁻¹²-13 SEC)

T 2 2.24 2.78 1.05 USHIDA 80 EMUL + NEU MIDERAND 12/81
 T 1 (1.4) AMNAR 80 HYPER + NEU MIDERAND 1/82

34 F#-(1202J) PARTIAL DECAY MODES

P1 F+ INTO ETA P1-
 P2 F+ INTO ETA ANYTHING
 P3 INTO ETA P1- P1+ P1-
 P4 F+ INTO ETA PRIME P1- P1+ P1-
 P5 F+ INTO RHO- PHI

DECAY MASSES

5884 139
 5880 0
 5884 1394 139
 974 1394 1394 139
 76941019

34 F#-(1202J) BRANCHING RATIOS

R1 F+ INTO I(ETA 710-)/ETA ANYTHING (P1)/(P2)
 R1 A (3.091 10.06) BRANDELIK 79 DASP E+ E- ECH=4.2GEV 4/82
 R1 A DENSITOMETER INCONSISTENT WITH PARTIDGE 81 (CRYSTAL BALL) 4/82

R2 F+ INTO ETA P1- P1+ P1- (P3)
 R2 300 SEEN ASTON 81 OEGE GAMMA P->F+ 1/82
 R3 300 SEEN ASTON 81 OEGE GAMMA P->F+ 1/82

R3 F+ INTO ETA PRIME P1- P1+ P1- (P4)
 R4 30 SEEN ASTON 81 OEGE GAMMA P->F+ 1/82

R4 F+ INTO RHO- PHI (P5)
 R4 30 SEEN ASTON2 81 OEGE GAM P->F+ 1/82

REFERENCES FOR F#-(1202J)

BRANDELIK 77 PL 70B 132
 BRANDELIK 79 PL 80B 412
 AMNAR 80 PL 948 118
 USHIDA 80 PRL 45 1053
 ASTON 81 PL 100B 91
 ASTON2 81 NP B108 205
 PARTIDGE 81 PRL 47 760

BRANDELIK + (IACH+CEBY+AMB+PI4+YOKY)
 BRANDELIK + (IACH+CEBY+AMB+PI4+YOKY)
 + (IACH+PNA+SEN+ITEP+CRIC+YTH+MB+5)
 + (IACH+PNA+KORE+SECO+MCCI+MAGO+OSU+OKAY+)
 (DNN+CEN+EPOL+GL+SL+LNC+MC+SL+LAL+PMP+)
 (DNN+CEN+EPOL+GL+SL+LNC+MC+SL+LAL+PMP+)
 PARTIDGE+PECK+ (CI+HARV+PRI+STAN+SLAC)

REVIEWS

G.H. TRELLING (LJL) (80)

TRELLING 81 PRPL 75 57

B

39 BOTTOM MESON B15200. JP

SEE ALSO THE LISTING FOR BOTTOM HADRON SEARCHES.
 NEEDS CONFIRMATION, NOT ENTERED INTO TABLE.

39 B MASS (MEV)

M 5240 TO 5270 ANDREWS 80 CLEO DIRECT E+ AT UPS1451 4/82
 M 5275 OR LESS FINOCCHI 80 CUSP UPS1451 THRESHOLD 4/82

39 B PARTIAL DECAY MODES

DECAY MASSES

P1 B INTO ELECTRON ANYTHING
 P2 B INTO MUON ANYTHING
 P3 B INTO E+ E- ANYTHING
 P4 B INTO MU MU- ANYTHING
 P5 B INTO KAON ANYTHING

39 B BRANCHING FRACTIONS

R1 B INTO ELECTRON ANYTHING/TOTAL (P1)
 R1 A 0.13 0.04 BEBER 81 CLEO DIRECT E+ AT UPS1451 4/82
 R1 A THE STATISTICAL AND SYSTEMATIC ERRORS ARE EACH +0.03. 4/82
 R1 B 0.136 0.050 SPENCER 81 CUSP DIRECT E+ AT UPS1451 4/82
 R1 B THE STATISTICAL AND SYSTEMATIC ERRORS ARE +0.025 AND +0.03. 4/82
 R1 AB THE ELECTRON ENERGY SPECTRA IN BOTH BEBER 81 AND SPENCER 81 FAVOR 4/82
 R1 AB B-TIME OVER B-TIME QUARK TRANSITIONS. 4/82
 R1 AVG 0.133 +/- 0.029 AVERAGE ERROR INCLUDES SCALE FACTOR OF 1.01

R2 B INTO MUON ANYTHING/TOTAL (P2)
 R2 0.094 0.036 CHADWICK 81 CLEO DIRECT MU AT UPS1451 4/82

R3 B INTO (E+ E- ANYTHING)/TOTAL (P3)
 R3 0.05 OR LESS CL=0.90 BEBER 81 CLEO E- E+ AT UPS1451 4/82

R4 B INTO (MU+ MU- ANYTHING)/TOTAL (P4)
 R4 0.017 OR LESS CL=0.90 CHADWICK 11 CLEO E+ E- AT UPS1451 4/82

R5 B INTO (KAON ANYTHING)/TOTAL (P5)
 R5 C SEEN BRODY 82 CLEO KADONS AT UPS1451 4/82
 R5 C ASSUMING UPS1451 -> B BRAN. A TOTAL OF 3.38+0.34+0.08 KADONS 4/82
 R5 C PER UPS1451 DECAY IS FOUND THE SECOND ERROR IS SYSTEMATIC. IN 4/82
 R5 C THE CONTEXT OF THE STANDARD B-DECAY MODEL, THIS LEADS TO A VALUE 4/82
 R5 C FOR B-QUARK -> C-QUARK/B-QUARK -> ALL1 OF 1.09+0.33+0.13. 4/82

REFERENCES FOR BOTTOM MESON B15200I

ANDREWS 80 PRL 45 214
 FINOCCHI 80 PRL 45 222
 BEBER 81 PRL 46 88
 CHADWICK 81 PRL 47 171
 BRODY 82 PRL 48 1070

(EON+HARV+ITH+SYR+RICE+RUTG+WANO)
 FJNOCCHIRO+GIAMMINI+
 (HARV+ITH+SYR+RICE+RUTG+WANO+CORN)
 (GANC+RICE+RUTG+SYR+HARV+CORN+MIT+ETH)
 (IACH+CEBY+AMB+PI4+YOKY+)
 (CI+HARV+PRI+STAN+SLAC)

B

16 PROTON(938, J=1/2) I=1/2

16 PROTON MASS (MEV)

M (938.256) (1.0055) COHEN 65 RVUE 7/66
 M (938.252) (1.00523) TAYLOR 69 RVUE 7/70
 M 938.2740 (0.0027) COHEN 73 RVUE USING NEW E/H 5/74

16 ANTI-PROTON MASS (MEV)

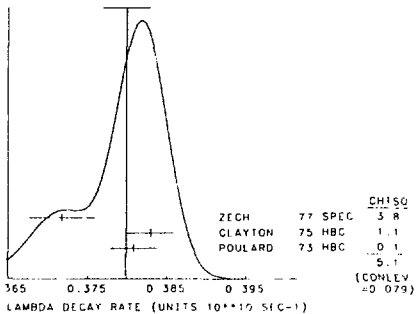
M1 938.3 D.5 BARBERO TO CNTR 12/79
 M1 938.179 0.058 HU 75 CNTR EXOTIC ATOMS 12/79
 M1 938.228 0.069 ROBERSON 77 CNTR 12/79
 M1 938.30 0.13 ROBERTS 78 CNTR 6/78
 M1
 M1 AVG 938.216 0.036 AVERAGE ERROR INCLUDES SCALE FACTOR OF 1.01

Data Card Listings

For notation, see key at front of Listings.

Stable Particles

WEIGHTED AVERAGE = 0.3799 ± 0.0029
 ERROR SCALED BY 1.6



RA	LAMBDA INTO EP	MEAS/TOTAL	(UNITS 10**4)	EP	(P1/P2)
RA	1	11.01	CR L 55	ALSTON	43 HBC
RA	1	11.01	CR L 55	KEARNS	64 FBC
RA	6	BEIN 1.3	AND 6.0	LIND	69 HBC
RA	3	1.3	0.7	LIND	64 PVUF
RA	2	1.5	1.2	RONDE	56 HBC
RA	9	1.4	0.8	CANTER	71 HBC
RA	14	1.4	0.5	BAGGETT 72	HBC STOP P-O 7771
RA					STOP K- 8772
RA	AVG	1.57	0.35	AVVERAGE (ERRR) INCLUDES SCALE FACTOR OF 1.03	

1R LAMBDA DECFN PARAMETERS
 RELATED TEXT SECTION VI D AND APPENDIX I

A-	ALPHA LAMBDA	(LAMBDA INTO PI-	PROTON)	63 ENT	LAMBDA FROM PI-
A-	1156	0.67	0.27	ERIKSS	63 HBC
A-		13.663	(0.022)	REGG	64 PVUF INCLUDES ABOVE 5/65
A-	10130	0.649	0.017	JOHNSON	67 PVUF
A-	11079	ED 7471	(0.383)	ALSTON	67 PVUF 68
A-	3520	7.67	0.06	DAUER	69 HBC FROM PI- DECAY 6/64
A-	10925	3.640	0.23	CLELAND	72 CSPM LAMBDA FROM PI-
A-	8500	0.584	0.046	ASTURP	75 SPEC LAMBDA FROM PI-
A-	AVG	1.313	0.024	AVVERAGE (ERRR) INCLUDES SCALE FACTOR OF 1.13	

1P LAMBDA - ANTILAMBDA/RV/L MEAN LIFE DIFFERENCE

01	0.044	0.085	BAGIER	67 HBC	2.4 PRAB P	R/67
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1R LAMBDA MAGNETIC MOMENT ENAGHEENS 0.98.76 MEVE

WM	-1.5	0.5	COOL	67 DSPK	
WM	0.9	0.8	KEARNS	63 HBC	
WM	0553	-1.39	0.72	ANDERSON	64 HBC
WM	151	-1.5	0.28	CHARRIET	65 EMUL
WM	40	-0.523	(0.21)	BARROV	71 EMUL PRO/TA. RESULT 2/72
WM	1300	-0.66	0.07	DANIELS	71 EMUL 46 FIFD+2006 6/71
WM	8668	-0.73	0.18	HILL	71 DSPK 10/71
WM	57	-0.65	0.28	BAROV	72 EMUL INCLUDES BARROV 71
WM	1.24	-0.57	0.05	BUNCE	76 SPEC 1/78
WM	3506	-0.549	0.37	HULL	77 SPEC 1/78
WM	34	-0.613	0.067	SCHWACHING	78 SPEC 1/79
WM	2008	-0.606	0.015	COX	81 SPEC 12/81
WM	AVG	-0.6133	0.0044	AVVERAGE (ERRR) INCLUDES SCALE FACTOR OF 1.03	

1B LAMBDA ELECTRIC DIPOLE MOMENT (UNITS 10**16 E CM)

MEASURE VALUE IMPLIES VIOLATION OF T AND P

EDM	11.01	FR LESS CL+05	GISSON	61 EMUL	7/72
EDM 6	11.01	FR LESS CL+05	BARON	71 EMUL	7/72
EDM 9	15.0131818	LESS CL+05	ANDERSON	81 SPEC	7/72
EDM B	BARON	MEASURE 15.0479310**15 E CM			7/72
EDM P	BARON P 1	MEASURE 1-3.047410**15 E CM			1/82

1R LAMBDA PARTIAL DECAY MODES

	DECAY MODES
P1	LAMBDA INTO PROTON PI-
P2	LAMBDA INTO NEUTRON PI0
P3	LAMBDA INTO PROTON M+ NEUTRON
P4	LAMBDA INTO PROTON M- NEUTRON
P5	LAMBDA INTO PROTON PI- GAMMA

1R LAMBDA BRANCHING RATIOS

R1	LAMBDA INTO EP	PI- (N+V) PI0	IPIII/(PI+P2)	
R1	7.627	0.031	CRAMFORD 59 HBC	
R1	0.65	0.09	COLUMELA 50 HBC	
R1 U	0.2693	(0.217)	ANDERSON 61 HBC	
R1	993	0.643	0.016	MUMPHREY 61 HBC
R1 U	4736	0.628	0.007	DOTTE 69 HBC
R1	4572	0.644	0.008	BALTAJ 71 HBC P-O AT REST 6/71
R1 U	ANDERSON RESULT	NOT PUBLISHED. EVENTS ADDED TO DOYLE SAMPLE.		7/71
R1	AVG	0.6199	0.0064	AVVERAGE (ERRR) INCLUDES SCALE FACTOR OF 1.03
R1	FIT	0.6419	0.0049	FROM FIT (ERRR) INCLUDES SCALE FACTOR OF 1.03

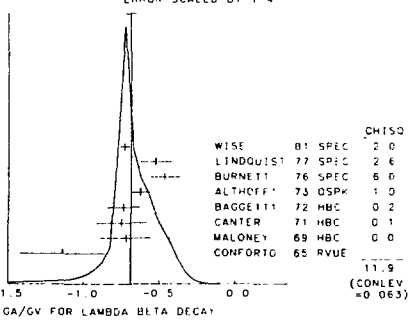
1R LAMBDA BRANCHING RATIOS

R2	LAMBDA INTO EP	PI0/PIEP	PI- (N+V) PI0	IPIII/(PI+P2)
R2	0.43	0.09	FISLE	57 HBC
R2	0.43	0.14	CRAMFORD	59 HBC
R2	0.28	0.38	BAELIN	60 HBC
R2	0.29	0.38	BRIDW	63 HBC
R2	75	0.291	0.036	CHRETIEN 63 HBC
R2	AVG	0.306	0.074	AVVERAGE (ERRR) INCLUDES SCALE FACTOR OF 1.01
R2	FIT	0.3581	0.0049	FROM FIT (ERRR) INCLUDES SCALE FACTOR OF 1.01

1R LAMBDA INTO EP - MEAS/TOTAL (UNITS 10**4)

R3	LAMBDA INTO EP	MEAS/TOTAL	(UNITS 10**4)	(P1)/(PI+P2)		
R3	0	15	(0.21)	0.51	MUMPHREY 61 PVUE	
R3	0	12	0.1	11.21	HUBERT 62 FBC	
R3	N	150	(0.82)	(0.12)	EY 63 FBC K- AT REST	
R3	N	102	(0.78)	(0.12)	10.33	BAELIN 64 FBC K- AT 1.45 -VEVC
R3	N	20	(0.58)	0.24	LIND 65 HBC	
R3	N	143	(0.70)	(0.08)	MALONEY 69 HBC	
R3	N	86	(0.78)	0.16	CANTER 71 HBC K-O AT REST	
R3	N	218	(0.88)	(0.10)	LINDQUIST 71 DSPK PI- P TO K0 LAP	
R3	N	THESE VALUES HAVE BEEN CHANGED BY US INTO PARTIAL RATIOS			7/72	
R3	N	MEASURE IS THE DIRECTLY MEASURED QUANTITY. SEE 85 BELOW			3/72	
R3	N	LOM STATISTICS EXPERIMENTS. NOT AVERAGED			7/70	

WEIGHTED AVERAGE = -0.690 ± 0.034
 ERROR SCALED BY 1.4



Data Card Listings

For notation, see key at front of Listings.

Stable Particles Σ^+

19 SIGMA+ BRANCHING RATIOS

Table with 4 columns: R#, SIGMA+ INTO (NEUTRON P+), SIGMA+ INTO (MUON P+), SIGMA+ INTO (LAMBDA E+). Includes rows for parameters like MURPHY, CHANG, BARLIFAUT, etc.

AVERAGE ERROR INCLUDES SCALE FACTOR OF 1.01

Table with 4 columns: R#, SIGMA+ INTO (NEUT P+), SIGMA+ INTO (MUON P+), SIGMA+ INTO (LAMBDA E+). Includes rows for parameters like BAZIN, BARRSH, etc.

AVERAGE ERROR INCLUDES SCALE FACTOR OF 1.01

Table with 4 columns: R#, SIGMA+ INTO (NEUT P+), SIGMA+ INTO (MUON P+), SIGMA+ INTO (LAMBDA E+). Includes rows for parameters like GERSHWIN, ANGLE, etc.

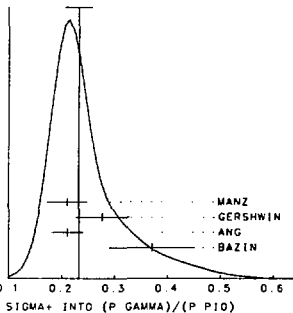
AVERAGE ERROR INCLUDES SCALE FACTOR OF 1.01

Table with 4 columns: R#, SIGMA+ INTO (NEUT P+), SIGMA+ INTO (MUON P+), SIGMA+ INTO (LAMBDA E+). Includes rows for parameters like GERSHWIN, ANGLE, etc.

AVERAGE ERROR INCLUDES SCALE FACTOR OF 1.21

Table with 4 columns: R#, SIGMA+ INTO (NEUT P+), SIGMA+ INTO (MUON P+), SIGMA+ INTO (LAMBDA E+). Includes rows for parameters like ANGLE, etc.

WEIGHTED AVERAGE = 0.232 +/- 0.025
ERROR SCALED BY 1.2



CHI/SQ

80 HBC 0.3

69 HBC 0.8

69 HBC 0.5

65 HBC 3.0

4.6

(CONV LQ = 0.207)

Table with 4 columns: R#, (SIGMA+ INTO N E-), (SIGMA+ INTO N E-), (SIGMA+ INTO N E-). Includes rows for parameters like STOP K, STOP K, STOP K, etc.

19 SIGMA+ DECAY PARAMETERS

RELATED TEST SECTION VJ D AND APPENDIX I

Table with 4 columns: A#, ALPHA/ALPHA FOR SIGMA+, (SIGMA+ TO P1), (SIGMA+ TO P10). Includes rows for parameters like COPIA, TRIP, BANGERTER, etc.

Table with 4 columns: A#, ALPHA FOR SIGMA+, (SIGMA+ INTO PROTON), (SIGMA+ INTO PROTON). Includes rows for parameters like BEALL, BANGERTER, etc.

Table with 4 columns: A#, PHA ANGLE (SIGMA+ INTO N P1), (SIGMA+ INTO N P1), (SIGMA+ INTO N P1). Includes rows for parameters like BANGERTER, BANGERTER, etc.

Table with 4 columns: A#, PHA ANGLE (SIGMA+ INTO N P10), (SIGMA+ INTO N P10), (SIGMA+ INTO N P10). Includes rows for parameters like BANGERTER, BANGERTER, etc.

Table with 4 columns: A#, PHA ANGLE (SIGMA+ INTO N P10), (SIGMA+ INTO N P10), (SIGMA+ INTO N P10). Includes rows for parameters like BANGERTER, BANGERTER, etc.

Table with 4 columns: A#, PHA ANGLE (SIGMA+ INTO N P10), (SIGMA+ INTO N P10), (SIGMA+ INTO N P10). Includes rows for parameters like BANGERTER, BANGERTER, etc.

Table with 4 columns: A#, PHA ANGLE (SIGMA+ INTO N P10), (SIGMA+ INTO N P10), (SIGMA+ INTO N P10). Includes rows for parameters like BANGERTER, BANGERTER, etc.

Table with 4 columns: A#, PHA ANGLE (SIGMA+ INTO N P10), (SIGMA+ INTO N P10), (SIGMA+ INTO N P10). Includes rows for parameters like BANGERTER, BANGERTER, etc.

Table with 4 columns: A#, PHA ANGLE (SIGMA+ INTO N P10), (SIGMA+ INTO N P10), (SIGMA+ INTO N P10). Includes rows for parameters like BANGERTER, BANGERTER, etc.

Table with 4 columns: A#, PHA ANGLE (SIGMA+ INTO N P10), (SIGMA+ INTO N P10), (SIGMA+ INTO N P10). Includes rows for parameters like BANGERTER, BANGERTER, etc.

Data Card Listings

For notation, see key at front of Listings.

Stable Particles

 Σ^-

F- PHI ANGLE (SIN(PHI)/COS(PHI)+BETA/GAMMA) (DEGREES)
 F= 0 1006 (+22.1) 130.3 BERLEY 67 HBC K-P TO SIG- P+ 11/67
 F= 1308 14 16.5 BANGERTI 69 HBC
 F= C1092 +5. 23. BERLEY TO HBC NEUTRON RESCATT. 11/69
 F= C CHANGED FROM -5 TO +5 TO AGREE WITH SIGN CONVENTION
 F= AVG 10.9 14.6 AVERAGE ERROR INCLUDES SCALE FACTOR OF 1.01

Note on $\Sigma^- \rightarrow \Lambda e \bar{\nu}$

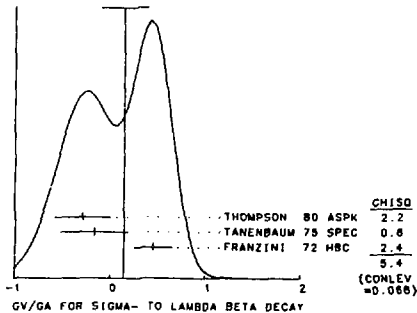
(by J. A. Thompson, University of Pittsburgh)

The decay $\Sigma^- \rightarrow \Lambda e \bar{\nu}$ is of special interest because its form is predicted by the strong form of CVC and is not sensitive to the current octet assumptions or SU_3 structure constants which enter into Cabibbo's predictions for the other hyperon decays. For $\Delta S = 0$ transitions, the weak interaction vector current is related to the electromagnetic current through a multiplicative constant, set by neutron beta decay, and an isospin rotation.

The decay $\Sigma^0 \rightarrow \Lambda \gamma$ (the isospin-rotation analogue of $\Sigma^- \rightarrow \Lambda e \bar{\nu}$) is mediated predominantly through the magnetic interaction, assuming there are no inhomogeneities in the Σ^0 , Λ charge distributions. Thus we expect the g_{WM} term [$g_{WM} \sim \frac{\mu_N \Lambda}{\sqrt{2}} \sim \frac{\sqrt{3}}{2} \mu_N$ (by SU_3)] to dominate the vector part of the weak current. The strong CVC predictions are thus: $g_V/g_A = 0$ and $g_{WM} \sim 1.6$.

AV GV/GA FOR SIGMA TO LAMBDA BETA DECAY (TEXT SEC VI D.1 FOR SIGM CONV)
 AV PREDICED TO BE ZERO BY CONSERVED VECTOR CURRENT THEORY.
 AV VALUES AVERAGED ASSUME CVC-SUB3 WEAK MAGNETISM TERM. 11/67
 AV FB 45 10.311 10.301 BARASH 67 HBC USING SIG- 11/67
 AV FS 51 10.71 10.41 BALTAY 69 HBC 4/69
 AV FS 81 140.22 10.281 EISELEZ 69 HBC 10/68
 AV F S 186 0.45 0.20 FRANZINI 72 HBC USTC SIG- 1/72
 AV 55 -0.17 0.35 TANENBAUM 75 SPEC OHL HYPERON BEAM 12/75
 AV 114 -0.29 0.29 THOMPSON 80 ASPK OHL HYPERON BEAM 1/82
 AV B BARASH AT MEASURED ABSOLUTE VALUE.
 AV S SIGN CHANGED TO AGREE WITH OUR CONVENTION.
 AV F FRANZINI 72 INCLUDES EVENTS OF BARASH 67. EISELEZ 69. BALTAY 69. 1/73
 AV AVG 0.14 0.24 AVERAGE ERROR INCLUDES SCALE FACTOR OF 1.63
 (SEE IDEOGRAM BELOW)

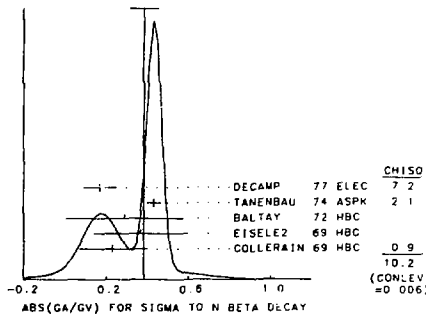
WEIGHTED AVERAGE = 0.14 ± 0.24
 ERROR SCALED BY 1.6



WM GV/GA FOR SIGMA TO LAMBDA BETA DECAY
 WM VALUES QUOTED ASSUME THE CVC PREDICTION GV=0.
 WM 186 2.4 2.1 FRANZINI 72 HBC USING SIG- 1/82
 WM 55 3.5 4.5 TANENBAUM SPEC OHL HYPERON BEAM 1/82
 WM 114 1.75 3.5 THOMPSON 80 ASPK OHL HYPERON BEAM 1/82
 WM AVG 2.2 1.7 AVERAGE ERROR INCLUDES SCALE FACTOR OF 1.01

AV1 GA/GV FOR SIGMA TO NEUTRON BETA DECAY (TEXT SEC VI D.1 FOR SIGM CONV)
 AV1 57 10.051 10.231 10.327 GERSHWIN 68 HBC POLARIZED SIGMA+ 6/68
 AV1 45 49.19 0.20 0.17 GERSHWIN 69 HBC POLARIZED SIGMA+ 10/69
 AV1 63 -0.33 0.30 0.85 BARGET TO HBC K-P AT 400 MEV/L 10/70
 AV1 45 -0.4 0.52 1.5 ELLIS 72 HBC POLARIZED SIGMA+ 10/71
 AV1 E (+40.103 10.113) ELLIS 72 PVUE SUN 11KFL-1+COL1 10/71
 AV1 E (-0.277 10.133) ELLIS 72 PVUE SUN 11KFL-1+COL1 10/71
 AV1 E ELLIS 72 INCLUDES THE MAXIMUM LIKELIHOODS OF COLLIM-1: 6% 3/72
 AV1 E EISELEZ 69. GERSHWIN 69. ELLIS 72. AND GETS TWO POSSIBLE VALUES. 3/72
 AV1 AVG 0.13 0.17 AVERAGE ERROR INCLUDES SCALE FACTOR OF 1.01
 AV2 ABSOLUTE VALUE OF GA/GV FOR SIGMA TO NEUTRON BETA DECAY
 AV2 49 0.23 0.18 COLLERAIR 69 HBC NEUTRON SCATTER 10/69
 AV2 33 0.37 0.26 0.19 EISELEZ 69 HBC NEUTRON SCATTER 10/69
 AV2 35 0.26 0.28 0.29 BALTAY 72 HBC NEUTRON SCATTER 6/72
 AV2 3507 0.435 0.035 TANENBAUM 74 ASPK H.C. HYPERON BEAM 11/77
 AV2 519 0.17 0.07 0.09 DECAP 77 ELEC 10/76
 AV2 AVG 0.385 0.070 AVERAGE ERROR INCLUDES SCALE FACTOR OF 2.33
 (SEE IDEOGRAM BELOW)

WEIGHTED AVERAGE = 0.385 ± 0.070
 ERROR SCALED BY 2.3



REFERENCES FOR SIGMA-
 BROWN 58 CERN CONF 270 BROWN, GLASER, GRAVES, PERL, CRONIN + (MICH)
 EISLER 58 NC SERIO 10 350 EISELER, PASSI, CONVERSE + (COLU, OHL, BGM, PISA)
 BARASH 61 PR 124 121 BARASH, OBER, WASSON, NICOLAS, SMITH + (MICH)
 CHIESA 81 MC 15 1171 A N CHIESA, B. QUASSI, L. RINAUDO (MILAN)
 HUMPHREY 62 PR 127 1305 H E HUMPHREY, P ROSS (MICH)
 TRIPP 62 PR 9 66 R J TRIPP, R WATSON, H FERRO, LUZZI (LBNL)
 BARASH 63 PR 11 26 W N BARASH, J N DYER, H M HECKMAN (LRL)
 BURNSTEIN 64 PR 13 60 BURNSTEIN, DRY, KEMDE, SECH, JOHNSON (UMD)
 COUANT 64 PR 136 9 291 COUANT, FILTHUTH + (CERN, HEIDENHORN, LBNL)
 MILLER 64 PR 11 282 MILLER, STANBRO, BEAUGUET + (LBNL, MIT, BERG)
 HUMPHY 64 PR 136 0 188 C THOMPSON, HUMPHY (MICH, CERN)
 NAUENBERG 64 PR 12 679 NAUENBERG, SCHMIDT, HARATECH + (COLU, PUTG, GPPN)
 BAZIN 65 PR 140 8 1358 BAZIN, PLANDS, SCHMIDT + (PREM, PUTG, COLU)
 DOSCH 65 PL 14 244 DOSCH, ENGELMANN, FILTHUTH, HEPP, KLUGER, THEISS (MICH)
 CHUNG, YUN CHANG (LUMBERT)
 SCHMIDT 65 PR 140 8 1328 BANGERTER, GALTIERI, BERG, MURPHY + (LUMBERT)
 BANGERTER 66 PR 17 495 BANGERTER, GALTIERI, BERG, MURPHY + (LUMBERT)
 CHUNG YUN CHANG (LUMBERT)
 CHEN 66 PR 151 1081 + LACH, SANDMEISS, TAIT, VEHROREN + (YALE + BNL)
 BARASH 67 PR 19 181 BARASH, DAY, GLASER, KEMDE, KOPF + (MAD, VAND)
 BERLEY 67 PR 19 579 BERLEY, HERTZBACH, KUDLER + (BNL, MASA, YALE)
 BIERMAN 68 PR 20 1459 BIERMAN, KONUSU, NAUENBERG + (PRINCETON)
 GERSHWIN 68 PR 20 1270 GERSHWIN, ALSTON-GARNJOST, BANGERTER + (LRL)
 HEPP 68 ZPHV 214 71 V. HEPP, H. SCHLEICH (HEIDELBERG)
 WHITESTEAD 68 MC 56A 537 H. WHITESTEAD, J. GOLLER (FOBERLIN)
 ANG 1 69 ZPHV 223 103 ANG, EISELEZ, ENGELMANN, FILTHUTH + (MICH)
 69 ZPHV 228 131 + BERNHARD, EISELEZ, ENGELMANN, FILTHUTH + (MICH)
 BARGETT 69 PR 23 244 BARGETT, KEMDE, SNOW (UNIV HARV, VAND)
 BALTAY 69 PR 22 615 BALTAY, FRANZINI, NEUMAN, MORTON + (COLL, STON)
 BANGERTER 69 ZPHV 19244 ROGER ODELL, BANGERTER (THEISS) (LRL)
 BANGERTI 69 PR 187 1821 BANGERTER, GARNJOST, GALTIERI, GERSHWIN + (LRL)
 BARLOTTA 69 NP D16 153 BARLOTTA, UDO, BELLEFON, GRANET + (SACL, CERN, HEID)
 COLLERAIR 69 MC 23 194 COLLERAIR, DAY, GLASER, KOPF + (LBNL, HARV, OHL)
 EISELEZ 69 ZPHV 221 3 EISELEZ, ENGELMANN, FILTHUTH, FOHLISCH, HEPP (MICH)
 EISELEZ 69 ZPHV 223 487 EISELEZ, ENGELMANN, FILTHUTH, FOHLISCH + (MICH)
 GERSHWIN 69 UCAL-19246 LAWRENCE KEMETH GERSHWIN (THEISS) (LRL)
 BERLEY 70 PR D1 2015 +YAMINI, HERTZBACH, KUDLER + (BNL, MASA, YALE)
 BODREF 70 PR D2 4 BODREF, TAIT, WELLS, BERLY + (BNL, MASA, YALE)
 EISELE 70 ZPHV 238 372 +FILTHUTH, HEPP, PRESSER + ZECH (MIDELBERG)
 BARBER 71 LNC 1 37 + BARBER, COLLAR, IZZO + (SACL + BNL + BHO + EPDL)
 COLE 71 PR D4 631 +LEE-FRANZINI, LOVELESS + BALTAY + (STON, COLU)
 HEBERT HORTON (COLUMBIA)
 DODGE 71 NP 853 493 ALSO 69 NEVIS-179 THEISS LOVE, BEL-GRADE, REIL, RYKUS, DUBTIN, MAN'S COLLAR

Stable Particles

Data Card Listings

$\Sigma^0, \Sigma^0, \Xi^-$

For notation, see key at front of Listings.

BALTY 72 PR 05 1569 ... +PEIWMAN,FRANZENI,NEWMAN,VEH ... BERLIM+BELGRADE+BRUC+DUBLIN+LOUC+HARSAM ...

BROWN 57 PR 10B 1036 ... J BROWN, O GLASER, N PERL ... M NIETO ...

Σ^0

21 SIGMA(1133, JP1/2, -) 1-1

01 N SEE NOTE PRECEDING LAMBDA MASS LISTINGS ... 01 18 4.75 0.1 BURNSTEIN 64 HBC ...

01 N SEE NOTE PRECEDING LAMBDA MASS LISTINGS ... 01 200 76.63 0.28 ...

21 SIGMA MEAN LIFE (UNITS 100ns - 10 SEC) ... T (E=14 ER LESS) DAVIS 62 EWAL ...

21 SIGMA PARTIAL DECAY MODES ... P1 SIGMA INTO LAMBDA GAMMA ...

21 SIGMA BRANCHING RATIOS ... R1 SIGMA INTO LAMBDA E+ E- TOTAL ...

REFERENCES FOR SIGMA

FEINBERG 58 PR 109 1019 ... G. FEINBERG ... DAY 62 PR 127 0059 ...

22 XI-MASS (MEV) ... M H 1113317.0 12.21 ... M H 1811317.9 1.91 ...

22 ANTI-XI+ MASS (MEV) ... M1 113322.01 11.33 ... M1 S 121321.71 1.61 ...

2. 1X1-1 - (ANTI-XI+) MASS DIFFERENCE (MEV) ... DM 1.0 1.1 CHIEN 66 HBC ...

22 XI- MAGNETIC MOMENT (MAGNETONS, 936.26 MEV) ... MM 2724 -0.1 2.1 ...

22 XI- MEAN LIFE (UNITS 10ns - 10 SEC) ... T H 11 13.51 13.41 ... T H 18 11.281 10.43 ...

22 ANTI-XI- MEAN LIFE (UNITS 10ns - 10 SEC) ... TL S 5 11.91 10.55 ... TL S 12 11.93 10.55 ...

22 XI- PARTIAL DECAY MODES ... P1 XI- INTO LAMBDA P1- ... P2 XI- INTO NEUTRON E- NEUTRINO ...

Stable Particles

ν BOUNDS, HEAVY LEPTON SEARCHES

Data Card Listings

For notation, see key at front of Listings.

NEUTRINO BOUNDS FROM ASTROPHYSICS AND COSMOLOGY

SEE THE NOTE ON NEUTRINOS BY R.L. SHERKIN IN THE ELECTRON NEUTRINO SECTION NEAR THE BEGINNING OF THESE DATA CARD LISTINGS.

Note on ν Mass Limits

These limits apply to m_{tot} given by

$$m_{tot} = \sum_{j=1}^n \left(\frac{g_{\nu_j}}{2} \right) m_{\nu_j}$$

where n is the number of neutrino species and g_{ν_j} is the number of independent components in the neutrino field; $g_{\nu_j} = 4$ for Dirac neutrinos; $g_{\nu_j} = 2$ for chiral Majorana neutrinos.

- * NU MASS EVB
- == LIMIT ON TOTAL MASS, SUMMED OVER ALL SPIN COMPONENTS, OF EFFECTIVELY STABLE NEUTRINOS m_{tot} , THESE WITH MEAN LIVES NORMALLY EQUAL TO OR GREATER THAN THE AGE OF THE UNIVERSE (UNITS EV)
- == A (130.1) OR LESS COSM 77 COSM
- == A (100.0) OR LESS LEP 77 COSM
- == A (100.0) OR LESS OLIVE 81 COSM 4/82*
- == A LOWER 72 AND LEP 77 HAVE BEEN GENERALIZED TO APPLY TO MIXTURES AS DEFINED ABOVE, WHERE ONE ALIENS EITHER DIRAC OR MAJORANA NEUTRINO.
- == A THESE SPEAKS ASSUMED DIRAC NEUTRINOS.
- == FOR OTHER LIMITS, SEE SATO 77, KYOSHIKI 77, DICUS 78, HUI 79, AND ZELDovich 80.
- == ANALYSES OF n^+ FLIGHT RATIOS AND DYNAMICS OF GALAXIES AND CLUSTERS, ARE CONSISTENT WITH PRESENCE OF DARK MATTER WHICH, IF ASSUMED TO BE NEUTRINOS, WOULD EMPLOY A NON-ZERO VALUE OF THE SUM OF THE MASSES OF NEUTRINOS WITH LIFETIMES SUFFICIENTLY LONG TO AFFECT THESE GALAXIES AND CLUSTERS. DIFFERENT ANALYSES FAVOR SOMEWHAT DIFFERENT VALUES OF THE SUM OF MASSES, THIS IS ALSO CONSISTENT WITH MODELS OF GALAXY FORMATION, ACCORDING TO O.SCHRAMM (PRINC. COMM.) IF ONE ALIENS EITHER DIRAC OR MAJORANA GALAXY FORMATION, THE RANGE OF MASSES IS PROBABLY FROM A TO 100 EV.
- == HOWEVER ADIABATIC n WOULD FAVOR m_{tot} MASS NEUTRINOS. (UNITS EV)
- == VIKENITE 79 COSM ISOTHERMAL 1/82*
- == OBER 3D ROUD 81 COSM ADIABATIC 2/82*
- == S-D 10 DAVIS 81 COSM ACIA, DECAVING NUS 1/82*
- == S-D SCHEMAM 81 COSM ISOTHERMAL 1/82*
- == A TREATISE TO STATES THAT THE DARK MATTER CANNOT BE PUON OR ELECTRON
- == A NEUTRINOS, OF HOWEVER HIGH MASS OR ANY NEUTRAL LEPTON LESS MASSIVE THAN 1 MEV.
- == 1/82*
- == 2/82*
- == 1/82*
- == 1/82*
- == 1/82*
- == 1/82*
- == 1/82*
- == 1/82*
- == 1/82*
- == 1/82*
- == 1/82*

- * NU RADIATIVE MEAN LIFE / MASS (UNITS SEC/PA)
- T COSM 77 COSM 1/82*
- T DICUS 77 COSM 1/82*
- T GOLDMAN 77 COSM 4/82*
- T FALK 78 COSM 1/82*
- T COSM 79 COSM 1/82*
- T GOLDMAN 79 COSM 1/82*
- T DEFIJULA 80 COSM 1/82*
- T STECKER 80 COSM 1/82*
- T HENRY 81 COSM 1/82*
- T KIMBLE 81 COSM 1/82*
- T TURNER 81 COSM 4/82*

- A POSSIBLE LIMITS ON NUMBER OF LIGHT ν ABOUT 1 MEV
- == 1-1) COMPOSITE NU TYPES
- == NUMBER COUPLING WITH FULL WEAK STRENGTH
- == (7) OR LESS SHVARTSMAN 69 COSM 1/82*
- == (4) OR LESS STEIGMAN 77 COSM 1/82*
- == (4) OR LESS YANG 79 COSM 1/82*
- == CRITICISM OF BOUND STECKER 80 COSM 1/82*
- == MAYBE NB FROM BOUND OLIVE 81 COSM 1/82*
- == (4) OR LESS TURNER 81 COSM 1/82*
- == CRITICISM OF BOUND RANA 82 COSM 1/82*
- == OR LESS YANG 82 COSM 1/82*
- == SEE, WHEREVER OLIVE 81 CRITIQUE AND STECKER 81 REPLY.
- == UNCERTAINTIES AND CRITICISMS COME FROM DIFFERING ESTIMATES OF LOWER LIMIT ON BARYON DENSITY OF THE UNIVERSE AND UPPER LIMIT ON THE PRIMORDIAL n PFLUM=ABUNDANCE.
- == 1/82*
- == 4/82*
- == 4/82*

- == NUMBER COUPLING WITH LESS THAN FULL WEAK STRENGTH
- == A (7) OR LESS STEIGMAN 79 COSM 1/82*
- == A LIMIT VARIES WITH STRENGTH OF COUPLING. 4/82*

- A MAGNETIC MOMENT OF SUFFICIENTLY LIGHT NU (UNITS EV/GAUSS)
- == 10⁻²⁶-197000 LESS SUTHERLAND 76 COSM FOP MNUKID REV. 1/82*
- == 5 USES SUTHERLAND 76 EQ.3 WITH F1/3 FROM THEIR TABLE AS MODIFIED TO 1/82*
- == 5 APPLY AS A LIMIT ON ANY ONE NEUTRINO SPECIES INDIVIDUALLY. 1/82*

REFERENCES FOR NU BOUNDS FROM ASPD. AND COSM.

SHVARTSMAN 69 JETPL 9 184 COSMIX 72 PR 29 669 SUTHERLAND 76 PR D13 2700	V. F. SHVARTSMAN 1 MOSUJ P. COSMIX, J. M. CLELLAND 1 UC83 SUTHERLAND, N. G. FLOWERS, + 1 PENN+COLU+NYU
COSMIX 77 PR 39 714 ALSO 79 COSMIX	R. COSMIX 1 IMPINTATFRC
DICUS 77 PR 39 148 GOLDMAN 77 PR D16 2256 HUI 77 PR 676 166 SATO 77 PR 6 1275 STEIGMAN 77 PR D63 2077 VYSOVSKEY 77 JETPL 28 188	J. R. DICUS, F. M. KOLB, V. L. TEPLITZ 1 TEXAS+PPI T. GOLDMAN, G. J. STEPHENSON 1 IASL R. W. LEE, S. M. HENRYBER 1 FINELE+SRM K. SATO, M. KOPRAYASHI 1 KYUJ G. J. STEIGMAN, D. P. SCHRAMM, J. J. OAKN 1 EVALF, CMFC, CIFF VYSOVSKEY, O. LOUGH, ZELDovich 1 IETFP
DICUS 78 PR 017 1525 FALK 78 PL 70B 511 ALSO 78 APR 271 1015	K. DOLB, TEPLITZ, WAGNER 1 TEXAS+PPI+STAJ S. FALK, D. P. SCHRAMM 1 CHEIC GUNN, LEE+ 1 CIT+CAMP+WMC+CHEIC+ALEE
COSMIX 76 PR D15 2219 GOLDMAN 76 PR D19 2215 HUI 76 PR 616 166 STEIGMAN 76 PR L 7 239 TREATISE 76 PR 62 407 YANG 76 APR 227 607 ALSO 76 STEIGMAN	R. COSMIX 1 IETFP T. GOLDMAN, G. J. STEPHENSON 1 IASL G. HUI, M. A. OLIVE, J. J. OAKN 1 FMSITERM+MEII G. STEIGMAN, M. OLIVE, D. P. SCHRAMM 1 IETFP S. TREATISE, J. J. OAKN, 1 ICLC+M+CAH YANG, S. CHARMAM, STEIGMAN, ACCD 1 ECHIC+ALE+VIRG FOOTNOTE 4
DEFIJULA 80 PR L 65 947 STECKER 80 PR L 65 1469 ALSO 81 NU 81 CPM+HAWAI STECKER 80 PR L 44 1237 ZELDovich 80 SJNP 71 864	A. DE RIJULA, S. L. GLASHOW 1 MIT+HARV F. W. STECKER 1 IASL F. W. STECKER 1 EPDOL, V. I. P. 1241 1 NAS1 F. W. STECKER 1 NAS1 ZELDovich, ALYPPIN, KHLIPOV, CHECHETKIN 1 IETFP
OLIVE 81 NU 81 CPM+HAWAI DAVIS 81 PR 46 516 HENRY 81 PR L 47 618 KIMBLE 81 PR L 46 80	J. R. OLIVE, A. S. SZLAP 1 UC83+CHEIC M. DAVIS, M. LEE, R. HODD, E. WITTENBERG, P. HUI 1 EPJ R. C. HENRY, P. O. FELDMAN 1 IJMUJ P. KIMBLE, S. BOMYER, P. JAKOBSEN 1 UC83
OLIVE 81 APR 246 557 OLIVE 81 PR 46 516 SCHRAMM 81 APR 243 141 STECKER 81 PR 46 517 TURNER 81 NU 81 CPM+HAWAI	D. P. SCHRAMM, STEIGMAN, TURNER, YANG+ 1 CHEIC+BERTI A. S. OLIVE, M. S. TURNER 1 IETFP D. P. SCHRAMM, G. STEIGMAN 1 CHEIC+BERTI F. W. STECKER 1 IASL M. S. TURNER 1 UC83+CHEIC
RANA 82 PR L 46 206 YANG 82 EPL PREPRINT	N. C. RANA 1 IETFP +TURNER+STEIGMAN, SCHRAMM, OLIVE 1 CHEIC+BERTI

HEAVY LEPTON SEARCHES

Data on the τ^\pm (1785) heavy lepton are listed in a separate section above, following the e and μ listings.

The following section contains information on searches for heavy leptons of other types and searches for the τ^\pm in collisions other than e^+e^- .

Several types of heavy leptons (that is, non-strongly-interacting fermions other than e and μ) have been proposed. In the Data Card Listings we distinguish four types.^{1,2} Each has a corresponding antiparticle with opposite charge and lepton number. For convenience we omit writing the antiparticles in the following descriptions. The four types are:

Sequential Leptons (L_\pm, ν_L)

Such a pair is assumed to have its own separately strictly conserved lepton number $n_L = +1$. This means that the radiative decays

$$\begin{array}{l}
 L^- \rightarrow e^- \gamma \\
 L^- \rightarrow \mu^- \gamma
 \end{array}
 \}
 \text{ are forbidden,}$$

while the weak decays (assuming n_L sufficiently massive)

Data Card Listings

For notation, see key at front of Listings.

Stable Particles HEAVY LEPTON SEARCHES

$$\left. \begin{array}{l} L^- + \nu_L e^- \bar{\nu}_e \\ L^- + \nu_L \mu^- \bar{\nu}_\mu \\ L^- + \nu_L \text{ hadrons} \end{array} \right\} \text{ are allowed.}$$

There could be an increasing mass sequence of such pairs. It is frequently assumed that the neutrinos are massless.

Decay rates are assumed calculable from conventional weak interactions theory. For L^- mass between 1 and 3 GeV, the branching fraction to each of the two leptonic modes should be roughly 10% to 20%. For L^- mass above 1 GeV, the mean life should be $\leq 10^{-12}$ sec, too short to be observed in a track chamber.¹

Paraleptons (E^+, E^0) and (M^+, M^0). These pairs have the same lepton numbers as the opposite-charge ordinary leptons, i.e., e^- and μ^- , respectively. Radiative decays are again forbidden and decays similar to those allowed for L^- are allowed here, e.g.,

$$M^+ \rightarrow \nu_\mu e^+ \bar{\nu}_e$$

$$\text{or } M^+ \rightarrow \nu_\mu \mu^+ \bar{\nu}_\mu$$

However, the lightest member is not stable as is the case for sequential leptons, so that bizarre decay schemes such as (assuming $m_{E^0} < m_{M^+}$)

$$E^+ \rightarrow E^0 \mu^+ \bar{\nu}_\mu$$

$$\downarrow$$

$$e^+ \bar{\nu}_e$$

are allowed.

Heavy leptons of this type (and/or a neutral intermediate boson Z^0) are desired in unified gauge theories of weak and electromagnetic interactions to cancel unphysical high energy behavior in such processes as $e^+e^- \rightarrow W^+W^-$.

Ortholeptons (F^- and \bar{N}^-). These have the same lepton numbers as e^- and μ^- , respectively. They may or may not have associated neutral leptons. Radiative decays are allowed in addition to weak modes similar to those of sequential leptons. The radiative mode can dominate or can be relatively unimportant depending on the model.⁴ Decays such as

$$F^- \rightarrow e^- + \text{hadrons}$$

are also allowed.

Long-Lived Penetrating Particles.

Heavy leptons could have long mean lives under certain circumstances. For example, if $m_{\nu_L} > m_{\nu_e}$, then L^- , the sequential lepton, is completely stable since its lepton number is conserved.

Experimental results. The results are summarized in the Data Card Listings below. Mass limits for sequential leptons are listed in subsection MS, while all other types are listed together in subsection M.

The Listings also contain cross-section upper limits reported as results of unsuccessful searches. We no longer list cross sections for anomalous $e\mu$ events in e^+e^- collisions. These cross sections are consistent with coming from $e^+e^- \rightarrow \tau^+\tau^-$ where the τ^+ (1785) is assumed to be a spin-1/2 Dirac point particle with a mass about 1785 MeV.

References

1. M. L. Perl and P. Rapidis, SLAC-PUB-1496 (October 1974).
2. C. H. Llewellyn Smith, Invited paper presented at the Royal Society Meeting on New Particles and New Quantum Numbers, 11 March 1976, Oxford Ref. 33/76.
3. J. D. Bjorken and C. H. Llewellyn Smith, Phys. Rev. D7, 887 (1973).
4. F. Wilczek and A. Zee, Nucl. Phys. B106, 461 (1976).

SEE PERL B1 FOR A REVIEW

PROPERTIES OF THE τ (1785) HEAVY LEPTON AND ITS ASSOCIATED NEUTRINO ARE LISTED SEPARATELY ABOVE FOLLOWING THE E AND M LISTINGS. THE FOLLOWING SECTION CONTAINS INFORMATION ON SEARCHES FOR HEAVY LEPTONS OF OTHER TYPES AND SEARCHES FOR τ (1785) COLLISIONS OTHER THAN e^+e^- . WE LIST MASS LIMITS AND CROSS SECTION UPPER LIMITS REPORTED AS NEGATIVE SEARCH RESULTS. WE NO LONGER LIST CROSS SECTIONS FOR THE ESTABLISHED PROCESS $e^+e^- \rightarrow \tau^+\tau^-$ AS WAS DONE IN OUR 1977 SUPPLEMENT.

MS HEAVY LEPTON MASS LIMITS

LIMITS APPLY ONLY TO HEAVY LEPTON TYPE GIVEN IN COMMENT AT RIGHT ON DATA CARD. SEE REVIEW ABOVE FOR DESCRIPTION OF TYPES.		3/77
IN COMMENTS BELOW: ALL BEAMS ARE MU TYPE NEUTRINO OR ANTI-NEUTRINO.		3/77
L=L+MUF, N STAND FOR SEQUENTIAL LEPTON, PARA-ELECTRON, PARA-MUON, ORTHO-ELECTRON, ORTHO-MUON RESPECTIVELY.		3/77
MS	SEQUENTIAL HEAVY LEPTON MASS LIMITS (GeV)	
MS A	(13.) ER MODE AZIMOV 80 μ^- SEQUENTIAL (L)	2/82*
MS B	10. ER MODE CL-95 BARBER 80 CNTR μ^- SEQUENTIAL (L)	9/81*
MS C	NONE AGEV TO 14.5GEV CL-95 BERGER 81 PLUT μ^- SEQUENTIAL (L)	1/82*
MS D	NONE BELOW 15.5 GEV CL-95 BRANDELIK B1 TASS μ^- SEQUENTIAL (L)	1/82*
MS A	AZIMOV 80 ESTIMATED PROBABILITIES FOR $\mu\mu$ TYPE EVENTS IN $e^+e^- \rightarrow$	1/82*
MS A	L- L- DEDUCING SEMI-HADRONIC DECAY MULTIPLICITIES OF L FROM e^+e^-	2/82*
MS A	ANIMULATION DATA AT MCFE/231MVL OBTAINED ABOVE LIMIT COMPARING	2/82*
MS A	THESE WITH e^+e^- DATA (BRANDELIK 80, PL 928 1993).	2/82*
MS B	BARBER 80 LOCKED FOR $e^+e^- \rightarrow L^+L^-$ (L- \rightarrow NEU(L)*X) WITH MARK-J AT	9/81*
MS B	DESY-PETRA.	9/81*
MS C	BERGER 81 IS DESY DORTS AND PETRA EXPT. LOOKING FOR $e^+e^- \rightarrow L^+L^-$.	1/82*
MS D	BRANDELIK 81 IS DESY PETRA EXPT. LOOKING FOR $e^+e^- \rightarrow L^+L^-$.	1/82*

Stable Particles
HEAVY LEPTON SEARCHES

Data Card Listings

For notation, see key at front of Listings.

Table of experimental results for heavy lepton searches, including entries for BEHMEND, BETHOUNE, BOWDITT, etc., with columns for experiment name, parameters, and results.

Table of experimental results for heavy lepton searches, including entries for CLARK 91, HAYES 82, NEU 1, etc., with columns for experiment name, parameters, and results.

Stable Particles

FREE QUARK SEARCHES

Data Card Listings

For notation, see key at front of Listings.

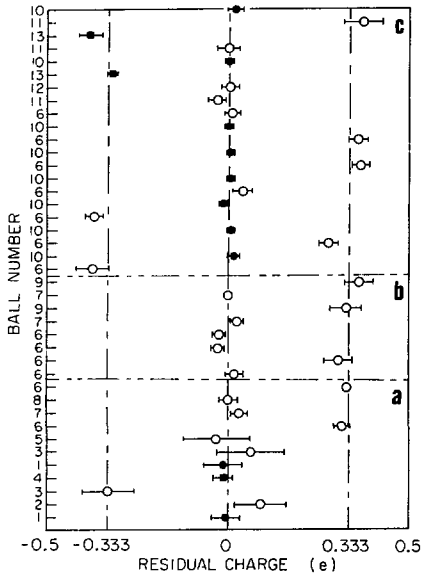


Fig. 1. Residual charge measurements in chronological order from LARUE 77 (a), 79 (b), and 81 (c). Ball radius 140 (open circles), 116 (solid squares), and 98 μm (solid circles). Figure adapted from LARUE 81: La Rue, Phillips, and Fairbank, *Phys. Rev. Lett.* **46**, 967 (1981).

$10^{-11} \text{ cm}^{-2} \text{ ster}^{-1} \text{ sec}^{-1}$. Cross-section upper limits established from proton accelerator experiments and calculations based on production models¹ imply that free quarks, if they exist, have a mass greater than about 5 GeV. Mass limits from photon and electron beam searches are slightly lower, but more reliable, depending only on the QED calculations for quark pair production. Limits on free quark concentrations in stable matter vary enormously depending on the source of matter and the technique.

The largely negative result of quark searches does not prove that free quarks do not exist, but indicates that they are hard to find. De Rujula, Giles, and Jaffe² have considered the question of

unconfined quarks in a framework of a renormalizable, spontaneously broken version of QCD, and conclude that: (1) production cross sections are small, (2) interaction cross sections with nucleons are very large, and (3) the physical masses of quarks are probably very large. On this basis, primordial quarks would be expected to be nonintegrally charged, superheavy nucleon complexes.

We group quark searches by experimental technique - proton beams, photon beams, neutrino beams, electron beams, e^+e^- annihilations, cosmic rays, and stable matter. Proton beam experiments generally measure quark production cross sections (we quote these in section C), differential cross-section ratios (sections AF and RPI), or differential cross sections (sections IC and D). The photon beam experiment measures cross section per equivalent quanta (section DG), and the neutrino experiment measures the ratio of quark events over total events (section NEU). Searches with electron beams may measure differential cross sections (section G) and set limits on the quark mass (section M). Searches in e^+e^- annihilation present the ratio to the μ -pair cross section (section EE). Cosmic ray experiments measure quark flux (section F), and searches in stable matter measure quark concentration (section RHO). Most of the accelerator and cosmic ray experiments have searched for fractionally charged particles, but some have searched for massive stable particles which would have low velocity. The latter searches are usually sensitive to a range of charges and may appear in the section below on Other New Particle Searches.

We have relied heavily on the review of L. W. Jones³ for data prior to April 1977.

References

1. T. K. Gaisser and F. Halzen, *Phys. Rev.* **D11**, 3157 (1975).
2. A. de Rujula, R. C. Giles, and R. L. Jaffe, *Phys. Rev.* **D17**, 285 (1978).
3. L. W. Jones, *Rev. Mod. Phys.* **69**, 717 (1977).

Data Card Listings

For notation, see key at front of Listings.

Stable Particles
FREE QUARK SEARCHES

Table with multiple columns containing experimental data for quark production and cross sections. Includes entries such as 'QUARK PRODUCTION CROSS SECT. FROM PROTON BEAM EXPTS.', 'QUARK PROD. DIFF. CROSS SEC. FROM PROTON BEAM EXPTS.', and 'QUARK PRODUCTION IN NEUTRINO BEAMS'. Lists various quarks like charm, strange, and bottom, and their associated experiments and results.

Data Card Listings

For notation, see key at front of Listings. FPEE QUARK, MAGNETIC MONOPOLE SEARCHES

Stable Particles

BARTON 67 HASL 29 67
 BATHOM 67 PL 258 163
 BUHLER-1 67 NC 404 209
 BUHLER-2 67 NC 404 209
 FOSS 67 PL 258 163
 GOMEZ 67 PRL 18 1022
 KASHA 67 PRL 18 1022
 STOWER 67 PL 164 1059

BELLAMY 68 PR 166 1391
 BODRHOZE 68 NC 853 241
 BRACINSKI 68 JEP 27 91
 BRITKOWE 68 NC 474 250
 FRANZINI 68 PRL 21 1013
 GARIBOLDI 68 PR 166 1391

HAYASHI 68 JIP 46 5734
 KASHAI 68 PR 172 1297
 KASHA 68 PRL 20 217
 KASHA 68 JIP 46 5734
 RANK 68 PR 176 1435

ALLRAY 69 CC 644 75
 APT POUJ 69 DL 298 245
 APT POUJ 69 DL 308 276
 CARRINS 69 DP 186 205
 COOK 69 68 189 2052
 JAPUSHIM 69 DP 178 2058
 MCCUSKER 69 PL 23 654

BOSIA 70 NC 064 167
 CHU 70 PRL 24 1540
 CHU 70 PL 250 1540
 ELBERT 70 NP 927 217
 KASSISMER 70 PRL 24 1537
 KOTLER 70 DP 876 216
 MORPURGO 70 JIN 74 65

APT POUJ 71 NP 827 374
 CHEN 71 NC 28 447
 LILANK 71 PRL 24 1541
 MACEY 71 PRL 24 482

BUCHANANP 72 PR 76 1211
 BOMM 72 PRL 28 326
 BUTT-BOD 72 DL 408 843
 COX 72 DP 120 1211
 GROUCH 72 DP 35 2627
 DARDU 72 NC 18 317
 EVANS 72 PRL 28 326
 FOWLER 72 JPA 5 565

ALPER 73 PL 468 245
 ASHTON 73 JPS 6 577
 HICKS 73 NC 2 1272A
 LEIPNER 73 PL 31 1272A

CLARK 74 DP 010 2121
 GALEK 74 DP 09 1869
 JACQUE 74 JPS 26 3290
 ZESH 74 PRL 32 854

ALBROW 75 NP 807 180
 BAGAN 75 NP 8101 363
 HAZEL 75 NP 8101 363
 JOYANUM 75 PL 568 105
 KANSOR 75 NC 274 132

BALDIN 76 JUP 22 784
 HATCHEK 76 NC 314 543
 SILVERMAN 76 DP 014 316

AMEYASU 77 PRL 39 113
 BASILE 77 NC 404 41
 BLAND 77 PRL 30 365
 GALL THOM 77 PRL 38 1255
 LARUE 77 PRL 38 1011

ALSCO 79 LARUE
 MULLER 77 SCIENCE 174 521
 GORDONDU 77 JEP 45 457

HASILEE 78 NC 454 171
 RASTLEZ 78 NC 454 291
 RYND 78 PRL 20 289
 RYND 78 PL 728 484
 LUND 78 RAO, M2 263 19
 RUIE 78 DP 017 1646
 SCHIFFER 78 DP 017 1224
 YOKA 78 PR 018 641

RYND 79 PRL 43 1298
 RYND 79 DP 018 1298
 LARUE 79 PRL 42 147
 RYND 79 PRL 42 147
 GORDONDU 79 JEP 49 575
 STEVENSON 79 DP 020 47

HARTTEL 80 JPH 40 206
 BASILE 80 LNC 29 21
 RUSSELL 80 NP 1134 1
 WEPINELL 80 PL 948 433
 HODGES 80 PL 948 427

HODGES 81 PRL 47 1151
 ALSON 81 PRL 48 97
 LARUE 81 PRL 46 967
 WEISS 81 PL 1014 479

BAITOV 74 JUP 11 444
 JONES 76 NP 40 73
 LUYMS 80 JPS 32-40 78/80

BAITOV 74 JUP 11 444
 JONES 76 NP 40 73
 LUYMS 80 JPS 32-40 78/80

MAGNETIC MONOPOLE SEARCHES

Note on Magnetic Monopole Searches

(by W. P. Trower, Virginia Polytechnic Institute and State University)

Although the idea that magnetic monopoles might exist is suggested by the usual formulation of Maxwell's equations, no observed phenomenon requires them for its explanation.¹ Monopoles first became interesting with the assertion that a single monopole anywhere in the universe would result in electric charge quantization everywhere.² The sole predicted property of this monopole was the magnitude of its least magnetic charge $e/2\alpha$, the Dirac charge. A monopole of two Dirac charges has been suggested,³ and an electrically charged monopole, a dyon, proposed.⁴ Observed pure multi-photon showers were attributed to virtual monopole pair-production/annihilation.⁵ Monopoles become indispensable to many gauge theories, and most grand unification theories require a monopole of mass $> 5 \times 10^{15}$ GeV. Estimates of the monopole number density at the earth have been made.⁶

Experiments to detect monopoles have essentially been based on either ionization or magnetic induction. Ionization experiments have relied on the fact that an elementary relativistic magnetic charge would produce more ionization than a relativistic electrical charge. Massive monopoles would have, however, lower velocities, $\beta \sim 10^{-3}$, and would thus reduce the prospects for ionization measurements. The theory of monopole energy loss at these velocities is currently confused, but progress is being made.⁷ It is however likely that prospects for ionization identification of massive monopoles will remain diminished.

Induction experiments measure the value of the monopole magnetic charge by detecting changes in magnetic flux induced when a monopole passes through a superconductor (CABRERA 75). These measurements are independent of monopole electric charge, mass, and velocity.

REVIEW ARTICLES

ALAN HODGES 158P
 L. W. JONES 118C
 L. LUYMS, Talk at US NP Forum, Ann. Meeting

Stable Particles

MAGNETIC MONOPOLE SEARCHES

With the recent attribution of large mass to the monopole, searches in matter are less appealing as the ferromagnetic trapping energies become comparable to terrestrial gravitational binding. Accelerators do not possess sufficient energy to produce real monopoles. Cosmic ray searches hold increased promise. Possibly monopole evidence will be obtained from direct astrophysical observation.

The following compilation with the indicative experimental limits should be used as a directory to the literature.

References

- J. D. Jackson, CERN-77-17 (1977).
- P. A. M. Dirac, Proc. Royal Soc. of London **A133**, 60 (1931).
- J. Schwinger, Phys. Rev. **144**, 1087 (1966).
- J. Schwinger, Science **165**, 757 (1969).
- M. A. Ruderman and D. Zwanziger, Phys. Rev. Lett. **22**, 146 (1969).
- G. Lazarides, Q. Shafi, and T. F. Walsh, Phys. Lett. **100B**, 21 (1981).
- See, e.g., S. F. Ahlen and K. Kinoshita, to be published in Phys. Rev. D (1982).

Data Card Listings

For notation, see key at front of Listings.

F	MONOPOLE (EVENTS)	F/HX (%/CM ² HR)	COSMIC RAYS	MASS (GEV)	CHARGE (G)
F	A	< 2	MALKUS 51 EMUL	< 1-3	4/82*
F	B	< 5	CAITHERS 66 ELEC	< 1	2/75
F	C	< 1	FLEISCHER 69 PLAS	> 2	2/76
F	D	< 1	PRICE 75 PLAS	> 200	2/75
F	E	< 2	BARTLETT 81 PLAS		2/82*
F	F	< 2	BONNARD 81 GCM		2/82*
F	G	< 1	KINOSHITA 81 PLAS	> 1	2/82*
F	H	< 2	ULLMANN 81 CNTR	> E+17	2/79
F	I	< 1	CABRERA 82 INTR		4/82*
F	A	NALUS 51	SEA LEVEL ACCELERATOR/CONCENTRATOR.		4/82*
F	B	CAITHERS 66	5FA LEVEL CONCENTRATOR.		12/75
F	C	FLEISCHER 69	OBSDIAN AND NICA.		12/75
F	D	PRICE 75	ANNUNCES EVENT.		11/76
F	E	BARTLETT 81	FINAL BC MAGNET. LENAX.		7/82*
F	F	BONNARD 81	INTERSTELLAR FLUX INFERRED FROM NEUTRON STARS.		2/82*
F	G	KINOSHITA 81	CR-30 ON WHITE MOUNTAIN CA. BETA > 0.02.		2/82*
F	H	ULLMANN 81	SENSITIVE TO VERTICAL VELOCITIES 100-350 KM/SEC.		2/79*
F	I	CABRERA 82	SENSITIVE ONLY TO MAGNETIC CHARGE.		4/82*

D	MONOPOLE (EVENTS)	DENSITY	MASS (GEV)	CHARGE (G)	
D	A	< 2	GOTO 63 EMUL	< 1-3	12/75
D	B	< 2	PEIKHOW 63 CNTR		12/75
D	C	< 2	FLEISCH 69 PLAS	< 1-120	4/82*
D	D	< 2	FLEISCH 69 PLAS	> 0	12/76
D	E	< 2	SCHATTEN 70 ELEC		4/77
D	F	< 2	KOLM 69 CNTR	< 140	12/75
D	G	< 2	ROSS 73 INDO	> .05	2/76
D	H	< 2	CABRERA 75 CNTR	> .04	5/82*
D	I	< 2	CARRIGAN 76 INDO		1/77
D	J	< 2	BRODERICK 79		12/81*
D	D	A	GOTO 63	MONODRIFT MOUNTAIN MAGNETITE AND TWO METEORITES.	12/75
D	D	B	PEIKHOW 63	SIMULTE-ALIN METEORITE.	12/75
D	C	FLEISCHER 69	MN NODULES OCEAN SEDIMENTS LAST 16 E+6 YEARS.		2/82*
D	C	FLEISCHER 69	MN EARTH CRUST.		12/75
D	E	SCHATTEN 70	LUNAR MAGNETIC WARE.		4/77
D	E	KOLM 71	OREO SNOWWATER.		2/76
D	G	ROSS 73	LUNAR DUST.		4/76
D	H	CABRERA 75	ELEVEN TERRESTRIAL MATERIALS.		2/76
D	I	CARRIGAN 76	IMP AND SEMI-MET.		1/77
D	J	BRODERICK 79	42-CM ABSORPTION IN NEUTRAL GALACTIC WOODEN.		12/81*

REFERENCES FOR MAGNETIC MONOPOLE SEARCHES

MALKUS 51	PP 03 894	MALKUS	(ICHC)
BREKEND 69	PP 11 604	BREKEND	(LBL)
FIDEKARD 61	NC 22 657	FIDEKARD	(ICRNL)
AMBLD 63	NC 28 773	AMBLD	(ICRNL)
GOTO 63	PP 13 287	GOTO	(ICRNL)
PEIKHOW 63	PP 09 49 87	PEIKHOW	(ICRNL)
RUF 63	PP 12 2326	RUF	(ICRNL)
CAITHERS 66	PP 14 170	CAITHERS	(ICRNL)
FLEISCH 69	PP 171 2270	FLEISCH	(ICRNL)
FLEISCHER 69	PP 184 1133	FLEISCHER	(ICRNL)
FLEISCHER 69	PP 184 1338	FLEISCHER	(ICRNL)
ALSO 10	JAP 41 58	ALSO	(ICRNL)
SCHATTEN 70	PP 03 2241	SCHATTEN	(ICRNL)
KOLM 71	PP 04 1284	KOLM	(ICRNL)
GUREVICH 72	PL 368 546	GUREVICH	(ICRNL)
ALSO 12	JETP 34 917	ALSO	(ICRNL)
CARRIGAN 73	PP 08 3717	CARRIGAN	(ICRNL)
ROSS 73	PP 08 698	ROSS	(ICRNL)
ALSO 10	SCI 167 701	ALSO	(ICRNL)
ALSO 11	PP 04 3260	ALSO	(ICRNL)
CARRIGAN 74	PP 010 387	CARRIGAN	(ICRNL)
GURKE 75	PL 60R 113	GURKE	(ICRNL)
CABRERA 75	PP 08 D 2515	CABRERA	(ICRNL)
CARRIGAN 76	PP 08 1 270	CARRIGAN	(ICRNL)
ALSO 71	PP 03 50	ALSO	(ICRNL)
BERNHARD 75	PP 011 3009	BERNHARD	(ICRNL)
GIACOMELLI 75	CM 28 21	GIACOMELLI	(ICRNL)
PRICE 75	PL 35 4E7	PRICE	(ICRNL)
ALSO 75	LBL 4260	ALSO	(ICRNL)
ALSO 75	LBL 4289	ALSO	(ICRNL)
ALSO 75	PP 08 1 270	ALSO	(ICRNL)
ALSO 75	PL 35 1107	ALSO	(ICRNL)
ALSO 76	LBL 4289	ALSO	(ICRNL)
ALSO 77	PL 38 129	ALSO	(ICRNL)
ALSO 78	PP 018 1392	ALSO	(ICRNL)
CARRIGAN 76	PP 217 1823	CARRIGAN	(ICRNL)
DELL 76	INC 15 268	DELL	(ICRNL)
STEVENS 76	PP 214 207	STEVENS	(ICRNL)
ZRELOW 76	CJJP 82 1305	ZRELOW	(ICRNL)
CARRIGAN 78	PP 017 1254	CARRIGAN	(ICRNL)
HOFMANN 78	NC 23 357	HOFMANN	(ICRNL)
RODRIGUEZ 79	PP 019 1546	RODRIGUEZ	(ICRNL)
BARTLETT 81	PP 02 612	BARTLETT	(ICRNL)
BONNARD 81	PP 02 23 323	BONNARD	(ICRNL)
KINOSHITA 81	PP 02 4 1127	KINOSHITA	(ICRNL)
ULLMANN 81	PP 04 7 789	ULLMANN	(ICRNL)
CABRERA 82	PP 04 7 1378	CABRERA	(ICRNL)
KINOSHITA 82	PP 04 7 1378	KINOSHITA	(ICRNL)

BIBLIOGRAPHIES

STEVENS 73	VPI-PPP-73-5	D.M. STEVENS	(VPI)
CARRIGAN 77	FERMILAB-77-42	R.A. CARRIGAN, JR.	(FERM)
CRAVAN 81	FERMILAB-81-37	R.E. CRAVAN, N.P. TOWER, P.A. CARRIGAN, VPI	(FERM)

Data Card Listings

For notation, see key at front of Listings.

Stable Particles

CHARM SEARCHES

CHARM SEARCHES AND EVIDENCE

Data on specific charmed states are listed in separate sections in the appropriate places in the Data Card Listings: D , F , and Λ_c - Stable Particles; D^* , F^* - Mesons; Σ_c - Baryons.

Evidence for charm not directly relatable to a given state is listed in this section. Neutrino-induced dilepton events are summarized. Short-lived tracks in emulsions are also dealt with.

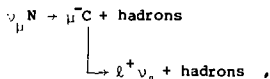
Tri-muon production in neutrino interactions is summarized in the Other New Particle Searches section.

Neutrino-induced Dilepton Events

Many neutrino experiments have now observed dilepton events. These data are summarized in subsections Y, V0, V0A, and VAP. Bubble chamber experiments have observed neutrino-induced $\mu^- e^+$ events associated with strange particle production in the reaction

$$\nu N \rightarrow \mu^- e^+ K^0 (\text{or } \Lambda) + \text{anything.}$$

Production of charmed particles (C) in neutrino interactions would be expected to give rise to such events via the mechanism



where the Cabibbo-favored transition would predict a strange particle among the hadrons. Thus the appearance of neutrino-induced opposite-sign dilepton events, $\mu^- e^+$ events, and associated strange particles can be understood via the charm mechanism. (For another potential explanation see the Heavy Lepton Searches section above.)

Recent experiments show that like-sign dileptons, on the other hand, are produced much in excess of theoretical expectations. See section Y below.

Short-Lived Tracks in Emulsions

The mean life of a weakly decaying charmed meson or baryon of mass M (in GeV) is expected to be in the range¹

$$\tau = (10^{-11} \text{ to } 10^{-13} \text{ sec}) \times 1/M^5$$

with a corresponding mean path length for lab mo-

mentum p (in GeV/c) of

$$\ell = \frac{pc\tau}{M} = (30\mu \text{ to } 3000\mu) \times p/M^6.$$

Thus even at Fermilab energies, the decays of such particles are hard to observe directly in most bubble, streamer, or other chambers, so emulsion is often used. We list data for these experiments in subsections CC and EM below.

Recent experiments using special bubble chambers, emulsions, and silicon detectors have been able to identify the particular charm state with some degree of confidence and to estimate the mean life. These measurements are listed in the separate sections associated with the particular states.

Charm Searches

Experimental evidence for charm production has now been accumulated in various reactions.

Sections CP and CPI include several types of evidence for associated charm production in pN and πN collisions: the prompt μ and ℓ signals; low-mass $\mu^+ \mu^-$ pairs with missing energy; opposite-sign μe events; observations of D and Λ_c in the hadronic final states; as well as hadronic charm decay in association with a lepton trigger. Observations of prompt muons in beam dump experiments are listed in section BD.

Sections CG, MU, and D include evidence in photon, muon, and neutrino beam experiments. Charmed baryon production in e^+e^- reactions is listed in section CE; further information can be obtained from the Listings.

References

1. M. K. Gaillard, B. W. Lee, and J. L. Rosner, Rev. Mod. Phys. **47**, 277 (1975).

PROPERTIES OF THE CHARMED D , D^* , F , F^* , Λ_{CB}/C^* , AND Σ_{CB}/C^* STATES ARE LISTED IN SEPARATE SECTIONS. THE FOLLOWING SECTION CONTAINS INFORMATION ON SEARCHES FOR OTHER CHARMED PARTICLE STATES AND SEARCHES FOR THE ABOVE STATES IN NEW COLLISION PROCESSES.

CE	CHARMED BARYON PRODUCTION IN (e^+e^-) COLLISIONS (CMM*2)	
CE A	PECCOLO 77 S*AG	1/78
CE B	(5.8C-3510R LESS CL=90 FERGUSON 78 S*AG	2/79
CE C		
CE A	PICCOLO 77 LDC AT INCLUSIVE P*AR AND LAMBDA PROD IN 3.7-7.6 GEVEM	1/78
CE A	E+E- AT SLAC. F*INDS SHARP RISE IN C'S BEV 4.4 AND 5 GEV. EVIDENCE	1/78
CE A	FOR PROD OF CHARMED BARYON IN THAT REGION.	1/78
CE B	FERGUSON 78 F*IND INCREASE IN ANTI SIGMA- PROD BY E+ E- (SLAC)	2/79
CE B	BETWEEN 4 AND 7 GEV (ECS) OF DELTA-R (ASIG=-1)D.12+-0.05 CONSISTENT	2/79
CE B	WITH CHARMED BARYON PRODUCTION MODELS. LIMIT IS ON	12/79
CE B	C.S.*BRIALAMBAC -> ASIG* P1=0 P1-1.	12/79

Stable Particles

BOTTOM HADRON, TOP HADRON SEARCHES

Data Card Listings

For notation, see key at front of Listings.

ALL 851A 79 PL 878 287
 ANGEL IN1 79 PL 820 337
 ANGEL IN2 79 PL 840 150
 ANTIPOD 79 JETP 30 143
 ARNEISE 79 PL 808 290
 ATVA 79 PL 43 414

BARANDV 79 PL 818 261
 BARANDV2 79 SJP 29 622
 RAUER 79 PR 43 1953
 BDDPST 79 SJP 29 46
 BERGE 79 PL 818 85

BLITZSCH 79 PL 808 108
 BRANDEN 79 PR 120 337
 BRDM 79 PR 43 410
 CHILMING 79 PL 838 136
 ALSO 79 PR 8151 29
 CTEUS 79 PR 42 1438

DEGROD 79 PL 808 103
 DIAMANT7 79 PR 43 1774
 DIAMANT 79 PL 818 261
 DRJARD 79 PL 810 250
 DRJARD2 79 PL 818 452

FUCHI 79 NC 113 8
 FUCHI2 79 PL 858 135
 GIBONI 79 PL 858 437

ALSO 79 THE SIS - HARVARD
 KOMAR 79 SJP 29 40
 KOMAR2 79 SJP 29 50
 LOKHAN 79 PL 858 463
 RA 79 PR 818 261
 SAMAYANA 79 PR 220 1037

ADAMOVIC 80 PL 898 427
 ALLISON 80 PL 938 509
 BARDON 80 PR 418 250
 ARNEISE 80 PL 948 527
 ASTON 80 PL 948 133
 AUBERT 80 PL 948 56
 RUBERT 80 PL 948 101

BALLAG 80 PL 898 423
 BALLAG2 80 PL 921 569
 BARLOTTA 80 PR 418 250
 BRODEK 80 PR 922 1513

CLARK 80 PR 45 882
 CLARK2 80 PR 45 1465

ALSO 81 PP 224 551
 FANDRUK 80 PR 418 250
 FRITZE 80 PL 960 427

JACOES 80 PR 221 1206
 RITCHIE 80 PR 44 230
 SACHS 80 PR 418 250
 SOJKAS 80 PR 44 266

ADAMOVIC 81 PL 998 271
 ADEV 81 PL 1028 285
 ADEV 81 PL 1017 365
 AMOSOV2 81 PL 1008 151
 ARNEISE 81 PL 1048 409
 AUBERT 81 PL 1008 419

BALLAG 81 PR 24 7
 BASILE 81 PL 938 730
 BASTIE 81 PL 938 457
 BIL 81 LMC 31 990
 GRASSLER 81 PL 998 150

IRJON 81 PL 998 455
 JONKER 81 PL 1078 241
 NISHIKAW 81 PR 44 1555
 NOLSTON 81 JETP 31 232
 BERWINDT 81 PR 223 1889

REVIEWS REFERRED TO IN DATA CARDS

GAISSER 76 PR 214 3153
 T.-G. GAISSER, F. HALZEN (BART+MISC)

BOTTOM HADRON SEARCHES

SEE ALSO THE LISTINGS FOR THE B52500 AND THE LAMBDA/8015000.

HD CROSS SECTION FOR THE PRODUCTION OF A CB, ANTI-B PAIR IN HADRONIC REACTIONS (E=2) NUCLEON.

HD
 HD A B E-39 OR LESS CL=90 COLEMAN 80 SPEC 225 GEV/C P1-NUCL 8/81
 HD B 50 E-39 OR LESS CL=90 DIAMANT 80 SPEC 400 GEV/C P E 8/81
 HD C 30 E38--1.21 <35 DIAMANT 81 SF 500 E38--1.21 LAMBDA E X 4/82
 HD D DRJARD 82 SF 6 P 63 P E 8/81

HD A COLEMAN QD LOOK FOR P1-NUCLEON -> B BAR ANNYTHING, FOLLOWED BY 8/81
 HD B A B -> J/PSI X, ANTI-B -> MU X (AND THE CHARGE CONJUGATE DECAYS). 8/81
 HD B X AND P FE -> MU MU MU X. TABLE 2 OF THEIR PAPER GIVES CS LIMITS 8/81
 HD B FOR SEVERAL FINAL STATES. 8/81

HD C BASILE 81 IS P P AT 62 GEV ECM AT THE CERN-ISR. A G-STANDARD- 4/82
 HD D DESIRATION PEAK IS SEEN IN THE K-0 RE-FLAS SPECTRUM. 4/82
 HD C TRIGG. NO. ON AN AND KEEPING ONLY THE K-0 P1 IN THE QD REGION. 4/82
 HD D SEE THE CERN-ISR LISTING FOR THE MASS OF THE ABOVE 4/82
 HD C CROSS SECTION AFTER THE ABOVE AND OTHER CUTS ARE MADE. 4/82
 HD D DRJARD 82- WITH SAME ACCELERATOR, ENERGY, AND DETECTOR AS BASILE 4/82
 HD D 1. SEE NO LAMBDA/80 PEAK, AND CLAIM THAT NEITHER EXPERIMENT HAS 4/82
 HD D THE SENSITIVITY TO SEE SUCH A SIGNAL. 4/82

BE BOTTOM HADRON PRODUCTION IN E+ E- COLLISIONS. SEE ALSO THE 8/81
 BE LISTING FOR THE B52500. 8/81

BE A ALLI 79 JETP 30 143 8/81
 BE B BARTEL 80 JADE ECM=27.35 GEV 4/82
 BE C BERGER 80 PLUT ECM=12-31.6 GEV 8/81

BE A ALLI 79 ANALYZE SPHERICITY AND THRUST DATA TAKEN BY PLUTO AND TASSC. 6/81
 BE A AT DESY-PETRA. SUGGEST ONSET OF A (B-QUARK ANTI-B+QUARK) THRESHOLD 6/81
 BE C IN THE ENERGY RANGE 9.4<ECM<10 GEV. 6/81

BE B BARTEL 80 IS JADE COLLABORATION AT DESY-PETRA. AN UPPER LIMIT OF 4/82
 BE D 2E-5 SEC IS GIVEN FOR THE B-MESON LIFETIME, ASSUMING ITS MASS IS 4/82
 BE E 5 GEV. 4/82

BE C BERGER 80 MEASURES INCLUSIVE MUON PRODUCTION BY PLUTO AT DESY-PETRA. 8/81
 BE G C (E-QUARK -> MU J) IS BETWEEN 3 AND 0.33. 8/81

MU BOTTOM HADRON PROD. CROSS SECTION IN MU NUCLEON INTERACTION (E=2). 2/82
 MU SEE ALSO SECTION MU IN CHARM SEARCHES. 2/82

MU A L2E-35 CR MASS CL=90 AUBERT 81 SPEC 250GEV MU+ ON FE 2/82
 MU B AUBERT 81 IS MU+ N AT THE CERN SPS. OBSERVE TWO MUON+MU- AND 2/82
 MU C ONE MU+MU- W/NO-SIGN EVENTS. 2/82

Y BOTTOM HADRON PRODUCTION IN NEUTRINO NUCLEON -> 2 LEPTONS ANYTHING 12/81
 Y AND ANTI-NEUTRINO NUCLEON -> 2 LEPTONS ANYTHING. SEE ALSO SECTIONS 12/81
 Y Y, V0, AND VAP IN CHARM SEARCHES. 12/81

Y A Z1 MU+MU- ARNEISE 80 H+BC NEUBAR 12/81
 Y B 2 MU+ (M>5.6GEV) AMOSOV 81 H+BC NEUBAR INDOUC 2/82
 Y C 6 P+M+ MU+MU- (5.6<M<5.6GEV) BALLAG 81 HYDR NEUBAR INDOUC 2/82

Y A ARNEISE 80 IS AT THE CERN-SPS WITH BEBE EXPOSED TO A WIDE-RANG 12/81
 Y A ANTI-NEUTRINO BEBE. OBSERVED PEAK AT AROUND 6 GEV IN D1MUD 12/81
 Y A EVENTS MAY BE DUE TO BOTTOM BARYON PRODUCTION. 12/81

Y B AMOSOV 81 AT FNAL OBSERVED 2 MU+ EVENTS WITH M>5.6GEV, PERHAPS 2/82
 Y B DUE TO B-BARON PRODUCTION AND DECAY. RATIO TO CHARGED CURRENT IS 2/82
 Y C (2<M<3-1.6GEV). 2/82

Y C BALLAG 81 AT FNAL OBSERVE EXCESS IN SAME MASS REGION AS 2/82
 Y C ARNEISE 80. 2/82

 REFERENCES FOR BOTTOM HADRON SEARCHES

ALLI 79 PL 838 335
 ARNEISE 81 PL 948 527
 BARTEL 80 JPHY C 6 295
 BERGER 80 PR 45 1533
 COLEMAN 80 ANALY 1313
 DIAMANT 80 PR 44 507

AMOSOV 81 PL 1008 151
 AUBERT 81 PL 1048 419
 BALLAG 81 PR 24 7
 BASILE 81 LMC 31 97
 DRJARD 82 PL 1088 367

ALLI+KORNER,WILLROD+KRAMER (DESY+HAMB)
 EBART+BRUNHILLER+LIMPHOLD+HUELS+SACL+LUDC
 JADE C.(DESY+HAMB+HEID+MCHS+LANC+RHEL+TOKY)
 C.(DESY+HAMB+HEID+MCHS+LANC+RHEL+TOKY)
 PLUTO C.(DESY+HAMB+HEID+MCHS+LANC+RHEL+TOKY)
 DIAMANT+BERGER+DESY+HAMB+INDISIG+WUPPI
 COLEMAN+IRJON+NEUBAR C. (FNAL+MICH)
 PLUTANO+BERGER+DISHAM C. (STANSLD-CAT)

TOP HADRON SEARCHES

TE TOP HADRON PRODUCTION IN E+ E- COLLISIONS

TE A NONE ECM=22-31.6 GEV BARTEL 79 JADE B 8/81
 TE B NONE ECM=22-31.6 GEV BARTEL2 79 JADE C5 8/81
 TE C NONE BARBER 79 WKJ J C50<CT5 12/81
 TE D NONE ECM=22-31.6 GEV BARBER 79 PLUT B C50<CT5 MU 8/81
 TE E NONE ECM=30-36 GEV BARBER 80 WRJ M C47 MU 8/81
 TE F NONE ECM=22-31.6 GEV BERGER 80 PR 45 1533 8/81
 TE G NONE ECM=33-8 GEV BARTEL 81 JADE B 8/81

TE COMMENTS

TE ALL ABOVE MEASUREMENTS ARE DONE AT DESY-PETRA. THE LAST COLUMN 8/81
 TE SPECIFIES MEASURED QUANTITIES. 8/81

TE B BARTEL2 79 OBSERVE NO SIGNIFICANT ACCUMULATION OF SPHERICAL EVENTS. 8/81

TE D BERGER 79 FIND R=3.80+0.22 WHICH ALONG WITH SPHERICITY AND THRUST 8/81
 TE D BEHAVIORS IS AGAINST OPEN TOP ANTI-TOP CHANNEL BELOW 30GEV. FINAL 8/81
 TE D MUONS ARE ALSO CONSISTENT WITH EXPECTATION WITHOUT TOP QUARK STATE. 8/81

TE E BARBER 80 FIND NO EVIDENCE FOR AN OPEN TOP ANTI-TOP THRESHOLD IN R, 8/81
 TE F FURTHER DIST. AND INCLUSIVE MUONS. ENERGY SCAN IN THE RANGE 29.0<ECM 8/81
 TE E <31.6GEV REVEALS NO HADRON RESONANCE CORRESPONDING TO A TOP-QUARK 8/81
 TE E ANTI-TOP-QUARK BOUND STATE. 8/81

TE F BERGER 80 MEASURES INCLUSIVE MUONS WITH MOMENTUM >1.4 GEV/C. AGREE 8/81
 TE F WITH EXPECTED SEMILEPTONIC DECAYS FROM CHARMED AND BOTTOM MESONS. 8/81

TE G BARTEL 81 MEASURES INCLUSIVE MUONS WITH MOMENTUM >1.4 GEV/C. AGREE 2/82
 TE G WITH EXPECTED SEMILEPTONIC DECAYS FROM CHARMED AND BOTTOM MESONS. 2/82

 REFERENCES FOR TOP HADRON SEARCHES

BARTEL 79 PL 888 171
 BARTEL2 79 PL 898 136
 HARR 79 PL 848 483
 BERGER 79 PL 848 413
 BARBER 80 PR 44 1722
 BERGER 80 PR 45 1533
 BARTEL 81 PL 998 277

JADE C.(DESY+HAMB+HEID+MCHS+LANC+RHEL+TOKY)
 JADE G.(DESY+HAMB+HEID+MCHS+LANC+RHEL+TOKY)
 HARR+COLEMAN
 PLUTO C.(DESY+HAMB+HEID+MCHS+LANC+RHEL+TOKY)
 HARR+J.COLEMAN (FAC+DESY+MICH+RHEP)
 PLUTO C.(DESY+HAMB+HEID+MCHS+LANC+RHEL+TOKY)
 CORDS+DESY+HAMB+HEID+LANC+MCHS+RHEL+TOKY

Data Card Listings

Stable Particles

For notation, see key at front of Listings.

OTHER STABLE PARTICLE SEARCHES

OTHER STABLE PARTICLE SEARCHES

We collect here those searches which do not fit neatly into one of the above search categories. These include axion searches (section AX), supersymmetric partner searches (SYM), trimuon and four-lepton production in neutrino and anti-neutrino reactions (T, FL), and heavy particle searches in accelerator experiments (BE, CH, CS, D, ICH, RPI, BD) and in cosmic rays (F).

Table of search listings under 'OTHER STABLE PARTICLE SEARCHES', including entries for AX (axion production), BE (accelerator experiments), and F (cosmic rays).

Table of search listings for ZENHNER 81, EDWARDS 82, and S, including details on transition rates and experimental conditions.

Table of search listings for FRIMUON PRODUCTION, BARISS 77, and HANSL2 78, detailing neutrino and muon interactions.

Table of search listings for FOUR-LEPTON PRODUCTION, BARISS 77, and M 2E+ E+ MU+, including muon pair production details.

Table of search listings for DI- AND TRI-MUON PRODUCTION, HANSL2 78, and DEGRAD 70, focusing on muon multiplicities.

Table of search listings for HEAVY PARTICLE PRODUCTION CROSS-SECTION, HANSL2 78, and RPI 82, detailing detector searches.

Table of search listings for HEAVY PARTICLE PRODUCTION CROSS SECTION, HANSL2 78, and RPI 82, including muon pair production.

Stable Particles
OTHER STABLE PARTICLE SEARCHES

Table with columns for experiment codes (CS, C, D, etc.), descriptions of searches, and dates. Includes entries like 'HEAVY PARTICLE PRODUCTION CROSS-SECTION', 'LONGLIVED HEAVY PARTICLE INVARIANT C.S.', and 'CROSS-SEC FOR PROD AND CAPT OF LONG-LIVED MASSIVE PARTICLES'.

Data Card Listings
For notation, see key at front of Listings.

Table with columns for experiment codes (F, C, D, etc.), descriptions of data card listings, and dates. Includes entries like 'HEAVY PARTICLE FLUX IN COSMIC RAYS', 'LIGH (BETWEEN MU AND E MASSES) PARTICLE MASS/UNITS-ELECTRON MASSES', and 'HIGHLY IONIZING PARTICLE FLUX'.

Data Card Listings

For notation, see key at front of Listings.

Stable Particles

OTHER STABLE PARTICLE SEARCHES

BLOUSOV 60 JEP 11 1143
 GORUNOV 60 JEP 11 51
 COWARD 63 PR 131 1782
 DORFAN 65 PRL 14 999
 JONES 67 PR 164 1586

BJORNBOE 68 NC B53 241
 BINON 69 PL 308 510

ANTIPON1 71 PL 348 164
 ANTIPON2 71 NP 831 235
 DMRDO 72 NC 9A 319
 TONMAR 72 JPA 5 565
 ALPER 73 PL 468 265
 LEIPNER 73 PRL 31 1226

APPEL 74 PRL 32 428
 FRANKEL 74 PR D9 1932
 YOCK 74 NP 876 175

ALBROW 75 NP 897 189
 BLAGOV 75 YAD. FIZ. 21 300
 FRANKEL 75 PR D12 2561
 JOVANDVI 75 PL 568 105
 YOCK 75 NP 886 216

ALEKSEE1 76 SJMP 22 531
 ALEKSEE2 76 SJMP 23 639
 BALDIN 76 SJMP 22 264
 BRJATORE 76 NC 31A 255
 GUSTAFSD 76 PRL 37 474

BARISH 77 PRL 38 577
 BENVENUT 77 PRL 26 1110
 BLETZACK 77 PRL 38 1241
 CHANG 77 PRL 39 529
 HOLDER 77 PL FOR 303

ALBRICHT 78 PR D18 108
 ALIBRAN 78 PL 748 134
 ASHBYVAN 78 PL 798 497
 BELLETTI 78 PL 768 223
 BENVENUT 78 PRL 40 488
 BMT 78 PRAM 10 115
 BOSETTI 78 PL 748 149

CARROLL 78 PRL 41 727
 CUTTS 78 PRL 41 363
 DONNELLY 78 PR D18 1607
 ALSO 76 PRL 37 315
 ALSO 74 PRL 33 119

HANSLI 78 PL 748 139
 HANSLI 78 PL 778 114
 ALSO 78 NP B142 381

REFERENCES FOR OTHER NEW PARTICLE SEARCHES

+RUSKOV,TAMM,CERENKOV (LEBO)
 +SPRIGDONOV,CERENKOV (LEBO)
 +ZETTELHAB,LYTCHAK,PIETSON (STFAN)
 +EADES,LEDERMAN,LEE,TING (COLUJ)
 (IMCH+MSL+BL+UCLA+MHN+CSU+COLU+MIRA)

+JANGARD,HANSEN,CHATTERGEE (DBNR+BERN)
 DUTEIL,KACHANDV,KHROVDOR,KUTYEN (SERP)

+DENISOV,DOMSKOV,GORIN,KACHANDV (SERP)
 +DENISOV,DOMSKOV,GORIN,KACHANDV (SERP)
 DARDO,MAVARRA,PENEGGO,SITTE (TORIJ)
 TOMBAR,MANANNA,SREERAKYAH (TFIF)
 (CERN+L+EPFL+UOH+ANEL+SOM+BERG+DUC)
 +LARSEN,SESSONS,SMITH,WILLIAMS (BNL+YALE)

+BOUDOUIN,GAENES,LEDERMAN,PAAR (COLU+FNAL)
 +FRAT,RESYANIS,YANG,NEZIRIC (PENN+FNAL)
 P.C.M. YOCK (UNIV OF AUCKLAND)

+BARBER,BENZI(CERN+DARE+PDM+LANC+MCHS+UTRE)
 +DORNA+MRA+ASHOVA,STREISSCHROVA (LEBO)
 +FRAT,RESYANIS,YANG,NEZIRIC (PENN+FNAL)
 JOVANDVI(ETHAMI+KACH+CERN+GEN+HARR+TORIJ)
 P.C.M. YOCK (UNIV OF AUCKLAND+SLAC)

ALEKSEEV,ZAITSEV,KALININA,KRUQLOV (JINR)
 ALEKSEEV,ZAITSEV,KALININA,KRUQLOV (JINR)
 +BERTOGOROV,VISHNEVSKI(LG(SHREVTICH+JINR)
 +MORDO,PIAZOLI,MANUSCICH (ICCT+FRAS+FREE)
 GUSTAFSON,AYRE,JONES,LONGO,MURPHY (RICH)

+BARTLETT,BODEN,BADW (CIT+FNAL+ROCK)
 BENVENUTI,CLINE (FNAL+HARR+PENN+RUTG+MISC)
 BLETZACKER,NIEH,SONI (STON+UCSB)
 CHEN,VAN GEWIKEN (RSU+FNAL)

+KNOBLOCH,MAY (CERN+DORT+HEID+SACL+BGNA)
 +SMITH,VERHASEREN (FNAL+SOM+PURD)
 RACH+BAR+BERG+BRUK+CERN+EQULM+LLA+ONS+
 +PESTEIN,FARKHUTOINDV (ITEP+SERP)
 +FORINI,ZANOTTI (ELB)
 BENVENUTI (FNAL+HARR+PENN+RUTG+MISC)
 +MANNA MURTY (TFIF)
 +DEENA (CERN+BDN+CERN+LOIC+DZF+SACL)

+CHANG,JOHNSON,RYCJA,KI (BNL+PRIN)
 +DULUDE (EROS+FNAL+ILL+BAR+MIT+HARR)
 +FREEDMAN,LVETL,PECCI,SCHWARTZ (STAN)
 REINES,GURN,SOBEL
 GURN,REINES,SOBEL
 +HOLDER+KNOBLOCH(CERN+DORT+HEID+SACL+BGNA)
 +HOLDER+KNOBLOCH(CERN+DORT+HEID+SACL+BGNA)
 HANSL+HOLDER (CERN+DORT+HEID+SACL+BGNA)

HOLDER 78 PL 738 105
 LOWLESS 78 PL 788 305
 MICELMAC 78 LNC 23 441
 MORI 78 PRL 40 432
 VIDAL 78 PL 778 344
 VIERTEL 78 LNC 22 235
 VYSOTSSK 78 JETPL 27 502

ARBITAGE 79 NP B150 67
 BACHT 79 JPD 5 L13
 BECHIS 79 PRL 42 1511
 BENVENUT 79 PRL 42 1024
 BOZDOL 79 NP B159 363
 GALAPRIC 79 PR D20 2708
 GOTEUS 79 PRL 42 1438

DEGRODT 79 PL 858 131
 DISHAM 79 PL 958 142
 GOODMAN 79 PR D19 2972
 SMITH 79 PR B149 525
 ZIMMETS 79 SJMP 29 517

BARBER 80 PRL 45 1304
 BARTLE 80 JPPV 66 255
 BUSSIERE 80 NP B174 1
 FAISSNER 80 PL 968 201

JACQUES 80 PR D21 1206
 LEBRITTO 80 PL 898 271
 SOUKAS 80 PRL 44 564
 YOCK 80 PR D27 61

ASANO 81 PL 1028 159
 DIEHLVAD 81 PL 1058 239
 FAISSNER 81 JPHY C10 95
 FAISSNER 81 PL 1038 234
 KIN 81 PL 1058 59

KINDSHIT 81 PR D28 1707
 LOSCOCO 81 PL 1028 209
 ULLMAN 81 PRL 47 285
 YOCK 81 PR D23 1207
 ZEMDOR 81 PL 1048 494

EDWARDS 82 PRL 48 903
 KINDSHIT 82 PRL 48 71

+KNOBLOCH,MAY (CERN+DORT+HEID+SACL+BGNA)
 +DEENAA (MISCOLLE MUC+FNAL+HARR+HARR)
 MICELRACHEP,PCNTGCOMO (JINR)
 BENVENUTI (FNAL+HARR+PENN+RUTG+MISC)
 +HERR,LECEDERMAN,SWOBER (COLU+FNAL+STON+UCB)
 +HARR,SCHACHER (BERN)
 VYSOTSSKI+IMST+APPL+MATH+USSR ACRA. SCI.)

+BENZ,BOBBINK+ CERN+DARE+PDM+MCHS+JTP+CEHT)
 BMT,GOPALARRISHNA,CUPTA,TEHMA (TFIF)
 +DORNA (UOH+COLU+L+R+L+BEH+SDA)
 BENVENUTI (FNAL+HARR+DZU+PENN+RUTG+MISC)
 +BUSSIERE,ELACOMELLI (EGH+CERN+APP+SACL)
 CALAPRICE,DUMFORD,KOZES,WILLER (ORIN)
 +DISSBURG,FINE,LEE,SODLUSKY (COLU+ILL+BNL)

+HANSL+HOLDER (CERN+DORT+HEID+SACL+BGNA)
 +DIKMAN+BERGER,FAESSLER,KLIU (SLAC+CIT)
 HELLSWOLH,ITG,MAGFALL,STONAN (ELM)
 +BENNETT (FNAL)
 ZIMMETS+L. SKOVEN (UNIV)

+BECKER,BEY (CERN+DORT+HEID+SACL+BGNA)
 JADF (CERN+HARR+HEID+LANC+MCHS+MEL+YORK)
 +GIACOMELLI,LESOVY (EGH+SACL+LAPP)
 +FRENZEL,HEINIGOS,PREUSSGER,SAMM,SAMM+IAC)

+KALEKAR,MILLER,PLAND (RUTG+STEV+COLU)
 LEBRITTO,MICAL,MEL,HEISSINGS (EGH+BNL+NSF)
 +HENDERSON,MENG,EBERGMAN,ENL (APV+ORNL+PENN)
 YOCK (RUC)

+KIUTANI,KUROKAWA (KEK+OK+OSAK)
 +GOLLOVITZ,KONS,FANTOV,KOMSTANTINOV (SERP)
 FAISSNER,FRENZEL,ORTH,HANSL,ORTH+MAM+IAC)
 FAISSNER,FRENZEL,HEINIGOS,PREUSSGER (IAC)
 G.L.RIM,CH,STANN (IAC)

K.KINDSHITA,P.B. PRILE (UCB)
 +SOLAR,GALIK,INDRSTOTTE (MICH+PENN+BNL)
 J.D.ULLMAN (IAC)
 P.C.M. YOCK (IAC)
 ZEMDOR (ETH)

+BARTLETT,PECK (CIT+HARR+PENN+STAN+SLAC)
 KINDSHITA,OSPEL,FRYBAY,UCB+SACL)

Mesons

$\pi^{\pm}, \pi^0, \eta, \rho(770)$

S=0, C=0 MESON STATES

π^{\pm}

0 CHARGED PION(1140.JPG=0-) I=1
SEE STABLE PARTICLE DATA CARD LISTINGS

π^0

9 NEUTRAL PION(135.JPG=0-) I=1
SEE STABLE PARTICLE DATA CARD LISTINGS

η

14 ETAF549.JPG=0-) J=0
SEE STABLE PARTICLE DATA CARD LISTINGS

$\rho(770)$

9 RHO(770.JPG = 1-) [1]

Note on the ρ^0 Mass and Width

Because of the broadness of the ρ meson, its shape is not very well described by a Breit-Wigner formula, which is a narrow-resonance approximation. Although most experimental distributions can be well described by a relativistic Breit-Wigner formula with a P-wave width and an additional shape parameter (PISUT 68), the resulting resonance parameters will clearly depend on this model. A consistent set of such determinations (PISUT 68, ESTABROOKS 74, BARTALUCCI 78, WICKLUND 78, HEYN 80) yields $m_{\rho^0} = (769 \pm 3)$ MeV, $\Gamma_{\rho^0} = (154 \pm 3)$ MeV.

Attempts have been made to determine the ρ pole position in a more model-independent way (LANG 79, BOHACIK 80). It is comforting that these determinations agree perfectly with the above mass average (LANG 79, however, finds a somewhat smaller width).

Independent support comes from an SU(4)-generalization of the Gell-Mann-Okubo mass formula, which in the limit of exactly ideal mixing can be written (MONTONEN 75)

$$\rho = \frac{2(\psi K^* - \phi D^*) + \omega(D^* - K^*)}{2(\psi - \phi) - (D^* - K^*)}$$

Since the masses of the vector mesons on the right-hand side have all been determined to a much better precision than that of the ρ , they can be

Data Card Listings

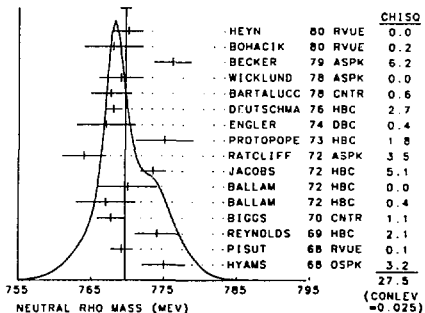
For notation, see key at front of Listings.

used to determine that the mass of the $\rho = (768 \pm 2)$ MeV. The theoretical error due to non-ideal mixing is only of the order of ± 1.5 MeV.

Thus we conclude that the mass of the ρ is (769 ± 3) MeV, somewhat lower than quoted in the 1976-80 editions.

9 RHO MASS (MEV)		
NO LENDOR LIST S-WAVE BREIT-WIGNER FITS, PBAR P DATA WITH HIGH COMBINATORIAL BACKGROUND, AND INSIGNIFICANT OR DOUBTFUL DATA. SEE ALSO THE MINI-REVIEW ABOVE.		
M	CHARGED ONLY	
R	(760.0) (5.0)	CARNOY 66 HBC + 3.5 P1+P TCUT 4
R	(768.0) (5.0)	OLIEDER 65 HBC + 3.5 P1+P 6/78
R	(765.0) (5.0)	ALFF-STEE 66 HBC + 2.3 P1+P 6/78
R	(760.0) (5.0)	MAGOPIANI 66 HBC + 3.0 P1+P 6/78
R	(765.0) (5.0)	MAGOPIANI 66 HBC + 2.14 P1+P TCUT12 9/78
R	2775 (753.5) (10.5)	JACOBS 66 HBC + 2.391-P TCUT 20 6/68
R	(775.0) (5.0)	JAMES 66 HBC + 2.1 P1+P TCUT2.5 8/68
R	(749.0) (3.0)	WEST 66 HBC + 2.1 P1+P 10/66
R	(768.0) (5.0)	MILLER 67 HBC + 2.7 P1+P TCUT20 9/66
R	(773.0) (2.0)	BACON 68 HBC + 2.8 P1+P CUT13 7/69
R	2 900 767. 6.5	EISSNER 67 HBC + 4.2 P1+P CUT10 1/73
X	9650 766.8 1.5	PISUT 68 RVUE 1.7-3.2P1-CT10 6/68
X	650D 766. 7.	BYERLY 73 OSPK + 5. P1-P 2/74
A	756.8 1.4	
A	756.8 1.4	AVERAGE ERROR INCLUDES SCALE FACTOR OF 1.01
NO NEUTRAL ONLY		
R	300 (760.0) (10.0)	AGOLEWS 63 HBC 0 3.5 P1+P
R	500 (770.0) (10.0)	GOLDSHNER 64 HBC 0 3.7 P1+P
R	(750.0) (5.0)	ALFF-STEE 66 HBC 0 2.3 P1+P 6/66
R	(775.0) (5.0)	MAGOPIANI 66 HBC 0 3.0 P1+P 6/66
R	(773.0) (5.0)	MAGOPIANI 66 HBC 0 2.1 P1+P TCUT 12 2/67
R	4207 (759.0) (7.0)	JACOBS 66 HBC + 2.391-P TCUT 20 6/68
R	(765.0) (8.0)	JAMES 66 HBC 0 2.1 P1+P 6/68
R	(760.0) (3.0)	WEST 66 HBC + 2.1 P1+P 10/66
R	4000 (765.0) (5.0)	ASBURY 2 67 CNTR D GAMMA + P8 1/73
R	(768.0) (2.0)	BACON 67 HBC 0 1.7 P1+P 9/67
R	(761.0) (3.0)	MUE 67 HBC 0 2.4 P1+P 7/67
R	(770.0) (4.0)	MILLER 67 HBC 0 2.7 P1+P TCUT20 9/66
R	(775.0) (2.0)	ARMENISE 68 DDC 0 5.1 P1+P 6/68
R	(768.0) (2.0)	MELANDER 69 RVUE 0 2.7 P1+P 1/73
R	(765.0) (10.0)	ALVENSEL 70 CNTR C GAMMA + P CUT.01 1/73
P	(775.0) (5.0)	GLUDING 73 CNTR P 2/74
H	(778.0) (2.0)	HYAMS 73 ASPK 0 17.1 P1+P, P1+P1- 1/74
M	(770.0) (9.0)	ESTABROOK 74 RVUE 0 17 P1+P, P1+P1- 12/75
G	(775.0) (2.0)	GRAYER 74 ASPK 0 12.1 P1+P, P1+P1- 2/74
D	(776.0) (10.0)	RODS 75 RVUE D PHASE SHIFTS 12/75
GM	(769.5) (0.7)	LANG 79 RVUE 0 17.1 P1+P, P1+P1- 3/82*
M	2250 775.0 3.0	HYAMS 68 DSPK 0 11.2 P1+P 9/68
M	13300 769.2 1.5	PISUT 68 RVUE 0 1.7-3.2P1-CT10 6/68
M	1700 774.0 3.0	REYNOLDS 68 HBC 0 2.26 P1+P 12/70
M	140K 767.7 1.9	BIGGS 70 CNTR 0 PHOTOPROD. 1/73
M	1930 767.0 4.0	BALLAM 72 HBC 0 2.8 GAMMA + P 1/73
M	2430 770.0 4.0	BALLAM 72 HBC 0 4.7 GAMMA + P 1/73
M	11200 773.5 1.7	JACOBS 72 HBC 0 2.8 P1+P 3/73
M	8000 768.0 3.0	RATCLIFF 72 ASPK 0 15. P1+P TCUT.3 2/74
M	32000 775.0 4.0	PROTOPOPE 73 HBC 0 1.1 P1+P TCUT.4 2/74
M	6100 767. 4.	ENGLER 74 DBC 0 6. P1+P, P1+P1- 4/78
M	7600 768.0 1.0	DEUTSCHMA 76 HBC 0 16. P1+P 4/78
M	767.6 2.7	BARTALUCCI 78 CNTR 0 8.9 P1+P, P1+P1- 12/77
M	769.0 3.0	WICKLUND 78 ASPK 0 3.4.6 P1+P1- 4/78
M	776.1 2.6	BECKER 79 ASPK 0 17.1 P1+P POLARIZ 12/79
M	768.0 4.0	BOHACIK 80 RVUE 0 12.1 P1+P, P1+P1- 1/82*
M	770 8.0	HEYN 80 RVUE 0 P10N FORM FACTOR 9/81*
M	769.65 1.4	AVERAGE ERROR INCLUDES SCALE FACTOR OF 1.41 (SEE IDEOGRAM BELOW)

WEIGHTED AVERAGE = 769.65 ± 0.74
ERROR SCALED BY 1.4



Data Card Listings

For notation, see key at front of Listings.

Mesons

$\rho(770)$

---NOTES---
 M A FROM FIT OF 3-PARAMETER RELATIVISTIC P-WAVE BREIT WIGNER TO TOTAL
 M B MASS DISTRIBUTION
 M C HEYN 80 INCLUDES ALL SPACELIKE AND TIMELIKE FIP1
 M D VALUES UNTIL 1978
 M E FROM PD E EXTRAPOLATION
 M F ENERGY-DEPENDENT ANALYSIS OF BATON 70, HYAMS 73,
 M G PHOTOPRODUCED 73 PHASE SHIFTS.
 M H PURE P-WAVE SYSTEM.
 M I FROM FIT OF 3-PARAMETER RELATIVISTIC BREIT-WIGNER TO
 M J HELICITY ZEQ PART OF P-WAVE INTENSITY, BECKER 70
 M K INCLUDES DATA OF GRAVER 74.
 M L FROM PHASE SHIFT ANALYSIS OF GRAVER 74 DATA.
 M M FROM PHOTO PRODUCTION, MODEL DEPENDENT.
 M N INCLUDED IN PISUT 68 RVUE
 M O PHASE SHIFT ANALYSIS, SYSTEMATIC ERRORS ADDED CORRESPONDING
 M X TO SPREAD OF DIFFERENT FITS.
 M Z MASS ERRORS ENLARGED BY US TO WIDTH/SORTIM-SEE *R* TYPED NOTE

9 (RHOO) - (RHOD) MASS DIFFERENCE (MEV)
 O A 3400 -5.5 5.0 FOSTER 68 HBC -0.8 PAR P AT REST 12/78
 O B 2250 2.4 2.1 PISUT 68 RVUE P1 N TO RMO N 6/78
 O C 3000 -4.0 4.0 REYNOLDS 69 HBC -0.226 PI-P 12/78
 O D
 O E AVG 0.3 2.2 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.3)

D A FROM QUOTED MASSES OF CHARGED AND NEUTRAL MESSAGES
 9 RHO WIDTH (MEV)
 W
 W NO LINGER LIST 5-WAVE BREIT-WIGNER FITS, PHAR P DATA WITH HIGH
 W COMBINATIONAL BACKGROUND, AND INSIGNIFICANT OR DOUBTFUL DATA.
 W SEE FURTHER MINI-REVIEW ABOVE.

CHARGED ONLY
 W R (177.0) (20.0) CARMONY 64 HBC + 3.5 PI+P,TCUT 4
 W R (100.0) 66 HBC + 2-3 PI+P 6/66
 W R (127.0) (5.0) BLEISEN 68 HBC -3.5 PI-P 6/66
 W R (150.0) (20.0) MAGOPIANI 68 HBC -3.0 PI-P 6/66
 W R (135.0) (20.0) MAGOPIANZ 68 HBC -2.1 PI-,TCUT13 9/67
 W R (132.0) (20.0) JACOB 68 HBC -2.3PI+T,CUT 20 6/68
 W R (147.0) (19.0) JAMES 66 HBC + 2.1 PI-,TCUT2,5 8/66
 W R (149.0) (13.0) WEST 66 HBC -2.1 PI-P 10/66
 W R (153.0) (13.0) MILLER 67 HBC -2.7 PI-,T CUT20 9/66
 W R (150.0) (5.0) BATON 68 HBC -2.8 PI-P 7/69

NEUTRAL ONLY
 W R 900 146. 13. EISENER 67 HBC -4.2 PI-,T CUT10 9/67
 W A 9550 148.2 4.1 PISUT 68 RVUE -1.5-2PI-,CT10 6/68
 W X 6500 146. 13.0 OBERLY 73 HBC 5.0 PI-P 2/74
 W AVG 147.8 3.7 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)

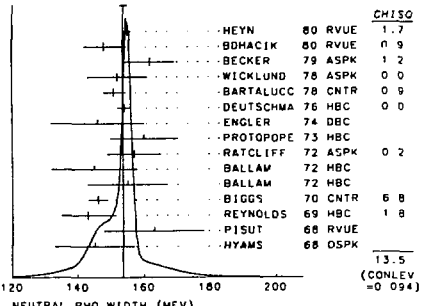
BRANCHING RATIOS
 R1 RHO INTO 4PI/2PI 1P21/1P11
 R1 RHO-- INT0 (PI+ PI- PI+ PI- P10) / (PI+ PI0)
 R1 10.0021R LESS FERBEL 66 HBC -- PI-P 8 RVUE 2.5 10/66
 R1 0.0035 0.008 JAMES 66 HBC 2.1 PI+P 11/66
 R1 RHO INTO (PI+ PI- PI+ PI-) / (PI+ PI-)
 R1 10.0018R LESS JAMES 66 HBC 0 2.1 PI+P 6/66
 R1 10.0021R LESS CHUNG 68 HBC 0 3.2,4,2 PI-P 7/67
 R1 10.0021R LESS CL+..0 HUSON 68 HBC 0 16.0 PI-P 1/71
 R1 10.0015R LESS CL+..0 ERPE 69 HBC 0 2.5-5.8 GAMMA P 10/67

---NOTES---
 M A FROM FIT OF 3-PARAMETER RELATIVISTIC P-WAVE BREIT WIGNER TO TOTAL
 M B MASS DISTRIBUTION
 M C HEYN 80 INCLUDES ALL SPACELIKE AND TIMELIKE FIP1
 M D VALUES UNTIL 1978
 M E FROM PD E EXTRAPOLATION
 M F ENERGY-DEPENDENT ANALYSIS OF BATON 70, HYAMS 73,
 M G PHOTOPRODUCED 73 PHASE SHIFTS.
 M H PURE P-WAVE SYSTEM.
 M I FROM FIT OF 3-PARAMETER RELATIVISTIC BREIT-WIGNER TO
 M J HELICITY ZEQ PART OF P-WAVE INTENSITY, BECKER 70
 M K INCLUDES DATA OF GRAVER 74.
 M L FROM PHASE SHIFT ANALYSIS OF GRAVER 74 DATA.
 M M FROM PHOTO PRODUCTION, MODEL DEPENDENT.
 M N INCLUDED IN PISUT 68 RVUE
 M O PHASE SHIFT ANALYSIS, SYSTEMATIC ERRORS ADDED CORRESPONDING
 M X TO SPREAD OF DIFFERENT FITS.
 M Z MASS ERRORS ENLARGED BY US TO **WIDTH/SORTIM**-SEE *R* TYPED NOTE

9 (RHOO) - (RHOD) MASS DIFFERENCE (MEV)
 O A 3400 -5.5 5.0 FOSTER 68 HBC -0.8 PAR P AT REST 12/78
 O B 2250 2.4 2.1 PISUT 68 RVUE P1 N TO RMO N 6/78
 O C 3000 -4.0 4.0 REYNOLDS 69 HBC -0.226 PI-P 12/78
 O D
 O E AVG 0.3 2.2 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.3)

D A FROM QUOTED MASSES OF CHARGED AND NEUTRAL MESSAGES
 9 RHO WIDTH (MEV)
 W
 W NO LINGER LIST 5-WAVE BREIT-WIGNER FITS, PHAR P DATA WITH HIGH
 W COMBINATIONAL BACKGROUND, AND INSIGNIFICANT OR DOUBTFUL DATA.
 W SEE FURTHER MINI-REVIEW ABOVE.

WEIGHTED AVERAGE = 153.7 ± 1.0
 ERROR SCALED BY 1.3



9 RHO PARTIAL DECAY MODES
 P1 RHO INTO 2PI 130+130
 P2 RHO INTO 4PI 139+139
 P3 RHO INTO PI GAMMA 139+0
 P4 RHO INTO E+ E- 5+5
 P5 RHO INTO PI+ PI- (VIOLATES G) 139+548
 P6 RHO INTO PI+ PI- (VIOLATES G) 100+100
 P7 RHO INTO PI+ PI- (VIOLATES G) 139+139+134
 P8 RHO INTO ETA GAMMA 548+0

9 RHO PARTIAL WIDTHS (KEV)
 W1 RHO INTO (PI GAMMA) 135.0 (10.0)
 W1 67. T. GORBI 74 OSCP - 73, PI+,PI-,PIDA 12/75
 W1 67. T. BERG 80 SPC - PI+,PI-,PID A 1/62*

9 RHO BRANCHING RATIOS
 R1 RHO INTO 4PI/2PI 1P21/1P11
 R1 RHO-- INT0 (PI+ PI- PI+ PI- P10) / (PI+ PI0)
 R1 10.0021R LESS FERBEL 66 HBC -- PI-P 8 RVUE 2.5 10/66
 R1 0.0035 0.008 JAMES 66 HBC 2.1 PI+P 11/66
 R1 RHO INTO (PI+ PI- PI+ PI-) / (PI+ PI-)
 R1 10.0018R LESS JAMES 66 HBC 0 2.1 PI+P 6/66
 R1 10.0021R LESS CHUNG 68 HBC 0 3.2,4,2 PI-P 7/67
 R1 10.0021R LESS CL+..0 HUSON 68 HBC 0 16.0 PI-P 1/71
 R1 10.0015R LESS CL+..0 ERPE 69 HBC 0 2.5-5.8 GAMMA P 10/67

Note on the e^+e^- and $\mu^+\mu^-$ Decays

Extraction of a ratio for $\rho^0 \rightarrow e^+e^-$ is complicated by interference with ω decay. In photoproduction, $\gamma A \rightarrow e^+e^- A$, there is substantial interference between the allowed $(\rho^0, \omega) \rightarrow e^+e^-$ decays. The interference in the colliding-beam reaction $e^+e^- \rightarrow \pi^+\pi^-$ is due to G-parity-violating mixing of the overlapping ρ^0 and ω resonances; it alters the results for the rate $\Gamma(\rho^0 \rightarrow e^+e^-)$ only by a small amount. Therefore at present we average only the values from the $e^+e^- \rightarrow \pi^+\pi^-$ experiments.

The same comment applies to the decay $\rho^0 \rightarrow \mu^+\mu^-$.

R3 RHO INTO(E+ E-)/(PI+PI-) (UNITS 10^4*-1) 1P41/1P11
 R3 SEE MINI-REVIEW ABOVE.
 R3 P 94 10.6-1 [0.14] ASBURY 1 67 CNTR PHOTO PRODUCTION 9/67
 R3 H 10.6-5 [0.11] (0.5) HERTZBACH 67 OSCP ASSUME SU(3)+MIXING 10/66
 R3 A 33 10.5-3 [0.11] ASTVAGTU 68 OSCP ASSUME SU(3)+MIXING 6/68
 R3 D 10.5-2 [0.10] AUSENER 69 OSCP ESTE-COLLID-BEAM 9/68
 R3 F 10.4-9 [0.12] (0.15) BIGGS 70 CNTR PHOTO PRODUCTION 6/70
 R3 0.4-1 0.05 BEMKASAS 72 OSCP E+E- COLLID-BEAM 12/72
 R3
 R3 AVG 0.4-78 0.045 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)
 R3 A NOT SEPARATED FROM OMEGA DECAY. ERROR STATISTICAL ONLY.
 R3 F ASSUMED RHO WIDTH 140 MEV. ERROR STATICAL ONLY.
 R3 H NOT SEPARATED FROM OMEGA DECAY.
 R3 P POSSIBLY LARGE RHO-OMEGA INTERFERENCE

Mesons

 $\omega(783)$, $\eta'(958)$

COYNE 71 NP B 32 333
 CIEADE 71 PRL 23 1749
 MATTHEWS 71 PRL 26 400
 MOFFETT 71 NP B 26 349

AGUILAR 72 NP 6 29
 IPEL 72 PRL 41 B 234
 BASILE 72 PHL CONF. PACO153
 BENKASAS 72 PL 39 B 289
 BENKASAS 72 PL 42 B 507
 BENKASAS 72 PL 42 B 511
 BROWN 72 PL 42 B 117
 DAKIN 72 PR D 6 2231
 FISENER 72 PR D 5 15
 NA-LIFF 72 PL 38 B 1355
 BORNSTEIN 72 NP B 5 1499

BINNIE 73 PR D 8 2789
 AURMS 73 PR D 7 1310

ESTABROOK 74 NP B 81 70
 GREGORIO 74 MC 20 A 43T
 CARRER 74 NP B 3 505
 OREN 74 NP B 71 189

OWMS 75 NP B 90 1
 KALBFLEISCH 75 PR D 11 987
 ROOS 75 NP B 97 165

BRANDEN 76 NP B 104 413
 76 PR D 14 29
 ALSO 73 BINNIE

ANDREWS 77 PRL 38 154
 BARTKE 77 NP B 118 304
 GESSARD 77 NP B 124 392
 HILKOWSKI 77 PRL 6 B 191
 LYONS 77 NP B 125 207
 ROOS 77 LHC 19 415

APFELDOR 78 NP B 133 245
 COOPER 78 NP B 141 209
 JUENER 78 NP B 76 512
 WICKLUND 78 PRL 17 1157

BENKMEIER 79 NP B 150 268
 DZHELYAD 79 PL 84 B 143

COOPER 80 NP B 172 13
 ROOS 80 LHC 27 321

DZHELYAD 81 PL 102 B 296

+BUTLER, FANG-LANDU, MACNAUGHTON (LRL)
 +COOPER, HINES, ALLISON (ANL, ORF)
 +RENTICE, VOOM, CARROLL, WALKER, (INTD, WISC)
 +BINGHAM, FRETTER, BALLAN, LELAND, UCHS, SLAC, TUFT

AGUILAR, BENITEZ, CHUNG, ESNEA, SAMIOS (BNL)
 +HUSLANDER, HULLER, BERTOLUCCI, + (ANL, PFS)
 +BOLLINI, BROGLI, GOLLINI, DALPIAZ, FRABETTI, + (CERN)
 +COSME, JEAN-MARIE, JULLIEN, LAPLANCHE, + (ORSAS)
 +COSME, JEAN-MARIE, JULLIEN, LAPLANCHE, + (ORSAS)
 +COSME, JEAN-MARIE, JULLIEN, LAPLANCHE, + (ORSAS)
 +DOMINGUEZ, HOLLADAY, HODD, BENSTEIN, + (LRL)
 +HAUSER, KREISLER, WISLICE (FRANCO)
 EISENBERG, BULLMAN, DAGAN, + (RHOD, SLAC, TEL)
 +BULOS, CARMIGNEI, ALUO, LETH, LYNN, + (SLAC)
 BORNSTEIN, D'AMBURO, KALBFLEISCH, + (BNL, MITCH)

+CARR, DEBENHAM, DUANE, GARNUTT, + (DIC, SHMP)
 +CHONDOS, KEM, MANDELKERN, PRICE, SCMULTZ (UCI)

ESTABROOKS, NYANS, JONES, BLUM, (CERN, NPJ)
 M.A., GREGORIO (ICTP-TRIESTE)
 +AYRES, DIEBOLD, GREENE, PAULICK, + (ANL)
 +COOPER, FALGOUT, ALLISON, + (ANL, ORF)

+INSON, STACEY, BELL, DALE, (EIR, DURHAM, SHEL)
 KALBFLEISCH, STRAND, CHAPMAN (BNL, MITCH)
 M, ROOS (MELS)

BRANDENBURG, CARMIGNEI, CASHMIRE, DAVIDSON, (SLAC)
 +BINNIE, CARR, DEBENHAM, GARNUTT, + (DIC, SHMP)

+FUKUSHIMA, HARVEY, LORKOWICZ, + (ANL, ORF)
 +144CH, BERL, +BORN, CERN, CRAC, + (CERN, ANL, ORF)
 GESSARD, + (JBN, AF, +ZERN, MIA, + (DIC, +P, +V))
 +JONES, JAMES, ENGLISH, + (CERN, HANST, ANL, ORF)
 +COOPER, CLARK (ORF)
 M, ROOS (MEL, INK 1)

VAN APFELDOR, GRUMENAN, HARTING, + (ZEEN)
 +GURTO, NOVAKET, + (FIS, CERN, COE, +MAD)
 +RIBES, RUMPF, BERTRAND, DIZOT, CHASE, + (LAL)
 +AYRES, DIEBOLD, GREENE, KRAMER, PAULICK, (ANL)

BENKMEIER, EISENBERG, + (EPOL, CERN, COE, +LAL)
 DZHELYAD, HILKOWSKI, GOLLICK, GRIFFIN, + (ISRP)

+DELCOURT, ESCHSTRUTH, FULDA, + (LAL)
 +PELLINEM (MELS)

DZHELYAD, GOLDOVKI, KONSTANTINOV, + (ISRP)

 $\eta(958)$

2 ETA PRIME (958, JPC=0-+) I=0

Note on the J^P Assignment of $\eta'(958)$

From the Dalitz plot analyses of the $\eta' \rightarrow \pi\pi\pi$ and $\eta' \rightarrow \pi^+\pi^-\gamma$ decays and from the observation of an $\eta' \rightarrow \gamma\gamma$ decay mode, all assignments except $J^{PC} = 0^{++}$ and 2^{++} are excluded. The Dalitz plot analyses favor spin 0, but cannot rule out spin 2. The indication of anisotropy in the decay of very forward-produced η' (KALBFLEISCH 73) has not been confirmed by BALTAY 74, thus again favoring spin 0, but still not ruling out spin 2 (LEDNICKY 77).

Two recent analyses, however, seem to have finally established the spin 0 assignment of the η' .

CERRADA 77 perform a partial-wave analysis of the $\pi\pi\pi$ system produced in the reaction $K^0 p \rightarrow \eta' \Lambda$, taking into account the η' and Λ joint decay angular correlations. They conclude that J^P is unambiguously 0^- (see also DELAGUILA 77).

ROUSSARIE 77 analyze a large sample of events from the reaction $\pi^- p \rightarrow \eta' n$ at beam momenta just above threshold. They verify that the η' is produced in a relative S-wave state, and thus the Adair condition is satisfied by their total sample of some 1800 events. The decay angular distribution of the η' is consistent with isotropy, and thus ROUSSARIE 77 conclude that η' cannot be 2.

Data Card Listings

For notation, see key at front of Listings.

2 ETA PRIME MASS (MEV)

M	3415	957.1	1.	REITTERER 09 HEC	1.7-2.7 K-P	9769
M	535	957.4	1.1	PASILEI 71 CNTR	1.6 P1-P, M, RD	11771
M	535	957.4	1.4	BASILE 71 CNTR	1.6 P1-P, M, RD	11771
M	1414	958.2	0.5	DAMBURG 73 HEC	2.2 K-P, LAM, RD	2774
M	400	958.1	1.5	JACOBS 73 HEC	2.9 K-P, LAM, RD	1774
M	957.46	0.33		DUANE 74 MMS	P1-P, N, M	1774

M AVG 957.57 0.25 AVERAGE ERROR INCLUDES SCALE FACTOR OF 1.0

2 ETA PRIME WIDTH (MEV)

M	1000	0.28	0.10	BINNIE 79 MMS	O P1-P, N, M	12770
---	------	------	------	---------------	--------------	-------

2 ETA PRIME PARTIAL DECAY MODES

DECAY MODES

P1	ETA PRIME	INTO	P1+ P1- ETA			
			(P1N)	ETAS DECAY INTO ALL NEUTRALS		
			(P1C)	ETAS DECAY CHARGED		
P2	ETA PRIME	INTO	P10 P10 ETA			
			(P2N)	ETAS DECAY INTO ALL NEUTRALS		
			(P2C)	ETAS DECAY CHARGED		
P3	ETA PRIME	INTO	P1+ P1- GAMMA			
			(P3I)	INCLUDING RHO GAMMA		
P4	ETA PRIME	INTO	GAMMA GAMMA			
P5	ETA PRIME	INTO	OMEGA GAMMA			
P6	ETA PRIME	INTO	RHO GAMMA			
P10	ETA PRIME	INTO	P1+ P1- E-			
P11	ETA PRIME	INTO	2 P1			
P12	ETA PRIME	INTO	3 P1			
P13	ETA PRIME	INTO	4 P1			
P14	ETA PRIME	INTO	5 P1			
P15	ETA PRIME	INTO	6 P1			
P16	ETA PRIME	INTO	P10 E+ E-	(VIBLATES C IN BORN APPROX.)		
P17	ETA PRIME	INTO	E+ E+ E-	(VIBLATES C IN BORN APPROX.)		
P18	ETA PRIME	INTO	P10 P10	(VIBLATES C IN BORN APPROX.)		
P19	ETA PRIME	INTO	P10 OMEGA	(VIBLATES C)		
P20	ETA PRIME	INTO	MU+ MU- GAMMA			
P21	ETA PRIME	INTO	P10 MU+ MU-			
P22	ETA PRIME	INTO	P10 MU+ MU-			

FITTED PARTIAL DECAY MODE BRANCHING FRACTIONS

The matrix below is derived from the error matrix for the fitted partial decay mode branching fractions, P_i , as follows: The diagonal elements are P_i^2 , where $P_i = \sqrt{b_i} \delta P_i$, while the off-diagonal elements are the normalized correlation coefficients $(\delta P_i \delta P_j) / (b_i b_j)$. For the definition of the individual P_i , see the listings above; only those P_i appearing in the matrix are assumed in the fit to be nonzero and are thus constrained to add to 1.

	P 1	P 2	P 3	P 4	P 5
P 1	.4267+-.0174				
P 2	-.0471	.2245+-.0204			
P 3	-.2807	-.0018	.3005+-.0160		
P 4	-.0934	-.1086	-.0042	.0167+-.0016	
P 5	-.0614	-.2054	-.1439	-.0276+-.0054	

Note on $\eta'(958)$ Branching Fractions

In our calculation of the branching fractions of the $\eta'(958)$, we use the decay modes $\pi\pi\pi$ (including $\pi\pi^0\pi^0$), $\rho^0\gamma$, $\omega\gamma$, and $\gamma\gamma$. It is assumed that the rate $\eta \rightarrow$ neutrals is 71.0%.

In the fit we do not use the constraint

$$R = \frac{\Gamma(\eta' \rightarrow \pi^+\pi^-\pi^0)}{\Gamma(\eta' \rightarrow \pi^0\pi^0\pi^0)} = 2$$

from I-spin conservation. The result of the fit is in agreement with it: $R = 1.8 \pm 0.2$.

2 ETA PRIME PARTIAL WIDTHS (KEV)

M1	ETA PRIME	INTO	(GAMMA GAMMA)	(G)
M1	C 23	45.8	(2.3)	ABRAMS 79 SMAG E-E, E-E- RHO GA 12770
M1	C			THE SYSTEMATIC ERROR HAS BEEN ADDED LINEARLY.

2 ETA PRIME BRANCHING RATIOS

SEE MINI-REVIEW ABOVE.

R1	ETA PRIME	INTO	(P1+ P1- ETA NEUTRAL DEC.) / TOTAL (P1N)	
R1	201	0.314	0.026	REITTERER 09 HEC 1.7-2.7 K-P 9769
R1				
R1	FIT	0.303	0.022	FROM FIT ERROR INCLUDES SCALE FACTOR OF 1.01

Data Card Listings

For notation, see key at front of Listings.

Mesons

 $\delta(980)$ $\delta(980)$

36 DELTA(980, JP=0-+ 1-1)

The quantum numbers of the $\delta(980)$ resonance are: $I^G = 1^-$ from its production in $D^0 + \delta\pi$, from its $\eta\pi$ decay, and from the absence of a $\pi\pi$ decay; and $J^P = 0^+$ from the absence of a 3π or $\rho\pi$ decay (LIPKIN 69, GRASSLER 77) and from the decay distributions of the $\eta\pi$ decay. With these quantum numbers the $\delta(980)$ is expected to couple to the $I=1$ $K\bar{K}$ system, too, and to explain the nearby $K\bar{K}$ threshold enhancement (ASTIER 67).

In the scalar $SU(3)$ nonet, the $\delta(980)$ has been a cause of problems. The near degeneracy with the $S^*(975)$ has been difficult to understand: if the S^* is mainly an $s\bar{s}$ system and the δ mainly a $u\bar{d}$ system, naive mass rules predict $m(S^*) - m(\delta) \approx 200$ MeV. Similarly, naive mass rules predict $m(K) - m(\delta) \approx 100$ MeV, whereas the K tends to appear near 1350 MeV.

Another problem has been the ratio of the widths: $\Gamma(K)/\Gamma(\delta)$ is experimentally about 6, whereas $SU(3)$ symmetry predicts 1.4. A solution to this problem came from coupled-channel analyses of the $\eta\pi$ and the $K\bar{K}$ channels (MORGAN 75, FLATTE 76). Above the $K\bar{K}$ threshold strong absorption makes the apparent width shrink, whereas below threshold strong analyticity effects reduce the apparent width.

Recent multi-channel analyses (ACHASOV 79,80, BRAMON 80, TORNVIST 82) confirm this picture, taking also the $\eta'\pi$ channel into account. Moreover the most complete analysis (TORNVIST 82) solves all the above problems at the same time.

Thus the picture of the $\delta(980)$ that emerges is a quark-antiquark state with large $q\bar{q}q\bar{q}$ components in the form of virtual two-meson states. The physical masses of all the scalar mesons are strongly influenced by the number of nearby 2-meson thresholds. Thus the $\delta(980)$ feels strongly the $\eta\pi$, $K\bar{K}$, and $\eta'\pi$ thresholds, whereas the $\kappa(1350)$ is far away from the $K\pi$ and $K\eta'$ thresholds and couples very weakly to the $K\eta$ threshold. This is the reason for the observed large $\kappa-\delta$ mass difference. The small $S^*-\delta$ mass difference, on the other hand, is due to the isoscalar meson feeling the $K\bar{K}$ threshold a factor of 2 times stronger than the isovector meson does.

Thus the conventional $q\bar{q}$ model is sufficient, and no $q\bar{q}q\bar{q}$ multiplet is necessary to explain the observed scalar mesons (TORNVIST 82). See also the mini-reviews under $\epsilon(1300)$ and $\kappa(1350)$ and Appendix IIC.

36 DELTA(980) MASS (MEV)					
M	ETA PI FINAL STATE ONLY.	APPROX.			
M	10 (980.1)	CHUNG 5	68 HBC	- 3.2 PI-P	5/70
M	80 (975.0)	DEFIOX	68 HBC	+ 1.2 PB P,ETA PI	11/77
M	15 (980.0)	WILLER	69 HBC	+ 4.5 K-M,ETA PI	7/69
M	21 (942.0)	BARDAQIN	71 HBC	+ 8 PI+P,P DO PI	11/77
M	(982.1)	ACHASOV1	80 RVUE		9/81*
M	30 980.0	ARMAR	68 HBC	+ 4.5K-ETA PI	2/73
M	20 970.0	BARNES	69 HBC	+ 4.5 K-P,ETA PI	9/67
M	980.0	CAMPBELL	69 DBC	+ 2.7 PI+ D	1/73
M	150 972.0	DEFIOX	72 HBC	+ 0.7 PBAR P,7 PI	1/73
M	70 989.0	WELLS	75 HBC	+ 3.1-6 K-P,ETA PI	11/77
M	80 981.0	GAY	76 HBC	+ 4.2 K-P,ETA PI	11/77
M	80 977.0	GRASSLER	77 HBC	+ 10 PI+P,ETA PI	11/77
M	47 980.0	COMFORTO	78 CSX*	+ 4.5 PI-P,P X-	4/78
M	30 976.0	CORDEN	78 DMG*	+ 12-15PI-P,ETA PI	4/78
M	R 145 980.0	CURTU	79 HBC	+ 4.2 K-P,ETA PI	12/79
M	R 500 980.	EVANGELIS	81 DMG*	12 PI-P,ETA3PI1	1/82*
M	AVG	983.4	2.1	AVERAGE ERROR INCLUDES SCALE FACTOR OF 1.01	
M	C	SYSTEMATIC ERROR 6 MEV DUE TO ENERGY CALIBRATION ADDED			
M	B	COUPLED CHANNEL ANALYSIS WITH FINITE WIDTH CORRECTIONS,SEE MINIREV.			
M	R	FROM DI(285) DECAY			
M	K	K BAR ONLY. SEE THE TYPED NOTE ABOVE			
M	A	100(100-3)	7.0	SYSTEMATIC ROSENFELD 65 RVUE	8/66
M	A	100(100-1)	11.0	ASTIER 67 HBC	12/71
M	A	310 970.	6.	DE RILLY 80 HBC	+ 1.2-2 PB P,D DMG 6/81*
M	A	ASTIER 67 INCLUDES DATA OF BARLOW 67,COMFORTO 67,ARRETEROS 65.			

36 DELTA(980) WIDTH (MEV)					
M	ETA PI FINAL STATE ONLY.	APPROX.			
M	80 (125.0)	DEFIOX	68 HBC	+ 1.2 PB P,ETA PI	11/77
M	20 (50.0)	BARNES	69 HBC	+ 4.5 K-P,ETA PI	11/77
M	150 (30.1)	DEFIOX	72 HBC	+ 0.7 PBAR P,7 PI	2/74
M	70 (16.0)	WELLS	75 HBC	+ 3.1-6 K-P,ETA PI	11/77
M	30 80.0	ARMAR	68 HBC	+ 4.5K-ETA PI	2/73
M	40 15.	CAMPBELL	69 DBC	+ 2.7 PI+ D	1/73
M	15 07.0	WILLER	69 HBC	+ 4.5 K-M,ETA PI	2/74
M	21 31.0	BARDAQIN	71 HBC	+ 8 PI+P,P DO PI	2/74
M	F 80 TO 300	FLATTE	76 RVUE	+ 2.2 K-P,ETA PI	11/77
M	N 55.0	GAY	76 HBC	+ 4.2 K-P,ETA PI	11/77
M	44.0	GRASSLER	77 HBC	+ 10 PI+P,ETA PI	11/77
M	47 60.0	COMFORTO	78 D5P*	+ 4.5 PI-P,P X-	4/78
M	D 50 86.0	CORDEN	78 DMG*	+ 12-15PI-P,ETA PI	4/78
M	R 145 90.0	CURTU	79 HBC	+ 4.2 K-P,ETA PI	11/77
M	B 103 TO 262	ACHASOV1	80 RVUE		9/81*
M	R 500 82.0	EVANGELIS	81 DMG*	12 PI-P,ETA3PI1	1/82*
M	B (500.1) APPROX.	TORNVIST	82 RVUE		1/82*
M	AVG	53.7	6.7	AVERAGE ERROR INCLUDES SCALE FACTOR OF 1.01	
M	B	COUPLED CHANNEL ANALYSIS WITH FINITE WIDTH CORRECTIONS,SEE MINIREV.			
M	F	USING A TWO CHANNEL RESONANCE PARAMETRIZATION OF GAY TO DATA.			
M	N	THE ERROR IN THE PAPER IS WRONGLY QUOTED AT ONE POINT			
M	R	FROM DI(285) DECAY			
M	K	K BAR ONLY. SEE THE TYPED NOTE ABOVE			
M	A	143 (57.0)	13.0	SYSTEMATIC ROSENFELD 65 RVUE	8/66
M	A	100 (25.1) APPROX.	ASTIER	67 HBC	+ SEE NOTE A ABOVE 9/67
M	A	(120.2) APPROX.	MORGAN	75 RVUE	1.2 PBAR P 12/75
M	A	ASTIER 67 INCLUDES DATA OF BARLOW 67,COMFORTO 67,ARRETEROS 65.			
M	M	FROM COUPLED CHANNEL FIT TO DUBUC 72 DATA			

36 DELTA(980) PARTIAL DECAY MODES		DECAY MASSES	
P1	DELTA(980) INTO ETA PI	548+134	
P2	DELTA(980) INTO RHO PI	769+139	
P3	DELTA(980) INTO K KBAR	979+149	
P4	DELTA(980) INTO PI ETA PRIME	139+957	

36 DELTA(980) BRANCHING RATIOS				
R1	DELTA(980) INTO (RHO PI)/ETA PI	[P23/P1]		
R1	(0.25) OR LESS CL-7.0	ARMAR	70 HBC	+ 4.5L,5K-ETA PI 5/70
R2	DELTA(980) INTO (K KBAR)/ETA PI	[P31/P1]		
R2	L (0.25) (0.8)	DEFIOX	72 HBC	+ 0.7 PBAR P 11/77
R2	L SEEN	GAY	76 HBC	+ 4.2 K-P,ETA PI 11/77
R2	(0.7) (0.3)	CORDEN	78 DMG*	+ 12-15PI-P 4/78
R2	(0.75) TO *2	ACHASOV1	80 RVUE	
R2	B	EVANGELIS	81 DMG*	12 PI-P,ETA3PI1 1/82*
R2	B	COUPLED CHANNEL ANALYSIS WITH FINITE WIDTH CORRECTIONS,SEE MINIREV.		
R2	L	FROM THE DECAY OF DI(285).		

REFERENCES FOR DELTA(980)

TURKOT 63 SIENNA CONF 1 661 *COLLINS+FUJII,KERP 18M+PITTSBURGH
 ARRETEROS 65 PL 17 344 ARRETEROS,EDWARDS, JACOBSEN + (CERN+CDF)
 BRUSH 65 PR 139 B 1659 *FRANZINI,KIRSCH,WILLER,STEINBERGER (CERN)
 KIENZLE 65 PL 10 438 *NAGL+LEVRAT,LEFFERVIS + (CERN)
 ROSENFELD 65 DFXRD CCF 56 A H ROSENFELD (LRL--RVUE)

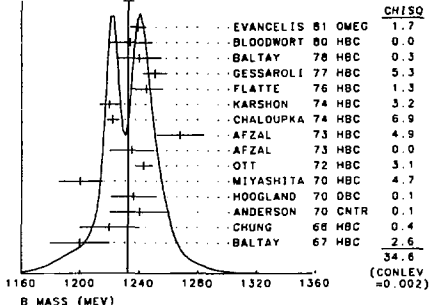
Mesons
B(1235), rho(1250)

B(1235)

11 B(1235,JPG=1) I=1

Table with columns: #, B MASS (MEV), #, B MASS (MEV), #, B MASS (MEV), #, B MASS (MEV). Lists various meson decays and their parameters.

WEIGHTED AVERAGE = 1232.5 +/- 3.4
ERROR SCALED BY 1.6



CHISO table listing decay modes and their contributions to the total fit.

Table with columns: #, B MASS (MEV), #, B MASS (MEV), #, B MASS (MEV), #, B MASS (MEV). Lists various meson decays and their parameters.

11 B WIDTH (MEV)
M D FROM FIT OF THE MASS SPECTRUM
M FIT REQUIRES AN ADDITIONAL JP=1- RESONANCE
M AT 1256 MEV, WIDTH 129 MEV.

Table with columns: #, B INTO OMEGA+ PI, #, B INTO K BAR PI, #, B INTO PI+ PI, #, B INTO PI- PI, #, B INTO ETA PI, #, B INTO K BAR PI. Lists partial decay modes.

11 B PARTIAL DECAY MODES

Table with columns: #, B INTO OMEGA+ PI, #, B INTO K BAR PI, #, B INTO PI+ PI, #, B INTO PI- PI, #, B INTO ETA PI, #, B INTO K BAR PI. Lists partial decay modes.

11 B RATIO FOR B(1235) INTO OMEGA PI
R10 O/S RATIO FOR B(1235) INTO OMEGA PI

Table with columns: #, B INTO OMEGA+ PI, #, B INTO K BAR PI, #, B INTO PI+ PI, #, B INTO PI- PI, #, B INTO ETA PI, #, B INTO K BAR PI. Lists partial decay modes.

Table with columns: #, B INTO K BAR PI / OMEGA PI, #, B INTO PI+ PI / OMEGA PI, #, B INTO PI- PI / OMEGA PI, #, B INTO ETA PI / OMEGA PI, #, B INTO K BAR PI / OMEGA PI. Lists partial decay modes.

REFERENCES FOR B
ABDLINS, LANDER, MEHLHOP, YAGER, UCCS
BONHAR, DODOO, BRACHEN, BRACHEN, HANS, DITZ, MPM
ADENHOLZ, BERLIN, BORN, BORN, HANBORN, LOIC, WAPIN
CARMONY, LANOER, NINOLETSKY, CHUNG, YAGER, UCCS
GOLDHABER, S. GOLDHABER, KADISH, SHEN
SEVERN, EMANUELE, COLUCCI, PELLER
HARDY, HESSE, JACZ, MILLER
HOGES, KADE, KIM, LARSEN, SANDER, VELODE
KARL, BELL, FORMAN, FERREL
MONTANET, MONTANET, CHUNG, YAGER, UCCS
ANDERSON, TO, PR, D, 1 27
CASON, TO, PR, D, 1 851
ERDFEEV, TO, SMP, 11 420
HONES, TO, PR, D, 2 827
HODGLOWD, TO, PR, 3 8 631
MIYASHITA, YON, KROG, KOPPEL, HAN, ICHIO, D
POL, TO, NP, B, 25 109
WERBROUCK, TO, LHC, 4 1267
DEVYNS, TO, PR, 27 1614
FRENKEL, TO, NP, B, 47 61
ITTSER, TO, D, 72, LBL-1547
SISTERSON, HARRISON, HEYDA, JOHNSON, HARWARD
APFAL, TO, LNC, 3 15 81
ARZENISE, TO, NC, 17, A, 707
ARZENISE, TO, LNC, B, 425
ARNDT, TO, PR, 3 14 1971
CASON, TO, NP, B, 64 14
CHUNG, TO, NP, B, 47, B, 226
CHEN, TO, NP, B, 0 23
BALLAC, TO, NP, 0 6 375
CHALOUPEK, TO, PL, 518 407
KARSHON, TO, PR, 10 310 608
CHUNG, TO, NP, 19 211 2426
OUBOVIKO, TO, SJMP, 20 229
FLATTE, TO, PL, 6 14 225
GESSAROLI, TO, NP, B, 126 372
BALTAY, TO, PR, D, 17 62
GAVLLET, TO, PL, 78 B, 158
BLOODWORT, TO, LHC, 21 555
EVANGELIS, TO, NP, B, 178 197
MONG, TO, PR, 46 974

Table with columns: #, B INTO OMEGA+ PI, #, B INTO K BAR PI, #, B INTO PI+ PI, #, B INTO PI- PI, #, B INTO ETA PI, #, B INTO K BAR PI. Lists partial decay modes.

rho(1250)
EVIDENCE NOT COMPELLING. OMITTED FROM TABLE.
SEE ALSO TO BNC PAPER, 1971-REVUE.

Table with columns: #, B INTO OMEGA+ PI, #, B INTO K BAR PI, #, B INTO PI+ PI, #, B INTO PI- PI, #, B INTO ETA PI, #, B INTO K BAR PI. Lists partial decay modes.

Table with columns: #, B INTO OMEGA+ PI, #, B INTO K BAR PI, #, B INTO PI+ PI, #, B INTO PI- PI, #, B INTO ETA PI, #, B INTO K BAR PI. Lists partial decay modes.

Table with columns: #, B INTO OMEGA+ PI, #, B INTO K BAR PI, #, B INTO PI+ PI, #, B INTO PI- PI, #, B INTO ETA PI, #, B INTO K BAR PI. Lists partial decay modes.

Data Card Listings

For notation, see key at front of Listings.

Mesons A₁(1270)

both Deck background and a resonance.

The resonance parameters of the A₁ resonance are obtained by fitting the data (intensity and relative phases) to a phenomenological amplitude containing direct resonance production and a coherent Deck background which is rescattered through the resonance (BOWLER 75, BASDEVANT 77). In the context of this model-dependent analysis, the Deck background is responsible for making the peak of the 1¹S⁰ intensity occur some 110 MeV below the most likely resonance mass.

We take the mass values for the A₁ from these reactions (1240 ± 80 MeV, DANKOWYCH 81; 1280 ± 30 MeV, DAUM 81). Note, however, the result reported in a study of a backwardly produced 3π system in the reaction K⁺p → Σ⁺n⁺π⁺ (1041 ± 13 MeV, GAVILLET 77). Based on a rather small statistical sample, GAVILLET 77 fitted the 1⁺S partial-wave intensity in terms of a single Breit-Wigner function disregarding a possible background component in the 1⁺S wave.

Table with columns for mass (M), A₁ mass (A), resonance width (W), resonance phase (P), and reference (R). Rows include BOWLER 75, GAVILLET 77, DAUM 81, and AVERAGE.

USES DATA OF ANTIPROTON 73, ASCOLI 74, TABAK 74, THOMPSON 74.
USES THE MODEL OF BOWLER 75.
USES MULTICHANNEL FITCHISON-BOULER MODEL.
USES DATA FROM DAUM 80 AND DANKOWYCH 81.
PRODUCED IN K- BACKWARD SCATTERING.

Table with columns for mass (M), A₁ width (A), resonance phase (P), and reference (R). Rows include BOWLER 75, GAVILLET 77, DAUM 81, and AVERAGE.

USES DATA OF ANTIPROTON 73, ASCOLI 74, TABAK 74, THOMPSON 74.
USES ANTIPROTON 73 DATA, WE SELECT SOLUTION B OF BASDEVANT 77.
USES THE MODEL OF BOWLER 75.
USES MULTICHANNEL FITCHISON-BOULER MODEL.

Table with columns for partial decay modes (P1, P2, P3), A₁ partial decay modes (A), decay masses (M), and reference (R).

Table with columns for A₁ branching ratios (A), A₁ branching ratios (A), reference (R), and branching ratios (B).

REFERENCES FOR A₁
BELLINI 63 NC 29 896
BELLINI, FIORINI, HERZ, NEGRI, RATTI (MILAN)

ADPHQZ 64 PL 10 225
GOLDBER 64 PPL 12 346 A
LANDER 64 PPL 13 346 A
...
REFERENCES FOR A₁
BELLINI 63 NC 29 896
BELLINI, FIORINI, HERZ, NEGRI, RATTI (MILAN)

Mesons η(1275), D(1285)

Data Card Listings

For notation, see key at front of Listings.

η(1275)

37 ETAL1275.JPG=0-1+1 0
SEEN IN PHASE SHIFTS ANALYSIS OF THE ETA PI+ PI- SYSTEM WITH PI+ PI- IN AN S-WAVE (STANTON 79). WAIT CONFIRMATION; OMITTED FROM TABLE.

37 ETAL1275) MASS (MEV)

Table with columns M, R, (1275.), APPROX., STANTON, 79, CNTR, 0, B, API-, P, ETA, ZPI, 12/79

37 ETAL1275) WIDTH (MEV)

Table with columns M, R, 170., APPROX., STANTON, 79, CNTR, 0, B, API-, P, ETA, ZPI, 12/79

37 ETAL1275) PARTIAL DECAY MODES

Table with columns P1, P2, ETAL1275) INTO DELTA PI, ETAL1275) INTO ETA PI+ PI-, DECAY MASSES, 989 139, 546 139+ 139

37 ETAL1275) BRANCHING RATIOS.

Table with columns R1, R2, ETAL1275) INTO DELTA PI, LARGE, STANTON, 79, CNTR, 0, B, API-, P, ETA, ZPI, 12/79

REFERENCES FOR ETAL1275)

STANTON 79 PRL 42 346 #BDEKRN+DANKOMY# + (DSU+CAR+MGI+INTO) JP

D(1285)

B D(1285.JPG=1+1+1 0

B D MASS (MEV)

Table with columns M, R, S, 500(1280.), 34(1271.0), 46(1275.0), 1283.0, 1290., 1273.0, 1285., 1303.0, 1283.0, 150 1292., 180 1296., 120 1279.0, 85 1295.0, 320 1282.0, 200 1288.0, 31 1275.0, 103 1283.0, 1278., 1283.0, AVG, FROM PHASE SHIFT ANALYSIS OF ETA PI+PI- SYSTEM. S SEEN IN THE MISSING MASS SPECTRUM.

B D WIDTH (MEV)

Table with columns M, R, 135.0, 110.0, 160., 144.0, 150., 180 146., 500 137., 34 153.0, 113.0, APPROX., D-ANDLAW 67 HBC, 1.0-4.2 PI- P, CARMBELL 69 DBC, 2.7 PI+ D, LORSTAD 69 HBC, 0.7 PB, P, 4.5-BDDY, BARODIN 71 HBC, 8 PI+ P, D+6P, ROESERICK 71 HBC, 16.0 PI+ P, 5 PI, DUBUC 72 HBC, 1.2 PBAR, P, 2KAP, 12/72, MACASCH 78 HBC, 74-76 PB, P, KXP, 4/78, GURTU 79 HBC, 4.2 K-, P, ETA, ZPI, 12/79, BROMBERG 80 SPEC, 100 PI+P, 2KPI, 1/82, DIMISI 80 HBC, 4 PI-, P, K, RB, PI, 12/79, EVANGELIS 81 DMEG, 12 PI-, P, ETASPIP, 1/82

B D WIDTH ERRORS ENLARGED BY US TO 4*MIN(0.5ORTIN). SEE KEY TYPED NOTE

P FROM PHASE SHIFT ANALYSIS OF ETA PI+PI- SYSTEM. R RESOLUTION NOT UNFOLDED. S SEEN IN THE MISSING MASS SPECTRUM.

U UNFOLDED BY DDBRZMSKI 71

B D PARTIAL DECAY MODES

Table with columns P1, P2, D INTO K, KBAR, PI, D INTO PI, PI, RHO, D INTO DELTA, PI, D INTO 2PI+, 2PI-, D INTO K, KBAR, P INTO D INTO 4D)

FITTED PARTIAL DECAY MODE BRANCHING FRACTIONS

The matrix below is derived from the error matrix for the fitted partial decay mode branching fractions, P_i, as follows: The diagonal elements are P_i^2, where P_i = sqrt(OP/OP^2), while the off-diagonal elements are the normalized correlation coefficients (OP_i OP_j) / (OP^2 OP_j^2). For the definitions of the individual P_i, see the listings above; only those P_i appearing in the matrix are assumed in the fit to be nonzero and are thus constrained to add to 1.

Table with columns P, 3, P, 7, -1177.4, -0266, 2931, -4883+, -0586, -0152, -1941, +4000+, -0711

B D BRANCHING RATIOS

THE D BRANCHING RATIOS FIT IS MADE WITH THE ASSUMPTION THAT THE D INTO PI DECAY IS ALWAYS VIA DECAY INTO PI+ PI- PAIRS I.E., RHO PI PI1.

Table with columns R1, R2, D INTO PI+ PI- RHO / I, K, KBAR, PI, D INTO K, KBAR, PI, I, ETA, PI, PI1, D INTO I, K, KBAR, PI, I, ETA, PI, PI1

Table with columns R2, R3, R4, R5, R6, R7, D INTO I, K, KBAR, PI, I, ETA, PI, PI1, D INTO I, K, KBAR, PI, I, ETA, PI, PI1, D INTO I, K, KBAR, PI, I, ETA, PI, PI1

Table with columns R2, R3, R4, R5, R6, R7, D INTO I, K, KBAR, PI, I, ETA, PI, PI1, D INTO I, K, KBAR, PI, I, ETA, PI, PI1, D INTO I, K, KBAR, PI, I, ETA, PI, PI1

Table with columns R2, R3, R4, R5, R6, R7, D INTO I, K, KBAR, PI, I, ETA, PI, PI1, D INTO I, K, KBAR, PI, I, ETA, PI, PI1, D INTO I, K, KBAR, PI, I, ETA, PI, PI1

Table with columns R3, R4, R5, R6, R7, D INTO I, K, KBAR, PI, I, ETA, PI, PI1, D INTO I, K, KBAR, PI, I, ETA, PI, PI1, D INTO I, K, KBAR, PI, I, ETA, PI, PI1

Table with columns R4, R5, R6, R7, D INTO I, K, KBAR, PI, I, ETA, PI, PI1, D INTO I, K, KBAR, PI, I, ETA, PI, PI1, D INTO I, K, KBAR, PI, I, ETA, PI, PI1

Table with columns R4, R5, R6, R7, D INTO I, K, KBAR, PI, I, ETA, PI, PI1, D INTO I, K, KBAR, PI, I, ETA, PI, PI1, D INTO I, K, KBAR, PI, I, ETA, PI, PI1

Table with columns R6, R7, D INTO I, K, KBAR, PI, I, ETA, PI, PI1, D INTO I, K, KBAR, PI, I, ETA, PI, PI1, D INTO I, K, KBAR, PI, I, ETA, PI, PI1

Table with columns R7, R8, D INTO I, K, KBAR, PI, I, ETA, PI, PI1, D INTO I, K, KBAR, PI, I, ETA, PI, PI1, D INTO I, K, KBAR, PI, I, ETA, PI, PI1

NOTE THAT CODEN 78 HBC AND GRASSLER 77 HBC 12-1591-P 4/78

REFERENCES FOR D

Table with columns D-ANDLAW 65 PL 17 347, MULLER 65 PL 14 1074, BARLOW 67 NC 50 A 701, DAHL 67 PR 162 1377, D-ANDLAW 68 NP 8 5 693, DEFOIX 68 PL 28 B 353, CARMBELL 69 PRL 22 1274, DONALD 69 NP 8 11 551, DEFOIX 69 NP 8 44 125, OTHINOMS 69 PL 29 B 520, ANWAR 70 PR 42 8 9, BRODWIN 71 PR 94 2711, ROESERICK 71 PL 34 B 659, GOLDBERG 71 LMC 1 627, BERENYI 72 NP 8 37 621, CHAPMAN 72 NP 8 42 1, TAYING 78 NP 8 359 327, DUBUC 72 NP 8 46 428, THUN 72 PRL 28 1733, VUILLEMIN 75 LNC 14 165, HELLS 75 NP 8 101 333, HANDLER 76 NP 8 110 173, FULLEMME 76 NP 8 334 823, GRASSLER 77 NP 8 121 189, CODEN 78 NP 8 144 253, TAYING 78 NP 8 359 327, MACASCH 78 NP 8 135 203, GURTU 79 NP 8 151 181, STANTON 79 PRL 42 346, BROMBERG 80 PP 2 12 1513, DE BILLY 80 NP 8 176 1, DIMISI 81 NP 8 169 1, EVANGELIS 81 NP 8 178 197

Table with columns D INTO K, KBAR, PI, D INTO PI, PI, RHO, D INTO DELTA, PI, D INTO 2PI+, 2PI-, D INTO K, KBAR, P INTO D INTO 4D)

Data Card Listings

For notation, see key at front of Listings.

 $\epsilon(1300)$

14 EPS1LOH11300, JPC(0+-) 1+0

S-Wave $\pi\pi$ and $K\bar{K}$ Interactions

In this note we discuss information on the non-strange $I_{G}^{J}PC = 0^{+}0^{++}$ partial wave (S wave) coupled to the $\pi\pi$ and $K\bar{K}$ systems.

Near the $\pi\pi$ threshold the S wave shows no resonant behavior. For a discussion of the relevant scattering lengths and various resonance-like kinematic effects, see our 1978 edition.

Up to the ρ meson mass region, the phase shift δ_0^0 is (qualitatively) uniquely determined: it rises monotonically and reaches 60° to 70° near 700 MeV (SONDEREGGER 69, BATON 70, BAILLON 72, CARROLL 72, FRENKIEL 72, GAIDOS 72, PROTOPOESCU 73, HYAMS 73, OCHS 73, ENGLER 74, ESTABROOKS 74,75. GRAYER 74).

In the early phase-shift analyses two solutions for δ_0^0 were found (the "up-down ambiguity") in the 700 to 900 MeV region. The "up" solution corresponds to an ϵ resonance under the ρ meson with mass and width similar to the ρ meson, the $\epsilon(800)$. The "down" solution is characterized by an approximately energy-independent phase shift of almost 90° , showing no resonant behavior. This ambiguity was considered resolved in favor of the "down" solution by the observation of a very rapid decrease in the modulus of the S-wave amplitude between 900 MeV and the $K\bar{K}$ threshold, followed by a sharp drop in the elasticity. δ_0^0 is -90° at about 900 MeV and reaches -180° around 990 MeV (FLATTE 72, GAIDOS 72, HYAMS 73, BINNIE 73, ENGLER 74). However, the region is complicated by the simultaneous presence of the S^* resonance and the opening of the $K\bar{K}$ channel, permitting almost discontinuous jumps from one solution to another.

Without polarization information, the reaction $\pi\pi \rightarrow \pi\pi\pi$ cannot be analyzed unambiguously due to the fact that there are more helicity amplitudes than observables (see, e.g., DONOHUE 75). Thus one is obliged to make some supplementary assumptions.

An amplitude analysis (ESTABROOKS 74) of the largest $\bar{\pi}p$ (unpolarized) $\rightarrow \pi^+\pi^-\pi^0$ experiment (HYAMS 73, GRAYER 74) still finds both the "up"

and "down" solutions. This analysis assumes both spin coherence (the unnatural-parity-exchange, s-channel helicity amplitudes are nucleon spin-flip, i.e., no A_1 -like exchange) and phase coherence (the S-wave amplitude and the unnatural-parity-exchange, meson helicity-zero P-wave amplitude have the same phase). These assumptions may tend to bias the results (MORGAN 74, DONOHUE 75,79).

The advent of π^+p (polarized) $\rightarrow \pi^+\pi^0n$ data (BECKER 79) has made both the spin coherence and phase coherence assumptions unnecessary. Analyzing their data in a model-independent way, BECKER 79 also find both the "up" and the "down" solutions.

The reaction $\pi^+p \rightarrow \pi^+\pi^-\Delta^{++}$ has been analyzed in the region 660 to 860 MeV (OWENS 76, DONOHUE 79) and in the region 600 to 920 MeV (GELFAND 78), using all the information carried by the Δ^{++} decay. The conclusion from both analyses is that the $\epsilon(800)$ of the "up" solution cannot be ruled out.

The only way to rule out a narrow ϵ under the ρ meson (the "up" solution) is to study the $\pi^0\pi^0$ system. With the exception of one experiment (BISWAS 81), all the $\pi^0\pi^0$ experiments agree that no such narrow resonance is present and that the "down" solution describes the data well (DEINET 69, SONDEREGGER 69, SHIBATA 70, BENSINGER 71, APEL 72,79, BRAUN 73, SKUJA 73, RIESTER 75, GRIVAZ 76, DAVID 77, BORREANI 79,81). The phase shifts of BISWAS 81 lie much lower than all others in the 300-700 MeV region, thus requiring a sudden phase motion in the ρ region to match the "down" solution above the ρ .

The region of elastic $\pi\pi$ scattering is known to extend to about 990 MeV, near the $K\bar{K}$ threshold (BATON 70, CARROLL 72, PROTOPOESCU 73, HYAMS 73, OCHS 73). Beyond 1 GeV we therefore have to consider the two channels $\pi\pi$ and $K\bar{K}$, and beyond 1100 MeV the $\eta\eta$ channel also opens up. In addition, the solutions have inherent ambiguities related to the Barrelet zeros of the amplitudes. Thus HYAMS 75 find four solutions in the region 1.0 to 1.8 GeV, ESTABROOKS 74 find eight solutions, and CORDEN 79, extending the $\pi\pi$ analysis to 2.08 GeV, find another eight solutions.

Mesons

 $\epsilon(1300)$

In the past many of these solutions have been ruled out by imposing continuity in various ways, as well as analyticity and unitarity (FROGGATT 75,77, COMMON 76, MARTIN 78).

Unfortunately, the polarization information (BECKER 79) has not yet been fully analyzed. One notes that a model-independent partial-wave analysis (BECKER 79) agrees qualitatively with solutions β and β' (of MARTIN 78).

The β and β' amplitudes describe the experimental moments in each bin without any explicit smoothing; they are analytic in s and approximately analytic in $\cos\theta$. They take into account all waves up to $\ell = 4$. The β solution has a highly elastic S wave, whereas the S wave of solution β' is somewhat inelastic (MARTIN 78). The unique solution of FROGGATT 77, which has explicit smoothness built in and which takes account only of $\ell \leq 3$ waves, is rather similar to β . However, it has problems with unitarity, apparently because of the neglected G wave (MARTIN 78).

The S wave is clearly resonant in the data of BECKER 79. In the 1150 to 1400 MeV region both the S-P and S-D phase differences show the presence of a broad resonance, and the intensity of the S wave confirms this by exhibiting a peak at about 1300 MeV with a width of about 300 MeV; see Fig. 1(a).

The amplitude analysis of the $\pi^-p \rightarrow \pi^+\pi^-n$ experiment of CORDEN 79 has two preferred solutions which are close to β and β' , giving some support for an $\epsilon(1300)$. Also the S wave in the $\pi^0\pi^0$ system tends to confirm the $\epsilon(1300)$ by staying near its unitarity limit around 1200 MeV (APEL 79).

Independent evidence for the $\epsilon(1300)$ comes from studies of the $K\bar{K}$ systems. In the reaction $\pi^-p \rightarrow K_S^0 K_S^0 n$, the S wave exhibits a large intensity in the 1300 MeV region (WETZEL 76, LOVERRE 80), with some evidence for a bump. Moreover, the γ_0^2 moment shows a large negative excursion indicating S-D interference (CASON 76, WETZEL 76, POLYCHRONAKOS 79, GOTTESMAN 80, LOVERRE 80). The main problem is the isospin of the bump: if OPE were the only mechanism, I=0 would be assured. However, an I=1 non-OPE contribution in the same region cannot be excluded. Moreover, the I=1

Data Card Listings

For notation, see key at front of Listings.

K^-K^0 system does show some peaking (MARTIN 79), so one will possibly have to disentangle two resonances in the $K_S^0 K_S^0$ bump.

In agreement with this, the K^+K^- system produced in π^-p , π^+n , and π^-p (polarized) scattering clearly shows the S wave peaking at 1300 MeV [see Fig. 1(b), (c)]; again, both I=0 and I=1 may be present (PAWLICKI 77, MARTIN 79, COHEN 80, COSTA 80, GORLICH 80, WICKLUND 80).

To get from phase shifts to resonance parameters and $q\bar{q}$ -composition one has to make coupled-channel analyses. It has become evident that it is not enough to consider the channels $\pi\pi$ and $K\bar{K}$ only; $\eta\eta$ must also be included. Two such fully unitarized analyses (ACHASOV 79,80,81, TORNVIST 82) finally appear to bring order to the scalar SU(3) nonet. TORNVIST 82 is more general in considering simultaneously all the pseudoscalar-pseudoscalar meson pairs coupling to the full scalar multiplet, but confirms well the earlier results of ACHASOV 79,80.

The picture emerging (TORNVIST 82) is that of a dominantly $q\bar{q}$ -system with large $q\bar{q}q\bar{q}$ components in the form of virtual two-meson bound

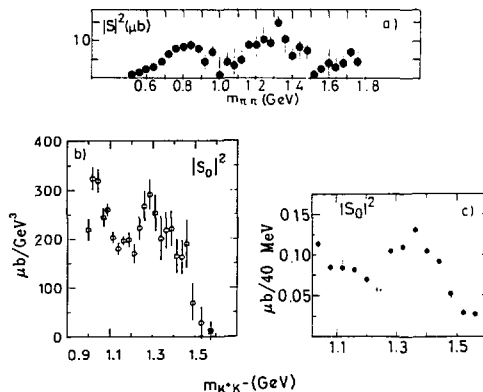


Fig. 1. (a) The absolute intensity (in μb) of the $\pi^+\pi^-$ S wave in 40 MeV bins (without dividing by the bin size), as given by the "down" solution of BECKER 79. (b) Absolute intensity (in $\mu\text{b}/\text{GeV}^3$) of the K^+K^- S wave, as given by the favored solution of COHEN 80, for $|t| < 0.08 \text{ GeV}^2$. (c) The absolute intensity (in $\mu\text{b}/40 \text{ MeV bin}$) of the K^+K^- S wave, as given by the favored solution of GORLICH 80.

Mesons

$\epsilon(1300)$, $\pi(1300)$, $A_2(1320)$

Table with columns for meson name, PDG number, and various parameters (e.g., mass, width, branching ratios). Includes entries for DMS, P1300, KISSER, SKUJA, BASEVANT, BAR-NIR, DINDHUE, FROGGATT, FUJII, HAYES, JONES, MORGAN, PASCUAL, etc.

TORNOVIS 82 HU-TFT-82-1 N.A. TORNOVIS I HELSI

$\pi(1300)$

Table for $\pi(1300)$ showing mass (MEV) and width (MEV) for various experiments (M, F, P, A, N). Includes parameters like AARON, BELLEINI, BONISINI, DAUM.

Table for $\pi(1300)$ showing mass (MEV) and width (MEV) for various experiments (M, F, P, A, N). Includes parameters like AARON, BELLEINI, BONISINI, DAUM.

Data Card Listings For notation, see key at front of Listings.

Table for $\epsilon(1300)$ showing mass (MEV) and width (MEV) for various experiments (M, F, P, A, N). Includes parameters like AARON, BELLEINI, BONISINI, DAUM.

Table for $\pi(1300)$ showing mass (MEV) and width (MEV) for various experiments (M, F, P, A, N). Includes parameters like AARON, BELLEINI, BONISINI, DAUM.

Table for $A_2(1320)$ showing mass (MEV) and width (MEV) for various experiments (M, F, P, A, N). Includes parameters like AARON, BELLEINI, BONISINI, DAUM.

Table for $A_2(1320)$ showing mass (MEV) and width (MEV) for various experiments (M, F, P, A, N). Includes parameters like AARON, BELLEINI, BONISINI, DAUM.

Table for $A_2(1320)$ showing mass (MEV) and width (MEV) for various experiments (M, F, P, A, N). Includes parameters like AARON, BELLEINI, BONISINI, DAUM.

WEIGHTED AVERAGE = 1310.4 ± 1.4 ERROR SCALED BY 1 ±

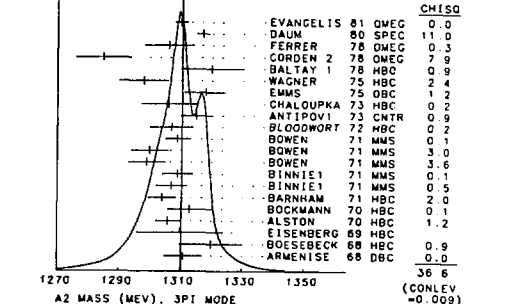


Table for $A_2(1320)$ showing mass (MEV) and width (MEV) for various experiments (M, F, P, A, N). Includes parameters like AARON, BELLEINI, BONISINI, DAUM.

M S SYSTEMATIC ERROR IN MASS SCALE SUBTRACTED M P FROM A FIT TO JP=0- ESILON PI PARTIAL WAVE.

Data Card Listings

For notation, see key at front of Listings.

Mesons A2(1320)

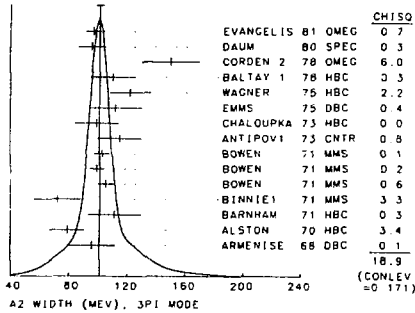
12 A2 MASS (MEV), ETA P1 MODE

Table with columns for mass measurements, key parameters (KEY, CONFORTO, DELFOSSO, B1 SPEC, etc.), and error percentages. Includes an average error and a note about systematic mass error.

12 A2 WIDTH (MEV), SPI MODE

Table listing width measurements for various experiments (ADMINSIE, ALSTON, BARNHAM, etc.) and an average error including a scale factor of 1.21.

WEIGHTED AVERAGE = 101.2 +/- 2.5
ERROR SCALED BY 1.2



FITTED PARTIAL DECAY MODE BRANCHING FRACTIONS

The matrix below is derived from the error matrix for the fitted partial decay mode branching fractions, P, as follows: The diagonal elements are P_i^2 / P_i^2, where P_i = sqrt(P_i^2) / (1/2) P_i. For the definition of the individual P_i, see the listings above, only those P_i appearing in the matrix are assumed in the fit to be nonzero and are thus constrained to add to 1.

Table with columns for P 1, P 2, P 3, P 4. Contains numerical values for branching fractions and their associated errors.

12 A2 PARTIAL WIDTHS

Table listing partial widths for different decay modes such as A2 INTO GAMMA GAMMA (K9V), F R D M RHO P1 CECAY MODE, etc.

12 A2 BRANCHING RATIOS

Table listing branching ratios for A2 (CHARGED ONLY) INTO K BAR K (RHO) P1, A2 INTO LEV LEV (RHO) P1 + K BAR + ETA P1, etc.

Table listing branching ratios for A2 INTO ETA PRIME P1 / TOTAL, A2 INTO ETA PRIME P1 / RHO P1, etc.

12 A2 WIDTH (MEV), CHARGED K BAR MODE

Table listing width measurements for charged K bar mode from experiments like GRAYER, FOLEY, MARGULIES, etc.

12 A2 WIDTH (MEV), ETA P1 MODE

Table listing width measurements for eta p1 mode from experiments like KEY, CONFORTO, DELFOSSO, etc.

12 A2 INTD DECAY MODES

Table listing decay modes such as A2 INTO RHO P1, A2 INTO K BAR, A2 INTO OMEGA P1, etc.

Table listing branching fractions for A2 INTO K BAR (RHO) P1 + K BAR + ETA P1, A2 INTO LEV LEV (RHO) P1, etc.

Table listing branching ratios for A2 INTO ETA PRIME P1 / TOTAL, A2 INTO ETA PRIME P1 / RHO P1, etc.

Table listing branching fractions for A2 INTO OMEGA P1, A2 INTO K BAR (RHO) P1, etc.

REFERENCES FOR A2

ADERHULZ 64 PL 12 226 [LACHENBERG, HENNINGSON, BRONKHORST, OJIC & HERR]
CHUNG 64 PRL 12 821 [OHN, HALL, MESS, ALB, REISSCH, LER]
GOLDNER 64 DURNA CONF 1 A80 G GOLDNER & S GOLDNER, (DALLMAN, SHENLICK)
LAWLER 64 PRL 12 1348 [CAY, GOLDNER, BRONKHORST, HADLER, HENNINGSON, OJIC & HERR]
LONDEBER 64 PRL 12 1348 [CAY, GOLDNER, BRONKHORST, HADLER, HENNINGSON, OJIC & HERR]

Mesons

A₂(1320), E(1420)

Data Card Listings
For notation, see key at front of Listings.

Table listing meson data for A2(1320) and E(1420). Columns include meson name, mass, width, and various experimental parameters.

Table listing meson data for A2(1320) and E(1420) (continued). Columns include meson name, mass, width, and various experimental parameters.

E(1420)

The E(1420) has JPC = 1++ (DIONISI 80) just like the isoscalar companion D(1285). Thus the D and E normally both appear in the same experiments. Recently a new resonance, the I(1440), has been claimed in the decay J/psi -> gamma K* K- (SCHARRE 80,81), on top of the E. It differs from the E in three respects: (a) No D(1285) is produced. (b) It has JPC = 0+-. (c) Its decay is dominantly O(980)n, unlike the E(1420).

We note that the early pp annihilation experiment at rest (BAILLON 67) agrees with (a) and (b). We do not open a new entry for the I(1440) for this edition, but we key BAILLON 67 and SCHARRE 80,81 separately under the E(1420).

Table listing meson data for E(1420) and I(1440). Columns include meson name, mass, width, and various experimental parameters.

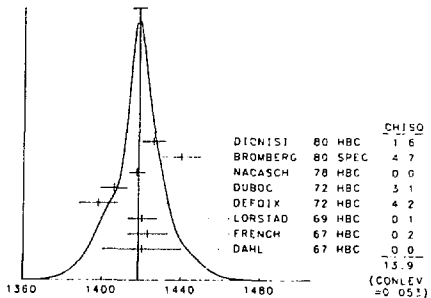
Data Card Listings

For notation, see key at front of Listings.

Mesons

E(1420), f(1515)

WEIGHTED AVERAGE = 1418.4 ± 3.5
 ERROR SCALED BY 4



13 E MASS (MEV)

6 E WIDTH (MEV)

6 E PARTIAL DECAY MODES

6 E BRANCHING RATIOS

6 E INTO K BAR PI (K BAR PI) / (K BAR PI)

6 E INTO PI PI (PI PI)

6 E INTO DELTA PI (DELTA PI)

6 E INTO ETA PI (ETA PI)

6 E INTO 4 PI (4 PI)

6 E INTO (K BAR K) / (K BAR)

6 E INTO (PI PI) / (PI PI)

6 E INTO (ETA PI) / (ETA PI)

6 E INTO (DELTA PI) / (DELTA PI)

6 E INTO (4 PI) / (4 PI)

6 E INTO (K BAR K) / (K BAR)

6 E INTO (PI PI) / (PI PI)

6 E INTO (ETA PI) / (ETA PI)

6 E INTO (DELTA PI) / (DELTA PI)

6 E INTO (4 PI) / (4 PI)

6 E INTO K BAR PI (K BAR PI) / (K BAR PI)

6 E INTO PI PI (PI PI)

6 E INTO DELTA PI (DELTA PI)

6 E INTO ETA PI (ETA PI)

6 E INTO 4 PI (4 PI)

6 E INTO (K BAR K) / (K BAR)

6 E INTO (PI PI) / (PI PI)

6 E INTO (ETA PI) / (ETA PI)

6 E INTO (DELTA PI) / (DELTA PI)

6 E INTO (4 PI) / (4 PI)

REFERENCES FOR E

BAILLON 67 NC 50A 393

BARSH 67 PR 156 1399

DAHL 67 PR 163 1377

ALSO 65 PR 14 1074

FRENCH 67 NC 52A 439

FOSTER 68 NP 8 B 174

BETTINI 69 NC 6 4 1038

LORSTAD 69 NC 16 6 43

DEVONS 71 PR 2 1614

CHAPMAN 72 NP 8 42 1

DEFDIX 72 NP 8 44 125

DUBOC 72 NP 8 46 429

VUILLEMI 75 LNC 14 165

REFERENCES FOR F

13 F PRIME MASS (MEV)

13 F PRIME WIDTH (MEV)

13 F PRIME PARTIAL DECAY MODES

13 F PRIME BRANCHING RATIOS

13 F PRIME INTO (K BAR K) / (K BAR)

13 F PRIME INTO (PI PI) / (PI PI)

13 F PRIME INTO (ETA PI) / (ETA PI)

13 F PRIME INTO (DELTA PI) / (DELTA PI)

13 F PRIME INTO (4 PI) / (4 PI)

13 F PRIME INTO (K BAR K) / (K BAR)

13 F PRIME INTO (PI PI) / (PI PI)

13 F PRIME INTO (ETA PI) / (ETA PI)

13 F PRIME INTO (DELTA PI) / (DELTA PI)

13 F PRIME INTO (4 PI) / (4 PI)

13 F PRIME PARTIAL DECAY MODES

13 F PRIME BRANCHING RATIOS

13 F PRIME INTO (K BAR K) / (K BAR)

13 F PRIME INTO (PI PI) / (PI PI)

13 F PRIME INTO (ETA PI) / (ETA PI)

13 F PRIME INTO (DELTA PI) / (DELTA PI)

13 F PRIME INTO (4 PI) / (4 PI)

13 F PRIME PARTIAL DECAY MODES

13 F PRIME BRANCHING RATIOS

13 F PRIME INTO (K BAR K) / (K BAR)

13 F PRIME INTO (PI PI) / (PI PI)

13 F PRIME INTO (ETA PI) / (ETA PI)

13 F PRIME INTO (DELTA PI) / (DELTA PI)

13 F PRIME INTO (4 PI) / (4 PI)

Mesons

f(1515), ρ'(1600)

Table with columns for experiment name (e.g., R3, R4), parameters (like F PRIME INTD, OR LESS), and results (like 67 HBC, 4.6, 5.0 K-P). Includes entries for R3, R4, R5, R6, R7, R8.

REFERENCES FOR F PRIME table listing various physics papers and authors such as BARNES, CRENNELL, ABRAMS, etc., with their respective publications.

ρ'(1600) 65 PHO PRIME(1600; PG=1) I=1

The ρ'(1600) has been seen in the ρ0π+π- final state, in photoproduction (BINGHAM 72, DAVIER 73, SCHACHT 74, ALEXANDER 75, LEE 75, ATIYA 79, RICHARD 79, BARBER 80, ASTON 81), in e+e- annihilation (BARBARINO 72, CONVERSI 74, CORDIER 79, COSME 79, BACCI 80, DELCOURT 81), in electroproduction (KILLIAN 80), and in a π+d experiment (DIBIANCA 81). If the π+π- subsystem were in an S wave, as has often been assumed, one would also expect to see the ρ' decaying into ρ0π0π0. This has, however, not been seen (ATKINSON 82). Thus the most likely decay chain is ρ' → A1 π + ρππ → 4π.

For the determination of the ρ'(1600) parameters we turn to its relatively rare π+π- decays, which do not have the problems of the above decay chain. The π+π- final state has been produced in πp interactions (HYAMS 73, BECKER 79), in photoproduction (ATIYA 79, ASTON 80), and with weaker evidence in e+e- annihilation (reviewed by GENSHIN 78, HEYN 80). The mass and width in these experiments are consistently 1600 MeV and 300 MeV, respectively. Note, however, that these parameters are the results of very simplified analyses which

Data Card Listings

For notation, see key at front of Listings.

are not adequate for such a broad resonance. An attempt to determine the ρ' pole position in a more model-independent way (LANG 79) from the HYAMS 73 data yields a mass at 1660 MeV.

The elusive ρ' (1250) has recently been reclaimed in the diffractively photoproduced ω0 system (ASTON 80, BARBER 80). However, the J^P determinations are complicated by the simultaneously present B(1235) resonance. In addition, other dynamical effects obscure the interpretation of the ρ' (1250) as a resonance.

Table with columns for parameters (M, H, R, P, M, P, M, X) and results (like 120.1, 102.0, 10.0, 24.0, 2.0, 1568.9, etc.). Includes a key for parameters like HYAMS, PROGATTI, MURFIN, etc.

Table with columns for parameters (M, H, R, P, M, X) and results (like 180.0, 300.0, 50.0, 100.0, 100.0, etc.). Includes a key for parameters like HYAMS, PROGATTI, MURFIN, etc.

Data Card Listings

For notation, see key at front of Listings.

Mesons

$\rho(1600), \theta(1640)$

65 RHO PRIME PARTIAL DECAY MODES

Table with columns: P1, P2, P3, P4, P5, P6, P7, P8, P9, P10, P11, P12. Rows contain decay mode notations and associated mass values.

65 RHO PRIME PARTIAL WIDTHS (KEV)

Table with columns: W1, W1 D, W1 D. Rows contain width values and model dependencies.

65 RHO PRIME BRANCHING RATIOS

Table with columns: R1, R1 S, R3, R3 S, R4, R4 D, R5, R5 C, R5 P, R5 S, R5 F, R5 H, R5 R, R5 U, R6, R7, R7 U, R8, R8 B, R8 U, R9, R9 D, R9 U. Rows contain branching ratio notations and values.

65 RHO PRIME G(I)G(E)E(I)G(TOTAL) (KEV)

Table with columns: G2, G2 P, G2 P, G2, G2 AVG, G5, G5 M, G10, G10 M, G12, G12 M. Rows contain G(I)G(E)E(I)G(TOTAL) values and model dependencies.

REFERENCES FOR RHO PRIME

Table with columns: ALVENS, BRAUN, BUDS. Rows contain references for RHO PRIME.

Table with columns: BACCI, BARBARI, BARTOLI, BINGHAM, BIRLOTTI, DIEBOLD, EISENBERG, JALOWY, SMADJA, CHERADINI, CHUNG, DAUER, DEWET, DUCHS, LAMONTAGN, PASC, BALLAM, BERNABE, CHALOUPEK, CORVETSI, ESTABROD, GRAYER, HIRSCHFEL, SCHACHT, ALEXANDER, ALLES, CHUNG, ESTABROD, GRAYER, HIRSCHFEL, SCHACHT, BASSOMPI, CECOMINI, JOHNSON, BIJONNE, COSTA, FROGGATT, GESSAROLI, GENSHI, G. MARTIN, ATTYA, BACCI, BIRLOTTI, COSME, LANG, RICHARD, ASTON, BARBER, BACCI, BIRLOTTI, MEY, KELLIAN, O'DONNELL, PENSO, ASTON, BACCI, DELCOURT, DEFECA, ATRINSON. Rows contain various meson decay data and references.

$\theta(1640)$

68 THE(1640) JP(2-)- I(0)
SEE IN JP(1) INTO GAMMA TETA. THEREFORE C-+.
TETA DECAYS INTO 2 TAU, THEREFORE I(0+).
JP(2-) IS PREFERRED OVER 0+, HIGHER SPINS NOT
STUDIED. NEEDS CONFIRMATION. OMITTED FROM TABLE.

Table with columns: M, W. Rows contain data for THE(1640) MASS (MEV) and THE(1640) WIDTH (MEV).

REFERENCES FOR THE(1640)

Table with columns: EDWARDS. Rows contain references for THE(1640).

Mesons

 $\phi(1850)$, $X(1850)$, $K(1935)$ Data Card Listings
For notation, see key at front of Listings.

54 PHI(1850) BRANCHING RATIOS

R1	PHI(1850) INTD	(K+K-)/TOTAL	(P21)/(P1)	
R1	0.8	0.4	ALHARRAN 81 OMEG	8.25 K-P, LAM X 1/82*

REFERENCES FOR PHI(1850)

ASTON	80 PL 92 B 219	(B0NN+CERN+EPOL+GLAS+LANCH+MC+SORS+PARIS+)
ALHARRAN	81 PL 101 B 357	(BARI+BDNN+CERN+CLAS+MICH+LBNP)
DELCOUR	81 BOW COME - 205	(DSM87)

X(1850)

FROM AN AMPLITUDE ANALYSIS OF THE K+K- SYSTEM SEEN IN P+P INTO K+K- AT 10 GEV/C. NEEDS CONFIRMATION. OMITTED FROM TABLE.

38 X(1850) MASS (MEV)

M	A	1857.0	35.0	24.0	COSTA	80 OMEG	0 10 PI-P, K+ K- N	1/82*
M <td>A</td> <td colspan="7">ERROR INCREASED BY SPREAD OF TWO SOLUTIONS.</td>	A	ERROR INCREASED BY SPREAD OF TWO SOLUTIONS.						

38 X(1850) WIDTH (MEV)

M	A	185.0	102.0	139.0	COSTA	80 OMEG	0 10 PI-P, K+ K- N	1/82*
M <td>A</td> <td colspan="7">ERROR INCREASED BY SPREAD OF TWO SOLUTIONS.</td>	A	ERROR INCREASED BY SPREAD OF TWO SOLUTIONS.						

29 K(1850) PARTIAL DECAY MODES

P1	X(1850) INTD	K+ K-	DECAY MASSES
			493+ 493

38 X(1850) BRANCHING RATIOS

R1	X(1850) INTD	(K+K-)/TOTAL	(P1)	
R1	SEEN		COSTA	80 OMEG 0 10 PI-P, K+ K- N 1/82*

REFERENCES FOR X(1850)

COSTA	80 NP B 175 402	(BARI+BDNN+CERN+CLAS+LIVP+MILAN+IENI)
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S(1935)

A narrow enhancement called the S(1935) has been observed in the antiproton-proton total cross section (CARROLL 74, CHALOUKPA 76, BRUCKNER 77, SAKAMOTO 79).

This observation is not confirmed by more recent experiments (ALLEN 80, KAMAE 80, JASTRZEMBSKI 81, LOWENSTEIN 81), or the effect is found to be smaller in magnitude and larger in width (HAMILTON 80).

No significant signal is observed for a narrow S(1935) in backward antiproton-proton elastic scattering (GARNJOST 79), nor in the charge-exchange cross section (GARNJOST 75, CHALOUKPA 76, HAMILTON 80).

The first positive observations of the S(1935) in hadro-production experiments (DAUM 80) are not confirmed by more recent ones (DAUM 81). In view of this situation, the only observation which remains unchallenged is the one of RICHARD 79 in a photoproduction experiment. One should wait for confirmation before taking the S(1935) as a well-established narrow resonance.

31 S MASS (MEV)

M	S	CHANNEL	NBAR	N	CLINE	70 HBC	0 -25-74	PBAR P	2/72
M <td>C</td> <td>(1940)</td> <td>(8.)</td> <td></td> <td>BENVENUTI</td> <td>71 HBC</td> <td>0 -1-8</td> <td>PBAR P</td> <td>2/72</td>	C	(1940)	(8.)		BENVENUTI	71 HBC	0 -1-8	PBAR P	2/72
M <td>B</td> <td>(1989)</td> <td>1.</td> <td></td> <td>CARROLL</td> <td>74 CNTR</td> <td>0</td> <td>S CHAN, PBAR P, D</td> <td>12/75</td>	B	(1989)	1.		CARROLL	74 CNTR	0	S CHAN, PBAR P, D	12/75
M <td>S</td> <td>1932</td> <td>2.</td> <td></td> <td>D'ANDOLU</td> <td>75 HBC</td> <td>0 -175-750</td> <td>PBAR P</td> <td>12/75</td>	S	1932	2.		D'ANDOLU	75 HBC	0 -175-750	PBAR P	12/75
M <td>C</td> <td>(1942)</td> <td>(5+)</td> <td></td> <td>KALOGERO</td> <td>75 DPC</td> <td></td> <td>PBAR N ANNH</td> <td>12/75</td>	C	(1942)	(5+)		KALOGERO	75 DPC		PBAR N ANNH	12/75
M <td>Z</td> <td>(1974+1)</td> <td>(2,6)</td> <td>(1,4)</td> <td>CHALOUKPA</td> <td>76 HBC</td> <td>0</td> <td>PBAR P TOT, ELAS</td> <td>12/75</td>	Z	(1974+1)	(2,6)	(1,4)	CHALOUKPA	76 HBC	0	PBAR P TOT, ELAS	12/75
M <td>S</td> <td>1935.0</td> <td>1.0</td> <td></td> <td>BRUCKNER</td> <td>77 SPEC</td> <td>0</td> <td>4+-85</td> <td>PBAR P</td>	S	1935.0	1.0		BRUCKNER	77 SPEC	0	4+-85	PBAR P
M <td>S</td> <td>1939.0</td> <td>3.0</td> <td></td> <td>SAKAMOTO</td> <td>79 HBC</td> <td>0</td> <td>-37-73</td> <td>PH P</td>	S	1939.0	3.0		SAKAMOTO	79 HBC	0	-37-73	PH P
M <td>S</td> <td>1935.5</td> <td>1.0</td> <td></td> <td>ASTON</td> <td>80 OMEG</td> <td>0</td> <td>PBAR P, P</td> <td>1/82*</td>	S	1935.5	1.0		ASTON	80 OMEG	0	PBAR P, P	1/82*
M <td>A</td> <td>1949.1</td> <td>(10.)</td> <td></td> <td>DEFDIX</td> <td>80 HBC</td> <td>0</td> <td>PBAR P, SPI</td> <td>1/80</td>	A	1949.1	(10.)		DEFDIX	80 HBC	0	PBAR P, SPI	1/80
M <td>N</td> <td>1939.0</td> <td>2.0</td> <td></td> <td>HAMILT02</td> <td>80 CNTR</td> <td>0</td> <td>S CHAN, PBAR P</td> <td>12/79</td>	N	1939.0	2.0		HAMILT02	80 CNTR	0	S CHAN, PBAR P	12/79

PRODUCTION EXPERIMENTS

M	A	361(942.0)	11.0	DAUM	80 CNTR	93 P, PB P + X	12/79
M <td>A</td> <td>1935.3</td> <td>1.0</td> <td colspan="4">AVERAGE ERROR INCLUDES SCALE FACTOR OF 1.71 (SEE IDEOGRAM BELOW)</td>	A	1935.3	1.0	AVERAGE ERROR INCLUDES SCALE FACTOR OF 1.71 (SEE IDEOGRAM BELOW)			

M A FROM ENERGY DEPENDENCE OF SPI CROSS-SECTION, 10+- FROM OBSERVATION

M B OF OMEGA AND DECAY. P+ AND P-1, AS 71 P1 ALSO SEEN.

M C SEEN AS A BUMP IN THE PBAR P - KS KL CROSS SECTION WITH JPC=1-+, NOT SEEN BY GARNJOST 72 WITH EQUAL STATISTICS.

M D FROM ENERGY DEPENDENCE OF FAR BACKWARD ELASTIC SCATTERING.

M E SOME INDICATION OF ADDITIONAL STRUCTURE.

M F 1+0 FAVORED, J=0 OR 1, SEEN IN TOTAL PBAR P TOTAL CROSS-SECTION.

M G PRIMARILY FROM ANNIH. REACTIONS. NOT SEEN IN PBAR P TOTAL AND

M H ANNIH. CROSS SECTIONS.

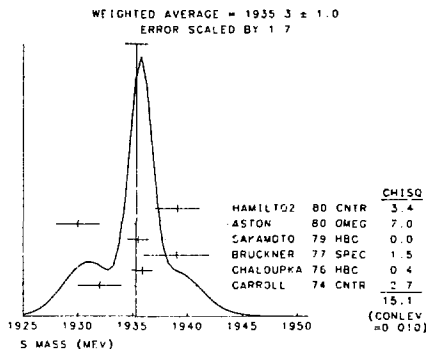
M I SEEN IN S CHARGED MODE. NOT SEEN BY BOWEN 73 WITH 6X STATISTICS.

M J NARROW BUMP SEEN IN TOTAL PBAR P, D CROSS-SECTIONS. 150SPM UNCERTAIN

M K NOT SEEN BY PEAR P CEE BY GARNJOST 75, CHALOUKPA 76. INTEGRATED

M L CROSS-SECTION 3X LARGER THAN BRUCKNER 77.

M M NOT SEEN BY ALBERT 70 WITH COMPARABLE STATISTICS.



31 S WIDTH (MEV)

M	S	CHANNEL	NBAR	N	CLINE	70 HBC	0 -25-74	PBAR P	2/72
M <td>C</td> <td>(1940)</td> <td>(8.)</td> <td></td> <td>BENVENUTI</td> <td>71 HBC</td> <td>0 -1-8</td> <td>PBAR P</td> <td>2/72</td>	C	(1940)	(8.)		BENVENUTI	71 HBC	0 -1-8	PBAR P	2/72
M <td>B</td> <td>(1989)</td> <td>1.</td> <td></td> <td>CARROLL</td> <td>74 CNTR</td> <td>0</td> <td>S CHAN, PBAR P, D</td> <td>12/75</td>	B	(1989)	1.		CARROLL	74 CNTR	0	S CHAN, PBAR P, D	12/75
M <td>S</td> <td>1932</td> <td>2.</td> <td></td> <td>D'ANDOLU</td> <td>75 HBC</td> <td>0 -175-750</td> <td>PBAR P</td> <td>12/75</td>	S	1932	2.		D'ANDOLU	75 HBC	0 -175-750	PBAR P	12/75
M <td>C</td> <td>(1942)</td> <td>(5+)</td> <td></td> <td>KALOGERO</td> <td>75 DPC</td> <td></td> <td>PBAR N ANNH</td> <td>12/75</td>	C	(1942)	(5+)		KALOGERO	75 DPC		PBAR N ANNH	12/75
M <td>Z</td> <td>(1974+1)</td> <td>(2,6)</td> <td>(1,4)</td> <td>CHALOUKPA</td> <td>76 HBC</td> <td>0</td> <td>PBAR P TOT, ELAS</td> <td>12/75</td>	Z	(1974+1)	(2,6)	(1,4)	CHALOUKPA	76 HBC	0	PBAR P TOT, ELAS	12/75
M <td>S</td> <td>1935.0</td> <td>1.0</td> <td></td> <td>BRUCKNER</td> <td>77 SPEC</td> <td>0</td> <td>4+-85</td> <td>PBAR P</td>	S	1935.0	1.0		BRUCKNER	77 SPEC	0	4+-85	PBAR P
M <td>S</td> <td>1939.0</td> <td>3.0</td> <td></td> <td>SAKAMOTO</td> <td>79 HBC</td> <td>0</td> <td>-37-73</td> <td>PH P</td>	S	1939.0	3.0		SAKAMOTO	79 HBC	0	-37-73	PH P
M <td>S</td> <td>1935.5</td> <td>1.0</td> <td></td> <td>ASTON</td> <td>80 OMEG</td> <td>0</td> <td>PBAR P, P</td> <td>1/82*</td>	S	1935.5	1.0		ASTON	80 OMEG	0	PBAR P, P	1/82*
M <td>A</td> <td>1949.1</td> <td>(10.)</td> <td></td> <td>DEFDIX</td> <td>80 HBC</td> <td>0</td> <td>PBAR P, SPI</td> <td>1/80</td>	A	1949.1	(10.)		DEFDIX	80 HBC	0	PBAR P, SPI	1/80
M <td>N</td> <td>1939.0</td> <td>2.0</td> <td></td> <td>HAMILT02</td> <td>80 CNTR</td> <td>0</td> <td>S CHAN, PBAR P</td> <td>12/79</td>	N	1939.0	2.0		HAMILT02	80 CNTR	0	S CHAN, PBAR P	12/79

PRODUCTION EXPERIMENTS

M	A	16.0	APPROX.	DAUM	80 CNTR	93 P, PB P + X	12/79
M <td>A</td> <td colspan="6">AVERAGE MEANINGLESS SCALE FACTOR = 1.91</td>	A	AVERAGE MEANINGLESS SCALE FACTOR = 1.91					

M SEE FOOTNOTES UNDER S MASS ABOVE

31 S PARTIAL DECAY MODES

P1	S	1470	PBAR P	DECAY MASSES
				936+ 939

REFERENCES FOR S(1935)

CLINE	68 PRL 23 1268	+EGLISH+REEDER+TERRELL+TWITY (MISCONSIN)
CLINE	70 PREPRINT	D.CLINE, J.ENGLESH, O.O. REEDER (MISC1)
ALSO	10 KIEV CONF.	ASTIER RAPPORTEUR T.1
BENVENUTI	71 PRL 23 289	BENVENUTI, CLINE, NUTZ, REEDER, SCHERER (MISC)
PINSKY	71 PRL 27 1548	STEPHEN S. PINSKY (UTAH-BARNOWIE)
BIZARRI	72 PR D 160	+GUDIMOV, MARZANO, CASTELLI, + (IRON+MIST)
BOWEN	71 PRL 25 890	+EARLES, FAISLEP, BLIEDEN, + (MEAS+STON)
BOWEN	73 PRL 30 332	+CONDON, DONAHUE, MANDELKERN, PRICE, + (UC)3
BURNS	73 PR D 1286	M. KENZLE (ICEN)
KENZLE	73 PR D 7 3520	
BURNS	74 NC 204 463	+CONDON, MANDELKERN, PRICE, SCHULTZ (UC)3
CARROLL	74 PRL 32 247	+CHANG+KITA+LI+MAZUR+MICHAEL (BNL)

Data Card Listings

For notation, see key at front of Listings.

Mesons
S(1935), δ(2030), h(2040), π(2050)

ARASHIAN 75 PDL 34 09E
D-ANDLAW 75 PL 58F 223
DEBILIE 75 WLLIAMS CONF.
DORNACHT 75 WC 29 4 20T
GARAJOSTI 75 PDL 35 1605
KALCIGERO 75 PDL 34 2473
WEINGART 75 PDL 34 1291
...
WAGNER 7A LONDON CONF.
APEL 75 PL 57B 359
BLUM 75 PL 57A 405
ANTIPOV 77 NP B 119 45
...
WAGNER 7A LONDON CONF.
APEL 75 PL 57B 359
BLUM 75 PL 57A 405
ANTIPOV 77 NP B 119 45
...
WAGNER 7A LONDON CONF.
APEL 75 PL 57B 359
BLUM 75 PL 57A 405
ANTIPOV 77 NP B 119 45
...

δ(2030)
17 DELTA2030.JPG+++ I=1
SEE IN PARTIAL WAVE ANALYSIS OF THE K KBAR AND
P1+P1-PI0 SYSTEMS.
NEEDS CONFIRMATION, OMITTED FROM TABLE.
...
17 DELTA(2030) MASS (MEV)
...
17 DELTA(2030) WIDTH (MEV)
...
17 DELTA(2030) PARTIAL DECAY MODES
...
17 DELTA(2030) BRANCHING RATIOS
...
17 DELTA(2030) INT0 K KBAR
...
17 DELTA(2030) INT0 PI+ PI- PI0
...
17 DELTA(2030) INT0 K KBAR
...
17 DELTA(2030) INT0 PI+ PI- PI0
...

h(2040)
16 H(2040.JPG+++ I=0
...
16 H MASS (MEV)
...
16 H WIDTH (MEV)
...
16 H PARTIAL DECAY MODES
...
16 H BRANCHING RATIOS
...
REFERENCES FOR H(2040)
...
16 H BRANCHING RATIOS
...
REFERENCES FOR H(2040)
...
16 H BRANCHING RATIOS
...
REFERENCES FOR H(2040)
...
16 H BRANCHING RATIOS
...
REFERENCES FOR H(2040)
...

17 DELTA(2030) PARTIAL DECAY MODES
...
17 DELTA(2030) BRANCHING RATIOS
...
17 DELTA(2030) INT0 K KBAR
...
17 DELTA(2030) INT0 PI+ PI- PI0
...
17 DELTA(2030) INT0 K KBAR
...
17 DELTA(2030) INT0 PI+ PI- PI0
...
17 DELTA(2030) INT0 K KBAR
...
17 DELTA(2030) INT0 PI+ PI- PI0
...
17 DELTA(2030) INT0 K KBAR
...
17 DELTA(2030) INT0 PI+ PI- PI0
...

π(2050)
43 PI(2050.JPG+++ I=1
PREVIOUSLY CALLED A4.
NEEDS CONFIRMATION, OMITTED FROM TABLE.
...
43 PI(2050) MASS (MEV)
...
43 PI(2050) WIDTH (MEV)
...
43 PI(2050) PARTIAL DECAY MODES
...
43 PI(2050) BRANCHING RATIOS
...
43 PI(2050) INT0 PI+ PI- PI0
...
43 PI(2050) INT0 G G
...
43 PI(2050) BRANCHING RATIOS
...
43 PI(2050) INT0 PI+ PI- PI0
...
43 PI(2050) INT0 G G
...
43 PI(2050) BRANCHING RATIOS
...
43 PI(2050) INT0 PI+ PI- PI0
...

Mesons

$\pi(2050)$, $\pi(2100)$, T AND U REGIONS, $\rho(2150)$ For notation, see key at front of Listings.

REFERENCES FOR $\pi(2050)$

HUSON 68 PL 28 B 208 +LUBATTI,BELLINI,BINGHAM,+ (ORSA+MLA+LBL)
 BEMGRAD 71 NP B 33 357 +DUFFY,COOLING,+ (CERN+ETH+LOIC+NL)
 CLAYTON 72 NP B 47 81 +MASON,NIR,HEAD,RIGOPPULOS,+ (LIVP+PATR)
 BASTIEN 73 JPD54A CCF. 73 +DUNN,HARRIS,LUBATTI,BINGHAM,+ (SEAT+UCR)
 OREN 74 NP B71 154 +COOPER,FIELDS,RHINES,WHITMORE,+ (ANL+OZF)
 DEUTSCHMANN 75 NP 899 397 DEUTSCHMANN,+ (ARBCCHM COLLABORATION)
 KALELKA 75 THESIS NEVIS 207 M.S.,KALELKA (COLU)
 ANTIPOV 77 NP B 119 45 +BUSNELLO,DAMGAARD,KIENZLE+ (CERN+SERP)
 BALTAY 77 PR 39 551 +CAUTIS,KALELKA (COLUMBIA) JP
 CAUTIS 77 THESIS NEVIS 221 C.-J.CAUTIS (CDU+UB) JP
 BALTAY 78 PR D 17 52 +CAUTIS,COHEN,CSDRHA,KALELKA+ (CD+UP+SI)

 $\pi(2100)$

20 $\pi(2100)$, JPC=2-1-1
 SEEN IN THE π ND π 11, EPSILON π 11 AND ρ π 11
 JP = 2- WAVES OF THE DIFFRACTIVELY PRODUCED π π
 SYSTEM. NEEDS CONFIRMATION. OMITTED FROM TABLE.

20 $\pi(2100)$ MASS (MEV)

M	L	2100.	150.	DAUM	BI CNTR	63.94	PI-	P, π 31	1/82*
M	L	FROM A TWO RESONANCE FIT TO FOUR 2-0+ WAVES.							

20 $\pi(2100)$ WIDTH (MEV)

M	L	651.	50.	DAUM	BI CNTR	63.94	PI-	P, π 31	1/82*
M	L	FROM A TWO RESONANCE FIT TO FOUR 2-0+ WAVES.							

20 $\pi(2100)$ PARTIAL DECAY MODES

P1	$\rho(2100)$	INTO π PI	DECAY MODES
01	$\rho(2100)$	INTO π PI	13% 130 130
02	$\rho(2100)$	INTO π ND π 11	70% 130
03	$\rho(2100)$	INTO π π 11	12% 130
04	$\rho(2100)$	INTO ρ S1100 π 11	150% 130

20 $\pi(2100)$ BRANCHING RATIOS

R1	$\pi(2100)$	INTO π ND π 11 (ALL π PI)	BRANCHING RATIOS
01	L	0.19 0.05	DAUM BI CNTR 63.94 PI- P 1/82*
02	L	0.36 0.09	DAUM BI CNTR 63.94 PI- P 1/82*
03	L	0.45 0.07	DAUM BI CNTR 63.94 PI- P 1/82*
04	L	0.35 0.23	DAUM BI CNTR 63.94 PI- P 1/82*

R L FROM A TWO RESONANCE FIT TO FOUR 2-0+ WAVES.

REFERENCES FOR $\pi(2100)$

DAUM BI NP B 182 269 +HERTZBERGER+(ANST+CERN+CRAC+IMP+INDOFR+HREL)

Note on T and U Regions

The observation of broad enhancements at 2190 and 2350 MeV comes from $\bar{p}p$ total cross-section measurements (ABRAMS 67), $\bar{p}p$ annihilation measurements (ALSPECTOR 73), $\bar{p}p$ elastic cross-section measurements (COUPLAND 77), and $\bar{p}p$ charge-exchange cross-section measurements (CUTTS 78). The mass regions centered around 2190 MeV and 2350 MeV have been called T and U regions, respectively.

Searches for resonances in exclusive $\bar{p}p$ annihilation channels which could be coupled to the enhancements observed in the $\bar{p}p$ total cross-section and in $\bar{p}p$ elastic scattering have been unsuccessful, except for the two-body annihilation

Data Card Listings

channels $\pi^+\pi^-$ and $\pi^0\pi^0$, where partial-wave analyses have shown that resonances are formed in the 2100-2500 MeV mass region (CARTER 77, DULUDE 78, MARTIN A 80, MARTIN B 80). We have listed the results on these analyses under the headings $c(2150)/T_0(2150)$ for the I=0, JP=2+ wave; $\rho(2150)/T_1(2190)$ for the I=1, JP=1- wave; $\rho(2250)$ for the I=1, JP=3- wave; $c(2300)/U_0(2300)$ for the I=0, JP=4+ wave; and $\rho(2350)/U_1(2400)$ for the I=1, JP=5- wave.

Various structures coupled to $\bar{p}p$ and observed in production experiments are listed under the heading $\bar{N}N(1400-3600)$.

 $\rho(2150)$

22 $\rho(2150)$, JPC=1-1-1
 THIS ENTRY WAS PREVIOUSLY CALLED $\pi(2150)$.
 CONTAINS ONLY RESULTS FROM PRODUCTION EXPERIMENTS. FOR
 PRODUCTION EXPERIMENTS SEE THE NEAR $\pi(1400-3600)$ ENTRY.
 SEE ALSO π -U MINI-REVIEWS.
 OMITTED FROM TABLE.

32 $\rho(2150)$ MASS (MEV)

M	PBAR P	INTO π 11	PI	PI
M	P	(2100.0) APPROX.	MARTIN A 80 PVUE	1/82*
M	P	(2120.0) APPROX.	MARTIN B 80 PVUE	1/82*
M	P	I=1, JP=1- FROM SIMULTANEOUS ANALYSIS OF P \bar{P} \rightarrow π 1- π 1 AND π 0 π 0		
M	S	CHANNEL NUCLEON ANTI-NUCLEON		
M	B	2190. 10.	ABRAMS 70 CNTR	S CHANNEL PBAR N 1/73
M	I	2193. 2.	ALSPECTOR 73 CNTR	S CHANNEL PBAR P 1/74
M	E	2155.0 15.0	COUPLAND 77 CNTR	0.1-2.4PB-P, PB-P 12/77
M	I	(2150.0) APPROX.	CUTTS 78 CNTR	0.7-5. PB-P, AND N 12/78
M	AVERAGE MEANINGLESS SCALE FACTOR = 1.0			
M	B	SEEN AS BUMP IN I=1 STATE. SEE ALSO COOPER 68.		
M	E	PEASLEE 75 CONFIRM PBAR P RESULTS OF ABRAMS 70 AND NARROW STRUCTURE		
M	E	FROM A FIT TO THE TOTAL ELASTIC CROSS SECTION.		
M	I	1 ISOSPINS 0 AND 1 NOT SEPARATED		

32 $\rho(2150)$ WIDTH (MEV)

M	PBAR P	INTO π 11	PI	PI
M	P	1200.0) APPROX.	MARTIN A 80 PVUE	1/82*
M	P	(2200.0) APPROX.	MARTIN B 80 PVUE	1/82*
M	P	I=1, JP=1- FROM SIMULTANEOUS ANALYSIS OF P \bar{P} \rightarrow π 1- π 1 AND π 0 π 0		
M	S	CHANNEL NUCLEON ANTI-NUCLEON		
M	B	(185.) APPROX.	ABRAMS 70 CNTR	S CHANNEL PBAR N 7/67
M	I	98. 8.	ALSPECTOR 73 CNTR	S CHANNEL PBAR P 1/74
M	E	134.0 75.0	COUPLAND 77 CNTR	0.7-2.4PB-P, 0-P 12/77
M	AVERAGE MEANINGLESS SCALE FACTOR = 1.0			
M	B	SEE NOTE B ABOVE.		
M	E	FROM A FIT TO THE TOTAL ELASTIC CROSS SECTION.		
M	I	1 ISOSPINS 0 AND 1 NOT SEPARATED		

REFERENCES FOR $\rho(2150)$

ABRAMS 67 PR 18 1239 +COUL,GIACOMELLI,KYCIA,LEONTIC,LI,+ (BNL)
 COOPER 68 PR 20 1659 +HYMAN,MANNER,MUGRAGE,VOYVODIC (ANL)
 BRICHAN 69 PL 29 B 451 +FERRO-LUZZI,BIARDI,+ (CERN+CAEN+SACL)
 ABRAMS 70 PR D 1 1517 +COUL,GIACOMELLI,KYCIA,LEONTIC,LI,+ (BNL)
 BACON 71 NP B 32 46 +BUFFERWORTH,NIELER,PHILAN,+ (RHEN-LIVP)
 FIELDS 71 PR 27 3749 +COOPER,RHINES,ALLISON (ANL+OZF)
 YDM 71 PR 26 722 +BARISH,CARDILL,LEONKVIC+ (CIT+BNL+RUC)
 ALEXANDER 72 NP B 45 29 ALEXANDER,BAR-NIR,BEVARY,DAGAN,+ (TEL)
 DONALD 72 PL 40 B 586 +GALLETTY,EDWARDS,DE BILLY,+ (LIVP+LMP)
 ALSPECTOR 73 PR 30 511 ALSPECTOR,COHEN,CVIJANGVIC,+ (FRUC+UPNA)
 BACON 73 PR D 7 577 +BUFFERWORTH, (RHEN-LIVP)
 BEGLINI 73 PR 15 A 303 +GARRINOST,REGI,+ (PADDO+LIP+ST+TOR)
 DONALD 73 NP B 61 353 +EDWARDS,CERRINS,BRIAND,DOBOS,+ (LIVP+LMP)
 NICHOLSON 73 PR D 7 2572 NICHOLSON,DELOURNE,CARROLL,+ (CIT+RUC+BNL)
 BERTANZA 74 NC 23A 209 +RIGI,CASATI,LARTICIA,+ (PISA+PADO+TOR)
 HYAMS 74 NP B 73 202 +JONES,WEILHANNER,OLSON,+ (CERN+MPI)
 DONNACHE 75 NC 26 A 317 A.DONNACHE,P.R.THOMAS (FRANCKSTERT)
 EISENHAN 75 NP B 96 109 EISENHAN,ELER,ELER,ELER,+ (LONDON+LIVP+DARE+HREL)
 HAYDLER 75 NP B 101 35 +JACQUES,JONES,PANOUILLAS,+ (FRUC+ST+VALRA)
 HUESMAN 75 NC 25A 91 +GARRINOST,ROSS,+ (LBU+PADDO+ST+TOR)
 PEASLEE 75 PL 57B 189 +DEBARZO,GUERRERO,+ (CERN+BAR+BRNO+MIT)

Data Card Listings

For notation, see key at front of Listings.

Mesons

$\rho(2150)$, $\epsilon(2150)$, $\rho(2250)$, $\epsilon(2300)$

GAY 76 NC 31 A 593
ZENAMY 76 NP B 103 537

CARTER 1 77 PL 67 B 117
CARTER 2 77 PL 67 B 122
CARTER 3 77 NP B 127 202
COUPLAND 77 PL 71 B 400
JONES 77 NP B 116 476
MONTANET 77 BOSTON CONF. 260

CARTER 1 78 NP B 132 176
CARTER 2 78 NP B 141 467
CUTTS 78 PR D 17 16

MARTIN 79 PL 86 B 93

MARTIN A 80 NP B 169 216
MARTIN B 80 NP B 176 355

JEANNERET,BOGDANSKI,+NEUC+LUS+LTP+LNPJ
+MING MOUNTAIN,SMITH (MSU)
+COUPLAND,EISENHANDLER,ASTBURY,+ILOOM+RHELI JP
A.A.CARTER (LLOOM) JP
+COUPLAND,ATKINSON,ARNISON,(LOOM+DRE+RHELI) JP
+EISENHANDLER,GIBSON,ASTBURY,+ (ILOOM+RHELI) JP
M.D.JONES,P.R.JPLANO (RUTGZ)
L.MONTANET (ICERN)
A.A.CARTER (LLOOM) JP
A.A.CARTER (LLOOM) JP
+GODO,GRANNIS,GREEN,LEE,PITTMAN+STON+MISCI (LLOOM) JP
A.D. MARTIN,M.R. PENNINGTON (DURH) JP
A.D. MARTIN,M.R. PENNINGTON (DURH) JP
B.R.MARTIN,D.MORGAN (ILOUC+RHELI) JP

$\epsilon(2150)$

42 EPSILON(2150).JP(2+1) 1=0
THIS ENTRY HAS PREVIOUSLY CALLED TOI2150.
CONTAINS ONLY RESULTS FROM FORMATION EXPERIMENTS, FOR
PRODUCTION EXPERIMENTS SEE THE NEAR N1400-3600 ENTRY.
SEE ALSO S,T,U MINI-REVIEWS.
OMITTED FROM TABLE.

42 EPSILON(2150) MASS (MEV)

M PBAR P INTO PI PI
M L (2150.0) APPROX. DULUDEZ 78 DSPK 1.-2.PB P.PIOPD 12/78
M K (2150.0) APPROX. MARTIN A 80 RVUE 1/82*
M P (2170.0) APPROX. MARTIN B 80 RVUE 1/82*

M L (G₀)+J₀+2* FROM PARTIAL WAVE AMPLITUDE ANALYSIS
M P 1=0,+J₀+2* FROM SIMULTANEOUS ANALYSIS OF P B → PI-PI AND P10 P10

M S CHANNEL PBAR P OR NBAR N
M I 2193. 2. ALSPECTOR 73 CNTR S CHANNEL PBAR N 1/74
M E 1 2155.0 15.0 COUPLAND 77 CNTR O -J-2.4PB-P,PB-P 12/77
M I 2150.0) APPROX. CUTTS 78 CNTR -07-3. PB P,NB N 12/78

M AVERAGE MEANINGLESS (SCALE FACTOR = 1.0)
M E FROM A FIT TO THE TOTAL ELASTIC CROSS SECTION.
M I ISOSPINS 0 AND I NOT SEPARATED

42 EPSILON(2150) MASS (MEV)

M PBAR P INTO PI PI
M L (250.0) APPROX. DULUDEZ 78 DSPK 1.-2.PB P.PIOPD 12/78
M K (250.0) APPROX. MARTIN A 80 RVUE 1/82*
M P (250.0) APPROX. MARTIN B 80 RVUE 1/82*

M L (G₀)+J₀+2* FROM PARTIAL WAVE AMPLITUDE ANALYSIS
M P 1=0,+J₀+2* FROM SIMULTANEOUS ANALYSIS OF P B → PI-PI AND P10 P10

M S CHANNEL PBAR P OR NBAR N
M I 78. 75.0 ALSPECTOR 73 CNTR S CHANNEL PBAR P 1/74
M E 1 135.0 B. COUPLAND 77 CNTR O -J-2.4PB-P,PB-P 12/77

M AVERAGE MEANINGLESS (SCALE FACTOR = 1.0)
M E FROM A FIT TO THE TOTAL ELASTIC CROSS SECTION.
M I ISOSPINS 0 AND I NOT SEPARATED

REFERENCES FOR EPSILON(2150)

FIELDS 71 PR 27 1749 (ANL+DF)
VDM 71 PR 26 922 +BARISH,CAROLL,LOBROVICZ+ (CIT+BNL+RUCN)
ODNALD 77 PL 40 B 586 +GALLELY,EDWARDS,DE BILLY,+ (LTP+LNPJ)
ALSPECTOR 73 24 30 511 +ALBERT JR.,COHEN,CVJANOVICH,+ (RUTG+UPNJ)
RACON 73 PR D 7 577 +BUTENMUTH, (RHELI+LVP)
ODNALD 77 PR D 01 1333 +EDWARDS,GIBSON,BRIANO,RUBOC,+ (LTP+LNPJ)
NICHOLSON 73 PR D 7 2572 NICHOLSON,DELORE,CARROLL,+ (CIT+RCH+BNL)

GAY 76 NC 31 A 593
COUPLAND 77 PL 71 B 400
CUTTS 78 PR D 17 16
DULUDEZ 78 PL 79 B 129
DULUDEZ 78 PL 79 B 335
MARTIN 79 PL 86 B 93
BONCHECK 80 LC 28 21
MARTIN A 80 NP B 169 216
MARTIN B 80 NP B 176 355

JEANNERET,BOGDANSKI,+NEUC+LUS+LTP+LNPJ
+EISENHANDLER,GIBSON,ASTBURY,+ (LOOM+RHELI) JP
+GODO,GRANNIS,GREEN,LEE,PITTMAN+STON+MISCI
+LANDAU,MASSIMO,PEASELE+ (BROMNIT+EARL) JP
+LANDAU,MASSIMO,PEASELE+ (BROMNIT+EARL) JP
A.D. MARTIN,M.R. PENNINGTON (DURH) JP
A.D. MARTIN,M.R. PENNINGTON (DURH) JP
B.R.MARTIN,D.MORGAN (ILOUC+RHELI) JP

$\rho(2250)$

44 RHO(2250).JP(3+1) 1=1
CONTAINS ONLY RESULTS FROM FORMATION EXPERIMENTS, FOR
PRODUCTION EXPERIMENTS SEE THE NEAR N1400-3600 ENTRY.
SEE ALSO S,T,U MINI-REVIEWS.
OMITTED FROM TABLE.

44 RHO(2250) MASS (MEV)

M PBAR P INTO PI PI OR KB K
M J (2150.0) APPROX. CARTER 1 77 CNTR O -J-2.4PB P,P1P1 12/77
M K (2140.0) APPROX. CARTER 2 78 CNTR O -J-2.4PB P,K+K+ 12/78
M E 1 (2300.0) APPROX. MARTIN A 80 RVUE 1/82*
M P (2250.0) APPROX. MARTIN B 80 RVUE 1/82*

44 RHO(2250) MASS (MEV)

M PBAR P INTO PI PI OR KB K
M J (2150.0) APPROX. CARTER 1 77 CNTR O -J-2.4PB P,P1P1 12/77
M K (2140.0) APPROX. CARTER 2 78 CNTR O -J-2.4PB P,K+K+ 12/78
M E 1 (2300.0) APPROX. MARTIN A 80 RVUE 1/82*
M P (2250.0) APPROX. MARTIN B 80 RVUE 1/82*

M J 1=1,+J₀+3* FROM AMPLITUDE ANALYSIS.
M K 1=0,+J₀+3* FROM BARRELET ZERO'S ANALYSIS.
M P 1=1,+J₀+3* FROM SIMULTANEOUS ANALYSIS OF P B → PI-PI AND P10 P10

M S CHANNEL NUCLEON ANTI NUCLEON

M B 2190. 10. ABRAMS 70 CNTR S CHANNEL PBAR N 1/73
M I 2195. 2. ALSPECTOR 73 CNTR S CHANNEL PBAR P 1/74
M E 1 2155.0 15.0 COUPLAND 77 CNTR O -J-2.4PB-P,PB-P 12/77
M I (2190.0) APPROX. CUTTS 78 CNTR -07-3. PB P,NB N 12/78

M AVERAGE MEANINGLESS (SCALE FACTOR = 1.0)
M B SEEN AS BUMP IN I=1 STATES. SEE ALSO COOPER 80.
M B PEASLEE 75 CONFERN PBAR P RESULTS OF ABRAMS 70,ND NARROW STRUCTURE
M E FROM A FIT TO THE TOTAL ELASTIC CROSS SECTION.
M I ISOSPINS 0 AND I NOT SEPARATED

44 RHO(2250) MIDD (MEV)

M PBAR P INTO PI PI OR KB K
M J (2200.0) APPROX. CARTER 1 77 CNTR O -J-2.4PB P,P1P1 12/77
M K (1150.0) APPROX. CARTER 2 78 CNTR O -J-2.4PB P,K+K+ 12/78
M E 1 (2200.0) APPROX. MARTIN A 80 RVUE 1/82*
M P (2250.0) APPROX. MARTIN B 80 RVUE 1/82*

M J 1=1,+J₀+3* FROM AMPLITUDE ANALYSIS.
M K 1=0,+J₀+3* FROM BARRELET ZERO'S ANALYSIS.
M P 1=1,+J₀+3* FROM SIMULTANEOUS ANALYSIS OF P B → PI-PI AND P10 P10

M S CHANNEL NUCLEON ANTI NUCLEON

M B (85. 1) APPROX. ABRAMS 70 CNTR S CHANNEL PBAR N 7/67
M I 98. 8. ALSPECTOR 73 CNTR S CHANNEL PBAR P 1/74
M E 1 135.0 75.0 COUPLAND 77 CNTR O -J-2.4PB-P,PB-P 12/77

M AVERAGE MEANINGLESS (SCALE FACTOR = 1.0)
M B SEE NOTE P ABOVE.
M E FROM A FIT TO THE TOTAL ELASTIC CROSS SECTION.
M I ISOSPINS 0 AND I NOT SEPARATED

REFERENCES FOR RHO(2250)

ABRAMS 67 PR 18 1209 +COUPLAND,EISENHANDLER,ASTBURY,+ (LOOM+RHELI) JP
COOPER 68 PR 20 1359 +HYMAN,MANNER,MUSGRAVE,VOYVODIC (ANL)
ABRAMS 70 PR D 1 1517 +COUPLAND,EISENHANDLER,ASTBURY,+ (LOOM+RHELI) JP
FIELDS 71 PR 27 1749 +COOPER,RHINES,ALLISON (ANL+DF)
VDM 71 PR 26 922 +BARISH,CAROLL,LOBROVICZ+ (CIT+BNL+RUCN)

ALSPECTOR 73 24 30 511 +ALBERT JR.,COHEN,CVJANOVICH,+ (RUTG+UPNJ)
RACON 73 PR D 7 577 +BUTENMUTH, (RHELI+LVP)
ODNALD 77 PR D 01 1333 +EDWARDS,GIBSON,BRIANO,RUBOC,+ (LTP+LNPJ)
NICHOLSON 73 PR D 7 2572 NICHOLSON,DELORE,CARROLL,+ (CIT+RCH+BNL)

BERTANZA 74 NC 23A 209 +BGI,CASAL,LARICIA,+ (FISA+PADO+TUSI)
ZENAMY 76 NP B 103 537 +MING MOUNTAIN,SMITH (MSU)

CARTER 1 77 PL 67 B 117 +COUPLAND, EISENHANDLER,ASTBURY,+ (LOOM+RHELI) JP
CARTER 2 77 PL 67 B 122 A.A.CARTER (LLOOM) JP
CARTER 3 77 NP B 127 202 +COUPLAND,ATKINSON,ARNISON+(LOOM+DRE+RHELI) JP
MONTANET 77 BOSTON CONF. 260 L.MONTANET (ICERN)
CARTER 1 78 NP B 132 176 A.A.CARTER (LLOOM) JP
CARTER 2 78 NP B 141 467 A.A.CARTER (LLOOM) JP
CUTTS 78 PR D 17 16 +GODO,GRANNIS,GREEN,LEE,PITTMAN+STON+MISCI (LLOOM) JP
MARTIN 79 PL 86 B 93 A.D. MARTIN,M.R. PENNINGTON (DURH) JP
MARTIN A 80 NP B 169 216 A.D. MARTIN,M.R. PENNINGTON (DURH) JP
MARTIN B 80 NP B 176 355 B.R.MARTIN,D.MORGAN (ILOUC+RHELI) JP

$\epsilon(2300)$

41 EPSILON(2300).JP(4+1) 1=0
THIS ENTRY HAS PREVIOUSLY CALLED UOI2300.
CONTAINS ONLY RESULTS FROM FORMATION EXPERIMENTS, FOR
PRODUCTION EXPERIMENTS SEE THE NEAR N1400-3600 ENTRY.
SEE ALSO S,T,U MINI-REVIEWS.
OMITTED FROM TABLE.

41 EPSILON(2300) MASS (MEV)

M PBAR P INTO PI PI OR KB K
M J (2310.0) APPROX. CARTER 1 77 CNTR O -J-2.4PB P,P1P1 12/77
M K (2340.0) APPROX. CARTER 2 78 CNTR O -J-2.4PB P,K+K+ 12/78
M E 1 (2300.0) APPROX. MARTIN A 80 RVUE 1/82*
M P (2300.0) APPROX. MARTIN B 80 RVUE 1/82*

M J 1=0,+J₀+4* FROM AMPLITUDE ANALYSIS.
M K 1=0,+J₀+4* FROM BARRELET ZERO'S ANALYSIS
M P 1=0,+J₀+4* FROM SIMULTANEOUS ANALYSIS OF P B → PI-PI AND P10 P10

M S CHANNEL PBAR P OR NBAR N
M I 2375. 10. ABRAMS 70 CNTR S CHANNEL PBAR N 1/71
M E 1 (2350. 1) ALSPECTOR 73 CNTR S CHANNEL PBAR P 1/74
M I (2295.0) APPROX. COUPLAND 77 CNTR O -J-2.4PB-P,PB-P 12/77
M I (2380.0) APPROX. CUTTS 78 CNTR -07-3. PB P,NB N 12/78

M E FROM A FIT TO THE TOTAL ELASTIC CROSS SECTION.
M I ISOSPINS 0 AND I NOT SEPARATED

Mesons J/ψ(3100)

Data Card Listings For notation, see key at front of Listings.

TO J/PSI(3100) BRANCHING RATIOS

FOR THE BRANCHING RATIOS R1 - R4. SEE ALSO THE PARTIAL WIDTHS ABOVE, AND (PARTIAL WIDTHS)/R1 BELOW.

Table containing branching ratios for J/ψ(3100) decays into various mesons and baryons. Columns include particle IDs (R1, R2, etc.), decay channels (e.g., J/PSI(3100) INTO E+ E-), branching ratios, and statistical errors.

Table containing branching ratios for J/ψ(3100) decays into baryon pairs (PBAR P) and other particles. Columns include particle IDs (R31, R32, etc.), decay channels (e.g., J/PSI(3100) INTO P(BAR P)), branching ratios, and statistical errors.

Data Card Listings

Mesons

For notation, see key at front of Listings.

J/ψ(3100), χ(3415)

J/PSI INTO P BAR ETA PRIME/TOT (UNITS 10**3) (P50)
R62 19 1.8 0.6 PERUZZI 78 SWAG E-E-P PR 1-2PJ 4/78
R63 J/PSI INTO EN NBAR PI+ PI-/TOTAL (UNITS 10**3) (P60)
R63 5 3.8 3.6 BESCH 81 BDM A E-E- 1/82**
R64 J/PSI INTO E(S1GM+ S1GMABAR)/TOT (UNITS 10**3) (P61)
R64 3 2.4 2.6 BESCH 81 BDM A E-E- 1/82**
R J FINAL STATE 2(XI)PI(2)
A ASSUMING ANGULAR DISTRIBUTION I(L+COS^2(TXETA)*R2)
R RADIATIVE DECAYS
R71 J/PSI(3100) INTO E2 GAMMA/TOT (UNITS 10**3) (P70)
R71 10 5.3 DR LESS CL+0.90 BARTEL 77 CNTR E-E- 4/77
R72 J/PSI(3100) INTO E(PID GAMMA)/TOTAL (UNITS 10**3) (P72)
R72 10 9.33 DR LESS CL+0.79 OAKSF E- E- 12/79
R73 J/PSI(3100) INTO ETA GAMMA/TOTAL (UNITS 10**3) (P73)
R73 21 1.3 0.4 BARTEL 77 CNTR E-E- 3 GAMMA 1/77
R74 0.02 0.10 BRANDELI 79 DASP E- E- 12/79
R75 0.08 0.19 KINGSMAN 82 CBAL E- E- 3 GAMMA 4/82**
R76 SYSTEMATIC ERROR ADDED LINEARLY BY US.
R77 - + - + - + -
R78 AVERAGE ERROR INCLUDES SCALE FACTOR OF 1.0)
R78 AVG 0.88 0.086
R79 J/PSI(3100) INTO E(TA PRIME GAM)/TOT (UNITS 10**3) (P74)
R79 13(2) DR LESS CL+0.90 MURTAG 76 FRAG E-E- 4/77
R79 57 (24) (0.7) BARTEL 77 CNTR E-E- 2 GAMMA RHD 1/77
R79 6 2.4 0.6 BRANDELI 79 DASP E-E- 3 GAMMA 12/79
R79 6 3.8 1.3 SCHARNE 79 SWAG E-E- 2 P10 GAMMA 12/79
R79 FROM THE INCLUSIVE GAMMA DECAY SPECTRUM
R79 3 SCHARNE 79 SWAG E-E- 2 P10 GAMMA 12/79
R79 4.1 0.9 KINGSMAN 82 CBAL E- E- 4/82**
R79 SYSTEMATIC ERROR ADDED LINEARLY BY US.
R79 - + - + - + -
R79 AVERAGE ERROR INCLUDES SCALE FACTOR OF 1.0)
R79 AVG 3.55 0.46
R80 J/PSI(3100) INTO J13 GAMMA/TOT (UNITS 10**3) (P71)
R80 10 7.8 DR LESS CL+0.90 PARTRIDGE 80 CNTR E-E- 3 GAMMA 12/79
R81 J/PSI(3100) INTO [GAMMA + 2 DR MORE NEUTRALS]/TOTAL (UNITS 10**3) (P71)
R81 7.0 2.0 BARTEL 77 CNTR E-E- 10**3 1/77
R81 J/PSI(3100) INTO E(GAMMA TOT)/TOT (UNITS 10**3) (P76)
R81 35 2.0 0.7 ALEXANDER 78 PLUT 0 E-E- 4/78
R81 30 1.2 0.4 BRANDELI 79 DASP E-E- P10+GAMMA 4/78
R81 RE-STARTED BY US. TO TAKE ACCOUNT OF SPECTRA OF 3 E-E- TRANSITIONS
R81 178 1.48 0.55 EDWARDS 82 CBAL E-E- 2 P10 GAMMA 2/82**
R81 SYSTEMATIC ERROR ADDED LINEARLY BY US.
R81 - + - + - + -
R81 AVERAGE ERROR INCLUDES SCALE FACTOR OF 1.0)
R81 AVG 1.51 0.35
R82 J/PSI(3100) INTO E(PI PRIME GAM)/TOT (UNITS 10**3) (P77)
R82 5 (0.23) CR LESS CL+0.90 ALEXANDER 78 PLUT E-E-KK+KAM 4/78
R82 4 (0.34) DR LESS CL+0.90 BRANDELI 79 DASP E-E- P10+GAMMA 12/79
R82 ASSUMING ISOTROPIC PRODUCTION AND DECAY OF THE P PRIME AND ISOSPIN.
R84 J/PSI(3100) INTO (D11285) GAMMA/TOT (P78)
R84 10.11 DR LESS CL+0.90 PERUZZI 78 SWAG E-E-P PR SWMR 4/78
R85 J/PSI(3100) INTO (D11285) GAMMA/TOT (P78)
R85 0 10.00 DR LESS CL+0.90 SCHARNE 80 SWAG E-E- 2/81**
R85 USING BRIDGE INFO K BAR PI+0.12
R86 J/PSI(3100) INTO [E1420] GAMMA/TOT (P79)
R86 E (0.0055) 10.00221 SCHARNE 80 SWAG E-E- 2/81**
R86 F USING BRIDGE INFO K BAR PI+0.05 SEE MINI-REVIEW UNDER [E1420]
NOTE THAT EVIDENCE INDICATES THAT THIS IS NOT E JPC=1++
[E1420] HENCE BR INTG K BAR PI+ UNKNOWN. SEE [E1420] MINI-REVIEW

TO J/PSI(3100) G(I)+GE+E-1/G(TOTAL) (KEV)

THIS COMBINATION OF A PARTIAL WIDTH WITH THE PARTIAL WIDTH INTO E-E- AND WITH THE TOTAL WIDTH IS OBTAINED FROM THE INTEGRATED CROSS-SECTION INTO CHANNEL(I) IN THE E-E- ANNIHILATION. WE ONLY LIST DATA NOT HAVING BEEN USED TO DETERMINE THE PARTIAL WIDTH G(I) OR THE BRANCHING RATIO G(I)/TOTAL.

Table with columns for ID, G(E-E) * GE-E-1/G(TOTAL), and decay mode details.

Table with columns for ID, G(SI) * GE(E)-1/G(TOTAL), and decay mode details.

SEE THE BRANCHING RATIOS AND PARTIAL WIDTHS ABOVE.

REFERENCES FOR J/PSI(3100)

CHRISTEN TO PRL 25 1523
CHRISTENSON, HICKS, LEDERMAN + COLU(BN+K+ERN)
ABRAMS 75 PRL 33 1453
ASH 74 NCL 11 1705
AUGUSTIN 74 PRL 33 1450
AUGUSTIN 74 PRL 33 1456
BACCI 74 PRL 33 1430
FOR CARAN
BALDINI 74 NCL 11 710
BARBIELL 74 NCL 11 710
BRÄUNSCHE 74 PRL 538 303

ANDREWS 75 PRL 34 231
AUBERT 75 NP B 891
BACCI 75 NCL 12 269
BALDINI 75 PRL 508 471
BRÄUNSCHE 75 PRL 538 303
BRÄUNSCHE 75 STANFORD SYMP.113 C
DEMPARD (PISA+FRASCATI)
BOYER,FAISLER,GANELICH,GETTNER + INEAS
BRONKHORST,DEGROOT,ROOZ,CLAVER,DEH
BRÄUNSCHE 75 PRL 538 303
BUSSER 75 PRL 50 842
CAMERINI 75 PRL 538 303
CRIGGEE 75 PRL 538 499
EDWARDS 75 PRL 50 842
DASPI 75 PRL 50 842
POSSIDIO 75 NCL 14 78
FORD 75 PRL 34 604
GITTELMAN 75 PRL 53 1616
GRICE 75 PRL 30 1357
HEINZ 75 STANFORD SYMP.07
JACOBSON,DEBASTA,STRAVATA FRAS
HEINZ 75 STANFORD SYMP.07
JACOBSON,DEBASTA,STRAVATA FRAS
ANAPOL 75 PRL 34 1040
KNAPP 75 PRL 34 1046
LIBERMAN 75 STANFORD SYMP.55
ALEX,BRONSTEIN + COLU(BN+K+ERN+ILL+FNAL)
MARTIN 75 PRL 34 288
PREPOSTI 75 STANFORD SYMP.241
SIMPSON 75 PRL 35 499
WIK 75 STANFORD SYMP.06
YENNIE 75 PRL 34 239
BACCI 76 NCL 12 269
BALDINI 76 PRL 64 843
BRÄUNSCHE 76 PRL 63 487
BUSSER 76 NP B 140 899
JEAN-MAR 76 PRL 30 201
KINGSMAN 76 TELLIS CONF. 246
PIERRE 76 TELLIS CONF. 246
SMYTER 76 PRL 36 1415
BARTEL 77 PL 66 B 489
BIDDICK 77 PRL 38 1324
BURMESTER 77 PL 37 303
COWEN 77 PL 68 B 56
FELDMAN 77 PL 33 295
KANEVICH 77 PL 68 B 104
YAMADA 77 HAMB. CONF. P. 69
ALEXANDER 78 PL 72 B 403
BESCH 78 PL 78 B 347
BRANDELI 78 PL 78 B 292
PERUZZI 78 PL 87 1720)
BRANDELI 79 ZPH C 1 233
KIKR 79 SLAC 42 619
LEONDIJE 79 FERMI CONF. 524
SCHARNE 79 PRL 207-2321
ALSO 79 LBL 4502
SCHARNE 80 PRL 44 712
SCHARNE 80 PRL 97 529
MORRISON,LEICHHEN,LEISHEN,NEFTIN
ALSO BL VAD. PHYS. 34 147)
HOLTZ ET AL.
FUDSI
BESCH 81 ZPH C 8 1
FESTERMANN,LOHN,KOWALSKI + (BOVI+DESY+MANZ)
BARATE 82 MPRIO WORKSHOP
ALSO 82 CEPN-EP82-15
LEONDIJE, BARATE + ISGL(LDC+SIS+INDI
LEONDIJE 82 NP D 68 PUB.1
KONIGSMAN 82 MPRIO CONF.

χ(3415)

56 CHIE(3415, JP00+)=1=0

OBSERVED IN THE RADIATIVE DECAY OF PSI(3368S) INTO
CHI(3415) GAMMA. THEREFORE C+=. THE OBSERVED DECAY INTO PI+ PI-
OR K+ K- IMPLIES G+=. JP=0+2---. THE ANGULAR DISTRIBUTION IS
CONSISTENT WITH J=0, JP ANOMALOUS EXCLUDED BY PI+ PI- AND
K+ K- DECAYS. JP=0+0, JP ANOMALOUS EXCLUDED BY FELDMAN 77 3

Table with columns for ID, G(SI) * GE(SI) * MASS (MEV), and decay mode details.

56 CHIE(3415) PARTIAL DECAY MODES

Table with columns for ID, CHI(3415) INTO P+ PI-, and decay mode details.

56 CHIE(3415) BRANCHING RATIOS

Table with columns for ID, CHI(3415) INTO 12 GAMMA/TOT, and decay mode details.

Mesons

$\chi(3415)$, P_c or $\chi(3510)$, $\chi(3555)$

Data Card Listings

For notation, see key at front of Listings.

Table with columns for particle name (e.g., $\chi(3415)$), fit type (e.g., INTG), parameters (e.g., Γ , θ), and fit quality (e.g., χ^2/N). Includes fit results for various mesons like $\chi(3415)$, $\chi(3510)$, and $\chi(3555)$.

REFERENCES FOR $\chi(3415)$

FELOMAN 76 PRL 35 821 +JEAN-MARIE, SADOULET, VANHUCCI, + (LBE+SLAC)
ALSO 75 PRL 35 1189 (ERRATA)
TANENBAU 75 PRL 35 1323 TANENBAUM, WHITAKER, ARMS+ (LBE+SLAC)
WIK 75 STANFORD SMP.69 B.W.WIK (DESY)

Table with columns for particle name (e.g., $\chi(3555)$), fit type (e.g., INTG), parameters (e.g., Γ , θ), and fit quality (e.g., χ^2/N). Includes fit results for $\chi(3555)$ and other mesons.

REFERENCES FOR $\chi(3555)$

DASP 76 PL 378 407 BRAUNSCHWEIG, KONTIGS, + (AACH+DESY+MPI+MINK)
FELOMAN 75 STANFORD SMP.39 C.J.FELOMAN (SLAC)
HEINITZE 75 STANFORD SMP.97 HEINITZE (MIEDELBERG)
SIMPSON 75 PRL 35 862 +BERON, FORD, HILGER, HOFSTADTER, + (SLAC+MENN)
TANENBAU 75 PRL 35 1323 TANENBAUM, WHITAKER, ARMS+ (LBE+SLAC)
WIK 75 STANFORD SMP.69 B.W.WIK (DESY)

P_c or $\chi(3510)$

55 PC DR CH(3510).JPD+1+1 J=0
OBSERVED IN THE RADIATIVE SCALED DECAY
OF THE $\psi(3686)$ INTO PC GAMMA. PC INTO
 $J/PSI(3100)$ GAMMA. THEREFORE, C++

$\chi(3555)$

57 CH(3555).JPD+2+1 J=0
OBSERVED IN RADIATIVE DECAY OF $\psi(3100)$ INTO
 $\chi(3555)$ GAMMA. THEREFORE C++ THE OBSERVED DECAY INTO $\chi(3100)$
AND $\psi(1300)$
AND IS EXCLUDED BY THE ANGULAR DISTRIBUTION IN THE HADRONIC
DECAYS. JP=2 EXCLUDED BY ANGULAR DISTRIBUTION IN THE $\psi(3686)$
DECAY. JP=2 PREFERRED FELOMAN 77, DREGLIA 82

Table with columns for particle name, fit type, parameters, and fit quality. Includes fit results for $\chi(3510)$ and $\chi(3555)$.

Table with columns for particle name, fit type, parameters, and fit quality. Includes fit results for $\chi(3555)$.

55 PC PARTIAL DECAY MODES

Table with columns for particle name, fit type, parameters, and fit quality. Includes fit results for $\chi(3510)$ and $\chi(3555)$.

57 CH(3555) MASS (MEV)

Table with columns for particle name, fit type, parameters, and fit quality. Includes fit results for $\chi(3555)$.

Data Card Listings

For notation, see key at front of Listings.

Mesons

$\chi(3555)$, $\eta_c(3590)$, $\psi(3685)$

57 CHII35551 PARTIAL DECAY MODES

	DECAY MODES	DECAY MASSES
P1	CHII35551 INT0 P1+ P1-	1396 1390
P2	CHII35551 INT0 P1+ P1-	4000 4000
P3	CHII35551 INT0 2(P1+ P1-)	1390 1390 1390 1390
P4	CHII35551 INT0 3(P1+ P1-)	
P5	CHII35551 INT0 P1+ P1- K+ K-	1390 1390 4930 4930
P6	CHII35551 INT0 J/PSI(3100) GAMMA	3096+ 0
P7	CHII35551 INT0 2 GAMMA	0+ 0
P8	CHII35551 INT0 P1+ P1- P BAR	1390 1390 938+ 938
P9	CHII35551 INT0 RHOD P1+ P1-	769 1390 1390
P10	CHII35551 INT0 K+(8020) K0+ P1-/-	891+ 4030 1390
P11	CHII35551 INT0 P BAR	938+ 938
P12	CHII35551 INT0 J/PSI(3100) P1+ P1- P TO	3096+ 1390 1390 1390

57 CHII35551 BRANCHING RATIOS

R1	CHII35551 INT0 (2 GAMMA)/TOTAL (0.0001008 LESS CL=0.9) YAMADA	77 DASP	E+ E-,3 GAMMA	12/77
R2	CHII35551 INT0 2(P1+ P1-)/TOTAL 0.023 0.005	TANENBAUM 78 SMAG	PSI(3685)TO GAM CHI	12/78
R3	CHII35551 INT0 (P1+ P1- K+ K-)/TOTAL 0.020 0.005	TANENBAUM 78 SMAG	PSI(3685)TO GAM CHI	12/78
R4	CHII35551 INT0 3(P1+ P1-)/TOTAL 0.012 0.000	TANENBAUM 78 SMAG	PSI(3685)TO GAM CHI	12/78
R5	CHII35551 INT0 (P1+ P1- AND K+ K-)/TOTAL 0.0026 0.0011	TANENBAUM 78 SMAG	PSI(3685)TO GAM CHI	12/78
R6	CHII35551 INT0 (P1+ P1+ P BAR)/TOTAL 0.0035 0.0014	TANENBAUM 78 SMAG	PSI(3685)TO GAM CHI	12/78
R7	CHII35551 INT0 (J/PSI(3100) GAMMA)/TOTAL (0.30) (0.14)	BIDDICK 77 CNTR	PSI(3685)TO GAM CHI	12/77
R7 T		BARTEL 78 CNTR	PSI(3685)TO GAM CHI	4/78
R7 B		BARTEL 78 CNTR	PSI(3685)TO GAM CHI	12/78
R7 T		TANENBAUM 78 SMAG	PSI(3685)TO GAM CHI	12/78
R7 T		BRANDL 79 DASP	PSI(3685)TO GAM CHI	12/79
R7 T		MIR 80 SMAG	PSI(3685)TO GAM CHI	9/81*
R7 T	479 0.170 0.030	OREGLIA 82 CBAL	PSI(3685)TO GAM CHI	2/82*
R7 AVG	0.157 0.017	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)		

R8	CHII35551 INT0 RHOD P1+ P1-1/TOTAL 0.0071 0.0042	TANENBAUM 78 SMAG	PSI(3685)TO GAM CHI	12/78
R9	CHII35551 INT0 K+(8020) K0+ P1- -1/TOTAL 0.0050 0.0029	TANENBAUM 78 SMAG	PSI(3685)TO GAM CHI	12/78
R10	CHII35551 INT0 (P1+ P1-1)/TOTAL (UNITS 10^+3) K+ 2.0 L+1	BRANDL 79 DASP	PSI(3685)TO GAM CHI	12/79
R11	CHII35551 INT0 (K+ K-1)/TOTAL (UNITS 10^+3) L+ 1.6 L- 2	BRANDL 79 DASP	PSI(3685)TO GAM CHI	12/79
R12	CHII35551 INT0 (P BAR)/TOTAL (UNITS 10^+3) (1.0) OR LESS CL=0.90	BRANDL 79 DASP	PSI(3685)TO GAM CHI	12/79
R13	CHII35551 INT0 (J/PSI P1+ P1-)/TOTAL (0.015) OR LESS CL=0.90	BARATE 81 SPEC	100 PI-BE-ZP120	1/82*
R T	ESTIMATED USING PSI(3685) TO GAMMA CHII(35551)/TOTAL=0.074			
R T	THE ERRORS DO NOT CONTAIN THE UNCERTAINTY IN THE PSI(3685) DECAY.			

REFERENCES FOR CHII(3555)

FELDMAN 75 PRL 35 821 *JEAN-MARIE SADDOUTI,VANUCCI,+ (LBL+SLAC)
 ALSO 75 PRL 35 189 (ERRATA)
 TANENBAUM 75 PRL 35 1323 TANENBAUM,WHITAKER,ABRAMS,+ (LBL+SLAC)
 TRILLING 76 STANFORD SYMP. 437 G. H. TRILLING (LBL)
 WHITAKER 76 PRL 37 1596 *TANENBAUM,ABRAMS,ALAM,BOYARSKI,+(SLAC+LBL)
 BIDDICK 77 PRL 38 1324 *BURNETT+ (UCSD+UND+PAV+PRIN+SLAC+STAN)
 FELDMAN 77 PL 33 C 285 *PERL (LBL+SLAC)
 YAMADA 77 WASH. CONF. P. 69 YAMADA (DESY+TOKYO)
 BARTEL 78 PRL 79 B 492 DIFFMANN,DUINKER,OLSSON,DE MEIJER+DESEY+HEIDI (SMAG)
 SPIETZER 78 NYOTA SUM.JUST.47 H. SPIETZER (SMAG)
 TANENBAUM 78 PR D 17 1731 TANENBAUM,ALAM,BOYARSKI,+ (SLAC+LBL)
 ALSO 82 PRIVATE COMM. G.H.TRILLING
 BRANDL 79 PR D C 1 233 BRANDL IN.CORDS,+FACH+DESEY+HAMB+MIPIM+TOKYO
 BRANDL 79 NP D 160 426 BRANDL IN.CORDS,+FACH+DESEY+HAMB+MIPIM+TOKYO
 MIR 79 PRL 42 619 *GODDARD,ALVERSON,+IFNAL+MVRILL+DOF+TUFJ
 HIEMEL 80 PRL 44 920 *ABRAMS,ALAM,BLOCKER,+ (LBL+SLAC)
 ALSO 82 PRIVATE COMM. G.H.TRILLING (LBL+JCD)
 BARATE 81 PR D 24 2994 *ASTBURY,HCEMENT,+ (SLAC+LOIC+SMP+GERM+IND)
 ALSO 82 CERN-EP/82-15 *LEMOINE,BONAMY,+ (SLAC+LOIC+SMP+IND)
 OREGLIA 82 PR D 10 BE PUB. *BLOOD,BULOS,+ (SLAC+CIT+HARV+PRIN+STAN)
 ALSO 82 PRIVATE COMM. M.OREGLIA (IFJ)

$\eta_c(3590)$

59 ETA C(3590)JPC= 1-1-
 OBSERVED IN THE RADIATIVE DECAY OF PSI(3685) INTO
 ETA C(3590) GAMMA. THEREFORE, C-+. NEEDS CONFIRMATION.
 OMITTED FROM TABLE.

Evidence for the $\eta_c(3590)$ is based on the observation of a monochromatic gamma line in the inclusive photon spectrum for $\psi(3685)$ decays (EDWARDS 82). No exclusive decay modes are known at this time. A signal had been reported for a state at similar mass in the decay $\psi(3685) \rightarrow \Upsilon\Upsilon/\psi(3100)$ (BARTEL 78), but this signal was not confirmed in an experiment with higher statistics (OREGLIA 82).

59 ETA C(3590) MASS (MEV)
 M A 3594.0 5.0 EDWARDS 82 CBAL E+E-,GAM INCL 1/82*

M A ASSUMING MASS OF PSI(3685) = 3686 MEV.

59 ETA C(3590) WIDTH (MEV)
 W (8.0) OR LESS CL=0.95 EDWARDS 82 CBAL E+E-,GAM INCL 1/82*

59 ETA C(3590) PARTIAL DECAY MODES

P1 ETA C(3590) INTO HADRONS

59 ETA C(3590) BRANCHING RATIOS

R1 ETA C(3590) INTO HADRONS (P1)

R1 SEEN EDWARDS 82 CNTR E+E-,GAM INCL 1/82*

REFERENCES FOR ETA C(3590)
 BARTEL 78 PRL 79 B 492 *DIFFMANN,DUINKER,OLSSON,+ (DESY+HEIDI)
 PORTER 81 SLAC SUM.CONF.1355 *EDWARDS,+ (CIT+HARV+PRIN+STAN+SLAC)
 EDWARDS 82 PRL 46 70 *PRIN+DOF+PECK,+ (CIT+HARV+PRIN+STAN+SLAC)
 OREGLIA 82 PR D 10 BE PUB.1 *BLOOD,BULOS,+ (SLAC+CIT+HARV+PRIN+STAN)

$\psi(3685)$

71 PSI(3685)JPC=1--1=0

71 PSI(3685) MASS (MEV)
 WE USE INDEPENDENT MEASUREMENTS OF THE J/PSI(3100) MASS, THE PSI(3685) MASS, AND THE MASS DIFFERENCE TO PERFORM A CONSTRAINED FIT.

M S	3680.3	37.	BRIDGE 75 PLUT	E+E-	2/75
M R	(3684.)	15.1	LUTH 75 SMAG	E+E-	1/76
M S	3684.	3.	PREPOST 75 SPEC	Z1, GAMMA D	1/76
M	140(3683.0)	160.0	LEMOINE 79 GOLF	D 150 PI-BE-ZMU	12/79
M F	3686.	3.	BRANDL 79 OASP	E+ E-	12/79
M	+13 3686.0C	0.10	ZOLENITZ 80 DLYA	E+E- COLL.BEAMS	9/81*
M AVG	3686.00D	0.100	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)		
M FIT	3686.00	0.10	FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.0)		2/82*
M F	FROM A SIMULTANEOUS FIT TO E+ E- MU+ MU- AND HADRONS CHANNELS				
M F	ASSUMING GEM E-1 = GEMU MU-1				
M R	REDUPIANT WITH DATA IN MASS DIFFERENCE BELCW				
M S	ERROR OF ABOUT 1 PER CENT FROM THE UNCERTAINTY IN CALIBRATION OF THE BEAM ENERGY.				

For notation, see key at front of Listings.

Mesons

ψ(3685)

Table with 11 columns: parameter name, value, unit, and error. Row 71: PSI(3685) - J/PSI(3100) MASS DIFFERENCE (MEV). Values include 586.7, 1.0, 1.6, 2.5, 1.6, 0.99, 0.13.

Table with 11 columns: parameter name, value, unit, and error. Row 71: PSI(3685) WIDTH (KEV). Values include 226, 56, 202, 57, 202, 57, 202, 57.

Table with 11 columns: parameter name, value, unit, and error. Row 71: PSI(3685) PARTIAL DECAY MODES. Values include 0.1, 0.2, 0.4, 0.1, 0.2, 0.4, 0.1, 0.2, 0.4.

Table with 11 columns: parameter name, value, unit, and error. Row 71: PSI(3685) PARTIAL DECAY MODES (continued). Values include 139, 139, 139, 139, 139, 139, 139, 139.

Table with 11 columns: parameter name, value, unit, and error. Row 71: PSI(3685) PARTIAL WIDTHS (KEV). Values include 0, 0, 0, 0, 0, 0, 0, 0.

Table with 11 columns: parameter name, value, unit, and error. Row 71: PSI(3685) BRANCHING RATIOS. Values include 0.1, 0.2, 0.4, 0.1, 0.2, 0.4, 0.1, 0.2, 0.4.

Table with 11 columns: parameter name, value, unit, and error. Row R2: PSI(3685) INTO MU+ MU- 1/TOTAL. Values include 0.077, 75, 0.005, 1.776.

Table with 11 columns: parameter name, value, unit, and error. Row R2: PSI(3685) INTO (KARONS)/TOTAL. Values include 0.1, 0.2, 0.4, 0.1, 0.2, 0.4, 0.1, 0.2, 0.4.

Table with 11 columns: parameter name, value, unit, and error. Row R2: PSI(3685) INTO (MU+ 1/ETA+). Values include 0.1, 0.2, 0.4, 0.1, 0.2, 0.4, 0.1, 0.2, 0.4.

Table with 11 columns: parameter name, value, unit, and error. Row R2: PSI(3685) INTO (J/PSI(3100) PI0). Values include 0.1, 0.2, 0.4, 0.1, 0.2, 0.4, 0.1, 0.2, 0.4.

Table with 11 columns: parameter name, value, unit, and error. Row R2: PSI(3685) INTO (J/PSI(3100) ETAs). Values include 0.1, 0.2, 0.4, 0.1, 0.2, 0.4, 0.1, 0.2, 0.4.

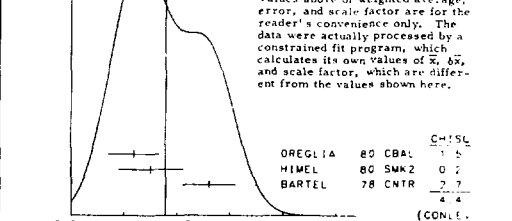


Table with 11 columns: parameter name, value, unit, and error. Row R2: PSI(3685) INTO (J/PSI(3100) PI0)/TOTAL. Values include 0.1, 0.2, 0.4, 0.1, 0.2, 0.4, 0.1, 0.2, 0.4.

Table with 11 columns: parameter name, value, unit, and error. Row R2: PSI(3685) INTO (2PI+ PI-). Values include 0.1, 0.2, 0.4, 0.1, 0.2, 0.4, 0.1, 0.2, 0.4.

Mesons

$\psi(3770)$, $\psi(4030)$, $\psi(4160)$, $\psi(4415)$

Data Card Listings

For notation, see key at front of Listings.

53 PS1(3770) PARTIAL WIDTHS (KEV)
 W1 PS1(3770) INTG E+E- (G1)
 W1 R 0.37 0.09 RAPIDS 77 SMAG 0 E+E- 12/77
 W1 0.18 0.06 RACIND 78 DLCO 0 E+E- 4/78
 W1 0.276 0.050 SCHINDLER 80 SMAG E+E- 1/62*
 W1 R SEE ALSO R2 BELOW
 W1 AVG 0.257 0.046 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.3)

53 PS1(3770) BRANCHING RATIOS
 R1 PS1(3770) INTG (D DBAR)/TOTAL (P2)
 DOMINANT PERUZZI 77 SMAG E+E- D DBAR 12/77
 R2 PS1(3770) INTG (E+ E-)/TOTAL (UNITS 10⁻⁵) (P1)
 RAPIDS 77 SMAG 0 E+E- 12/77

REFERENCES FOR PS1(3770)

PERUZZI 77 PRL 39 1301 *PICCOLO, FELDMAN, PERL, + L., AC, LBL, NMS + (MAMA)
 ->1015 77 PRL 39 526 *DOBBI, LUKE, PERL, + (STAN) SLAC + LBL + NMS + (MAMA)
 BALCINO 78 PRL 40 671 *BAUMGARTEN, BIRKWOOD, + (SLAC + STAN) UCLA + UC1
 SCHINDLER 80 PR D 21 2716 SCHINDLER, STEIGRIST, ALAM, BOYARSKI + (SLAC + LBL)

$\psi(4030)$

72 PS1(4030, JPC=1-1) 1-
 SEEN CLEARLY SEPARATED FROM THE PS1(4160)
 BY DASP AND CONFIRMED WITH LESS STATISTICS BY PLUTO
 SEEN ALSO BY MARK 1, DELCO AND THE CRYSTAL BALL
 (KIRKBY 79).

72 PS1(4030) MASS (MEV)
 M 4028.0 2.5 GOLDBAER 77 SMAG E+E- 12/77
 M 4040.0 10.0 BRANDELIK 78 DASP E+E- 4/78
 M 4028.7 2.6 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.2)

72 PS1(4030) WIDTH (MEV)
 M 52.0 10.0 BRANDELIK 78 DASP E+E- 4/78

72 PS1(4030) PARTIAL DECAY MODES
 P1 PS1(4030) INTG D DBAR DECAP MASSES
 1809+1809
 P2 PS1(4030) INTG D DBAR AND D BAR D 2007+1804
 P3 PS1(4030) INTG D DBAR 2007+2007
 P4 PS1(4030) INTG J/PSI(3100) HADRONS
 P5 PS1(4030) INTG E+ E- -5. 105
 P6 PS1(4030) INTG MU+ MU- 105+ 105

72 PS1(4030) PARTIAL WIDTHS (KEV)
 W1 PS1(4030) INTG E+E- (G5)
 0.75 0.15 BRANDELIK 78 DASP E+E- 12/78

72 PS1(4030) BRANCHING RATIOS
 R1 PS1(4030) INTG (D DBAR)/(D DBAR + D BAR D) (P1)/(P2)
 D CO GOLDBAER 77 SMAG 0 E+E- 12/77
 R2 PS1(4030) INTG J/PSI(3100) HADRONS (P4)
 LOOKED FOR BURMESTER 77 PLUT E+E- 4/77
 R3 PS1(4030) INTG (D DBAR)/(D DBAR + D BAR D) (P3)/(P2)
 12.0 GOLDBAER 77 SMAG 0 E+E- 12/77
 R4 PS1(4030) INTG (E+ E-)/TOTAL (UNITS 10⁻⁵) (P5)
 FELDMAN 77 SMAG E+E- 12/77

REFERENCES FOR PS1(4030)

AUGUSTY 75 PRL 34 764 *BOYARSKI, ABRAMS, BRIGGS + (SLAC + LBL)
 SACI 75 PL 588 481 *BIDOLI, PEMS, STELL + (ORNL + FRAS)
 BOYARSKI 75 PRL 34 762 *BREIDENBACH, ABRAMS, BRIGGS, + (SLAC + LBL)
 EPSD110 75 PL 588 478 *FELICE, PERUZZI, + (FRAS + NAPL + PADU + ROMA)
 PERUZZI 76 PRL 37 565 *PICCOLO, FELDMAN, NGUYEN, WESS, + (SLAC + LBL)
 BURMESTE 77 PL 66 B 395 *CRIEGEE, DEHME + (CESY + AMB + SIEG + MUPP)
 GOLDBAER 77 PL 69 D 505 *GOLDBAER, WESS, ABRAMS, ALAM + LUTN, + (LBL + SLAC)
 FELDMAN 77 PL 33 C 26 *PERL (LBL + SLAC)
 LUTH 77 PL 70 B 120 *PIERRE, ABRAMS, ALAM, BOYARSKI, + (LBL + SLAC)
 BRANDELIK 78 PL 76 B 361 BRANDELIK, CORDS + (AACH + DESY + AMB + MPIM + TOKY)
 ALSO 79 ZPHY C 1 233 BRANDELIK, CORDS, + (AACH + DESY + AMB + MPIM + TOKY)
 KIRKBY 79 FERMILAB SYMP. 107 J. KIRKBY RAPPORTEUR (SLAC)

$\psi(4160)$ 25 PS1(4160, JPC=1-1) 1-
 SEEN CLEARLY SEPARATED FROM THE PS1(4030)
 BY DASP AND CONFIRMED WITH LESS STATISTICS BY PLUTO.
 MARK 1, DELCO AND THE CRYSTAL BALL SEE A PROMINENT
 SHOULDER BUT NO SEPARATION (KIRKBY 79).

25 PS1(4160) MASS (MEV)
 M 4159.0 20.0 BRANDELIK 78 DASP E+E- 4/78

25 PS1(4160) WIDTH (MEV)
 M 78.0 20.0 BRANDELIK 78 DASP E+E- 4/78

25 PS1(4160) PARTIAL DECAY MODES
 P1 PS1(4160) INTG E+ E- DECAP MASSES
 -5. 105

25 PS1(4160) PARTIAL WIDTHS (KEV)
 W1 PS1(4160) INTG E+E- (G1)
 0.77 0.23 BRANDELIK 78 DASP E+E- 12/78

REFERENCES FOR PS1(4160)

BURMESTE 77 PL 66 B 395 *CRIEGEE, DEHME + (CESY + AMB + SIEG + MUPP)
 BRANDELIK 78 PL 76 B 361 BRANDELIK, CORDS + (AACH + DESY + AMB + MPIM + TOKY)
 KIRKBY 79 FERMILAB SYMP. 107 J. KIRKBY RAPPORTEUR (SLAC)

$\psi(4415)$

73 PS1(4415, JPC=1-1) 1-
 73 PS1(4415) MASS (MEV)
 M 4414. 7. SIEGRIST 76 SMAG E+E- 2/76
 M (4400.) APPROX. KNIES 77 PLUT 0 E+E- MU+ MU- 12/77
 M 4417.0 10.0 BRANDELIK 78 DASP E+E- 4/78
 M 4415.0 5.7 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)

73 PS1(4415) WIDTH (MEV)
 M 33. 10. SIEGRIST 76 SMAG E+E- 2/76
 M 66.0 15.0 BRANDELIK 78 DASP E+E- 4/78
 M 43.2 15.2 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)

73 PS1(4415) PARTIAL DECAY MODES
 P1 PS1(4415) INTG E+ E- DECAP MASSES
 -5. 105

73 PS1(4415) PARTIAL WIDTHS (KEV)
 W1 PS1(4415) INTG E+E- (G1)
 0.49 0.13 BRANDELIK 78 DASP E+E- 12/78

73 PS1(4415) BRANCHING RATIOS
 R1 PS1(4415) INTG (E+ E-)/TOTAL (UNITS 10⁻⁵) (P5)
 1.3 .3 SIEGRIST 76 SMAG E+E- 2/76
 R2 PS1(4415) INTG HADRONS/TOTAL (P4)
 R2 SIEGRIST 76 SMAG E+E- 12/77

REFERENCES FOR PS1(4415)

*ABRAMS, BOYARSKI, BREIDENBACH, + (LBL + SLAC)
 BURMESTE 77 PL 66 B 395 *CRIEGEE, DEHME + (CESY + AMB + SIEG + MUPP)
 KNIES 77 HAMBURG SYMP. 93 *KNIES HAMBURG TALK ON PLUTO COLL. S. (DESY)
 LUTH 77 PL 70 B 120 *PIERRE, ABRAMS, ALAM, BOYARSKI, + (LBL + SLAC)
 BRANDELIK 78 PL 76 B 361 BRANDELIK, CORDS + (AACH + DESY + AMB + MPIM + TOKY)

Data Card Listings
For notation, see key at front of Listings.

Mesons
T(9460), T(10020)

T(9460)

49 UPSILON(9460) JPC=1- 1 +

Table with columns for mass (MEV), fixed target experiments, and systematic errors. Includes data for 49 UPSILON(9460) MASS (MEV) and systematic error added linearly by us.

49 UPSILON(9460) WIDTH (KEV)

Table with columns for width (KEV), fixed target experiments, and systematic errors. Includes data for 49 UPSILON(9460) WIDTH (KEV) and systematic error added linearly by us.

49 UPSILON(9460) PARTIAL DECAY MODES

Table with columns for partial decay modes, intensity (INTO), and branching ratios. Includes data for 49 UPSILON(9460) PARTIAL DECAY MODES.

49 UPSILON(9460) PARTIAL WIDTHS (KEV)

Table with columns for partial widths (KEV), intensity (INTO), and branching ratios. Includes data for 49 UPSILON(9460) PARTIAL WIDTHS (KEV) and systematic error added linearly by us.

49 UPSILON(9460) BRANCHING RATIOS

Table with columns for branching ratios, intensity (INTO), and total width (TOTAL). Includes data for 49 UPSILON(9460) BRANCHING RATIOS and systematic error added linearly by us.

REFERENCES FOR UPSILON(9460)

List of references for UPSILON(9460) including authors like Finkler, Goldberger, and Lederman, and publication details.

T(10020)

52 UPSILON(10020) JPC=1- 1 +

Table with columns for mass (MEV), fixed target experiments, and systematic errors. Includes data for 52 UPSILON(10020) MASS (MEV) and systematic error added linearly by us.

52 UPSILON(10020) WIDTH (KEV)

Table with columns for width (KEV), fixed target experiments, and systematic errors. Includes data for 52 UPSILON(10020) WIDTH (KEV) and systematic error added linearly by us.

52 UPSILON(10020) PARTIAL DECAY MODES

Table with columns for partial decay modes, intensity (INTO), and branching ratios. Includes data for 52 UPSILON(10020) PARTIAL DECAY MODES.

52 UPSILON(10020) PARTIAL WIDTHS (KEV)

Table with columns for partial widths (KEV), intensity (INTO), and branching ratios. Includes data for 52 UPSILON(10020) PARTIAL WIDTHS (KEV) and systematic error added linearly by us.

52 UPSILON(10020) BRANCHING RATIOS

Table with columns for branching ratios, intensity (INTO), and total width (TOTAL). Includes data for 52 UPSILON(10020) BRANCHING RATIOS and systematic error added linearly by us.

REFERENCES FOR UPSILON(10020)

List of references for UPSILON(10020) including authors like Finkler, Goldberger, and Lederman, and publication details.

Mesons

T(10350), T(10570), K+, K0, K+(892)

Data Card Listings
For notation, see key at front of Listings.

T(10350)

46 UPSILON(10350) JPC=1- 1+
48 UPSILON(10350) MASS (GEV)
M MASS 10.347 0.010 FROM UPSILON(9460) MASS AND MASS DIFFERENCE BELOW

48 UPSILON(10350) PARTIAL DECAY MODES
P1 UPSILON(10350) INTO MU+ MU- DECAY MASSES
P2 UPSILON(10350) INTO E+ E- 105+ 105 -5+ -5

48 UPSILON(10350) PARTIAL WIDTHS (KEV)
W1 UPSILON(10350) INTO IE+ E-1/TOTAL (G2)
W1 S 0.44 0.09 ANDREWS 80 CLEO E+E-HADRONS 9/81
W1 D 0.41 0.09 BOWLINGER 80 CUSB E+E-HADRONS 9/81

REFERENCES FOR UPSILON(10350)
COBB 77 PL 72 B 273 +1MATA,FADJAN,GOLDBERG,+BNL+CERN+SYR+YALE|
HERB 77 PRL 30 252 +HBR,LIEDERMAN,APPEL,ITD+ | CCLU+FNAL+STON|
KAPLAN 78 PRL 40 435 +APPEL,HERB,HOM,LEDERMAN,+ | STON+FNAL+CCLU|

T(10570)

47 UPSILON(10570) JPC=1- 1+
47 UPSILON(10570) MASS (GEV)
M MASS 10.569 0.010 FROM UPSILON(9460) MASS AND MASS DIFFERENCE BELOW
47 UPSILON(10570) WIDTH (MEV)
W D 19.8 10.5 ANDREWS 80 CLEO E+E-HADRONS 9/81
W S 12.6 6.0 FINDERMIA 80 CUSB E+E-HADRONS 9/81

47 UPSILON(10570) PARTIAL DECAY MODES
P1 UPSILON(10570) INTO MU+ MU- DECAY MASSES
P2 UPSILON(10570) INTO E+ E- 105+ 105 -5+ -5

47 UPSILON(10570) PARTIAL WIDTHS (KEV)
W1 UPSILON(10570) INTO IE+ E-1/TOTAL (G2)
W1 S 0.24 0.08 ANDREWS 80 CLEO E+E-HADRONS 9/81
W1 D 0.32 0.09 FINDERMIA 80 CUSB E+E-HADRONS 9/81

REFERENCES FOR UPSILON(10570)
ANDREWS 80 PRL 45 238 + [CERN+HAR+ITHAL+ENR+CORN+MFG+UCR+VANVAND]
FINDERMIA 80 PRL 45 222 FINDERMIA,GRAMINI,BOWLINGER,+CCLU+FNAL+STON|

S=1, C=0 MESON STATES

K+
10 CHARGED K(498) JP=0- 1-1/2
SEE STABLE PARTICLE DATA CARD LISTINGS

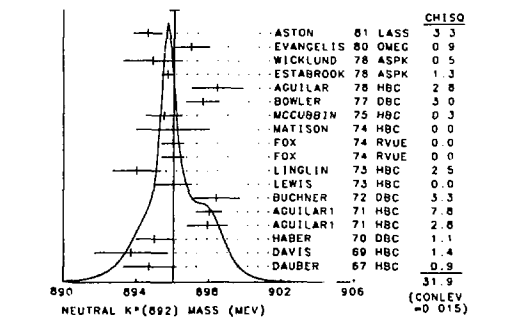
K0
11 NEUTRAL K(498) JP=0- 1-1/2
SEE STABLE PARTICLE DATA CARD LISTINGS

K+(892)
18 K(892) JP=1- 1-1/2
18 K(892) MASS (MEV)

Table listing meson masses and decay widths for K+(892). Columns include particle name, mass (MeV), width (keV), and reference.

Table listing meson masses and decay widths for various mesons including K0, K+(892), and others. Columns include particle name, mass (MeV), width (keV), and reference.

WEIGHTED AVERAGE = 896.05 +/- 0.25
ERROR SCALED BY 1.4



Data Card Listings

For notation, see key at front of Listings.

Mesons

 $Q_1(1280), \kappa(1350)$

REFERENCES FOR D112280

ARMENYER 64 DUNN CONF 1 577
 ALSO 64 DUNN CONF 1 617
 ARMENYER 64 PL 207
 ALSO 64 PP 149 1095

ARNEFLA 64 PL 16 124

SIEN 64 PP 17 776
 ALSO 66 EPJAT COMMUNICATIONS GROUP

HASSONMEYER 64 PL 268 30
 BERLINCHEM 64 PL 18 1087
 CHENWELL 67 PP 19 44
 DE BARNA 67 MC 404 376
 ALSO PRIVATE COMMUNICATION BY B. JONGEJANS
 GOLDMADY 67 PP 19 976

HARTICH 68 PP 88 9
 DENSCHER 68 PP 20 1194

ALEXANDER 69 PP 13 503
 ANDREWS 69 PP 22 731
 ASHIER 69 PP 18 10 65
 BARBARO 69 PP 22 1207
 BETTINI 69 MC 62 A 103B
 CHEN 69 PL 298 433
 CHUNG 69 PP 182 1443
 COLLEY 69 MC 6 50 520
 FAIN 69 PP 8 304
 FRIEDMAN 69 UCRL-10860
 HERRER 69 PP 189 2023

ARMAN 70 PP 3 2433
 HARTICH 70 PP 9 201
 SCHLEIPER 70 PL 31 318
 KASBERG 70 PP 1 1 96

BARNHAM 71 PP 825 45
 DENSCHER 71 PP 9 20 1194
 GOLMAN 71 PP 3 7610
 GELFAND 71 PP 26 105

ANDERSON 72 PP 2 1 1823
 BINGHAM 72 PP 8 45 157
 BRANDENBURG 72 PP 9 45 157
 BRENDEN 72 PP 28 932
 CHENWELL 72 PP 8 45 157
 JAMES 72 PP 2 2568
 FRIEDMAN 72 PP 8 5 505
 KESSELHORN 72 PP 8 45 157
 FRATI 72 PP 6 2361
 HARTICH 72 PP 8 45 157

BARLUGH 73 PP 8 49 74
 BINGHAM 73 PP 8 45 157
 DE JONGE 73 PP 8 45 157
 JONES 73 PP 8 45 157
 LEVINE 73 PP 8 45 157
 MAYER 73 PP 2 1 275

ANGELIDOU 74 MC 278 45
 HERRER 74 PP 874 693
 KROUSSIN 74 PP 20 108 338
 DEUTSCHMAN 74 PL 408 338

ANEPPOV 75 PP 886 311
 MAYER 75 PP 892 227
 ZORE 75 MC 17 265
 MELVILL 75 PL 55 B 245
 DUNNODDIE 75 PP 891 189
 TITERS 75 PP 894 311
 OTTER 75 PP 853 365
 OTTER 75 PP 896 24
 TREVY 75 PP 897 29

HALDEVAN 76 PP 37 937
 BOEL 76 PP 3 14 2998
 HERRER 76 PP 3 736
 BRANDENBURG 76 PP 76 703
 OTTER 76 PP 1 106 77
 VERGEST 76 PL 62 A 471

CAPONEGGI 77 PP 6 127 500
 CAPONEGGI 77 PL 68 287

WISCH 78 PL 74 8 207
 KALLIEM 78 PL 17 265
 JINCH 78 PP 1 172 401

SASDEVANT 74 PP 3 10 246
 MAZZUCATO 74 PP 1 154 512
 PERSEFT 74 PP 1 154 265

SACCHI 80 PP 8 162 189
 DEWISS 80 PP 8 162 11
 FAIN 80 PP 2 21 42
 BRIFING 80 PP 8 153 1
 MAZZUCATO 80 PP 1 171

SHAM 81 PP 1 107 1
 SHEN 81 PP 1 103 1
 HERRER 81 PP 1 9 9

 $\kappa(1350)$ 1C KAPPA(1350, J^{PC}=0⁺) I=1/2

The isodoublet S-wave K π phase shift δ_0^1 is compatible with elastic unitarity up to the K_1^0 threshold. It grows monotonically, reaching 90°

at about 1350 MeV (MERCER 71, BINGHAM 72, FIRESTONE 71,72, MATISON 72,74, GALTIERI 73, YUTA 73, FOX 74, BAKER 75, LAUSCHER 75, BOWLER 77, ESTABROOKS 78, ASTON 81). The ambiguous "up" solution in the region of the $K^*(892)$ has been ruled out conclusively (MATISON 72,74, GALTIERI 73, BOWLER 77, ESTABROOKS 78).

The first inelastic two-body threshold is K_1^0 , which, however, is very weakly coupled to the κ , in accordance with the SU(3) prediction for the $\langle K_1^0 \kappa \rangle$ coupling.

In the energy range 1350-1450 MeV this phase shift exhibits rapid motion (BOWLER 77, ESTABROOKS 78, MARTIN 78), which has been taken as an indication for a κ resonance at 1500 MeV (ESTABROOKS 79). However, the phase-shift behavior can also be understood as a cusp effect due to the nearby K_1^0 threshold at 1455 MeV (TORNVIST 82). Above this energy the inelasticity due to K_1^0 is important. The phase shift can be fitted up to about 1500 MeV in a unitary coupled-channel analysis with proper analytic structure with one resonance, the $\kappa(1350)$, without background (TORNVIST 82).

This supports earlier interpretations (FIRESTONE 71,72, FRATI 72, ROUGE 72, CORDS 73, LAUSCHER 75, MORGAN 75, ENGELEN 78, ASTON 81) that the κ mass indeed is where δ_0^1 passes through 90°.

At still higher energies some evidence for a second scalar resonance exists (ASTON 81).

See also Appendix IIC.

19 KAPPA MASS (MEV)

M	C	1325.5 APPROX.	ESTABROOKS 78 ASPK	13 +- P	12/77
		1350.0 APPROX. <td>MARTIN 78 SPEC <td>10 +- 0.45 P</td> <td>12/79</td> </td>	MARTIN 78 SPEC <td>10 +- 0.45 P</td> <td>12/79</td>	10 +- 0.45 P	12/79
		1378.1 APPROX. <td>BING 70 RWU <td>0</td> <td>1/82</td> </td>	BING 70 RWU <td>0</td> <td>1/82</td>	0	1/82
		1400.1 APPROX. <td>ASTON 81 LASS <td>0 11 +- 0.4 P</td> <td>1/82</td> </td>	ASTON 81 LASS <td>0 11 +- 0.4 P</td> <td>1/82</td>	0 11 +- 0.4 P	1/82
		1370.0 <td>TORNVIST 82 RWU <td>13 +- 0</td> <td>1/82</td> </td>	TORNVIST 82 RWU <td>13 +- 0</td> <td>1/82</td>	13 +- 0	1/82
			FROM ELASTIC K π PARTIAL WAVE ANALYSIS (SEE KAPPA MINT-DIEN I)		
			POLE EXTRAPOLATION USING FIRESTONE 72 AND MATISON 74 DATA.		

19 KAPPA WIDTH (MEV)

M	C	200-300 APPROX.	ESTABROOKS 78 ASPK	13 +- P	12/77
		1540.1 (100.1) <td>LANG 70 RWU <td>0</td> <td>1/82</td> </td>	LANG 70 RWU <td>0</td> <td>1/82</td>	0	1/82
		2250.1 APPROX. <td>ASTON 81 LASS <td>0 11 +- 0.4 P</td> <td>1/82</td> </td>	ASTON 81 LASS <td>0 11 +- 0.4 P</td> <td>1/82</td>	0 11 +- 0.4 P	1/82
		645.1 OR MORE <td>TORNVIST 82 RWU <td>13 +- 0</td> <td>1/82</td> </td>	TORNVIST 82 RWU <td>13 +- 0</td> <td>1/82</td>	13 +- 0	1/82
			FROM ELASTIC K π PARTIAL WAVE ANALYSIS (SEE KAPPA MINT-DIEN I)		
			POLE EXTRAPOLATION USING FIRESTONE 72 AND MATISON 74 DATA.		

14 KAPPA PARTIAL DECAY WIDTHS

P1	KAPPA INT0 K 01	DECAY WIDTHS
P2	KAPPA INT0 K 01A	40% 448
P3	KAPPA INT0 K 01A PRIME	40% 957

16 KAPPA BRANCHING RATIOS

M	C	KAPPA INT0 IN DEFTOTAL	ESTON 81 LASS	0 11 +- 0.4 P <th>1/82</th>	1/82
		10.051 APPROX.	TORNVIST 82 RWU	14 +- 0	1/82
		10.933 APPROX.			

Data Card Listings

For notation, see key at front of Listings.

Mesons

$K_2(1400)$, $K'(1400)$, $K'(1430)$

DAUM 81 NP 8 187 1 +HERTZBERGER+AMST+CERN+CRAC+MPI+DXF+RHEL I
COTTER 81 NP 181 1 (ARGHBER+LDF+CEM+HEBER+BEL+CEC+EM+STOHI)
KODEBACK 81 ZPHY C 9 9 +SODGEN+ARMY+TERDS.+ (CERN+DEF+MADR+STOHI)

K'(1400)

21 K PRIME(1400, JP=0) I=1/2

OBSERVED IN K PI PI PARTIAL-WAVE ANALYSIS.
NOT SEEN BY VERGEEST 79.
WAIT CONFIRMATION. OMITTED FROM TABLE.

M A (1400.) APPROX. BRANDEBU 78 ASPK + 13 K+-P,K PI P 12/77
M (1400.) APPROX. DAUM 81 CNTR - 63 K+-P,K 2PI P 1/82*

M A COUPLED MAINLY TO K EPSILON. DECAY INTO K*(892) PI SEEN.
M

21 K PRIME WIDTH (MEV)
V A (1200.) APPROX. BRANDEBU 78 ASK + 13 K+-P,K PI P 12/77
W (1200.) APPROX. DAUM 81 CNTR - 63 K+-P,K 2PI P 1/82*

M A COUPLED MAINLY TO K E-SILON. DECAY INTO K*(892) PI SEEN.
M

21 K PRIME PARTIAL DECAY MODES

P1 K PRIME INTO K*(892) PI 891+ 139
P2 K PRIME INTO K A0 493+ 769
P3 K PRIME INTO KAPPA PI 1350+ 159

21 K PRIME PARTIAL WIDTHS (MEV)

W1 K PRIME INTO K*(892) PI (109.) APPROX. DAUM 81 CNTR 63 K+-P,K 2PI P 1/82*
W2 K PRIME INTO K A0 (134.) APPROX. DAUM 81 CNTR 63 K+-P,K 2PI P 1/82*
W3 K PRIME INTO KAPPA PI (117.) APPROX. DAUM 81 CNTR 63 K+-P,K 2PI P 1/82*

REFERENCES FOR K PRIME
BRANDEBU 78 PRL 36 1239 BRANDEBU 78 CARNEGIE, CASIMIRO, DAVIER + SLACI JP
VERGEEST 79 NP 8 156 265 +JONGEJANS, DIONISI + (ENIJA+AMST+CERN+DXF)

DAUM 81 NP 8 187 1 +HERTZBERGER+AMST+CERN+CRAC+MPI+DXF+RHEL I

K'(1430)

22 K*(1430, JP=2+) I=1/2

WE CONSIDER THAT PHASE-SHIFT ANALYSES PROVIDE MORE
RELIABLE DETERMINATIONS OF THE MASS AND WIDTH.
SEE RHD(77)01 NTMI-REVIEW.

22 K*(1430) MASS (MEV)

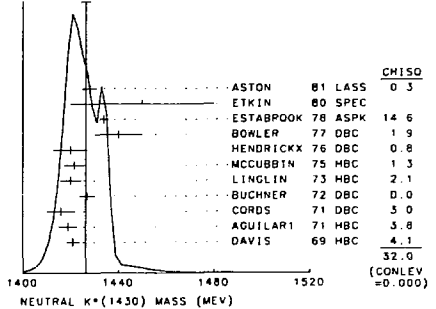
M CHARGED ONLY, WITH FINAL STATE K PI
D 39 1423. 11.0 BASSANO 67 HBC - 4.6, 5.0, K+-P, KPI - 12/75
D 63 1427.0 12.0 SCHWEINGR 66 HBC - 5.5 K+-P (K PI) 12/77
D 220 1416.0 4.0 CERMELLI 68 DBC - 3.9, 4.4, K+-P 1/71
D 60 1414. 13.0 LIND 69 HBC + 9.4 K+-P(KO) 12/77
M 1400 1420.0 3.1 AGUILARI 71 HBC - 3.9, 4.6, K+-P 1/71
M D 225 1425. 10.0 BARNHAM 71 HBC + K+-P(KO) P 1/75
M B 1428.0 4.6 MARTIN 78 SPEC - 10 K+-P, K5 PI P H
M B 1423.8 4.6 MARTIN 78 SPEC - 10 K+-P, K5 PI P H
D 5781 1448.5 5.0 TOLSON R SPEC + K+-P, K5 PI P 1/82*
D 2921 1431.8 15.6 DELFOSSE 81 SPEC - K+-P, K+-PI P 1/82*
M 995 1423.0 1.0 TOAFF 81 HBC - 8.5 K+-P(KO) PI - P 1/82*
M AVG 1422.0 1.9 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)

M NEUTRAL ONLY
W 2200 1421. 2.6 DAVIS 69 HBC 0 12, K+-P (K PI) 9/69
W 1800 1419. 4.7 AGUILARI 71 HBC 0 3.9, 4.6, K+-P 1/71
W 600 1416. 6.0 CORDS 71 DBC 0 9.4, K+-P, K+-PI - P 2/72
W 1100 1427. 3.1 BUCHNER 72 DBC 0 4.0, K+-P, K+-PI - P 1/72
M C 1420.1 4.3 LINGLIN 73 HBC 0 2.13 K+-P, K+-PI - 12/75
M 800 1421.6 4.2 MCCUBBIN 75 HBC 0 1.6, K+-P, K+-PI - P 12/75
M E 1423.0 13.0 BUCHNER 77 DBC 0 2.13 K+-P, K+-PI - P 1/74
M 300 1420.0 5.0 HENDRIX 76 DBC 8.25 K+-P, K+-PI 7/77
M P 1440.0 10.0 BOWLER 77 DBC 0 5.5, K+-P, K+-PI - P 12/77
M P 1436.0 2.0 ESTABROOK 78 ASPK 0 13K+-P, K+-PI - P 3/82*
M 1450. 3.0 ETKIN 80 SPEC 0 6, K+-P, K+-PI - I-N 3/82*
M 1428. 3.1 ASTON 81 LASS 0 11 K+-P, K+-PI - N 1/82*

M AVG 1426.4 2.0 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)
15SEE IDEOGRAM BELOW I

M B SYSTEMATIC ERROR ADDED BY US.
M C FROM POLE EXTRAPOLATION, USING WORLD K*0 DST
M D ERKIDS ENLARGED BY US TO GAMMA/SORTINI. SEE TYPED NOTE ON K*0
M E SEE MORE RECENT PARTIAL WAVE ANALYSIS (ETKIN 80)
M F FROM PHASE SHIFT ANALYSIS.
M W NUMBER OF EVENTS IN PEAK REEVALUATED BY US

WEIGHTED AVERAGE = 1426.4 ± 2.0
ERROR SCALED BY 1.9



22 K*(1430) WIDTH (MEV)

W CHARGED ONLY, WITH FINAL STATE K PI
W 1400 96.7 15.1 12.5 AGUILARI 71 HBC - 3.9, 4.6, K+-P 11/71
W 96.5 3.8 MARTIN 78 SPEC + 10 K+-P, K5 PI P 12/78
W 97.7 4.0 MARTIN 78 SPEC - 10 K+-P, K5 PI P 12/78
W 0 519 119.9 20.0 DELFOSSE 81 SPEC - K+-P, K+-PI - P 1/82*
W 0 292 96.0 27.5 DELFOSSE 81 SPEC - K+-P, K+-PI - P 1/82*
W 935 65.0 11.0 TOAFF 81 HBC - 6.5 K+-P, K+-PI - P 1/82*
W AVG 97.0 2.6 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)

M NEUTRAL ONLY
W 2200 101. 10.0 DAVIS 69 HBC 0 12, K+-P (K PI) 9/69
W 1800 116.6 10.3 15.5 AGUILARI 71 HBC 0 3.9, 4.6, K+-P 1/71
W 0 1100 109. 4.0 BUCHNER 72 DBC 0 4.0, K+-P, K+-PI - P 1/72
M C 161.0 (14.0) TOLSON R SPEC + K+-P, K5 PI P 1/82*
W 800 116.0 18.0 MCCUBBIN 75 HBC 0 1.6, K+-P, K+-PI - P 12/75
W 1170.0 (120.0) BOWLER 77 DBC 0 5.5, K+-P, K+-PI - P 12/77
M P 96.0 5.0 LINGLIN 73 HBC 0 2.13 K+-P, K+-PI - P 1/74
M P 140. 30.0 ETKIN 80 SPEC 0 6, K+-P, K+-PI - I-N 3/82*
W 98. 8.0 ASTON 81 LASS 0 11 K+-P, K+-PI - N 1/82*

M AVG 101.7 3.5 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)

M C FROM POLE EXTRAPOLATION, USING WORLD K*0 DST
M D ERKIDS ENLARGED BY US TO GAMMA/SORTINI. SEE K*0 TYPED NOTE.
M F FROM PHASE SHIFT ANALYSIS.
M W NUMBER OF EVENTS IN PEAK REEVALUATED BY US

22 K*(1430) PARTIAL DECAY MODES

P1 K*(1430) INTO K PI 891+ 139
P2 K*(1430) INTO K A0 493+ 769
P3 K*(1430) INTO K A0 493+ 769
P4 K*(1430) INTO K OMEGA 493+ 769
P5 K*(1430) INTO K EFA 493+ 769
P6 K*(1430) INTO K*(892) PI 891+ 139
P7 K*(1430) INTO K OMEGA PI 493+ 769+ 139

FITTED PARTIAL DECAY MODE BRANCHING FRACTIONS

The matrix below is derived from the error matrix for the fitted partial decay mode branching fractions, P, as follows: The diagonal elements are P_i ± 6P_i, where P_i = √(BP_iBP_i), while the off-diagonal elements are the normalized correlation coefficients (BP_iBP_j)/(BP_iBP_j). For the definitions of the individual P_i, see the listings above; only those P_i appearing in the matrix are assumed in the fit to be nonzero and are thus constrained to add to 1.

	P 1	P 2	P 3	P 4	P 5	P 6	
P 1	-4.479	-0.226					
P 2	-0.842	-2.458	-0.701				
P 3	-0.109	-4.180	-0.879	-0.100			
P 4	-1.587	-2.289	-1.375	-0.415	-0.152		
P 5	-4.441	-4.415	-3.130	-0.918	-0.470	-0.266	
P 6	-39.13	-4.761	-3.126	-0.987	-1.155	-1.294	-0.255

22 K*(1430) BRANCHING RATIOS

R1 K*(1430) INTO K PI/TOTAL (PI)
R1 P 0.49 0.32 ESTABROOK 78 ASPK + 13K+-P, K+-PI 12/77
R1 P 0.43 0.01 ASTON 81 LASS 0 11 K+-P, K+-PI - N 1/82*
R1 P FROM PHASE SHIFT ANALYSIS.
R1 - - - - -
R1 0.442 0.026 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 2.7)
R1 FIT 0.429 0.029 FROM FIT (ERROR INCLUDES SCALE FACTOR OF 2.7)
R2 K*(1430) INTO K*(892) PI/K PI + K PI P1 (P2/P1+P3+P4+P5+P6)
R2 OR 10.475 10.133 BADIER 65 HBC - 3.0 0.0 1/78
R2 0 10.475 10.101 BASSANO 67 HBC - 4.0, 4.6, 5.0 K+-P 10/67

Mesons

$K^*(1780)$, $K^*(2060)$

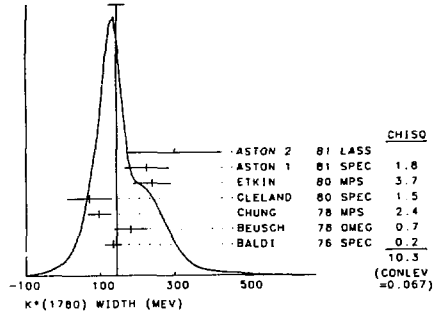
Data Card Listings

For notation, see key at front of Listings.

wave of the $K\pi\pi$ system produced in the charge-exchange process $K^-\pi \rightarrow K^0\pi^+\pi^-$ also shows resonance-like behavior at ~1800 MeV (BEUSCH 78, ETKIN 80). Since the mass values quoted for the $K\pi$ and $K\pi\pi$ modes are not significantly different, it seems natural to consider them as alternative decay modes of a single resonance.

There appears to be some disagreement in the values of the width obtained using the $K\pi$ channel. The measured values tend to become larger when the number of angular moments included in the fit increases. For the time being the observed discrepancies seem to originate from the explicit parameterization of the experimental distributions rather than from the data themselves.

WEIGHTED AVERAGE = 144.2 ± 21.1
ERROR SCALED BY 1.4



60 $K^*(1780)$ MASS (MEV)

EXPT	MODE	VAL	ERR	UNIT	DATE
M	M	1779.0	11.0	BALDI	76 SPEC + 10 $K^-\pi^0\pi^+\pi^0$ 12/77
M	A	1776.0	28.0	BRANDENB	76 ASPK 013 $K^-\pi^+\pi^-\pi^0$ 12/75
M	D	1812.0	28.0	BEUSCH	78 OMEG 10K- $\pi^0\pi^+\pi^0$ 4/78
M	D	1786.0	8.0	CHUNG	78 MPS 0 $K^-\pi^+\pi^-\pi^0$ 8 DEV 1/78
M	D	1754.0	9.0	CLELAND	80 SPEC +- 50 $K^-\pi^+\pi^0$ 2/81
M	D	1550.0	50.0	ETKIN	80 MPS 0 & $K^-\pi^0\pi^+$ 1/82
M	J	1786.0	15.0	ASTON 1	81 SPEC 011 $K^-\pi^+\pi^0$ 2/81
M	K	1755.0	25.0	ASTON 2	81 LASS 0 11 $K^-\pi^+\pi^0$ 1/82
M	K	190 1762.0	9.0	TOAFF	81 HBC - 6.5 $K^-\pi^+\pi^0$ 1/82
M	AVG	1772.7	5.5	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.3) (SEE IDEOGRAM BELOW)	

M A CONFIRMED BY PHASE SHIFT ANALYSIS OF ESTABROOKS 77, YIELDS JP=3-
M J FROM A FIT TO γ 0 MOMENT.
M K FROM ENERGY INDEPENDENT PWA.
M M FROM A FIT TO γ 0 MOMENT, JP=3- FOUND.

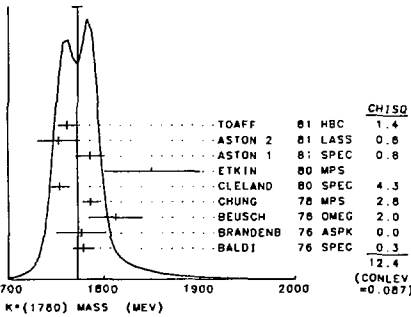
60 $K^*(1780)$ PARTIAL DECAY MODES

MODE	DECAY	MASS
P1	$K^*(1780)$ INTO $K^+\pi$	4934 139
P2	$K^*(1780)$ INTO $K^0(1920)\pi$	891 139
P3	$K^*(1780)$ INTO $K^0\pi$	4934 769
P4	$K^*(1780)$ INTO $K^0(1430)\pi$	1434 139
P5	$K^*(1780)$ INTO $K^0\pi^+\pi^-$	4934 139 139
P6	$K^*(1780)$ INTO $K^0(1920)\pi$	12754 769

60 $K^*(1780)$ BRANCHING RATIOS

MODE	BR	ERR	DATE	EXPT	MODE	BR	ERR
R4	0.19	0.02	ESTABROO	78 ASPK	D 13	$K^-\pi^+\pi^0$	12/77
R4	0.16	0.01	ASTON 2	81 LASS	0 11	$K^-\pi^+\pi^0$	1/82
R4	0.166	0.012	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.3)				

WEIGHTED AVERAGE = 1772.7 ± 5.5
ERROR SCALED BY 1.3



REFERENCES FOR $K^*(1780)$

CARMON	71 PRL 27 1160	*CORDS, CLOPP, ERIN, MEIERE, + (Purdue)
FIRESTONE	71 PL 36 B 513	*GOLDHABER, LISSAUER, TRILLING (LBL)
AGUILAR	73 PRL 30 672	*CHUNG, EISEN, PROTOPOESCU, SANTOS, + (BNL)
WALUCH	73 PR D 8 2837	*FLATY, FRIEDMAN (LBL)
BALDI	76 PL 63 B 344	*BOEHRINGER, DRSZ, MINGERDUEHLER, + (Geneva) JP
BRANDENB	76 PL 80 B 478	*BRANDENBURG, CARNegie, CASIMIRO, DAVIER, + (SLAC) JP
SPIND	76 PL 40 B 379	*BURLUTAU, PALER, CHLUDER, + (Saclay) (HEP) JP
BNL	77 PR D 126 31	*DELTON, DRAKE, WILLIAMS (BNL) JP
CARMON	77 PR D 16 1251	*CLOPP, LANDER, MEIERE, YEN, + (Purdue) JP
GRASSLER	77 PR D 125 189	*KLUEDER, + (Frankfurt) (BNL) (CERN) (Louvain) JP
BEUSCH	78 PL 74 B 282	*BERMAN, KONES, OSTER, + (IERN) (AAHE) JP
CHUNG	78 PL 40 355	*ETKIN, FLATY, INO, + (BNL) (BRAN) (CERN) (MSA) (PENN) JP
ESTABROO	78 PR D 153 640	*ESTABROOKS, CARNegie, + (MONT) (CARL) (DURHAM) (SLAC) JP
ALSO	78 PR D 17 658	*ESTABROOKS, CARNegie, + (MONT) (CARL) (DURHAM) (SLAC) JP
CLELAND	80 PL 97B 465	*DORSAZ, MARTIN, REF., + (IPITT) (GEVAL) (AUS) (DUR) JP
EINGELN	80 PR D 167 61	*JONGEANS, OUDINS, + (MIL) (JAN) (CERN) (ORF) JP
ETKIN	80 PR D 22 42	*FOLY, LINDENBAUM, KRAPER, + (BNL) (CERN) JP
ASTON 1	81 PL 99 B 502	*DUNNODDIE, DURHAM, FEGUTH, + (SLAC) (CARL) (DUR) JP
ASTON 2	81 PL 106 B 235	*CARNegie, DUNNODDIE, DURHAM, ECKLUND, (DUR) (TAS) JP
TOAFF	81 PR D 23 1500	*MUSGRAVE, ANNA, OAVIS, ECKLUND, + (BNL) (KANS) JP

$K^*(2060)$
35 $K^*(2060)$ JP=3- 1/2
ORITED FROM TABLE.

60 $K^*(1780)$ WIDTH (MEV)

EXPT	MODE	VAL	ERR	UNIT	DATE
M	M	135.0	22.0	BALDI	76 SPEC + 10 $K^-\pi^0\pi^+\pi^0$ 12/77
M	E	1270.1	170.1	BRANDENB	76 ASPK 013 $K^-\pi^+\pi^-\pi^0$ 12/75
M	D	181.0	44.0	BEUSCH	78 OMEG 10K- $\pi^0\pi^+\pi^0$ 4/78
M	D	96.0	31.0	CHUNG	78 MPS 0 $K^-\pi^+\pi^-\pi^0$ 8 DEV 1/78
M	D	70.0	60.0	CLELAND	80 SPEC +- 50 $K^-\pi^+\pi^0$ 2/81
M	D	240.0	50.0	ETKIN	80 MPS 0 & $K^-\pi^0\pi^+$ 1/82
M	J	225.0	60.0	ASTON 1	81 SPEC 011 $K^-\pi^+\pi^0$ 2/81
M	K	300.0	170.0	ASTON 2	81 LASS 0 11 $K^-\pi^+\pi^0$ 1/82
M	K	190 180.1	APPRX.	TOAFF	81 HBC - 6.5 $K^-\pi^+\pi^0$ 1/82
M	AVG	144.2	21.1	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.4) (SEE IDEOGRAM BELOW)	

M D ERRORS ENLARGED BY US TO 4% (GAMMA/SORTIN). SEE K* TYPED NOTE.
M E ESTABROOKS 77 FIND THAT BRANDENBURG 76 DATA ARE CONSISTENT
M H WITH ITS NEW WIDTH-AVERAGE.
M J FROM A FIT TO γ 0 MOMENT.
M K FROM ENERGY INDEPENDENT PWA.
M M FROM A FIT TO γ 0 MOMENT, JP=3- FOUND.

35 $K^*(2060)$ MASS (MEV)

EXPT	MODE	VAL	ERR	UNIT	DATE
M	A	488 3115.	46.	CARMON	77 HBC 0 9 $K^0\pi^+\pi^0$ 12/78
M	B	2024.	20.	CLELAND	80 SPEC +- 50 $K^-\pi^+\pi^0$ 1/82
M	B	2023.	10.	CLELAND	80 SPEC +- 50 $K^-\pi^+\pi^0$ 1/82
M	C	2092.	21.	ASTON 1	81 LASS 011 $K^-\pi^+\pi^0$ 1/82
M	D	2070.	100.	40.	ASTON 2 81 LASS 011 $K^-\pi^+\pi^0$ 1/82
M	AVG	2036.7	16.2	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 2.0) (SEE IDEOGRAM BELOW)	

M A FROM A FIT TO γ 0, 1, 2 AND γ 0, 1 MOMENTS.
M B FROM A FIT TO 0 MOMENTS.
M C FROM A FIT TO γ 0, 1, 2 AND γ 0, 1 MOMENTS.
M D FROM ENERGY INDEPENDENT PWA.

Data Card Listings

For notation, see key at front of Listings. $K^*(2060)$, $K^*(2200)$, D^+ , D^0 , $D^{*+}(2010)$, $D^{*0}(2010)$

Mesons

35 $K^*(2060)$ WIDTH (MEV)
 W 300. 200. CARMONY 77 HBC 0 9 K⁺ K⁺ PIONS 12/79
 W A 130. 65. CLELAND 80 SPEC -- 50 K⁺ P₁ P 1/82
 W B 324. 95. CLELAND 80 SPEC -- 50 K⁺ P₁ P 1/82
 W C 205. 70. 55. ASTON 1 81 LASS 011 K⁺ P₁ P 1/82
 W D 240. 500. 100. ASTON 2 81 LASS 011 K⁺ P₁ P 1/82
 W E 249.2 45.5 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.51)

W A FROM A FIT TO Y(6,0), Y(6,2) AND Y(8,0) MOMENTS.
 W B FROM A FIT TO 9 MOMENTS.
 W C FROM A FIT TO Y(6,0), Y(6,2) AND Y(8,0) MOMENTS.
 W D FROM ENERGY INDEPENDENT PWA.

35 $K^*(2060)$ PARTIAL DECAY MODES
 P1 $K^*(2060)$ INTO K P1 DECAY MASSES 493+ 139

35 $K^*(2060)$ BRANCHING RATIOS
 R1 $K^*(2060)$ INTO K P1 / TOTAL (P1)
 Q1 0.07 0.01 ASTON 2 81 LASS 0 11 K⁺ P₁ P 1/82

REFERENCES FOR $K^*(2060)$
 CARMONY 71 PRL 27 1160 *CORDS, CLOPP, ERMIN, METEKE, + (PUND)UCD(IN)
 CARMONY 77 PRD 16 1251 *CLOPP, LANDER, METEKE, YEEN, + (PUND)UCD(U)
 BROBERG 80 PR D 22 1513 *MAGGERTY, ABRAMS, DEJERBEK, FINKEL, ILLIC(IN)
 CLELAND 80 PL 978 465 *KORSZAJ, MARTIN, NEF, + (PITT)GEVA(LAUS+DURN)JP
 ASTON 1 81 PL 99 B 502 *WOODDIE, DURKIN, FIEGUTH, (SLAC)CARL+OTT(J)
 ASTON 2 81 PL 106 B 235 *CARGNIE, WOODDIE, DURKIN, (SLAC)CARL+OTT(J) JP

$K^*(2200)$ 40 $K^*(2200, J^P = 1^-)$
 THIS ENTRY CONTAINS VARIOUS PEAKS IN STRANGE MESON SYSTEMS REPORTED IN THE 2100-2500 MEV REGION AS WELL AS ENHANCEMENTS SEEN IN ANTIHYPERON NUCLEON MASS SPECTRA, OMITTED FROM THE TABLE.

40 $K^*(2200)$ MASS (MEV)
 M C 201(240.) (120.) LISSAUER 70 HBC 9, K⁺ P 11/71
 M C 12200.) APPROX. SLATTERY 71 RVUE 8-13 K⁺ P 11/71
 M C 371(247.) (4.) CHLIAPIK 79 HBC + K⁺ TO LAM-BAR P 1/80
 M D 2235. 50. BAUBILLIE 81 HBC - 8, K⁺ LAM PBAR 1/82
 M Q 2246. 30. CLELAND 81 SPEC -- 50 K⁺ LAM PBAR 1/82
 M P 2320. 30. CLELAND 81 SPEC -- 50 K⁺ LAM PBAR 1/82
 M R 2490. 20. CLELAND 81 SPEC -- 50 K⁺ LAM PBAR 1/82

AVERAGE MEANINGLESS (SCALE FACTOR = 1.0)

M C COMPIATION OF (ANTI)HYPERON-NUCLEON MASS IN K⁺ P B.-13. GEV/C
 M P JP= FROM MOMENTS ANALYSIS.
 M Q JP= FROM MOMENTS ANALYSIS.
 M R JP= FROM MOMENTS ANALYSIS.

40 $K^*(2200)$ WIDTH (MEV)
 W 20 (80-1) (20.) LISSAUER 70 HBC 9, K⁺ P 11/71
 W C (200.) APPROX. SLATTERY 71 RVUE 8-13 K⁺ P 11/71
 W Q (200.) APPROX. CHLIAPIK 79 HBC + K⁺ TO LAM-BAR P 1/80
 W Q (200.) APPROX. BAUBILLIE 81 HBC - 8, K⁺ LAM PBAR 1/82
 W Q 216. 30. CLELAND 81 SPEC -- 50 K⁺ LAM PBAR 1/82
 W P (250.) APPROX. CLELAND 81 SPEC -- 50 K⁺ LAM PBAR 1/82
 W R (250.) APPROX. CLELAND 81 SPEC -- 50 K⁺ LAM PBAR 1/82

M C COMPIATION OF (ANTI)HYPERON-NUCLEON MASS IN K⁺ P B.-13. GEV/C
 M P JP= FROM MOMENTS ANALYSIS.
 M Q JP= FROM MOMENTS ANALYSIS.
 M R JP= FROM MOMENTS ANALYSIS.

REFERENCES FOR $K^*(2200)$
 ALEXANDE 68 PRL 20 755 ALEXANDER, FIRESTONE, GOLDHABER, SHEN (LRL)
 LISSAUER 70 NP B 18 491 *ALEXANDER, FIRESTONE, GOLDHABER (LBL)
 SLATTERY 71 UR-875-332(PREP) P, SLATTERY, A REVIEW OF STRANGE MESONS(ROCM)
 CHLIAPIK 79 NP B 158 253 CHLIAPIKOV, GERDUKOV (CERN)BELG+RONS

BAUBILLIE 81 NP B 183 1 BAUBILLIER, + (BRUN)CEHN+GLAS+MSU+PNP1 JP
 CLELAND 81 NP B 184 1 *NEF, MARTIN, + (PITT)GEVA(LAUS+DURN) JP

$C = \pm 1$ MESON STATES

D^+ 31 CHARGED $D^+(1869, J^P = 1^-)$
 SEE STABLE PARTICLE DATA CARD LISTINGS

D^0 32 NEUTRAL $D^0(1869, J^P = 0^-)$ 1-1/2
 SEE STABLE PARTICLE DATA CARD LISTINGS

$D^{*+}(2010)$ 62 CHARGED $D^{*+}(2010, J^P = 1^-)$ 1-1/2

62 CHARGED $D^{*+}(2010)$ MASS (MEV)
 M G (2006.) (3.) GOLDHABE 77 SWAG -- E+E- 12/77
 M P (2008.6) (11.0) PERUZZI 77 SWAG -- E+E- 12/77
 M MASS 2010.1 0.7 FROM DD MASS (TRILLING BL RVUE) AND MASS DIFFERENCE BELOW

M G FROM SIMULTANEOUS FIT TO D^{*+}, D_0, D_1 , AND D_2 , NOT INDEPENDENT OF FELDMAN 77 MASS DIFFERENCE BELOW.
 M P PERUZZI 77 MASS NOT INDEPENDENT OF FELDMAN 77 MASS DIFFERENCE BELOW AND PERUZZI 77 DD MASS VALUE.

62 $(D^{*+}) - (D^0)$ MASS DIFFERENCE (MEV)
 DM 30 (45.3) 0.5 FELDMAN 77 SWAG $D^+ - D^0$ DD P1+ 12/77
 DM 2 (45.2) 0.6 BLIETSCH 79 RENC NEUTRINO P 12/79
 DM (145.5) APPROX. AVERY 80 SPEC GAMMA A 1/82
 DM 60 (45.5) 0.3 FITCH 81 SPEC P1-A 1/82
 DM 4
 DM 4 AVG (45.41) 0.4 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)

62 $(D^{*+}) - (D^0)$ MASS DIFFERENCE (MEV)
 E4 P 2.6 1.8 PERUZZI 77 SWAG -- E+E- 12/77
 E4 P NOT INDEPENDENT OF FELDMAN 77 MASS DIFFERENCE ABOVE, PERUZZI 77
 E4 P DD MASS, AND GOLDHABER 77 DD MASS.
 E4 M
 E4 DMASS 2.9 1.3 FROM $(D^{*+}) - (D^0)$ AND $(D^{*0}) - (D^0)$ MASS DIFFERENCES

62 CHARGED $D^+(2010)$ WIDTH (MEV)
 M 18 (20.0) OR LESS PERUZZI 76 SWAG -- E+E- PSI(4030) 1/77
 M 30 (2.0) CR LESS CL, WD PERUZZI 77 SWAG $D^+ - D^0$ DD P1+ 12/77

62 CHARGED $D^+(2010)$ PARTIAL DECAY MODES
 P1 $D^+(2010)$ INTO DD P1+ 1869+ 139
 P2 $D^+(2010)$ INTO DD GAMMA 1869+ 0
 P3 $D^+(2010)$ INTO DD P10 1869+ 134

P $D^+(2010)$ MODES ARE CHARGE CONJUGATES OF ABOVE MODES

62 CHARGED $D^+(2010)$ BRANCHING RATIOS
 R1 $D^+(2010)$ INTO DD P1+ / TOTAL (P1) GOLDHABE 77 SWAG -- E+E- 12/77
 R1 G 0.6 0.15 ASSUMING THAT ISOSPIN IS CONSERVED IN THE DECAY
 R1 G
 R2 $D^+(2010)$ INTO DD GAMMA / TOTAL (P2)
 R2 0.08 0.0 KIRBY 79 RVUE E+E- 12/79
 R3 $D^+(2010)$ INTO DD P10 / TOTAL (P3)
 R3 G 0.28 0.09 KIRBY 79 RVUE E+E- 12/79

REFERENCES FOR CHARGED $D^+(2010)$
 PERUZZI 76 PRL 37 569 *PICCOLI, FELDMAN, NGUYEN, MISS, + (SLAC+LBL)
 FELDMAN 77 PRL 38 1313 *PERUZZI, PICCOLI, ABRAMS, ALAMI, (SLAC+LBL)
 PERUZZI 77 PRL 39 1201 *PICCOLI, FELDMAN, PERL, (SLAC+LBL) (MS+MNS)
 GOLDHABE 77 PRL 69 B 503 *MISS, ABRAMS, ALAMI, BOVARSKI, + (LBL+SLAC)
 BLIETSCH 79 PRL 86 B 108 BLIETSCH, + (CERN+SDM+CERN) (M+D)FI
 KIRBY 79 BATAVIA CONF. 107 J. KIRBY (SLAC)
 AVERY 80 PRL 44 1309 *MISS, (MCKLEY, ATTYA, + (ILL+FNAL+COLUM)
 FITCH 81 PRL 46 761 *GEVA, CAVALLI, MAY, + (IP) (MS+SLAC+TORONTO)
 TRILLING 81 PRD 75 57 G, M, TRILLING (LBL+UCR)

$D^{*0}(2010)$ 61 NEUTRAL $D^{*0}(2010, J^P = 1^-)$ 1-1/2
 3 CONSISTENT WITH 1. VALUE RULED OUT (NGUYEN 77).

61 NEUTRAL $D^{*0}(2010)$ MASS (MEV)
 M G (2006.) (1.5) GOLDHABE 77 SWAG E+E- 12/77
 M G FROM SIMULTANEOUS FIT TO D^{*+}, D_0, D_1 , AND D_2 .
 M MASS 2007.2 2.1 FROM DD MASS (TRILLING BL RVUE) AND MASS DIFFERENCE BELOW

Mesons

D⁰(210), F[±], F⁰(2140), B, EXOTIC MESONS

For notation, see key at front of Listings.

61 (F0) - (DD) MASS DIFFERENCE (MEV)
 DR G 142.7 1.7 GOLDHABER 77 SHAG 0 E+- 3/82
 DR G 142.2 2.0 SADRIZIN 80 TMBL 0 DM TO DO P10 3/82
 DR G FROM SIMULTANEOUS FIT TO D⁰, D⁰±, D[±], AND D[±].
 DM DM
 DM AVG 142.5 1.3 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)

61 NEUTRAL D*(2101) WIDTH (MEV)
 M 15.1 OR LESS GOLDHABER 76 SHAG E+- TO D⁰ 3/77

61 NEUTRAL D*(2101) PARTIAL DECAY MODES
 P1 D⁰(2101) INTO DO P10 DECAY MASSES 1864± 134
 P2 D⁰(2101) INTO DO GAMMA 1864± 0
 P D⁰(2101)BAR MODES ARE CHARGE CONJUGATES OF ABOVE MODES

61 NEUTRAL D*(2101) BRANCHING RATIOS
 R1 D⁰(2101) INTO (DO GAMMA)/(DO P10 + DO GAMMA) (P21)/(P1+P2) 12/77
 R1 G 0.45 0.15 GOLDHABER 77 SHAG E+-
 R1 G WE QUOTE THE NORMAL FIT VALUE FROM TABLE 1. THE ISD-SPIN
 R1 G CONSTRAINED FIT IS NOW KNOWN TO GIVE A DO GAMMA FRACTION WHICH IS 1
 R1 G TOO LARGE. SEE DETAILS IN FOOTNOTE 21 OF FELDMAN 77 REVIEW.

R2 D⁰(2101) INTO (DO P10)/TOTAL (P1) 12/70
 R2 G 0.55 0.15 KIRKBY 79 RVUE E+ -

REFERENCES FOR NEUTRAL D*(2101)

- GOLDHABER 76 PRL 37 255 GOLDHABER,PIERRE,ABRARS,ALAM,+ (LDB+SLAC)
- GOLDHABER 76 SLAC CONF. 370 G.GOLDHABER (AVAIL. AS LBL-5534) (LDB+SLAC)
- GOLDHABER 77 PL 69 B 503 GOLDHABER,ABRARS,ALAM,+ (LDB+SLAC)
- ALSO 77 BANF SUPP. INST 75 G.J.FELDMAN (SLAC)
- NGUYEN 77 PRL 39 262 CHISS,ABRARS,ALAM,BOVARSKI,+ (LDB+SLAC)
- KIRKBY 79 BATAVIA CONF. 107 J. KIRKBY (SLAC)
- SADRIZIN 80 MADISON CONF. 68 SADRIZINSKI,+ (PRINCE+HARRV+SLAC+STAN)
- TRILLING 81 PRL 75 57 G.H.TRILLING (LDB+UCB)

F[±]

24 F-(120,0)J[±] 1 ±
 SEE STABLE PARTICLE DATA CARD LISTINGS

F⁰(2140)

24 F*(2140)J⁰ 3 1 ±
 OMITTED FROM TABLE.

74 F[±] MASS (MEV)
 M 2140.0 60. BRANDELIK 77 DASP +- E+E-,PI 3 GAMMA 12/77

74 (F[±]) - (F0) MASS DIFFERENCE (MEV)
 DM 110. 46. BRANDELIK 79 DASP +- E+E-,F GAMMA 12/79

74 F[±] PARTIAL DECAY MODES
 P1 F[±] INTO F GAMMA DECAY MASSES 2021± 0

74 F[±] BRANCHING RATIOS
 R1 F[±] INTO (F GAMMA)/TOTAL (P1) 12/77
 R1 PROBABLY SEEN BRANDELIK 77 DASP E+E-

REFERENCES FOR F*(2140)

- BRANDELIK 77 PL 70 B 132 BRANDELIK,CORDS,+(AACH+DESY+AMB+MPI+TDKY)
- BRANDELIK 78 PL 76 B 361 BRANDELIK,CORDS,+(AACH+DESY+AMB+MPI+TDKY)
- BRANDELIK 79 PL 8D B 412 BRANDELIK,CORDS,+(AACH+DESY+AMB+MPI+TDKY)

B=±1 MESON STATE

B

39 BOTTOM MESON (B(5200)J[±])
 SEE STABLE PARTICLE DATA CARD LISTINGS

EXOTIC MESON STATES

EXOTICS

50 EXOTICS

THE PURPOSE OF THIS ENTRY IS TO PROVIDE A LIST OF REFERENCES FOR EXOTIC MESON SEARCHES (SEE MAIN TEXT, SEC. 3 AND TABLE 1), AS WELL AS THEORETICALLY BASED SUGGESTIONS FOR EXPERIMENTS. NOTE THAT LEPKIN 73 PROPOSES EXPERIMENTS WHICH ARE CONCLUSIVE EVEN IF NEGATIVE RESULTS ARE OBTAINED.

REFERENCES FOR EXOTICS

REPORTS ON SEARCHES

- ROSENFELD 68 PHILA.CCNP.P.455 A.H.ROSENFELD (LRL)
- DODD 69 PR 177 1491 J.DUDERMAN, PALMER, SAMIOS (BNL)
- CHO 70 PL 32 B 409 *DEBRICK,JOHNSON,MUSGRAVE,+ (AMNHMS+KANS)
- GILCUMBE 70 PL 33 B 373 G.GILCUMBE,+ (BCN+SLAC+MSTR+ND+POLL)
- LVS 70 PR 0 2 2225 J.L.VAN PIPER (FZJ)
- ROSENER 70 EXP.MESCH SPECTROSCOPY,ED. C.BALTY AND A.H.ROSENFELD,P.499
- BUMH 72 NP 0 37 421 *CLTIE,TERRILL (MCSOONSIN)
- COHEN 73 NP 8 53 1 *FERRELL,SLATTERY,WEINER (ROCHESTER)
- DURUSOV 73 PL 45 B 517 *BAUBILLIER,GEORGE,ARMENTISE,+ (LND+PAR)
- AIAM 74 PL 57B 227 *BRABSON,GALLOWAY,+ (IND+PURD+SLAC+VAND)
- COHEN 74 BOSTON D.COHEN REVIEW TALK (COLU)
- DREN 74 NP 8 171 189 *CODDERS,FIELDS,PHINES,WHITMORE,+ (AM+DFI)
- BALTY 75 PL 37B 293 *CAUTES,COHEN,KAELER,PISELLI,+ (COLU+BRNG)
- DAVIS 75 NP 896 476 *AMAR,KRDPAC,YARGER,+ (FANS+CAC+ANL)
- BRUNDIR 76 PL 64 B 107 BRUNDERS,BRUN,FLURI,+ (FREEBURG+SLAC+ETH)
- BOUCRIOT 77 NP 8 121 251 *NAVARCH,RIEVE,+ (LAL+CEBN+CEDEF+MADR+STH)
- HODGLAND 77 NP 8 126 109 *GAYLER,HYAMS,BLUM,DETL,+ (MSTR+CEBN+MPI)
- HODGLAND 77 NP 8 126 109 *GAYLER,HYAMS,BLUM,DETL,+ (MSTR+CEBN+MPI)
- MOSES 77 NP 8 129 78 F.J.MOSES (FEF)
- ALAM 78 PRL 40 1385 *BAGGETT,BAGLIN,BONAMY,+ (IND+PURD+SLAC+VAND)
- ARNSTADN 78 PL 77 B 467 *ARNSTADN,FRANKE,MEHES,BIENLIE,+ (SLAS+DESY)
- LEMOIGNE 79 BATAVIA CONF. 524 *ABOLINS,BARATE,+ (SACL+LOIC+SAMP+IND)
- KOOIJMAN 80 PRL 45 316 *ARENTON,AYRES,DIEBOLD,MAV+ (ANL+FEF)
- AGUILAR 81 ZPHY C 6 109 *ALBAJAR,SJOGREN,+ (CERN+CEDEF+MADR+STH)
- APPEL 81 NP 8 193 269 *AUGENSTEIN,BERTOLUCCI,DOMSROV,+ (SENP+CEBN)
- BIONTA 81 PRL 46 970 *CARROLL,ROBERTSON,+ (BNL+CEBN+PML+SMSS)
- EVANGELISTA 81 NP 8 178 197 *EVANGELISTA,IBARI,BONNACCI,RODRIGUEZ,+ (MILA)
- FRANKE 81 PL 107 B 301 *RUGHES,COLEY,ARNSTRONG,+ (CLAS+BRN+CEBN)
- IRVING 81 NP 8 193 1 *MADYERS,AGUILAR,+ (CERN+CEDEF+MADR+STH)

SUGGESTIONS FOR SEARCHES

- ROSENER 68 PRL 21 950, L468 J.L.ROSENER (LTEL-VIUI)
- ROSENER 70 EXP.MESON SPECTROSCOPY,ED. C.BALTY AND A.H.ROSENFELD,P.499
- LIPKIN 73 PL 43 B 307 D.FATMAN,G.GOLDHABER,Y.ZAKH (CERN)
- LIPKIN 73 PR 0 7 2262 H.J.LIPKIN (ARGONN+FNAL)
- HCLNGREN 76 PL 77 B 394 *ANNINGTON (STH+CEBN)

Data Card Listings

For notation, see key at front of Listings.

Baryons

N's and Δ 's

Note on N's and Δ 's

I. Introduction

For this edition, three N's and one Δ have been promoted to the Baryon Table, and several new resonances that are not yet established appear in the Data Card Listings. Table I.1 lists all the entries in the Listings, and gives our evaluation of the status of each, both overall and channel by channel.

The masses, widths, and branching fractions of the N and Δ resonances come mainly from a few comprehensive partial-wave analyses. There are also several other recent analyses based on more limited data sets and/or energy ranges. Production and total-cross-section experiments can be valuable in establishing the existence of high mass bumps, but at lower energies these experiments have limited statistics compared to formation experiments, and it is seldom clear which of several states having nearly the same mass is being observed.

Even when there are good scattering data, there are two main problems in obtaining reliable resonance parameters from partial-wave analyses. First, there is sometimes disagreement among analyses on the partial-wave amplitudes themselves. This obviously depends on the quality and quantity of the scattering data and on the procedures used to determine the amplitudes from the data. Secondly, even if smooth curves were available for the amplitudes, there would still be some parametrization-dependent uncertainty about the values of the usual Breit-Wigner resonance parameters. From a theoretical standpoint, the most unambiguously defined resonance parameter is the pole position and residue, and it has been found in practice that, given sufficiently precise partial-wave amplitudes, these quantities can be extracted in a stable and parametrization-independent way, in spite of the fact that an extrapolation away from the physical region is required. This point has been discussed in detail with regard to the $\Delta(1232)$ in earlier editions of this review.^{1,2} Pole parameters have now been determined for many of the N and Δ resonances, and these are included in the Data Card Listings. In most cases, we specify pole parameters by giving the real and imaginary parts

TABLE I.1. STATUS OF Δ RESONANCES

THOSE WITH AN OVERALL STATUS OF *** OR **** ARE INCLUDED IN THE MAIN BARYON TABLE. THE OTHERS AWAIT CONFIRMATION.

PARTICLE	LIJ	STATUS	STATUS AS SEEN IN --																		
			OVERALL	TOTAL*	PI	N	ETA	K	LAM	K	SEG	PI	DE	CAN	N	OTHER					
N(1930)	P11	****	****	*													EPS	N			
N(1440)	P11	****	****	*														RHO	N		
N(1520)	D13	****	****	*	*	*	*	*	*	*	*	*	*	*	*	*	*		RHO	N	
N(1535)	S11	****	****	*	*	*	*	*	*	*	*	*	*	*	*	*	*		RHO	N	
N(1540)	P13	*	*	*	*	*	*	*	*	*	*	*	*	*	*	*	*		RHO	N	
N(1650)	S11	****	****	*	*	*	*	*	*	*	*	*	*	*	*	*	*		RHO	N	
N(1675)	D15	****	****	*	*	*	*	*	*	*	*	*	*	*	*	*	*		RHO	N	
N(1680)	F15	****	****	*	*	*	*	*	*	*	*	*	*	*	*	*	*		RHO	N	
N(1700)	D13	****	****	*	*	*	*	*	*	*	*	*	*	*	*	*	*		RHO	N	
N(1710)	P11	****	****	*	*	*	*	*	*	*	*	*	*	*	*	*	*		RHO	N	
N(1720)	P13	****	****	*	*	*	*	*	*	*	*	*	*	*	*	*	*		RHO	N	
N(1990)	F17	**	**	*	*	*	*	*	*	*	*	*	*	*	*	*	*		RHO	N	
N(2000)	F15	**	**	*	*	*	*	*	*	*	*	*	*	*	*	*	*		RHO	N	
N(2080)	D13	**	**	*	*	*	*	*	*	*	*	*	*	*	*	*	*		RHO	N	
N(2100)	S11	*	*	*	*	*	*	*	*	*	*	*	*	*	*	*	*		RHO	N	
N(2100)	P11	*	*	*	*	*	*	*	*	*	*	*	*	*	*	*	*		RHO	N	
N(2190)	G17	****	****	*	*	*	*	*	*	*	*	*	*	*	*	*	*		RHO	N	
N(2200)	D15	**	**	*	*	*	*	*	*	*	*	*	*	*	*	*	*		RHO	N	
N(2220)	H19	****	****	*	*	*	*	*	*	*	*	*	*	*	*	*	*		RHO	N	
N(2250)	G19	****	****	*	*	*	*	*	*	*	*	*	*	*	*	*	*		RHO	N	
N(2400)	I11	**	**	*	*	*	*	*	*	*	*	*	*	*	*	*	*		RHO	N	
N(2700)	K13	*	*	*	*	*	*	*	*	*	*	*	*	*	*	*	*		RHO	N	
N(2800)	G19	*	*	*	*	*	*	*	*	*	*	*	*	*	*	*	*		RHO	N	
N(3030)	*	**	**	*	*	*	*	*	*	*	*	*	*	*	*	*	*		RHO	N	
N(3245)	*	*	*	*	*	*	*	*	*	*	*	*	*	*	*	*	*		RHO	N	
N(3490)	*	*	*	*	*	*	*	*	*	*	*	*	*	*	*	*	*		RHO	N	
N(3755)	*	*	*	*	*	*	*	*	*	*	*	*	*	*	*	*	*		RHO	N	
DEL(1232)P33	****	****	****	F															RHO	N	
DEL(1550)P31	**	**	**	D																RHO	N
DEL(1600)P33	****	****	****	R																RHO	N
DEL(1620)P31	****	****	****	B																RHO	N
DEL(1700)D33	****	****	****	I																RHO	N
DEL(1900)F31	****	****	****	D																RHO	N
DEL(1905)F35	****	****	****	E																RHO	N
DEL(1910)P31	****	****	****	N																RHO	N
DEL(1920)P33	****	****	****	E																RHO	N
DEL(1930)I35	****	****	****	F																RHO	N
DEL(1940)D33	*	*	*	R																RHO	N
DEL(1950)F37	****	****	****	R																RHO	N
DEL(2150)S31	*	*	*	B																RHO	N
DEL(2160)	*	*	*	I																RHO	N
DEL(2200)G37	**	**	**	D																RHO	N
DEL(2300)H39	**	**	**	E																RHO	N
DEL(2350)I35	*	*	*	E																RHO	N
DEL(2400)F37	*	*	*	N																RHO	N
DEL(2400)G39	*	*	*	F																RHO	N
DEL(2420)H31	****	****	****	G																RHO	N
DEL(2750)I313	**	**	**	R																RHO	N
DEL(2850)	**	**	**	B																RHO	N
DEL(2850)K315	**	**	**	I																RHO	N
DEL(3230)	**	**	**	D																RHO	N

**** GOOD, CLEAR, AND UNMISTAKABLE.
 *** GOOD, BUT IN NEED OF CLARIFICATION OR NOT ABSOLUTELY CERTAIN.
 ** NEEDS CONFIRMATION.
 * WEAK.
 + ATTRIBUTED TO THE STATE CLOSEST TO WHERE THE CROSS SECTION PEAKS.

of the pole position and residue. It should be noted that these real and imaginary parts tend to be highly correlated. In particular, the absolute value of the residue is often better determined than is the phase. For further discussion, see the relevant references, e.g., NOGOVA 73, SPEARMAN 74, BALL 75, LICHTENBERG 75, VASAN 76, LONGACRE 77, ZIDELL 78, CUTKOSKY 79, MIROSHNICHENKO 79, ZIDELL 80, and CUTKOSKY 80.

The following sections discuss various recent developments in experimental N and Δ spectroscopy. For a discussion of earlier results, see our 1980 edition³ and the reviews of R.L. Kelly,⁴ R. Koch,⁵ and A.J.G. Hey and R.L. Kelly.⁶

References for Section I

1. Particle Data Group, Rev. Mod. Phys. **43**, S114 (1971).

2. Particle Data Group, Phys. Lett. 39B, 103 (1972).
3. Particle Data Group, Rev. Mod. Phys. 52, S175 (1980).
4. R.L. Kelly, in Proceedings of the 14th International Conference on Baryon Resonances (Toronto, 1980), ed. N. Isgur, p. 149.
5. R. Koch, in Proceedings of the Conference on Low and Intermediate Energy Kaon-Nucleon Physics (Rome, 1980), p. 1.
6. A.J.G. Hey and R.L. Kelly, Physics Reports C (to be published).

II. Two-Body Partial-Wave Analyses and New Resonances

(by R.E. Hendrick, St. Bonaventure University)

Several new partial-wave analyses have appeared and several older analyses have been extended to include new scattering data since our 1980 edition.¹ For $\pi N \rightarrow \pi N$ reactions, we have included results of new analyses by CUTKOSKY 80, ZIDELL 80, KOCH 80, and CHEW 80. CUTKOSKY 80 includes new elastic cross-section and polarization data and extends the mass range of the CUTKOSKY 79 analysis up to 2500 MeV. This analysis reports several new baryon states and confirms several tentative states in the 2000-2400 MeV region. A brief discussion of the revised status of these N and Δ states is given below. ZIDELL 80 performs an energy-dependent partial-wave analysis from threshold to 350 MeV. Deviations from isospin invariance are reported, and new masses, widths, and pole positions of the Δ^0 and Δ^{++} are given. KOCH 80 performs an energy-independent analysis over a similar low-energy range, but finds little evidence of isospin invariance violation. CHEW 80 uses a Barrelet-zero technique to analyze $\pi^+ p$ elastic data between 1550 and 2100 MeV and reports several new S31 and P31 states, but disagrees with the parameters of the two 4-star S31 and P31 resonances established by other analyses.

Two new inelastic analyses have been reported: an energy-dependent analysis of $\pi^- p \rightarrow \Lambda K^0$ below 1900 MeV by MUSETTE 80, and an energy-dependent analysis of the three isospin-coupled reactions $\pi^- p \rightarrow \Sigma^- K^+$, $\pi^- p \rightarrow \Sigma^0 K^0$, and $\pi^+ p \rightarrow \Sigma^+ K^+$ up to 2000 MeV by LIVANOS 80. A number of resonance masses, widths, and πN to $K\Sigma$ branching ratios have been included from this analysis.

This year, four new N and Δ resonances have been promoted from 2-star to 3-star status and added to the Baryon Table: the D13 N(2080), D15 N(2200), S31 Δ(1900), and P33 Δ(1920). The D35 Δ(1930), already listed in the Baryon Table, has been promoted from 3-star to 4-star status.

A number of weak, higher-mass states have been added to the N and Δ Listings. These include a third P11 at 2100 MeV, a second D33 at 1940 MeV, a third S31 at 2150 MeV, a second D35 at 2350 MeV, and a second F37 at 2400 MeV. The G37 Δ(2200) has been separated from the Δ(2160) listing and has been given a separate listing with 2-star status. The H39 Δ(2300) has been promoted from 1-star to 2-star status. In addition, a number of nominal resonance masses have been changed from our last edition to put them in better agreement with recent results.

Figure II.1 shows the partial-wave amplitudes obtained by HOEHLER 79 and by CUTKOSKY 80.

Reference for Section II

1. Particle Data Group, Rev. Mod. Phys. 52, S1 (1980).

III. The $\pi N \rightarrow \pi\pi N$ Channel

(by R.L. Crawford, University of Glasgow)

A general $\pi N \rightarrow \pi\pi N$ event may be described by the center-of-mass energy W, three angles α , β , and γ , and two sub-energies, say $w_{\pi N}$ and $w_{\pi\pi}$. Unlike 2 + 2-body reactions, fits of $\pi N \rightarrow \pi\pi N$ distributions at single values of W cannot be parametrized in terms of a set of constants without introducing some assumptions into the analysis. All fits to $\pi N \rightarrow \pi\pi N$ use the isobar model, which notes that almost all such events lie in quasi-2-body bands in the Dalitz plot. Thus it is assumed that any purely 3-body interaction is negligible and that the reaction proceeds by the formation of quasi-2-body intermediate states.

The basic form used for the amplitudes is

$$T(\pi N \rightarrow \pi\pi N) = \sum [T_{\Delta\pi}^{JILL'}(W) \cdot BW_{\Delta}(w_{\pi N}) \cdot X_{\Delta\pi}^{JILL'} + T_{N^*\pi}^{JILL'}(W) \cdot BW_{N^*}(w_{\pi N}) \cdot X_{N^*\pi}^{JILL'} + T_{\rho N}^{JILL'}(W) \cdot BW_{\rho}(w_{\pi\pi}) \cdot X_{\rho N}^{JILL'} + T_{\epsilon N}^{JILL'}(W) \cdot BW_{\epsilon}(w_{\pi\pi}) \cdot X_{\epsilon N}^{JILL'}]$$

Data Card Listings

For notation, see key at front of Listings.

Baryons

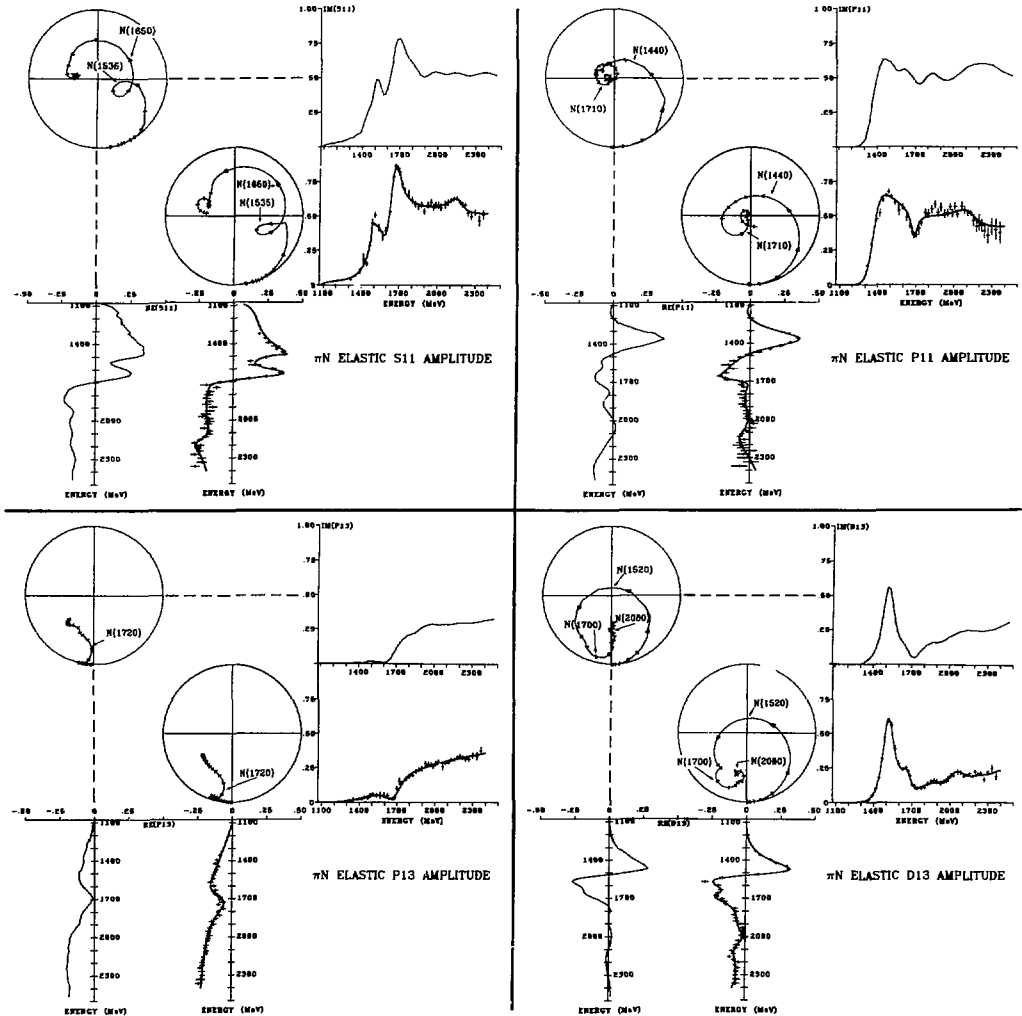
N's and Λ 's

Fig. II.1(a). The $L=2I=2J = S_{11}, P_{11}, P_{13},$ and D_{13} partial-wave amplitudes for πN elastic scattering. The upper plot for each amplitude is from HOEHLER 79 and the lower one is from CUTKOSKY 80. In the Argand plots, the ticks are at integral multiples of 50 MeV, and the established resonances are shown at their nominal positions. The real and imaginary parts of the amplitudes as functions of energy are shown projected in alignment with the Argand plots (in the projections of the CUTKOSKY 80 amplitudes, the "data points" are results of energy-independent fits, and the curves are from an energy-dependent fit to join them).

Baryons

N's and Δ 's

Data Card Listings
For notation, see key at front of Listings.

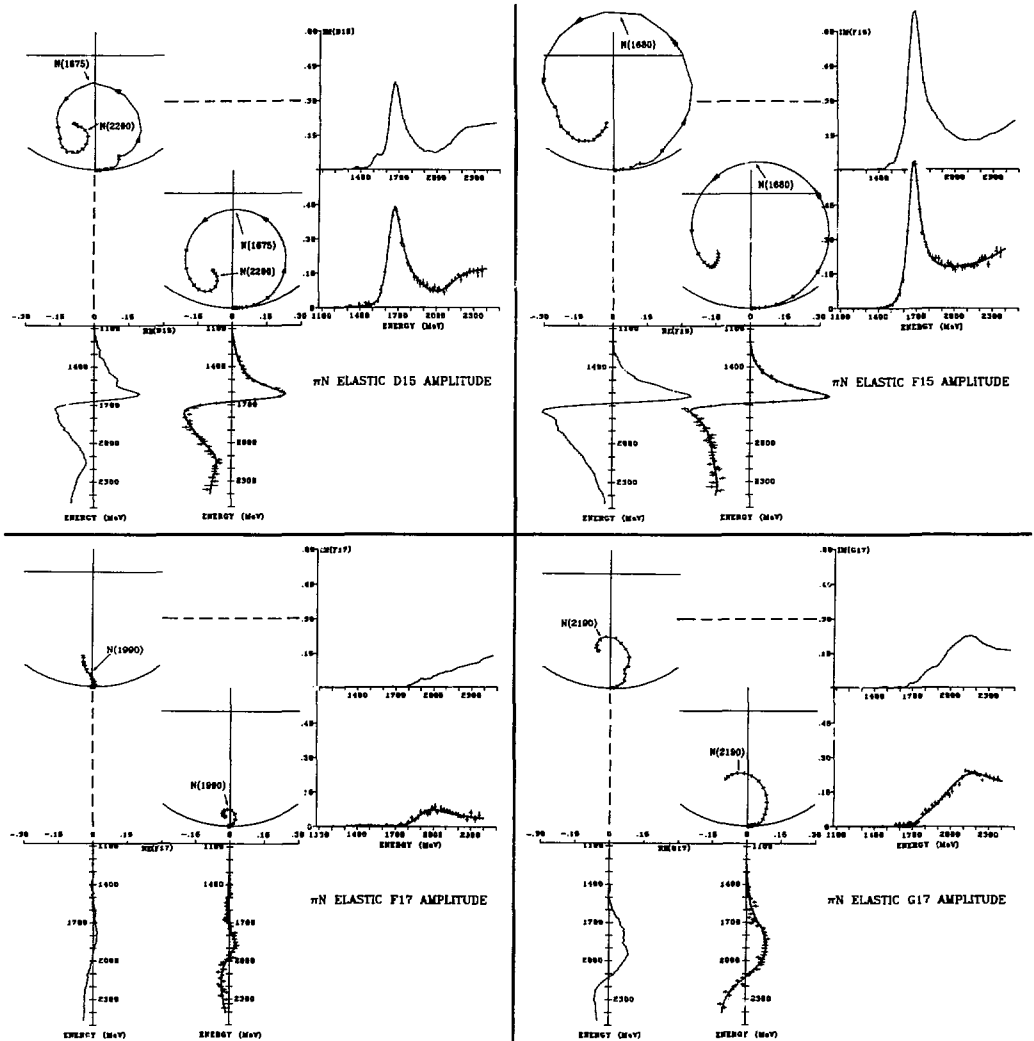


Fig. II.1(b). The L-2I-2J = D15, F15, F17, and G17 partial-wave amplitudes for πN elastic scattering. The upper plot for each amplitude is from HOEHLER 79 and the lower one is from CUTKOSKY 80. In the Argand plots, the ticks are at integral multiples of 50 MeV, and the established resonances are shown at their nominal positions. The real and imaginary parts of the amplitudes as functions of energy are shown projected in alignment with the Argand plots (in the projections of the CUTKOSKY 80 amplitudes, the "data points" are results of energy-independent fits, and the curves are from an energy-dependent fit to join them).

Data Card Listings
 For notation, see key at front of Listings.

Baryons
 N's and Δ 's

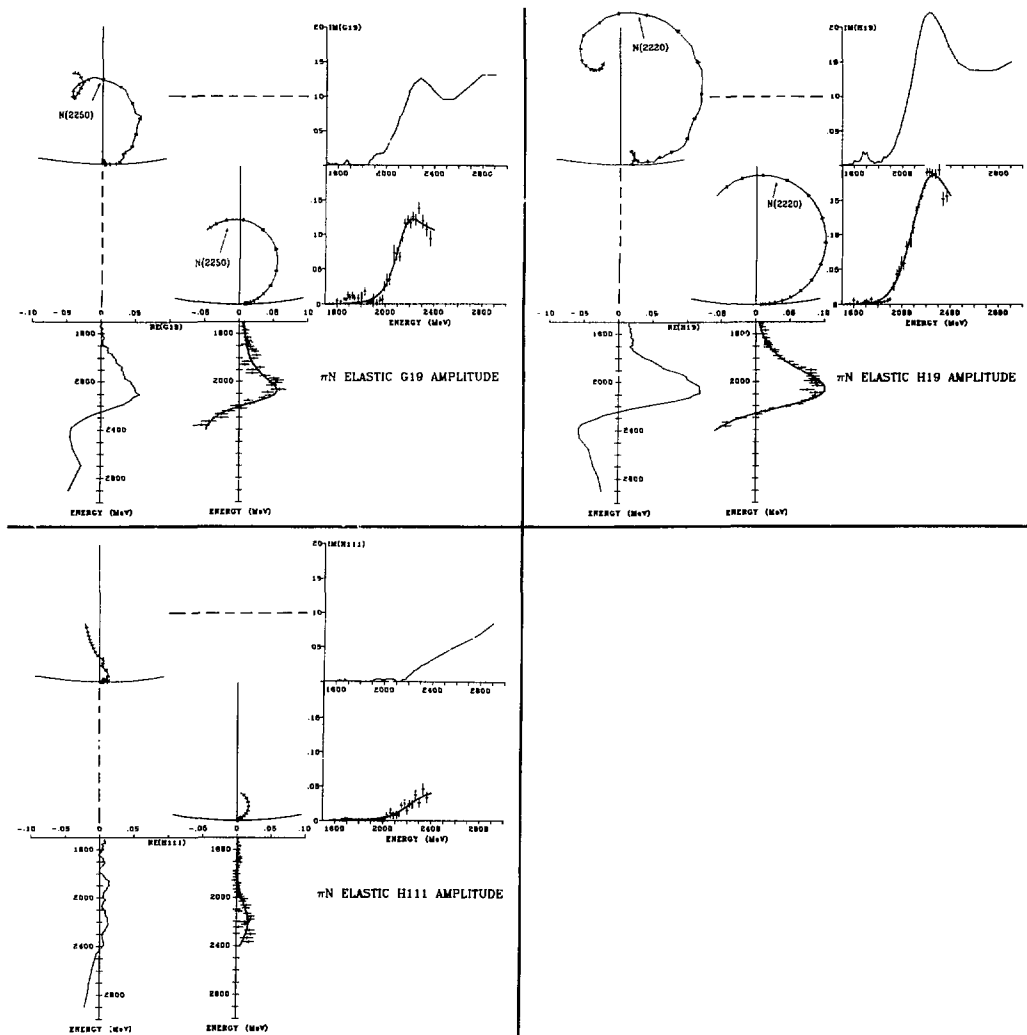


Fig. II.1(c). The $L=2I=2J = G19, H19, \text{ and } H11$ partial-wave amplitudes for πN elastic scattering. The upper plot for each amplitude is from HOEHLER 79 and the lower one is from CUTKOSKY 80. In the Argand plots, the ticks are at integral multiples of 50 MeV, and the established resonances are shown at their nominal positions. The real and imaginary parts of the amplitudes as functions of energy are shown projected in alignment with the Argand plots (in the projections of the CUTKOSKY 80 amplitudes, the "data points" are results of energy-independent fits, and the curves are from an energy-dependent fit to join them).

Baryons

N's and Δ 's

Data Card Listings
For notation, see key at front of Listings.

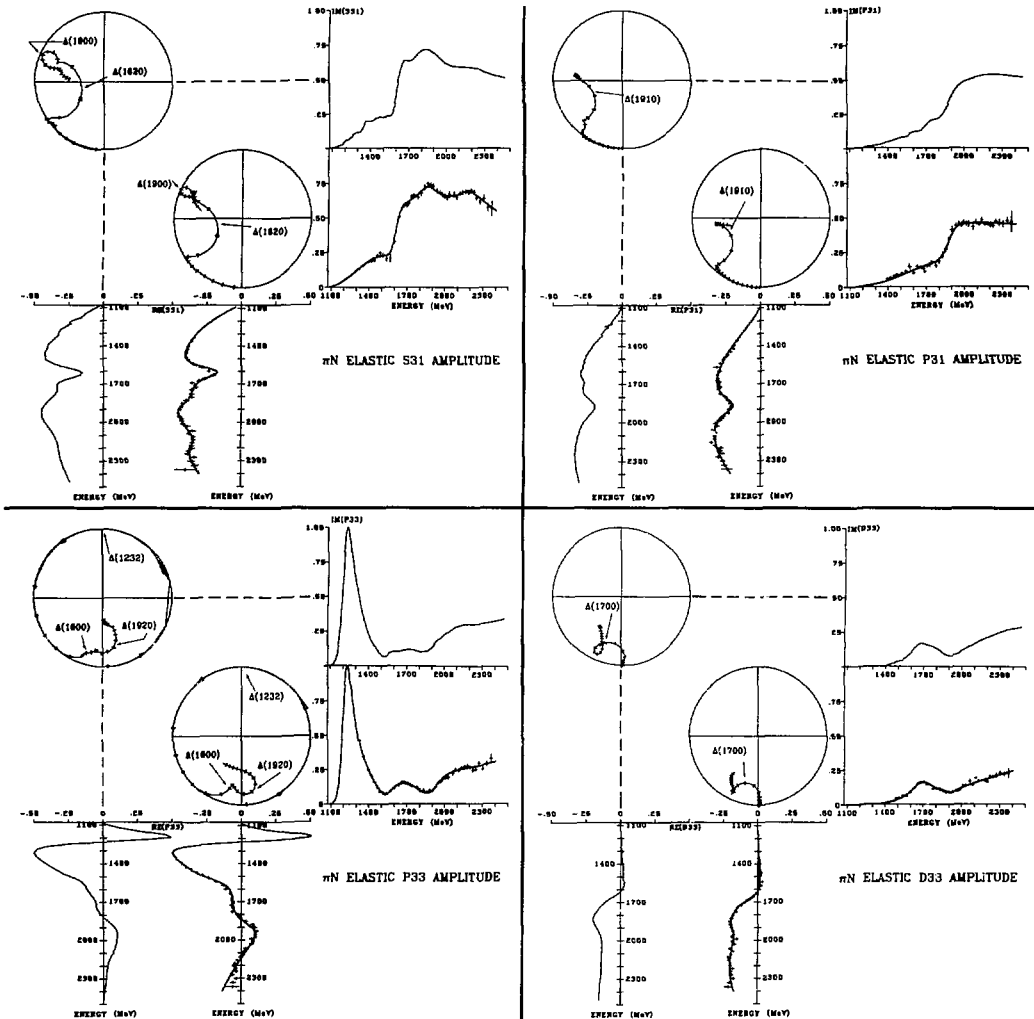


Fig. II-1(d). The L-2I-2J = S31, P31, P33, and D33 partial-wave amplitudes for nN elastic scattering. The upper plot for each amplitude is from HOEHLER 79 and the lower one is from CUTKOSKY 80. In the Argand plots, the ticks are at integral multiples of 50 MeV, and the established resonances are shown at their nominal positions. The real and imaginary parts of the amplitudes as functions of energy are shown projected in alignment with the Argand plots (in the projections of the CUTKOSKY 80 amplitudes, the "data points" are results of energy-independent fits, and the curves are from an energy-dependent fit to join them).

Data Card Listings

For notation, see key at front of Listings.

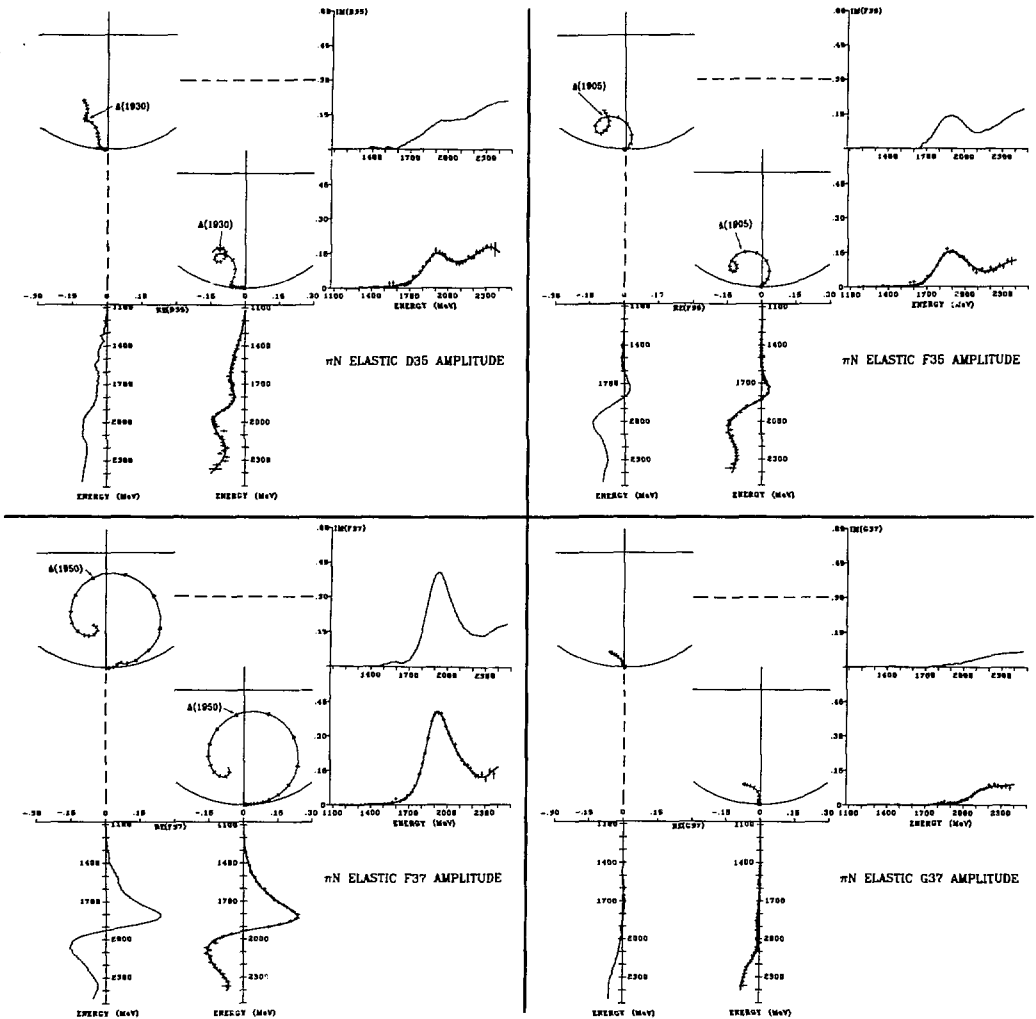
Baryons
N's and Δ 's

Fig. 11.1(e). The $L=2L-2J = D35, F35, F37,$ and $G37$ partial-wave amplitudes for πN elastic scattering. The upper plot for each amplitude is from HOEHLER 79 and the lower one is from CUTKOSKY 80. In the Argand plots, the ticks are at integral multiples of 50 MeV, and the established resonances are shown at their nominal positions. The real and imaginary parts of the amplitudes as functions of energy are shown projected in alignment with the Argand plots (in the projections of the CUTKOSKY 80 amplitudes, the "data points" are results of energy-independent fits, and the curves are from an energy-dependent fit to join them).

Baryons

N's and Δ 's

Data Card Listings

For notation, see key at front of Listings.

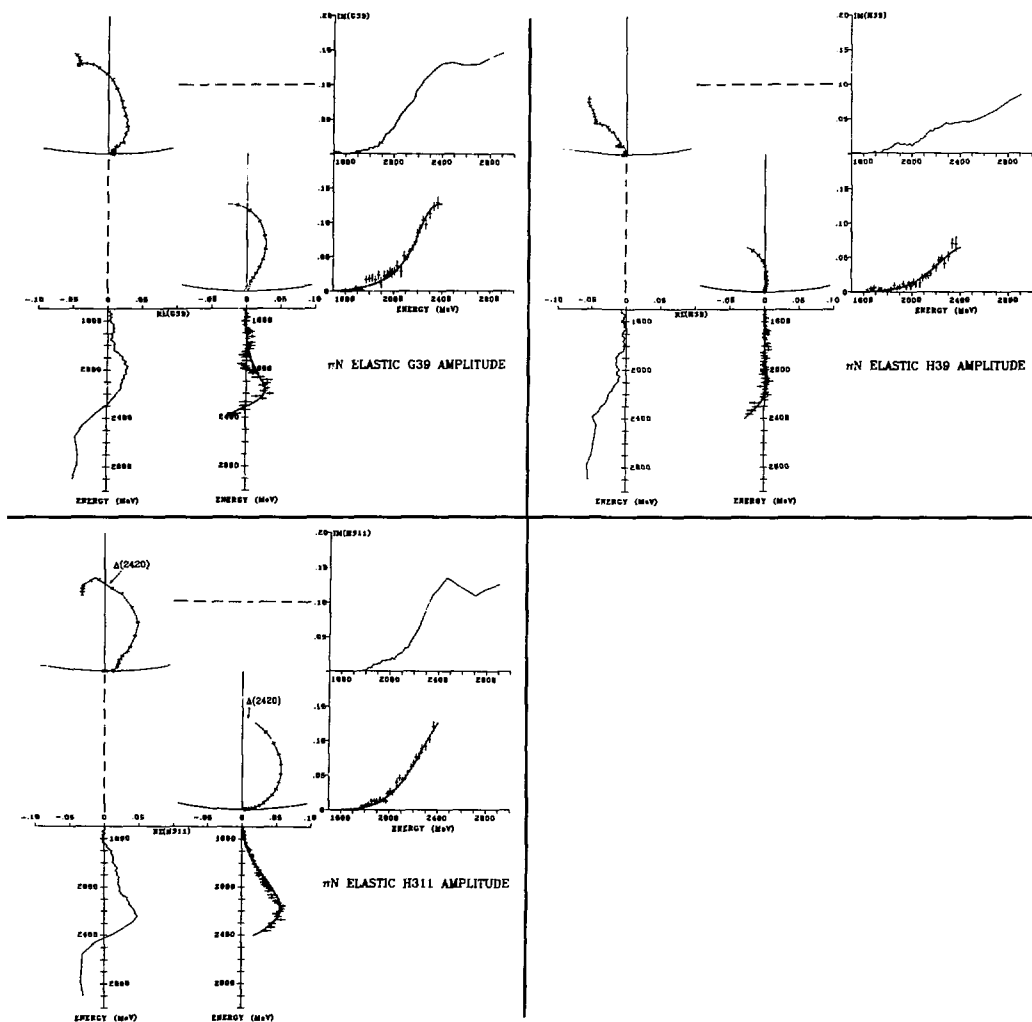


Fig. II-1(f). The L-2I-2J = G39, H39, and H31 partial-wave amplitudes for πN elastic scattering. The upper plot for each amplitude is from HOEHLER 79 and the lower one is from CUTKOSKY 80. In the Argand plots, the ticks are at integral multiples of 50 MeV, and the established resonances are shown at their nominal positions. The real and imaginary parts of the amplitudes as functions of energy are shown projected in alignment with the Argand plots (in the projections of the CUTKOSKY 80 amplitudes, the "data points" are results of energy-independent fits, and the curves are from an energy-dependent fit to join them).

Data Card Listings

For notation, see key at front of Listings.

Baryons

N's and Δ 's

where in present analyses Δ is the $\Delta(1232)$, N^* is the $N(1440)$, ρ is the $\rho(770)$, and ϵ is the S-wave, $I=0$ $\pi\pi$ enhancement (not all these isobars are included in every analysis). Here, BW denotes the appropriate Breit-Wigner or corresponding 2-body amplitude from πN or $\pi\pi$ analyses, and X is a well-defined function containing all the angular information. The decays of the resonances that are formed in the reaction are described by the partial-wave amplitudes $T_{\Delta\pi}^{JILL'}$, etc., where J is the total angular momentum and I is the total isospin of the state formed, and L and L' are the orbital angular momenta in the initial 2-body and the final quasi-2-body states. So $\vec{J} = \vec{L} + \vec{S} = \vec{L}' + \vec{S}'$, where S and S' are the initial and final total spins. For the ρN amplitude, it is necessary to add the suffix $2S'$, equal to 1 or 3, to indicate the total ρN spin. The partial-wave amplitudes are often denoted by $\Delta \cdot L \cdot L' \cdot 2I \cdot 2J$, $\rho_{2S} \cdot L \cdot L' \cdot 2I \cdot 2J$, etc.

The Listings give the results from four analyses.

LONGACRE 75 (LBL-SLAC) is an analysis of 200K $\pi^- p \rightarrow \pi^- \pi^0 p$, $\pi^- p \rightarrow \pi^- \pi^+ n$, and $\pi^+ p \rightarrow \pi^+ \pi^0 p$ events for $1300 \leq W \leq 2000$ MeV. Approximate unitarity constraints are imposed using a simplified K-matrix formalism that links the $\pi\pi N$ channel to the πN channel. This gives smooth solutions and eliminates the overall phase ambiguity at each energy. The $\Delta\pi$, ρN , and ϵN intermediate states are included. Couplings and T-matrix pole positions are given for 14 resonances.

LONGACRE 77 (Saclay) is a coupled-channel analysis similar to LONGACRE 75 that fits 100K events for $1380 \leq W \leq 1740$ MeV. The couplings and pole positions are found for 16 resonances, including a P13 $N(1540)$ and a P31 $\Delta(1550)$ suggested for the first time by this analysis.

NOVOSELLER 78 (CIT) is an analysis of $\pi^- p \rightarrow \pi^- \pi^0 p$, $\pi^- p \rightarrow \pi^- \pi^+ n$, and $\pi^+ p \rightarrow \pi^+ \pi^0 p$ for $1630 \leq W \leq 1990$ MeV based on the LBL-SLAC energy-independent analysis.¹ Again the $\Delta\pi$, ρN , and ϵN states are used, but the resonances are fitted using a simple Breit-Wigner amplitude rather than the K-matrix formalism of LONGACRE 75. Single-pion exchange with $\pi\pi$ rescattering is used to calculate

the higher partial waves (taking account of a criticism made of earlier analyses), and it is concluded that this improves the fit above 1800 MeV and helps eliminate the phase ambiguity. Another study of the importance of single-pion exchange has been made by Aaroe et al.,² who also find that it can give important corrections to the angular dependence. NOVOSELLER 78 gives two solutions, the second including the effects of pion exchange. They are given in the Listings as fits to LONGACRE 75 and NOVOSELLER 78.

BARNHAM 80 (Imperial College) is an analysis of 44K $\pi^+ p \rightarrow \pi^+ \pi^0 p$ and $\pi^+ p \rightarrow \pi^+ \pi^+ n$ events for $1440 \leq W \leq 1700$ MeV. It thus concerns only the Δ resonances, and it uses data that were not available to the other analyses. The intermediate states are $\Delta\pi$, ρN , and $\pi N(1440)$, the last being necessary to account for the difference between the $\pi^+ \pi^0 p$ and $\pi^+ \pi^+ n$ cross sections. Also included is the effect of single-pion exchange leading to the S-wave $\pi\pi$ state with $I = 2$. The phase ambiguity is resolved by requiring that the $\pi\Delta$ amplitude for the D33 $\Delta(1700)$ have a Breit-Wigner phase. The parameters are found for four resonances, including the P31 $\Delta(1550)$, but since some of the data were also used by LONGACRE 77 it is not clear that this resonance is being confirmed. However, there is now some evidence that it has also been seen in single-pion photoproduction.³

There have also been analyses of $\pi^- p \rightarrow \pi^- \pi^+ n$ for $W \leq 1400$ MeV,^{4,5} a range dominated by the low energy tail of the P11 $N(1440)$, with ϵN as the dominant intermediate state. The energy range of these analyses is too low for them to determine decay coupling constants, and they do not appear in the Listings.

It is difficult to assess the systematic uncertainties of these analyses or of the quoted couplings of the resonances to the isobar states. However, those that are indicated in the Listings as being well determined by LONGACRE 77 do in general agree, at least in sign, with the values from the other analyses, although some of the D_{33} couplings have not been measured elsewhere. The Imperial College group also claims clear measurements of the signs of ρ_{SS31} , ρ_{DD33} , and N^*_{PP33} .

Baryons N's and Δ 's

Data Card Listings

For notation, see key at front of Listings.

All existing isobar models can be criticized for neglecting possible sub-energy dependence of the partial-wave amplitudes, which makes them possibly inconsistent with unitarity.⁶ This has been studied by Aitchison and Brehm,⁷ who derive an isobar expansion that is consistent with Bose symmetry and with sub-energy analyticity and unitarity. The resulting coupled integral equations are suitable for both dynamical and phenomenological studies of $\pi N \rightarrow \pi N$. They estimate the sub-energy corrections to the isobar model and conclude that such corrections may not be significant for existing isobar fits but that they could become so with better data.⁸ A rough estimate of these corrections has also been made by the Imperial College group,⁹ who also find that they are small.

References for Section III

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IV. Photoproduction and Compton Scattering

(by R.L. Crawford, University of Glasgow)

The γN couplings of an N or Δ resonance can be studied in any photon-induced formation process in which the coupling of the resonance to the final state is well known. In practice, this limits accurate sources for such information to partial-wave analyses of single-pion photoproduction, where

there is a large amount of data and where the final state is well known from $\pi N \rightarrow \pi N$ partial-wave analyses. Recently, however, some couplings have been obtained from Compton scattering on protons. All photoproduction analyses rely heavily on $\pi N \rightarrow \pi N$ analyses for the existence, masses, and widths of the resonances. In only a few photoproduction analyses are the masses and widths treated as free parameters. However, the photoproduction results for the masses and widths are of interest since they give access to the charge +1 states.

The most important analyses of single-pion photoproduction are reviewed below. The formalism has been described in an earlier edition of this Review,¹ to which the reader is referred for additional information. There are three basic methods of analysis. All have to cope with the stability problems of having four independent complex spin amplitudes at any energy and angle, but only six (and frequently fewer) independent experimental measurements.

(a) Simple isobar model: This is the simplest form of energy-dependent analysis. The partial waves are parametrized as a smooth background to which Breit-Wigner resonances are added. Usually the electric but not the magnetic Born terms are included explicitly to reproduce the forward peak in charged pion production. This method is sufficiently flexible to give excellent fits to the data, but there are in principle difficulties concerning the uniqueness of the solution due to the large number of waves involved. This problem is eliminated by the form of the parametrization, but it is not clear how this may bias the solution.

The most extensive analysis of this type is METCALF 74 (for references in this form, see the Data Card Listings), which is an extension of the earlier Walker analysis.² It fits $\gamma p \rightarrow \pi^+ n$, $\gamma p \rightarrow \pi^0 n$, and $\gamma n \rightarrow \pi^+ p$ from the first to the fourth resonance regions. FELLER 76 is a similar analysis which does not fit $\gamma n \rightarrow \pi^+ p$, but uses data that were not available to METCALF 74. Other isobar analyses are ROSSI 73, HEMMI 73, HEMMI 73, BENEVENTANO 74, KRIVETS 74, TAKEEDA 80, and BRATASHEVSKIJ 80; these have used relatively small and sometimes restricted data sets, usually in association with a particular experiment. ISHII 80 is an

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For notation, see key at front of Listings.

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isobar analysis of proton Compton scattering for laboratory photon energies from 550 MeV to 950 MeV, covering the second resonance region.

(b) Fixed-t dispersion relations (FTDR): This technique uses the apparent resonance dominance of the photoproduction amplitudes to get a relatively simple parametrization of their imaginary parts. Fixed-t dispersion relations are then used to calculate the real parts; sometimes a few but usually no additional free parameters are introduced. Since there are fewer free parameters, the probability of having multiple solutions is reduced, and the requirements of analyticity are automatically satisfied. However, the method is inflexible compared to the isobar model and gives poorer fits. Also, as has been described in NOELLE 78 and elsewhere,³ the divergence of the partial-wave expansions for the dispersion integrals does not allow the use of data at all angles above the third resonance region. Some but not all analyses satisfy the constraints of unitarity and time reversal invariance as given by Watson's theorem.⁴

FTDR analyses have been made by groups at Berkeley (MOORHOUSE 73, KNIES 74, and MOORHOUSE 74), at Lancaster (DEVENISH 73 and DEVENISH2 74), at Glasgow (CRAWFORD 75, BARBOUR 76, BARBOUR 78, and CRAWFORD 80), at Yerevan (AZNAURYAN 77), and at Tokyo (ARAI 80 and FUJII 81). NOELLE 78 is a hybrid analysis incorporating FTDR in a coupled-channel isobar calculation.

(c) Energy independent analyses: These evaluate the partial waves by making fits at a set of essentially single energies, and are thus the least biased of all analyses. It is necessary to use Watson's theorem to fix the complex phases of many of the partial waves and thus to get a unique solution. Due to the onset of inelasticity, this becomes difficult above the first resonance region, and only BERENDS 77 extends into the second resonance region. This analysis gets significantly different couplings from the other analyses for the D13 N(1520).

New analyses in the Data Card Listings: ARAI 80 is an FTDR analysis of $\gamma p + \pi^+ n$, $\gamma p + \pi^0 p$, and $\gamma n + \pi^+ p$ using 7768 data points for energies from 1168 to 2078 MeV and for momentum transfers t out

to -1.6 (GeV/c)². A K-matrix formalism is used to parametrize the resonances, and a Regge formalism is used for the high energy part of the dispersion integrals.

CRAWFORD 80 is an extension of earlier Glasgow analyses to include more data and more resonances. The imaginary parts of the partial waves are parametrized as Breit-Wigner amplitudes plus background in the S and P waves. The high energy parts of the dispersion integrals are parametrized using a Regge formalism, and data are fitted at all energies for $\gamma p + \pi^+ n$, $\gamma p + \pi^0 p$, and $\gamma n + \pi^+ p$. A total of 8838 data points are used for energies up to 2.5 GeV and for t out to -1.5 (GeV/c)². Evidence is found to suggest that the P31 Δ(1550) is photoproduced, but it is not conclusive.

BRATASHEVSKIJ 80 is an isobar analysis of $\gamma p + \pi^0 p$ using recoil proton polarization data from Kharkov and is essentially a variation of METCALF 74.

TAKEDA 80 is similarly a variation of METCALF 74 to fit new recoil proton polarization data in $\gamma n + \pi^+ p$.

FUJII 81 is based on the ARAI 80 analysis and includes new neutron target data.

Resonance couplings and errors in the Data Card Listings: The Listings give the results of all recent and extensive analyses. If no error is given, only a unique result is quoted. The Berkeley analyses and CRAWFORD 75 give for errors the spread of solutions around a central value. The Lancaster group gives for errors the change of value that is required to increase the "best possible χ^2 " by 1%. METCALF 74, FELLER 76, AZNAURYAN 77, and ARAI 80 quote similar errors. In BARBOUR 78 and CRAWFORD 80, the systematic differences due to the different methods of analysis are considered to be at least as important as the purely statistical errors, and the errors quoted are obtained by comparing with other analyses as well as from the spread of parameters over a number of fits. Thus there is often a wide variation in what the errors quoted in the Listings mean.

Table IV.1 gives a compilation of the γN decay couplings from METCALF 74, KNIES 74, MOORHOUSE 74, DEVENISH2 74, BARBOUR 76, BARBOUR 78, ARAI 80, and CRAWFORD 80. The errors quoted are a combination

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Table IV.1. A compilation of measured ΥN decay couplings and of predictions of the quark model. Sources are given in the text.

Resonance	Target	Helicity	Couplings ($\text{GeV}^{-1/2} \times 10^{-3}$)		
			Partial-Wave Analyses	Quark Model Predictions	
P ₁₁	N(1440)	p	1/2	-70 ± 9	-5 to -50
		n	1/2	$+42 \pm 19$	+4 to +38
D ₁₃	N(1520)	p	1/2	-17 ± 11	-41 to +6
			3/2	$+166 \pm 7$	+95 to +174
		n	1/2	-69 ± 14	-23 to -52
			3/2	-136 ± 14	-102 to -144
S ₁₁	N(1535)	p	1/2	$+67 \pm 15$	+68 to +147
		n	1/2	-78 ± 29	-83 to -119
S ₁₁	N(1650)	p	1/2	$+45 \pm 17$	-9 to +95
		n	1/2	-23 ± 33	-45 to +4
D ₁₅	N(1675)	p	1/2	$+13 \pm 8$	0 to +12
			3/2	$+22 \pm 12$	0 to +16
		n	1/2	-37 ± 24	-31 to -55
			3/2	-54 ± 24	-44 to -78
F ₁₅	N(1680)	p	1/2	-13 ± 10	-7 to +24
			3/2	$+132 \pm 15$	+47 to +154
		n	1/2	$+29 \pm 15$	-32 to +27
			3/2	-28 ± 16	-25 to +2
D ₁₃	N(1700)	p	1/2	-20 ± 12	-7 to +9
			3/2	$+1 \pm 17$	-12 to +33
		n	1/2	0 ± 50	-15 to +25
			3/2	$+8 \pm 42$	-17 to -76
F ₁₁	N(1710)	p	1/2	$+3 \pm 15$	-7 to -37
		n	1/2	$+9 \pm 30$	-21 to +29
P ₁₃	N(1720)	p	1/2	$+60 \pm 36$	-133 to +74
			3/2	-34 ± 26	-65 to +46
		n	1/2	-6 ± 25	-23 to +57
			3/2	-26 ± 90	-61 to +12
F ₁₇	N(1990)	p	1/2	$+17 \pm 35$	-8 to -10
			3/2	0 ± 20	-10 to -13
		n	1/2	-82 ± 25	-18 to -19
			3/2	-86 ± 40	-23 to -25
G ₁₇	N(2190)	p	1/2	-43 ± 40	
			3/2	$+131 \pm 100$	
		n	1/2	-64 ± 45	
			3/2	-60 ± 130	
F ₃₃	$\Delta(1232)$	p	1/2	-141 ± 6	-94 to -127
			3/2	-258 ± 8	-162 to -220
F ₃₃	$\Delta(1600)$	p	1/2	$+2 \pm 30$	-61 to +2
			3/2	-3 ± 30	-107 to +4
S ₃₁	$\Delta(1620)$	p	1/2	$+23 \pm 38$	+43 to +86
D ₃₃	$\Delta(1700)$	p	1/2	$+109 \pm 31$	+78 to +106
			3/2	$+73 \pm 35$	+70 to +105
F ₃₅	$\Delta(1905)$	p	1/2	$+33 \pm 10$	+10 to +44
			3/2	-38 ± 19	-41 to +15
F ₃₁	$\Delta(1910)$	p	1/2	-24 ± 15	-16 to +15
D ₃₅	$\Delta(1930)$	p	1/2	-46 ± 50	-17
			3/2	$+13 \pm 50$	-24
F ₃₇	$\Delta(1950)$	p	1/2	-73 ± 12	-25 to -48
			3/2	-85 ± 13	-32 to -69

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For notation, see key at front of Listings.

of the statistical errors from these analyses and the systematic differences between them. Also shown for comparison are the range of predictions for the couplings from recent quark model calculations.⁵⁻⁸ While the quantitative agreement between the quark models and the measured couplings is not yet good, there is qualitative agreement to the extent that at least one quark model gives the correct sign in the cases where the sign has been clearly measured.

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V. Electroproduction

(by F. Foster, University of Lancaster)

Introduction: Since this review of resonance electroproduction was last revised (1978), several new data sets on virtual photoproduction of single pions have become available. The DESY group has completed its detailed survey of the whole resonance region at values of Q^2 from 0.6 to 3.0 GeV^2 . The new data on π^+ and π^0 production and an analysis using the fixed- t dispersion relation methods of Devenish and Lyth¹ are available in internal DESY reports,^{2,3} and those on η production are available in reports^{2,4} and a publication.⁵ Data at smaller Q^2 continue to appear from the Bonn group, for η production⁶ at $Q^2 = 0.4 \text{ GeV}^2$ and for π^+ and π^0 production⁷ at $Q^2 = 0.3 \text{ GeV}^2$. The Lancaster-Manchester group working at Daresbury has also completed new data sets at 0.5 and 1.0 GeV^2 using both hydrogen^{8,9} and deuterium as targets.¹⁰⁻¹²

Analysis of the new data in terms of the appropriate resonance multipole amplitudes has been

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facilitated by the availability of the Devenish-Lyth fitting program, which uses fixed-t dispersion relation constraints.¹ Ch. Gerhardt has analyzed much of the data available to 1978,¹³ while Davenport and Morris have analyzed recent Daresbury data¹⁴ and V. Gerhardt has analyzed the recent DESY data.³ There are now quite consistent results for, in addition to the P33 $\Delta(1232)$ and the S11 N(1535) resonance multipoles, the D13 N⁺(1520) and the F15 N⁺(1680) multipoles for Q^2 up to 3.0 GeV², and there are some clear indications of the behavior of the D33 $\Delta^+(1700)$. Analysis of the Daresbury "neutron target" data¹² has also given good results for the S11 N⁰(1535) and D13 N⁰(1520) multipoles at $Q^2 = 0.5$ GeV² and allowed significant tests of the generalized SU(6)_W or single-quark excitation models. On the theoretical side, Alcock et al.¹⁵ have presented an excellent review of all recent quark models for resonance electroproduction, as well as their own proposals for suitable quark models on a proper relativistic basis. They conclude that the single-quark transition model can reproduce qualitatively the observed behavior of the orbit-flip, spin-flip, and spin-orbit parameters for the excitation of members of the [70,1⁻] multiplet.

First resonance region: Excitation of the P33 $\Delta(1232)$ resonance proceeds mainly via the magnetic M_{1+} multipole; the electric E_{1+} multipole is consistent with zero as is expected from SU(6) and the quark model. The scalar multipole has a small but significant negative value with S_{1+} being between 5 and 10% of M_{1+} , and as Q^2 increases the background amplitudes E_{0+} and M_{1-} account for an increasing fraction of the single-pion cross section. The new DESY data^{2,3} at $Q^2 = 3.0$ GeV² give $E_{1+}/M_{1+} \approx 0.05 \pm 0.05$ and $S_{1+}/M_{1+} \approx -0.12 \pm 0.05$ at 1230 MeV, and the background ratios are $E_{0+}/M_{1+} \approx 0.5$ and $M_{1-}/M_{1+} \approx -0.3$ near resonance. From its measurements of M_{1+} , the DESY group also obtains a value for the transition form factor $G_M^+(Q^2)$ at $Q^2 = 3.0$ GeV² that is much less subject to systematic error than were the earlier "single arm" measurements of the total cross section. The results confirm that $G_M^+(Q^2)$ falls more rapidly with Q^2 than does the nucleon "dipole" form factor. It is also clear that at the high values of Q^2 the Devenish-Lyth program has

great difficulty in fitting the shape of the M_{1+} variation with mass, and also in accommodating the observed background excitation, particularly above 1300 MeV. An unfortunate consequence is that the program cannot give unambiguous results for the excitation of the P11 N(1440) Roper resonance.

The [70,1⁻] resonances: Table V.1 gives the results of the various dispersion-relation fits for the low-lying negative-parity resonances S11 N(1535), D13 N(1520), S11 N(1650), and D33 $\Delta(1700)$ from the [70,1⁻] multiplet. It is difficult to

Table V.1. Results of recent fits to π^0 , π^+ , and η data for members of the [70,1⁻] multiplet.

Resonance	Multipole (π^+)	Multipole $\mu b^{1/2}$		
		$Q^2=0.6$	1.0	
$Q^2 = 0.3$ GeV ² Bonn (Refs. 7, 17)	D13 N ⁺ (1520)	E_{2-}	0.34 ± 0.03	
		M_{2-}	0.58 ± 0.06	
	S11 N ⁺ (1535)	E_{0+}	0.60 ± 0.07	
	S11 N ⁺ (1650)	E_{0+}	0.33 ± 0.04	
	$Q^2 = 0.5$ GeV ² Lancaster- Manchester (Ref. 14)	D13 N ⁺ (1520)	E_{2-}	0.37 ± 0.13
			M_{2-}	0.53 ± 0.07
		S_{2-}	0.10 ± 0.10	
D33 $\Delta^+(1700)$		E_{2-}	0.23 ± 0.13	
		M_{2-}	-0.17 ± 0.03	
		S_{2-}	-0.05 ± 0.05	
$Q^2 = 1.0$ GeV ² Lancaster- Manchester (Ref. 14)	D13 N ⁺ (1520)	E_{2-}	0.08 ± 0.02	
		M_{2-}	0.43 ± 0.05	
		S_{2-}	0.0	
	D33 $\Delta^+(1700)$	E_{2-}	0.18 ± 0.07	
		M_{2-}	-0.06 ± 0.02	
		S_{2-}	-0.04	

Resonance	Multipole (π^+)	DESY data (Ref. 3)			
		$\mu b^{1/2}$			
		$Q^2=0.6$	1.0	2.0	3.0 GeV ²
S11 N(1535)	E_{0+}	0.52	0.51	0.39	0.32
	S_{0+}	-0.30	0.00	-0.11	-0.10
D13 N(1520)	E_{2-}	0.13	-0.04	-0.11	-0.08
	M_{2-}	0.47	0.35	0.17	0.063
	S_{2-}	0.06	0.00	0.00	0.00

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assign uncertainties to these results, but where this has been done it reflects the spread observed in independent fits to the same experimental data. Also, the S11 N(1535) multipoles may be regarded as being very well determined since they are extracted from η electroproduction data of high accuracy. The low- Q^2 data from Bonn were analyzed for the helicity 1/2 couplings only, and the multipoles here have been extracted by combining the helicity 3/2 couplings from the analysis of Ch. Gerhardt.¹³ The trends in the data are clear: the E_{0+} multipole falls off very slowly so that the S11 N(1535) is the dominant component of the second resonance at $Q^2 > 1.0 \text{ GeV}^2$. Using only η production cross sections and resonance peak heights from total-cross-section measurements, the following results are obtained:³

Q^2 (GeV^2)	σ_{peak} (μb)	S11 (μb)	D13 (μb)
2.0	15.0 ± 1.0	10.07 ± 0.56	5.0 ± 1.2
3.0	7.7 ± 1.1	6.85 ± 0.37	0.88 ± 1.2

It has also been demonstrated that the effect is unlikely to be due to scalar/longitudinal photon excitation, since the following results are obtained for σ_L/σ_T using η production cross sections at different values of the photon polarization:

Q^2 (GeV^2)	σ_L/σ_T	Reference
0.4	0.23 ± 0.14	6
0.6	0.25 ± 0.23	4
1.0	-0.13 ± 0.16	4

In contrast, the electric multipoles for the D13 and D33 resonances fall rather rapidly with Q^2 . This is particularly noticeable for the D13 N(1520), where E_{2-} falls from $1.05 \mu\text{b}^{1/2}$ at $Q^2 = 0$ to less than $0.1 \mu\text{b}^{1/2}$ at $Q^2 = 1.0 \text{ GeV}^2$, while the magnetic multipole M_{2-} falls much more slowly. The effect is associated with the well-known rapid changes in excitation of the resonances from mainly helicity 3/2 at $Q^2 = 0$ to mainly helicity 1/2 at $Q^2 > 1 \text{ GeV}^2$. More detailed discussions may be found elsewhere (e.g., Ref. 17). An alternative and more instructive point of view is to compare the Q^2 variation of the single-quark-transition-model (SQTM) parameters A, B, and C which describe the

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excitation of all the members of the $[70, 1^-]$ multiplet: A and B correspond to orbit- and spin-flip excitation, and C corresponds to simultaneous spin and orbit excitation. The SQTM works well in the photoproduction limit where the parameters are necessary and sufficient to describe all the excitation¹⁶ and A, B, and C are roughly in the ratio 9:4:4. The relations between A, B, and C and the dominant multipole amplitude are given in Ref. 17, and it is straightforward to determine A, B, and C as functions of Q^2 using the results outlined in Table V.1. Figure V.1 shows these results, and it is certain that at high Q^2 the spin-flip amplitude B dominates. This behavior is expected from most explicit quark models,^{15,17} and Fig. V.1 also shows the results of a recent quark model calculation¹⁸ based on the successful photoproduction model of Kubota and Ohta.¹⁹ Two points may be mentioned: (1) agreement between the explicit model and the data could only be achieved by including an effective quark "form factor;" and (2) at high Q^2 the spin-orbit term C seems to go significantly negative in contradiction to quark model expectations.^{15,20}

The SQTM may be further tested by comparing the predictions for the neutron-target multipoles of $I=1/2$ resonances calculated from the observed values of A, B, and C with the results of the Daresbury deuterium-target experiments. This has been done with some degree of success by the Daresbury group,^{10-12,17} and their results are given in Table V.2.

The P11 N(1440): Interest in this enigmatic resonance persists because two rather different theoretical approaches both predict that the M_{1-} multipole sign will change as Q^2 increases from zero. In photoproduction, M_{1-} has the wrong sign according to the simplest quark models, but consideration of the large spin-orbit contribution to the excitation in an explicit model¹⁹ gives rough agreement with experiment. As in the $[70, 1^-]$ multiplet, spin-orbit excitations become less important as Q^2 increases, so the multipole should revert quickly to its nominal sign.¹⁸ The same effect is predicted by Gavela²¹ using a "bag" model for electroproduction. Experimentally the situation is still far from clear, but such sign changes

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are certainly not ruled out.¹⁷

The F15 N(1680): All analyses of the data near the third resonance give consistent values for

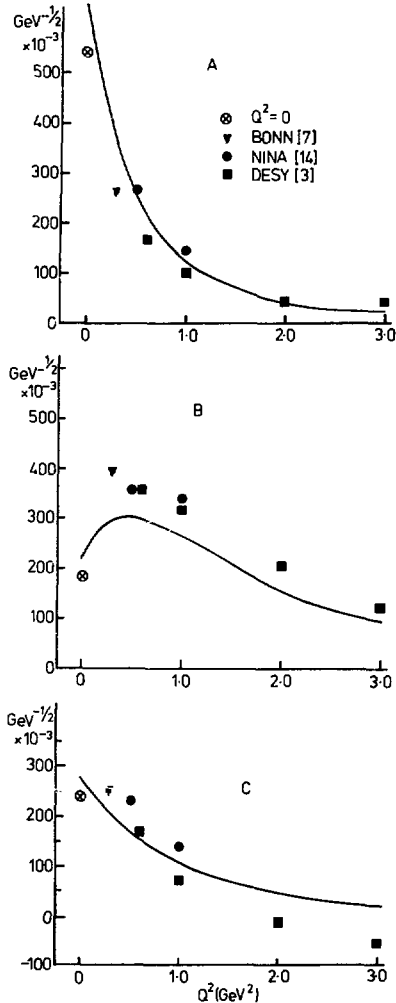


Fig. V.1. The Q^2 variation of the orbit-flip, spin-flip, and spin-orbit-flip parameters A, B, and C for the $[70, 1^-]$ multiplet. The curves are from Ref. 18.

the E_{3-} and M_{3-} multipoles for this $[56, 2^+]$ multiplet member. Table V.3 shows the results, and again the electric multipole falls much more rapidly than the magnetic multipole. This clearly agrees with the quark model expectation that spin-flip amplitudes must dominate at high Q^2 .

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Table V.2. A comparison of neutron-target results at $Q^2=0.5 \text{ GeV}^2$ (Ref.12) with SQTM predictions.

Resonance	Multipole (π^-)	Data $\mu\text{b}^{1/2}$	SQTM
S11 $N^0(1535)$	E_{0+}	-0.11	-0.46
D13 $N^0(1520)$	E_{2-}	-0.55	-0.43
	M_{2-}	-0.16	-0.07
D33 $\Delta(1700)$	E_{2-}	-0.13	-0.20
	M_{2-}	+0.029	+0.036
D13 $N(1700)$	E_{2-}	-0.034	-0.093
	M_{2-}	+0.016	-0.037
S11 $N(1650)$	E_{0+}	-0.09	-0.34

Table V.3. The F15 N(1680) resonance amplitudes.

Q^2 (GeV^2)	Amplitudes (μb) ^{1/2}			Ref.
	E_{3-}	M_{3-}	S_{3-}	
0.0	0.47	0.16	--	
0.3	0.16 ± 0.02	0.27 ± 0.03	--	7
0.5	0.12 ± 0.10	0.27 ± 0.01	-0.02 ± 0.02	14
0.6	0.13	0.17	0.01	3
1.0	0.01	0.23 ± 0.02	0.01 ± 0.01	14
1.0	0.07	0.15	0.00	3
2.0	0.013	0.083	0.00	3
3.0	-0.084	0.038	0.00	3

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Table listing baryon data for N(1440) with columns for name, mass, spin, and other properties. Includes entries like N(1440) 1/2+, N(1440) 3/2+, etc.

Table listing baryon data for N(1440) with columns for name, mass, spin, and other properties. Includes entries like N(1440) 1/2+, N(1440) 3/2+, etc.

THE FOLLOWING ARE THEORETICAL PAPERS CONCERNING THE N(1440) STATES... SCHWARTZ 66 PP 152 1325 J SCHWARTZ (LL)

(PI) N ENHANCEMENTS PEAKING WELL BELOW 1400 MEV... FUKUNAGA 80 HBC + P1+P TO PI P1 N 2/82+

1440 MEV REVISION PRODUCTION EXPERIMENTS

UNDER THIS HEADING WE INCLUDE ALL BUMPS WHICH LIE WELL BELOW 1500 MEV... AND DELTA LISTINGS FOR A DISCUSSION OF PRODUCTION EXPERIMENTS.

91 N(1440) PARTIAL DECAY MODES (PROD. EXP.)

Table showing partial decay modes for N(1440) with columns for mode, branching ratio, and other parameters.

91 N(1440) MASS (MEV) (PROD. EXP.)

Table showing mass measurements for N(1440) with columns for experiment, mass value, and error.

91 N(1440) BRANCHING RATIOS (PROD. EXP.)

Table showing branching ratios for N(1440) with columns for decay mode, ratio, and other parameters.

REFERENCES FOR N(1440) (PROD. EXP.)

Table listing references for N(1440) production experiments with columns for author, year, and journal.

91 N(1440) WIDTH (MEV) (PROD. EXP.)

Table showing width measurements for N(1440) with columns for experiment, width value, and error.

PAPERS NOT REFERRED TO IN DATA CARDS

Table listing papers not referred to in data cards with columns for author, year, and journal.

Data Card Listings

For notation, see key at front of Listings.

Baryons
N(1520), N(1535)

N(1520) PARTIAL DECAU MODES (PROD. EXP.)
 # N(1520) PARTIAL DECAU MODES (PROD. EXP.)

P1 #N(1520) INTO PI N 139 938
 P2 #N(1520) INTO N(3/2) 1272 139
 P3 #N(1520) INTO N PI 1 438 139 139
 P4 #N(1520) INTO INU NEUTRON PI 1 438 139
 P5 #N(1520) INTO PROTON PI 1 438 139 139
 P6 #N(1520) INTO N ETA 938 448
 P7 #N(1520) INTO N PI(1/2), I=0 538 139 139
 P8 #N(1520) INTO N RHO 538 769

N(1520) BRANCHING RATIOS (PROD. EXP.)

R1 #N(1520) INTO IN PI/TOTAL (PI) 1
 R2 #N(1520) INTO IN PI/TOTAL PRODUCTION EXPERIMENTS
 R3 #N(1520) INTO IN PI/TOTAL
 R4 #N(1520) INTO IN PI/TOTAL
 R5 #N(1520) INTO IN PI/TOTAL
 R6 #N(1520) INTO IN PI/TOTAL
 R7 #N(1520) INTO IN PI/TOTAL
 R8 #N(1520) INTO IN PI/TOTAL
 R9 #N(1520) INTO IN PI/TOTAL

REFERENCES FOR #N(1520) (PROD. EXP.)
 +ROPELL 67 4 67 237 ALLES-BORELLI, FRENCH, P.F.I.R., NICHEJDA (CERN)
 ALEXANDER 67 20 1274 ALLEN, BENARY, CZAPKA, FINE (MANNHEIM)
 BASSOMPIE 67 PL 208 462 BASSOMPIE, LEBLANC (CERN), BRUJLLET
 LEE 67 PL 559 1136 +DESS, ROE, SINCLAIR, VANDEP VELDE (MICH)

ANDERSON 70 PL 2 689 +BESER, ALLEDY, COLLINS+ (BVL, CERN)
 AMALDI 71 PL 346 435 +BIVACASTELLI, BOSIO, GI SANTIA ROMA (CERN)
 SILLS 71 PL 27 442 +MAGLICH, HORNEM, SAVINES, STEINBERG (RTG)
 NORSE 71 PR 04 132 +DM, WALKER, CARROLL, LYNCH + (MESC+ATD) JJ

EDLSTEIN 72 PR 55 1073 EDLSTEIN, CARLIGNI, MECHE, MAHOLI, VIGNALONI (CERN)
 JOHNSON 72 PR 488 477 +MOLLERUS, +JACOBSEN, NORHE, HALLS, OSLO, STOKHJPP
 KUN 72 PL 428 487 +FUNG, KEMAN, PODE-SKLAM, SIKEN (CERN) JJ
 COOPER 74 PR 875 259 COOPER, SEIDL, VANDERVELDE (MICH)

BRUHN 75 PR 305 481 +GERBER, MAUER, NICHOLSON, SCHIBY (ESTR, LDP) JJ
 MUSGRABE 75 PR 367 565 +RETERS, SCHRIEBER, MITT MOREL, YUTS (ANU)
 STRACHMAN 75 PR 808 120 +STRACHMAN, BRAUN, GERBER, MAURER (LDP, STRAB)
 ALSO 76 PL 107 330 STRACHMAN
 WEBB 75 PL 458 331 +TRILLING, TELEGI + (AACH+KLAU+UCR+CE) JJ
 ALSO 75 PL 558 335 +TRILLING, TELEGI + (AACH+KLAU+UCR+CE) JJ

BOUGE 77 PL 608 115 +DE ROSSY, FLEURY, RIVET+ (LDP+KDEF+LLOI)
 FERREER 78 PR 3142 77 +BOUGEUT, D. ALMAGRE+ (LLOI+CERN+DEF+EP) JJ JJ

PAPERS NOT REFERRED TO IN DATA CARDS

RUSHERD 76 PR 13 1835 +S, SHRODKE, RAJA, ANSPODE, CARTER, HALE (CRAW) JJ
 DONT1910U D, SONT1910U (CERN) JJ
 BIEL 76 PR 30 307 +BESSEL, SCLAFER+ (RGCH+MFNL+ACH+HNS)
 ALSO 76 PL 36 464, 507 BIEL, KEMAN, FERBER+ (RGCH+MFNL+ACH+HNS)
 CHMONDICH 76 PR 17 1713 +CARROLL, CHALOUPOUK, BALLAM+ (SLAC+SLAC+HNS) JJ
 GOSPIE 76 PR 3443 3005 +CAWELL, SFRORZA+KONTAR+ (CERN+PA) JJ
 GOSPIE 79 PR 8161 14 +CONTAR, FRATERNALI, LILIVAN+ (CERN+PA) TRST

N(1535)

N(1535) JP=1/2- I=1/2 S₁₁
 THIS RESONANCE IS WELL ESTABLISHED.
 IT COUPLES STRONGLY TO THE ETA NUCLEON CHANNEL.

N(1535) MASS (MEV)
 HENDRY 65 PVUE ETA N 511 PI N 9/66
 MICHAEL 66 PVUE FITS BARREYRE S11 7/66
 UCHIMAYA 64 PVUE FITS N ETA N 9/66
 FITTING GIVES TWO SOLUTIONS. PROBLEMS MATCHING OF PHASE SHIFT
 BARREYRE 68 PVUE PHASE-SHIFT ANAL 11/67
 WHERE CROSS SECTION IS GREATEST = EVERALL FIT
 DONNACHI 68 PVUE PHASE-SHIFT ANAL 8/68
 DELCOURT 69 CNTR PHOTO PDDUCT. 8/69
 AYED 70 IPWA P-N ANAL SOL 8 8/69
 HICKS 73 MPWA GAM-D+ETA 0 9/73
 CRANFORD 75 DPWA PI N PHOTO PDD. 1/78
 LONGACRE 75 IPWA PI N TO 2P1 N 11/75
 BARBOUR 76 DPWA PI N PHOTO PDD. 1/78
 BERENDS 77 IPWA PI-N PHOTOPROD. 1/78
 BHANDARI 77 DPWA O USES ETA N CUSP 1/78
 LONGACRE 77 IPWA PI N TO 2P1 N 11/77
 ALL LONGACRE PARAMETERS ARE FROM SOLUTION 52, EXCEPT FOR THE POLE
 POSITION WHICH IS FROM SOLUTIONS 51 AND 51
 SUPPOSES BARBOUR 76.

N(1540), 120.1, 120.1, 120.1, 120.1
 HENDRY 65 PVUE 9/66
 MICHAEL 66 PVUE 7/66
 UCHIMAYA 64 PVUE 9/66
 BARREYRE 68 PVUE 11/67
 DONNACHI 68 PVUE 8/68
 DELCOURT 69 CNTR PHOTO PDDUCT. 8/69
 AYED 70 IPWA P-N ANAL SOL 8 8/69
 HICKS 73 MPWA GAM-D+ETA 0 9/73
 CRANFORD 75 DPWA PI N PHOTO PDD. 1/78
 LONGACRE 75 IPWA PI N TO 2P1 N 11/75
 AYED 76 IPWA 11/77
 BARBOUR 76 DPWA PI N PHOTO PDD. 1/78
 BERENDS 77 IPWA PI-N PHOTOPROD. 1/78
 BHANDARI 77 DPWA O USES ETA N CUSP 1/78
 LONGACRE 77 IPWA PI N TO 2P1 N 11/77
 BARBOUR 78 DPWA PI-N PHOTOPROD. 3/79
 KUTSKY 79 DPWA O PI-N TO ETA N 12/79
 CRANFORD 80 DPWA PI N PHOTO PDD. 12/81
 KUTSKY 80 IPWA PI N TO PI N 1/82

N(1535) WIDTH (MEV)

HENDRY 65 PVUE 9/66
 MICHAEL 66 PVUE 7/66
 UCHIMAYA 64 PVUE 9/66
 BARREYRE 68 PVUE 11/67
 DONNACHI 68 PVUE 8/68
 DELCOURT 69 CNTR PHOTO PDDUCT. 8/69
 AYED 70 IPWA P-N ANAL SOL 8 8/69
 HICKS 73 MPWA GAM-D+ETA 0 9/73
 CRANFORD 75 DPWA PI N PHOTO PDD. 1/78
 LONGACRE 75 IPWA PI N TO 2P1 N 11/75
 AYED 76 IPWA 11/77
 BARBOUR 76 DPWA PI N PHOTO PDD. 1/78
 BERENDS 77 IPWA PI-N PHOTOPROD. 1/78
 BHANDARI 77 DPWA O USES ETA N CUSP 1/78
 LONGACRE 77 IPWA PI N TO 2P1 N 11/77
 BARBOUR 78 DPWA PI-N PHOTOPROD. 3/79
 KUTSKY 79 DPWA O PI-N TO ETA N 12/79
 CRANFORD 80 DPWA PI N PHOTO PDD. 12/81
 KUTSKY 80 IPWA PI N TO PI N 1/82

N(1535) REAL PART OF POLE POSITION (MEV)

LONGACRE 75 IPWA PI N TO 2P1 N 11/75
 BHANDARI 77 DPWA O USES ETA N CUSP 1/78
 LONGACRE 77 IPWA PI N TO 2P1 N 11/77
 KUTSKY 79 IPWA PI N TO PI N 12/79
 KUTSKY 80 IPWA PI N TO PI N 1/82

N(1535) -PI MAG PART OF POLE POSITION (MEV)

LONGACRE 75 IPWA PI N TO 2P1 N 11/75
 BHANDARI 77 DPWA O USES ETA N CUSP 1/78
 LONGACRE 77 IPWA PI N TO 2P1 N 11/77
 KUTSKY 79 IPWA PI N TO PI N 12/79
 KUTSKY 80 IPWA PI N TO PI N 1/82

N(1535) REAL PART OF ELASTIC POLE RESIDUE (MEV)

BHANDARI 77 DPWA O USES ETA N CUSP 1/78
 KUTSKY 79 IPWA PI N TO PI N 12/79
 KUTSKY 80 IPWA PI N TO PI N 1/82

N(1535) IMAG PART OF ELASTIC POLE RESIDUE (MEV)

BHANDARI 77 DPWA O USES ETA N CUSP 1/78
 KUTSKY 79 IPWA PI N TO PI N 12/79
 KUTSKY 80 IPWA PI N TO PI N 1/82

N(1535) PARTIAL DECAU MODES

DECAU MASSES
 #N(1535) INTO PI N 139 938
 #N(1535) INTO N PI 1 438 139 139
 #N(1535) INTO N ETA 938 448
 #N(1535) INTO N ERSLON 1232 130
 #N(1535) INTO N RHO 938 769
 #N(1535) INTO GAM N, HELICITY=+1/2 0 939
 #N(1535) INTO GAM N, HELICITY=1/2 0 939
 #N(1535) INTO N RHO, S=1/2, S=WAVE 938 769

N(1535) BRANCHING RATIOS

HENDRY 65 PVUE (PI) 9/66
 MICHAEL 66 PVUE 9/66
 UCHIMAYA 64 PVUE SEE NOTE ON MASS 9/66
 BARREYRE 68 PVUE PI N TO ETAB+C 11/67
 DONNACHI 68 PVUE 8/68
 DELCOURT 69 CNTR 8/69
 AYED 70 IPWA P-N ANAL SOL 8 8/69
 HICKS 73 MPWA GAM-D+ETA 0 9/73
 CRANFORD 75 DPWA PI N PHOTO PDD. 1/78
 LONGACRE 75 IPWA PI N TO 2P1 N 11/75
 AYED 76 IPWA 11/77
 BARBOUR 76 DPWA PI-N PHOTOPROD. 1/78
 BERENDS 77 IPWA PI-N PHOTOPROD. 1/78
 BHANDARI 77 DPWA O USES ETA N CUSP 1/78
 LONGACRE 77 IPWA PI N TO 2P1 N 11/77
 KUTSKY 79 IPWA PI N TO PI N 12/79
 CRANFORD 80 DPWA PI N PHOTO PDD. 12/81
 KUTSKY 80 IPWA PI N TO PI N 1/82

Baryons

N(1535), N(1540)

Data Card Listings

For notation, see key at front of Listings.

Table of baryon data for N(1535) and N(1540). Includes columns for particle ID (R2, R3, etc.), name (N(1535) INTO IN ETA), mass (M), width (G), and various parameters (P2, P3, P4). Includes a section for gamma-decay amplitudes (REFRAC) and a section for photoproduction (PHOTON DECAY AMPLITUDES).

Table of baryon data for N(1540). Includes columns for particle ID (R2, R3, etc.), name (N(1540) INTO IN ETA), mass (M), width (G), and various parameters (P2, P3, P4). Includes a section for photoproduction (PHOTON DECAY AMPLITUDES) and a section for scattering lengths (SCATTERING LENGTHS).

Data Card Listings

For notation, see key at front of Listings.

Baryons

N(1540), N(1650)

Table listing data for N(1540) REAL PART OF POLE POSITION (MEV). Includes columns for ID, description, and values. Rows include 109, 108, and 107.

Table listing data for N(1540) BRANCHING RATIOS. Includes columns for ID, description, and values. Rows include 01, 02, 03, 04, 05, 06, 07, 08.

Table listing data for N(1540) REFERENCES FOR N(1540). Includes columns for reference name and values. Rows include LLNAGRE, ALSG, and DOUBEAU.

Table listing data for N(1540) REAL PART OF POLE POSITION (MEV). Includes columns for ID, description, and values. Rows include 9, 10, 11, 12, 13, 14, 15, 16, 17, 18, 19, 20, 21, 22.

Table listing data for N(1540) -2P1/4 PART OF POLE POSITION (MEV). Includes columns for ID, description, and values. Rows include 01, 02, 03, 04, 05, 06, 07, 08, 09, 10, 11, 12.

Table listing data for N(1540) REAL PART OF ELASTIC POLE RESIDUE (MEV). Includes columns for ID, description, and values. Rows include 01, 02, 03, 04, 05, 06, 07, 08, 09, 10, 11, 12.

N(1650)

Table listing data for N(1650) MASS (MEV). Includes columns for ID, description, and values. Rows include 01, 02, 03, 04, 05, 06, 07, 08, 09, 10, 11, 12, 13, 14, 15, 16, 17, 18, 19, 20, 21, 22, 23, 24, 25, 26, 27, 28, 29, 30, 31, 32, 33, 34, 35, 36, 37, 38, 39, 40, 41, 42, 43, 44, 45, 46, 47, 48, 49, 50, 51, 52, 53, 54, 55, 56, 57, 58, 59, 60, 61, 62, 63, 64, 65, 66, 67, 68, 69, 70, 71, 72, 73, 74, 75, 76, 77, 78, 79, 80, 81, 82, 83, 84, 85, 86, 87, 88, 89, 90, 91, 92, 93, 94, 95, 96, 97, 98, 99, 100.

Table listing data for N(1650) REAL PART OF ELASTIC POLE RESIDUE (MEV). Includes columns for ID, description, and values. Rows include 01, 02, 03, 04, 05, 06, 07, 08, 09, 10, 11, 12, 13, 14, 15, 16, 17, 18, 19, 20, 21, 22, 23, 24, 25, 26, 27, 28, 29, 30, 31, 32, 33, 34, 35, 36, 37, 38, 39, 40, 41, 42, 43, 44, 45, 46, 47, 48, 49, 50, 51, 52, 53, 54, 55, 56, 57, 58, 59, 60, 61, 62, 63, 64, 65, 66, 67, 68, 69, 70, 71, 72, 73, 74, 75, 76, 77, 78, 79, 80, 81, 82, 83, 84, 85, 86, 87, 88, 89, 90, 91, 92, 93, 94, 95, 96, 97, 98, 99, 100.

Table listing data for N(1650) BRANCHING RATIOS. Includes columns for ID, description, and values. Rows include 01, 02, 03, 04, 05, 06, 07, 08, 09, 10, 11, 12, 13, 14, 15, 16, 17, 18, 19, 20, 21, 22, 23, 24, 25, 26, 27, 28, 29, 30, 31, 32, 33, 34, 35, 36, 37, 38, 39, 40, 41, 42, 43, 44, 45, 46, 47, 48, 49, 50, 51, 52, 53, 54, 55, 56, 57, 58, 59, 60, 61, 62, 63, 64, 65, 66, 67, 68, 69, 70, 71, 72, 73, 74, 75, 76, 77, 78, 79, 80, 81, 82, 83, 84, 85, 86, 87, 88, 89, 90, 91, 92, 93, 94, 95, 96, 97, 98, 99, 100.

Table listing data for N(1650) WIDTH (MEV). Includes columns for ID, description, and values. Rows include 01, 02, 03, 04, 05, 06, 07, 08, 09, 10, 11, 12, 13, 14, 15, 16, 17, 18, 19, 20, 21, 22, 23, 24, 25, 26, 27, 28, 29, 30, 31, 32, 33, 34, 35, 36, 37, 38, 39, 40, 41, 42, 43, 44, 45, 46, 47, 48, 49, 50, 51, 52, 53, 54, 55, 56, 57, 58, 59, 60, 61, 62, 63, 64, 65, 66, 67, 68, 69, 70, 71, 72, 73, 74, 75, 76, 77, 78, 79, 80, 81, 82, 83, 84, 85, 86, 87, 88, 89, 90, 91, 92, 93, 94, 95, 96, 97, 98, 99, 100.

Table listing data for N(1650) BRANCHING RATIOS. Includes columns for ID, description, and values. Rows include 01, 02, 03, 04, 05, 06, 07, 08, 09, 10, 11, 12, 13, 14, 15, 16, 17, 18, 19, 20, 21, 22, 23, 24, 25, 26, 27, 28, 29, 30, 31, 32, 33, 34, 35, 36, 37, 38, 39, 40, 41, 42, 43, 44, 45, 46, 47, 48, 49, 50, 51, 52, 53, 54, 55, 56, 57, 58, 59, 60, 61, 62, 63, 64, 65, 66, 67, 68, 69, 70, 71, 72, 73, 74, 75, 76, 77, 78, 79, 80, 81, 82, 83, 84, 85, 86, 87, 88, 89, 90, 91, 92, 93, 94, 95, 96, 97, 98, 99, 100.

Baryons
N(1650), N(1675)

Table with columns for particle ID (e.g., R9, R10), name (e.g., N(121650)), and properties (e.g., spin, parity, mass). Includes entries for various baryon resonances.

N(1675)
FOR DEFINITION OF GAMMA-NUCLEON DECAY AMPLITUDES, SEE MINI-REVIEW PRECEDING THE BARYON LISTINGS.

Table listing decay amplitudes for N(1675) into gamma + p or gamma + n, with columns for amplitude values and phase angles.

Table listing decay amplitudes for N(1675) into pi + p or pi + n, with columns for amplitude values and phase angles.

Table listing decay amplitudes for N(1675) into K + p or K + n, with columns for amplitude values and phase angles.

Table listing decay amplitudes for N(1675) into eta + p or eta + n, with columns for amplitude values and phase angles.

Table listing decay amplitudes for N(1675) into pi + p or pi + n with specific helicity states, with columns for amplitude values and phase angles.

Data Card Listings
For notation, see key at front of Listings.

Table listing baryon resonances with columns for name (e.g., R94), mass (M), spin-parity (J^P), and other properties.

PARAS NOT REFERRED TO IN DATA CARDS
BARYON DECAY TABLES...

N(1675)
1-15 RESONANCE IS WELL ESTABLISHED.

Table listing mass measurements for N(1675) with columns for mass (M), width (Gamma), and references.

Table listing width measurements for N(1675) with columns for width (Gamma), mass (M), and references.

Table listing real parts of pole positions for N(1675) with columns for real part (M), imaginary part (Gamma), and references.

Table listing imaginary parts of pole positions for N(1675) with columns for imaginary part (Gamma), real part (M), and references.

Table listing real parts of elastic pole residues for N(1675) with columns for real part (M), imaginary part (Gamma), and references.

Table listing imaginary parts of elastic pole residues for N(1675) with columns for imaginary part (Gamma), real part (M), and references.

Table listing various baryon resonances with columns for name, mass, spin-parity, and other properties.

Data Card Listings

For notation, see key at front of Listings.

Baryons N(1675)

14 N(121675) PARTIAL DECAY MODES

Table with columns: P1, P2, P3, P4, P5, P6, P7, P8, P9, P10, P11, P12, P13. Rows list decay modes like INTD PI N, INTD N ETA, INTD LAMBDA K, etc.

Table with columns: A2, A3, A4. Rows list decay modes like BARBOJ, ARAJ, DEVENIS, etc.

14 N(121675) BRANCHING RATIOS

Table with columns: P1, P2, P3, P4, P5, P6, P7, P8, P9, P10, P11, P12, P13. Rows list branching ratios for various decay channels.

Table with columns: A2, A3, A4. Rows list branching ratios for various decay channels.

14 N(121675) BRANCHING RATIOS (continued)

Table with columns: P1, P2, P3, P4, P5, P6, P7, P8, P9, P10, P11, P12, P13. Rows list branching ratios for various decay channels.

Table with columns: A2, A3, A4. Rows list branching ratios for various decay channels.

14 N(121675) BRANCHING RATIOS (continued)

Table with columns: P1, P2, P3, P4, P5, P6, P7, P8, P9, P10, P11, P12, P13. Rows list branching ratios for various decay channels.

Table with columns: A2, A3, A4. Rows list branching ratios for various decay channels.

14 N(121675) BRANCHING RATIOS (continued)

Table with columns: P1, P2, P3, P4, P5, P6, P7, P8, P9, P10, P11, P12, P13. Rows list branching ratios for various decay channels.

Table with columns: A2, A3, A4. Rows list branching ratios for various decay channels.

14 N(121675) BRANCHING RATIOS (continued)

Table with columns: P1, P2, P3, P4, P5, P6, P7, P8, P9, P10, P11, P12, P13. Rows list branching ratios for various decay channels.

Table with columns: A2, A3, A4. Rows list branching ratios for various decay channels.

14 N(121675) BRANCHING RATIOS (continued)

Table with columns: P1, P2, P3, P4, P5, P6, P7, P8, P9, P10, P11, P12, P13. Rows list branching ratios for various decay channels.

Table with columns: A2, A3, A4. Rows list branching ratios for various decay channels.

14 N(121675) BRANCHING RATIOS (continued)

Table with columns: P1, P2, P3, P4, P5, P6, P7, P8, P9, P10, P11, P12, P13. Rows list branching ratios for various decay channels.

Table with columns: A2, A3, A4. Rows list branching ratios for various decay channels.

14 N(121675) BRANCHING RATIOS (continued)

Table with columns: P1, P2, P3, P4, P5, P6, P7, P8, P9, P10, P11, P12, P13. Rows list branching ratios for various decay channels.

Table with columns: A2, A3, A4. Rows list branching ratios for various decay channels.

14 N(121675) BRANCHING RATIOS (continued)

Table with columns: P1, P2, P3, P4, P5, P6, P7, P8, P9, P10, P11, P12, P13. Rows list branching ratios for various decay channels.

Table with columns: A2, A3, A4. Rows list branching ratios for various decay channels.

Data Card Listings

For notation, see key at front of Listings.

Baryons N(1680), N(1700)

Table with columns: ID, N(1712(1680)) INTP, GAM, P, HELICITY, I=3/2, CEV=...-1/23, and various particle names like DEVENISH, MOORHUIS, KNIES, METCALF, etc.

Table with columns: ID, N(1721(1680)) INTO, GAM, P, HELICITY, I=1/2, CEV=...-1/23, and various particle names like DEVENISH, MOORHUIS, KNIES, METCALF, etc.

Table with columns: ID, N(1721(1680)) INTP, GAM, N, HELICITY, I=3/2, CEV=...-1/23, and various particle names like DEVENISH, MOORHUIS, KNIES, METCALF, etc.

AVERAGE MEANINGLESS (SCALE FACTOR = 1.0)

REFERENCES FOR N(1721(1680))

Table listing various baryon resonances with columns for name, mass, width, and other properties. Includes entries like BRANDON, HEUSCH, BAREYRE, DONNACHIE, etc.

PAPERS NOT REFERRED TO IN DATA CARDS. CROUCH 65 DESY CONF 244, DERADO 65 AUSTIN CONF 244, JUIE 65 SCL 15 468, MERLO 66 P ROY SOC 239 489, etc.

N(1700) 18 N(1721(1700)) J=3/2-1 I=1/2 D13 THIS RESONANCE IS WELL ESTABLISHED.

Table with columns: ID, N(1721(1700)) MASS (MEV), and various particle names like DONNACHIE, WAGNER, DEANS, etc.

AVERAGE MEANINGLESS (SCALE FACTOR = 1.0)

Table with columns: ID, N(1721(1700)) WIDTH (MEV), and various particle names like DEANS, LANGRENE, LANGRACE, etc.

AVERAGE MEANINGLESS (SCALE FACTOR = 1.0) SEE THE NOTES ACCOMPANYING THE MASSES Q-1010.

Table with columns: ID, N(1721(1700)) REAL PART OF POLE POSITION (MEV), and various particle names like LANGRACE, LANGRACE, etc.

Table with columns: ID, N(1721(1700)) -20% MAG PART OF PCLE POSITION (MEV), and various particle names like LANGRACE, LANGRACE, etc.

Table with columns: ID, N(1721(1700)) REAL PART OF ELASTIC PCLE RESIDUE (MEV), and various particle names like CUTKOSKY, LANGRACE, etc.

Table with columns: ID, N(1721(1700)) -20% MAG PART OF ELASTIC PCLE RESIDUE (MEV), and various particle names like CUTKOSKY, LANGRACE, etc.

Table with columns: ID, N(1721(1700)) I=3/2, CEV=...-1/23, and various particle names like ARAI, CRAWFORD, etc.

Data Card Listings

For notation, see key at front of Listings.

Baryons N(1700)

Table with columns for particle ID, mass, and various physical parameters. Includes entries for particles like BRAUN, CAVALLI, MUSARRA, STRACHNA, WEBB, MERTON, APPLE, etc.

Table titled '20 N(1217) WIDTH MEAS (PROD. EXP.)' containing particle data and decay widths. Includes entries for ALBORELLI, GALLWAY, BENNETTI, RHODE, CIRBA, COPPER, KERNELL, KUZNETSOV, WELLMANN, KHALID, BALLM, ROSEBEC, ELLIS, MORSE, KUSHNOROKI, ERLESTEIN, KARSHON, LAMSA, RONAT, ABE, DAVIDSON, LICHTMAN, ATHERTON, CAVALLI, MUSARRA, CAVALLI, MUSARRA, WEBB, MERTON, APPLE, ERLESTEIN, HEINEN, ROUGE, FERRER, GODOARD, APELOORN, FUKUNAGA.

Table titled '20 N(1217) PARTIAL DECAY MODES (PROD. EXP.)' showing decay channels and branching ratios. Includes entries for CAVALLI, MUSARRA, KERNELL, KUZNETSOV, WELLMANN, AMALDI, DAVIDSON, LICHTMAN, ATHERTON, BARNES, GLOBELI, BRAUNI, CAVALLI, MUSARRA, WEBB, APELOORN, FUKUNAGA.

Table titled '25 N(1217) BRANCHING RATIOS (PROD. EXP.)' providing detailed branching ratios for various decay modes. Includes entries for CAVALLI, MUSARRA, BARNES, KERNELL, WELLMANN, MORSE, ALBORELLI, GALLWAY, BENNETTI, ROUGE, FERRER, GODOARD, APELOORN, FUKUNAGA, BARNES, KERNELL, WELLMANN, MORSE.

Table with columns for particle ID, mass, and various physical parameters. Includes entries for N(1217)0 INTO IN P11 IN P11, N(1217)0 INTO IN NEUTRON P1+P1 P1- P1-, N(1217)0 INTO IN SIG I/FLAMB K1, N(1217)0 INTO IN P1+P1+P1-.

Table with columns for particle ID, mass, and various physical parameters. Includes entries for N(1217)0 INTO IN P1+P1+P1-, N(1217)0 INTO IN P1+P1+P1-, N(1217)0 INTO IN P1+P1+P1-.

Table titled 'REFERENCES FOR N(1217)0 (PROD. EXP.)' listing scientific references and authors. Includes names like KRAMEP, ALEXANDER, ALBORELLI, LEE, ALMEIDA, GALLWAY, KAYAS, BARNES, ROUGE, AMERSON, CIRBA, COPPER, CRENELL, KUZNETSOV, WELLMANN, AMALDI, DAVIDSON, LICHTMAN, ATHERTON, BARNES, GLOBELI, BOESEBEC, LAMSA, RUSHDROE, DELSTEIN, JOHNSON, KARSHON, LAMSA, OH, RONAT.

Table with columns for particle ID, mass, and various physical parameters. Includes entries for ALBORELLI, GALLWAY, BENNETTI, RHODE, CIRBA, COPPER, KERNELL, KUZNETSOV, WELLMANN, AMALDI, DAVIDSON, LICHTMAN, ATHERTON, BARNES, GLOBELI, BRAUNI, CAVALLI, MUSARRA, WEBB, APELOORN, FUKUNAGA, ATHERTON, BARNES, GLOBELI, BOESEBEC, LAMSA, RUSHDROE, DELSTEIN, JOHNSON, KARSHON, LAMSA, OH, RONAT, ABE, ATHERTON, APPLE, HEINEN, ROUGE, MERZ, GAGGER, KULLANDER, BOUQUET, ALBARGNE, GYER, GORDON, LIAI, WOODPH, APELOORN, HARTING, MCELROY, FUKUNAGA, MERLO, LISAUER, ANSOLOV, LINDICKY, CARNEY, DEKREBELT, RUSHDROE, SOTIRIOU.

Data Card Listings

For notation, see key at front of Listings.

Baryons

N(1710), N(1720)

Table with columns for particle ID, mass, width, and various parameters. Includes entries for N(1710) and N(1720) with various production methods and decay modes.

Table with columns for particle ID, mass, width, and various parameters. Includes entries for N(1710) and N(1720) with various production methods and decay modes.

Table with columns for particle ID, mass, width, and various parameters. Includes entries for N(1710) and N(1720) with various production methods and decay modes.

Table with columns for particle ID, mass, width, and various parameters. Includes entries for N(1710) and N(1720) with various production methods and decay modes.

Table with columns for particle ID, mass, width, and various parameters. Includes entries for N(1710) and N(1720) with various production methods and decay modes.

Table with columns for particle ID, mass, width, and various parameters. Includes entries for N(1710) and N(1720) with various production methods and decay modes.

Table with columns for particle ID, mass, width, and various parameters. Includes entries for N(1710) and N(1720) with various production methods and decay modes.

Table with columns for particle ID, mass, width, and various parameters. Includes entries for N(1710) and N(1720) with various production methods and decay modes.

Table with columns for particle ID, mass, width, and various parameters. Includes entries for N(1710) and N(1720) with various production methods and decay modes.

Table with columns for particle ID, mass, width, and various parameters. Includes entries for N(1710) and N(1720) with various production methods and decay modes.

Table with columns for particle ID, mass, width, and various parameters. Includes entries for N(1710) and N(1720) with various production methods and decay modes.

Table with columns for particle ID, mass, width, and various parameters. Includes entries for N(1710) and N(1720) with various production methods and decay modes.

Table with columns for particle ID, mass, width, and various parameters. Includes entries for N(1710) and N(1720) with various production methods and decay modes.

Table with columns for particle ID, mass, width, and various parameters. Includes entries for N(1710) and N(1720) with various production methods and decay modes.

Baryons

N(1720), N(1990)

Data Card Listings

For notation, see key at front of Listings.

Table of baryon data cards for N(1720) and N(1990). Columns include particle ID, mass, width, and various fit parameters like helicity and isospin.

FOR DEFINITION OF GAMMA-NUCLEON DECAUPLY MODES, SEE N(1720) PAPER PRECEDING THE BARYON LIST.

Table of baryon data cards for N(1720) with helicity 1/2 and isospin 1/2. Includes fit parameters and average meanless scale factors.

Table of baryon data cards for N(1720) with helicity 3/2 and isospin 1/2. Includes fit parameters and average meanless scale factors.

Table of baryon data cards for N(1720) with helicity 1/2 and isospin 1/2. Includes fit parameters and average meanless scale factors.

Table of baryon data cards for N(1720) with helicity 3/2 and isospin 1/2. Includes fit parameters and average meanless scale factors.

REFERENCES FOR N(1720): DONNACHI, R G. KIRSOPP, C. LOVELACE... RUSH, J E RUSH... BOTKE, J C BOTKE... DEANS, G R DEANS... AYED, T R KIEV CONK... CARREAS, T O CARREAS... DARRIES, T O DARRIES...

Table of baryon data cards for N(1990) and other related particles. Columns include particle ID, mass, width, and various fit parameters.

PAPER NOT REFERRED TO IN DATA CARDS

Table of baryon data cards for N(1990) with helicity 1/2 and isospin 1/2. Includes fit parameters and average meanless scale factors.

Table of baryon data cards for N(1990) with helicity 3/2 and isospin 1/2. Includes fit parameters and average meanless scale factors.

Table of baryon data cards for N(1990) with helicity 1/2 and isospin 1/2. Includes fit parameters and average meanless scale factors.

Table of baryon data cards for N(1990) with helicity 3/2 and isospin 1/2. Includes fit parameters and average meanless scale factors.

REFERENCES FOR N(1990): DONNACHI, R G. KIRSOPP, C. LOVELACE... RUSH, J E RUSH... BOTKE, J C BOTKE... DEANS, G R DEANS... AYED, T R KIEV CONK... CARREAS, T O CARREAS... DARRIES, T O DARRIES...

Data Card Listings
For notation, see key at front of Listings.

Baryons

N(1990), N(2000), N(2080)

17 N#1/2119001 EMAG PART OF ELASTIC POLE RESIDUE (MEV)
CUTKOSKY 79 1PMA PI N TO PI N 12/79
CUTKOSKY 80 1PMA PI N TO PI N 1/82*

17 N#1/2119001 PARTIAL DECAY MODES
DECAY MASSES
P1 N#1/2119001 INTO PI N 139+ 938
P2 N#1/2119001 INTO N PI PI 938+ 1394 139

17 N#1/2119001 BRANCHING RATIOS
R1 N#1/2119001 INTO (PI N)/TOTAL (P1)
R1 3 (0.29) KIASOPP 58 1PMA PHASE SHIFT ANAL 10/69
R1 7 (0.15) ALMEHED 72 1PMA 2/72

17 N#1/2119001 INTO (ETA N)/TOTAL (P3)
R2 B PANMETRIZAT ION USED COULD BE IN DANGER OF DOUBLE COUNTING
R3 N#1/2119001 FROM GAMMA PROTON TO K LAMBDA
R4 N#1/2119001 FROM GAMMA PROTON TO ETA PROTON

17 N#1/2119001 PHOTON DECAY AMPLITUDE (N#1-1)
FOR DEFINITION OF GAMMA-NUCLEON DECAY AMPLITUDES, SEE MINI-REVIEW PRECEDING THE BARYON LISTINGS.

A1 N#1/2119001 INTO GAM P, HELICITY=1/2 (GEV**1/2) 1/76
A1 (+0.013) BARBOUR 76 0PMA PI N PHOTOPROD. 3/76
A1 (-0.001) CRAWFORD 80 0PMA PI N PHOTOPROD. 12/81*

REFERENCES FOR N#1/2119001
DONNACH1 68 PL 268 161 A DONNACHIE- R G KIRSOPP, C LOWE/LACE (CERN) IJP
R G KIRSOPP (EDIN)

ALMEHED 72 PD 840 157 +LOWLACE (RUTG) IJP
DEANS 72 PD 6 1906 DEANS, JACOBS, LYONS, MONTGOMERY (SOUTH FLA.) IJP
MICKS 73 HW 7 2614 +DEANS, JACOBS, LYONS+ (LEARN) ORN., (SOUTH FLA.) IJP

DEANS 69 PR 177 2623 S R DEANS (UNIV S FLORIDA)
AYED 70 PL 318 598 +BARREVE, VELLE (SACLAY) IJP
LEA 71 PL 218 25 +COMAN, BOSON, MATHI-MORE+ (INHEL, BRISTOL) IJP

N(2000)
D6 N#1/212000, JP#5/2+ 1/1/2
F15

06 N#1/2120001 MASS (MEV)
M 7 (12175.3) ALMEHED 72 1PMA 2/72
M 1 (1930.1) DEANS 72 0PMA GAM P=K L.M.SDL 3 9/73

06 N#1/2120001 WIDTH (MEV)
M 7 (150.1) ALMEHED 72 1PMA GAM P=K L.M.SDL 3 2/72
M 1 (112.1) DEANS 72 0PMA GAM N=K SIG.SDL 3 9/73

06 N#1/2120001 PARTIAL DECAY MODES
DECAY MASSES
P1 N#1/2120001 INTO PI N 139+ 938
P2 N#1/2120001 INTO GAM P, HELICITY=3/2 1115+ 497

06 N#1/2120001 BRANCHING RATIOS
R1 N#1/2120001 INTO (PI N)/TOTAL (P1)
R1 7 (0.25) ALMEHED 72 1PMA 2/72
R1 1 (0.04) DEHNLEP 70 1PMA PI N TO PI N 12/79

06 N#1/2120001 FROM GAMMA PROTON TO K LAMBDA
R2 N#1/2120001 FROM GAMMA PROTON TO K SIGMA
R3 N#1/2120001 FROM PI N TO K SIGMA

ALMEHED 72 NP 840 157 +LOWLACE (PUTG) IJP
DEANS 72 PD 6 1906 DEANS, JACOBS, LYONS, MONTGOMERY (SOUTH FLA.) IJP
LANGBEIN 73 NP 853 251 LANGBEIN, WAGNER (MUMPI) IJP

N(2080)
16 N#1/212080, JP#1/2-1 1/1/2
D13

THERE IS SOME EVIDENCE THAT TWO RESONANCES EXIST IN THIS MASS BETWEEN 1800 AND 2200 MEV. SEE CUTKOSKY 79 AND CUTKOSKY 80. BOTH ARE LISTED HERE AWAITING INDEPENDENT CONFIRMATION OF THE LOWER MASS STATE.

16 N#1/2120801 MASS (MEV)
M 3 (2057.0) DONNACH1 68 8VUE PHASE-SHIFT ANAL 6/69
M 3 (2030.) DONNACH2 68 8VUE PHAS. SHIFT-CERN 10/69

DEANS 70 PD 6 1906 DEANS, JACOBS, LYONS, MONTGOMERY (SOUTH FLA.) IJP
MICKS 73 HW 7 2614 DEANS, JACOBS, LYONS+ (LEARN) ORN., (SOUTH FLA.) IJP
LANGBEIN 73 NP 853 251 LANGBEIN, WAGNER (MUMPI) IJP

Baryons

N(2080), N(2100)

Data Card Listings

For notation, see key at front of Listings.

Table 1: 16 N*(212080) WIDTH (MEV) with columns for mass, width, and various physical parameters.

Table 2: 16 N*(212080) REAL PART OF POLE POSITION (MEV) with columns for real and imaginary parts.

Table 3: 16 N*(212080) -2*IMAG PART OF POLE POSITION (MEV) with columns for real and imaginary parts.

Table 4: 16 N*(212080) REAL PART OF ELASTIC POLE RESIDUE (MEV) with columns for real and imaginary parts.

Table 5: 16 N*(212080) IMAG PART OF ELASTIC POLE RESIDUE (MEV) with columns for real and imaginary parts.

16 N*(212080) PARTIAL DECAY MODES

Table 6: Partial decay modes for N*(212080) listing decay channels and branching ratios.

16 N*(212080) BRANCHING RATIOS

Table 7: Branching ratios for N*(212080) into various decay channels.

16 N*(212080) BRANCHING RATIOS (continued)

Table 8: Further branching ratios for N*(212080) including resonance parameters.

16 N*(212080) BRANCHING RATIOS (continued)

Table 9: Final branching ratios for N*(212080) including resonance parameters.

Table 10: 16 N*(212080) PHOTON DECAY AMPLITUDES (-1/2) with columns for FOP definition and amplitudes.

Table 11: REFERENCES FOR N*(212080) listing various scientific papers and authors.

Table 12: REFERENCES FOR N*(212080) (continued) listing additional scientific references.

Table 13: REFERENCES FOR N*(212080) (continued) listing further scientific references.

Table 14: REFERENCES FOR N*(212080) (continued) listing more scientific references.

16 N*(212100) MASS (MEV)

Table 15: Mass measurements for N*(212100) from various experiments.

16 N*(212100) WIDTH (MEV)

Table 16: Width measurements for N*(212100) from various experiments.

16 N*(212100) REAL PART OF POLE POSITION (MEV)

Table 17: Real part of pole position for N*(212100) from various experiments.

16 N*(212100) IMAG PART OF ELASTIC POLE RESIDUE (MEV)

Table 18: Imaginary part of elastic pole residue for N*(212100) from various experiments.

16 N*(212100) PARTIAL DECAY MODES

Table 19: Partial decay modes for N*(212100) listing decay channels and branching ratios.

N(2100) ANY STRUCTURE IN THIS MASS ABOVE 1000 MEV IS LISTED HERE

Data Card Listings
For notation, see key at front of Listings.

Baryons
N(2100), N(2190)

04 N*(212100) BRANCHING RATIOS
R1 N*(212100) INTO EPI N(TOTAL)
R1 10.53 ALMEHD 72 IPMA 1P11 2/72
R1 10.151 AYED 76 IPMA 1P1 N TO PI N 11/77
R1 0.09 0.05 HOEHLER 79 IPMA 1P1 N TO PI N 12/79
R1 0.18 0.08 CUTKOSKY 80 IPMA 1P1 N TO PI N 1/82*

REFERENCES FOR N*(212100)
ROYCHOUD 71 NP 827 125 R K ROYCHODHURJ, B H PANDEEN (DURHI)JP
ALMEHD 72 NP 840 157 *LOVELACE (LUND, ROTG)JP
76 CEH-N(1521) AYED (THES15) (SCAL11)JP
HOEHLER 79 HANDBOOK OF PI-N SCATTERING, PHYSIC DATEN VOL.12-1 **KAISER, KOCH, PIETARINEN (KARLSRUHE I)JP
ALSD 80 TORONTO CONF 3 R, KDC-H (KARLSRUHE I)JP
CUTKOSKY 80 TORONTO CONF 19 *FORSYTH, BABCOCK, KELLY, HENRICK (CARNABLI)JP
SAXON 80 NP B162 522 *BAKER, MELL, BLI SZEIT, BLODMORITZ, HUEL-BRISII)JP

MA 76 PRD 13 3027 E. M.G.L. SHAW (OREG-HUCI)JP
PAPERS NOT REFERRED TO IN DATA CARDS

N(2100) 132 N*(212100), Jp=1/2+1-1/2 P-11
132 N*(212100) MASS (MEV)
M 2950. 70. HOEHLER 79 IPMA PI N TO PI N 1/82*
M 2125. 75. CUTKOSKY 80 IPMA PI N TO PI N 1/82*

AVERAGE MEANINGLESS ISCALE FACTOR = 1.01

132 N*(212100) WIDTH (MEV)
M 220. 30. HOEHLER 79 IPMA PI N TO PI N 1/82*
M 260. 100. CUTKOSKY 80 IPMA PI N TO PI N 1/82*

132 N*(212100) REAL PART OF POLE POSITION (MEV)
RE 2120. 40. CUTKOSKY 80 IPMA PI N TO PI N 1/82*

-2*IMAG PART OF POLE POSITION (MEV)
IM 240. 80. CUTKOSKY 80 IPMA PI N TO PI N 1/82*

132 N*(212100) REAL PART OF ELASTIC POLE RESIDUE (MEV)
RER 11. 7. CUTKOSKY 80 IPMA PI N TO PI N 1/82*

132 N*(212100) IMAG PART OF ELASTIC POLE RESIDUE (MEV)
IMR 8. 6. CUTKOSKY 80 IPMA PI N TO PI N 1/82*

132 N*(212100) PARTIAL DECAY MODES
P1 N*(212100) INTO PI N DECAY MASSES 139+ 935

132 N*(212100) BRANCHING RATIOS
R1 N*(212100) INTO EPI N(TOTAL) (PI)
R1 0.10 0.09 HOEHLER 79 IPMA PI N TO PI N 1/82*
R1 0.12 0.02 CUTKOSKY 80 IPMA PI N TO PI N 1/82*

REFERENCES FOR N*(212100)
HOEHLER 79 HANDBOOK OF PI-N SCATTERING, PHYSIC DATEN VOL.12-1 **KAISER, KOCH, PIETARINEN (KARLSRUHE I)JP
ALSD 80 TORONTO CONF 3 R, KDC-H (KARLSRUHE I)JP
CUTKOSKY 80 TORONTO CONF 19 *FORSYTH, BABCOCK, KELLY, HENRICK (CARNABLI)JP

2100 MEV REGION - PRODUCTION EXPERIMENTS
114 N*(212100), Jp= 1-1/2 PRODUCTION EXPERIMENTS
RESONANCE-LIKE BUMP OBSERVED IN PP TO PI N (PI*) AT GERN 15R (DE KERRET 76). THE ENHANCEMENT SHOWS UP MORE CLEARLY WHEN EVENTS CORRESPONDING TO TRANSFERAL DECATS OF THE (N PI*) SYSTEM ARE SELECTED, CONTRARY TO WHAT WOULD BE EXPECTED FOR A DIFFRACTIVE-LIKE EFFECT.

114 N*(212100) MASS (MEV) (PROD. EXPERIMENTS)
M (2100.) DE KERRET 76 15R * PI N PI*+N45GLY 1/78

114 N*(212100) PARTIAL DECAY MODES (PROD. EXP.)
P1 N*(212100) INTO PI N DECAY MASSES 139+ 493

REFERENCES FOR N*(212100) PKCD. EXPERIMENTS
DEKERRET 76 PL 638 477. #63 *MAY, REGLER, GRANDT* (GERN+HAMB+IPP+VIEN)

N(2190) 71 N*(212190), Jp=7/2-1-1/2 G12
THIS RESONANCE IS WELL ESTABLISHED.

71 N*(212190) MASS (MEV)
M (2190.0) DIDDEHS 63 CNTR PI- P TOTAL 1/77
M (2190.0) HOHLER 64 PVUE DATA + DISR PEEL 10/71
M (2190.0) APPROX YOKOSAWA 66 CNTR PE- P DSIG + PUL 7/66
M (2205.0) APPROX DONNACHI 68 PVUE PHASE-SHIFT ANAL 6/68
M (2200.0) LEA 69 CNTR PI-P ELASTIC 8/69
M 2180. 25. ANDERSON 70 HNS - PI- P TO PI- HNS 2/71
M (2195.0) AYED 70 IPMA 1P1 N TO PI N 1/77
M (2200.0) HULL 70 IPMA SMALL ANGLE PI-P 1/77
M (2180.-1) 150.01 AMALDI 71 CNTR P P AT 24 DEVI 10/71
M (2180.) BRANKEN 71 DPMA 3/72
M (2200.) ROYCHOUD 71 DPMA 3/72
M (2200.) ALMEHD 72 IPMA 2/72
M (2190.) DIT 72 NPRA O PI-P BKWD ELSTIC 2/73
M (2208.) HICKS 73 MPMA GAM P-ETA P 9/73
M 1 DM N STATES FROM TABLE VII OF HICKS 73 ARE INCLUDED IN LISTINGS 9/73
M 1 M AND M ARE FURTHER SOLUTION C2, R= SORTIG)M WITH G FROM TABLE VII. 9/73
M (2208.) 129. ABE 74 * Pap-Spax, JCBN PK 6/75
M (2181.-) AYED 76 IPMA 1/77
M (2177.) BARBUDR 78 DPMA O PI-N PHOTOPROD. 3/79
M (2140.) HENDRY 78 IPMA PI N TO PI N 12/79
M (2140.) BAKER 79 DPMA O PI- P TO ETA N 12/79
M (2150.) 100. CUTKOSKY 79 IPMA PI N TO PI N 12/79
M (2162.) 12. HOEHLER 79 IPMA PI N TO PI N 12/79
M (2098.0) CRAWFORD 80 DPMA PI N PHOTOPROD. 12/81
M 2200. 70. CUTKOSKY 80 IPMA PI N TO PI N 1/82*
M (2180.-) SAKON 80 DPMA O PI- P TO K LAM 12/79
M AVERAGE MEANINGLESS ISCALE FACTOR = 1.01

71 N*(212190) WIDTH (MEV)
M (220.0) DIDDEHS 63 CNTR 7/66
M (220.0) HOHLER 64 PVUE 7/66
M (220.0) APPROX YOKOSAWA 66 CNTR 6/66
M (224.0) DONNACHI 68 PVUE 6/66
M 275. 70. ANDERSON 70 HNS - PI- P TO PI- HNS 2/71
M (225.0) HULL 70 IPMA SMALL ANGLE PI-P 1/77
M (2190.) ALMEHD 72 IPMA 2/72
M (2180.) HICKS 73 MPMA GAM P-ETA P 9/73
M (2143.) AYED 76 IPMA 1/77
M (222.) BARBUDR 78 DPMA O PI-N PHOTOPROD. 3/79
M (220.) HENDRY 78 IPMA PI N TO PI N 12/79
M (2140.) BAKER 79 DPMA O PI- P TO ETA N 12/79
M (2150.) 100. CUTKOSKY 79 IPMA PI N TO PI N 12/79
M (2162.) 12. HOEHLER 79 IPMA PI N TO PI N 12/79
M (2098.0) CRAWFORD 80 DPMA PI N PHOTOPROD. 12/81
M (2180.-) SAKON 80 DPMA O PI- P TO K LAM 12/79
M AVERAGE MEANINGLESS ISCALE FACTOR = 1.71
SEE THE NOTES ACCOMPANYING THE MASSES QUOTED.

71 N*(212190) REAL PART OF POLE POSITION (MEV)
RE (2111.) CUTKOSKY 79 IPMA PI N TO PI N 12/79
RE 2100. 50. CUTKOSKY 80 IPMA PI N TO PI N 1/82*

-2*IMAG PART OF POLE POSITION (MEV)
IM (308.) CUTKOSKY 79 IPMA PI N TO PI N 12/79
IM 400. 160. CUTKOSKY 80 IPMA PI N TO PI N 1/82*

71 N*(212190) REAL PART OF ELASTIC POLE RESIDUE (MEV)
RER (24.) CUTKOSKY 79 IPMA PI N TO PI N 12/79
RER 22. 14. CUTKOSKY 80 IPMA PI N TO PI N 1/82*

71 N*(212190) IMAG PART OF ELASTIC POLE RESIDUE (MEV)
IMR (-12.) CUTKOSKY 79 IPMA PI N TO PI N 12/79
IMR -11. 20. CUTKOSKY 80 IPMA PI N TO PI N 1/82*

11 N*(212190) PARTIAL DECAY MODES
P1 N*(212190) INTO PI N DECAY MASSES 139+ 935
P2 N*(212190) INTO LAMBDA K 115+ 493
P3 N*(212190) INTO N PI N 938+ 139+ 139
P4 N*(212190) INTO GAM P, HELICITY 3/2 0+ 939
P5 N*(212190) INTO GAM P, HELICITY 1/2 0+ 939
P6 N*(212190) INTO GAM P, HELICITY 3/2 0+ 939
P7 N*(212190) INTO GAM P, HELICITY 1/2 0+ 939
P8 N*(212190) INTO ETA N 548+ 939
P9 N*(212190) INTO SIGMA K 403+ 1189

Baryons
N(2190), N(2200)

Data Card Listings
For notation, see key at front of Listings.

Table with columns for particle ID, mass, width, and various parameters. Includes entries for N(212190) and N(212200).

71 N(212190) PHOTON DECAY AMPLITUDE=1/21
FOR DEFINITION OF GAMMA-NUCLEON DECAY AMPLITUDES, SEE MINI-REVIEW PRECEDING THE BARYON LISTINGS.

Table of photon decay amplitudes for N(212190) and N(212200) with columns for particle ID, mass, width, and amplitude values.

REFERENCES FOR N(212190)

Table of references for N(212190) listing authors, journal names, and page numbers.

PAPERS NOT REFERRED TO IN DATA CARDS

Table of references not referred to in data cards, listing authors and journal information.

N(2200) 05 N(212200) JP=5/2-1 1/2 D15

Table of data for N(2200) with columns for particle ID, mass, width, and various parameters.

Table with columns for particle ID, mass, width, and various parameters.

Table with columns for particle ID, mass, width, and various parameters.

Table with columns for particle ID, mass, width, and various parameters.

Table with columns for particle ID, mass, width, and various parameters.

Table with columns for particle ID, mass, width, and various parameters.

05 N(212200) BRANCHING RATIOS

Table of branching ratios for N(212200) with columns for particle ID, mass, width, and ratio values.

REFERENCES FOR N(212200)

Table of references for N(212200) listing authors, journal names, and page numbers.

2200 MEV REGION - PRODUCTION EXPERIMENTS

111 N(212200, JP=1/2- PRODUCTION EXPERIMENTS
WE LIST HERE BUMPS OBSERVED IN THE RANGE 1900-2500 MEV.

Table of production experiments for N(212200) with columns for particle ID, mass, width, and various parameters.

Data Card Listings

For notation, see key at front of Listings.

Baryons

N(2200), N(2220), N(2250)

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111 N(22200) WIDTH (MEV) (PROD. EXP.)
W 125. 70. APPLE 77 SPEC + P P TO R EP DIA 1/78
W 75. 50. APPLE 77 SPEC + P P TO R EP DIA 1/78
W D 45 160. 30. SUGAHARA 79 HBC +D P1-P2 AT 4.5 GEV 12/79
W E 34 20. 30. SUGAHARA 79 HBC +D P1-P2 AT 4.5 GEV 12/79
M AVERAGE MEANINGLESS (SCALE FACTOR = 1.9)
SEE THE NOTES ACCOMPANYING THE MASSES QUOTED ABOVE.

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REFERENCES FOR N(22200)

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AMR101 71 PL 348 495 +RIANCASTELL,ROSSO,MAFFIEEA (SANTICORN)
APPLE 77 LMC 18 167 +ASH,CHENG,COYNE,GROSSMAN (PRIN+PAVIA)
SUGAHARA 79 UC 524 373 +SUZUKI,FOKUMI,KABEKI,ICHIMI,UCHIHA+ (KEK)

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N(2220)

90 N(222200) J^{PC} = 2⁺ 1⁺ 1/2⁻

H₁₉

THIS RESONANCE IS WELL ESTABLISHED.

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90 N(222200) MASS (MEV)
M (2200.) APPROX. BUSZA 6T ESPK LEG.POLYN. EVAL 2/71
M 6 (2221.0) AYED 70 IPWA 1/77
M b FROM ENER. DEP. FIT OF ARGAND DIAGRAM
M 12245.0 J HULL 70 MPWA SMALL ANGLE PI-P 1/77
M 12245.0 J HULL 70 MPWA SMALL ANGLE PI-P 1/77
M 2300. 100. HENDY 78 MPWA PI N TO PI N 12/79
M (2295.1) BAKER 79 DPWA 0 PI- P TO ETA N 12/79
M (2295.1) CUTKOSKY 79 IPWA PI N TO PI N 12/79
M 2205. 10. HUEHLER 79 IPWA PI N TO PI N 12/79
M 2239. 40. CUTKOSKY 80 IPWA PI N TO PI N 1/82*
M AVERAGE MEANINGLESS (SCALE FACTOR = 1.0)

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90 N(222200) WIDTH (MEV)
W 6 (228.0) AYED 70 IPWA 1/77
W (232.0) HULL 70 MPWA SMALL ANGLE PI-P 1/77
W (247.) AYED 76 IPWA 1/77
W 450. HENDY 78 MPWA PI N TO PI N 12/79
W (450.1) HENDY 78 MPWA PI N TO PI N 12/79
W (450.1) CUTKOSKY 79 IPWA PI N TO PI N 12/79
W 365. 30. HUEHLER 79 IPWA PI N TO PI N 12/79
W 500. 150. CUTKOSKY 80 IPWA PI N TO PI N 1/82*
M AVERAGE MEANINGLESS (SCALE FACTOR = 1.0)

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90 N(222200) REAL PART OF POLE POSITION (MEV)
RE (2189.) 90. CUTKOSKY 79 IPWA PI N TO PI N 12/79
RE (2160.) 90. CUTKOSKY 80 IPWA PI N TO PI N 1/82*

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90 N(222200) -2MAG PART OF POLE POSITION (MEV)
IM (1400.) 90. CUTKOSKY 79 IPWA PI N TO PI N 12/79
IM (480.) 100. CUTKOSKY 80 IPWA PI N TO PI N 1/82*

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90 N(222200) REAL PART OF ELASTIC POLE RESIDUE (MEV)
REP (137.) CUTKOSKY 79 IPWA PI N TO PI N 12/79
REP 32. 20. CUTKOSKY 80 IPWA PI N TO PI N 1/82*

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90 N(222200) IMAG PART OF ELASTIC POLE RESIDUE (MEV)
IMR (1-21.) CUTKOSKY 79 IPWA PI N TO PI N 12/79
IMR -57. 20. CUTKOSKY 80 IPWA PI N TO PI N 1/82*

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90 N(222200) PARTIAL DECAY MODES
P1 N(222200) INTO PI N 3394 938
P2 N(222200) INTO N ETA 5394 548
P3 N(222200) INTO LAMBDA K 11154 407

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90 N(222200) BRANCHING RATIOS
R1 N(222200) INTO (PI N) TOTAL (PI)
R1 6 (0.345) AYED 70 IPWA 1/77
R1 (0.345) HULL 70 MPWA SMALL ANGLE PI-P 1/77
R1 10.2.1 AYED 76 IPWA 1/77
R1 0.1. HENDY 78 MPWA PI N TO PI N 12/79
R1 (0.2.0) HENDY 78 MPWA PI N TO PI N 12/79
R1 (0.2.0) CUTKOSKY 79 IPWA PI N TO PI N 12/79
R1 0.1. HUEHLER 79 IPWA PI N TO PI N 12/79
R1 0.15 0.03 CUTKOSKY 80 IPWA PI N TO PI N 1/82*
R1 AVERAGE MEANINGLESS (SCALE FACTOR = 1.1)

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R2 N(222200) FROM PI N TO K LAMBDA SORT(EP#P3) 12/79
R2 NOT SEEN SAKON 80 DPWA 0 PI- P TO K LAM 12/79
R3 N(222200) FROM PI N TO ETA N SORT(EP#P2) 12/79
R3 10.034 BAKER 79 DPWA 0 PI- P TO ETA N 12/79

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REFERENCES FOR N(222200)

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BUSZA 6T UC 524 371 +DAVIS,DUFF,HEYMAN,HEMION + (LOND+LOND)
AYED 70 KJEF CONF M AYED,P BAREYFF, G HILLET (SACLAT)
HULL 70 PR 02 1763 J HULL, R CACKER (TSU)

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AYED 76 CEA-N-1921 ARED (MESI) (SACLAT)
HENDY 78 PL 41 222 ARCHIBALD W. HENDY (INDO+ELL)
ALSO 80 TORONTO CONF 113 ARCHIBALD W. HENDY (INDO)

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BAKER 79 NP 8156 93 +BROWN,CLARK,DAVIES,DEPAIGET,EVANS+ (RH)ELL
CUTKOSKY 79 PR 20 2839 +FORSYTH,HENDRICK,KELLY (CAN+ELL)
HUEHLER 79 HANNOVER DE PI-N SCATTERING. PMSIC DISEN VOL.12-1 +KATSER,KOCH,PIETARINEN +KARLSRUHE JJP
R,KOCH (KARLSRUHE JJP)
ALSO 80 TORONTO CONF 3

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CUTKOSKY 80 TORONTO CONF 39 +FORSYTH,BABCOCK,KELLY,HENDRICK (CAN+ELL)
SAKON 80 NP 8162 522 +BAKER,BELL,BLISSETT,BLODDGORTH+(FEL+BTS)JJP

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PAPERS NOT REFERRED TO IN DATA CAPS

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AYED 70 PL 318 598 +BAREYRE-VILLET (SACLAT)
M 76 PR 13 3327 E. M.L.G. L. SWAN (EDORE+UC)JJP

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N(2250)

113 N(222500) J^{PC} = 2⁺ 1⁺ 1/2⁻

G₁₉

THIS RESONANCE IS WELL ESTABLISHED.

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113 N(222500) MASS (MEV)
M (2133.) AYED 76 IPWA 1/78
M 2200. 100. HENDY 78 MPWA PI N TO PI N 12/79
M (2200.) CUTKOSKY 79 IPWA PI N TO PI N 12/79
M 2268. 15. HUEHLER 79 IPWA PI N TO PI N 12/79
M (2250.) CUTKOSKY 80 IPWA PI N TO PI N 1/82*
M AVERAGE MEANINGLESS (SCALE FACTOR = 1.0)

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113 N(222500) WIDTH (MEV)
W (101.) AYED 76 IPWA 1/78
W 350. 100. HENDY 78 MPWA PI N TO PI N 12/79
W (130.1) 40. CUTKOSKY 79 IPWA PI N TO PI N 12/79
W 300. 40. HUEHLER 79 IPWA PI N TO PI N 12/79
W (400. 120. CUTKOSKY 80 IPWA PI N TO PI N 1/82*
M AVERAGE MEANINGLESS (SCALE FACTOR = 1.0)

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113 N(222500) REAL PART OF POLE POSITION (MEV)
RE (2129.) 90. CUTKOSKY 79 IPWA PI N TO PI N 12/79
RE 2150. 50. CUTKOSKY 80 IPWA PI N TO PI N 1/82*

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113 N(222500) -2MAG PART OF POLE POSITION (MEV)
IM (190.) 90. CUTKOSKY 79 IPWA PI N TO PI N 12/79
IM (50.) 100. CUTKOSKY 80 IPWA PI N TO PI N 1/82*

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113 N(222500) REAL PART OF ELASTIC POLE RESIDUE (MEV)
REP (115.) 90. CUTKOSKY 79 IPWA PI N TO PI N 12/79
REP 15. 7. CUTKOSKY 80 IPWA PI N TO PI N 1/82*

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113 N(222500) IMAG PART OF ELASTIC POLE RESIDUE (MEV)
IMR (1-17.) 90. CUTKOSKY 79 IPWA PI N TO PI N 12/79
IMR -15. 9. CUTKOSKY 80 IPWA PI N TO PI N 1/82*

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113 N(222500) PARTIAL DECAY MODES
P1 N(222500) INTO PI N 1394 538
P2 N(222500) INTO LAMBDA K 11154 407
P3 N(222500) INTO ETA N 4394 548

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113 N(222500) BRANCHING RATIOS
R1 N(222500) INTO (PI N) TOTAL (PI)
R1 17.091 AYED 76 IPWA 1/78
R1 7.09 HENDY 78 MPWA PI N TO PI N 12/79
R1 13.1.0.0 CUTKOSKY 79 IPWA PI N TO PI N 12/79
R1 0.12 0.02 HUEHLER 79 IPWA PI N TO PI N 12/79
R1 0.10 0.02 CUTKOSKY 80 IPWA PI N TO PI N 1/82*
R1 AVERAGE MEANINGLESS (SCALE FACTOR = 1.01)

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R2 N(222500) FROM PI N TO K LAMBDA SORT(EP#P2) 12/79
R2 NOT SEEN SAKON 80 DPWA 0 PI- P TO K LAM 12/79
R3 N(222500) FROM PI N TO ETA N SORT(EP#P3) 12/79
R3 (5.041) BAKER 79 DPWA 0 PI- P TO ETA N 12/79

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REFERENCES FOR N(222500)

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AYED 76 CEA-N-1921 ARED (MESI) (SACLAT)
HENDY 78 PL 41 222 ARCHIBALD W. HENDY (INDO+ELL)
ALSO 80 TORONTO CONF 113 ARCHIBALD W. HENDY (INDO)

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BAKER 79 NP 8156 93 +BROWN,CLARK,DAVIES,DEPAIGET,EVANS+ (RH)ELL
CUTKOSKY 79 PR 20 2839 +FORSYTH,HENDRICK,KELLY (CAN+ELL)
HUEHLER 79 HANNOVER DE PI-N SCATTERING. PMSIC DISEN VOL.12-1 +KATSER,KOCH,PIETARINEN +KARLSRUHE JJP
R,KOCH (KARLSRUHE JJP)
ALSO 80 TORONTO CONF 3

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CUTKOSKY 80 TORONTO CONF 39 +FORSYTH,BABCOCK,KELLY,HENDRICK (CAN+ELL)
SAKON 80 NP 8162 522 +BAKER,BELL,BLISSETT,BLODDGORTH+(FEL+BTS)JJP

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Baryons

N(2600), N(2650), N(2700), N(2800)

Data Card Listings
For notation, see key at front of Listings.

N(2600) 120 N(1/2)2600, JP=1(2-) 1-1/2 **I₁₁**

120 N(1/2)2600 MASS (MEV)

M	2790.	100.	HENDRY	78	PPWA	PI N TO PI N	12/79
M	2577.	50.	HOEHLER	79	IPWA	PI N TO PI N	12/79

AVERAGE MEANINGLESS (SCALE FACTOR = 1.11)

120 N(1/2)2600 WIDTH (MEV)

M	900.	100.	HENDRY	78	PPWA	PI N TO PI N	12/79
M	400.	100.	HOEHLER	79	IPWA	PI N TO PI N	12/79

AVERAGE MEANINGLESS (SCALE FACTOR = 3.51)

120 N(1/2)2600 PARTIAL DECAY MODES

PI	N(1/2)2600	INTO	PI N	DECAY MASSES	139+ 93P
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120 N(1/2)2600 BRANCHING RATIOS

R1	N(1/2)2600	INTO	PI N	PI TOTAL	(P1)	12/79	
R1	0.05	0.01	HENDRY	78	PPWA	PI N TO PI N	12/79
R1	0.05	0.01	HOEHLER	79	IPWA	PI N TO PI N	12/79

AVERAGE MEANINGLESS (SCALE FACTOR = 1.31)

REFERENCES FOR N(1/2)2600

HENDRY 78 PRL 41 222 ARCHIBALD W. HENDRY (IND+LBL)JJP
 ALSO 80 TORONTO CONF 113 ARCHIBALD W. HENDRY (IND)JJP
 HOEHLER 79 HANDBOOK OF PI-N SCATTERING, PHYSIK DATEN VOL.12-1
 *KRAISER,ROCH,PIETARINEN /KARLSRUHE JJP
 ALSO 80 TORONTO CONF 3 P.ROCH

2650 MEV REGION - MISCELLANEOUS EXPERIMENTS

72 N(1/2)2650, 1 1/2 PRODUCTION EXPERIMENTS

ROYCHODHURY 71 CLAIM F15(2400) AND GE1(2400) TO BE POSSIBLE RESONANCES. BRANSDEN 71 FIND THE POSSIBLE RESONANT CANDIDATES S11(2500) AND M(1/2)2500. RECENT PI N PAPER ESTABLISH THE EXISTENCE OF A JP=1/2- STATE IN THIS REGION, BUT THE POSSIBILITY THAT THERE ARE ALSO OTHER STATES REMAINS. SEE THE MMT-REVIEW PRECEDING THE N AND DELTA LISTINGS.

72 N(1/2)2650 MASS (MEV) (PROD. EXP.)

M	2790.01	64	CNTR	PI PHOTOPROD.	
M	2660.01	HOHLER	64	PVUE DATA + DISP REL	
M	2690.01	APPROX	WAMLIG	64	CPSE O PI-P CN EX
M	2833.01		BARGER	66	FIT TOTAL + CN EX
M	2649.0	10.0	CITRON	66	CNTR PI+- P TOTAL

72 N(1/2)2650 WIDTH (MEV) (PROD. EXP.)

M	6100.01	64	CNTR	ALVAREZ	64	CNTR	7/66
M	6200.01	HOHLER	64	PVUE	7/66		
M	6425.01	BARGER	66	FIT	11/67		
M	363.0	20.0	CITRON	66	CNTR	7/66	

72 N(1/2)2650 PARTIAL DECAY MODES (PROD. EXP.)

P1	N(1/2)2650	INTO	PI N	DECAY MASSES	139+ 93B
P2	N(1/2)2650	INTO	LAMBDA K	1135+ 497	
P3	N(1/2)2650	INTO	PI P	93B+ 139+ 139	

72 N(1/2)2650 BRANCHING RATIOS (PROD. EXP.)

R1	N(1/2)2650	INTO	PI N	PI TOTAL	(P1)	12/79
R1	ONLY (J=1/2)PI N	MEASURED	FOR THIS STATE			
R1	B	10+450 (0.018)	BARGER	66	PVUE TOTAL + CN EX	11/67
R1	W	0+26 (0.028)	CITRON	66	CNTR TOTAL CROSS-SEC.	11/67
R1	B	(0.30)	BARGER	66	PVUE USES KORNYANOS67	11/67
R1	B	USES REGGE AMP+RESON. TO CALCULATE DEF. CROSS SECTIONS AT 180 DEGREE				
R1	B	FOR CRITICISM OF THIS METHOD, SEE DOLEN 68.				
R1	D	(0.24)	DIMKEN	67	PVUE USES KORNYANOS66	11/67
R1	D	USES ONLY RESONANCES TO CALCULATE DEF. CROSS SECTIONS AT 180 DEGREE				
R1	D	(0.06)	KORNYANOS	67	CNTR PI-P AT 180 DEG.	11/67

REFERENCE. FOR N(1/2)2650 (PROD. EXP.)

ALVAREZ 64 PRL 32 710 *BAR-YAM, KERN, LUCKEY, OSBORNE. * (MIT+CEA)
 HOHLER 64 PL 12 144 * HOHLER, J. GEISLER (KARLSRUHE) I
 WAMLIG 64 PRL 13 103 *MANWELL, SODICKSON, FACKLER, WARD. * (MIT)
 BARGER 66 PR 151 1123 V BARGER, M OLSSON (WISC)
 CITRON 66 PR 144 1101 *GALBRITHE, POCIAL, LEONTIC, PHILLIPS. * (IUN) I
 BARGER 67 PR 155 1792 V BARGER, O CLINE (WISC) P
 DIMKEN 67 PRL 18 798 F W DIMKEN (WICH)
 KORNYANOS 67 PR 144 1651 KORNYANOS, KRISCH, OFALLON. * (MICH) I P

PAPER NOT REFERRED TO IN DATA CARDS

RAACKE 67 NK 514 761 J RAACKE, M VVERT (KARLSRUHE-NPSBY)J-L
 DOLEN 68 PR 166 1768 R DOLEN, J HORN, C SCHMID (ICIT)
 WAMLIG 68 PR 166 1515 W A WAMLIG, J MANWELL (MIT, PISA)
 FINAL VERSION OF DATA USED IN WAMLIG 64. IN CONJUNCTION WITH
 CITRON 66 134L CROSS SECTIONS, THIS CHARGE EXCHANGE DATA GIVES
 COMPLEX ELASTIC SCATTERING AMPLITUDE AT 0 DEGREE.

BRANSDEN 71 NP 874 511 *ODDEN (ORNL)JJP
 ALSO 70 NP 816 461 *ROYCHODHURY, PERRIN, BRANSDEN (ORNL)JJP
 ROYCHODH 71 NP 827 125 P K ROYCHODHURY, R H BRANSDEN (ORNL)JJP

N(2700)

121 N(1/2)2700, JP=1(2-) 1-1/2

K₁₁₃

121 N(1/2)2700 MASS (MEV)

M	3000.	100.	HENDRY	78	PPWA	PI N TO PI N	12/79
M	2612.	45.	HOEHLER	79	IPWA	PI N TO PI N	12/79

AVERAGE MEANINGLESS (SCALE FACTOR = 3.51)

121 N(1/2)2700 WIDTH (MEV)

M	900.	150.	HENDRY	78	PPWA	PI N TO PI N	12/79
M	150.	50.	HOEHLER	79	IPWA	PI N TO PI N	12/79

AVERAGE MEANINGLESS (SCALE FACTOR = 3.51)

121 N(1/2)2700 PARTIAL DECAY MODES

PI	N(1/2)2700	INTO	PI N	DECAY MASSES	139+ 93B
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121 N(1/2)2700 BRANCHING RATIOS

R1	N(1/2)2700	INTO	PI N	PI TOTAL	(P1)	12/79	
R1	0.07	0.02	HENDRY	78	PPWA	PI N TO PI N	12/79
R1	0.04	0.01	HOEHLER	79	IPWA	PI N TO PI N	12/79

AVERAGE MEANINGLESS (SCALE FACTOR = 1.31)

REFERENCES FOR N(1/2)2700

HENDRY 78 PRL 41 222 ARCHIBALD W. HENDRY (IND+LBL)JJP
 ALSO 80 TORONTO CONF 113 ARCHIBALD W. HENDRY (IND)JJP
 HOEHLER 79 HANDBOOK OF PI-N SCATTERING, PHYSIK DATEN VOL.12-1
 *KRAISER,ROCH,PIETARINEN /KARLSRUHE JJP
 ALSO 80 TORONTO CONF 3 P.ROCH

N(2800)

122 N(1/2)2800, JP=0(2-) 1-1/2

G₁₉

122 N(1/2)2800 MASS (MEV)

M	2792.	100.	HOEHLER	79	IPWA	PI N TO PI N	12/79
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122 N(1/2)2800 WIDTH (MEV)

M	240.	100.	HOEHLER	79	IPWA	PI N TO PI N	12/79
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122 N(1/2)2800 PARTIAL DECAY MODES

PI	N(1/2)2800	INTO	PI N	DECAY MASSES	139+ 93B
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122 N(1/2)2800 BRANCHING RATIOS

R1	N(1/2)2800	INTO	PI N	PI TOTAL	(P1)	12/79	
R1	0.02	0.015	HOEHLER	79	IPWA	PI N TO PI N	12/79

REFERENCES FOR N(1/2)2800

HOEHLER 79 HANDBOOK OF PI-N SCATTERING, PHYSIK DATEN VOL.12-1
 *KRAISER,ROCH,PIETARINEN /KARLSRUHE JJP
 ALSO 80 TORONTO CONF 3 P.ROCH

Data Card Listings

For notation, see key at front of Listings. $N(>3000)$, $N(3030)$, $N_1(3245)$, $N(3690)$, $N_1(3755)$

Baryons

>3000 MEV REGION - FORMATION EXPERIMENTS

128 $N_1(21>3000)$ I=1/2

WE LIST HERE I=1/2 RESONANCES WITH MASS GREATER THAN ABOUT 3.0 GEV WHICH HAVE BEEN SEEN IN A SINGLE PARTIAL WAVE ANALYSIS ONLY. ALL RESONANCES WHICH HAVE BEEN OBSERVED IN >1 ANALYSES AT ABOUT THE SAME MASS ARE GIVEN A SEPARATE LISTING WITH THE APPROPRIATE QUANTUM NUMBERS.

128 $N_1(21>3000)$ MASS (MEV)

M	3500.	200.	HENDRY	78 HPWA	P1 N	LL15	12/79
M	3400.	200.	HENDRY <td>78 HPWA <td>P1 N <td>ML17 <td>12/79</td> </td></td></td>	78 HPWA <td>P1 N <td>ML17 <td>12/79</td> </td></td>	P1 N <td>ML17 <td>12/79</td> </td>	ML17 <td>12/79</td>	12/79
M	4100.	200.	HENDRY <td>78 HPWA <td>P1 N <td>ML19 <td>12/79</td> </td></td></td>	78 HPWA <td>P1 N <td>ML19 <td>12/79</td> </td></td>	P1 N <td>ML19 <td>12/79</td> </td>	ML19 <td>12/79</td>	12/79
M	AVERAGE MEANINGLESS (SCALE FACTOR = 1.5)						

128 $N_1(21>3000)$ WIDTH (MEV)

M	1300.	200.	HENDRY	78 HPWA	P1 N	LL15	12/79
M	1000.	200. <td>HENDRY <td>78 HPWA <td>P1 N <td>ML17 <td>12/79</td> </td></td></td></td>	HENDRY <td>78 HPWA <td>P1 N <td>ML17 <td>12/79</td> </td></td></td>	78 HPWA <td>P1 N <td>ML17 <td>12/79</td> </td></td>	P1 N <td>ML17 <td>12/79</td> </td>	ML17 <td>12/79</td>	12/79
M	1900.	200. <td>HENDRY <td>78 HPWA <td>P1 N <td>ML19 <td>12/79</td> </td></td></td></td>	HENDRY <td>78 HPWA <td>P1 N <td>ML19 <td>12/79</td> </td></td></td>	78 HPWA <td>P1 N <td>ML19 <td>12/79</td> </td></td>	P1 N <td>ML19 <td>12/79</td> </td>	ML19 <td>12/79</td>	12/79
M	AVERAGE MEANINGLESS (SCALE FACTOR = 1.2)						

128 $N_1(21>3000)$ PARTIAL DECAY MODES

P1	$N_1(21>3000)$	INTO	P1 N	DECAY MASSES
				139+ 938

128 $N_1(21>3000)$ BRANCHING RATIOS

R1	$N_1(21>3000)$	INTO	PI N/TOTAL	PI11	12/79		
R1	3.055	0.02	HENDRY	78 HPWA	P1 N	LL15	12/79
R1	0.040	0.015	HENDRY	78 HPWA	P1 N	ML17	12/79
R1	0.030	0.015	HENDRY	78 HPWA	P1 N	ML19	12/79
R1	AVERAGE MEANINGLESS (SCALE FACTOR = 1.0)						

REFERENCES FOR $N_1(21>3000)$

HENDRY 78 PRL 4 1 222 ARCHIBALD W. HENDRY IIND4LQ11JP
ALSO BO TORONTO CONF 113 ARCHIBALD W. HENDRY IIND11JP

**N(3030)
BUMPS**
73 $N_1(213030)$, JP= 1 I=1/2 PRODUCTION EXPERIMENTS

M	(3030.0)	(3030.0)	HOHLER	66 RVUE	DATA + DISP REL	7/66
M			CITRON	66 CNTR	P1+ P TOTAL	7/66

M	(400.7)	CITRON	66 CNTR	7/66
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73 $N_1(213030)$ PARTIAL DECAY MODES (PROD. EXP.)

P1	$N_1(213030)$	INTO	P1 N	DECAY MASSES
P2	$N_1(213030)$	INTO	P1 P1	938+ 139+ 139

73 $N_1(213030)$ BRANCHING RATIOS (PROD. EXP.)

R1	$N_1(213030)$	INTO	PI N/TOTAL	PI11	11/67		
R1	ONLY 13+12/14 P1 N/TOTAL	MEASURED FOR THIS STATE					
R1	B	(0.088)	(0.016)	BARGER	66 RVUE	TOTAL + CH. ETC.	11/67
R1	B	(0.040)		CITRON	66 CNTR	TOTAL CROSS-SEC.	11/67
R1	B	(0.12)		BARGER	67 CNTR	USES KORMANYOS68	11/67
R1	P	USES RIDGE AMP. RESPONSE. TO CALCULATE DIFF. CROSS SECTIONS AT 180 DEGRE					
R1	B	FOR CRITICAL PH OF THIS METHOD. SEE OLEEN 68.					
R1	D	(0.016)		DIXON	67 RVUE	USES KORMANYOS67	11/67
R1	D	USES ONLY REFERENCE TO CALCULATE DIFF. CROSS SECTIONS AT 180 DEGREES					

REFERENCES FOR $N_1(213030)$ (PROD. EXP.)

HCHLER 66 PRL 12 140 G HOHLER, J REESECKE (KARLSRUHE) I
BARGER 66 PR 151 1123 V BARGER, M OLSSON (MINE) I
CITRON 66 PP 146 1101 *GALBRAITH, NYC, ILENTIC, PHILLIPS, + (MINE) I
BARGER 67 PP 155 1702 V BARGER, D CLINE (MINE) P
DIKHEM 67 PRL 1F 765 F N DIKHEM (MINE) I

PAPERS NOT REFERRED TO IN DATA CARDS

KORMANYOS 67 PP 164 1661 *KORMANYOS, KRISCH, OFALLON, + (MICH,ANL) P
OLEEN 68 PP 166 1768 O OLEEN, O HORN, C SCHMID (CITI)

 **$N_1(3245)$
BUMPS**

74 $N_1(213245)$, JP= 1 PRODUCTION EXPERIMENTS
EXISTENCE NOT CONCLUSIVELY ESTABLISHED. I-SPIN NOT DETERMINED, BUT THE NARROW WIDTH PRECLUDES IDENTIFICATION WITH THE $N_1(213230)$. OMITTED FROM TABLE.

M	3245.0	10.0	KORMANYOS	67 CNTR	PI-P	180 DEG EL	6/68
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M	135.0	0 LE55	KORMANYOS	67 CNTR	6/68
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74 $N_1(213245)$ PARTIAL DECAY MODES (PROD. EXP.)

P1	$N_1(213245)$	INTO	P1 N	DECAY MASSES
				139+ 938

74 $N_1(213245)$ BRANCHING RATIOS (PROD. EXP.)

R1	$N_1(213245)$	INTO	PI N/TOTAL	PI11	6/68
R1	J IS NOT KNOWN. FOLLOWING IS 139+1210+PI N/TOTAL				
R1	J IS NOT KNOWN. FOLLOWING IS 139+1210+PI N/TOTAL				

REFERENCES FOR $N_1(213245)$ (PROD. EXP.)

KORMANYOS 67 PP 164 1661 *KORMANYOS, KRISCH, OFALLON, + (MICH,ANL) P

 **$N(3690)$
BUMPS**

75 $N_1(213690)$, JP= 1 I=1/2 PRODUCTION EXPERIMENTS
A BUMP SEEN IN THE INVARIANT MASS OF A VERY COMPLICATED STATE (N + SEVEN P'S). SO AS EVIDENCE FOR A NEW RESONANCE IT IS NOT CONCLUSIVE. NOT INCLUDED IN TABLE.

M	3690.0	10.0	BARTKE	67 HBC	+ PI+P	B PROGS	8/67
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M	51.0	30.0	BARTKE	67 HBC	8/67
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75 $N_1(213690)$ PARTIAL DECAY MODES (PROD. EXP.)

P1	$N_1(213690)$	INTO	N + 7 P'S	DECAY MASSES
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REFERENCES FOR $N_1(213690)$ (PROD. EXP.)

BARTKE 67 PL 24R 118 *CZYZEMSKI, DANYSZ, + (CRACOW,ORSAY) I

 **$N_1(3755)$
BUMPS**

76 $N_1(213755)$, JP= 1 PRODUCTION EXPERIMENTS
A SMALL PEAK IN THE (P P DBAR) INVARIANT MASS FROM 0.4 BEV/C P+ P TO P+ P P DBAR EVENTS. AS EVIDENCE FOR A NEW RESONANCE IT IS NOT CONCLUSIVE. OMITTED FROM TABLE.

M	3755.0	8.0	ERLICH	68 HBC	+ P1+ P P DBAR	6/68
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M	43.0	20.0	ERLICH	68 HBC	6/68
---	------	------	--------	--------	------

76 $N_1(213755)$ PARTIAL DECAY MODES (PROD. EXP.)

P1	$N_1(213755)$	INTO	P1+ P P DBAR	DECAY MASSES
				139+ 938+ 938+ 938

REFERENCES FOR $N_1(213755)$ (PROD. EXP.)

ERLICH 68 PRL 20 686 R ERLICH, R PLANDJ, B WHITTAKER (RUTGERS)

Baryons

$\Delta(1232)$

Data Card Listings

For notation, see key at front of Listings.

S=0 I=3/2 NUCLEON STATES (Δ)

$\Delta(1232)$

33 N#3/2(1232, 1/2⁺3/2⁺) 1⁺3/2

P₃₃

THIS RESONANCE IS WELL ESTABLISHED. SEE CARTER 71 AND CARTER 73 FOR P₁₁ CROSS-SECTION DATA IN THIS REGION.

33 N#3/2(1232) MASS (MEV)

Table with columns for mass values and parameters like ROPER, ALMEHD, CHENG, etc. Includes sub-headers like (1232.1), (1235.1), (1243.21), etc.

Table with columns for mass values and parameters like BERENDS, CRAWFORD, BARBOUR, etc. Includes sub-headers like (1231.0), (1231.7), (1231.2), etc.

Table with columns for mass values and parameters like OLSSON, CARTER, PEDRONI, etc. Includes sub-headers like (1236.0), (1236.5), (1237.1), etc.

37 N#3/2(1232) WIDTH (MEV)

Table with columns for width values and parameters like ROPER, ALMEHD, CHENG, etc. Includes sub-headers like (120.1), (112.21), (145.0), etc.

Table with columns for width values and parameters like OLSSON, CARTER, PEDRONI, etc. Includes sub-headers like (120.0), (111.1), (111.5), etc.

Table with columns for width values and parameters like CRAWFORD, BARBOUR, MIRDSHNIC, etc. Includes sub-headers like (120.2), (111.0), (131.1), etc.

Table with columns for width values and parameters like OLSSON, CARTER, PEDRONI, etc. Includes sub-headers like (119.6), (114.7), (114.9), etc.

33 (N#0) - (N#++) MASS DIFFERENCE (MEV)

Table with columns for mass difference values and parameters like OLSSON, CARTER, PEDRONI, etc. Includes sub-headers like (10.451), (1.3), (1.9), etc.

33 (N#0) - (N#++) WIDTH DIFFERENCE (MEV)

Table with columns for width difference values and parameters like CARTER, PEDRONI, etc. Includes sub-headers like (6.5), (10.3), (6.6), etc.

33 N#3/2(1232) REAL PART OF POLE POSITION (MEV)

Table with columns for real part of pole position values and parameters like MICHAEL, RALL, POG, etc. Includes sub-headers like (1214.), (1211.), (1211.0), etc.

37 N#3/2(1232) -IMAG PART OF POLE POSITION (MEV)

Table with columns for imaginary part of pole position values and parameters like MICHAEL, RALL, POG, etc. Includes sub-headers like (152.), (150.), (152.0), etc.

33 N#3/2(1232) ABSOLUTE VALUE OF ELASTIC POLE RESIDUE (MEV)

Table with columns for absolute value of elastic pole residue values and parameters like BALL, CUTKOSKY, VASAN, etc. Includes sub-headers like (153.), (53.), (52.43TD), etc.

33 N#3/2(1232) PHASE OF ELASTIC POLE RESIDUE (RADIANS)

Table with columns for phase of elastic pole residue values and parameters like BALL, CUTKOSKY, VASAN, etc. Includes sub-headers like (-0.011), (-0.82), (-0.8231TD), etc.

33 N#3/2(1232) PHASE OF N#13/2(1) PHOTOPRODUCTION MULTIPLE AMPLITUDE POLE RESIDUE

Table with columns for phase of nucleon resonance multiple amplitude pole residue values and parameters like MIP, MIRDSHNIC, etc. Includes sub-headers like (10.451), (1.3), (1.9), etc.

33 N#3/2(1232) MAGNETIC MOMENT (NUCLEAR MAGNETONS)

Table with columns for magnetic moment values and parameters like MIP, MIRDSHNIC, etc. Includes sub-headers like (44.71TD), (146.7), (146.7), etc.

Data Card Listings
For notation, see key at front of Listings.

Baryons
A(1232)

33 N#3/212321 PARTIAL DECAY MODES
P1 N#3/212321 INTO M PI
P2 N#3/212321 INTO M GAMMA
P3 N#3/212321 INTO M PI P1
P4 N#3/212321 INTO GAM NUCLEON, HELICITY=1/2
P5 N#3/212321 INTO GAM NUCLEON, HELICITY=3/2

33 N#3/212321 BRANCHING RATIOS
R1 N#3/212321 INTO (M GAMMA)/(M PI) (PERCENT) (P2)/(P1)
R2 0.55 0.02 DALITZ 66 PVUE 7/60
R3 0.53 0.025 BERENDS 71 IPWA PHOTOPROD. ANAL. 10/71
R4 AVERAGE MEANINGLESS (SCALE FACTOR = 1.0)

FOR DEFINITION OF GAMMA-NUCLEON DECAY AMPLITUDES, SEE MINI-REVIEW PRECEDING THE BARYON LISTINGS.

N#3/212321 INTO GAM NUCLEON, HELICITY=1/2 (GEV**2)
-0.142 0.006 DEVENISH 73 DPWA PI M PHOTOPROD. 2/73
-0.138 0.006 KNIES 74 DPWA PI M PHOTOPROD. 2/74
-0.141 0.006 METCALF 74 DPWA PI M PHOTOPROD. 2/74

N#3/212321 INTO GAM NUCLEON, HELICITY=3/2 (GEV**2)
-0.262 0.015 DEVENISH 73 DPWA PI M PHOTOPROD. 2/73
-0.259 0.016 MOODHUS 74 DPWA PI M PHOTOPROD. 2/74
-0.253 0.007 KNIES 74 DPWA PI M PHOTOPROD. 2/74

REFERENCES FOR N#3/212321
M G OLSSON (EMSC)
L O ROSE, R N WRIGHT, B T FELD (LRMETH) JP
DALITZ, SUTHERLAND (EXFOR)
CONTAINS REFERENCES TO EARLIER WORK ON DELTA PHOTOPRODUCTION.

ALMEHED 72 NP 840 157
BALL 72 PRL 28 1343
BARBER 72 NP 117 290
CARTER 73 NP 858 370
CHENG 73 PD 2249
DEVENISH 73 PL 478 573
FELDER 73 PL 438 499
GODVA 73 NP 881 438
GODVA 73 NP 881 438
TSHANG 73 NP 855 445
KNIES 74 PRD 9 2680
METCALF 74 NP 876 253
MOODHUS 74 PRD 9 1
SPERMAN 74 PRD 10 1840
BALL 75 PRD 11 1171
BERENDS 75 NP 884 342
CRAWFORD 75 NP 897 125
CRAWFORD 75 SAMP 20 1840
LICHTEHN 75 LMC 12 1216
OYED (THIS IS)
R. BALLADIN, L. CRAWFORD
CAMPBELL, SHANWAL (60)SUC14)H1JP
FUKUSHIMA, HORIKAWA, KAJIYAMA (MAGDOVA)SAR1)JP
S. S. VASAN (CAN1)JP
S. S. VASAN (CAN1)JP
AZNAURYA 77 EPL-2644571-77
BARBOUR 78 NP 8141 253
NEPENS 78 PRD 18 3911
WAGLE 78 PD 19 1107
PEDRONI 78 NP 4300 321
ZIDELL 78 LMC 21 140

MOELNER 79 MANDSOK OF PI-M SCATTERING, PHYSIK DAHEM WIL-17-1
ALSO 80 TORONTO CONF 3
MADSON 79 SAMP 29 54
ARAI 80 TORONTO CONF 19 1, #P41
CRAWFORD 80 TORONTO CONF 107 R.L. CRAWFORD
KROSKY 80 TORONTO CONF 91 *FORSTYH, BARDOCK, KELLY, HENDRICK
KROCK 80 NP 4336 331 R. KROCK, L. PETERMANN
ZIDELL 80 PD 01 1255 V.S. ZIDELL, R.A. ANNOT, L.D. ROEPER

PAPERS NOT REFERRED TO IN DATA CAPDS
DONNACHIE 68 PL 268 761 DONNACHIE, LOVELACE, KIPSOD (GERM)
FONDA 75 PRD 0 354 FONDA, GHEORGIU, SHAW (ICTP-TRIESTE)SR15JP
HENEY 74 PP 09 302 HENEY, KANE (RICHI)JP
OLSSON 74 LMC 10 333 OLSSON (EDM)JP
SUZUKI 74 NP 808 413 SUZUKI, KUROKAWA, KONO (TOKYO)
GANEKO 75 SAMP 22 522 *KRIEVS, MROSHNICHENKO, NIKIFOROV *SKEVI)JP
NIGRO 75 NP 884 201 NIGRO, SPILLANTINI, VALENTE (PRADO, FRIS)
GANEKO 76 SAMP 24 284 *KRIEVS, MROSHNICHENKO, NIKIFOROV *SKEVI)JP
GANEKO 76 SAMP 24 594 *GANEKO, KRIEVS, MOLESNIKOV *SKEVI)JP
ZABEV 76 SAMP 24 70 ZABEV, KUZNETSOV, STUKOV (FRS)JP

1232 MEV REGION - PRODUCTION EXPERIMENTS

81 N#3/212321 (J3P1/2) 1+3/2 PRODUCTION EXPERIMENTS
SEE THE MINI-REVIEW PRECEDING THE N AND DELTA LISTINGS
FOR A DISCUSSION OF PRODUCTION EXPERIMENTS.

N#3/212321 MASS (MEV) (PROD. EXP.)
4 1221. 9. ANDERSON 70 HMC PI-0 TO PI-M 2/71
1227.0 7.0 ELLIS 71 CNTR 865 PP 3.7, GEV 10/71
1222.7 5.3 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)
1232.0 (6.0) FERRO-LUZ 65 HBC ** K P TO K P PI-0 7/66
1236.0 (14.4) DEANS 66 PVUE ** REAP TOTAL 7/66
1226.0 (2.0) GIDAL 66 DRC ** D TO A (RD)PI 7/70
1226.0 (2.0) HAGER 70 DBC K-D TO A (RD)PI 7/70
1226.0 (2.0) COLTON 72 HBC ** PP TO P1+M TCEV 1/73
1226.0 (2.0) COLTON 72 HBC ** TO P1+P TCEV 1/73
1226.0 (2.0) COLTON 72 HBC ** TO P1+P1-PI-PI 1/73
1231.0 (13.1) LEWIS 73 HBC ** K P TO K P PI 1/76
1226.0 (5.0) HENNER 74 HBC ** PIP TO P1 P 4/75
1226.0 (5.0) LICHMAN 74 HBC ** P1-P TO P1 P 4/75
1226.0 (5.0) BRAUN 75 BC PBAR P AND D,S, T 1/75
1226.0 (12.0) ATHERTON 76 HBC PBAR P 5.7 GEV 2/77
840 1226. 3. GGG12 78 ISP ** PP TO DELTA P 1/80
840 1226. 3. APELODUN 79 HBC ** PBAR P 5.7 GEV 12/79
840 1226. 3. ALLEN 80 HBC ** M P TO M P PI-0 2/82
1225.6 (1.1) AVERAGE (ERROR INCLUDES SCALE FACTOR OF 2.5)
1226.0 (1.0) BRAUNI 75 HBC PBAR P 5.7 GEV 11/75
1226.0 (1.5) APPLE 77 SPEC ** P TO P (N PI+) 1/78
1226.0 (1.0) COPPER 74 HBC D D P CEE 4/75
1226.0 (1.0) BRAUNI 75 HBC PBAR P 5.7 GEV 11/75
1226.0 (5.0) BRAUNI 75 BC PBAR P AND D,S, T 1/75
1226.0 (1.1) COPPER 74 HBC D D P CEE 4/75
1226.0 (1.0) BRAUNI 75 HBC PBAR P 5.7 GEV 11/75
1226.0 (5.0) BRAUNI 75 BC PBAR P AND D,S, T 1/75
1226.0 (1.1) GIDAL 66 DRC - TO P1+P1-PI-PI 7/66
1226.0 (1.0) GIDAL 72 HBC - TO P1+P1-PI-PI 1/73

81 (M*) - (M**)+ MASS DIFFERENCE (MEV) (PROD. EXP.)
D 7.9 6.8 GIDAL 66 DRC
D TR 5.9 3.1 DARND 91 DBC PI+M TO PI+PI-0 2/82
D AVG 6.2 2.8 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)

81 N#3/212321 WIDTH (MEV) (PROD. EXP.)
W 115. 5. ANDERSON 70 HMC PI-0 TO PI-M 2/71
W 103.0 7.0 ELLIS 71 CNTR 865 PP 3.7, GEV 10/71
W 141. 11. MSHGRAVE 75 HBC ** K P TO K P PI 11/75
W 112.0 (5.0) AVERAGE (ERROR INCLUDES SCALE FACTOR OF 2.0)
1125.0 (130.0) FERRO-LUZ 65 HBC ** 7/66
1121.0 (114.0) DEANS 66 PVUE ** 7/66
1126.0 (18.0) GIDAL 66 DRC ** D TO A (RD)PI 7/70
1126.0 (2.0) HAGER 70 DBC K-D TO A (RD)PI 7/70
1126.0 (2.0) COLTON 72 HBC ** PP TO P1+M TCEV 1/73
1126.0 (2.0) COLTON 72 HBC ** TO P1+P TCEV 1/73
1126.0 (2.0) COLTON 72 HBC ** TO P1+P1-PI-PI 1/73
1126.0 (11.0) LEWIS 73 HBC ** K P TO K P PI 1/76
1126.0 (16.0) HENNER 74 HBC ** PIP TO P1 P 4/75
1126.0 (16.0) LICHMAN 74 HBC ** P1-P TO P1 P 4/75
1126.0 (10.0) BRAUN 75 BC PBAR P AND D,S, T 1/75
1126.0 (11.0) ATHERTON 76 HBC PBAR P 5.7 GEV 2/77
1126.0 (11.0) GGG12 78 ISP ** PP TO DELTA P 1/80
1126.0 (11.0) APELODUN 79 HBC ** PBAR P 5.7 GEV 12/79
1126.0 (11.0) ALLEN 80 HBC ** M P TO M P PI-0 2/82
1126.0 (5.0) AVERAGE (ERROR INCLUDES SCALE FACTOR OF 2.4)
115.0 (1.0) BRAUNI 75 HBC PBAR P 5.7 GEV 11/75
1126.0 (60.0) APPLE 77 SPEC ** P TO P (N PI+) 1/78
109. 32. 25. COPPER 74 HBC D D P CEE 4/75
1126.0 (1.0) BRAUNI 75 HBC PBAR P 5.7 GEV 11/75
1126.0 (1.0) BRAUNI 75 BC PBAR P AND D,S, T 1/75
1126.0 (1.0) GIDAL 66 DRC - TO P1+P1-PI-PI 7/66
1126.0 (1.0) GIDAL 72 HBC - TO P1+P1-PI-PI 1/73

Baryons

 $\Delta(1232)$, $\Delta(1550)$, $\Delta(1600)$

Data Card Listings

For notation, see key at front of Listings.

REFERENCES FOR $\Delta(1232)$ 19DD. EXP.)

FERRO-LU 65 MC 36 1101
DEANS 66 PREP111
GIDRAL 66 PP 141 1761

ANDERSON 70 PP 295 699
MADER 70 PP 178 250
ELLIS 71 PL 27 442
COLTON 72 PP 86 95
ELLIS 73 PP 840 283

LOCHTER 74 PP 897 259
LICHTMAN 74 PP 881 31

BRUNU 75 PP 895 441
BRUNU 75 PP 395 553
NUSGRAVE 75 PP 887 367

ATHERTON 76 PP 9103 381
APPLE 77 LIC 18 167
CORGI 78 PP 8163 365
APPELDOORN 79 PP 8156 111
ALLEN 80 PP 8176 289
OKHNO 81 PL 1048 23

FERRO-LUZZI, GEORGE, + (CERN)
S R DEANS, W G MULLADAY (LANDED)
GIDRAL, A KERMAN, S KIM (LLBL)

*BLESEN, BL IEDEN, COLLINS + (BNL, CERN)
*SHAPIRO, MERILL, MONTAGNI (BNL)
*MAGLICH, ANDREW, WINNEN, SILVERMAN (MIT)
E COLTON, A KIRSOPP (LLBL)
*ALLEY, ANDREW, DANIEL, ISLAMI, FLOM, CLOIC, COFF

COOPER, SEIDL, VAN DER VEIDE (MIT)
LICHTMAN, ISMAS, CASON, KEMPE, MCGAHAN (NDNR)

*GERBER, MAURER, MICHALCH, SCHIBUY (STRB, LBNP)
*GERBER, MAURER, MICHALCH, SCHIBUY (STRB, LBNP)
*PEETERS, SCREINER, WHITMORE, YUTA (ANL)

*ATHERTON, FRENCH, SKURA, BOMM (CERN, PRG)
*ASHI, CHENG, COVINE, GROSSMAN (PRIN, PAU)
*CARWELL, ISIDORA, CONTI (CERN, PRG)
VAN APPELDOORN, MARTING, KOLTUZHINA (ANST)
*BLITSCHAU, + (CERN, HDON, GREN, MP, IMP, DDF)
*WOLFEN, HANVTSOVA, LOEBBEHE, VAN RINDY (ILN)

PAPERS NOT REFERRED TO IN DATA CARDS

ALEXANDER 73 PP 852 221
SEAPURE 73 PP 816 93
YERLAND 74 PP 875 97
ALSO 74 PL 518 187
ATHERTON 75 MC 308 505
STRACHMAN, BRAUN, GERBER, MAURER, C (LBNP, STROF)
WALTER 76 PP 447 33 80
SALANO 80 PP 447 33 187

ALEXANDER, BENARY + TEL. AVI + HEIDELBERG + DESY)
*TASCHNER, RUDOLF, CERN, HDON, GREN, MP, IMP, DDF)
BERLAND, MABER, HODDUS, MULSTIZER + (MIT)
BERLAND, MABER, HODDUS, MULSTIZER + (MIT)
ATHERTON, BRUNU, FRENCH, SKURA (CERN, PRG)
STRACHMAN, BRAUN, GERBER, MAURER, C (LBNP, STROF)
*BECKER, + (DELF, ALACH, CERN, SE, PAU, SAC, VINC)
*KROHNE, + (DELF, ALACH, CERN, SE, PAU, SAC, VINC)

 $\Delta(1550)$ 110 $\Delta(1550)$, JP=1/2+, I=3/2P₃₁

11/77

110 $\Delta(1550)$ MASS (MEV) 11/77

M 8 (1550.1) LONGACRE 77 IPWA P1 N TO 2P1 N 11/77
M 8 ALL LONGACRE 77 PARAMETERS ARE FROM SOLUTION 52, EXCEPT FOR THE POLE 11/77
M 8 POSITION WHICH IS FROM SOLUTIONS 51 AND C1 11/77
M (1525.) BARNHAM 80 IPWA ** P1 N TO 2P1 N 12/79
M (1506.0) CRAWFORD 80 DPWA P1 N PHOTOPROD. 12/81

110 $\Delta(1550)$ WIDTH (MEV) 11/77

M 8 (111.) LONGACRE 77 IPWA P1 N TO 2P1 N 11/77
M 4 (42.) BARNHAM 80 IPWA ** P1 N TO 2P1 N 12/79
M (137.0) CRAWFORD 80 DPWA P1 N PHOTOPROD. 12/81

110 $\Delta(1550)$ REAL PART OF POLE POSITION (MEV) 11/77

RE B 1554. OR 1553. LONGACRE 77 IPWA P1 N TO 2P1 N 11/77

110 $\Delta(1550)$ -2*IMAG PART OF POLE POSITION (MEV) 11/77

IM 8 105. OR 104. LONGACRE 77 IPWA P1 N TO 2P1 N 11/77

110 $\Delta(1550)$ PARTIAL DECAY MODES 11/77

DECAY MASSES

P1 $\Delta(1550)$ INTO P1 N 1394 938
P2 $\Delta(1550)$ INTO $\Delta(1232)$ P1 1232 139
P3 $\Delta(1550)$ INTO N RHO, S+3/2 938 769

110 $\Delta(1550)$ BRANCHING RATIOS 11/77

R1 $\Delta(1550)$ FCW P1 N TO $\Delta(1232)$ P1 SORT(F1=P1) 11/77
R1 B (4.2) LONGACRE 77 IPWA P1 N TO 2P1 N 11/77
R1 0.13 0.05 BARNHAM 80 IPWA ** P1 N TO 2P1 N 12/79

R2 $\Delta(1550)$ FCW P1 N TO N RHO, S+3/2 SORT(F1=P3) 11/77
R2 B (4.0) LONGACRE 77 IPWA P1 N TO 2P1 N 11/77
R2 0.11 0.05 BARNHAM 80 IPWA ** P1 N TO 2P1 N 12/79

110 $\Delta(1550)$ PHOTON DECAY AMPLITUDE=1/21

FOR DEFINITION OF GAMMA-NUCLEON DECAY AMPLITUDES, SEE MINI REVIEW PRECEDING THE BARYON LISTINGS.

A1 $\Delta(1550)$ INTO GAM NUCLEON, HELICITY=1/2, DEWPA=1/21
A1 (0.013) CRAWFORD 80 DPWA P1 N PHOTOPROD. 12/81

REFERENCES FOR $\Delta(1550)$

LONGACRE 77 PP 8122 463
MADER 70 PP 818 595
BARNHAM 80 PP 8168 243
CRAWFORD 80 TORONTO CONF

LONGACRE, DON BEAU (ISACILJJP)
DOLBEK, FRANCIS, REVELES, CADLEY (ISACILJJP)
BARNHAM, GILCKMAN, WIER, JEDRZEJCZAK (LIC)
R.L. CRAWFORD (GLAS)

 $\Delta(1600)$ 19 $\Delta(1600)$, JP=3/2+, I=3/2P₃₁

RECENT ANALYSES INDICATE AT LEAST ONE P33 RESONANCE IN THE 1600-1900 MEV REGION, ALTHOUGH THERE IS SOME DISAGREEMENT ON ITS MASS AND WIDTH. IN ADDITION TO THE STATE LISTED, THE ANALYSIS OF BARNHAM 80 IS CONSISTENT WITH THE PRESENCE OF ANOTHER P33 RESONANCE WITH MASS 1600 MEV AND WIDTH 80 MEV.

SEE ALSO THE $\Delta(121920)$, JP=3/2+ LISTINGS.

19 $\Delta(1600)$ MASS (MEV)

M 3 (1600.) DDMNACH2 68 RVUE PHAS-SHIFT-CERN1 10/69
M 3 KIRSOPP 68 RVUE PHAS-SHIFT ANAL 10/69
M 6 (1801.0) AYO 70 IPWA 1/71
M 6 ENER. DEP. FIT OF ARGAND DIAGRAM 1/71
M 7 (1680.) ALMEHEED 72 IPWA P1 N TO 2P1 N 2/72
M L 1900. OR 1640. LONGACRE 75 IPWA P1 N TO 2P1 N 11/75
M THE 2 SETS OF PARAMETERS ARE FROM METHODS 1 AND 2 OF LONGACRE 75. 11/75
M (1600.) AYO 76 IPWA P1 N TO 2P1 N 11/77
M 8 (1260.) LONGACRE 77 IPWA P1 N TO 2P1 N 11/77
M 8 ALL LONGACRE 77 PARAMETERS ARE FROM SOLUTION 52, EXCEPT FOR THE POLE 11/77
M 8 POSITION WHICH IS FROM SOLUTIONS 51 AND C1 11/77
M (1600.) (50.) CUTKOSKY 79 IPWA P1 N TO P1 N 12/79
M 1522. 13. MOELER 79 IPWA P1 N TO P1 N 12/79
M 4 (1609.1) BARNHAM 80 IPWA ** P1 N TO 2P1 N 12/79
M 1020. 50. CUTKOSKY 80 IPWA P1 N TO P1 N 1/82
M AVERAGE MEANINGLESS SCALE FACTOR = 1.51

19 $\Delta(1600)$ WIDTH (MEV) 11/77

M 3 (281.) DDMNACH2 68 RVUE PHAS-SHIFT-CERN1 10/69
M 3 (240.) KIRSOPP 68 RVUE PHAS-SHIFT ANAL 10/69
M 6 (496.0) AYO 70 IPWA 1/71
M 7 (120.) ALMEHEED 72 IPWA 2/72
M L 205. OR 300. LONGACRE 75 IPWA P1 N TO 2P1 N 11/75
M (120.) AYO 76 IPWA P1 N TO 2P1 N 11/77
M 8 (110.) LONGACRE 77 IPWA P1 N TO 2P1 N 11/77
M (1370.) (70.) CUTKOSKY 79 IPWA P1 N TO P1 N 12/79
M 223. 40. MOELER 79 IPWA P1 N TO P1 N 12/79
M 4 (250.) BARNHAM 80 IPWA ** P1 N TO 2P1 N 12/79
M 300. 100. CUTKOSKY 80 IPWA P1 N TO P1 N 1/82
M AVERAGE MEANINGLESS SCALE FACTOR = 1.01
SEE THE NOTES ACCOMPANYING MASSES QUOTED

19 $\Delta(1600)$ REAL PART OF POLE POSITION (MEV) 11/75

RE B (1600.) LONGACRE 75 IPWA P1 N TO 2P1 N 11/75
RE (154.) OR 1542. LONGACRE 77 IPWA P1 N TO 2P1 N 11/77
RE (1547.) BARNHAM 80 IPWA ** P1 N TO 2P1 N 12/79
RE 1550. 40. CUTKOSKY 80 IPWA P1 N TO P1 N 1/82

19 $\Delta(1600)$ -2*IMAG PART OF POLE POSITION (MEV) 11/75

IM 8 (323.) LONGACRE 75 IPWA P1 N TO 2P1 N 11/75
IM 170. OR 178. LONGACRE 77 IPWA P1 N TO 2P1 N 11/77
IM (230.) CUTKOSKY 79 IPWA P1 N TO P1 N 12/79
IM 200. 60. CUTKOSKY 80 IPWA P1 N TO P1 N 1/82

19 $\Delta(1600)$ REAL PART OF ELASTIC POLE RESIDUE (MEV) 12/79

RER (-18.) 6. CUTKOSKY 79 IPWA P1 N TO P1 N 12/79
RER -15. 6. CUTKOSKY 80 IPWA P1 N TO P1 N 1/82

19 $\Delta(1600)$ IMAG PART OF ELASTIC POLE RESIDUE (MEV) 12/79

IMR 1-11.1 8. CUTKOSKY 79 IPWA P1 N TO P1 N 12/79
IMR 8. 8. CUTKOSKY 80 IPWA P1 N TO P1 N 1/82

19 $\Delta(1600)$ PARTIAL DECAY MODES

DECAY MASSES

P1 $\Delta(1600)$ INTO P1 N 1394 938
P2 $\Delta(1600)$ INTO $\Delta(1232)$ P1 P-WAVE 493 1189
P3 $\Delta(1600)$ INTO $\Delta(1232)$ P1 P-WAVE 1232 139
P4 $\Delta(1600)$ INTO $\Delta(1232)$ P1 P-WAVE 1232 139
P5 $\Delta(1600)$ INTO N RHO, S+1/2, P-WAVE 938 769
P6 $\Delta(1600)$ INTO N RHO, S+3/2, P-WAVE 938 769
P7 $\Delta(1600)$ INTO $\Delta(1232)$ P1 1400 139

19 $\Delta(1600)$ BRANCHING RATIOS

R1 $\Delta(1600)$ INTO NIFOTAL (P1)

R1 (0.) DDMNACH2 68 RVUE PHAS-SHIFT-CERN1 10/69
R1 (0.08) KIRSOPP 68 RVUE PHAS-SHIFT ANAL 10/69
R1 (0.135) AYO 70 IPWA 1/71
R1 (0.1) ALMEHEED 72 IPWA 2/72
R1 (0.1) AYO 76 IPWA 11/77
R1 (0.20) (10.06) CUTKOSKY 79 IPWA P1 N TO P1 N 12/79
R1 (0.2) (0.06) MOELER 79 IPWA P1 N TO P1 N 12/79
R1 0.18 0.04 CUTKOSKY 80 IPWA P1 N TO P1 N 1/82
R1 AVERAGE MEANINGLESS SCALE FACTOR = 1.01

R2 $\Delta(1600)$ INTO $\Delta(1232)$ /TOTAL (P2)
R2 (0.0020)OR LESS FEUERBACH TO RYUE P1 P TO ** SIG 7/70
R2 1 ASSUME MASS= WIDTH, X(ELAST) OF DDMNACH2 68
R2 1 MODEL USED MAY DOUBLE COUNT.

Data Card Listings

For notation, see key at front of Listings.

Baryons

 $\Delta(1600)$, $\Delta(1620)$

R3 #43/211600 FROM PI N TO K SIGMA SORTIP#P21 11/75
 R3 2 (0.006)10.042 DEANS 75 DPWA PI N TO K SIGMA 11/75
 R3 2 RANGE GIVEN IS FROM FOUR BEST SOLUTIONS
 R3 2 DEANS'S DISAGREES WITH P1+ P TO K+ SIGMA+ DATA OF WINNIK77 11/75
 R3 2 AROUND 1920 MEV. 11/75

R4 #43/211600 FROM PI N TO #3/2(1232) PI F-WAVE SORTIP#P31 11/75
 R4 L (-0.36)19 -0.30 LONGACRE 75 IPWA PI N TO ZPI N 11/75
 R4 8 (-0.34) LONGACRE 75 IPWA PI N TO ZPI N 11/77
 R4 8 LONGACRE 77 CONSIDER THIS COUPLING TO BE WELL DETERMINED
 R4 M (0.31) NOVOSELLE 78 IPWA PI N TO ZPI N 3/79
 R4 M BM FIT TO LONGACRE 75 IPWA, PHASE IS NEAR -90 DEGREES. 3/79
 R4 N (0.27) NOVOSELLE 78 IPWA PI N TO ZPI N 3/79
 R4 N BM FIT TO NOVOSELLE 78 IPWA, PHASE IS NEAR -90 DEGREES. THIS FIT 3/79
 R4 N ASSUMES THE MASS IS NEAR 1900 MEV. 12/79
 R4 4 -0.24 0.05 BARNHAM 80 IPWA ** PI N TO ZPI N 12/79
 R4 4 THE COUPLING SIGNS (WHERE DETERMINED) OF BARNHAM 79 HAVE BEEN 12/79
 R4 4 CHANGED TO AGREE WITH THE CONVENTION OF LONGACRE 77. 12/79

R5 #3/2(11600) FROM PI N TO #3/2(1232) PI F-WAVE SORTIP#P41 11/77
 R5 6 (1+0.07) LONGACRE 77 IPWA PI N TO ZPI N 11/77

R6 #3/2(1600) FROM PI N TO RHO,S=1/2,P-WAVE SORTIP#P51 11/77
 R6 8 (-0.01) LONGACRE 77 IPWA PI N TO ZPI N 11/77

R7 #3/2(1600) FROM PI N TO RHO,S=3/2,P-WAVE SORTIP#P61 11/77
 R7 8 (-1.01) LONGACRE 77 IPWA PI N TO ZPI N 11/77

R8 #3/2(11600) FROM PI N TO #1/2(1440) P1 SORTIP#P71 12/79
 R8 4 (-0.23 0.04 BARNHAM 80 IPWA ** PI N TO ZPI N 12/79

19 #3/2(1600) PHOTON DECAY AMPL(GEV**1/2)

FOR DEFINITION OF GAMMA-NUCLEON DECAY AMPLITUDES, SEE MINI-REVIEW PRECEDING THE BARYON LISTINGS.

A1 #3/2(11600) INTO GAM NUCLEON, HELICITY=1/2 (GEV**1/2)
 A1 0.016 0.051 DEVENISH 73 DPWA PI N PHOTOPROD. 2/74
 A1 -0.033 0.037 DEVENISH 74 DPWA PI N PHOTOPROD. 4/75
 A1 0.003 0.015 KNIES 74 DPWA PI N PHOTOPROD. 2/74
 A1 0.0 0.038 METCALF 74 DPWA PI N PHOTOPROD. 2/74
 A1 0.0 0.020 FELLER 76 DPWA PI N PHOTOPROD. 2/77
 A1 +0.006 0.016 AZNAURVAN 77 DPWA P10 PHOTOPROD. 12/79
 A1 -0.110 0.014 AZNAURVAN 77 DPWA P10 PHOTOPROD. 2 12/79
 A1 0.000 0.030 BARBOUR 78 DPWA PI-N PHOTOPROD. 3/79
 A1 -0.005 0.020 CRAWFORD 80 DPWA PI N PHOTOPROD. 12/81
 A1
 A1 AVERAGE MEANINGLESS ISCALE FACTOR = 2.73

A2 #3/2(11600) INTO GAM NUCLEON, HELICITY=3/2 (GEV**1/2)
 A2 0.076 0.064 DEVENISH 73 DPWA PI N PHOTOPROD. 2/74
 A2 -0.034 0.022 DEVENISH 74 DPWA PI N PHOTOPROD. 2/75
 A2 0.0 0.038 METCALF 74 DPWA PI N PHOTOPROD. 2/74
 A2 0.0 0.015 FELLER 76 DPWA PI N PHOTOPROD. 2/77
 A2 -0.133 0.042 AZNAURVAN 77 DPWA P10 PHOTOPROD. 12/79
 A2 -0.088 0.019 AZNAURVAN 77 DPWA P10 PHOTOPROD. 2 12/79
 A2 0.000 0.045 BARBOUR 78 DPWA PI-N PHOTOPROD. 3/79
 A2 -0.909 0.020 CRAWFORD 80 DPWA PI N PHOTOPROD. 12/81
 A2
 A2 AVERAGE MEANINGLESS ISCALE FACTOR = 1.43

REFERENCES FOR #3/2(11600)

DOMMACH 40 VIENNA 139 DOMMACH DE BARPOTZEPUS TALK (GLAS)
 #P5QPP 68 THESES R G KIRSOPP (FEINI)

AYED 70 KIEV CONF A AYED,P BAREYRE, G VILLET (SACLII)P
 FEUERBACH 70 NP 168 85 FEUERBACHER/HOLLADAY (VANDERBILT)
 ALMENDED 72 NP 840 157 #ADVICE (LOND)RUTER (GLAS)
 DEVENISH 73 PL 477 53 DEVENTISH,RANKIN,LYTH (LOND)RANKIN(LI)P

DEVENISH 73 PL 528 227 DEVENTISH,LYTH,RANKIN (LOND)LANC(LI)P
 KNIES 74 PRD 9 2680 KNIES,MOHRDORF,OBERLACK (LBL,GLAS)P
 METCALF 74 NP 876 253 M J METCALF, P L WALKER (CITII)P

DEANS 75 NP 874 94 #MITCHELL,MONTEGEMER, V ISFLA,ALABAMA)I)P
 LONGACRE 75 PL 558 411 #ROSEFELD,LASINSKI,SHALUDA (LBL,SACL)I)P
 ALSO 78 PRD 117 195 LONGACRE,PALASINSKI,ROSEFELD (LBL,GLAS)

AYED 76 CEAN-1213 AYED (THESES) (SACLII)P
 FELLER 76 NP 8104 293 #OKUSHI,NA,HORI,KAMA,NA,IKAMA+(INADY+O5AK)I)P

AZNAURVA 77 EFF-264571-77 #KAGOPRO,BAGDOSARYAN (YEREVAN PHYSICS INST.)I)P
 LONGACRE 77 NP 8122 493 LONGACRE,ODLEBEAU (SACLII)P
 ALSO 76 NP 8108 365 DOLBEAUX,TRANTIS,S,MEVEU,CAOYET (SACLII)P

BARBOUR 78 NP 8141 253 BARBOUR,CRAWFORD,PARSONS (GLAS)
 NOVOSELLE 78 NP 8137 500 L E NOVOSELLE (ICAL TECH)I)P
 ALSO 78 NP 8137 500 L E NOVOSELLE (ICAL TECH)I)P

CUTKOSKY 70 PRD 25 2839 #FORSYTH,HENDRICK,KELLY (ICARN)DLB)I)P
 HOEHLER 79 HANDBUCH OF P1-N SCATTERING, PHYS IN DATEN VOL.12-1 #KALSER,WOCHE,PIETAREN (KARLSRUHE I)P
 ALSO 80 TORONTO CONF 3 R-KOCH (KARLSRUHE)I)P

BARNHAM 80 NP 8168 243 BARNHAM,GOLICKMAN,MIER-JEOR,JEJOWICZ (LOIC)
 CRAWFORD 80 TORONTO CONF 107 R L CRAWFORD (GLAS)
 CUTKOSKY 80 TORONTO CONF 14 #FORSYTH,BARCOCK,KELLY,HENORICK (ICARN)DLB)I)P

PAPERS NOT REFERRED TO IN DATA CARDS

AYED 70 NP 318 566 #BAREYRE,VILLET (SACLII)P
 HOEHLER 70 NP 178 321 #CASHMIRE (U, OXFORD)
 WINNIK 77 NP 8128 66 #TOMAF,REVEL,GOLDBERG,BEANY (HAIFI)

 $\Delta(1620)$

R2 #43/211620, JP=1/2-1 1/3/2

S₃₁

THIS RESONANCE IS WELL ESTABLISHED.

R2 #43/211620 MASS (MEV)

M 1 (1646.0) (12.0) DEVLIN 65 CVTR P1+ P-TOTAL ANAL 11/67
 M 1 (1645.0) BAREYRE 68 RVE P4-SHIFT ANAL 11/67
 M 1 WHERE CROSS SECTION IS GREATEST - EVERALL FIT
 M 3 (1635.0) DOMMACH 68 RVE P4-SHIFT ANAL 1/68
 M 6 (1634.0) AYO 70 IPWA P1 N TO ZPI N 1/71
 M 6 EMER. DEP. FIT OF ARGAND DIAGRAM
 M 4 (1637.0) DAVIES 70 RVE P-5 ANAL SOL A 8/69
 M 7 (1620.) ALMENDED 72 IPWA P1 N PHOTOPROD. 2/72
 M 1622.2 TO 1686. CRAWFORD 75 DPWA P1 N PHOTOPROD. 1/76
 M L 1625. CR 1600. LONGACRE 75 IPWA PI N TO ZPI N 11/75
 M (1623.) AYO 70 IPWA P1 N TO ZPI N 11/75
 M (1680.) BARBOUR 76 DPWA P1 N PHOTOPROD. 1/76
 M 8 (1580.) LONGACRE 77 IPWA PI N TO ZPI N 11/77
 M 8 ALL LONGACRE'S PARAMETERS ARE FROM SOLUTION 52, EXCEPT FOR THE POLE
 M 8 POSITION WHICH IS FROM SOLUTIONS 51 AND CL. 11/77
 M 2 (1662.) BARBOUR 78 DPWA PI-N PHOTOPROD. 3/79
 M 2 SUPERSEDES BARBOUR 76.
 M (1620.) (20.) CUTKOSKY 79 IPWA PI N TO PI N 12/79
 M 3010. 7. HOEHLER 79 IPWA PI N TO PI N 12/79
 M 5 BARNHAM 80 IPWA ** PI N TO ZPI N 1/82
 M A (1732.0) (6.) CHEW 80 DPWA ** P1+P TO P1+P 1/82
 M (1786.7) (1.) CHEW 80 DPWA ** P1+P TO P1+P 1/82
 M 48 CHEW 80 REPORTS TAG 531 RESONANCES AT SOMEWHAT HIGHER MASSES THAN
 M 48 OTHER ANALYSES. BOTH ARE LISTED HERE, LABELED A AND B FOR LOWER
 M AB AND HIGHER STATES. 1/82
 M (1657.0) CRAWFORD 80 DPWA PI N PHOTOPROD. 12/81
 M 1629. 20. CUTKOSKY 80 IPWA PI N TO PI N 1/82
 M AVERAGE MEANINGLESS ISCALE FACTOR = 1.01

R2 #43/211620 WIDTH (MEV)

M 1 (250.0) BAF TYPE 68 RVE 11/67
 M 3 (177.0) DOMMACH 68 RVE 6/69
 M 6 (142.0) AYO 70 IPWA P1 N TO ZPI N 1/71
 M 4 (141.0) DAVIES 70 RVE P-5 ANAL SOL A 8/69
 M 7 (140.) ALMENDED 72 IPWA P1 N TO ZPI N 2/72
 M 125. TO 214. CRAWFORD 75 DPWA P1 N PHOTOPROD. 1/76
 M L 160. CR 150. LONGACRE 75 IPWA PI N TO ZPI N 11/75
 M (161.) BARBOUR 76 DPWA P1 N PHOTOPROD. 11/77
 M (124.) BARBOUR 76 DPWA PI N PHOTOPROD. 11/77
 M 8 (120.) LONGACRE 77 IPWA PI N TO ZPI N 11/77
 M 2 (150.) BARBOUR 78 DPWA PI-N PHOTOPROD. 3/79
 M (147.) (20.) CUTKOSKY 79 IPWA PI N TO PI N 12/79
 M 2 (150.) HOEHLER 79 IPWA PI N TO PI N 12/79
 M 5 (120.) BARNHAM 80 IPWA ** PI N TO ZPI N 1/82
 M A (228.3) (16.) CHEW 80 DPWA ** P1+P TO P1+P 1/82
 M 8 (130.) (15.) CHEW 80 DPWA ** P1+P TO P1+P 1/82
 M (161.0) CRAWFORD 80 DPWA PI N PHOTOPROD. 12/81
 M 160. 20. CUTKOSKY 80 IPWA PI N TO PI N 1/82
 M SEE THE NOTES ACCOMPANYING THE MASSES QUOTED.
 M AVERAGE MEANINGLESS ISCALE FACTOR = 1.01

R2 #43/211620 REAL PART OF POLE POSITION (MEV)

RE (1583.) LONGACRE 75 IPWA PI N TO ZPI N 11/75
 RE 8 (1575. CR 1572. LONGACRE 77 IPWA PI N TO ZPI N 11/77
 RE (1597.) CUTKOSKY 79 IPWA PI N TO PI N 12/79
 RE 1600. 15. CUTKOSKY 80 IPWA PI N TO PI N 1/82

R2 #43/211620 -2IMAG PART OF POLE POSITION (MEV)

IM (143.) LONGACRE 75 IPWA PI N TO ZPI N 11/75
 IM 8 (119. DR 128. CUTKOSKY 79 IPWA PI N TO ZPI N 11/77
 IM (120.) CUTKOSKY 79 IPWA PI N TO ZPI N 12/79
 IM 120. 20. CUTKOSKY 80 IPWA PI N TO ZPI N 1/82

R2 #43/211620 REAL PART OF ELASTIC PCLE RESIDUE (MEV)

RER (-6.) CUTKOSKY 79 IPWA PI N TO ZPI N 12/79
 RER -5. CUTKOSKY 80 IPWA PI N TO ZPI N 1/82

R2 #43/211620 1MAG PART OF ELASTIC PCLE RESIDUE (MEV)

IMR (-15.) CUTKOSKY 79 IPWA PI N TO ZPI N 12/79
 IMR -14. 3. CUTKOSKY 80 IPWA PI N TO ZPI N 1/82

R2 #43/211620 PARTIAL DECAY MODES

P1 #43/211620 INTO PI N DECAY MASSES
 1304 938
 P2 #43/211620 INTO N PI P1 538+130+130
 P3 #43/211620 INTO GAM NUCLEON, HELICITY=1/2 0 938
 P4 #43/211620 INTO #3/2(1232) P1 1232+130
 P5 #43/211620 INTO RHO 938+769
 P6 #43/211620 INTO RHO,S=1/2,P-WAVE 938+769
 P7 #43/211620 INTO RHO,S=3/2,P-WAVE 938+769
 P8 #43/211620 INTO NHO,1/2(1440) P1 1440+130

Baryons
Δ(1620), Δ(1700)

Data Card Listings

For notation, see key at front of Listings.

Table with columns for resonance ID, energy, width, and various parameters. Includes entries for Δ(1620) and Δ(1700) with branching ratios.

82 N*3/2(1620) PHOTON DECAY AMPLITUDE=1/2
FOR DEFINITION OF GAMMA-NUCLEON DECAY AMPLITUDES, SEE MINI-REVIEW PRECEDING THE BARYON LISTINGS.

Table listing photonic decay amplitudes for various baryon resonances, including Δ(1620) and Δ(1700).

82 N*3/2(1620) PHOTON DECAY AMPLITUDE=1/2

Table listing photonic decay amplitudes for various baryon resonances, including Δ(1620) and Δ(1700).

Δ(1700) THIS RESONANCE IS WELL ESTABLISHED.

Table listing resonance parameters for Δ(1700), including energy, width, and various amplitudes.

10 N*3/2(1700) WIDTH (MEV)

Table listing widths for various baryon resonances, including Δ(1700).

10 N*3/2(1700) REAL PART OF POLE POSITION (MEV)

Table listing real parts of pole positions for various baryon resonances, including Δ(1700).

10 N*3/2(1700) -2*IMAG PART OF POLE POSITION (MEV)

Table listing imaginary parts of pole positions for various baryon resonances, including Δ(1700).

10 N*3/2(1700) REAL PART OF ELASTIC PCLE RESIDUE (MEV)

Table listing real parts of elastic pole residues for various baryon resonances, including Δ(1700).

10 N*3/2(1700) IMAG PART OF ELASTIC PCLE RESIDUE (MEV)

Table listing imaginary parts of elastic pole residues for various baryon resonances, including Δ(1700).

10 N*3/2(1700) PARTIAL DECAY MODES

Table listing partial decay modes for various baryon resonances, including Δ(1700).

10 N*3/2(1700) BRANCHING RATIOS

Table listing branching ratios for various baryon resonances, including Δ(1700).

Data Card Listings

Baryons

For notation, see key at front of Listings.

Δ(1700), Δ(1900)

82 N#3/211700 INTO GAM SIGMA TOTAL (P3)
83 1 (0.000) FOR LESS FEUERBACH 70 PV P TO K+ SIGA 7/70
84 1 ASSUME MASS, WIDTH, X(ELAST) OF DDNACHIE 68
85 1 MODEL USED W/ DOUBLE COUNT.

PAPERS NOT REFERRED TO IN DATA CARDS
OCNACHI 69 NP 108 433 A DDNACHIE, R KIRSOPP, C LUTVELACE (E84S+E01N)
AYED 70 PL 318 990 +BARREY+VILLET (ISLACTY)
EMLER 70 NP 118 331 +EASHORE (E0, OFFORD)
WINKER 77 NP B128 65 +TOAFF+REVL.GOLDBERG+RENY (EMF11)

Δ(1900)

S3

10 N#3/211700 PHOTON DECAY AMPL(GEV+-1/2)
FOR DEFINITION OF GAMMA-NUCLEON DECAY AMPLITUDES, SEE MEMO-
REVUE PRECEDING THE BARYON LISTINGS.
A1 N#3/211700 INTO GAM NUCLEON, HELICITY=1/2, I(GEV+-1/2)
A1 0.036 0.052 DEVENISH 73 DPMA P1 N PHOTOPROD. 2/74
A1 +0.068 0.042 MDRHOUS 73 DPMA P1 N PHOTOPROD. 2/73
A1 0.11 0.029 DEVENIS 72 DPMA P1 N PHOTOPROD. 6/75
A1 0.078 0.009 KNIES 74 DPMA P1 N PHOTOPROD. 2/74
A1 0.058 0.048 METCALF 74 DPMA P1 N PHOTOPROD. 2/74
A1 0.041 0.029 MDRHOUS 74 DPMA P1 N PHOTOPROD. 2/74
A1 +0.101 0.011 CRAWFORD 75 DPMA P1 N PHOTOPROD. 1/76
A1 +0.120 0.011 BARBOUR 76 DPMA P1 N PHOTOPROD. 1/76
A1 +0.072 0.033 FELLER 76 DPMA P1 N PHOTOPROD. 2/77
A1 +0.147 0.013 AZNAURAN 77 DPMA P10 PHOTOPROD. 1/77
A1 0.10 0.136 AZNAURAN 77 DPMA P10 PHOTOPROD. 2/77
A1 5 +0.130 0.037 BARBOUR 78 DPMA P1-N PHOTOPROD. 3/79
A1 +0.112 0.006 AKAI 80 DPMA P1 N PHOTOPROD. 12/81
A1 +0.058 0.006 AKAI 80 DPMA P1 N PHOTOPROD. 12/81
A1 0.123 0.022 CRAWFORD 80 DPMA P1 N PHOTOPROD. 12/81

TO N#3/211700 WIDTH (MEV)
+ 2 1140.1 LANGBEIN 73 IPMA P1 N-K SIGA-SOL 1 9/73
+ 2 1162.1 LANGBEIN 73 IPMA P1 N-K SIGA-SOL 2 9/73
+ 2 1107.1 AYED 76 IPMA P1 N TO P1 4 1/77
+ 2 1172.1 CUTSKOSKY 79 IPMA P1 N TO P1 4 12/79
+ 2 1179.1 MOELER 79 IPMA P1 N TO P1 4 12/79
+ 2 1193.1 CHEW 80 DPMA P1+P TO P1+P 1/82
+ 2 1170.1 DEWMS 75 DPMA P1 N PHOTOPROD. 12/81
+ 2 1172.1 CUTSKOSKY 80 IPMA P1 N TO P1 4 1/82
AVERAGE MEANINGLESS ISCALE FACTOR = 1.01
SEE THE NOTES ACCOMPANYING MASSES OUFED
TO N#3/211900 PEAL PART OF PCLE POSITION (MEV)
RE 1184.1 CUTSKOSKY 79 IPMA P1 N TO P1 4 12/79
RE 1070. 4.5. CUTSKOSKY 80 IPMA P1 N TO P1 4 1/82
TO N#3/211900 2-PMAG PART OF PCLE POSITION (MEV)
IN 1142.2 LANGBEIN 73 IPMA P1 N TO P1 4 12/79
IN 1171. 5.6. CUTSKOSKY 80 IPMA P1 N TO P1 4 1/82
TO N#3/211900 REAL PART OF ELASTIC PCLE RESIDUE (MEV)
RBP 17.1 CUTSKOSKY 79 IPMA P1 N TO P1 4 12/79
RER 1. 4. CUTSKOSKY 80 IPMA P1 N TO P1 4 1/82
TO N#3/211900 IMAG PART OF ELASTIC PCLE RESIDUE (MEV)
IMR 1-1.1 CUTSKOSKY 79 IPMA P1 N TO P1 4 12/79
IMR 1-1. 7. CUTSKOSKY 80 IPMA P1 N TO P1 4 1/82
TO N#3/211900 PARTIAL DECAY MODES
P1 N#3/211900 INTO P1 4 DECAT MASSES 1394.938
P2 N#3/211900 INTO K SIGMA 493.1189

REFERENCES FOR N#3/211700
DDNACHI 68 PL 268 161 A DDNACHIE, R G KIRSOPP, C LUTVELACE (E84S+E01N)
ALSD 68 VIEW 139 DDNACHIE +RAPPORTEUR'S TALK (E84S)
ALSD 68 THE 158 R G KIRSOPP (E84S)
AYED 70 KIEVA CONF R AYED-P BARREY, G VILLET (ISLACTY)
DAVIES 70 NP B21 359 A DAVIES FEUERBACHER+HOLLAAY (LVAENDREIT)
FEUERBACH 70 NP 168 85
ALMEFRO 72 NP 448 157 +DOLVELACE (LUNO,RTGJ11P)
DEVENISH 73 PL 478 55 DEVENISH+KANKIN+LYTH (LCOU+BOBON+LARC11P)
MDRHOUS 73 PL 422 361 MDRHOUS+OBERLACK (LGLS+LGLJ11P)
DEVENIS 72 PL 529 227 DEVENIS+LYTH+RANKIN EDSY+LANG,DDNACHI
KNIES 74 PR3 2 2080 MIES+MDRHOUS+OBERLACK (LGLS+LGLJ11P)
METCALF 74 NP 876 251 M+METCALF+L WALKER (CIT11P)
MDRHOUS 74 PR3 9 1 MDRHOUS+OBERLACK+ROSENFELD (E84S+LGLJ11P)
CRAWFORD 75 NP 807 125 R L CRAWFORD (LGLS11P)
DEANS 75 NP 804 90 MITCHELL+MONTGOMERY,+ ISFL+ALBA+KAI11P
GAIOS 75 PR3 12 2645 GAIOS+MILLER (LURD11P)
LANGRECE 76 PL 558 411 +E84S+LGLS+LINSINK+SPAC+2+ (LGLS+LGLJ11P)
MDRHOUS 76 PR3 17 1795 MDRHOUS+LINSINK+ROSENFELD (LGLS+LGLJ11P)
AYED 76 CF8-N-1721 AYED (THESS1) (ISLACTY)
BARBOUR 76 NP 811 358 L. W. BARBOUR+R. L. CRAWFORD (LGLS11P)
FELLER 76 NP B104 219 +FUKUSHI+MA+HORI+KANA+KAJI+KANA+INAGAWA+OSAKAI11P
AZNAURYA 77 EPI-264157-77 +AKOPDY+BADGASAPYAN IYEVANAN PHYSICS INST.11P (ISLACTY)
LANGRECE 77 NP B122 493 LANGRECE+JOLLEAU (ISLACTY)
ALSD 76 NP 1108 365 DD+BEAU+LENTIS+NEVEU+CADIE (ISLACTY)
BARBOUR 76 NP B141 255 +BARBOUR+CRAWFORD+PARSONS (LGLS)
CUTSKOSKY 79 PR3 20 2639 +FORSTYTH+HENDRICK+KELLY (CAR+N+LRL11P)
MOELER 79 +MAYORING, PHYSIK DATEN VOL.12-1 SCATTERING (LGLS)
+KASSER+KOCM+PIETRINGHE (KARLSRUHE 11P)
ALSD 80 TORONTO CONF 3 R,KOCH (KARLSRUHE11P)
ARAI 80 TORONTO CONF 93 R L ARAI (TOKY)
BARNHAM 80 NP B168 243 BARNHAM+GLICKMAN+MIER+JEDR+JEDNICE+2+ (LGC1)
CRAWFORD 80 TORONTO CONF 123 D. W. CRAWFORD (LGLJ11P)
CRAWFORD 80 TORONTO CONF 107 R L CRAWFORD (LGLJ11P)
CUTSKOSKY 80 TORONTO CONF 19 +FORSTYTH+BARCOCK+KELLY+HENDRICK (CAR+N+LRL11P)
LIVANOS 80 TORONTO CONF 35 +BAYON+COUTURES+ROCHONSKI+NEVEU (ISLACTY11P)

REFERENCES FOR N#3/211900
LANGBEIN 73 NP 852 251 LANGBEIN+WAGNER (EMFHC11P)
DEANS 75 NP 804 90 MITCHELL+MONTGOMERY,+ (E84S+LGLS+LGLJ11P)
AYED 76 CF8-N-1921 AYED (THESS1) (ISLACTY)

Data Card Listings

Baryons

For notation, see key at front of Listings.

Δ(1905), Δ(1910)

REFERENCES FOR N*3/2(1905)

DOHNACHI 68 PL 268 161
DONNACHIE, R G KIRSOPP, C LOVELACE (CBRN)JJP (GLAS)
ALSO 68 VIEWS 139 DONNACHIE RAPPOURTEUR'S TALK (EDIN)

AYED 70 KIEV CONF R AYED, P BAREYRE, G VILLET (SACLIIJJP)
DAVIES 70 NP B21 359 A DAVIES (GLAS)
FEUERBACHER+HOLLER+JAY IVANDEB+ORLJ (LRL)

ALMEHD 72 NP B40 157 +LOVELACE (LUND+RUTG)JJP
NEHANI 72 PRL 29 1534 +FUNG, KERNAN, SCHALK, + (UNC+ALB.)
LANGBEIN 73 NP B52 251 LANGBEIN+WAGNER (RHMCH)JJP
DEVENISSE 74 PL 528 227 DEVENISSE+LANC+BANKIN (DESY+LANC+BNM)JJP
KNIES 74 PRD 9 26BD KNIES,MOORHOUSE,+BERLACK (LBL+GLASS)JJP
METSALF 74 NP B76 253 M J METCALF,R L WALKER (CITIJJP)

CRAWFORD 75 NP 807 125 R L CRAWFORD (GLAS)JJP
DEANS 75 NP B96 90 +MITCHELL,MONTGOMERY,+ (SFLA,ALAB+M)JJP
LANGREAC 75 PL 558 415 +ROSENFELD,LASINSKI,SMADJAO+ (LBL+SCLC)JJP
ALSO 76 PRD 17 1795 LANGREAC+LASINSKI,ROSENFELD (LBL+SCLC)

AYED 76 CE-N-1921 AYED (THESIS) (SACLIIJJP)
BARBOUR 76 NP B111 358 I. M. BARBOUR,+ R. CRAWFORD (GLAS)JJP
CUTKOSKY 76 NP 27 445 CUTKOSKY+HENDRICK,+ KELLY (CBRN+LBJ)JJP
ALSO 76 DIFD7D CCMF+ 49 CUTKOSKY+HENDRICK,+ CHAD+ (CBRN+LBJ+BRS)JJP
AZNAURFA 77 EP1-26(197)-77 *KADPOV,BAGDASYAN (YEREVAN PHYSICS INST.)JJP

BARBOUR 78 NP B141 253 BARBOUR,CRAWFORD,PARSONS (GLAS)
NDVOSEL 78 NP B137 509 D. E. NDVOSEL (CAL TECH)JJP
ALSO 78 NP B137 445 B. E. NDVOSEL (CAL TECH)JJP

CUTKOSKY 79 PRD 20 2839 +FORSYTH,HENDRIK,+ KELLY (CBRN+LBJ)JJP
HOEHLER 79 HANDBOOK OF P1-N SCATTERING, PHYSI+ DATEN VOL.12-1 KUEISER,KOCH,PIETARINEN (KARLSRUHE)JJP
ALSO 80 TORONTO CDMF 3 R. KOCH

ARA1 80 TORONTO CDMF 93 J. L. ARA1 (TORO)
CHEW 80 TORONTO CDMF 123 D. N. CHEW (LRL)JJP
CRAWFORD 80 TORONTO CDMF 107 R. L. CRAWFORD (GLAS)
CUTKOSKY 80 TORONTO CDMF 19 +DEANS,+BAROCCO,+KELLY,+HENDRICK (CBRN+LBJ)JJP
LIVAMOS 80 TORONTO CDMF 35 +BATON,+COUTURES,+NCHOWSI,+NEU (SACLIIJJP)

PAPERS NOT REFER TO IN DATA CARDS

AYED 70 PL 318 598 +BAREYRE+VILLET (SACLIIJJP)
MINIKET 77 NP 8128 58 +ATONE+RHEL+COLE+DEGR.+BENNY (MATEJ)

Delta(1910) 12 N*3/2(1910, JP=1/2+) I 72 P31
THIS RESONANCE IS WELL ESTABLISHED.

12 N*3/2(1910) MASS (MEV)

Table with columns for resonance name, year, and mass value. Includes entries for 1295.0, 1300.0, 1305.0, etc.

12 N*3/2(1910) WIDTH (MEV)

Table with columns for resonance name, year, and width value. Includes entries for 1339.0, 1308.0, 1290.0, etc.

12 N*3/2(1910) REAL PART OF POLE POSITION (MEV)

Table with columns for resonance name, year, and real part of pole position. Includes entries for 1363.0, 1702.0, 1871.0, etc.

12 N*3/2(1910) -2PIRAG PART OF PCLE POSITION (MEV)

Table with columns for resonance name, year, and pole position. Includes entries for 1250.0, 172.0, 1200.0, etc.

12 N*3/2(1910) REAL PART OF ELASTIC PCLE RESIDUE (MEV)

Table with columns for resonance name, year, and real part of elastic pole residue. Includes entries for -0.3, -0.6, 0.0, etc.

12 N*3/2(1910) IMAG PART OF ELASTIC PCLE RESIDUE (MEV)

Table with columns for resonance name, year, and imaginary part of elastic pole residue. Includes entries for -0.27, -18.1, -2.0, etc.

12 N*3/2(1910) PARTIAL DECAY MODES

Table with columns for resonance name, year, and decay modes. Includes entries for N*3/2(1910) INTO PI, N*3/2(1910) INTO N+SIGMA, etc.

12 N*3/2(1910) BRANCHING RATIOS

Table with columns for resonance name, year, and branching ratios. Includes entries for 10.0, 0.18, 0.2, etc.

R1 AVERAGE MEANINGLESS (SCALE FACTOR = 1.0)
R2 12 N*3/2(1910) INTO IR SIGMA/TOTAL FEUERBACH 70 PVUE (P3)
R3 1 ASSUME MASS, WIDTH, KIELAST OF DONNACHIE 68
R4 1 MODEL USED MAY DOUBLE COUNT.

Table with columns for resonance name, year, and branching ratios. Includes entries for N*3/2(1910) FROM PI N TO K SIGMA, LANGBEIN 73 IPHA, etc.

Table with columns for resonance name, year, and branching ratios. Includes entries for N*3/2(1910) FROM PI N TO N RHO, S+3/2, LANGREAC 77 IPHA, etc.

R5 N EVIDENCE FOR THIS COUPLING IS WEAK. SEE NOVSELETT RE THIS. 12/79
R5 N COUPLING ASSUMES THE MASS IS NEAR 1920 MEV.

12 N*3/2(1910) PHOTON DECAY AMPLITUDE (1/2)

Table with columns for resonance name, year, and photon decay amplitude. Includes entries for N*3/2(1910) INTO GAMMA NUCLEON, HELICITY+1/2, DEVENISSE 74 DPWA, etc.

A1 AVERAGE MEANINGLESS (SCALE FACTOR = 2.9)

REFERENCES FOR N*3/2(1910)

DOHNACHI 68 PL 268 161 A DONNACHIE, R G KIRSOPP, C LOVELACE (CBRN)JJP
ALSO 68 VIEWS 139 DONNACHIE RAPPOURTEUR'S TALK (GLAS)
ALSO 68 THEISIS R G KIRSOPP (EDIN)

AYED 70 KIEV CONF R AYED, P BAREYRE, G VILLET (SACLIIJJP)
DAVIES 70 NP B21 359 A DAVIES (GLAS)
FEUERBACHER+HOLLER+JAY IVANDEB+ORLJ (LRL)

ALMEHD 72 NP B40 157 +LOVELACE (LUND+RUTG)JJP
LANGBEIN 73 NP B52 251 LANGBEIN+WAGNER (LRL+GLASS)JJP

DEVENISSE 74 PL 528 227 DEVENISSE+LANC+BANKIN (DESY+LANC+BNM)JJP
KNIES 74 PRD 9 26BD KNIES,MOORHOUSE,+BERLACK (LBL+GLASS)JJP
METSALF 74 NP B76 253 M J METCALF,R L WALKER (CITIJJP)

CRAWFORD 75 NP 807 125 R L CRAWFORD (GLAS)JJP
DEANS 75 NP B96 90 +MITCHELL,MONTGOMERY,+ (SFLA,ALAB+M)JJP
AYED 76 CE-N-1921 AYED (THESIS) (SACLIIJJP)
BARBOUR 76 NP B111 358 I. M. BARBOUR,+ R. CRAWFORD (GLAS)JJP
CUTKOSKY 76 NP 27 445 CUTKOSKY+HENDRICK,+ KELLY (CBRN+LBJ)JJP
ALSO 76 OXFORD CCMF+ 49 CUTKOSKY+HENDRICK,+ CHAD+ (CBRN+LBJ+BRS)JJP

Baryons

$\Delta(2200)$, $\Delta(2350)$, $\Delta(2400)$

Data Card Listings

For reference, see key at front of Listings.

135 N#3/212200) WIDTH (MEV)

M	400.	150.	HENDRY	78 IPNA	PI N TO PI N	1/82*
M	450.	100.	CUTKOSKY	79 IPNA	PI N TO PI N	1/82*
M	450.	100.	HOEHLER	79 IPNA	PI N TO PI N	1/82*
M	450.	100.	CUTKOSKY	80 IPNA	PI N TO PI N	1/82*

AVERAGE MEANINGLESS (SCALE FACTOR = 1.0)

135 N#3/212200) REAL PART OF POLE POSITION (MEV)

RE	(2094.)		CUTKOSKY	79 IPNA	PI N TO PI N	1/82*
RE	2100.	50.	CUTKOSKY	80 IPNA	PI N TO PI N	1/82*

135 N#3/212200) -2*IMAG PART OF POLE POSITION (MEV)

IM	(274.-)		CUTKOSKY	79 IPNA	PI N TO PI N	1/82*
IM	340.	80.	CUTKOSKY	80 IPNA	PI N TO PI N	1/82*

135 N#3/212200) REAL PART OF ELASTIC PCLE RESIDUE (MEV)

RE	12.		CUTKOSKY	79 IPNA	PI N TO PI N	1/82*
RE	3.	5.	CUTKOSKY	80 IPNA	PI N TO PI N	1/82*

135 N#3/212200) IMAG PART OF ELASTIC PCLE RESIDUE (MEV)

IM	(-1.-)		CUTKOSKY	79 IPNA	PI N TO PI N	1/82*
IM	-8.	3.	CUTKOSKY	80 IPNA	PI N TO PI N	1/82*

135 N#3/212200) PARTIAL DECAY MODES

PI	N#3/212200)	INTO PI N	DECAY MASSES	1394 938
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135 N#3/212200) BRANCHING RATIOS

R1	N#3/212200)	INTO (PI N)/TOTAL	(P1)				
R1	0.09	0.02	HENDRY	78 IPNA	PI N TO PI N	1/82*	
R1	(0.055)	0.02	CUTKOSKY	79 IPNA	PI N TO PI N	1/82*	
R1	0.05	0.02	HOEHLER	79 IPNA	PI N TO PI N	1/82*	
R1	0.06	0.02	CUTKOSKY	80 IPNA	PI N TO PI N	1/82*	

AVERAGE MEANINGLESS (SCALE FACTOR = 1.0)

REFERENCES FOR N#3/212200)

HENDRY 78 PRL 41 222 ARCHIBALD W. HENDRY (IND+LBL)JJP
 ALSO 80 TORONTO CONF 113 ARCHIBALD W. HENDRY (IND)JJP
 CUTKOSKY 79 PRO 21 2839 *FORSYTH,BARCOCK,KELLY (CARN+LBL)JJP
 HOEHLER 79 HANDBOOK OF PI-N SCATTERING, PHYSIK DATEN VOL.12-1 (CARLSTRUME IJP
 ALSO 80 TORONTO CONF 3 *KRAISER,KOCH,PIETARINEN /KARLSRUHE IJP
 R.KOCH
 CUTKOSKY 80 TORONTO CONF 19 *FORSYTH,BARCOCK,KELLY,HENDRICK (CARN+LBL)JJP

$\Delta(2300)$
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H₃₉

123 N#3/212300) MASS (MEV)

M	2450.	100.	HENDRY	78 IPNA	PI N TO PI N	12/79
M	2217.	80.	HOEHLER	79 IPNA	PI N TO PI N	12/79
M	(2204.-)	(13.4)	CHEW	80 IPNA **	PI#P TO PI#P	1/82*
M	2400.	125.	CUTKOSKY	80 IPNA	PI N TO PI N	1/82*

AVERAGE MEANINGLESS (SCALE FACTOR = 1.4)

123 N#3/212300) WIDTH (MEV)

M	500.	200.	HENDRY	78 IPNA	PI N TO PI N	12/79
M	300.	70.	HOEHLER	79 IPNA	PI N TO PI N	12/79
M	(32.-3)	(11.0)	CHEW	80 IPNA **	PI#P TO PI#P	1/82*
M	425.	150.	CUTKOSKY	80 IPNA	PI N TO PI N	1/82*

AVERAGE MEANINGLESS (SCALE FACTOR = 1.0)

123 N#3/212300) REAL PART OF POLE POSITION (MEV)

RE	2370.	80.	CUTKOSKY	80 IPNA	PI N TO PI N	1/82*
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123 N#3/212300) -2*IMAG PART OF POLE POSITION (MEV)

IM	420.	160.	CUTKOSKY	80 IPNA	PI N TO PI N	1/82*
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123 N#3/212300) REAL PART OF ELASTIC PCLE RESIDUE (MEV)

RE	9.	4.	CUTKOSKY	80 IPNA	PI N TO PI N	1/82*
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123 N#3/212300) IMAG PART OF ELASTIC PCLE RESIDUE (MEV)

IM	-3.	5.	CUTKOSKY	80 IPNA	PI N TO PI N	1/82*
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123 N#3/212300) PARTIAL DECAY MODES

PI	N#3/212300)	INTO PI N	DECAY MASSES	1394 938
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123 N#3/212300) BRANCHING RATIOS

R1	N#3/212300)	INTO (PI N)/TOTAL	(P1)				
R1	0.08	0.02	HENDRY	78 IPNA	PI N TO PI N	12/79	
R1	0.03	0.02	HOEHLER	79 IPNA	PI N TO PI N	12/79	
R1	(0.05)	0.02	CHEW	80 IPNA **	PI#P TO PI#P	1/82*	
R1	0.06	0.02	CUTKOSKY	80 IPNA	PI N TO PI N	1/82*	

AVERAGE MEANINGLESS (SCALE FACTOR = 1.3)

REFERENCES FOR N#3/212300)

HOEHLER 79 HANDBOOK OF PI-N SCATTERING, PHYSIK DATEN VOL.12-1 (CARLSTRUME IJP
 *KRAISER,KOCH,PIETARINEN /KARLSRUHE IJP
 R.KOCH
 ALSO 80 TORONTO CONF 3
 CHEW 80 TORONTO CONF 123 D.W.CHEW (LBL)JJP
 CUTKOSKY 80 TORONTO CONF 19 *FORSYTH,BARCOCK,KELLY,HENDRICK (CARN+LBL)JJP

$\Delta(2350)$
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D₃₅

134 N#3/212350) MASS (MEV)

M	2305.	26.	HOEHLER	79 IPNA	PI N TO PI N	1/82*
M	2400.	125.	CUTKOSKY	80 IPNA	PI N TO PI N	1/82*

AVERAGE MEANINGLESS (SCALE FACTOR = 1.0)

134 N#3/212350) WIDTH (MEV)

M	300.	70.	HOEHLER	79 IPNA	PI N TO PI N	1/82*
M	400.	150.	CUTKOSKY	80 IPNA	PI N TO PI N	1/82*

AVERAGE MEANINGLESS (SCALE FACTOR = 1.0)

134 N#3/212350) REAL PART OF POLE POSITION (MEV)

RE	2400.	125.	CUTKOSKY	80 IPNA	PI N TO PI N	1/82*
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134 N#3/212350) -2*IMAG PART OF POLE POSITION (MEV)

IM	400.	150.	CUTKOSKY	80 IPNA	PI N TO PI N	1/82*
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134 N#3/212350) REAL PART OF ELASTIC PCLE RESIDUE (MEV)

RE	5.	17.	CUTKOSKY	80 IPNA	PI N TO PI N	1/82*
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134 N#3/212350) IMAG PART OF ELASTIC PCLE RESIDUE (MEV)

IM	-14.	10.	CUTKOSKY	80 IPNA	PI N TO PI N	1/82*
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134 N#3/212350) PARTIAL DECAY MODES

PI	N#3/212350)	INTO PI N	DECAY MASSES	1394 938
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134 N#3/212350) BRANCHING RATIOS

R1	N#3/212350)	INTO (PI N)/TOTAL	(P1)				
R1	0.04	0.02	HOEHLER	79 IPNA	PI N TO PI N	1/82*	
R1	0.26	0.10	CUTKOSKY	80 IPNA	PI N TO PI N	1/82*	

AVERAGE MEANINGLESS (SCALE FACTOR = 1.0)

REFERENCES FOR N#3/212350)

HOEHLER 79 HANDBOOK OF PI-N SCATTERING, PHYSIK DATEN VOL.12-1 (CARLSTRUME IJP
 *KRAISER,KOCH,PIETARINEN /KARLSRUHE IJP
 R.KOCH
 ALSO 80 TORONTO CONF 3
 CUTKOSKY 80 TORONTO CONF 19 *FORSYTH,BARCOCK,KELLY,HENDRICK (CARN+LBL)JJP

$\Delta(2400)$
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F₃₇

133 N#3/212400) MASS (MEV)

M	2425.	80.	HOEHLER	79 IPNA	PI N TO PI N	1/82*
M	2350.	100.	CUTKOSKY	80 IPNA	PI N TO PI N	1/82*

AVERAGE MEANINGLESS (SCALE FACTOR = 1.0)

Data Card Listings

For notation, see key at front of Listings.

Baryons

$\Delta(2400)$, $\Delta(2420)$

133 N*3/212400 WIDTH (MEV)

#	330.	80.	HOEHLER	70 IPWA	PI N TO PI N	1/82*
#	333.	100.	CUTKOSKY	80 IPWA	PI N TO PI N	1/82*
#	AVERAGE MEANINGLESS (SCALE FACTOR = 1.0)					

133 N*3/212400 REAL PART OF POLE POSITION (MEV)

#	235.	100.	CUTKOSKY	80 IPWA	PI N TO PI N	1/82*
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133 N*3/212400 IMAG PART OF POLE POSITION (MEV)

#	760.	100.	CUTKOSKY	80 IPWA	PI N TO PI N	1/82*
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133 N*3/212400 REAL PART OF ELASTIC POLE RESIDUE (MEV)

RE	7.	13.	CUTKOSKY	80 IPWA	PI N TO PI N	1/82*
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133 N*3/212400 IMAG PART OF ELASTIC POLE RESIDUE (MEV)

IMP	-17.	8.	CUTKOSKY	80 IPWA	PI N TO PI N	1/82*
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137 N*3/212400 PARTIAL DECAY MODES

PI	N*3/212400	INTO	PI N	DECAY MASSES	1394 938
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133 N*3/212400 BRANCHING RATIOS

R1	N*3/212400	INTO	EPI N1/TOTAL	EPI3	1/77*	
R1	3.07	0.04	HOEHLER	70 IPWA	PI N TO PI N	1/82*
R1	3.08	0.04	CUTKOSKY	80 IPWA	PI N TO PI N	1/82*
R1	AVERAGE MEANINGLESS (SCALE FACTOR = 1.0)					

HOEHLER 79 HANDBOOK OF PI-N SCATTERING, PHYSICS DATEN VOL.12-1
KAIJSER,KOCH,PIETARINEN /KAPLUSHNE IJP
ALSO 80 TORONTO CONF 3 K,KOCH /KARLSRUHE IJP

CUTKOSKY 80 TORONTO CONF 19 #FORSYTH,BABCOCK,KELLY,HENDRICK /CARNAULBL IJP

Δ(2400) → **C₃₈**

124 N*3/212400 MASS (MEV)

#	12170.	100.	AYED	70 IPWA	PI N TO PI N	11/77
#	2220.	100.	HENDRY	78 IPWA	PI N TO PI N	12/79
#	2468.	50.	HOEHLER	79 IPWA	PI N TO PI N	12/79
#	2300.	100.	CUTKOSKY	80 IPWA	PI N TO PI N	1/82*
#	AVERAGE MEANINGLESS (SCALE FACTOR = 1.0)					

124 N*3/212400 WIDTH (MEV)

#	1205.	100.	AYED	70 IPWA	PI N TO PI N	11/77
#	450.	200.	HENDRY	78 IPWA	PI N TO PI N	12/79
#	485.	100.	HOEHLER	79 IPWA	PI N TO PI N	12/79
#	333.	100.	CUTKOSKY	80 IPWA	PI N TO PI N	1/82*
#	AVERAGE MEANINGLESS (SCALE FACTOR = 1.0)					

124 N*3/212400 REAL PART OF POLE POSITION (MEV)

PE	2260.	00.	CUTKOSKY	80 IPWA	PI N TO PI N	1/82*
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124 N*3/212400 IMAG PART OF POLE POSITION (MEV)

IM	320.	160.	CUTKOSKY	80 IPWA	PI N TO PI N	1/82*
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124 N*3/212400 REAL PART OF ELASTIC POLE RESIDUE (MEV)

RE	7.	4.	CUTKOSKY	80 IPWA	PI N TO PI N	1/82*
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124 N*3/212400 IMAG PART OF ELASTIC POLE RESIDUE (MEV)

IMP	-3.	3.	CUTKOSKY	80 IPWA	PI N TO PI N	1/82*
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124 N*3/212400 PARTIAL DECAY MODES

PI	N*3/212400	INTO	PI N	DECAY MASSES	1394 938
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124 N*3/212400 BRANCHING RATIOS

R1	N*3/212400	INTO	EPI N1/TOTAL	EPI3	12/79
R1	10.04	0.13	AYED	70 IPWA	11/77
R1	9.13	0.03	HENDRY	78 IPWA	12/79
R1	7.05	0.03	HOEHLER	79 IPWA	12/79
R1	7.05	0.02	CUTKOSKY	80 IPWA	1/82*
R1	AVERAGE MEANINGLESS (SCALE FACTOR = 1.0)				

REFERENCES FOR N*3/212400

AYED	76	18A-N-1021	AYED	17HE1151	ISAGLIIJP
HENDRY	78	PL 41 272	ARCHIBALD W. HENDRY		17HE1811JP
ALSO 80	TORONTO CONF 103	103	HOEHLER	79	IPWA
HOEHLER	79	HANDBOOK OF PI-N SCATTERING, PHYSICS DATEN VOL.12-1	KAIJSER,KOCH,PIETARINEN		17HE1011JP
ALSO 80	TORONTO CONF 3	3	K,KOCH		17HE152411JP

CUTKOSKY 80 TORONTO CONF 19 #FORSYTH,BABCOCK,KELLY,HENDRICK /CARNAULBL IJP

Δ(2420) **H₃₁₁**

THIS RESONANCE IS WELL ESTABLISHED.

BETH ROYCHOUDHURY 71 AND BRANSDEN 71 SEE A POSSIBLE RESONANCE IN THIS MASS REGION. IN ADDITION BRANSDEN 71 FIND A RESONANT P33 AT 2600 MEV.

84 N*3/212420 MASS (MEV)

#	12412.00	100.	AYED	70 IPWA	PI N TO PI N	1/71
#	12400.1	100.	BRANSDEN	71 IPWA	PI N TO PI N	3/72
#	12400.1	100.	ROSENTHAL	71 IPWA	PI N TO PI N	3/72
#	12440.1	100.	REY	74 IPWA	PI N TO PI N	10/74
#	12400.1	100.	AYED	70 IPWA	PI N TO PI N	11/77
#	2400.	60.	HENDRY	78 IPWA	PI N TO PI N	12/79
#	2416.	17.	HOEHLER	79 IPWA	PI N TO PI N	12/79
#	12358.1	100.	CHEW	80 IPWA	PI N TO PI N	1/82*
#	2430.	125.	CUTKOSKY	80 IPWA	PI N TO PI N	1/82*
#	AVERAGE MEANINGLESS (SCALE FACTOR = 1.0)					

84 N*3/212420 WIDTH (MEV)

#	1347.0	100.	AYED	70 IPWA	PI N TO PI N	1/71
#	1488.1	170.	REY	74 IPWA	PI N TO PI N	10/74
#	1298.1	100.	AYED	70 IPWA	PI N TO PI N	11/77
#	460.	100.	HENDRY	78 IPWA	PI N TO PI N	12/79
#	360.	270.	HOEHLER	79 IPWA	PI N TO PI N	12/79
#	1202.1	100.	CHEW	80 IPWA	PI N TO PI N	1/82*
#	450.	150.	CUTKOSKY	80 IPWA	PI N TO PI N	1/82*
#	AVERAGE MEANINGLESS (SCALE FACTOR = 1.2)					

84 N*3/212420 REAL PART OF POLE POSITION (MEV)

RE	2360.	100.	CUTKOSKY	80 IPWA	PI N TO PI N	1/82*
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84 N*3/212420 IMAG PART OF POLE POSITION (MEV)

IM	420.	100.	CUTKOSKY	80 IPWA	PI N TO PI N	1/82*
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84 N*3/212420 REAL PART OF ELASTIC POLE RESIDUE (MEV)

RE	50.	8.	CUTKOSKY	80 IPWA	PI N TO PI N	1/82*
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84 N*3/212420 IMAG PART OF ELASTIC POLE RESIDUE (MEV)

IMP	-9.	11.	CUTKOSKY	80 IPWA	PI N TO PI N	1/82*
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84 N*3/212420 PARTIAL DECAY MODES

PI	N*3/212420	INTO	PI N	DECAY MASSES	1394 938
PI	N*3/212420	INTO	SIGMA K		1197 443

84 N*3/212420 BRANCHING RATIOS

R1	N*3/212420	INTO	EPI N1/TOTAL	EPI3	1/71
R1	6	10.1131	AYED	70 IPWA	11/77
R1	7	10.41	REY	74 IPWA	10/74
R1	1	10.1571	0.0701	0.0351	REY 74 DETERMINES 13H1/21X ONLY, HE HAVE DIVIDED BY 6.
R1	1	REY 74 DETERMINES 13H1/21X ONLY, HE HAVE DIVIDED BY 6.			
R1	10	10.5	AYED	70 IPWA	11/77
R1	0.11	0.02	HENDRY	78 IPWA	12/79
R1	0.0	0.015	HOEHLER	79 IPWA	12/79
R1	10.23	100.	CHEW	80 IPWA	1/82*
R1	0.08	0.03	CUTKOSKY	80 IPWA	1/82*
R1	AVERAGE MEANINGLESS (SCALE FACTOR = 1.0)				

Baryons

$\Delta(2420)$, $\Delta(>2500)$, $\Delta(2750)$

Data Card Listings

For notation, see key at front of Listings.

REFERENCES FOR $\Delta(212420)$

AYED 70 KIEV CONF R AYED, P BAREVRE, G WILLET (SACL1)JP

BRANSDEN 71 NP 826 511 +DGOEN (DURM)JP
 ALSO 70 NP 816 461 ROYCHOUDHURY, PERRIN, BRANSDEN (DURM)JP
 ROYCHOUD 71 NP 827 125 R K ROYCHOUDHURY, D H BRANSDEN (DURM)JP
 DTT 72 PL 428 133 +TRISCHN, VAVRA, RECHARDSS+ (PGCI-STLO, IDH)JP
 ALSO 72 MCGILL THE515 J. VAVRA (MCGI) JP

KEY 74 PAL 32 908 REV, LEMOND, POIRIER, PRETZL (INDM+MPI)JP
 ALSO 74 PAL 33 250 REV, LEMOND, POIRIER, PRETZL (INDM+MPI)JP
 ALSO 75 PRD 11 1777 LENNOX, POIRIER, REV, SANDER+ (INDM+PRAL, VAN)JP

AYED 76 CE#-N-1921 AYED (THE515) (SACL1)JP
 HENDRY 78 PR 41 222 ARCHIBALD W. HENDRY (IND+L)JP
 ALSO 80 TORONTO CONF 113 ARCHIBALD W. HENDRY (IND)JP
 MOEHLER 79 HANDBOOK OF P1-N SCATTERING, PHYSIC DATEN VOL. 12-1 +KREISSER, KOCH, PIETARINEN (KARLSRUHE)JP
 ALSO 80 TORONTO CONF 3 R. KOCH (KARLSRUHE)JP

CHEW 80 TORONTO CONF 123 D. N. CHEW (L)JP
 CUTSKOSKY 80 TORONTO CONF 19 +FORSYTH, BARCOCK, KELLY, HENDRICK (CARN, LBL)JP

PAPERS NOT REFERRED TO IN DATA CARDS

DELANAY 67 PR 19 476 +BUCKLEY, DOBINSON, + (WESTFIELD, LOUC) JP
 AYED 70 PL 318 958 +BAREVRE+VILLET (SACLAV)

2420 MEV REGION - PRODUCTION AND σ_{TOTAL} EXP'TS

69 $\Delta(2420)$, JP = 1/2 PRODUCTION EXPERIMENTS

SEE THE NINI-REVIEW PRECEDING THE M AND DELTA LISTINGS FOR A DISCUSSION OF PRODUCTION EXPERIMENTS.

69 $\Delta(2420)$ MASS (MEV) (PRD., EXP.)

M	2360.0	140.0	DIOGENS	63 CNTR	P1+ P TOTAL	
M	1250.0		ALVAREZ	64 CNTR	P1 + Φ OTODRND	7/66
M	1240.0		HOEHLER	64 RVUE	DATA + DISP REL	
M	1240.0		WAGLIG	64 DSPK	O PI-P CH EX	
M	1245.0		BARGER	66 RVUE	TOTAL + CH EX	11/67
M	1245.0		WAGLIG	64 DSPK	O PI-P CH EX	
M	1245.0		CITRON	66 CNTR	P1+ P TOTAL	7/66

69 $\Delta(2420)$ WIDTH (MEV) (PROD., EXP.)

M	1200.0		DIOGENS	63 CNTR	P1+ P TOTAL	
M	1245.0 <td></td> <td>HOEHLER</td> <td>64 RVUE</td> <td>DATA + DISP REL</td> <td>7/66</td>		HOEHLER	64 RVUE	DATA + DISP REL	7/66
M	1275.0 <td></td> <td>BARGER</td> <td>66 RVUE</td> <td>TOTAL + CH EX</td> <td>11/67</td>		BARGER	66 RVUE	TOTAL + CH EX	11/67
M	310.0	20.0	CITRON	66 CNTR		7/66

69 $\Delta(2420)$ PARTIAL DECAY MODES (PROD., EXP.)

P1	$\Delta(2420)$	INTO	PI	N	DECAY MASSES	
P1	$\Delta(2420)$	INTO	PI	N	1394 938	
P2	$\Delta(2420)$	INTO	SIGMA K		11794 409	
P3	$\Delta(2420)$	INTO	$\Delta(1232)$	P1	1232+ 139	
P4	$\Delta(2420)$	INTO	NEUTRON	P1+ P1	9394 1394 139	

69 $\Delta(2420)$ BRANCHING RATIOS (PROD., EXP.)

R1	$\Delta(2420)$	INTO	PI	N	Φ TOTAL	(P1)	
R1	10.0071	APPROX	DIOGENS	63 CNTR	ASSUMING J=1/2	7/66	
R1	0.113	0.0036	CITRON	66 CNTR	ASSUMING J=1/2	7/66	
R1	10.1631		BARGER	67 FIT	ASSUMING J=1/2	11/67	
R1	10.01		DIOGENS	63 CNTR	ASSUMING J=1/2	11/67	
R1	10.01		CITRON	66 CNTR	ASSUMING J=1/2	11/67	

69 $\Delta(2420)$ PARTIAL DECAY MODES (PROD., EXP.)

R2	$\Delta(2420)$	INTO	PI	N	NEUTRON P1+ P1	Φ TOTAL+21	(P1PA)	
R2	10.0195	0.0048	GALLOWAY	68 RVUE			6/68	

REFERENCES FOR $\Delta(2420)$ (PROD., EXP.)

+JENKINS, KYCIA, RELEY (IBVLL) I
 +BIR-TAN, REBER, LUCKY, OSBORNE, + (INT, GER) I
 G MOHLER, J GIESECKE (KARLSRUHE) I
 +MANNELL, SODICKSON, FACKLER, WARD, + (HEIT) I
 V BARGER, M OLSSON (MUSC) I
 +GALBRAITH, KYCIA, LEHNTY, PHILLIPS, + (BNL) I

BARGER, O CLINE (MUSC) P
 F N DIRKEN (MUSC) P
 KOMMAYOS, KRISHN, DFALLON, + (MICH, ANL) P
 K F GALLOWAY (IND)ANL I

PAPERS NOT REFERRED TO IN DATA CARDS

BACKE 67 MC 514 761 J BACKE, M VERT (KARLSRUHE, OSVAJ-L
 DOBRONOL 67 PL 246 293 D BRANSDEN, GUSKOV, LIKHACHEV, + (DUBNA) P
 DOLEN 68 PR 166 1768 M DOLEN, O HOHN, E SCHMID (ICIT) I
 MAHLIG 68 PR 166 1515 M MAHLIG, I HANKE (MICH) I

FINAL VERSION OF DATA USED IN WAGLIG 64. IN CONJUNCTION WITH CITRON 66 TOTAL CROSS SECTIONS, THIS CHARGE EXCHANGE DATA GIVES COMPLEX ELASTIC SCATTERING AMPLITUDE AT 0 DEGREES.

>2500 MEV REGION - FORMATION EXPERIMENTS

127 $\Delta(2500)$ I=3/2

WE LIST HERE I=3/2 RESONANCES WITH MASS GREATER THAN ABOUT 2.5 GEV WHICH HAVE BEEN SEEN IN A SINGLE PARTIAL WAVE ANALYSIS ONLY. ALL RESONANCES WHICH HAVE BEEN OBSERVED IN Δ ANALYSIS AT ABOUT THE SAME MASS ARE GIVEN A SEPARATE LISTING WITH THE APPROPRIATE QUANTUM NUMBERS.

127 $\Delta(2500)$ MASS (MEV)

M	2850.	150.	HENDRY	78 MPWA	PI N	I311	12/79
M	3200.	200.	HENDRY	78 MPWA	PI N	K313	12/79
M	3300.	200.	HENDRY	78 MPWA	PI N	L317	12/79
M	3700.	200.	HENDRY	78 MPWA	PI N	M319	12/79
M	4100.	300.	HENDRY	78 MPWA	PI N	N321	12/79

M AVERAGE MEANINGLESS (SCALE FACTOR = 2.31)

127 $\Delta(2500)$ WIDTH (MEV)

W	700.	200.	HENDRY	78 MPWA	PI N	I311	12/79
W	1000.	300.	HENDRY	78 MPWA	PI N	K313	12/79
W	1100.	300.	HENDRY	78 MPWA	PI N	L317	12/79
W	1300.	400.	HENDRY	78 MPWA	PI N	M319	12/79
W	1600.	500.	HENDRY	78 MPWA	PI N	N321	12/79

M AVERAGE MEANINGLESS (SCALE FACTOR = 1.01)

127 $\Delta(2500)$ PARTIAL DECAY MODES

P1	$\Delta(2500)$	INTO	PI	N	DECAY MASSES	
P1	$\Delta(2500)$	INTO	PI	N	1394 938	

127 $\Delta(2500)$ BRANCHING RATIOS

R1	$\Delta(2500)$	INTO	PI	N	Φ TOTAL	(P1)	
R1	0.06	0.02	HENDRY	78 MPWA	PI N	I311	12/79
R1	0.045	0.02	HENDRY	78 MPWA	PI N	K313	12/79
R1	0.03	0.01	HENDRY	78 MPWA	PI N	L317	12/79
R1	0.028	0.01	HENDRY	78 MPWA	PI N	M319	12/79
R1	0.018	0.01	HENDRY	78 MPWA	PI N	N321	12/79

M AVERAGE MEANINGLESS (SCALE FACTOR = 1.01)

REFERENCES FOR $\Delta(212500)$

HENDRY 78 PR 41 222 ARCHIBALD W. HENDRY (IND+L)JP
 ALSO 80 TORONTO CONF 113 ARCHIBALD W. HENDRY (IND)JP

$\Delta(2750)$

125 $\Delta(2750)$, JP=1/2-1 I=3/2

125 $\Delta(2750)$ MASS (MEV)

M	2650.	100.	HENDRY	78 MPWA	PI N TO PI 4	12/79
M	2794.	80.	HOEHLER	78 IPWA	PI N TO PI 4	12/79

M AVERAGE MEANINGLESS (SCALE FACTOR = 1.11)

125 $\Delta(2750)$ WIDTH (MEV)

W	900.	100.	HENDRY	78 MPWA	PI N TO PI 4	12/79
W	350.	100.	HOEHLER	78 IPWA	PI N TO PI 4	12/79

M AVERAGE MEANINGLESS (SCALE FACTOR = 1.11)

125 $\Delta(2750)$ PARTIAL DECAY MODES

P1	$\Delta(2750)$	INTO	PI	N	DECAY MASSES	
P1	$\Delta(2750)$	INTO	PI	N	1394 938	

125 $\Delta(2750)$ BRANCHING RATIOS

R1	$\Delta(2750)$	INTO	PI	N	Φ TOTAL	(P1)	
R1	0.05	0.01	HENDRY	78 MPWA	PI N TO PI 4	12/79	
R1	0.04	0.015	HOEHLER	78 IPWA	PI N TO PI 4	12/79	

M AVERAGE MEANINGLESS (SCALE FACTOR = 1.01)

REFERENCES FOR $\Delta(2750)$

HENDRY 78 PR 41 222 ARCHIBALD W. HENDRY (IND+L)JP
 ALSO 80 TORONTO CONF 113 ARCHIBALD W. HENDRY (IND)JP
 MOEHLER 79 HANDBOOK OF P1-N SCATTERING, PHYSIC DATEN VOL. 12-1 +KREISSER, KOCH, PIETARINEN (KARLSRUHE)JP
 ALSO 80 TORONTO CONF 3 R. KOCH (KARLSRUHE)JP

Data Card Listings

For notation, see key at front of Listings. Δ(2850), Δ(2950), Δ(3230), EXOTIC NUCLEONS

Baryons

Δ(2850) BUMPS

85 N*3/2(2850) JP 1 1-3/2 PRODUCTION EXPERIMENTS
MASS (MEV) (PROD. EXP.)
W 1287.0, 0 HDHLER 64 PVLE DATA + CHG REL
R1 1292.0, 0 49977F WAREIG 66 CSW DATA + CHG F

86 N*3/2(2850) WIDTH (MEV) (PROD. EXP.)
W 1333.31 SARDADIN 66 HBC ** 7/66
R1 433.0 40.0 CITRON 66 ENTP 7/66
W 1493.1 126.1 126.1 REY 74 PPKA ** P1+ P 180 DEG CS 10/74

87 N*3/2(2850) PARTIAL DECAY MODES (PROD. EXP.)
DECAY MASSES
P1 N*3/2(2850) INTD P1 N 1394 938
P2 N*3/2(2850) INTD P1 P1 438 1394 130+ 139
P3 N*3/2(2850) INTD P1 P1 638 1394 130

88 N*3/2(2850) BRANCHING RATIOS (PROD. EXP.)
R1 N*3/2(2850) INTD P1 N/TOTAL (P1)
R1 ONLY 1241/234 P1 N/TOTAL MEASURED FOR THIS STATE
R1 B 10.224 10.2161 BARGER 66 PVUE TOTAL + CH EXC. 11/67
R1 B 0.261 0.248 CITRON 66 ENTP TOTAL CROSS-SEC. 11/67
R1 B USES RESC AMP. METHOD TO CALCULATE DIF. CROSS SECTIONS AT 180 DEGS

REFERENCES FOR N*3/2(2850) (PROD. EXP.)
G. HDHLER, J. GIESSECKE (KARLSRUHE) J
M. MANNELL, SODDERSON, FACKLER, MARO, + (MIT)
SARDADIN, DE WIT, HOSKINS, DANNEY, + (CERN)
V. BARGER, M. OLSSON (MILWAUKEE) I
GALBRAITH, KYCIA, LEDNICEK, PHILLIPS, + (BRL) I
V. BARGER, D. CLINE (MILWAUKEE) P
N. DIRKEN (MILWAUKEE) P
DOBROVOLSKI, GUSKOV, LITVACHEV, + (DUBNA) P
KORNYANS, KRISCH, OFALLEN, + (MICH+ANL) P
HALDORSE, JACOBSEN (OSL) J
REY, LENOX, POIRIER, PRETZL (INDAN+ANL) P
REY, LENOX, POIRIER, PRETZL (INDAN+ANL) P
LENOX, POIRIER, PEV, SANDER (INDAN+ANL) P

PAPERS NOT REFERRED TO IN DATA CARDS
BARKER 67 MC 518 761 J. BARKER, L. MYER (KARLSRUHE-ORSAY) J-I
DOLEN 68 PR 266 1769 R. DOLEN, D. MORN, C. SCHMID (CERN)
WAHLIG 68 PR 168 1315 M. A. WAHLIG, E. MANNELL (MIT-PHYS)
FINAL VERSION OF DATA USED IN WAHLIG 66, IN CONJUNCTION WITH CITRON ON TOTAL CROSS SECTIONS, THIS CHARGE EXCHANGE DATA GIVES COMPLEX ELASTIC SCATTERING AMPLITUDE AT 0 DEGREES.

Δ(2950)

K315

126 N*3/2(2950) JP=15/2+1 1+3/2
MASS (MEV)
W 2950, 100. HENDRY 78 WPKA P1 N TO P1 N 12/79
R1 2900, 100. HENDRY 78 WPKA P1 N TO P1 N 12/79
M AVERAGE MEANINGLESS (SCALE FACTOR = 1.0)

126 N*3/2(2950) WIDTH (MEV)
W 702, 200. HENDRY 78 WPKA P1 N TO P1 N 12/79
W 330, 100. HDHLER 79 WPKA P1 N TO P1 N 12/79
M AVERAGE MEANINGLESS (SCALE FACTOR = 1.7)

126 N*3/2(2950) PARTIAL DECAY MODES
DECAY MASSES
P1 N*3/2(2950) INTD P1 N 1394 938

126 N*3/2(2950) BRANCHING RATIOS
R1 N*3/2(2950) INTD P1 N/TOTAL (P1) 12/79
R1 0.03 0.01 HENDRY 78 WPKA P1 N TO P1 N 12/79
R1 0.04 0.02 HDHLER 79 WPKA P1 N TO P1 N 12/79
R1 AVERAGE MEANINGLESS (SCALE FACTOR = 1.0)

REFERENCES FOR N*3/2(2850)
HENDRY 78 PRL 41 727 ARCHIBALD W. HENDRY (L140+R1) J
ALSO 80 TORONTO CONF 113 ARCHIBALD W. HENDRY (IND) J
HOEHER 76 HANNOVER CONF OF P1-N SCATTERING, PHYSIC DATA VOL 12-1
KARLSRUHE JIP (KARLSRUHE) JIP
ALSO 80 TORONTO CONF 3 N. KOCH

Δ(3230) BUMPS

86 N*3/2(3230) JP 1 1-3/2 PRODUCTION EXPERIMENTS
MASS (MEV) (PROD. EXP.)
W 1330.0, 0 CITRON 66 ENTP P1+ P TOTAL 7/66
R1 1296.1 179.1 178.1 REY 74 PPKA ** P1+ P 180 DEG CS 10/74

86 N*3/2(3230) WIDTH (MEV) (PROD. EXP.)
W 1440.0, 0 CITRON 66 ENTP 7/66
R1 1243.1 1243.1 1243.1 REY 74 PPKA ** P1+ P 180 DEG CS 10/74

86 N*3/2(3230) PARTIAL DECAY MODES (PROD. EXP.)
DECAY MASSES
P1 N*3/2(3230) INTD P1 N 1394 938
P2 N*3/2(3230) INTD N P1 P1 438 1394 139

86 N*3/2(3230) BRANCHING RATIOS
R1 N*3/2(3230) INTD P1 N/TOTAL (P1)
R1 ONLY 1241/234 P1 N/TOTAL MEASURED FOR THIS STATE
R1 B 10.00 10.01 CITRON 66 ENTP TOTAL + CH EXC. 11/67
R1 B 0.003 0.1 DCL BARGER 67 ENTP USES CHARGE EXCHANGE 11/67
R1 B USES RESC AMP. METHOD TO CALCULATE DIF. CROSS SECTIONS AT 180 DEGS
R1 B FOR CALCULATION OF T IS METHOD, SEE DOLEN 68.
R1 D 10.251 10.251 DIRKEN 67 PVUE USES CHARGE EXCHANGE 11/67
R1 D USES RESONANCE TO CALCULATE DIF. CROSS SECTIONS AT 180 DEGREES
R1 D 10.451 10.091 10.133 REY 74 PPKA ** P1+ P 180 DEG CS 10/74

REFERENCES FOR N*3/2(3230) (PROD. EXP.)
V. BARGER, M. OLSSON (MILWAUKEE) I
GALBRAITH, KYCIA, LEDNICEK, PHILLIPS, + (BRL) I
V. BARGER, D. CLINE (MILWAUKEE) P
N. DIRKEN (MILWAUKEE) P
DOBROVOLSKI, GUSKOV, LITVACHEV, + (DUBNA) P
KORNYANS, KRISCH, OFALLEN, + (MICH+ANL) P
REY, LENOX, POIRIER, PRETZL (INDAN+ANL) P
LENOX, POIRIER, PEV, SANDER (INDAN+ANL) P

EXOTIC NUCLEONS - 1400-1700 MEV REGION

E(1400-1700) JP = 1+5/2
THIS IS NOT A COMPLETE LIST. WE TABULATE ONLY FROM 1970 ON.
IN A MISSING MASS EXPERIMENT, P1+ P TO P1- + + + +.
DIRULEV 71 FIND NO EVIDENCE FOR EXOTIC 11/5/21 RESONANCES IN THE MASS INTERVAL 1.2 TO 2.2 GEV.

E(1400-1700) MASS (MEV)
M A 29 1627, 12. PRICE 70 DBC -- K-D AT 4.916GEV/C 3/71
M C 25 1400, 12. ADVIVALI 79 HBC + + + + TO NP 4P1 2/82*
M C CROSS-SECTION IS QUOTED AS 24 MICROBARS. DECAYS TO DELTA+ P1+

E(1400-1700) WIDTH (MEV)
W B 29 30, OR LESS CL. 90 PRICE 70 DBC -- P1-P1+ PUMP 3/71
W C 25 43. ADVIVALI 79 HBC + + + + TO NP 4P1 2/82*

E(1400-1700) CROSS SECTION LIMITS (MICROBARN)
CS B 40, OR LESS BANNER 70 OSPA + + + P1+P1+ 9 GEV/C 7/70
CS B 1/5/2 LIMIT GIVEN ABOVE IS FOR MASS RANGE 1540-1750 MEV

REFERENCES FOR E(1400-1700)
CHEZE, HAMEL, FEIGER, ZACCONE + (SACLAVY)
PRICE 70 PL 338-533
ADVIVALI 79 SUMP 29 795

PAPERS NOT REFERRED TO IN DATA CARDS
AMMANN 71 PL 348 533
DIRULEV 71 SUMP 12 536
JONSON 71 PL 348 428
D. JONSON

Baryons

Z's, Z₀(1780)

Data Card Listings

For notation, see key at front of Listing.

Note on the S = +1 Baryon System

The evidence for S = +1 baryons was thoroughly reviewed in our 1976 edition,¹ and has been reviewed more recently by Kelly² and by Oades.³

In the last few years the results of several new experiments have appeared. In particular, the K⁺n charge-exchange polarization data of WATTS 80 (for references in this form, see the Listings that follow), the K⁺n elastic polarization data of ROBERTSON 80, and the K_{Lp}⁰ + K_{Sp}⁰ data of CAMERON 78, ENGLER 78, and CORDEN 79 provided badly needed new constraints on the I = 0 KN amplitudes.

Three new phase-shift analyses have been reported since our last edition. Energy-independent analyses were made by Martin and Oades⁴ in the I=0 and I=1 KN channels, and by WATTS 80 in the I=0 and I=1 channels simultaneously. CORDEN 82 made an energy-dependent analysis of the I=0 channel.

All analyses agree there is little real evidence for classical Breit-Wigner Z* resonances. The Argand diagrams for the P01 and P13 amplitudes exhibit counterclockwise movement, but there are no corresponding speed maxima. However, the parametrizations of MARTIN 75 and ARNDT 78 for the P13 amplitude do have unphysical-sheet poles. The recent low-energy K⁺p elastic polarization data of LOVETT 81 can be well described either by the analysis of ARNDT 78 (which used preliminary data of LOVETT 81), or by other analyses not requiring structure in the P13 amplitude.

The multi-quark states predicted by Jaffe⁵ in the MIT bag model should be most easily observed in the KN system. Jaffe and Low⁶ showed that such multi-quark states need not appear as classical resonances in hadron-hadron scattering. Rather they would manifest themselves as poles in the P-matrix, which connects the discrete quark states inside the bag with the physical scattering states outside. De Swart et al.⁷ have predicted the low-lying multi-quark states which might be seen in the KN channel, and Roiesnel⁸ (using the partial-wave analysis of MARTIN 75) and CORDEN 82 have found such P-matrix poles. However, the agreement with the predictions of de Swart et al. is not completely compelling.

References

1. Particle Data Group, Rev. Mod. Phys. **48**, S188 (1976).
2. R.L. Kelly, in Proceedings of the Meeting on Exotic Resonances (HUPD-7813) (Hiroshima, 1978), ed. I. Endo et al.
3. G.C. Oades, in Low and Intermediate Energy Kaon-Nucleon Physics (1981), ed. E. Ferrari and G. Violini.
4. B.R. Martin and G.C. Oades, in Proceedings of the IVth International Conference on Baryon Resonances (Toronto, 1980), ed. N. Isgur.
5. R.L. Jaffe, Phys. Rev. **D15**, 267 (1977).
6. R.L. Jaffe and F.E. Low, Phys. Rev. **D19**, 2105 (1979).
7. J.J. de Swart et al., in Proceedings of the IVth International Conference on Baryon Resonances (Toronto, 1980), ed. N. Isgur.
8. C. Roiesnel, Phys. Rev. **D20**, 1646 (1979).

S=1 I=0 EXOTIC STATES (Z₀)

Z ₀ (1780)		P ₀₁				
SEE THE MINI-REVIEW PRECEDING THIS LISTING.						
WILSON 72 AND GIACOMELLI 76 FIND SOME SOLUTIONS WITH RESONANT-LIKE BEHAVIOR IN THE P01 PARTIAL WAVE. THE EFFECT SEEN IN THE I=0 TOTAL CROSS SECTIONS. IF A RESONANCE, MUST NAME SPIN=1/2, BECAUSE THE INELASTIC CROSS SECTION IS VERY SMALL AND THE TOTAL CROSS SECTION IS ABOUT 440/KN*2.						

55 Z0(1780) MASS (MEV)						
M	1780.0	CONV	70 CNTR + KAP. D TOTAL	1/71		
M	0	DMELL	70 CNTR	KAP. D TOTAL	1/70	
M	11920.1	WILSON	72 PWA	KAP P01 WAVE	3/72	
M	1	ESTIMATE OF PARAMETERS FROM RW + QUADRATIC BACKGROUND FIT TO P01			3/72	
M	1	(1780.)	CARDOLL	73 CNTR	KN I=0 TCS.FIT 1	6/73
M	1	(1825.)	CARDOLL	73 CNTR	KN I=0 TCS.FIT 2	6/73
M	1	FIT 1+11 OF SINGLE L1 BW+BACKGROUND TO I=0 TCS FROM 0.4-1.1 GEV/C			4/73	
M	1	FIT 2+11 OF L=1 AND L=2 BWS TO SAME DATA SET (L1800) FOR L2 PART			6/73	
M	(1780.)	GIACOMELLI	76 PWA	.38-1.51 GEV/C	10/74	

55 Z0(1780) WIDTH (MEV)						
M	(565.0)	CONV	70 CNTR + KAP. D TOTAL	1/71		
M	(390.)	WILSON	72 PWA	KAP P01 WAVE	3/72	
M	1	(565.)	CARDOLL	73 CNTR	KN I=0 TCS.FIT 1	6/73
M	1	(845.)	CARDOLL	73 CNTR	KN I=0 TCS.FIT 2	6/73
M	(470.)	GIACOMELLI	76 PWA	.38-1.51 GEV/C	10/74	

55 Z0(1780) PARTIAL DECAY MODES						
P1	Z0(1780)	1/37	K N	DECAY MASSES	6/16 6/30	

55 Z0(1780) BRANCHING RATIOS						
R1	Z0(1780)	INTD IN %TOTAL	CONV	70 CNTR + KAP. D TOTAL	1/71	
R1	13.151	WILSON	72 PWA	KAP P01 WAVE	3/72	
R1	13.551	WILSON	72 PWA	KAP P01 WAVE	3/72	
R1	13.751	CARDOLL	73 CNTR	IF J=1/2, FIT 1	6/73	
R1	17.911	CARDOLL	73 CNTR	IF J=1/2, FIT 2	6/73	
R1	10.851	GIACOMELLI	76 PWA	.38-1.51 GEV/C	10/74	

REFERENCES FOR Z0(1780)						
CONV	70	DUSE	CONV	47	W L CONV	(44N)
ALSO	45	PL	109	76	ARANS, CONV, GIACOMELLI, WYLLI &	(50N)
ALSO	70	PR	31	10	CONV, GIACOMELLI, WYLLI & CONIT	(14N)
DMELL	70	DURE	45		J.D. DMELL	(180N)
WILSON	72	MO	842	445	WILSON 72 MO 842 445	(50N+GLAS+DM+P01)
CARDOLL	73	PL	458	531	CARDOLL 73 PL 458 531	(45N+GLAS+DM+P01)
GIACOMELLI	76	PP	471	134	GIACOMELLI 76 PP 471 134	(50N+GLAS+DM+P01+T01)JP

Data Card Listings

For notation, see key at front of Listings.

Baryons

Z₀(1780), Z₀(1865), Z₁(1900)

PAPERS NOT REFERRED TO IN DATA CARDS

LYNCH TO DUNE 9
HIRATA 71 NP 830 157
GOLDHABER, HALL, SEEGER, TRILLING, WOHL (LRLL)
JENKINS, KALBAH, PETERSEN * (LRJ12)
JOHNSON, WASSERLUND (LRJ12)

EXPERIMENTS MAINLY ABOUT ELASTIC CHANNELS --
GOLDHABER, CHENOVSKY, GOLDHABER * (LRLL)
RAY 69 PR 183 1185
KAMEIAGE 72 NH 1497 391
GIACOMELLI * (LRJ12)
ALSO 73 PRA 191 41-73
LONDON 74 PR 9 1569
ADAMS 75 NP 859 348

EXPERIMENTS MAINLY ABOUT INELASTIC CHANNELS --
GIACOMELLI * (LRJ12)
EDLSTEIN, PISK, JOSEPH (CAPN)
KAMEIAGE 77 NP 117 111
GLASSER 77 PD 15 1220
SARITZ 77 PR 15 1846
ROBERTSON, ASHURBY, GRAY, SHAW * (LRJ12)
WATTS 80 PL 958 123

Z₀(1865)

96 Z0(1865), JP 3/2-1 1-0
THIS EFFECT IS STRONGLY ASSOCIATED WITH THE ** N THRESHOLD. SEE HIRATA AND ANO 70 AND GIACOMELLI 73 REPORT PARTIAL WAVE ANALYSIS. HIRATA 73 CLAIMS A RESONANCE IN A MODEL DEPENDENT PHASE. SEE ALSO Z0(1780).

D₀₃

76 Z0(1865) MASS (MEV)
M 1 (1860.0) (110.0) CARTER 67 THEO DISPERSION FIT 8/67
M 1 (1860.0) (110.0) CODL 70 CNTR KAP. 0 TOTAL 8/67
M 1 (1860.0) (110.0) ARON 73 PRA 140 KN *-1.5 GFC 9/73
M 1 (1860.0) (110.0) CARROLL 73 CNTR KN *-1 TCS, FIT 2 9/73
M 1 1 FITTED TO L1 AND L2 RMS TO 1-0 TCS FROM 0.4-1.1 GEV/C. 9/73
M 1 SEE Z0(1780) FOR F1 AND L2 PART OF FIT. 9/73

76 Z0(1865) WIDTH (MEV)
M 1 (120.0) (150.0) CARTER 67 THEO 8/67
M 1 (160.0) (130.0) CODL 70 CNTR 8/67
M 1 (175.1) (125.1) ARON 73 PRA 140 KN *-1.5 GFC 9/73
M 1 (175.1) (125.1) CARROLL 73 CNTR KN *-1 TCS, FIT 2 9/73

76 Z0(1865) PARTIAL DECAY MODES
P1 Z0(1865) INTO K N 493 939
P2 Z0(1865) INTO K N K(1892) 938 901

76 Z0(1865) BRANCHING RATIOS
R1 Z0(1865) INTO K N / TOTAL 10.1% 10.1%
R1 1 (0.15%) (0.025) CODL 70 CNTR IF J=3/2 9/73
R1 1 (0.15%) (0.025) CARROLL 73 CNTR IF J=3/2, FIT 2 9/73
R2 Z0(1865) INTO K N K(1892) / TOTAL 10.2% 10.2%
R2 MAIN INELASTIC DECAY HIRATA 68 HBC 11/68

REFERENCES FOR Z0(1865)

A A CARTER (CAVEYDISH)
HIRATA, WOHL, GOLDHABER, TRILLING (LRLL)
CODL, GIACOMELLI, KYCIA, LEONTICIL * (LRJ12)
GIACOMELLI, KYCIA, LEONTICIL, LUMBERY * (LRJ12)
ABRAMS, CODL, GIACOMELLI, KYCIA, LIL * (LRJ12)
ARON, MICH, MOGAM, SRIVASTAVA (LASL)
KYCIA, MICHEL, MOCETTI, RAHN (LRJ12)

PAPERS NOT REFERRED TO IN DATA CARDS

HIRATA TO DUNE 479
ARON 71 NP 24 427
HIRATA - 71 NP 833 445
GIACOMELLI 72 NP 837 597
ALSO 72 NP 862 531

GOLDHABER, SEEGER, TRILLING, WOHL (LRLL)
ARON 71 NP 24 427
GOLDHABER, HALL, SEEGER, TRILLING, WOHL (LRLL)
GIACOMELLI * (LRJ12)
ALSO 73 PRA 191 41-73
LONDON 74 PR 9 1569
ADAMS 75 NP 859 348

S=1 I=1 EXOTIC STATES (Z₁)

Z₁(1900)

P₁₃

97 Z1(1900), JP 3/2-1 1-1
THIS EFFECT IS STRONGLY ASSOCIATED WITH THE ** DELTA THRESHOLD. SEE THE REVIEW PRECEDING Z0(1780)

97 Z1(1900) MASS (MEV)
M 1 (1902.0) (110.0) AVEO 70 PRA P13, SOL 1 6/70
M 1 (1890.0) (110.0) AVEO 70 PRA P13, SOL 1 6/70
M 1 (2030.0) (110.0) AVEO 70 PRA 511, SOL 1 6/70
M 1 THREE SIGMAS IN ORDER OF DECREASING SIGNIFICANCE, THROUGH AVEO 70
M 1 GIVE PARAMETERS, THEY CONCLUDE PERSONAL INTERPRETATION, SUBJECTIVE.
M 2 (1830.0) (110.0) BARNEIT 70 PRA P13, SOL 1 9/73
M 2 RESONANCE SIGNAL BARELY ABOVE BACKGROUND TO THE LARGE ENERGIES
M 2 IN THE AMPLITUDES RESULTING FROM THE ANALYSIS.
M 1 (1880.0) (110.0) ALBROD 71 PRA ** SOL, GAMMA 10/71
M K (1890.0) (110.0) KATO 71 PRA SOL 11(FIT 3) 10/71
M K (2040.0) (110.0) KATO 71 PRA SOL 11(FIT 1) 10/71
M K KATO 71 ESTIMATE RESONANCE PARAMETERS -- UPDATED PHASE SHEET 3/72
PUBLISHED IN MILLER 72.

97 Z1(1900) WIDTH (MEV)
M 1 (1520.0) (110.0) AVEO 70 PRA ** 6/70
M 1 (1890.0) (110.0) AVEO 70 PRA ** 6/70
M 1 (1575.0) (110.0) AVEO 70 PRA ** 6/70
M 2 (1120.0) (110.0) BARNEIT 70 PRA P13, SOL 1 11 6/70
M 1 (2040.0) (110.0) CODL 70 CNTR ** 3/71
M 1 (190.0) (110.0) ALBROD 71 PRA ** SOL, GAMMA 10/71
M K (220.0) (110.0) KATO 71 PRA SOL 11(FIT 3) 10/71
M K (260.0) (110.0) KATO 71 PRA SOL 11(FIT 1) 10/71
SEE THE NOTES ACCOMPANYING MASSES QUOTED.

97 Z1(1900) REAL PART OF POLE POSITION
RE 1 (1887.0) (110.0) ARNDT 74 OPA K P ELASTIC 4/75
RE 3 (2796.0) (110.0) ARNDT 76 OPA K P 3/75
SUPERSEDES ARNDT 74.

97 Z1(1900) -IMAGINARY PART OF POLE POSITION
IM 3 (1100.0) (110.0) ARNDT 74 OPA K P ELASTIC 4/75
ARNDT 76 OPA K P 3/75

97 Z1(1900) PARTIAL DECAY MODES
P1 Z1(1900) INTO K N 493 939
P2 Z1(1900) INTO N K(1892) 938 901

97 Z1(1900) BRANCHING RATIOS
R1 Z1(1900) INTO K N / TOTAL 10.1% 10.1%
R1 1 (0.15%) (0.025) CODL 70 CNTR IF J=3/2 9/73
R1 1 (0.15%) (0.025) CARROLL 73 CNTR IF J=3/2, FIT 2 9/73
R2 Z1(1900) INTO K N K(1892) / TOTAL 10.2% 10.2%
R2 MAIN INELASTIC DECAY HIRATA 68 HBC 11/68

SEE NOTES ACCOMPANYING THE MASSES QUOTED.

Z1(1900) INTO K N / TOTAL 10.1% 10.1%
Z1(1900) INTO K N K(1892) / TOTAL 10.2% 10.2%
NO EVIDENCE, SPEED HAS MIN. GRIFFITHS 72 HBC ** -9-1.5 GEV/C 3/72

REFERENCES FOR Z1(1900)

BOHNER, BRODM, GOS, GOLDHABER, SEEGER, * (LRLL)
CARTER 67 PRA 18 1077
A A CARTER (CAVEYDISH)
BARVEY, FELTSE, VILLET (SACLAY)
BARNEIT, GOLDMAN, LAASAMEN, STEINBERG (UMD)
ALSO 70 DUNE 445
BARNEIT, GOLDMAN, LAASAMEN, STEINBERG (UMD)
GIACOMELLI, KYCIA, LEONTICIL * (LRJ12)
ALSO 66 PR 17 102
CODL, GIACOMELLI, KYCIA, LEONTICIL * (LRJ12)

ALBROD 71 NP 830 273
ALSO TO DUNE 375
KATO 71 H. PRA, HORIZONTAL
ALSO TO DUNE 367
ALSO TO DUNE 24 615
ANDERSON, ALMEHEED, ... UD, WAGNER (CERN)
ERNE, SENS, WAGNER (CERN)
KATO 71 H. PRA, HORIZONTAL
ALSO TO DUNE 367
ALSO TO DUNE 24 615
KATO, KOEHLER, NOVEY, VOYDASOVA * (LRLL)
HIRATA, HUGHES * (LRJ12)
NOVEY, VOYDASOVA, CUFOSKY * (LRLL)
ARNDT, HACHMAN, ROBER, STEINBERG (UMD)
ARNDT, ROBER, STEINBERG (UMD)

Baryons

Z₁(1900), Z₁(2150), Z₁(2500)

Data Card Listings

For notation, see key at front of Listings.

 PAPERS NOT REFERRED TO IN Z₁ DATA CARDS
 TOTAL-CROSS-SECTION EXPERIMENTS ---
 BUGG 68 PR 202 2599 +GILLMORE, KNIGHT, + (RHIL, BERN, CAVE) I
 BOWEN 70 PR D2 2599 +CALDWELL, DIKREN, JENKINS, KALBACH, + (ARTZ) I
 EDWARDS 73 PR D2 272 +JENNINGS, KALBACH, PETERSEN + (FARRI+WECH)
 CARROLL 73 PR 439 531 +KYLE, JAL, + (TICHEL, MOORE, IT, RAHN) (BNL)
 A K-MATRIX ANALYSIS OF SOME OF THE EARLY K_{PP} DATA ---
 HITE 67 THESIS G E HITE (ILLINOIS)

THEORETICAL AND MODEL DEPENDENT ANALYSES
 CARRERAS 70 NP B19 349 R CARRERAS, A DONNACHIE (DARESBURY, MCHS)
 ALCOCK 73 NP B58 301 ALCOCK, COTTINGHAM (BRIS)JP
 ALCOCK 76 NP B40 73 ALCOCK, COTTINGHAM, DAVIS (BRIS)JP
 ALCOCK 78 J.PHYS. G 4 323 ALCOCK, COTTINGHAM, DAVIS (BRIS)JP
 ROLESNEL 79 PRD 20 1646 C. ROLESNEL (MIT)

EXPERIMENTS MAINLY ABOUT INELASTIC CHANNELS ---
 BLAND 68 UCRL-18123 THESIS R W BLAND (LRL)
 BLAND 69 NP B13 555 +BOWLER, BRODM, KADRY, GOLDHABER, + (LRL)
 BLAND 70 NP B18 537 +BOWLER, BRODM, GOLDHABER, (LRL)
 BLAND 69 AND US DE TO REPLACE BLAND 67 AND BLAND 68:
 HIRATAI 71 NP 833 445 +GOLDBACH, HALL, SEEGER, TRILLING, WOHLL (LBL)
 BRUNET 72 NP 837 114 BRUNET, NARODOUX, DANYSZ + (COFF+SACL+LOIC+LOND)
 HIRATAI 72 NP 838 173 +HIRATAI, HUGHES, LACROIX, EGNA, GLAS, ROMA, TRSTJJP
 LOKER 72 PR D2 2346 +BARISH, GOMEZ, DAVIES, SCHEIN, + (CTI, UCL)
 BERTHON 73 NP 863 96 LEVINS, ALLEN, JACOBS, DANYSZ, (LOND, COIC+COFF)
 LEWIS 73 NP 887 365 +MULLER, TRIANTIS, BERTHON, + (SACL+CERN)JP
 LESQUY 75 NP B99 346 +MULLER, TRIANTIS, BERTHON, (SACL+CERN)JP
 MUSGRAVE 75 NP B87 365 GIACOMELLI+MANDRILLI (BON+GLAS+ROMA+TRST)
 ARMITAGE 77 NP B123 111 +ASTON, DURODD, HELLISON, FITTON (MCHS+DARE)
 MUSGRAVE 77 PR D15 100 +SINDH, TRIVETI, BURNSTEIN, FU, PETRI + (LOND+TRST)
 SKELTY 77 PRD 15 1846 SKELTY, THOMPSON (BNL)

THE MAIN ELASTIC SCATTERING AND POLARIZATION EXPERIMENTS ---
 CARROLL 68 PR 21 1282 +FISCHER, LUNDBY, PHILLIPS, + (BNL, RDCB)
 ANDERS-1 69 PL 288 611 ANDERSON, DAUM, ERNE, LAGNAUX, + (CERN)
 ANDERS-2 69 PL 308 56 DAUM, ERNE, LAGNAUX, + (CERN)
 ASBURY 69 PR 23 194 +DOMELL, KATO, LUNDQUIST, NOVEY, + (ANL, UMD)
 BLAND 70 PR 259 618 R W BLAND, G GOLDHABER, G H TRILLING (LRL)
 BARBER 70 PL 328 214 +BARBER, DUFF, NEWMAN, JRIE + (LDC, ANELL)JP
 GIACOMELLI 70 NP 820 301 GIACOMELLI, GRIFFITHS, (EGNA, GLAS, ROMA, TRST)JP
 +BLAND, GOLDBACH, TRILLING (LRL)
 REBA 70 PR 24 160 +ROTHBERG, ETKINS, GLADIS, (YALE)JP
 ADAMS 71 PR D4 2637 +DAVIES, DOMELL, GRAVER, HATTERS + (BRN+MHEL)
 ERNICH 71 PL 47 70 +LASSON, STEINBERG, + (ORDAN+MCHS+LOND)
 ERNICH 71 PR 26 925 +ETKIN, GLADIS, HUGHES, KONOD, LU+MDRI + (YALE)
 WHITMORE 71 PR D3 1092 +BRAGMS, EISENSTEIN, KIM, CHALLORNA, + (ILL)
 ADAMS 72 NL PAPER 326 +COX, DAVIES, DOMELL, GRAY + (BRN+MHEL)
 CHARLES 72 PL 409 290 +CHOW, EDWARDS, GIBSON, + (BRIS, RHEL, SHMP)
 ALSO 72 NL PAPER 287 CHARLES, CHOW, EDWARDS + (BRIS+MHEL+SHMP)
 ADAMS 73 NP B42 29 +PERNEY, STEWART, THOMPSON, + (LDC, COFF, LOND)
 ADAMS 73 NP B40 36 ADAMS, COX, DAVIES, DOMELL, + (BRN+MHEL)JP
 BARBER 73 NP B41 47 BARBER, ERNE, BUSE, + DAVIES, DUFF + (LDC+MHEL)
 BARNETT 73 PR D 2751 BARNETT, LAASANEH + (UMD+MCHS+BNL)
 CHARLES 73 PURDUE CONF. 170 CHARLES, EDWARDS, + (SHMP+ARH+MHEL+BRIS)
 BURNSTEIN 74 PRD 10 276 BURNSTEIN, LEFERBER, PETERSEN, RUBIN + (ILL+UMD)
 CAMERON 74 NP B78 93 CAMERON, HERAT, JENNINGS, + GLAS, ROMA, TRST)JP
 YUTA 74 NP 881 189 YUTA, BOCK, MUSGRAVE, PEETERS, SCHREINER, + (BNL)
 ADAMS 75 PRD 11 1719 +BARNETT, GOLDMAN, LASCHON, + (UMD, ANL)
 ADAMS 75 NP B87 11 +CARTER, COOK, GLAS, GREEN (MASH)
 PATTON 75 PR 34 475 +BARLETTA, ERNICH, EFRAIM + (YALE, TONY, BNL)
 WATTS 80 PL 958 373 +HUGES, CARTER, COOK, ANH (LQDN+MELI)
 LOVEIT 81 PRD 23 1924 +HUGES, MISHINA, ZELLER (YALE+BNL+KYOTO)

PARTIAL-WAVE ANALYSES (SEE ALSO ADAMS 73, CAMERON 74, AND ARNOT 78)
 CARRERA 70 NP B23 525 R CARRERAS, A DONNACHIE (DARESBURY, MCHS)
 ALSO 70 DUC 447 DONNACHIE, RIRSOPP (DARESBURY, MCHS)
 ERNE 70 DJK 375 +SEMS, WAGNER (CERN)JP
 LEVINE 71 NP 826 413 +MARTIN, THOMPSON (RHEL, LDC)JP
 LORANCE 71 NP 838 141 +WAGNER (CERN)JP
 CUTOOSKY 72 NL PAPER 210 +HICKS, KELLY, SMITH, JOHNSON (CARN+ILL+ANL)
 ERNICH 72 NL PAPER 447 +ETKIN, GLODUS, HUGHES, LU, PATTON + (YALE)
 MARTIN 72 PREPRINT B A MARTIN (LDC)
 MARTIN 75 NP B96 413 CUTOOSKY, HICKS, KELLY, + (CARN+ILL+ANL)JP
 CUTOOSKY 76 NP 8102 139 GIACOMELLI+MANDRILLI (BON+GLAS+TRJ)
 GIACOMELLI 76 NP B110 67 GIACOMELLI+MANDRILLI (LQDN+ARH)
 MARTIN 80 TDMYD CONF. 355 B A MARTIN, G C GADES (LQDN+ARH)

EARLIER ANALYSES THAT DO NOT INCLUDE RECENT POLARIZATION DATA ---
 LEA 68 PR 185 1770 LEA, MARTIN, DADES (RHIL, BNL, CERN)
 CUTOOSKY 70 PR 21 2547 R E CUTOOSKY, B B DEO (CARNegie-MELLON I)

LATEST REVIEW TALKS AND PAPERS
 LEVINE 57 EPV LUND CONF 343 R LEVINE SETTI (RAPPORTEUR) (CHICAGO)
 GLODUS 70 DUC 447 G. GOLDBERG (REVIEWER) (LRL)
 DOMELL 72 NL REVIEW RAPPORTEUR TALK IN BANYON SESSION (BIPR)
 LOVELL 72 NL REVIEW RAPPORTEUR TALK (RUTG)
 DOMELL 73 PURDUE CONF. 157 DOMELL (LOND)
 CUTOOSKY 74 LONDON CONF 11-54 CUTOOSKY (CARN)
 KELLY 75 ANL-HEP-CO-75-58 REVIEW TALK IN BANYON SESSION (LBL)
 ANDERSON 75 PL 408 17 LEA (LBL)
 MARTIN 76 OFDRE CONF. 409 RAPPORTEUR TALK (LDC)
 KELLY 78 HUPB 70 134 WELCH, ON EXOTIC RESONANCES, MHDOSHMA (LBL)
 OADES 81 NRC CONF. 53 LON + INTERMEDIATE ENERGY KIN PHYSICS (IARR)

Z₁(2150) BUMPS

63 Z₁(2150), J_P = 1-1
 A SMALL BUMP IN TOTAL CROSS SECTION AT P=1.8 GEV/C

4 2150. 20. ABRAMS TO CNTR ++ K*P TOTAL 10/71

63 Z₁(2150) WIDTH (MEV)
 I (175.) ABRAMS TO CNTR + K*P TOTAL 10/71

63 Z₁(2150) PARTIAL DECAY MODES
 P1 Z₁(2150) INTO K N DECAY MASSES 493+ 938

63 Z₁(2150) BRANCHING RATIOS
 RI Z₁(2150) INTO K N / J TOTAL
 R1 J IS NOT KNOWN, THE FOLLOWING IS (J+1/2)*P1 (P2)
 Q1 10.04) ABRAMS TO CNTR + K*P TOTAL 10/71

REFERENCES FOR Z₁(2150)
 ABRAMS 70 PR D2 1937 +CDDI, GIACOMELLI, KYCIA, LEONTEICLI + (BNL)
 ALSO 67 PR 19 257 ABRAMS, COOL, GIACOMELLI, KYCIA, LEONTEICLI + (BNL)

Z₁(2500) BUMPS

94 Z₁(2500), J_P = 1-1
 A SMALL BUMP IN TOTAL CROSS SECTION AT P=2.7 GEV/C

64 Z₁(2500) MASS (MEV)
 M 2500. 20. ABRAMS TO CNTR ++ K*P TOTAL 10/71

64 Z₁(2500) WIDTH (MEV)
 W (160.) ABRAMS TO CNTR ++ K*P TOTAL 10/71

64 Z₁(2500) PARTIAL DECAY MODES
 P1 Z₁(2500) INTO K N DECAY MASSES 493+ 938

64 Z₁(2500) BRANCHING RATIOS
 RI Z₁(2500) INTO K N / J TOTAL (P2)
 R1 J IS NOT KNOWN, THE FOLLOWING IS (J+1/2)*P1
 R1 10.031 ABRAMS TO CNTR ++ K*P TOTAL 10/71

REFERENCES FOR Z₁(2500)
 ABRAMS 70 PR D2 1517 +COOL, GIACOMELLI, KYCIA, LEONTEICLI + (BNL)
 ALSO 67 PR 19 257 ABRAMS, COOL, GIACOMELLI, KYCIA, LEONTEICLI (BNL)

Z, CROSS SECTION LIMITS

SEE MINIREVIEW PRECEDING Z₀
 CS UNITS MICROBARS
 CS LESS THAN 50.
 CS A LESS THAN 0.2 +0.3 -0.1 ANDERSON 69 ASPK + PI-P TO +-Z= 10/69
 CS B ABOVE LIMIT FOR *1.2 TO 1.4 GEV CL 95 P.C.
 CS B LESS THAN 1.4 +1.9 -0.5 ANDERSON 69 ASPK + PI-P TO +-Z= 10/69
 CS B ABOVE LIMIT FOR *1.5 TO 2.5 GEV

REFERENCES FOR Z₁ CROSS SECTION LIMITS
 BASSOMPI 68 PL 27E 46R BASSOMPIERRE, + (CERN, BUKELLESI)
 ANDERSON 69 PL 29E 136 +BLESER, BLIEDEN, COLLINS, + (BNL, CARNEGIE)

PAPERS NOT REFERRED TO IN DATA CARDS
 TYSON 67 PR 19 255 +GREENBERG, HUGHES, LU, MINEHART, (YALE)
 MORI 68 PL 28E 152 +GREENBERG, HUGHES, LU, ROTHBERG, + (YALE)
 MORI 69 PR 185 1087 +GREENBERG, HUGHES, LU, MINEHART, + (YALE)
 MORI 69 REPLACES TYSON 67 AND MORI 68.

Data Card Listings

For notation, see key at front of Listings.

Baryons
A's and E'sNote on A's and E's

The number of established Y^* resonances is still slowly increasing. In 1978 two A's and one E were promoted to the Baryon Table, in 1980 none were promoted, and in this edition two A's have been promoted. There remain, however, a large number of proposed but unconfirmed resonances in the Data Card Listings. Table I is an attempt to evaluate the status, both overall and channel by channel, of each Y^* in the Listings; the evaluations are of course partly subjective. A blank indicates there is no evidence at all: either the relevant couplings are small or the resonance does not really exist. The Baryon Table includes only the well-established resonances (overall status 3 or 4 stars). Any particular one of the questionable resonances may disappear with the next analysis, but there are probably many new resonances underlying those already established.

None of the Y^* s proposed in the last few years couple strongly to the main 2-body decay channels $\bar{N}K$, $\Lambda\pi$, and $\Sigma\pi$, and thus they seldom appear in cross sections or invariant mass distributions. Rather, when the reactions $\bar{K}N + \bar{K}N$, $\bar{K}N + \Lambda\pi$, and $\bar{K}N + \Sigma\pi$ are partial-wave analyzed, some of the amplitudes are found to traverse small, more-or-less resonance-like counterclockwise circles. The question in each case is: Is this really a resonance, or is it an idle meander? Is the effect even real, or is it the result of imperfect data and analysis procedures?

I. Formation Experiments

(by G.P. Gopal, Rutherford and Appleton Laboratories)

Partial-wave analyses have been made mainly for the $\bar{N}K$, $\Lambda\pi$, and $\Sigma\pi$ channels, but there are also a few results for the $\bar{E}K$, $\Lambda\omega$, and some quasi-2-body channels. The earliest analyses covered, in the main, rather narrow mass ranges, usually the range of a single bubble chamber experiment. Although the amplitudes obtained often did not join smoothly with those from other analyses of the same channel performed in neighboring mass ranges, they gave useful and fairly reliable information about the strongly coupled resonances dominating the middle of the ranges covered. The more recent analyses

TABLE I. STATUS OF Y^* RESONANCES
THOSE WITH AN OVERALL STATUS OF *** OR **** ARE INCLUDED IN THE MAIN BARYON TABLE. THE OTHERS AWAIT CONFIRMATION.

PARTICLE	LJ	OVERALL STATUS	STATUS AS SEEN IN --						
			CR. SEC.	CHAP. N	LAMP. PI	SIG. PI	OTHER CHANNELS		
LAM11151	P01	****						WEAK TO N, PI	
LAM11451	S01	****							
LAM15222	D03	****	****	****	E	****			
LAM136001	P01	****			D	****		LAMP. PI, LAMP. OMEGA	
LAM13701	S01	****	****	****		****			
LAM16600	G03	****	****	****	I	****		LAMP. PI, SIG. PI	
LAM18001	S01	****			D	****		TO N, SIG. PI	
LAM18051	P01	****							
LAM18101	G04	DIAC				E			
LAM18001	PE					N		LAMP. PI	
LAM18001	F05	****	****	****	F	****		SIG. PI, SIG. PI	
LAM18301	D05	****	****	****	D	****		TO N, SIG. PI	
LAM18001	P03	****	****	****		****			
LAM12001					B			LAMP. OMEGA, N, K	
LAM20201	F07				F				
LAM21001	G07	****	****	****	D	****		LAMP. OMEGA, N, K	
LAM21101	F05	****	****	****	D	****		LAMP. OMEGA, N, K	
LAM23251	D04	*			F			LAMP. OMEGA	
LAM23501		****	****	****	L				
LAM25851		****	****	****					
SIG11931	P11	****						WEAK TO N, PI	
SIG13851	P13	****				****			
SIG14401	PE	*				*			
SIG15600	PE	****				****			
SIG15801	D13	**	**	**		*			
SIG16201	S11	****	****	****		****			
SIG16601	P11	****	****	****		****			
SIG16701	D13	****	****	****	****	****		SEVERAL OTHERS	
SIG16701	PE	**	**	**	**	**	**	SEVERAL OTHERS	
SIG18901	PE	*	*	*	*	*	*	LAMP. 2-PI	
SIG137501	S11	****	****	****	****	****	****	SIG. PI	
SIG11701	P11	DIAC							
SIG11751	D15	****	****	****	****	****	****	SEVERAL OTHERS	
SIG11801	P11	*	*	*	*	*	*	N, K	
SIG19151	F25	****	****	****	****	****	****	SIG. PI	
SIG19401	D13	****	****	****	****	****	****	QUEST. 7-800F	
SIG20001	S11	*	*	*	*	*	*	N, K, LAMP. PI	
SIG20901	F17	****	****	****	****	****	****	SEVERAL OTHERS	
SIG20701	F15	****	****	****	****	****	****		
SIG20801	P13	**	**	**	**	**	**		
SIG21001	G17	****	****	****	****	****	****		
SIG22901		****	****	****	****	****	****		
SIG24551		****	****	****	****	****	****		
SIG26701		****	****	****	****	****	****		
SIG3001		**	**	**	*	*	*		
SIG33701	PE	**						MULTI-PI	

**** GOOD, CLEAR, AND UNMISKABLE.

*** SOUND, BUT IN NEED OF CLARIFICATION OR NOT ABSOLUTELY CERTAIN.

** NEEDS CONFIRMATION.

* WEAK.

* ATTRIBUTED TO THE STATE CLOSEST TO WHERE THE CROSS SECTION PEAKS.

covering much wider mass ranges have found that using these established dominant states as input provides extremely useful constraints in determining the overall partial-wave amplitudes and thus in getting information about the less strongly coupled resonances. Besides covering wider mass ranges, some of the more ambitious analyses have treated several channels simultaneously so that unitarity constraints are automatically satisfied and only a single mass and width is obtained for each resonance. The early multi-channel analyses¹⁻³ included the channels $\bar{N}K$, $\Lambda\pi$, and $\Sigma\pi$, and covered the only mass range (1500 to 1900 MeV) with data of relatively good statistical accuracy, nearly all of it from a single bubble chamber experiment.⁴

The amount of bubble chamber data has since been considerably increased below,⁵ across,⁶ and above⁷ the range covered by that experiment, and there are now data with lower statistics up to 2500 MeV.⁸ More recently, counter experiments have made

Baryons

Λ 's and Σ 's

major contributions by measuring the $K^-p \rightarrow \bar{K}^0n$ total and differential cross sections at low energies,⁹ the K^-p polarizations down to 1630 MeV for the first time,¹⁰ the K^-p polarizations over the range 1700 to 1900 MeV with an order of magnitude increase in statistics,¹¹ the pure $I=1$ K^-n elastic angular distributions from 1600 to 1800 MeV¹² and from 1900 to 2300 MeV,¹³ and the 180° K^-p and 0° $\Sigma^-\pi^+$ differential cross sections from 1550 to 1900 MeV.¹⁴ There are also new data with good statistics on pure $I=1$ K^-p interactions in the $\Lambda\pi^+$ and, more importantly, the $\Sigma^0\pi^+$ channels.¹⁵⁻¹⁷ Finally, there are new $K^-n \rightarrow (\Sigma\pi)^-$ data over the range 1750 to 2200 MeV.¹⁸ All these new data provide sensitive tests of the overall correctness of existing competing partial-wave amplitudes. Some of the new data have yet to be incorporated into new analyses.

In the following, we discuss the more recent partial-wave analyses, comparing them with each other and with the new data.

The $\bar{N}K$ channel: The most recent analysis of this channel is an update¹⁹ of the old Rutherford Lab-Imperial College (RLIC 77) analysis.²⁰ As was its predecessor, it is a conventional energy-dependent analysis with the added constraint that the masses and widths of the resonances must be consistent with those determined in the inelastic channels analyzed previously -- $\Lambda\pi$, $\Sigma\pi$, $\Lambda(1520)\pi$, $\Sigma(1385)\pi$, and $\bar{N}K^*(892)$. With good angular distributions now available at lower energies,⁵ the analysis goes closer to threshold: the range covered is 1470 to 2170 MeV. With the exception of the very high-statistics charge-exchange differential cross-section measurements⁹ (which are in serious disagreement with both the earlier and the latest⁶ high-statistics bubble chamber measurements) and the backward elastic data,¹⁴ all the new $\bar{N}K$ data mentioned above have been included. As before, angular distributions (a total of 5110 data points) have been fitted directly. The new partial-wave amplitudes are not significantly different from the old RLIC 77 amplitudes for this channel. However, the inclusion of the $I=1$ K^-n data has removed some of the uncertainties in the $I=1$ \bar{K}^0 spectrum.

The LBL-Mc. Holyoke-CERN analysis²¹ covers the

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For notation, see key at front of Listings.

narrower range of 1500 to 1940 MeV and also includes most of the new data. It is an energy-dependent analysis using a unitary background parametrization in terms of scattering lengths. The cusp effects at the $\Lambda\eta$ and $\Sigma\eta$ thresholds are included by introducing a square-root singularity in the energy variation of the widths of the corresponding resonances. The inclusion of their own high-statistics charge-exchange data⁹ -- admittedly not in good agreement with bubble chamber measurements -- all but kills the less well-established resonances.

The University College, London (UCL) K -matrix energy-dependent analysis²² covers the range 1540 to 2000 MeV. The $\bar{N}K$ amplitudes are consistent with those of the other analyses over the greater part of this range. However, at the low-energy end there are considerable differences, indicating the lack of strong constraints from the dominant $D03$ $\Lambda(1520)$ state -- just below the range covered. The new K^-n angular distributions and K^-p polarization measurements are not very well described by the amplitudes from this analysis.

These analyses, all below 2200 MeV, are complemented by the College de France-Saclay (CdF-S) energy-dependent analysis⁸ covering the range 2070 to 2440 MeV. Besides the conventional polynomial parametrization of the background amplitudes, they also tried a more economical parametrization with constraints imposed from the duality hypothesis (that s -channel backgrounds come exclusively from the t -channel Pomeron exchange amplitude). With 30 fewer free parameters, the results are consistent with the more conventional approach.

The $\Sigma\pi$ channel: There are no new analyses of this channel. There is very little agreement, particularly in the lower partial waves, between the amplitudes of the two multi-channel analyses.^{20,22} The low energy $K^-p \rightarrow \Sigma^0\pi^+$ data^{15,16} are better explained by the RLIC 77 amplitudes than by the UCL amplitudes. At the higher mass end there is good continuity between the RLIC 77 amplitudes and those of the single-channel analysis of the CdF-S collaboration⁸ covering the range 2070 to 2440 MeV. The two strongly coupled resonances, $D03$ $\Lambda(1520)$ and $G07$ $\Lambda(2110)$, lying below and above the mass range covered by the UCL analysis, clearly provide strong

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For notation, see key at front of Listings.

Baryons
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constraints on the amplitudes.

The $\Lambda\pi$ channel: This pure $I=1$ channel, where simultaneous measurements of the polarization and differential cross section are possible, has been the subject of many wide-range energy-dependent and energy-independent analyses (For example, RLIC 77,²⁰ UCL,²² Baillon-Litchfield,²³ de Bellefont-Berthon,²⁴ and Van Horn²⁵). However, even the widespread use of the method of Barrelet zeroes has not helped to resolve the $I=1$ Y^* spectrum -- probably because most of the $I=1$ states do not couple strongly to the initial $\bar{N}K$ channel.

Quasi-2-body channels: The RLIC group has made energy-dependent analyses of the $\Lambda(1520)\pi$, $\Sigma(1385)\pi$, and $\bar{N}K^*(892)$ channels over the widest ranges for which data are available. The data are extracted from the appropriate 3-particle final states by making 4-variable fits to an incoherent superposition of quasi-2-body final states and 3-particle Lorentz-invariant phase space. The quality of the fits suggests a maximum model-dependent systematic uncertainty of 10%. The errors quoted on the Y^* couplings do not include this uncertainty. The $\Lambda\omega$ channel has been analyzed from threshold to 2440 MeV by the CdF-S collaboration.⁸

Figures: Argand plots of fifteen $S=1$ partial waves are shown in Fig. 1(a) through (k). The analyses shown were picked for illustrative purposes rather than on the basis of our judgment of their quality. For the $\bar{K}N$ channel, we show the amplitudes from RLIC 77²⁰ and from LBL-Mt. Holyoke-CERN,²¹ and for the $\Lambda\pi$ and $\Sigma\pi$ channels those from RLIC 77²⁰ and from UCL.²²

Errors on masses and widths: Often the errors quoted on resonance parameters from partial-wave analyses are only statistical, and the parameters can change by more than these errors when a different parametrization of the waves is used. Furthermore, the different analyses use more or less the same set of data, so it is not really appropriate to treat the different determinations of the resonance parameters as independent measurements or to average them together. In any case, the spread of the masses, widths, and branching fractions from the different analyses is certainly a

better indication of the uncertainties than are the quoted errors. In the Baryon Table, usually a range reflecting the spread of values obtained is given rather than a particular value with error.

For three states, the $\Lambda(1520)$, the $\Lambda(1820)$, and the $\Sigma(1775)$, there is enough information available to make an overall fit of the various branching fractions (see Sec. VII B of the main text). It is then necessary to use the quoted errors, but the errors obtained from the fit are not to be taken seriously.

II. Production Experiments

Partial-wave analyses of course separate partial waves, whereas a peak in a cross section or an invariant mass distribution usually cannot be disentangled from background and analyzed for its quantum numbers; and more than one resonance may be contributing to the peak. Results from partial-wave analyses and from production experiments are generally kept separate in the Listings, and in the Baryon Table results from production experiments are used only for the low mass states. The $\Sigma(1385)$ and $\Lambda(1405)$ of course lie below the $\bar{K}N$ threshold and everything about them comes from production experiments; and production and formation experiments agree quite well in the case of $\Lambda(1520)$ and results have been combined. There is some disagreement between production and formation experiments in the 1600-to-1700 MeV region: see the $\Sigma(1620)$ and $\Sigma(1670)$ notes for details.

The most interesting recent result from a production experiment is the observation of a narrow hyperon of mass 3170 MeV decaying to multi-strange final states containing five or six particles ($\Lambda\Sigma\bar{K}K + \text{pions}$, $\Xi K + \text{pions}$). The effect is seen in two high-statistics bubble chamber experiments studying K^+p interactions at 6.5 and 8.25 GeV/c.²⁶ The statistical significance of the peak is about 6.5 standard deviations, and the observed width is consistent with the experimental resolution. The modes with three strange particles possibly suggest an exotic (e.g., $qqqq\bar{q}$) state.

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(see also the Listings)

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Baryons

Λ 's and Σ 's

Data Card Listings

For notation, see key at front of Listings.

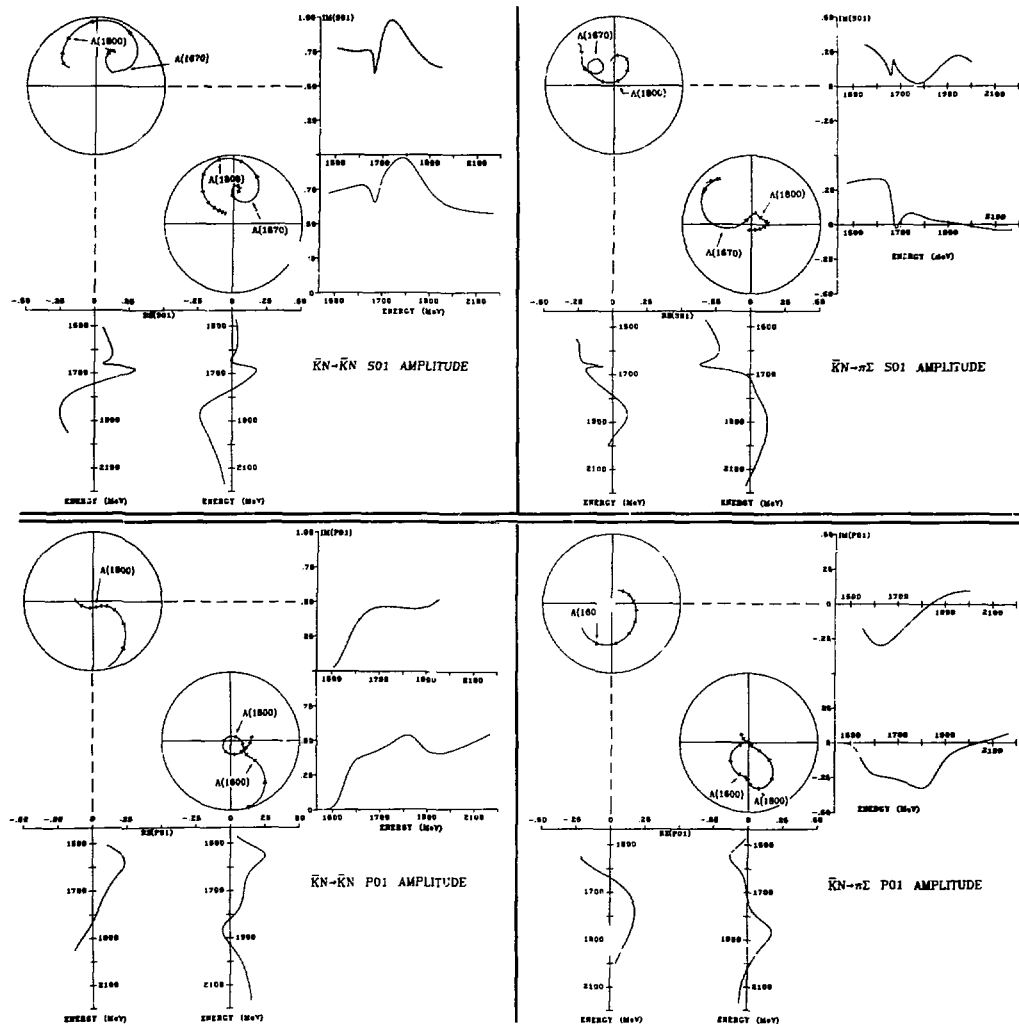


Fig. 1(a). The $L=1, J=2$ S01 and P01 partial-wave amplitudes for $\bar{K}N$ scattering in the elastic and E_{π} channels. The lower plot for each amplitude is from RLIC 77, the upper plots for the elastic amplitudes are from ALSTON 78, and the upper plots for the E_{π} amplitudes are from MARTIN 77. In the Argand plots, the ticks are at integral multiples of 50 MeV, and the established resonances are shown at their nominal positions [the S01 $A(1405)$ is of course below threshold and is not shown]. The real and imaginary parts of the amplitudes as functions of energy are shown projected in alignment with the Argand plots.

Data Card Listings

For notation, see key at front of Listings.

Baryons
 Λ 's and Σ 's

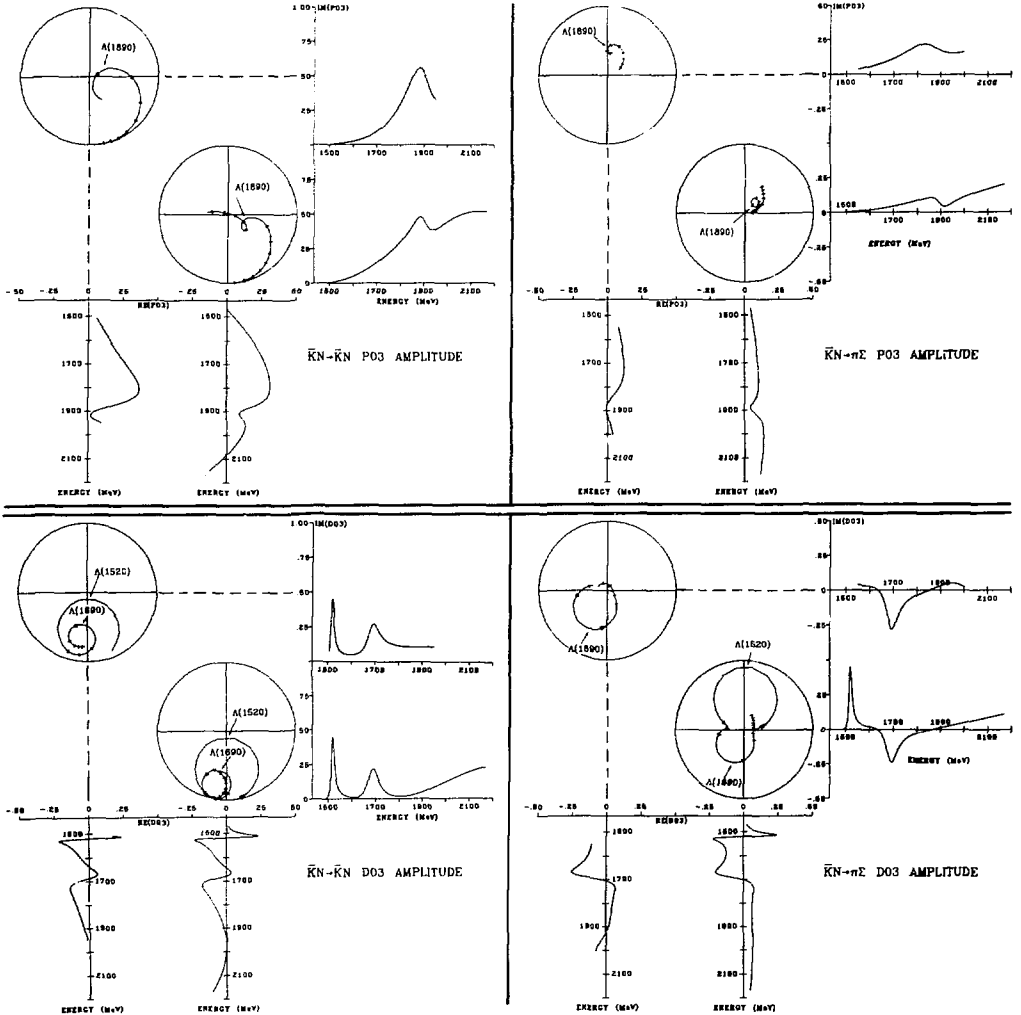


Fig. 1(b). The L-I-2J = P03 and D03 partial-wave amplitudes for $\bar{K}N$ scattering in the elastic and $\Sigma\pi$ channels. The lower plot for each amplitude is from RLIC 77, the upper plots for the elastic amplitudes are from ALSTON 78, and the upper plots for the $\Sigma\pi$ amplitudes are from MARTIN 77. In the Argand plots, the ticks are at integral multiples of 50 MeV, and the established resonances are shown at their nominal positions. The real and imaginary parts of the amplitudes as functions of energy are shown projected in alignment with the Argand plots.

Baryons

Λ 's and Σ 's

Data Card Listings

For notation, see key at front of Listings.

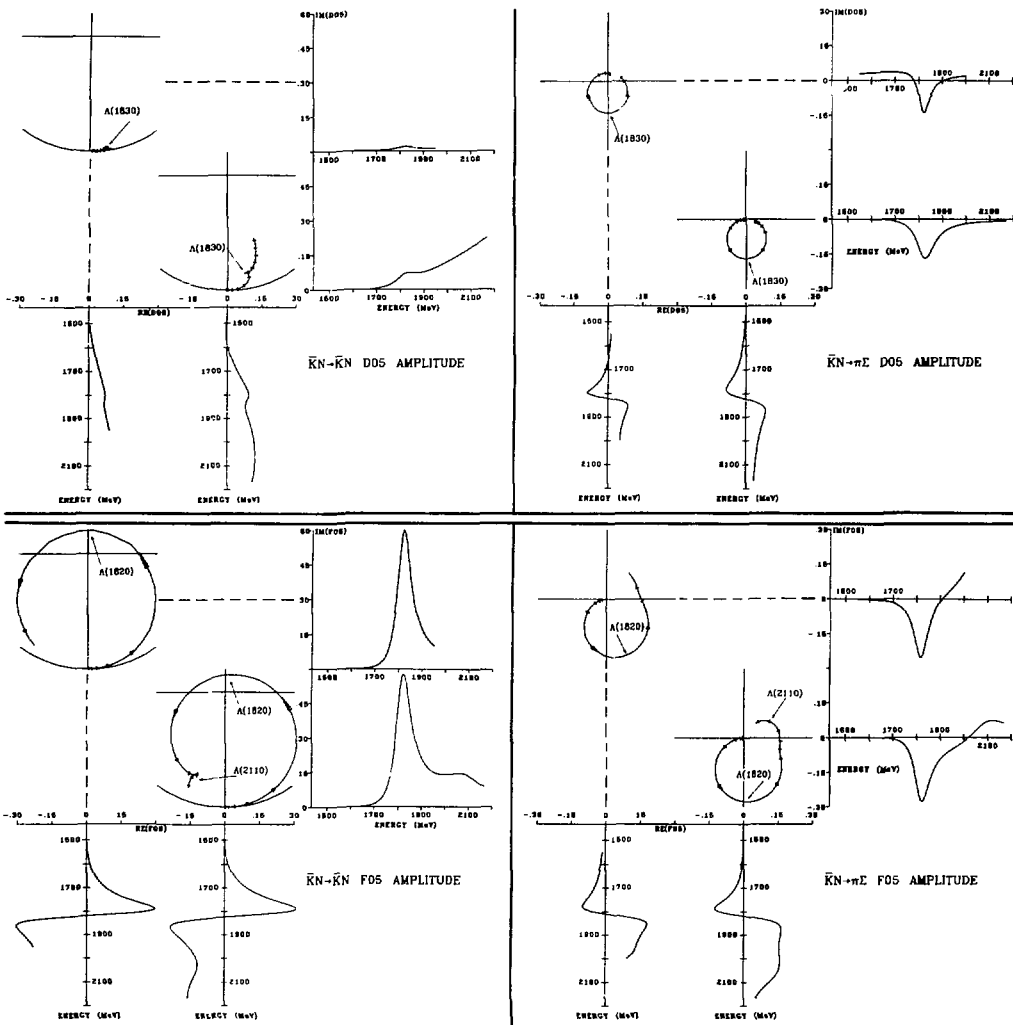


Fig. 1(c). The $L=I=2J = D05$ and $F05$ partial-wave amplitudes for $\bar{K}N$ scattering in the elastic and E_n channels. The lower plot for each amplitude is from RLIC 77, the upper plots for the elastic amplitudes are from ALSTON 78, and the upper plots for the E_n amplitudes are from MARTIN 77. In the Argand plots, the ticks are at integral multiples of 50 MeV, and the established resonances are shown at their nominal positions. The real and imaginary parts of the amplitudes as functions of energy are shown projected in alignment with the Argand plots.

Data Card Listings

For notation, see key at front of Listings.

Baryons

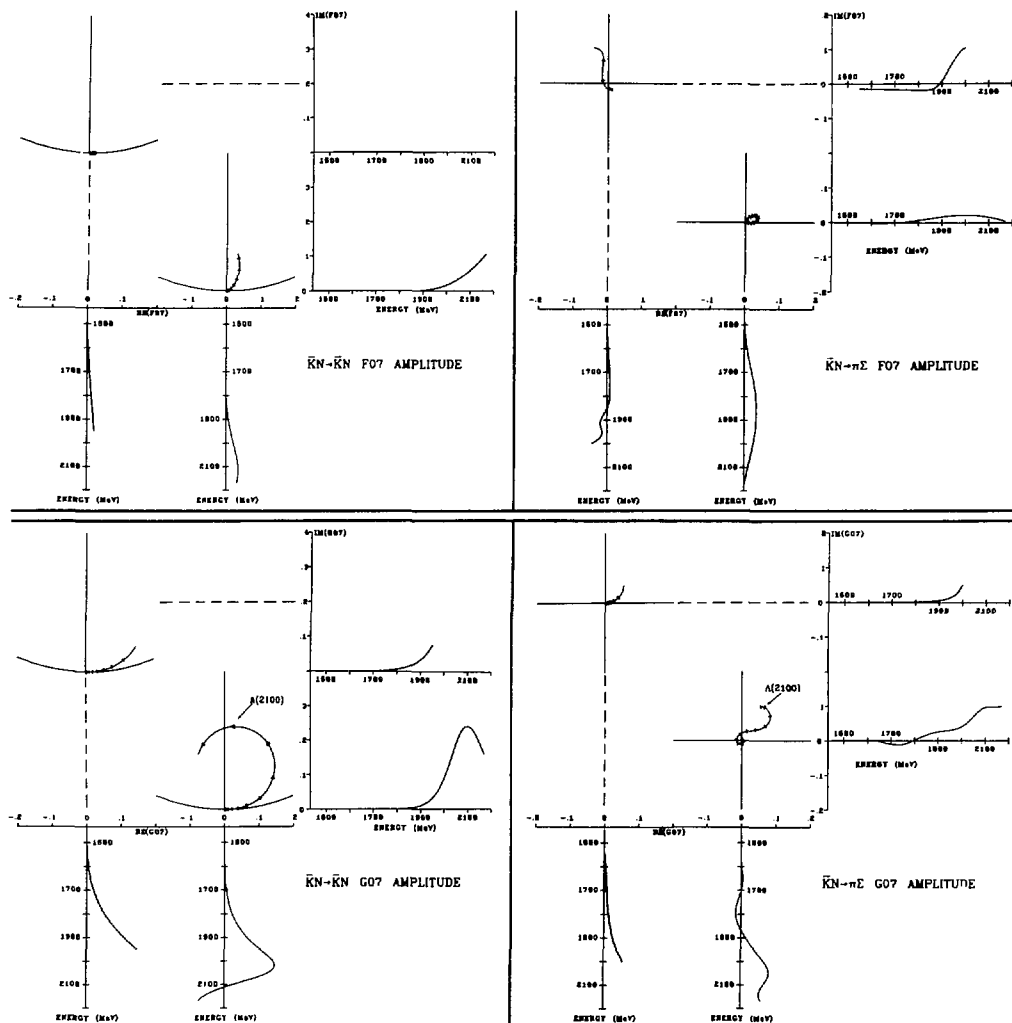
 Λ 's and Σ 's

Fig. 1(d). The $L=1, J=0$ partial-wave amplitudes for $\bar{K}N$ scattering in the elastic and $\Sigma\pi$ channels. The lower plot for each amplitude is from RLIC 77, the upper plots for the elastic amplitudes are from ALSTON 78, and the upper plots for the $\Sigma\pi$ amplitudes are from MARTIN 77. In the Argand plots, the ticks are at integral multiples of 50 MeV, and the established resonance is shown at its nominal position. The real and imaginary parts of the amplitudes as functions of energy are shown projected in alignment with the Argand plots.

Baryons

Λ 's and Σ 's

Data Card Listings
 For notation, see key at front of Listings.

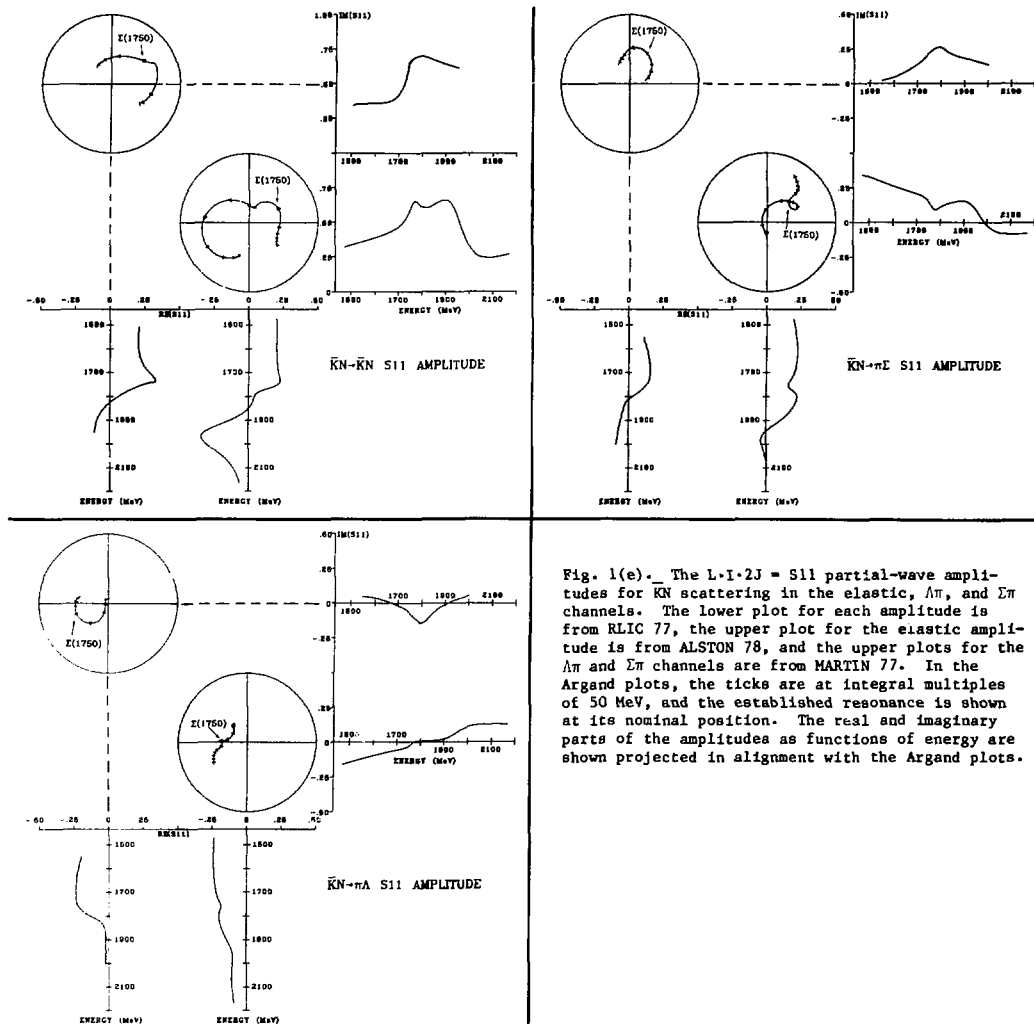


Fig. 1(e). The $L=1, J=2$ = S₁₁ partial-wave amplitudes for $\bar{K}N$ scattering in the elastic, $\Lambda\pi$, and $\Sigma\pi$ channels. The lower plot for each amplitude is from RLIC 77, the upper plot for the elastic amplitude is from ALSTON 78, and the upper plots for the $\Lambda\pi$ and $\Sigma\pi$ channels are from MARTIN 77. In the Argand plots, the ticks are at integral multiples of 50 MeV, and the established resonance is shown at its nominal position. The real and imaginary parts of the amplitudes as functions of energy are shown projected in alignment with the Argand plots.

Data Card Listings

For notation, see key at front of Listings.

Baryons

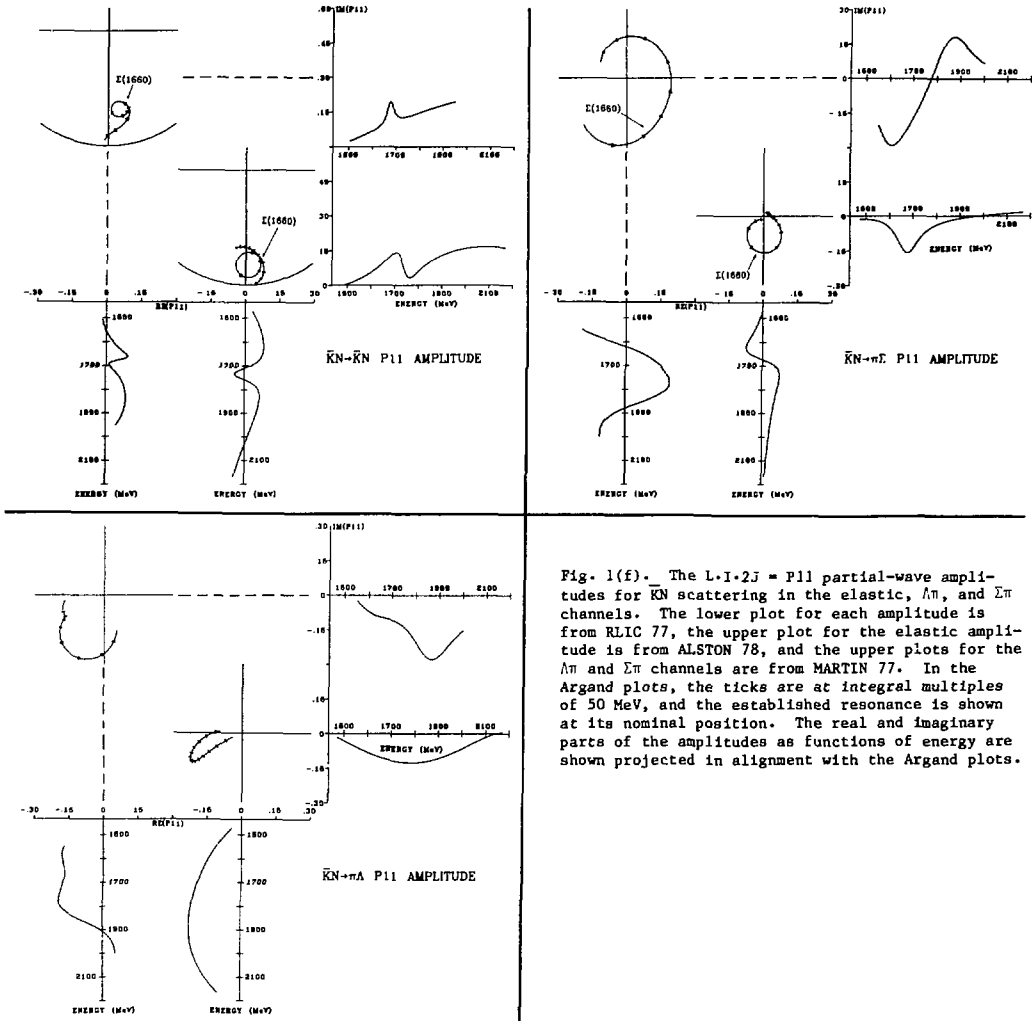
 Λ 's and Σ 's

Fig. 1(f). The $L=1, J=1$ partial-wave amplitudes for $\bar{K}N$ scattering in the elastic, $\Lambda\pi$, and $\Sigma\pi$ channels. The lower plot for each amplitude is from RLIC 77, the upper plot for the elastic amplitude is from ALSTON 78, and the upper plots for the $\Lambda\pi$ and $\Sigma\pi$ channels are from MARTIN 77. In the Argand plots, the ticks are at integral multiples of 50 MeV, and the established resonance is shown at its nominal position. The real and imaginary parts of the amplitudes as functions of energy are shown projected in alignment with the Argand plots.

Baryons

Λ 's and Σ 's

Data Card Listings

For notation, see key at front of Listings.

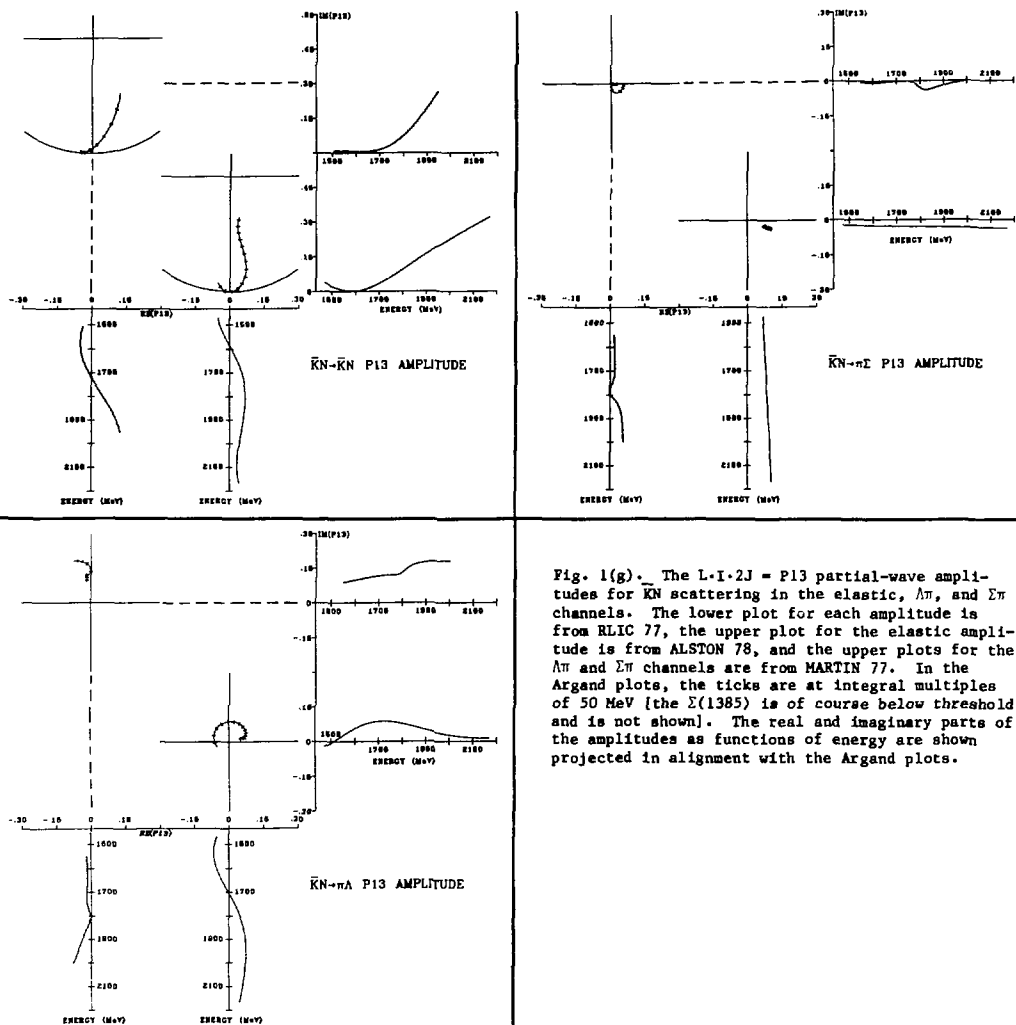


Fig. 1(g). The $L=1, J=2$ = P_{13} partial-wave amplitudes for $\bar{K}N$ scattering in the elastic, $\Lambda\pi$, and $\Sigma\pi$ channels. The lower plot for each amplitude is from ALSTON 78, the upper plot for the elastic amplitude is from ALSTON 78, and the upper plots for the $\Lambda\pi$ and $\Sigma\pi$ channels are from MARTIN 77. In the Argand plots, the ticks are at integral multiples of 50 MeV [the $\Sigma(1385)$ is of course below threshold and is not shown]. The real and imaginary parts of the amplitudes as functions of energy are shown projected in alignment with the Argand plots.

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For notation, see key at front of Listings.

Baryons
 Λ 's and Σ 's

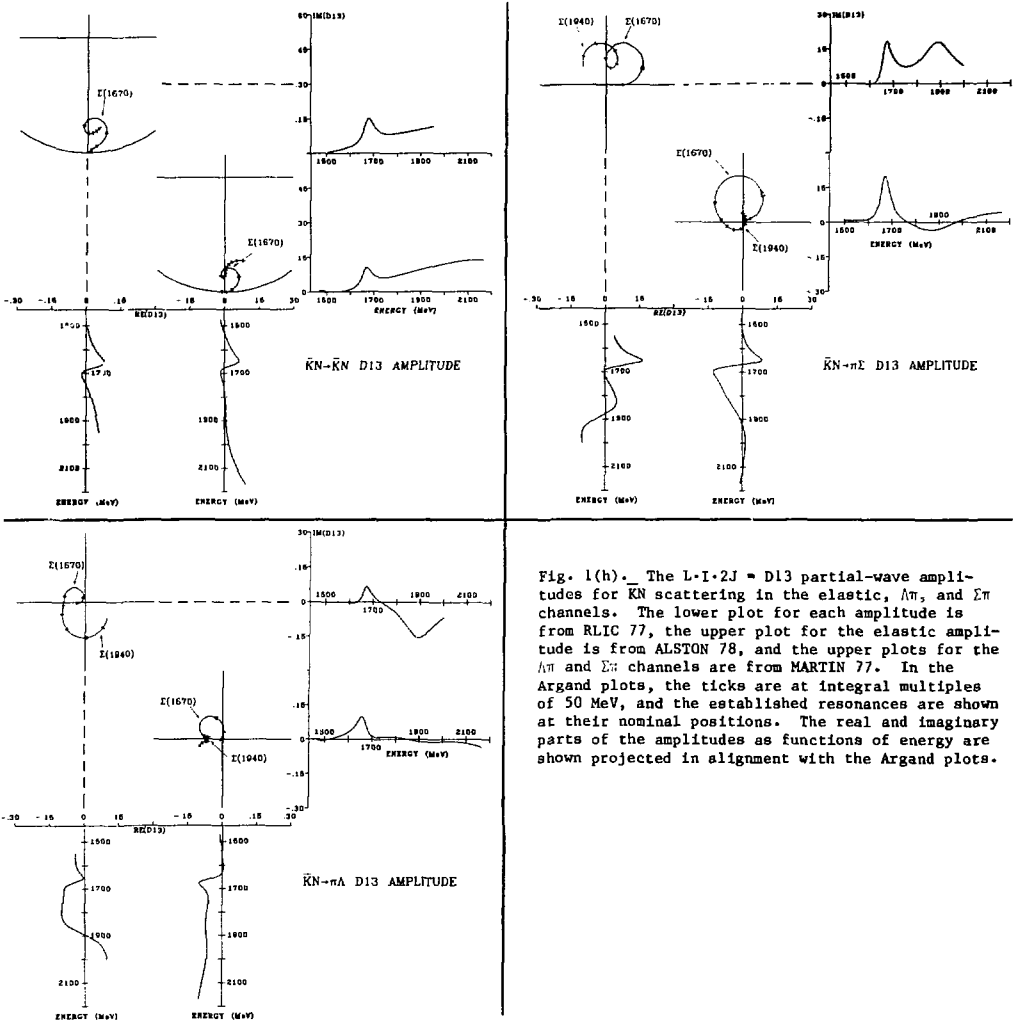


Fig. 1(h). The $L=I=2J=D_{13}$ partial-wave amplitudes for KN scattering in the elastic, $\Lambda\pi$, and $\Sigma\pi$ channels. The lower plot for each amplitude is from RLIC 77, the upper plot for the elastic amplitude is from ALSTON 78, and the upper plots for the $\Lambda\pi$ and $\Sigma\pi$ channels are from MARTIN 77. In the Argand plots, the ticks are at integral multiples of 50 MeV, and the established resonances are shown at their nominal positions. The real and imaginary parts of the amplitudes as functions of energy are shown projected in alignment with the Argand plots.

Baryons

Λ 's and Σ 's

Data Card Listings

For notation, see key at front of Listings.

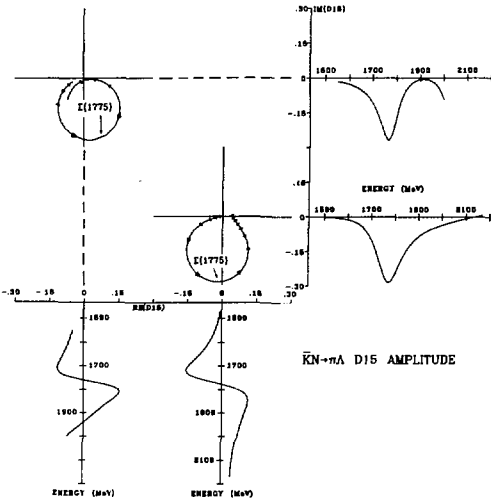
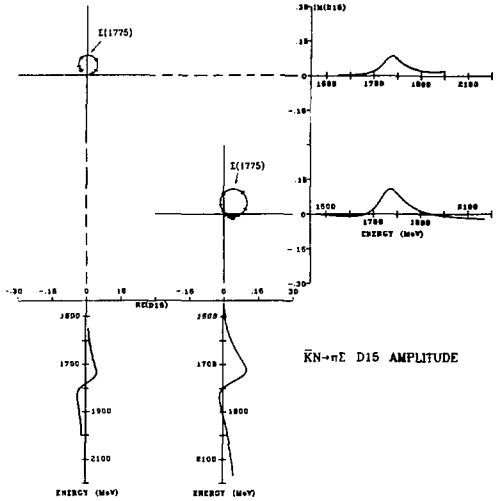
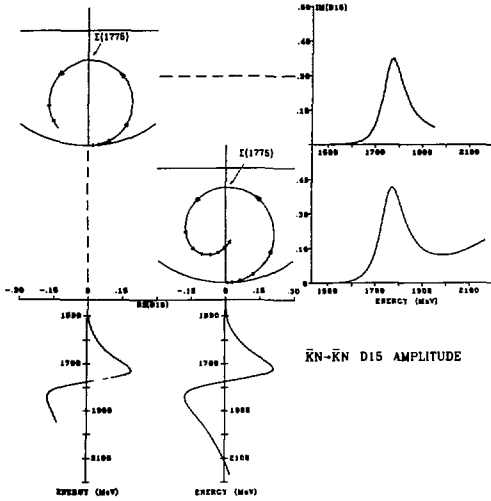


Fig. 1(i). The $L \cdot I \cdot 2J = D15$ partial-wave amplitudes for $\bar{K}N$ scattering in the elastic, $\Lambda\pi$, and $\Sigma\pi$ channels. The lower plot for each amplitude is from RLIC 77, the upper plot for the elastic amplitude is from ALSTON 78, and the upper plots for the $\Lambda\pi$ and $\Sigma\pi$ channels are from MARTIN 77. In the Argand plots, the ticks are at integral multiples of 50 MeV, and the established resonance is shown at its nominal position. The real and imaginary parts of the amplitudes as functions of energy are shown projected in alignment with the Argand plots.

Data Card Listings

For notation, see key at front of Listings.

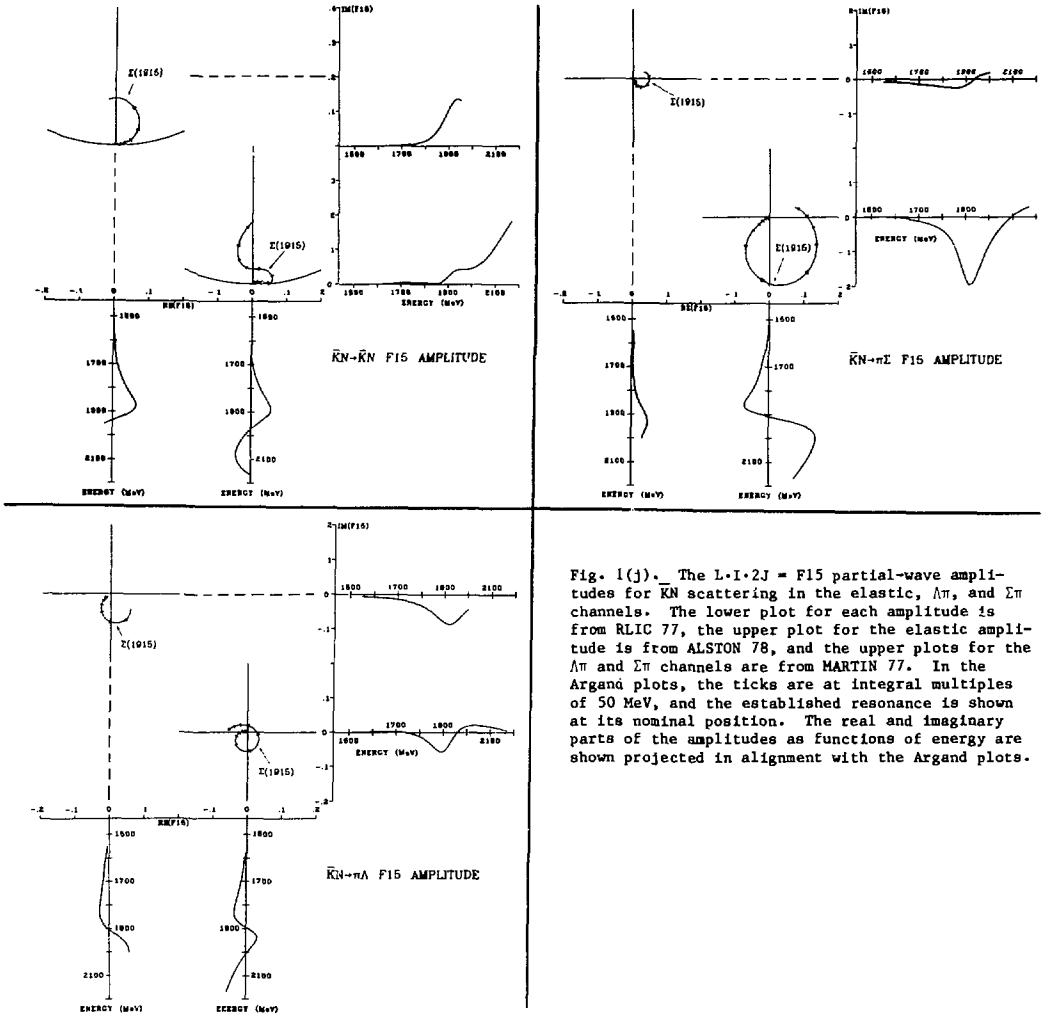
Baryons
 Λ 's and Σ 's

Fig. 1(j). The $L=1, J=2$ = F_{15} partial-wave amplitudes for $\bar{K}N$ scattering in the elastic, $\Lambda\pi$, and $\Sigma\pi$ channels. The lower plot for each amplitude is from RLIC 77, the upper plot for the elastic amplitude is from ALSTON 78, and the upper plots for the $\Lambda\pi$ and $\Sigma\pi$ channels are from MARTIN 77. In the Argand plots, the ticks are at integral multiples of 50 MeV, and the established resonance is shown at its nominal position. The real and imaginary parts of the amplitudes as functions of energy are shown projected in alignment with the Argand plots.

Baryons

Λ 's and Σ 's

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For notation, see key at front of Listings.

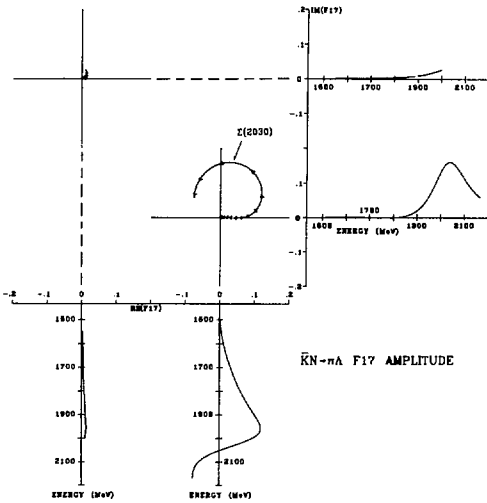
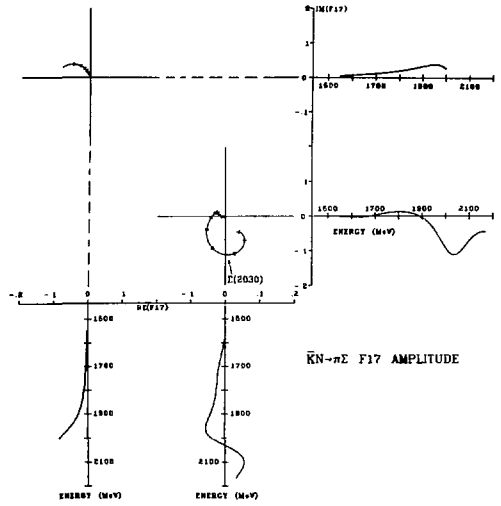
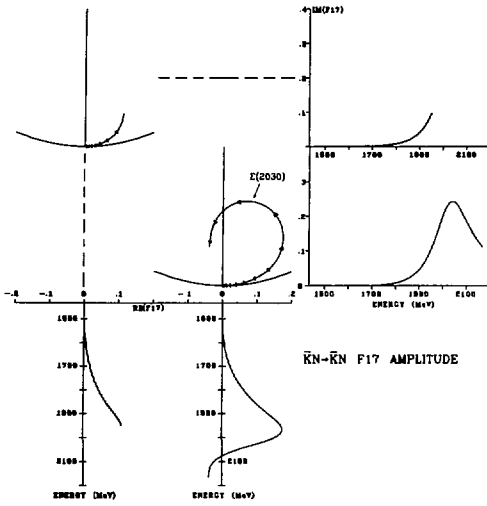


Fig. 1(k). The $L=1, I=2, J=1$ partial-wave amplitudes for $\bar{K}N$ scattering in the elastic, $\Lambda\pi$, and $\Sigma\pi$ channels. The lower plot for each amplitude is from ALSTON 78, and the upper plots for the $\Lambda\pi$ and $\Sigma\pi$ channels are from MARTIN 77. In the Argand plots, the ticks are at integral multiples of 50 MeV, and the established resonance is shown at its nominal position. The real and imaginary parts of the amplitudes as functions of energy are shown projected in alignment with the Argand plots.

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 $S=-1$ $I=0$ HYPERON STATES (Λ)**A**

1R LAMSDA(11110, JP=1/2+) 1-0
SEE STABLE PARTICLE DATA CARD LISTINGS

A(1405)**S₀₁**

THIS RESONANCE CAN BE IDENTIFIED WITH THE VIRTUAL Λ QUANTUM STATE IN THE Λ - \bar{N} SYSTEM FOUND IN THE ANALYSIS OF Λ CH ENERGY Λ - \bar{N} INTERACTION. WE LIST ONLY EXPERIMENTS SEPARATELY BELOW. WE USE ONLY PRODUCTION EXPERIMENTS FOR AVERAGING OF MASSES AND WIDTHS.

27 Y001405, JP=1/2- 1-0

M	11405.01	ALSTON	63 HRC	K-P	1.15 MEV/C
M	11410.01	ALEXANDER	62 HRC	P1-P	2.1 BEV/C
M	11405.01	ALSTON	62 HRC	K-P	1.12-1.9 BEV/C
M	11392.01	18.01	ENGLER	63 HRC	P1-P, D14 1.08
M	1400-0	24.0	MUSGRAVE	65 HRC	PHYS 9-34 REVZ 7766
M	87 1409.0	5.0	BIRMINGHAM	66 HRC	K-P 1.5
M	220 1409.0	5.0	GALTIERI	68 HRC	K-D 2.1-2.74EVCZ 6768
M	AVG	1402.4	3.5	AVERAGE ERROR INCLUDES SCALE FACTOR OF 1.01	

27 Y001405, WIDTH (MEV) (PROD. EXP.)

M	170.01	ALSTON	61 HRC		7766
M	75.0	ALEXANDER	62 HRC		
M	150.01 <td>ALSTON</td> <td>62 HRC</td> <td></td> <td></td>	ALSTON	62 HRC		
M	189.01	120.01	ENGLER	65 HRC	7766
M	60.0	20.0	MUSGRAVE	65 HRC	7766
M	67 50.0	10.0	BIRMINGHAM	66 HRC	K-P 1.5
M	120 35.0	8.0	GALTIERI	68 HRC	K-D 2.1-2.74EVCZ 6768
M	AVG	38.1	3.9	AVERAGE ERROR INCLUDES SCALE FACTOR OF 1.01	

27 Y001405, PARTIAL DECAY MODES (PURE, EXP.)

D1	Y001405	INTO SIGMA P1	1300-179
REFERENCES FOR Y001405 (PURE, EXP.)			
ALSTON	63 P1	0 698	*ALVAREZ, CERDAS, GODO, GORTALEZ + (P1) F 1
ALEXANDER	62 P1	8 447	ALEXANDER, KALB, LITSCHE, MILLER + (P1) F 1
ALSTON	62 CERN CONF	311	*ALVAREZ, FERRO-LUZZI, PIRESEVILI + (P1) F 1
ENGLER	65 P1	15 224	*ISKRAKRAMER, MELTZER, MUSGRAVE + (P1) F 1
MUSGRAVE	65 NC	35 735	*PIMENTAS + (P1) F 1
BIRMINGHAM	66 PP	152 1146	BIRMINGHAM, GLASCO, LILES, OBERDORF, RUTHERFORD
GALTIERI	68 PP	21 573	BARBARD-GALTIERI, CHAUDHURI + (P1) F 1

1405 MEV REGION: EXTRAPOLATIONS BELOW THRESHOLD

24 Y001405, JP=1/2- 1-0

EXTRAPOLATION BELOW THRESHOLD

SEE NOTE IN Y001405 PRODUCTION EXPERIMENTS. THE DIFFERENCES IN EXTRAPOLATING FROM THE PARTIAL DECAY TO THE RESONANCE LOCATION ARE DISCUSSED BY GALTIERI ET AL.

THE QUESTION OF WHETHER $\Lambda(1405)$ IS A Λ - \bar{N} BOUND STATE OR A CDD POLE (DALLIZ 70, RAJASEKARAN 72) HAS BEEN INVESTIGATED BY CLINE ET AL., MARTIN 71, GALTIERI 72, AND OBERDORF 72. THE LAST TWO PAPERS CONCLUDE THAT THE DATA CANNOT TELL THE DIFFERENCE.

THE $(\Lambda$ - $\bar{N})$ (PI SIGMA) COUPLING RATIO IS DISCUSSED BY OBERDORF 72.

24 Y001405, MASS (MEV)

M	1410.7	11.01	ENGLER <th>65 HRC</th> <th>0-FF-RANGE F11</th> <th>7766</th>	65 HRC	0-FF-RANGE F11	7766
M	1409.6	11.71	SAKETT	65 HRC <td>0-FF-RANGE F11 <td>7766</td> </td>	0-FF-RANGE F11 <td>7766</td>	7766
DATA OF SAKETT ARE USED IN FIT BY KITTLES						
M	1407.5	11.21	KITTLES	60 HRC <td>0-FF-RANGE F11 <td>7766</td> </td>	0-FF-RANGE F11 <td>7766</td>	7766
M	1403.0	13.01	ENGLER	67 HRC <td>K MATRIX FIT (KPI)</td> <td>8768</td>	K MATRIX FIT (KPI)	8768
M	1416.0	14.01	MARTIN	69 HRC <td>CONST. K MATRIX</td> <td>10769</td>	CONST. K MATRIX	10769
M	1421.01		MARTIN	70 PVI	CONST. K MATRIX	6770
M	1	1406.1	CHAD	73 DPVA	0-RNG. FIT, SCL B	9773
M	1					9773
SEE ALSO THE ACCOMPANYING PAPER BY THEM 54573.						

24 Y001405, WIDTH (MEV)

M	37.0	13.23	ENGLER <th>65 HRC</th> <th></th> <th>7766</th>	65 HRC		7766
M	28.7	14.11	SAKETT	65 HRC <td></td> <td>7766</td>		7766
M	34.1 <th>14.11</th> <td>KITTLES</td> <td>60 HRC <td></td> <td>7766</td> </td>	14.11	KITTLES	60 HRC <td></td> <td>7766</td>		7766
M	50.0	15.01	ENGLER	67 HRC <td>K MATRIX FIT (KPI)</td> <td>8768</td>	K MATRIX FIT (KPI)	8768
M	29.0	16.01	MUSGRAVE	69 HRC <td>CONST. K MATRIX</td> <td>10769</td>	CONST. K MATRIX	10769
M	120.01		MARTIN	70 PVI	CONST. K MATRIX	6770
M	1	1406.1	CHAD	73 DPVA	0-RNG. FIT, SCL B	9773
M	1					9773
ASYMMETRIC RESONANCE SHAPE, 1427.41 MEV BELOW RESONANCE, 14 MEV ABOVE. 9773						

Data Card Listings
For notation, see key at front of Listings.

Baryons
A(1800)

A(1800)

S01

77 Y00E1800, JP=177-1-10

THE FOLLOWING SHOWS A TABLE FOR EACH SECOND COMPONENT SEPARATED BY THE PRODUCTION NEW SECTION. THERE ARE ALSO SECTIONS WHICH SHOW THE MASS, WIDTH, AND COUPLING DETERMINATIONS.

Table with columns for mass (MEV), width (MEV), and other parameters. Includes rows for SPECTRA and AMPLITUDE.

Table with columns for mass (MEV), width (MEV), and other parameters. Includes rows for SPECTRA and AMPLITUDE.

Table with columns for mass (MEV), width (MEV), and other parameters. Includes rows for SPECTRA and AMPLITUDE.

Table with columns for mass (MEV), width (MEV), and other parameters. Includes rows for SPECTRA and AMPLITUDE.

Table with columns for mass (MEV), width (MEV), and other parameters. Includes rows for SPECTRA and AMPLITUDE.

Table with columns for mass (MEV), width (MEV), and other parameters. Includes rows for SPECTRA and AMPLITUDE.

A(1800)

P01

77 Y00E1800, JP=177-1-10

SEE THE MINI-REVIEW AT THE START OF THE ** LISTINGS.

THE EVIDENCE FOR THIS STATE IS SOMEWHAT CONFUSED, IT WAS FIRST SUGGESTED IN A PARTIAL WAVE ANALYSIS OF K800 N DATA BY THE BEHAVIOUR OF THE POI AMPLITUDE WHEN IT WAS DETERMINED TO BE A TWO-STATE-TO-ONE-STATE BACKGROUND TRANSITION.

ALMOST ALL THE DEBRIS ANALYSES CONTAIN A POI STATE, AND SOMETIMES INCLUDE THE MASS, WIDTH, AND BRANCHING RATIOS OBTAINED IN THE DIFFERENT ANALYSES VARY GREATLY. SEE ALSO THE Y00E1800POI LISTING.

Table with columns for mass (MEV), width (MEV), and other parameters. Includes rows for SPECTRA and AMPLITUDE.

Table with columns for mass (MEV), width (MEV), and other parameters. Includes rows for SPECTRA and AMPLITUDE.

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Table with columns for mass (MEV), width (MEV), and other parameters. Includes rows for SPECTRA and AMPLITUDE.

Table with columns for mass (MEV), width (MEV), and other parameters. Includes rows for SPECTRA and AMPLITUDE.

Table with columns for mass (MEV), width (MEV), and other parameters. Includes rows for SPECTRA and AMPLITUDE.

Table with columns for mass (MEV), width (MEV), and other parameters. Includes rows for SPECTRA and AMPLITUDE.

Baryons

$\Lambda(1800)$, $\Lambda(1820)$

Data Card Listings

For notation, see key at front of Listings.

177 127 Y0018001 WIDTH (MEV) 1770
 M 2 1817.0 (0.1) WIDTH 77 DPMW KBAR N MULTICANAL 1770

102 Y0018001 PARTIAL DECAY MODES 1770

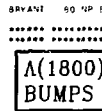
71 Y0018001 INT'L $\Lambda(1800)$ DEBRAY MASSES 630+493
 P2 Y0018001 INT'L $\Lambda(1800)$ 1197+130

102 Y0018001 BRANCHING RATIOS 1770

P1 Y0018001 INT'L LAMBDA PI/3 TOTAL (P1) 1770
 P1 C (70.04) (70.01) RLIC 77 DPMW KBAR N MULTICANAL 1770
 P1 C (KBAR N)/TOTAL FROM P1 L1 77 IS SUPERSEDED BY DPMW R0.
 R1 C 0.03 (70.23) GOPAL 80 DPMW KBAR N ELASTIC 12701*

P2 Y0018001 FROM KBAR N INTO SIGMA PI (SQRTP10P2) 1770
 P2 LESS THAN 0.04 RLIC 77 DPMW KBAR N MULTICANAL 1770

REFERENCES FOR Y0018001
 RLIC 77 NP 3319 167 TOTAL MODES FROM KBAR N (EMERSON) 110E+10411
 GOPAL 80 F0RCE3 COM 150 SUP GOPAL 150 F0RCE3 10411



39 Y0018001 PARTIAL DECAY MODES 1770

P1 Y0018001 INT'L $\Lambda(1800)$ DEBRAY MASSES 630+493
 P2 Y0018001 INT'L SIGMA PI 1197+130
 P3 Y0018001 INT'L SIGMA PI 1197+130
 P4 Y0018001 INT'L LAMBDA PI/3 TOTAL 548+1115
 P5 Y0018001 INT'L $\Lambda(1800)$ PI/3 TOTAL 130+1395
 P6 Y0018001 INT'L $\Lambda(1800)$ PI/3 TOTAL 130+1395

PRODUCION EXPERIMENTS
 LOCKMAN 78 OBSERVE A Λ STD. DEV. ENHANCEMENT IN THE
 LAMBDA PI/3 MASS SPECTRUM FROM THE REACTION
 PP \rightarrow LAMBDA PI/3 PI \rightarrow ANYTHING IN A Λ (SEN 159
 EXPERIMENT AT CM L. ENERGIES OF 43 AND 67 GEV.
 THE MAIN DECAY MODES APPEAR TO BE $\Lambda \rightarrow$ SIGMA PI AND
 $\Lambda \rightarrow$ LAMBDA PI (SEE THE ENTRY FOR Y0018001).
 THE LAMBDA PI IS WELL ESTABLISHED, BUT SINCE THE LAMBDA PI DECAY IS NOT
 OBSERVED, IT IS PROBABLY NEGLECTED.

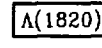
119 Y0018001 MASS (MEV) (PROD. EXP.)
 M 60 1907. 3. (LOCKMAN 78 DEF 0 PP TO L PI PI) 12770

110 Y0018001 WIDTH (MEV) (PROD. EXP.)
 M 60 74. 8. (LOCKMAN 78 DEF 0 PP TO L PI PI) 12770
 C 60 74. 8. (LOCKMAN 78 DEF 0 PP TO L PI PI) 12770
 C OBSERVED WIDTH CONSISTENT WITH EXPERIMENTAL RESOLUTION.

116 Y0018001 PARTIAL DECAY MODES (PROD. EXP.)
 P1 Y0018001 INT'L LAMBDA PI/3 TOTAL 1135+130+130 DEBRAY MASSES
 P2 Y0018001 INT'L $\Lambda(1800)$ PI/3 TOTAL 130+130
 P3 Y0018001 INT'L $\Lambda(1800)$ PI/3 TOTAL 130+130

119 Y0018001 BRANCHING RATIOS (PROD. EXP.)
 R1 Y0018001 L C (LAMBDA PI/3)/TOTAL (P1) 1770
 R1 SFR4 (LOCKMAN 78 SFC 0 PP TO L PI PI) 12770
 R2 Y0018001 INT'L $\Lambda(1800)$ PI/3 TOTAL (P2) 1770
 R2 SFR4 (LOCKMAN 78 REC 0 PP TO L PI PI) 12770
 R3 Y0018001 INT'L $\Lambda(1800)$ PI/3 TOTAL (P3) 1770
 R3 SFR4 (LOCKMAN 78 SFC 0 PP TO L PI PI) 12770

REFERENCES FOR Y0018001 (PROD. EXP.)
 LOCKMAN 78 GEN DPMW 78-01 (EMERSON, PANDEP, POSTER, SCHLEIN) IUCFA+SIELI



39 Y0018201 PARTIAL DECAY MODES 1770

71 Y0018201 INT'L $\Lambda(1820)$ DEBRAY MASSES 630+493
 P2 Y0018201 INT'L SIGMA PI 1197+130
 P3 Y0018201 INT'L SIGMA PI 1197+130
 P4 Y0018201 INT'L LAMBDA PI/3 TOTAL 548+1115
 P5 Y0018201 INT'L $\Lambda(1820)$ PI/3 TOTAL 130+1395
 P6 Y0018201 INT'L $\Lambda(1820)$ PI/3 TOTAL 130+1395

PRODUCION EXPERIMENTS
 LOCKMAN 78 OBSERVE A Λ STD. DEV. ENHANCEMENT IN THE
 LAMBDA PI/3 MASS SPECTRUM FROM THE REACTION
 PP \rightarrow LAMBDA PI/3 PI \rightarrow ANYTHING IN A Λ (SEN 159
 EXPERIMENT AT CM L. ENERGIES OF 43 AND 67 GEV.
 THE MAIN DECAY MODES APPEAR TO BE $\Lambda \rightarrow$ SIGMA PI AND
 $\Lambda \rightarrow$ LAMBDA PI (SEE THE ENTRY FOR Y0018001).
 THE LAMBDA PI IS WELL ESTABLISHED, BUT SINCE THE LAMBDA PI DECAY IS NOT
 OBSERVED, IT IS PROBABLY NEGLECTED.

39 Y0018201 BRANCHING RATIOS 1770
 R1 Y0018201 INT'L LAMBDA PI/3 TOTAL (P1) 1770
 R1 SFR4 (LOCKMAN 78 SFC 0 PP TO L PI PI) 12770
 R2 Y0018201 INT'L $\Lambda(1820)$ PI/3 TOTAL (P2) 1770
 R2 SFR4 (LOCKMAN 78 REC 0 PP TO L PI PI) 12770
 R3 Y0018201 INT'L $\Lambda(1820)$ PI/3 TOTAL (P3) 1770
 R3 SFR4 (LOCKMAN 78 SFC 0 PP TO L PI PI) 12770

REFERENCES FOR Y0018201 (PROD. EXP.)
 LOCKMAN 78 GEN DPMW 78-01 (EMERSON, PANDEP, POSTER, SCHLEIN) IUCFA+SIELI

71 1917.0 (0.1) WIDTH 77 DPMW KBAR N MULTICANAL 1770
 M 2 THE TWO ENTRIES FOR WIDTH 77 CORRESPOND TO EXTRAPOLATION OF SLOPE
 2 PARAMETERS FROM THE T-MATRIX PILE AND FROM A R.W. FIT, RESPECTIVELY.
 M 4 1822. (2.1) RLIC 77 DPMW KBAR N MULTICANAL 1770
 M 1819. (2.0) ALSTON 78 DPMW KBAR N ELASTIC 12701
 M 1823.0 (3.0) GOPAL 80 DPMW KBAR N ELASTIC 12701*
 M 4 EPROP STATIST. ONLY- NO ERROR DUE TO PARTICULAR P.W. ANAL. 1770

39 Y0018201 PARTIAL DECAY MODES 1770

P1 Y0018201 INT'L $\Lambda(1820)$ DEBRAY MASSES 630+493
 P2 Y0018201 INT'L SIGMA PI 1197+130
 P3 Y0018201 INT'L SIGMA PI 1197+130
 P4 Y0018201 INT'L LAMBDA PI/3 TOTAL 548+1115
 P5 Y0018201 INT'L $\Lambda(1820)$ PI/3 TOTAL 130+1395
 P6 Y0018201 INT'L $\Lambda(1820)$ PI/3 TOTAL 130+1395

PRODUCION EXPERIMENTS
 LOCKMAN 78 OBSERVE A Λ STD. DEV. ENHANCEMENT IN THE
 LAMBDA PI/3 MASS SPECTRUM FROM THE REACTION
 PP \rightarrow LAMBDA PI/3 PI \rightarrow ANYTHING IN A Λ (SEN 159
 EXPERIMENT AT CM L. ENERGIES OF 43 AND 67 GEV.
 THE MAIN DECAY MODES APPEAR TO BE $\Lambda \rightarrow$ SIGMA PI AND
 $\Lambda \rightarrow$ LAMBDA PI (SEE THE ENTRY FOR Y0018001).
 THE LAMBDA PI IS WELL ESTABLISHED, BUT SINCE THE LAMBDA PI DECAY IS NOT
 OBSERVED, IT IS PROBABLY NEGLECTED.

39 Y0018201 BRANCHING RATIOS 1770

R1 Y0018201 INT'L LAMBDA PI/3 TOTAL (P1) 1770
 R1 SFR4 (LOCKMAN 78 SFC 0 PP TO L PI PI) 12770
 R2 Y0018201 INT'L $\Lambda(1820)$ PI/3 TOTAL (P2) 1770
 R2 SFR4 (LOCKMAN 78 REC 0 PP TO L PI PI) 12770
 R3 Y0018201 INT'L $\Lambda(1820)$ PI/3 TOTAL (P3) 1770
 R3 SFR4 (LOCKMAN 78 SFC 0 PP TO L PI PI) 12770

REFERENCES FOR Y0018201 (PROD. EXP.)
 LOCKMAN 78 GEN DPMW 78-01 (EMERSON, PANDEP, POSTER, SCHLEIN) IUCFA+SIELI

AVERAGE SCALE FACTOR OF 1.01
 AVG MOD 0.2642 0.0082 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.01)
 FIT 0.2633 0.0082 FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.01)

39 Y0018201 BRANCHING RATIOS 1770

R1 Y0018201 INT'L LAMBDA PI/3 TOTAL (P1) 1770
 R1 SFR4 (LOCKMAN 78 SFC 0 PP TO L PI PI) 12770
 R2 Y0018201 INT'L $\Lambda(1820)$ PI/3 TOTAL (P2) 1770
 R2 SFR4 (LOCKMAN 78 REC 0 PP TO L PI PI) 12770
 R3 Y0018201 INT'L $\Lambda(1820)$ PI/3 TOTAL (P3) 1770
 R3 SFR4 (LOCKMAN 78 SFC 0 PP TO L PI PI) 12770

REFERENCES FOR Y0018201 (PROD. EXP.)
 LOCKMAN 78 GEN DPMW 78-01 (EMERSON, PANDEP, POSTER, SCHLEIN) IUCFA+SIELI

AVERAGE SCALE FACTOR OF 1.01
 AVG MOD 0.2642 0.0082 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.01)
 FIT 0.2633 0.0082 FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.01)

Baryons

A(1890), A(2000), A(2020)

Data Card Listings

For notation, see key at front of Listings.

Table with columns for particle type (A, FOT), energy (E), width (W), and various decay parameters (ARMENTEROS, BUGG, BICMAN, etc.).

Table titled '60 Y(401890) PARTIAL DECAY MODES' showing decay channels like INTIC N KBAR, INTO LAMBDA OMEGA, etc.

Table titled '60 Y(401890) BRANCHING RATIOS' listing branching ratios for various decay modes.

Table titled '60 Y(401890) INTIC N SIGMA PI' listing partial decay widths and branching ratios.

Table titled 'REFERENCES FOR Y(401890)' listing references for various baryons like ARMENTEROS, BUGG, BICMAN, etc.

Table titled 'PAPERS NOT REFERRED TO IN DATA CARDS' listing additional references.

Table titled 'A(2000)' with columns for energy (E), width (W), and decay parameters (GALTERI, BRANDSTETTER, etc.).

Table titled '60 Y(402000) MASS (MEV)' showing mass values for various decay channels.

Table titled '60 Y(402000) WIDTH (MEV)' showing width values for various decay channels.

Table titled '60 Y(402000) PARTIAL DECAY MODES' listing decay channels like INTIC N KBAR, INTO LAMBDA OMEGA, etc.

Table titled '60 Y(402000) BRANCHING RATIOS' listing branching ratios for various decay modes.

Table titled '60 Y(402000) INTIC N SIGMA PI' listing partial decay widths and branching ratios.

Table titled '60 Y(402000) INTIC N SIGMA PI' listing partial decay widths and branching ratios.

Table titled 'A(2020)' with columns for energy (E), width (W), and decay parameters (GALTERI, BRANDSTETTER, etc.).

Table titled '60 Y(402020) MASS (MEV)' showing mass values for various decay channels.

Table titled '60 Y(402020) WIDTH (MEV)' showing width values for various decay channels.

Table titled '60 Y(402020) PARTIAL DECAY MODES' listing decay channels like INTIC N KBAR, INTO LAMBDA OMEGA, etc.

Table titled 'REFERENCES FOR Y(402020)' listing references for various baryons like ARMENTEROS, BUGG, BICMAN, etc.

Data Card Listings

For notation, see key at front of Listings.

Baryons

$\Lambda(2350)$, $\Lambda(2585)$, Σ^+ , Σ^- , $\Sigma(1385)$

Table with 2 columns: P1 (Decay Masses) and P2 (Branching Ratios). Rows include data for $\Lambda(2350)$ and $\Lambda(2585)$.

REFERENCES FOR $\Lambda(2350)$ (EPJCL, EPJL)
REFERENCES FOR $\Lambda(2585)$ (EPJCL, EPJL)
REFERENCES FOR $\Sigma(1385)$ (EPJCL, EPJL)

$\Lambda(2585)$ BUMPS

Table of experimental data for $\Lambda(2585)$ bumps, including mass, width, and branching ratios.

S=-1 I=1 HYPERON STATES (Σ)

Table with 2 columns: P1 (Decay Masses) and P2 (Branching Ratios) for Σ^+ and Σ^- .

Table with 2 columns: P1 (Decay Masses) and P2 (Branching Ratios) for $\Sigma(1385)$.

Table with 2 columns: P1 (Decay Masses) and P2 (Branching Ratios) for $\Sigma(1385)$ and $\Sigma(1385)$.

Baryons

 $\Sigma(1560)$, $\Sigma(1580)$, $\Sigma(1620)$

R1 Y#11560) INTO SIGMA PI/(SIGMA PI + LAMBDA PI) (P2)/(P1+P2)
 R1 0.35 0.12 DIDNISI 78 HBC K-P TO P1+K KBAR 5/79
 R2 Y#11560) INTO LAMBDA PI/TOTAL (P1) 1/11
 R2 SEF4 LOCKMAN 78 SPEC PP TO L P1 PI X 12/79

REFERENCES FOR Y#11560) (PROD. EXP.)
 DIDNISI 78 PL 788 154 +HMYTEROS, DIAZ* (CERN+MST+JHEP+ORFII)
 LOCKMAN 78 CEN EPHE 78-01 +MEYER, RANDEY, POSNER, SCHLEIN* IUCLA+SA, LI

PAPERS NOT REFERRED TO IN DATA CARDS
 CARROLL 76 PPL 37 366 +CHIANG, KYCIA, LI, MAZUR, MICHAEL* (BNLII
 HEADMS 80 TORONTO CONF. 283 N.J. HEADMS (CIN)

$\Sigma(1580)$ **D₁₃**
 CO Y#11580, JP=3/2-1 I=1 4/75
 (OBSERVED IN K⁻ N⁺ I=1 TOTAL CS WITHOUT JP ASSIGNMENT AT
 ENLII 73; CARROLL 75; CARROLL 76) AND IN PWA OF K⁻ P⁺
 -> LAMBDA PI FOR CH ENERGIES=1560-1600 MEV BY
 LITCHEFIELD 74; LITCHEFIELD 74 FINDS JP=3/2- NOT SEEN B
 ENGLER 78 OR BY CAMERON 78 WITH LARGER STATISTICS I
 IN KLDUS P TO P1+ LAMBDA AND P1+ SIGMA0.

DD Y#11580) MASS (MEV) 4/75
 M L 1582. 4. LITCHEFIELD 74 DPWA 0 K⁻ P TO LAM PI 4/75
 M C 1583. CARROLL 76 DPWA I=1 TOTAL CS 2/77

CO Y#11580) WIDTH (MEV) 4/75
 M L 11. 4. LITCHEFIELD 74 DPWA 0 K⁻ P TO LAM PI 4/75
 M C (15.) CARROLL 76 DPWA I=1 TOTAL CS 2/77

DD Y#11580) PARTIAL DECAY MODES 4/75
 P1 Y#11580) INTO N KBAR 938+ 493
 P2 Y#11580) INTO LAMBDA PI 1135+ 139
 P3 Y#11580) INTO SIGMA PI 1189+ 139

CO Y#11580) BRANCHING RATIOS 4/75
 R1 Y#11580) INTO SIGMA PI/TOTAL (P1) 4/75
 R1 L +0.03 0.01 LITCHEFIELD 74 DPWA KBAR N MULTIC+NL 4/75
 R1 L MAIN EFFECT OBSERVED BY LITCHEFIELD 74 IS IN P1 LAMBDA FINAL STATE. 4/75
 R1 L KBAR N AND SIGMA PI COUPLINGS ALSO ESTIMATED FROM MULTICHANNEL FII 4/75
 R1 L INCLUDING TOTAL CROSS SECTION DATA (LI T33) 4/75
 R1 C TOTAL CROSS SECTION BUMP WITH I=1/2+K=06 SEEN BY CARROLL 76 2/77

R2 Y#11580) FROM KBAR N TO LAMBDA PI 4/75
 R2 L +0.10 0.02 LITCHEFIELD 74 DPWA D K⁻ P TO LAM PI 4/75
 R2 NOT SEEN CAMERON 78 HBC + KL P TO P1+ LAM 1/78
 R2 NOT SEEN ENGLER 78 HBC + KL P TO P1+ LAM 2/77

R3 Y#11580) FROM KBAR N TO SIGMA PI 4/75
 R3 L +0.03 0.04 LITCHEFIELD 74 DPWA KBAR N MULTIC+NL 4/75
 R3 NOT SEEN CAMERON 78 HBC + KL P TO P1+ SIGD 1/78
 R3 NOT SEEN ENGLER 78 HBC + KL P TO P1+ SIGD 2/77

REFERENCES FOR Y#11580)
 LITCHEFIELD 74 PL 518 509 LITCHEFIELD (CERN)JP
 CARROLL 76 PPL 37 366 +CHIANG, KYCIA, LI, MAZUR, MICHAEL* (BNLII)
 ENGLER 76 PL 638 231 +KEYES, KRAMEER, SCHLEIN, TANAKA* (CERN, ANLII)
 CAMERON 78 NP 8132 189 +CAPUZZI, IGGAN+EDIN+CLAS+PISAR+HELI
 ENGLER 78 PR D18 3061 +KEYES, KRAMEER, TANAKA+CHD* (CERN, ANLII)

PAPERS NOT REFERRED TO IN DATA CARDS
 CARROLL 73 APS BKLY MTG 208 CARROLL, CHIANG, KYCIA, LI, MAZUR, MICHAEL+ (BNLII)
 LJ 73 PURDUE CONF. 283 LI (BNLII)

Note on the $\Sigma(1620)$

This state was first suggested by CRENNELL 68 from evidence in the reaction $K^{\bar{n}} + \Sigma(1620) \rightarrow \bar{K} + \pi^-$ at 3.9 GeV/c, with the $\Sigma(1620)$ decaying into $\Lambda\bar{n}$. The situation is still far from clear.

MILLER 70 is a good review of the production experiments; there has been no new evidence from them since 1970. The evidence is only in the $\Lambda\bar{n}$ channel. CRENNELL 69 claimed the effect in this channel with no evidence seen in $\bar{K}N$ or $\bar{K}N\pi$. SABRE

Data Card Listings

For notation, see key at front of Listings.

70 studied the same reaction at 3.0 GeV/c with comparable statistics and saw no evidence in the $\Lambda\bar{n}$ channel; on the contrary, they believe the effect to be a spurious peak resulting from misidentified Σ^0 from the production of $\Sigma(1670)$ decaying into $\Sigma^0\pi^+$. AMMANN 70 studied the same reaction at 4.5 GeV/c and reported a state at 1640 MeV, again decaying only into $\Lambda\bar{n}$ (no evidence seen in the $\Sigma\pi$ or $\bar{K}N$ channels). HUNGERBUHLER 74 reported upper limits on production cross sections for a 25 GeV/c Σ^- beam.

CARROLL 76 measured the $K^{\bar{p}}$ and $K^{\bar{d}}$ total cross sections from 0.4 to 1.1 GeV/c and found three narrow (10-15 MeV wide) bumps in the $I = 1$ $\bar{K}N$ cross section at 1583, 1608, and 1633 MeV.

Several partial-wave analyses have found evidence for one or two fairly narrow states within about 50 MeV of the effect seen in production: see the entries for the D13 $\Sigma(1580)$, the S11 $\Sigma(1620)$, and the P11 $\Sigma(1660)$. However, the various analyses do not agree on widths or branching ratios.

In conclusion, clarification of the $\Sigma(1620)$ question probably must await more data and a more complete understanding of the entire region from 1600 to 1700 MeV. The closeness of the $\Sigma(1620)$ and $\Sigma(1670)$ masses suggests that the complications of the two regions may be related (see the "Note on the $\Sigma(1670)$ " below).

 $\Sigma(1620)$

32 Y#11620, JP=1/2-1 I=1 S11
 THE S11 STATE AT 1697 MEV REPORTED BY VANHORN 75 IS INTERMEDIATE IN MASS BETWEEN THE SIGMA(1620) AND SIGMA(1750). WE TENTATIVELY LIST IT UNDER SIGMA(1750). CARROLL 76 SEES TWO BUMPS IN THE I=1 TOTAL CROSS SECTIONS NEAR THIS MASS.

32 Y#11620) MASS (MEV)
 M (1620-) 4/75
 M 1630-0 (10.0) KEM 71 DPWA K-MATRIX ANAL. 3/71
 M L 1608. 5. LANGBEIN 72 1PWA MULTICHANNEL 12/72
 M H 1633. 10. CARROLL 76 DPWA I=1 TOTAL CS 2/77
 M J 1600-0 (10.0) MORRIS 78 DPWA K⁻ N TO LAM PI 3/79
 M I AN EQUALLY GOOD FIT IS OBTAINED WITHOUT INCLUDING THIS RESONANCE. 3/79
 M AVG 1633-0 10.0 AVERAGE ERROR INCLUDES SCALE FACTOR OF 2.2)

32 Y#11620) WIDTH (MEV)
 M (160-) 4/75
 M 65.0 (20.0) KEM 71 DPWA K-MATRIX ANAL. 3/71
 M L (15-) CARROLL 76 DPWA MULTICHANNEL 12/72
 M H (10-) MORRIS 78 DPWA I=1 TOTAL CS 2/77
 M J (87.0) (19.0) MORRIS 78 DPWA K⁻ N TO LAM PI 3/79

32 Y#11620) PARTIAL DECAY MODES
 P1 Y#11620) INTO N KBAR 938+ 493
 P2 Y#11620) INTO SIGMA PI 1189+ 139
 P3 Y#11620) INTO LAMBDA PI 1135+ 134

Data Card Listings
For notation, see key at front of Listings.

Baryons

Σ(1620), Σ(1660)

Σ(1620) BRANCHING RATIOS

Table with columns for data card ID, values, and labels. Includes entries for Σ(1620) into KΛ and Σ(1620) from KΛ.

REFERENCES FOR Σ(1620)

Table listing references for Σ(1620) with columns for author, year, journal, and title.

PAPERS NOT REFERRED TO IN DATA CARDS

Table listing additional references for Σ(1620) not included in the main reference list.

1620 MEV REGION - PRODUCTION EXPERIMENTS

Σ(1620) JP, 1 1-1 PRODUCTION EXPERIMENTS

SEE THE MINI-REVIEW AT THE START OF THE YR LISTINGS. THIS RESONANCE NEEDS CONFIRMATION... COLLASORATION AT 3.0 GEV... BRANCHING RATIOS THEY DO NOT ASSOCIATE WITH THE Σ(1620). SEE MILLER 79 FOR A REVIEW OF THESE COMPLEXES.

Σ(1660) MASS (MEV) (PROD. EXP.)

Table showing mass measurements for Σ(1660) with columns for data card ID, mass values, and labels.

Σ(1620) WIDTH (MEV) (PROD. EXP.)

Table showing width measurements for Σ(1620) with columns for data card ID, width values, and labels.

Σ(1620) PARTIAL DECAY MODES (PROD. EXP.)

Table showing partial decay modes for Σ(1620) with columns for data card ID, mode names, and branching ratios.

Σ(1620) BRANCHING RATIOS (PROD. EXP.)

Table showing branching ratios for Σ(1620) with columns for data card ID, mode names, and ratios.

REFERENCES FOR Σ(1620) (PROD. EXP.)

Table listing references for Σ(1620) production experiments with columns for author, year, journal, and title.

PAPERS NOT REFERRED TO IN DATA CARDS

Table listing additional references for Σ(1620) production experiments not included in the main reference list.

Σ(1660)

SEE THE MINI-REVIEW AT THE START OF THE YR LISTINGS.

Σ(1660) MASS (MEV)

Table showing mass measurements for Σ(1660) with columns for data card ID, mass values, and labels.

Σ(1660) WIDTH (MEV)

Table showing width measurements for Σ(1660) with columns for data card ID, width values, and labels.

Σ(1660) PARTIAL DECAY MODES

Table showing partial decay modes for Σ(1660) with columns for data card ID, mode names, and branching ratios.

Σ(1660) BRANCHING RATIOS

Table showing branching ratios for Σ(1660) with columns for data card ID, mode names, and ratios.

Table showing additional mass and width measurements for Σ(1660) with columns for data card ID, values, and labels.

Table showing additional production experiment data for Σ(1660) with columns for data card ID, values, and labels.

REFERENCES FOR Σ(1660)

Table listing references for Σ(1660) with columns for author, year, journal, and title.

Data Card Listings

For notation, see key at front of Listings.

Baryons

$\Sigma(1670)$

<p>REFERENCES FOR $\Psi(1670)$</p> <p>11 $\Psi(1670)$ PARTIAL DECAY MODES (PROD. EXP.)</p>	<p>11 $\Psi(1670)$ BRANCHING RATIOS (PROD. EXP.)</p>
<p>11 $\Psi(1670)$ INTG Λ BARR</p> <p>12 $\Psi(1670)$ INTG Λ BARR</p> <p>13 $\Psi(1670)$ INTG Σ BARR</p> <p>14 $\Psi(1670)$ INTG Σ BARR</p> <p>15 $\Psi(1670)$ INTG Σ BARR</p> <p>16 $\Psi(1670)$ INTG Σ BARR</p> <p>17 $\Psi(1670)$ INTG Σ BARR</p>	<p>11 $\Psi(1670)$ INTG Λ BARR</p> <p>12 $\Psi(1670)$ INTG Λ BARR</p> <p>13 $\Psi(1670)$ INTG Λ BARR</p> <p>14 $\Psi(1670)$ INTG Λ BARR</p> <p>15 $\Psi(1670)$ INTG Λ BARR</p> <p>16 $\Psi(1670)$ INTG Λ BARR</p> <p>17 $\Psi(1670)$ INTG Λ BARR</p>
<p>REFERENCES FOR $\Psi(1670)$</p> <p>11 $\Psi(1670)$ PARTIAL DECAY MODES (PROD. EXP.)</p>	<p>11 $\Psi(1670)$ BRANCHING RATIOS (PROD. EXP.)</p>
<p>11 $\Psi(1670)$ INTG Λ BARR</p> <p>12 $\Psi(1670)$ INTG Λ BARR</p> <p>13 $\Psi(1670)$ INTG Λ BARR</p> <p>14 $\Psi(1670)$ INTG Λ BARR</p> <p>15 $\Psi(1670)$ INTG Λ BARR</p> <p>16 $\Psi(1670)$ INTG Λ BARR</p> <p>17 $\Psi(1670)$ INTG Λ BARR</p>	<p>11 $\Psi(1670)$ INTG Λ BARR</p> <p>12 $\Psi(1670)$ INTG Λ BARR</p> <p>13 $\Psi(1670)$ INTG Λ BARR</p> <p>14 $\Psi(1670)$ INTG Λ BARR</p> <p>15 $\Psi(1670)$ INTG Λ BARR</p> <p>16 $\Psi(1670)$ INTG Λ BARR</p> <p>17 $\Psi(1670)$ INTG Λ BARR</p>
<p>REFERENCES FOR $\Psi(1670)$</p> <p>11 $\Psi(1670)$ PARTIAL DECAY MODES (PROD. EXP.)</p>	<p>11 $\Psi(1670)$ BRANCHING RATIOS (PROD. EXP.)</p>
<p>11 $\Psi(1670)$ INTG Λ BARR</p> <p>12 $\Psi(1670)$ INTG Λ BARR</p> <p>13 $\Psi(1670)$ INTG Λ BARR</p> <p>14 $\Psi(1670)$ INTG Λ BARR</p> <p>15 $\Psi(1670)$ INTG Λ BARR</p> <p>16 $\Psi(1670)$ INTG Λ BARR</p> <p>17 $\Psi(1670)$ INTG Λ BARR</p>	<p>11 $\Psi(1670)$ INTG Λ BARR</p> <p>12 $\Psi(1670)$ INTG Λ BARR</p> <p>13 $\Psi(1670)$ INTG Λ BARR</p> <p>14 $\Psi(1670)$ INTG Λ BARR</p> <p>15 $\Psi(1670)$ INTG Λ BARR</p> <p>16 $\Psi(1670)$ INTG Λ BARR</p> <p>17 $\Psi(1670)$ INTG Λ BARR</p>

$\Sigma(1670)$ BUMPS



Baryons

Σ(1670), Σ(1690), Σ(1750)

CARROLL 70 PH 37 806 ... <CHENG,VICALLI,PRUDR,MICHAEL> 18N131 ...

LIVONIE 65 PL 18 69 ... * (SACLAY,ERDOL,GLASSON,LOIC,OFF,HELI) JP ...

Σ(1690) BUMPS

SEE THE MINI-REVUE AT THE START OF THE ΞA LISTINGS. SEE NOTE PRECEDING Ξ(1670) LISTINGS; SEE IN PRO. EXPERIMENTS ONLY; MAIN DECAY MODE IS LAMBDA P1.

58 Ξ(1690) MASS (MEV) (PROD., EXP.) ... COLLEY 67 HBC + K-P 6 GEV/C 8/67 ...

59 Ξ(1690) WIDTH (MEV) (PROD., EXP.) ... COLLEY 67 HBC + K-P 6 GEV/C 8/67 ...

56 Ξ(1690) PARTIAL DECAY MODES (PROD., EXP.) ... DECEAY MASSES ...

57 Ξ(1690) BRANGLING RATIOS (PROD., EXP.) ... COLLEY 67 HBC + K-P 6 GEV/C 8/67 ...

58 Ξ(1690) BRANGLING RATIOS (PROD., EXP.) ... COLLEY 67 HBC + K-P 6 GEV/C 8/67 ...

59 Ξ(1690) BRANGLING RATIOS (PROD., EXP.) ... COLLEY 67 HBC + K-P 6 GEV/C 8/67 ...

REFERENCES FOR Ξ(1690) (PROD., EXP.) ... COLLEY 67 PH 246 649 ...

AGUILAR 70 PH 25 74 ... AGUILAR-BENITEZ,BARNES,BASSANO, (BN,SYN) ...

Data Card Listings

For notation, see key at front of Listings.

Σ(1750)

57 Ξ(1750), JP=1/2- 1/4 ... THERE IS EVIDENCE FOR THIS STATE IN MANY PARTIAL- ...

57 Ξ(1750) MASS (MEV) ... ABOUT SIGMA ETA THRESHOLD ...

57 Ξ(1750) WIDTH (MEV) ... ABOUT 80.0 ...

57 Ξ(1750) PARTIAL DECAY MODES ... DECEAY MASSES ...

57 Ξ(1750) BRANGLING RATIOS ... COLLEY 67 HBC + K-P 6 GEV/C 8/67 ...

57 Ξ(1750) BRANGLING RATIOS ... COLLEY 67 HBC + K-P 6 GEV/C 8/67 ...

AGUILAR 70 PH 25 74 ... AGUILAR-BENITEZ,BARNES,BASSANO, (BN,SYN) ...

Baryons
Σ(1915), Σ(1940)

Data Card Listings

For notation, see key at front of Listings.

46 Y*(1915) PARTIAL DECAY MODES

Table with columns for particle ID, decay mode, and mass/energy. Includes entries for Σ(1915) into nK and Σ(1915) into ΛK.

46 Y*(1915) BRANCHING RATIOS

Table of branching ratios for Σ(1915) decays into various channels like nK, ΛK, and ΣK. Includes theoretical predictions and experimental values.

Table of branching ratios for Σ(1915) decays into nK and ΣK. Includes theoretical predictions and experimental values.

Table of branching ratios for Σ(1915) decays into nK and ΣK. Includes theoretical predictions and experimental values.

REFERENCES FOR Y*(1915)
ARMENIERS, FERRO-LUZZI + (CERN, HELV. SAFLAY)
SMART 69 PR 160 1330
BERTHON TO NP 820 476
BERNHOFF TO NP 824 417
BRICKMAN TO NP 830 511
CDX 70 NP 819 61
GALTIERI TO QUÉBEC CONF 173
LITCHEF TO NP 823 476
CONFORTO TO NP 824 414
LITCHEF 71 NP 830 125
KANE 72 PR 85 1583
DEVENISH 74 NP 881 330
KANE 74 LBL 2549
BAILLON 75 NP 894 39
HEMINGWAY 75 NP 891 12
VATHRON 75 NP 887 145
A. J. VAN HORN
BELLEFON 76 NP 8109 120
CRODEM 76 NP 8104 382
CRODENI 76 NP 8125 41
DECLAIS 77 CERNA 17-16
MARTIN 77 NP 8127 349
ALSO 77 NP 8126 266
ALSO 77 NP 8126 285
RILEY 77 NP 8128 262
ALSTON 78 PR D18 172
ALSO 77 PR L 38 1007
CAMERON 78 NP 8143 180
G.P. COPAL
PAPERS NOT REFERRED TO IN DATA CARDS
SMART 66 PR L 556
SUPERSEDED BY SMART 69.
CONFORTO 68 NP 88 265
HARSEN, LASTINSKI + (CICHAGO, METEDEL)
SUPERSEDED BY CONFORTO 71.

1915 MEV REGION - PRODUCTION AND TOTAL EXP'TS

29 Y*(1915), Jπ = 1 1/2 PRODUCTION EXPERIMENTS
SEE THE MINI-REVIEW AT THE START OF THE Y* LISTINGS.

SEE THE NOTES TO THE Y*(1915) AND Y*(1940), WHICH IMMEDIATELY PRECEDE AND FOLLOW THIS ENTRY.
THE Y*(1915) ONLY CARRIES LISTINGS OF PERKS SEEN IN CROSS SECTION AND INVAARIANT MASS DISTRIBUTIONS. THE CROSS-SECTION PERKS ARE ALD57 CERTAINLY ASSOCIATED WITH THE Y*(1915) SINCE IN CROSS SECTION ANALYSES. THE INVARIANT-MASS PERKS SEEM MORE LIKELY TO BE ASSOCIATED WITH THE D13 Y*(1940).

29 Y*(1915) MASS (MEV) (PROD. EXP.)

Table of mass measurements for Σ(1915) from various experiments. Includes columns for experiment name, mass value, and error.

29 Y*(1915) 410TH (MEV) (PROD. EXP.)

Table of mass measurements for Σ(1915) from various experiments. Includes columns for experiment name, mass value, and error.

29 Y*(1915) PARTIAL DECAY MODES (PROD. EXP.)

Table of partial decay modes for Σ(1915) from various experiments. Includes columns for decay mode, experiment name, and branching ratio.

29 Y*(1915) BRANCHING RATIOS (PROD. EXP.)

Table of branching ratios for Σ(1915) from various experiments. Includes columns for decay mode, experiment name, and branching ratio.

REFERENCES FOR Y*(1915) (PROD. EXP.)
BOCK 65 PL 17 166
BUGG 68 PR 168 1486
BARNES 69 OBL 22 475
AGUILAR 70 PL 25 SE
BRICKMAN 70 PL 310 152
COOGL 70 PR D1 1887
DADD 72 OBL 29 1405
RIFEL 77 PR D1 2706
FERRER 81 NP 8178 373
*COOPER, FRENCH, KINSON, + (CERN, SAFLAY)
*GILKROMELL, KVICIA, LEONIC, LI, LUMBOY, + (BNL) I
*HOLMROD, KNIGHT, DAVIES + (BIRMINGHAM, MELLI)
*LEWIS, MOULTON, SMITHS + (BNL, SYRAC)
*AGUILAR-BENITEZ, BARNES, + (BNL, SYRAC)
*FERRO-LUZZI, PERREAU + (CERN, GENEVA, SAFLAY)
*GILKROMELL, KVICIA, LEONIC, LI, + (BNL) I
*BERNHOFF, GOLDBERG, HESS, TAYLOR, MASS (SEE THE PRECEDING ENTRY). THIS STATE IS NOT REQUIRED IN K-N NEUTRON TO (P) SIGMA+ ANALYSIS OF GOLF 77. KBAR N ANALYSIS (GOLF 80) WITH K- NEUTRON ELASTIC DATA DOES NOT REQUIRE THIS STATE.

Σ(1940)

28 Y*(1940), Jπ = 3/2- 1 1/2
SEE THE MINI-REVIEW AT THE START OF THE Y* LISTINGS.

SOME, NOT ALL, PARTIAL MASS ANALYSES SUGGEST A STATE IN THIS REGION. IT IS PERHAPS ASSOCIATED WITH THE BUMPS SEEN IN PRODUCTION EXPERIMENTS NEAR THE MASS (SEE THE PRECEDING ENTRY). THIS STATE IS NOT REQUIRED IN K-N NEUTRON TO (P) SIGMA+ ANALYSIS OF GOLF 77. KBAR N ANALYSIS (GOLF 80) WITH K- NEUTRON ELASTIC DATA DOES NOT REQUIRE THIS STATE.

28 Y*(1940) MASS (MEV)

Table of mass measurements for Σ(1940) from various experiments. Includes columns for experiment name, mass value, and error.

Data Card Listings

For notation, see key at front of Listings.

Baryons

 $\Sigma(3000)$, $\Sigma(3170)$, EXOTIC HYPERONS, Ξ 's

54 Y#1(2620) WIDTH (MEV) (PROD. EXP.)

W	(175.0)	R1.	ABRAMS	70 CNTR	K-P, D TOTAL	10/70
	221.		DIBIANCA	75 DBC	X1 < P1	1/76

54 Y#1(2620) PARTIAL DECAY MODES (PROD. EXP.)

P1	Y#1(2620)	INTC N	KBAR	DECAY MASSES
				938 + 493

54 Y#1(2620) BRANCHING RATIOS (PROD. EXP.)

P1	Y#1(2620)	INTC	KBAR	N/TOTAL	(P1)
R1		J IS NOT KNOWN. THE FOLLOWING IS J#1/2#0#1.			
P1		(0,32)	ABRAMS	70 CNTR	K-P, D TOTAL
R1		0.36	0.12	DIBIANCA	70 CNTR D TOTAL AND CH EXP

REFERENCES FOR Y#1(2620) (PROD. EXP.)

ABRAMS AT PRL 15 678 +COOL, GIACOMELLI, KYCIA, LEONTEGILI, + (BNL)
SUPERSEDED BY ABRAMS 70.
ABRAMS 70 PR D 1917 +COOL, GIACOMELLI, KYCIA, LEONTEGILI, + (BNL) I
BRICHMAN 70 PR 318 152 +FERRO LUZZI, MORENO+ (CECERN,CALIF,SLAC) I

DIBIANCA 75 YP 568 127 DIBIANCA,LEODORFER (CERN)

**$\Sigma(3000)$
BUMPS**



59 Y#1(3000, J# = 1) [4] PRODUCTION EXPERIMENTS

SEE THE MINI-REVIEW AT THE START OF THE Y# LISTINGS.

ENHANCEMENT IN LAMBDA P1 AND RAR N INVARIANT MASS SPECTRA AND IN MISSING MASS OF NEUTRALS DECAYING AGAINST K0. EVIDENCE NOT CONCLUSIVE. OMITTED FROM TABLE.

59 Y#1(3000) MASS (MEV) (PROD. EXP.)

M	(3000.0)	EMPLICH	66 HBC	D P1-P	7.01 BEVC	9/66
---	----------	---------	--------	--------	-----------	------

59 Y#1(3000) PARTIAL DECAY MODES (PROD. EXP.)

P1	Y#1(3000)	INTC N	KBAR	DECAY MASSES
P2 <td>Y#1(3000) <td>INTC LAMBDA P1</td> <td></td> <td>1115-1139</td> </td>	Y#1(3000) <td>INTC LAMBDA P1</td> <td></td> <td>1115-1139</td>	INTC LAMBDA P1		1115-1139

REFERENCES FOR Y#1(3000) (PROD. EXP.)

EMPLICH 66 PR 152 1194 R EMPLICH, W SELOFF, H YUTA (PENNSYLV) I

119 Y#1(3170, J# = 1) [1] PRODUCTION EXPERIMENTS

SEEN BY AMIRZADEH 79 AS A NARROW C... (G. DEV. ENHANCEMENT IN THE REACTION K⁻P⁻ -> Y# P1- USING DATA FROM TWO INDEPENDENT HIGH STATISTICS BUBBLE CHAMBER EXPERIMENTS AT 8.25 AND 6.5 GEV/C. THE DOMINANT DECAY MODES ARE INTO MULTI-BODY, MULTI-STRANGE FINAL STATES AND THE PRODUCTION IS VIA 1532 BARYON EXCHANGE. T-1 IS FAVORED.

14 NEED OF CONFIRMATION. OMITTED FROM TABLES.

**$\Sigma(3170)$
BUMPS**



118 Y#1(3170) MASS (MEV) (PROD. EXP.)

M	35 3170.	S.	AMIRZAD	79 HBC	K-P TO Y# P1-	12/79
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118 Y#1(3170) WIDTH (MEV) (PROD. EXP.)

M	C	35 (20.)	CR LESS	AMIRZAD	79 HBC	K-P TO Y# P1-	12/79
M <td>C <td></td> <td>OBSERVED WIDTH</td> <td>CONSISTENT WITH EXPERIMENTAL RESOLUTION.</td> <td></td> <td></td> <td></td> </td>	C <td></td> <td>OBSERVED WIDTH</td> <td>CONSISTENT WITH EXPERIMENTAL RESOLUTION.</td> <td></td> <td></td> <td></td>		OBSERVED WIDTH	CONSISTENT WITH EXPERIMENTAL RESOLUTION.			

118 Y#1(3170) PARTIAL DECAY MODES (PROD. EXP.)

P1	Y#1(3170)	INTC	LAMBDA K	KBAR	+ PIONS	DECAY MASSES
P2	Y#1(3170)	INTC SIGMA K <th>KBAR</th> <th>+ PIONS</th> <td></td> <td></td>	KBAR	+ PIONS		
P3	Y#1(3170)	INTC X1 <th>K</th> <th>+ PIONS</th> <td></td> <td></td>	K	+ PIONS		

118 Y#1(3170) BRANCHING RATIOS (PROD. EXP.)

R1	Y#1(3170)	INTC	LAMBDA K	KBAR	+ PIONS	TOTAL	(P1)
P1		SEEN	AMIRZAD	79 HBC	K-P TO Y# P1-		12/79
R2	Y#1(3170)	INTC SIGMA K <th>KBAR</th> <th>+ PIONS</th> <th>TOTAL</th> <th>(P2)</th>	KBAR	+ PIONS	TOTAL	(P2)	
P2		SEEN	AMIRZAD	79 HBC	K-P TO Y# P1-		12/79
R3	Y#1(3170)	INTC X1 <th>K</th> <th>+ PIONS</th> <th>TOTAL</th> <th>(P3)</th>	K	+ PIONS	TOTAL	(P3)	
P3		SEEN	AMIRZAD	79 HBC	K-P TO Y# P1-		12/79

REFERENCES FOR Y#1(3170) (PROD. EXP.)

AMIRZAD 79 PL 898 175 AMIRZADEH (BIRMINGHAM+GLASGOW+ORNL+CAMB) I
ALSO 80 TORONTO CONF. 263 J.B. KINOSHITA (BIRMINGHAM+GLASGOW+ORNL+CAMB) I

EXOTIC HYPERON CROSS SECTION LIMITS

THIS IS NOT A COMPLETE LIST. WE TABULATE ONLY FROM 1970 ON.

CS	UNITS	MICROBARNS
CS G	120	CR LESS GALTIERI 68 DBC -- K-M TO SS-PI-P10 7/70
CS G	ABOVE LIMIT	FOR MASS < 2.15 GEV AND WIDTH < 60 MEV 12.1 GEV/C K-1 7/70
CS A	(42.1)	OR LESS GALTIERI 68 DBC -- K-M TO SS-PI-P10 7/70
CS A	ABOVE LIMIT	FOR MASS < 2.3 GEV AND WIDTH < 120 MEV 12.7 GEV/C K-1 7/70
CS X	(4.7)	CP LESS CL+90 BRIEFEL 75 DBC -- K-D 2.87 GEV/C 3/79
CS Y	WIDTH < 40 MEV, K-M -> 131-PI-1 6*	
CS Y	(1.4)	CR LESS CL+90 BRIEFEL 75 DBC -- K-D 2.87 GEV/C 3/79
CS Y	WIDTH < 40 MEV, K-M -> 131-PI-1 6*	
CS Z	WIDTH < 60 MEV, K-M -> 141-PI-1 6* P10	
CS B	(8.0)	CP LESS CL+90 KATSUOFE 78 DEC -- K-D 2.87 GEV/C 3/79
CS B	WIDTH < 60 MEV, K-M -> 151(GM)-PI-1 P1*	
CS C	(13.3)	CR LESS CL+90 KATSUOFE 78 DEC -- K-D 2.87 GEV/C 3/79
CS D	(6.9)	CR LESS CL+90 KATSUOFE 78 DEC -- K-D 2.87 GEV/C 3/79
CS D	MASS > 2 GEV, WIDTH < 60 MEV, K-M -> 151(GM)-PI-1 P1*	
CS E	(17.7)	CR LESS CL+90 KATSUOFE 78 DEC -- K-D 2.87 GEV/C 3/79
CS E	MASS > 2 GEV, WIDTH < 120 MEV, K-M -> 151(GM)-PI-1 P1*	
CS F	(17.7)	CR LESS CL+90 KATSUOFE 78 DEC -- K-D 2.87 GEV/C 3/79
CS F	WIDTH < 60 MEV, K-M -> 151(GM)-PI-1 P10 P1*	
CS H	(12.3)	CR LESS CL+90 KATSUOFE 78 DEC -- K-D 2.87 GEV/C 3/79
CS H	WIDTH < 120 MEV, K-M -> 151(GM)-PI-1 P10 P1*	
CS I	(12.1)	CR LESS CL+90 KATSUOFE 78 DEC -- K-D 2.87 GEV/C 3/79
CS I	WIDTH < 60 MEV, K-M -> 151(GM)-PI-1 P1- P10	

REFERENCES FOR EXOTIC HYPERONS

GALTIERI 68 PRL 21 573 A,BRABAO-GALTIERI,CHADHA, + (CUL+SLAC)
BRIEFEL 78 PRL 12 1899 KATSUOFE, KATSUOFE, CHAN, + (ORNL+UNIV OF TORONTO)
KATSUOFE 78 PRL 18 16 KATSUOFE, CANTER, MANN, SCHNEPS, + (TUFT+BRANII)

Note on Ξ^* Resonances

The Ξ^* resonance situation has always been unsettled. This is because: (1) Ξ^* 's can only be produced as a part of a final state, and so the analysis is more complicated than if direct formation were possible; (2) they are produced with small cross sections (typically a few μ); and (3) the final states are topologically complicated and difficult to study with electronic techniques. Thus in the past our knowledge of Ξ^* resonances has come wholly from bubble chamber experiments, where the numbers of events are small.

Until fairly recently only the $\Xi(1530)$ could be considered to be really well established. However, the 1978 edition of this Review¹ saw a major improvement in the situation with the results of GAY 76 and HENINGWAY 77. The $\Xi(1820)$ and $\Xi(2030)$ were firmly established as narrow states (widths of about 20 MeV), and the spin of the $\Xi(1820)$ was found to be 3/2 (TEODORO 78).

Since then, however, little has changed, although there is some evidence for a new $\Xi(2370)$ (AMIRZADEH 80, HASSALL 81). There is probably at least one other Ξ^* in the 1850-2000 MeV region, and there are indications of several others above 2000 MeV. Indeed, there should be many Ξ^* 's below 2500 MeV, and the broad (and not completely established) $\Xi(1940)$ could well be a mixture of several of them.² For now we are forced to group together

Baryons

Ξ 's, Ξ^- , Ξ^0 , $\Xi(1530)$

disparate observations and await new results. The disagreements among experiments are shown in ideograms in the Listings.

Results from experiments using electronic techniques are now becoming available. BIAGI 81 used the CERN hyperon beam to study inclusive $\bar{K}\bar{K}$ and $\Xi\pi$ mass spectra from 102 and 135 GeV/c Ξ^- incident on hydrogen and deuterium. They saw a large $\Xi(1820)$ signal in $\bar{K}\bar{K}$ as well as a peak at about 1700 MeV, which might be associated with the threshold enhancement seen by DIONISI 78. The $\Xi(1940)$ appears as a broad bump in the $\Xi\pi$ mass spectrum, and there is a very clean $\Xi(1530)$ signal. Preliminary results from the Brookhaven multiparticle spectrometer were reported at the Toronto Conference (see Ref. 3). The $\Xi(1820)$ is clearly seen as a narrow object decaying to $\Xi(1530)\pi$.

The table below gives our evaluation of the present status of the Ξ^* resonances. For a detailed review, see MEADOWS 80.³

References

1. Particle Data Group, Phys. Lett. **75B**, 1 (1978).
2. R.-J. Hemingway, in Proceedings of the Topical Conference on Baryon Resonances (Oxford, 1976), ed. R.T. Ross and D.H. Saxon.
3. B.T. Meadows, in Proceedings of the IVth International Conference on Baryon Resonances (Toronto, 1980), ed. N. Isgur.

Data Card Listings

For notation, see key at front of Listings.

S=-2 I=1/2 HYPERON STATES (Ξ)

Ξ^-

22 XI-11321, JP=1/2 1 1/1/2
SEE STABLE PARTICLE DATA CARD LISTINGS

Ξ^0

23 XI011315, JP=1/2 1 1/1/2
SEE STABLE PARTICLE DATA CARD LISTINGS

$\Xi(1530)$

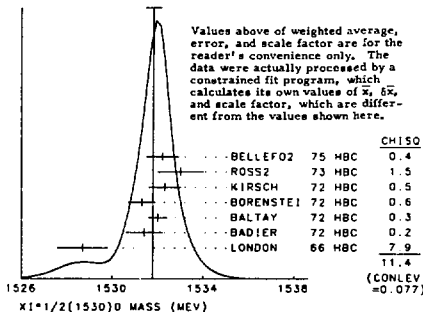
P_{13}

49 XI11/211530, JP=3/2+1 1/1/2
THIS IS THE ONLY WELL-ESTABLISHED XI⁰ WHOSE PROPERTIES ARE ALL AT LEAST REASONABLY WELL-KNOWN. SPI-U-PARTICLE 1/2+ IS FAVOURED BY THE DATA.
WE DO NOT USE DETERMINATIONS OF THE MASS AND THE WIDTH OF THIS STATE UNLESS THEY ARE ACCOMPANIED BY SOME DISCUSSION OF SYSTEMATICS AND RESOLUTION.

49 XI^{1/2}11530 MASS (MEV)

M MIXED CHARGES									
M 20(1535.1)									
M 55(1529.0)	(5.0)					BERNANZA	62 HBC	-0	K-P 2.3 GEV/C
M 11532.0)	(2.0)					QUERREU	82 HBC	-0	K-P 1.9 GEV/C
						BADIER	84 HBC	-0	K-P 3 GEV/C
M- NEGATIVE CHARGE ONLY									
M- 38 1835.7	3.2					LONDON	66 HBC	-K-P	2.24 GEV/C 7/66
M- 334(1534.7)	(1.1)					BALTAY	72 HBC	-K-P	1.79 GEV 1/73
M- 145 1936.2	1.9					KIRSCH	72 HBC	-K-P	2.20 GEV/C 2/72
M- 1535.3	2.0					ROSS2	73 HBC	-X1	KBAR PI (P1) 2/74
M- 48(1540.1)	(3.1)					BERNTHN	74 HBC	-QUAS1	2 BODY CS 10/74
M- 1534.9	1.2					BELLEFO2	75 HBC	-K-P	AT 8.25 GEV 11/75
M- AVG	1535.18	0.34							AVERAGE ERROR INCLUDES SCALE FACTOR OF 1.61
M- FIT	1534.97	0.43							FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.01)
M0 NEUTRAL CHARGE ONLY									
M0 76 1928.7	1.1					LONDON	66 HBC	0	K-P 2.24 GEV/C 7/66
M0 59 1531.4	0.8					BADIER	72 HBC	0	K-P AT 3.95 GEV/C 10/71
M0 1262 1932.0	0.6					BALTAY	72 HBC	0	K-P 1.79 GEV 1/73
M0 324 1931.3	0.6					BORNENSTEI	72 HBC	0	K-P 2.20 GEV/C 2/72
M0 286 1932.3	0.7					KIRSCH	72 HBC	0	K-P 2.20 GEV/C 2/72
M0 1533.0	1.0					ROSS2	73 HBC	0	X1 KBAR PI (P1) 2/74
M0 97(1533.6)	(1.4)					BERNTHN	74 HBC	0	QUAS1 2 BODY CS 10/74
M0 1932.2	0.7					BELLEFO2	75 HBC	0	K-P AT 8.25 GEV 11/75
M0 80(1527.1)	(0.1)					SIXEL	79 HBC	0	INCL. K-P 10 GEV 1/80
M0 100(1535.7)	(1.4)					SIXEL	79 HBC	0	INCL. K-P 16 GEV 1/80
M0 2700(1532.1)	(10.6)					RUBILLIETTE	81 HBC	0	K-P AT 8.25 GEV 2/84
M0 I FIT TO INCLUSIVE SPECTRUM. RESOLUTION (5 MEV) NOT UNFOLDED.									
M0 450(1530.1)	(1.1)					BIAGI	81 SPEC	-	HYPERON BEAM 2/82*
M0 AVG	1531.78	0.34							AVERAGE ERROR INCLUDES SCALE FACTOR OF 1.41
M0 FIT	1531.80	0.31							FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.31) (SEE IDEOGRAM BELOW I)

WEIGHTED AVERAGE = 1531.78 ± 0.34
ERROR SCALED BY 1.4



STATUS OF XI* RESONANCES
THOSE WITH AN OVERALL STATUS OF *** OR **** ARE INCLUDED IN THE MAIN BARYON TABLE. THE OTHERS ARE NOT CONFIRMED.
IN THE PAST WE HAVE LOWERED OUR STANDARDS FOR XI* RESONANCES AND TABULATED STATES EVEN THOUGH THEY HAD ONLY BEEN SEEN AT LOW LEVELS OF STATISTICAL SIGNIFICANCE. NOW THAT NEW HIGH STATISTICS DATA IS AVAILABLE, WE HAVE ADOPTED SOMEWHAT STRICTER CRITERIA.

STATUS AS SEEN IN ---

PARTICLE	LJ	OVERALL STATUS	XI	PI	LAN K	SIF. K	XI*	PI	OTHER CHANNELS
XI(1317)	P13	****							WEAK TO LAN PI
XI(1530)	P13	****	***						
XI(1630)			*						
XI(1680)					*		**		
XI(1820)	I3	***			*		**	**	
XI(1940)		**	**					**	
XI(1700)	I	***			*		**	**	
XI(1700)					*				3-BODY DECAYS
XI(2250)		**			*				3-BODY DECAYS
XI(2370)	I	**			*				3-BODY DECAYS
XI(2500)		*			*				3-BODY DECAYS

*** GOOD, CLEAR, AND UNDISPUTABLE.
**** GOOD, BUT IN NEED OF CLARIFICATION OR NOT ABSOLUTELY CERTAIN.
** NEEDS CONFIRMATION.
* WEAK.

Baryons

$\Xi(1680)$, $\Xi(1820)$

Data Card Listing

For notation, see key at front of Listings.

5 XI⁺1/2(1680) PARTIAL DECAY MODES

		DECAY MODES
P1	XI ⁺ 1/2(1680)	INTO SIGMA KBAR 1192+197
P2	XI ⁺ 1/2(1680)	INTO LAMBDA KBAR 1215+197
P3	XI ⁺ 1/2(1680)	INTO XI PI 1314+134
P4	XI ⁺ 1/2(1680)	INTO XI ⁺ 1/2(1530) PI 1933+134
P5	XI ⁺ 1/2(1680)	INTO XI PI PI (INCLUDING P4) 1314+134+134

5 XI⁺1/2(1680) BRANCHING RATIOS

R1	R2	R3	R4	R5
XI ⁺ 1/2(1680)	INTO (XI PI)/(SIGMA KBAR)	(P3)/(P1)		
(2.7)	(0.0)			
(3.1)	(1.1)			
NEUTRAL CHARGE				
NEGATIVE CHARGE				
XI ⁺ 1/2(1680)	INTO (XI PI)/(SIGMA KBAR)	(P3)/(P1)		
(10.0)	OR LESS			
XI ⁺ 1/2(1680)	INTO (XI PI)/(SIGMA KBAR)	(P3)/(P1)		
(7.0)	OR LESS			
XI ⁺ 1/2(1680)	INTO (XI PI)/(SIGMA KBAR)	(P3)/(P1)		
(10.0)	OR LESS			
XI ⁺ 1/2(1680)	INTO (XI PI)/(SIGMA KBAR)	(P3)/(P1)		
(10.0)	OR LESS			

REFERENCES FOR XI⁺1/2(1680)

SIGNISI	78 PL 805 145	*DIAZ, ARMENTEROS	CERN+ANST+IUM+OXI+J.P
BIAGI	81 DPH C9 305	* (8015+CAMB+GEVA+PEID+LAUS+LOM+MEL)	

$\Xi(1820)$

50 XI⁺1/2(1820), JP=3/2 1-1/2
WE LIST HERE EVERYTHING REPORTED IN THE MASS RANGE 1750-1875 MeV.

The clearest evidence for this state comes from GAY 76, who saw an 8 standard-deviation peak in ΛK^- as well as signals in $\Xi(1530)\pi$ and ΣK^- . The peak is narrow ($\Gamma = 21 \pm 7$ MeV), whereas earlier (and much smaller) experiments found widths of up to 100 MeV (see the Listings below).

A spin-parity analysis of the GAY 76 data, but with more events (TEODORO 78), favors spin 3/2 but cannot make a parity discrimination.

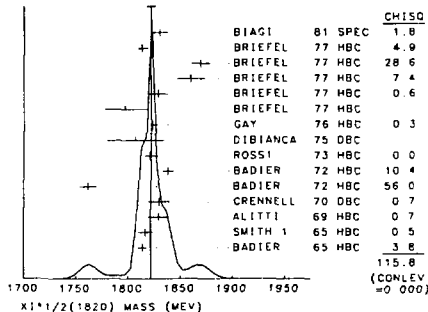
BIAGI 81 used the CERN hyperon beam to study Ξ^- interactions in hydrogen and deuterium. The diffractively produced ΛK^- system has a broad peak ($\Gamma = 72 \pm 20$ MeV) at 1830 MeV on top of a substantial background. There is also a smaller peak in the inclusive ΛK_S^0 spectrum.

Neither GAY 76 nor BIAGI 81 saw a peak in the $\Xi\pi$ channel. It is possible that $\Xi\pi$ peaks seen in this region by some lower-momentum experiments are at least partly due to the $\Xi(1940)$, with a shape distorted by the limited phase space available (SMITH 65). The situation is further confused because some experiments were forced to add several different channels together to overcome poor statistics (CRENNELL 70, BADIER 71).

10 XI⁺1/2(1820) MASS (MEV)

M	(1770.0)		MALTEINSI 63 FRC	-O K-PR 3.5 GEV/C
M	30 1814.0	4.0	BADIER 65 HBC	O LAMBDA K0BAR
M	29 1817.0	7.0	SMITH 2 65 HBC	O LAMBDA K0BAR
M	40 1830.0	10.0	ALITTI 69 HBC	- LAM. SIG KBAR
M	25 1830.0	10.0	CRENNELL 70 DBC	-O 3.6, 3.9 GEV/C
M	0		CRENNELL 70 DBC	-O 3.6, 3.9 GEV/C
M	0		CRENNELL TO DBC	-O 3.6, 3.9 GEV/C
M	28 1762.0	8.0	BADIER 72 HBC	-O XI PI
M	38 1839.0	5.0	BADIER 72 HBC	-O XI PI
M	1 30 1821	5.0	ROSSI 73 HBC	-O LAMBDA K-K0BAR
M	130 1823.0	2.0	GAY 76 HBC	-K- P AT 4.2 GEV
M	74 1797.0	19.0	BRIEFEL 77 HBC	-O XI PI
M	68 1797.0	9.0	BRIEFEL 77 HBC	-O XI PI
M	39 1865.0	14.0	BRIEFEL 77 HBC	-O SIGMA K0BAR
M	44 1870.0	9.0	BRIEFEL 77 HBC	-O LAMBDA K0BAR
M	57 1813.0	4.0	BRIEFEL 77 HBC	-O LAMBDA K-
M	1 300 1835	5.0	BIAGI 81 SPEC	- HYPERON BEAM
M	1		BIAGI 81 SPEC	- HYPERON BEAM

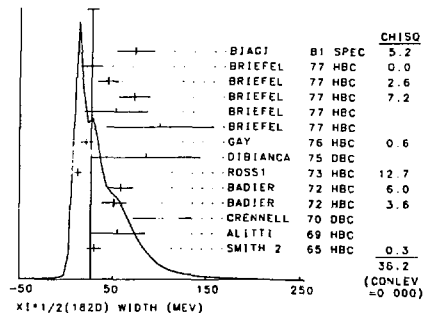
AVERAGE MEANINGLESS SCALE FACTOR = 3.11 (SEE LOGOGRAM BELOW)



10 XI⁺1/2(1820) WIDTH (MEV)

M	180.0	OR LESS	MALTEINSI 63 FRC	-O K-PR 3.5 GEV/C
M	112.0	(4.0)	BADIER 65 HBC	O LAMBDA K0BAR
M	30.0	7.0	SMITH 2 65 HBC	O LAMBDA K0BAR
M	5.0	40.0	20.0	ALITTI 69 HBC
M	105.0	38.0	24.0	CRENNELL 70 DBC
M	145.0	(36.0)	119.0	CRENNELL 70 DBC
M	11.0	13.0		BADIER 72 HBC
M	94.0	13.0		BADIER 72 HBC
M	1 30	12.0		ROSSI 73 HBC
M	130	95.0	58.0	DIBIANCA 75 DBC
M	74	99.0	57.0	GAY 76 HBC
M	68	52.0	34.0	BRIEFEL 77 HBC
M	19	72.0	17.0	BRIEFEL 77 HBC
M	44	44.0	11.0	BRIEFEL 77 HBC
M	67	26.0	11.0	BRIEFEL 77 HBC
M	1 300	72.0	20.0	BIAGI 81 SPEC

AVERAGE MEANINGLESS SCALE FACTOR = 2.21 (SEE LOGOGRAM BELOW)



Data Card Listings
For notation, see key at front of Listings.

Baryons
Ξ(1820), Ξ(1940)

50 XI*1/2(1820) PARTIAL DECAY MODES

Table with columns for decay modes (P1, P2, P3, P4) and corresponding branching fractions or parameters.

FITTED PARTIAL DECAY MODE BRANCHING FRACTIONS

The matrix below is derived from the error matrix for the fitted partial decay mode branching fractions, Pij, as follows: The diagonal elements are Pij*Pij, where Pij = sqrt(SPij/SPj) * Pij, while the off-diagonal elements are the normalized correlation coefficients (SPij/SPi)^(1/2) * (SPj/SPj)^(1/2) * Pij.

Small table showing correlation coefficients for parameters P1, P2, P3, and P4.

50 XI*1/2(1820) BRANCHING RATIOS

Main table listing branching ratios for various decay channels (R1-R52) with associated fit parameters and error estimates.

REFERENCES FOR XI*1/2(1820)

Bibliography listing references for the Xi*(1820) particle, including authors like Halstein, Badier, and others.

PAPERS NOT REFERRED TO IN DATA CARDS

Table listing papers not referred to in data cards, including authors like Smith, Merril, and others.

Ξ(1940)

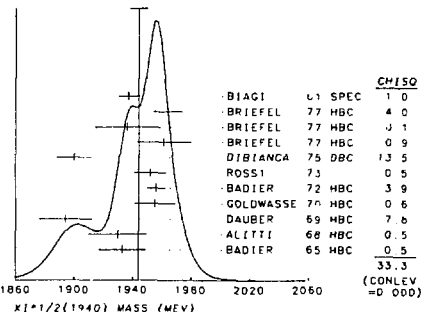
52 XI*1/2(1940) JP = 1 1/2

WE LIST UNDER XI(1940) EVERYTHING REPORTED IN THE MASS RANGE 1875-2000 MEV

52 XI*1/2(1940) MASS (MEV)

Table of mass measurements for Xi*(1940) with columns for mass (M), width (W), and various fit parameters.

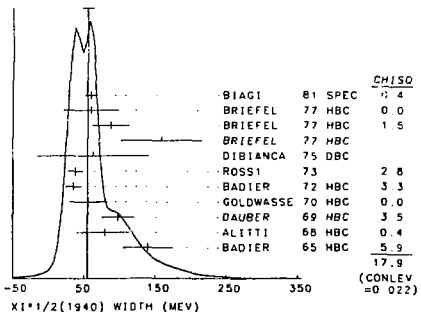
AVERAGE MEANINGLESS SCALE FACTOR = 1.81 (SEE IDEOGRAM BELOW)



52 XI*1/2(1940) WIDTH (MEV)

Table of width measurements for Xi*(1940) with columns for width (W), mass (M), and various fit parameters.

AVERAGE MEANINGLESS SCALE FACTOR = 1.51 (SEE IDEOGRAM BELOW)



Baryons

 Λ_c^+ , $\Sigma_c(2450)$, Λ_b^0 DIBARYONS

Data Card Listings

For notation, see key at front of Listings.

CHARMED BARYONS

 Λ_c^+

73 LAMBDA/C+(2282, JP=)

SEE STABLE PARTICLE DATA CARD LISTINGS

 $\Sigma_c(2450)$

104 SIGMA/C(2450, JP=)

THE SIGMA/C DECAYS TO LAMBDA/C P1, AND THE SCHEM IN MASSES HERE REFLECTS THAT IN MEASUREMENTS OF THE LAMBDA/C MASS (THE HIGHER MASSES ARE PRESENTLY FAVORRED). THE IMPRESSIVE AGREEMENT IN THE SIGMA/C-LAMBDA/C MASS DIFFERENCE SEPARATELY INDICATES THIS TO BE THE CASE, RATHER THAN THAT TWO STATES (THE SIGMA/C AND YVC/C) ARE BEING OBSERVED.

THIS PARTICLE IS AT ABOUT THE 2-AND-1/2-STAR LEVEL. A DEFINITIVE EXPERIMENT IS NEEDED.

104 SIGMA/C(2450) MASS (MEV)

M	1	2626.0	12.0	CAZZOLI	75 HBC	** NU P IN BNL 7-F1	3/77
M <td>1</td> <td>2450.0</td> <td>0.1</td> <td>KNAPP</td> <td>76 SPEE - GAMMA BE</td> <td></td> <td>11/81*</td>	1	2450.0	0.1	KNAPP	76 SPEE - GAMMA BE		11/81*
M <td>1</td> <td>2459.1</td> <td>0.0</td> <td>BARTSH</td> <td>77 DEC ** NU D IN 12-FE</td> <td></td> <td>3/77</td>	1	2459.1	0.0	BARTSH	77 DEC ** NU D IN 12-FE		3/77
M <td>6</td> <td>2475.0</td> <td>0.0</td> <td>BALTY</td> <td>79 HBC ** NU H-M IN 15-F1</td> <td>12/79</td> <td></td>	6	2475.0	0.0	BALTY	79 HBC ** NU H-M IN 15-F1	12/79	
M <td>1</td> <td>2457.0</td> <td>0.0</td> <td>CALICHO</td> <td>80 HBC ** NU P IN RESC-TST</td> <td>11/81*</td> <td></td>	1	2457.0	0.0	CALICHO	80 HBC ** NU P IN RESC-TST	11/81*	
M <td>1</td> <td>2454.0</td> <td>0.0</td> <td>BOSETTI</td> <td>82 HBC ** NU P IN RESC</td> <td>3/82*</td> <td></td>	1	2454.0	0.0	BOSETTI	82 HBC ** NU P IN RESC	3/82*	

104 (SIGMA/C)-(LAMBDA/C) MASS DIFFERENCE (MEV)

J	1	166.0	15.0	CAZZOLI	75 HBC	** NU P IN BNL 7-F1	11/81*
J <td>6</td> <td>165.0 <td>3.0</td> <td>BALTY</td> <td>79 HBC ** NU H-M IN 15-F1</td> <td>12/79</td> <td></td> </td>	6	165.0 <td>3.0</td> <td>BALTY</td> <td>79 HBC ** NU H-M IN 15-F1</td> <td>12/79</td> <td></td>	3.0	BALTY	79 HBC ** NU H-M IN 15-F1	12/79	
J <td>1</td> <td>168.0</td> <td>0.0</td> <td>CALICHO</td> <td>80 HBC ** NU P IN RESC-TST</td> <td>11/81*</td> <td></td>	1	168.0	0.0	CALICHO	80 HBC ** NU P IN RESC-TST	11/81*	
J <td>1</td> <td>166.0 <td>11.0</td> <td>BOSETTI</td> <td>82 HBC ** NU P IN RESC</td> <td>3/82*</td> <td></td> </td>	1	166.0 <td>11.0</td> <td>BOSETTI</td> <td>82 HBC ** NU P IN RESC</td> <td>3/82*</td> <td></td>	11.0	BOSETTI	82 HBC ** NU P IN RESC	3/82*	

104 SIGMA/C(2450) PARTIAL DECAY MODES

P1 SIGMA/C(2450) INTO LAMBDA/C+ P1 2282+139

REFERENCES FOR SIGMA/C(2450)

CAZZOLI 75 PPL 34 1125 *CHDPS,COMMONLY-LOUTIT+MURTHCH. + (BNL)
 KNAPP 76 PPL 37 862 *LE,LEUNG,SMITH, + COLU+PANA+ILL+PNAL
 HARTSH 77 PP 015 1 *LBRICK,DMRECK,MUSSARAVE, + (ANL+PURD)
 BALTY 79 PPL 62 1721 *CARDONALIS+FRONCH+FRS, + (COLU+BNL)
 CALICHO 80 PL 93B 521 *BAR+HBI+BRUFF+CERN+EROL+RHEL+ST+LOUC
 BOSETTI 82 PL 1096 236 *GRASSLER, + (MCH+BRONN+CERN+M+DAP)

THEORY AND REVIEW

DEERJUKA 75 PP 012 147 *GEORGI+GLASHOW (HARV)
 LEE 77 PP 015 157 *QUIFF+PODNER (HARV)
 TOSLING 81 PPL 10 57 C M TOSLING (BNL)

BOTTOM (BEAUTY) BARYON

 Λ_b^0

40 LAMBDA/B(5500, JP=)

SEE STABLE PARTICLE DATA CARD LISTINGS

Note on Dibaryon Resonances

(by L.D. Roper, Virginia Polytechnic Institute and State University)

The first modern theoretical discussion of dibaryon resonances was probably by Oakes.¹ The first experimental hint of them was in a Λp invariant mass distribution by Dahl,² and in NN partial-wave analyses by Arndt³ for the 1D_2 state in a pp partial-wave analysis. [The notation is $(2S+1)L_J$, where S is the total spin, L is the orbital angular momentum, and J is the total angular momentum.]

Dibaryon popularity rose dramatically in 1977 when strong energy dependence was unexpectedly observed in pp polarization experiments at Argonne.⁴ Also, in that year and the next, Hoshizaki claimed the existence of dibaryon resonances in a pp partial-wave analysis.⁵ In the same year Jaffe gave a detailed theoretical treatment of multi-quark states.⁶

Since the last issue of this Review some new pp and np polarization experiments⁷ have been reported which, along with other new and old NN scattering data, have been newly partial-wave analyzed.⁸ There are also partial-wave analyses (given in the Listings) of πd elastic, $\pi d + \pi p n$, $pp + \pi d$, and $yd + pn$ scattering. Most of these strongly indicate the existence of dibaryon resonances in the 1D_2 and 3F_3 nucleon-nucleon states, and some indicate possible resonances in the 1S_0 , 3S_1 , 3P_1 , 3P_2 , 3D_3 , 1F_3 , or 1G_4 states.

The Pauli principle restricts two nucleons to be in one of the following states:

$$I=0: ({}^3S_1, {}^3D_1), ({}^1P_1, {}^3D_2), ({}^3D_3, {}^3G_3), ({}^1F_3, {}^3G_4), \dots$$

$$I=1: ({}^1S_0, {}^3P_0, {}^3P_1, ({}^3F_2, {}^3F_2), {}^1D_2, {}^3F_3, ({}^3F_4, {}^3H_4), \dots$$

Here the states that couple together (same J^P) are enclosed in parentheses. Similarly, only certain quantum number sets are allowed for $\Lambda\Lambda$, etc.

In the Listings below, we separate the determinations of pole positions and Breit-Wigner parameters. If there is a resonance, the pole must occur on the second sheet for the elastic channel; it may be a bound state or resonance for inelastic channels. The advantages of pole parameters over Breit-Wigner parameters are discussed briefly in the "Note on N's and Δ 's."

The idea that exotic resonances are really 'pseudo-resonances'⁹ has recently taken a new turn.¹⁰ The idea is that box diagrams (e.g., involving $N\Delta$ in NN scattering) create resonance-like loops in the Argand diagram without resonance poles actually existing. The question is whether poles would be created when one unitarized the box diagrams in order to calculate physical scattering amplitudes. Kloeit and Tjon¹⁰ have recently shown that a model exists in which, indeed, that is the case. However, resonance hunters should definitely

Data Card Listings

For notation, see key at front of Listings.

Baryons
DIBARYONS

report pole positions rather than looping Argand diagrams in the future. All who still suggest that the NN 1D_2 or 3F_3 resonance-like structure is due to some dynamics other than resonances must take their case to the world collection of NN scattering data in the form of a detailed partial-wave analysis.

The dinucleon resonances also communicate with the Yd and πd channels. There is not much Yd data, and the multipole analysis does not yield much certainty about which dibaryons are involved. In the πd case, uncertainties abound and the partial-wave analysis yields poor fits compared to the analyses of NN data. Nevertheless, the existence of dibaryons appears to be corroborated by the Yd and πd data.

Only a few papers about strange dibaryon states have appeared since the last issue of this Review. It appears that the S=-1 dominant resonance is in the 3S_1 state, an SU(3) partner of the deuteron. An excellent review of the S=-1 situation is given by Dalitz,¹¹ who concludes that the S=-1 3S_1 resonance pole probably exists.

For a more detailed recent review of dibaryons, see Hoshizaki.¹²

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DIBARYONS

S=0 DIBARYONS

106 BARYON NUMBER 2, STRANGENESS 0 STATES

IN THIS SECTION WE USE THE FOLLOWING ABBREVIATIONS FOR TYPES OF ANALYSES--

GDPN GAMMA D -> P N PARTIAL-WAVE ANALYSIS RESONANCE
 PARAMETRIZATION
 NN FIT TO NN ELASTIC PARTIAL-WAVE ANALYSIS RESULTS
 NHF FIT TO NN FORWARD AMPLITUDES
 PID PC-D ELASTIC PARTIAL-WAVE ANALYSIS RESONANCE
 PARAMETRIZATION
 PID1 PC-D -> P1- P N AMPLITUDE ANALYSIS RESONANCE
 PARAMETRIZATION
 PP0D P, P -> P1-D PARTIAL-WAVE ANALYSIS RESONANCE
 PARAMETRIZATION

NN(2170)

I=1, 1D_2

106 B=2, S=0, I02 -- BREIT-WIGNER MASS (MEV)

BREIT-WIGNER MASS APPROXIMATELY EQUALS REIPOLE POSITION(s).

N	(2170.)	HOSHIZAKI 79 NH	I02 ASSUMED BCKGRND	1/82*
N	(2180.)	ARVIEUX 80 PID	I02	1/82*
N	A (2185.)	KAM0 80 PP0D I02		1/82*
N	(2170.)	HOSHIZAKI 81 PID1 I02		1/82*
N	(2140.)	KANAI 81 PID I02 SOL. B AND C		1/82*
N				
N		KAM0 80 DID ACT TRY FITS WITH FEWER THAN SIX RESONANCES.		1/82*
N		KANAI 81 FIT WITH NO RESONANCES WAS VERY POOR AND DID NOT TRY		1/82*
N		OTHER FITS WITH FEWER THAN FOUR RESONANCES.		1/82*

106 B=2, S=0, I02 -- BREIT-WIGNER WIDTH (MEV)

BREIT-WIGNER WIDTH APPROXIMATELY EQUALS 2 TIMES INIPOLE POSITION(s).

N	102, 105D.	HOSHIZAKI 79 NH	I02 ASSUMED BCKGRND	1/82*
N	A (134.)	KAM0 80 PP0D I02		1/82*
N	(175.)	HOSHIZAKI 81 PID1 I02		1/82*
N	(156.)	KANAI 81 PID I02 SOL. B		1/82*
N	B (154.)	KANAI 81 PID I02 SOL. C		1/82*

106 B=2, S=0, I02 -- BREIT-WIGNER ELASTICITY

BREIT-WIGNER ELASTICITY APPROXIMATELY EQUALS ABSOLUTE VALUE OF POLES/INIPOLE POS(s).

N1	(0.1)	HOSHIZAKI 79 NH	I02	1/82*
----	-------	-----------------	-----	-------

Data Card Listings

For notation, see key at front of Listings.

Baryons
DIBARYONS

107 B=2, S=1 -- BREIT-WIGNER WIDTH (MEV)

BREIT-WIGNER WIDTH APPROXIMATELY EQUALS 2 TIMES (MEV/PC POSITION).		COHN 64 LPJM Q=0		1/82*
M	I 20.	CLINE	68 LPJM 351 Q=1	1/82*
M	H 10, CA 1855	JAIN	69 LPJM Q=1	1/82*
M	J 100.	TAN	69 LPJM Q=1	1/82*
M	H 7.0	TAN	69 LPJM Q=1	1/82*
M	H 18.11	TAN	69 LPJM Q=1	1/82*
M	H 10.1	EASTWOOD	71 LPJM Q=1	1/82*
M	K 5.	SIM	71 LPJM Q=1	1/82*
M	L 20.6	SHARBAZI	73 LPJM Q=1 2325 PEAK	1/82*
M	H 171.11	SHARBAZI	73 LPJM Q=1 2251 PEAK	1/82*
M	H 159.1	SHARBAZI	73 LPJM Q=1	1/82*
M	H 5.9	BRUHM	77 LPJM Q=1	1/82*
M	O 125.1	GOVAL	80 SPJM Q=0	1/82*

107 B=2, S=1 -- REIP/PC POSITION (MEV)

REIP/PC POSITION APPROXIMATELY EQUALS BREIT-WIGNER MASS.				1/82*
RE	C (2132.)	NAGELS	70 BB 351 Q=1	1/82*
RE	C (2137.)	NAGELS	70 BB 351 Q=0	1/82*
RE	H (2129.)	DOSCH	80 LPJM 351 Q=1	1/82*
RE	O (2127.)	TAKAHASHI	80 BB 351 Q=0, 1+2	1/82*
RE	E (2148.)	TAKAHASHI	80 BB 1P Q=0, 1+2	1/82*
RE	C	NAGELS 70 REPORTS POLE POSITION FOR TWO DIFFERENT 351 CHARGE		1/82*
RE	C	STATES.		1/82*
RE	H	K = D TO P = LAMBDA P.		1/82*

107 B=2, S=1 -- IM/PC POSITION (MEV)

IM/PC POSITION APPROXIMATELY EQUALS BREIT-WIGNER WIDTH DIVIDED BY 2.				1/82*
IM	G (22.4)	NAGELS	70 BB 351 Q=1	1/82*
IM	C (22.6)	NAGELS	70 BB 351 Q=0	1/82*
IM	H (5.1)	DOSCH	80 LPJM 351 Q=1	1/82*
IM	H (5.1)	TAKAHASHI	80 BB 351 Q=0, 1+2	1/82*
IM	H (6.1)	TAKAHASHI	80 BB 1P Q=0, 1+2	1/82*

REFERENCES FOR B=2, S=1 STATES

COHN	64 PRL 22 668	H.D.COHN, W.H.BHATT, W.M. BUGG	(ORNL+TENN)
CLINE	68 PRL 20 1452	D.CLINE, R. LAURMAN, J. HARP	(MISC)
ALEXANDE	69 PRL 22 483	ALEXANDER, HALL, JEN, KALUS, KERNAN	(LBL+UC)
JAIN	69 PR IPT 1816	P.L.JAIN	(IUPUI)
TAN	69 PRL 23 395	T.H.TAN	(SLAC)
EASTWOOD	71 PR D3 2463	FERT, HEATHCOTE, ISLANS, (BIRM+EDIN+GLAS+LOIC)	
SIMS	71 PR D3 1192	M.FINELL, ALANRISHI, BRUCKER, LANUTTI	(FSU)
SHARBAZI	73 PR D3 39	B. SHARBAZI, A.V. TIMONOVA	(JINR)
SODHI	75 NP 807 403	A. SODHI, D. GOVAL	(DELH)
GOVAL	77 NP 8124 45	G. GRIM, H. HEPF, STROEBELE, THOEL	(HEFD+MPI)
GOVAL	78 PR D18 948	D. GOVAL, A. SODHI	(DELH)
NAGELS	79 PR D20 1033	M. NAGELS, S. P. JUKEN, J. DESHART	(MIL)
DOSCH	80 PR 3 245	H. DOSCH, I. STAMATESCU	(HEID)
GOVAL	80 PR 64 700	D. GOVAL, J. MISRA	(DELH)
TAKAHASHI	80 NP 8326 367	TAKAHASHI, TAMURA, KITAHARA, KUME	(TOYO)

PAPERS NOT REFERRED TO IN DATA CARDS

ALLEN	68 PR 173 1452	J.A. PIRQUE	(IPPEN)
BUNEL	70 PR D2 98	M.ASHOQ, S.HALIM	(FOM+HEID)
GOVAL	71 PR D3 1259	M. BRICE, F. FIELDS, H.M.G., H.F.FES	(MICH+ANL)
KROPP	71 PR D7 13	O.P. GOVAL	(DELH)
DOSCH	78 PR 718 4011	HALLBAUER, J. HALL, GARDOSCHEN, SPILLING	(LEIB)
MIZUNO	79 PR 2 1691	H.G. DOSCH, V. HEPF	(HEID)
ROSEN	79 PR 849 217	T. MIZUNO	(TOYO)
DIAGOSTI	81 PL 1048 330	M. WANDERFELDE, H. ELQET, WICKENS*	(LOND+BRUK)
KIMURA	81 PR 65 649	M. KIMURA, Y. TAMURA, Y. TERASHIMA*	(ICM+SAIG+VAND)
TOBER	81 NP 1362 405	G. TOBER, A. J. GIL, L. EISENBERG	(TOYO)

S=-2 DIBARYON

108 BARYON NUMBER 2, STRANGENESS -2 STATES

IN THIS SECTION WE USE THE FOLLOWING ABBREVIATIONS FOR MEASURED QUANTITIES--

XPM	XT-P INVARIANT MASS
LLIM	LAMBDA-LAMBDA INVARIANT MASS

108 B=2, S=-2 -- MASS (MEV)

M	D (2367.1) (4.1)	HELLIERE	72 LLIM Q=0 GAUSSIAN FIT	1/82*
M	D (2365.3) (9.6)	SHARBAZI	73 LLIM Q=0	1/82*
M	P (2440.1)	GOVAL	80 PRJM Q=0	1/82*
M	P	C = 0 TO K1 = P #0.		1/82*
M	Q	V B TO LAMBDA LAMBDA X AND P1 = P TO LAMBDA LAMBDA X FDP P 14		1/82*
M	R	CARROLL 12.		1/82*
M	R	GOVAL 80 ALSO SEES A SHOULDER AT 2360 MEV.		1/82*

108 B=2, S=0 -- WIDTH (MEV)

M	P (15.) (4.1)	HELLIERE	72 LLIM Q=0 GAUSSIAN FIT	1/82*
M	Q (47.0) (15.7)	SHARBAZI	73 LLIM Q=0	1/82*

REFERENCES FOR B=2, S=-2 STATES

HELLIERE	72 PL 398 671	HELLIERE, HAVENP*	(IPPEN+CEPNI+UT+LOIC)
SHARBAZI	73 PR 853 19	B. SHARBAZI, A.V. TIMONOVA	(IUPUI)
GOVAL	80 PR 321 607	D. GOVAL, J. MISRA, A. SODHI	(DELH)
CARROLL	78 PL 41 777	M. CHANG, J. MASON, P. C. IA, P. I.*	(MIL+DTN)

Appendix I

TEST OF $\Delta I=1/2$ RULE FOR HYPERON DECAYS

O. E. Overseht
University of Michigan

1. Nonleptonic decay Amplitudes

In this edition we again use the new convention for the amplitudes A and B adopted in 1973. Some theorists have suggested that dimensionless amplitudes are more useful to them than the ones appearing in the literature. Berge¹ used a convention with A and B in units of $\text{sec}^{-1/2}$. Samios² used a convention which gave A and B in units of $(\text{MeV}\cdot\text{sec})^{-1/2}$. Following is the convention suggested by Jackson³, which gives dimensionless A and B.

The effective Lagrangian density for nonleptonic hyperon decays ($B_1 \rightarrow B_2 + \pi$) can be written

$$L_{\text{eff}} = G\mu_c^2 \bar{\psi}_2(A+B\gamma_5)\psi_1\phi_\pi,$$

where $G=10^{-5}m_p^{-2}$ is a coupling constant characteristic of first-order weak decays, μ_c is the charged pion mass, and A and B are dimensionless complex numbers giving the relative amplitudes of the parity-violating and parity-conserving decays, respectively. The matrix γ_5 is to be taken in the Pauli form, $\gamma_5 = \begin{pmatrix} 0 & -1 \\ 1 & 0 \end{pmatrix}$. The invariant amplitude for the decay is

$$M = G\mu_c^2 [\bar{u}(p)(A+B\gamma_5)u(P)],$$

where P is the 4-momentum of the decaying hyperon of mass M, and p is the 4-momentum of the baryon decay product of mass m. With the normalization convention, $\bar{u}_i u_i = 2m_i$, the Pauli form of the matrix element in the rest frame of the decaying hyperon is

$$M = G\mu_c^2 \langle x_2 | \sqrt{2M(E+m)}A + \sqrt{2M(E-m)}B\hat{\sigma} \cdot \hat{q} | x_1 \rangle,$$

where E is the total energy of the final baryon and \hat{q} is a unit vector in the direction of motion of the final baryon. Comparison with Sec. VI D of the text shows that the amplitudes s and p defined there are proportional to A and B:

$$\frac{p}{s} = \left(\frac{E-m}{E+m} \right)^{1/2} \frac{B}{A} = \left[\frac{(M-m)^2 - \mu^2}{(M+m)^2 - \mu^2} \right]^{1/2} \frac{B}{A}.$$

Here μ is the mass of the pion entering the decay. The parameters α , β , and γ can therefore be expressed in terms of A and B, rather than s and p, if desired.

The decay rate for $B_1 \rightarrow B_2 + \pi$ is

$$\Gamma = \frac{G^2 \mu_c^4}{8\pi} q \left\{ \left[\frac{(M+m)^2 - \mu^2}{M^2} \right] |A|^2 + \left[\frac{(M-m)^2 - \mu^2}{M^2} \right] |B|^2 \right\},$$

where q is the c.m. momentum of the decay products. For reference, the dimensionless constant in this expression has the value $(G^2 \mu_c^4 / 8\pi) = 1.9488 \times 10^{-15}$.

Table I summarizes the amplitudes A and B for the nonleptonic decays of the Λ , Σ , and Ξ hyperons. These amplitudes have been calculated by using the experimental data for mean lives, branching ratios, and the decay asymmetry α given in the Stable Particle Table of this Review. Time-reversal invariance is assumed and final-state interactions are neglected, so A and B are taken to be relatively real. The subscript on the hyperon refers to the sign of the decaying pion. The statistical correlation coefficient

$$C_{AB} = \frac{\langle \Delta A \Delta B \rangle}{\sqrt{\langle \Delta A^2 \rangle \langle \Delta B^2 \rangle}}$$

is also given. The absolute signs of A and B have been assigned, using the following convention. Taking $A(\Lambda_0^0)$ as positive, the other S-wave decay amplitudes are chosen to give an approximate fit to the triangular relationships

$$\sqrt{2}A(\Sigma_0^+) + A(\Sigma_+^+) = A(\Sigma_-^-) \quad \text{and} \quad \sqrt{3}A(\Sigma_0^+) + A(\Lambda_0^0) = 2A(\Xi_-^-).$$

The signs of the B amplitudes relative to those of the corresponding A amplitudes are determined by the sign of the appropriate α decay parameter.

Table I

M	$\rightarrow m + \mu$	A	B	C_{AB}
Λ_0^0	$\rightarrow p + \pi^-$	1.47 ± 0.01	9.98 ± 0.24	-0.289
Λ_0^0	$\rightarrow n + \pi^0$	-1.07 ± 0.01	-7.14 ± 0.56	-0.740
Σ_+^+	$\rightarrow n + \pi^+$	0.06 ± 0.01	19.07 ± 0.07	-0.038
Σ_0^+	$\rightarrow p + \pi^0$	1.48 ± 0.05	-12.04 ± 0.58	0.982
Σ_-^-	$\rightarrow n + \pi^-$	1.93 ± 0.01	-0.65 ± 0.07	0.003
Ξ_0^0	$\rightarrow \Lambda + \pi^0$	1.55 ± 0.03	-5.56 ± 0.33	-0.148
Ξ_-^-	$\rightarrow \Lambda + \pi^-$	2.04 ± 0.01	-7.49 ± 0.28	0.237

2. Tests of the $\Delta I=1/2$ Rule(a) Λ Decay

For Λ decay the $\Delta I=1/2$ rule predicts that $\Gamma_0/\Gamma_- = 0.50$ and $\alpha_0 = \alpha_-$. In order to determine the magnitude of possible $\Delta I=3/2$ amplitudes present we write the linear expressions⁴ for the $\Delta I=3/2$ A- and B-wave amplitudes in terms of $\Delta\alpha$, where $\Delta\alpha$ is the measured value of α_0/α_- minus the predicted value, and in terms of $\Delta\Gamma$ similarly defined. Evaluating these we find

$$\Delta\alpha = -1.54 (A_3/A_1) + 1.61(B_3/B_1),$$

$$\Delta\Gamma = 1.84 (A_3/A_1) + 0.25(B_3/B_1).$$

Here the $\Delta I=3/2$ amplitudes are expressed relative to the $\Delta I=1/2$ amplitudes. The numerical values of the coefficients depend on the ratio B/A . The uncertainties in the coefficients are small compared to the uncertainties in $\Delta\alpha$ and $\Delta\Gamma$. Final-state πN interactions have been included in these relations but have a very small effect. From the Stable Particle Table,

$$\Delta\alpha = 0.006 \pm 0.066, \quad \Delta\Gamma = 0.058 \pm 0.012,$$

and hence

$$(A_3/A_1) = 0.027 \pm 0.008$$

and

$$(B_3/B_1) = 0.030 \pm 0.037.$$

The possible 3% $\Delta I=3/2$ A-wave amplitude is due to the disagreement of decay rates with prediction. At this level the results are sensitive to electromagnetic corrections. However, in Λ decay the phase space correction and the other radiative corrections appear to be about equal in magnitude and have opposite signs,^{5,6} and hence cancel each other in the correction to the decay rates.

(b) Ξ Decay

The analysis for Ξ decay is very similar to that for Λ decay. If the $\Delta I=1/2$ rule is valid, $\Gamma_0(\Xi^0)/\Gamma_-(\Xi^-) = 0.50$ and $\alpha_0 = \alpha_-$. For this case the expressions linear in $\Delta I=3/2$ A- and B-wave amplitudes are⁴

$$\Delta\alpha = 1.35(A_3/A_1) - 1.35(B_3/B_1),$$

$$\Delta\Gamma = -1.43(A_3/A_1) - 0.07(B_3/B_1).$$

From the Stable Particle Table,

$$\Delta\alpha = -0.05 \pm 0.06, \quad \Delta\Gamma = 0.066 \pm 0.020,$$

and we find

$$(A_3/A_1) = -0.046 \pm 0.014$$

and

$$(B_3/B_1) = -0.01 \pm 0.04.$$

(c) Σ Decay

The traditional test of the $\Delta I=1/2$ rule in Σ decay is that the amplitudes satisfy the relationship

$$\sqrt{2}\Sigma_0^+ + \Sigma_+^+ - \Sigma_-^- = 0.$$

Graphically this is equivalent to closing the Σ triangle when the amplitudes are plotted on A, B axes. Including $\Delta I \geq 3/2$ amplitudes in Σ decay analysis, the " Σ triangle" relationship becomes

$$\sqrt{2} A_0 + A_+ - A_- = -3\sqrt{2}/5 A_3 + \frac{2}{\sqrt{15}} A_5,$$

where A_3 and A_5 are $\Delta I=3/2$ and $\Delta I=5/2$ amplitudes, respectively. There is a similar equation for the B amplitudes. From Table I,

$$\sqrt{2} A_0 + A_+ - A_- = 0.22 \pm 0.09$$

and

$$\sqrt{2} B_0 + B_+ - B_- = 2.7 \pm 1.0.$$

If we neglect the $\Delta I=5/2$ amplitudes and assume all amplitudes to be real we can solve for possible $\Delta I=3/2$ amplitudes. The result is

$$\frac{A_3}{A_-} = -0.061 \pm 0.024$$

and

$$\frac{B_3}{B_+} = -0.074 \pm 0.027.$$

Thus for hyperon decay, present experimental data limit $\Delta I=3/2$ amplitudes to less than about 5%.

3. The Lee-Sugawara Relation

From Table I the Lee-Sugawara relation,^{7,8} $\sqrt{3}\Sigma_0^+ + \Lambda_0^0 - 2\Sigma_-^- = 0$, is satisfied to -0.05 ± 0.11 for the A amplitudes, and to 4.1 ± 1.8 for the B amplitudes.

References

1. J. P. Berge, in *Proceedings of the 13th International Conference on High-Energy Physics, Berkeley, (1966)* (University of California Press, Berkeley, 1967), p. 46.
2. N. P. Samios, International Conference on Weak Interactions, Argonne, (1965), p. 189.
3. J. D. Jackson, private communication (1973).
4. See O. E. Overseth and S. Pakvasa, *Phys. Rev.* **184**, 1663 (1969). The expression for Γ_0/Γ_- for Λ decay should read

$$\frac{\Gamma_0}{\Gamma_-} \approx \frac{1}{2} \left\{ 1 + 3\sqrt{2} \times \left[\frac{S_{11}S_{33}\cos(\delta_1 - \delta_3) + P_{11}P_{33}\cos(\delta_{11} - \delta_{31})}{S_{11}^2 + P_{11}^2} \right] \right\}.$$
5. See A. A. Belavin and I. M. Narodetsky, *Yadern. Fiz.* **8**, 978 (1968) [*Soviet J. Nucl. Phys.* **8**, 568 (1969)].
6. G. W. Intemann, private communication (1973).
7. See B. W. Lee, *Phys. Rev. Lett.* **12**, 83 (1964).
8. See H. Sugawara, *Prog. Theor. Phys.* **31**, 213 (1964).

Appendix II

A. SU(3) CLASSIFICATION OF BARYON RESONANCES

A. Barbaro-Galtieri
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It is established that a symmetry higher than SU(3) is necessary to classify the known baryon resonances. However, many higher-symmetry schemes have been proposed, and even for SU(6) various versions exist (for a review see Dalitz¹). Since it is not clear which one of these schemes best fits the data, we do not review them here, but we report once again fits of baryon states into SU(3) multiplets.

For the reader's convenience, we collect here the relevant formulae.

Exact SU(3) symmetry predicts that all the members of a multiplet should have the same mass and the same couplings for decays into other multiplets. It has been found, however, that the members of the octet of stable baryons lie within 20% of their mean mass; therefore a symmetry-breaking interaction has been introduced by Gell-Mann and Okubo independently.² In addition, for the isospin-0 vector mesons (ω and ϕ), an additional symmetry-breaking interaction has been introduced by Sakurai³ to take care of octet-singlet mixing. The relevant formulae for masses and decay rates are given below.

Mass Formulae

Broken SU(3) gives:

$$\text{Decuplet} \quad \Delta - \Sigma = \Sigma - \Xi = \Xi - \Omega \quad \text{GMO} \quad (1)$$

$$\text{Octet} \quad 2(N+\Xi) = 3\Lambda + \Sigma \quad \text{GMO} \quad (2)$$

$$\text{Octet-singlet mixing} \quad \left\{ \begin{array}{l} \sin^2\theta = \frac{\Lambda - M_g}{\Lambda - \Lambda'} \\ M_g = \frac{2(N+\Xi) - \Sigma}{3} \end{array} \right. \quad \text{Mixing angle} \quad (3)$$

Here GMO stands for the Gell-Mann-Okubo formula; the particle symbol indicates its mass. The formulae would be the same if squared masses were used. For the nonet case, Λ is the "mostly-octet" particle, Λ' is the "mostly-singlet" particle.

Decay Rates

In terms of a relativistically invariant matrix element T , the decay rate for 2-body decay of a resonance of mass M_R is

$$\Gamma \propto \frac{|T|^2 R_2}{M_R}, \quad (5)$$

where $R_2 = k/M_R$ is the 2-body phase space factor. Since the numerator is an invariant, and since Γ must transform as $1/E$, we introduce the denominator M_R .⁵

For meson decays (see below) the rates are calculated according to Eq. (5); for baryon resonance decays into $1/2^+$ baryons and 0⁻ mesons, one next takes into account the fact that spin sums in $|T|^2$

introduce another factor M_R , canceling the $1/M_R$. We are then left with

$$\Gamma = \frac{|T|^2 k}{M_R} M_N, \quad \text{for baryons} \quad (5')$$

$$= \frac{|T|^2 k}{M_R^2} M_N^2, \quad \text{for mesons.} \quad (5'')$$

The powers of the nucleon mass M_N or M_N^2 have been introduced so that we can treat $|T|$ as dimensionless.

$|T|^2$ contains centrifugal barrier factors, which we call B_ℓ . We then have

$$\left. \begin{array}{l} \text{Decuplet} \\ \text{Singlet} \end{array} \right\} \quad \Gamma = (c_g)^2 B_\ell(k) \frac{M_N}{M_R} k \quad (6)$$

$$\text{Octet} \quad \Gamma = (c_D g_D + c_F g_F)^2 B_\ell(k) \frac{M_N}{M_R} k \quad (7)$$

$$\left. \begin{array}{l} \text{Octet-} \\ \text{singlet} \\ \text{mixing} \end{array} \right\} \quad \begin{array}{l} \Lambda = G_8 \cos\theta + G_1 \sin\theta \\ \Lambda' = -G_8 \sin\theta + G_1 \cos\theta \end{array} \quad (8)$$

$$\text{with} \quad \begin{array}{l} G_8 = c_D g_D + c_F g_F \\ G_1 = c_1 g_1 \end{array} \quad (9)$$

Here c_i are the SU(3) coefficients with the sign convention adopted in this article [see note in the Table of SU(3) Isoscalar Factors and Fig. 2 in the text]. M_N is the nucleon mass, M_R is the resonance mass for which Γ is calculated, k is the center-of-mass momentum for the channel being considered, and g_i are the relevant couplings. For the case of singlet-octet mixing, Eq. (8) has to be used in conjunction with (6) and (7). G_8 and G_1 represent the couplings for the multiplet, and Λ and Λ' represent the couplings for the physical states.

The relation between g_D , g_F , and the parameter α is

$$\alpha = \left[1 + \frac{\sqrt{5}}{3} \frac{g_F}{g_D} \right]^{-1}. \quad (10)$$

Exact SU(3) predicts that the couplings g_i for all the members of a multiplet are the same; however, since the symmetry is broken for the masses, it is probably broken for the widths. In the case of the $3/2^+$ decuplet, broken SU(3) sum rules have been derived by Becchi,⁶ Gupta,⁷ and Konuma⁸ independently. The form derived by Gupta relates the g_i for the members of the decuplet by the relation

$$2(\Delta + \Xi) = 3\Sigma^*(\Lambda\pi) + \Sigma^*(\Sigma\pi), \quad (11)$$

where $\Sigma^*(\Lambda\pi)$ is the coupling for the $\Sigma(1385) \rightarrow \Lambda\pi$ decay and $\Sigma^*(\Sigma\pi)$ is the coupling for the decay $\Sigma(1385) \rightarrow \Sigma\pi$.

As mentioned in the text (Sec. IV B) the determination of the relative signs of resonant amplitudes can be useful in making an SU(3)

assignment of resonances. In fact the resonant amplitude $T \propto \sqrt{x_e x_i} \propto G_e G_i$, where the subscript e refers to the elastic channel and the G_e , G_i are the couplings of Eqs. (6) through (9). Assuming that all g_i are positive, the sign of the G_i are dependent upon the sign of the Clebsch-Gordan coefficients c_i . Once a sign convention is adopted (we use the Levi-Setti⁹ convention, see Fig. 2 in the text) and the signs for a Σ state ($I=1$) and a Λ state ($I=0$) of known SU(3) assignment have been chosen for reference, the signs of all the other amplitudes can be useful in determining multiplet assignments. For exact SU(3) all the decays of members of a decuplet have the same sign. For octets the relative sign depends upon the value of g_D/g_F and the mixing angle, as seen from Eqs. (7) through (9).

Fits to the Data

Fits of baryon decay rates within SU(3) can be found in, among others, papers by Tripp,^{10,11} Levi-Setti,⁹ Samios,¹² and Plane.¹³ More recent fits were made by Barbaro-Galtieri¹⁴ and Samios.¹⁵ A fit of the decay rates within SU(6)_w may be found in Litchfield et al.¹⁶ Analysis of the baryon mass spectrum using the quark shell model has been done by Jones et al.¹⁷ An analysis of baryon couplings in a quark model with chromodynamics has been done recently by Konjuk and Isgur.¹⁸

For our SU(3) analysis in fitting the data a choice for B_ℓ has to be made. Plane¹³ tried two forms for B_ℓ :

(a) The form $B_\ell = (kr)^6 D_\ell(kr)$, r being the radius of interaction and D_ℓ the polynomials in kr given by Blatt and Weisskopf.¹⁹ Usually r is taken to be 1 fermi.¹⁰

(b) The form $B_\ell = k \ell^{\ell}$.

However, for final results form (b) was chosen. A discussion of the differences among these two forms has been given by Barbaro-Galtieri.²⁰ As shown in Ref. 20, not only the values of the couplings, g_i , depend upon the form used for B_ℓ , but also the value obtained for the mixing angle. For the $3/2^-$ singlet, $\Lambda(1520)$, and the isospin-0 member of the octet, $\Lambda(1690)$, the mixing angles obtained in the two cases were

$$\theta_a = (-16.1^{+1.4}_{-1.3})^\circ, \theta_b = (-27.5^{+3.6}_{-3.4})^\circ,$$

in disagreement by a few standard deviations. However, if a radius of interaction of $r = 0.15$ fermi was used for form (a), the two values of θ agreed. This value of r does not fit resonance shapes when used in the Breit-Wigner resonant form. Samios¹⁵ used form (b) for B_ℓ .

Table I is a summary of the fits made by us (update of Barbaro-Galtieri¹⁴) using the barrier factor form (a) and exact SU(3). The values of the masses, widths, and amplitudes used in the fits are taken from this edition's Tables and Listings.

$1/2^-$ Nonet (Baryon-Eta Resonances)

For this nonet Eq. (7) was multiplied by the factor

$$\left[\frac{M_R - M_B}{\bar{M}_R - \bar{M}_B} \right]^2,$$

where M_B is the decay baryon and $\bar{M}_R - \bar{M}_B = 568$ MeV is the difference of the mean $1/2^-$ and $1/2^+$ baryon octet masses. This kinematic factor comes from PCAC arguments (i.e., the assumption that the axial

Table I. SU(3) baryon multiplets with two or more known members. Values of θ and α [defined by Eqs. (8) and (10)] are the result of fits made to all the measured 2-body decay rates of each multiplet.

J^P	Octet members ^a			Singlet	$\theta(\text{deg})^b$	α	
$1/2^-$	N(1535)	$\Lambda(1670)$	$\Sigma(1750)$	$[\Xi(1737)]$ $[\Xi(1850)]$	$\Lambda(1405)$	-31 ± 7 -22 ± 11	$\left\{ \begin{array}{l} 0.50 \pm 0.06 \\ 0.56 \pm 0.11 \end{array} \right\}^c$
$3/2^-$	N(1520)	$\Lambda(1690)$	$\Sigma(1670)$	$[\Xi(1819)]$	$\Lambda(1520)$	-21 ± 3	0.31 ± 0.04
$5/2^-$	N(1675)	$\Lambda(1830)$	$\Sigma(1775)$				1.18 ± 0.04
$5/2^+$	N(1680)	$\Lambda(1820)$	$\Sigma(1915)$	$[\Xi(2069)]$	$\Lambda(2110)$	22 ± 5	0.70 ± 0.03
	Decuplet members ^d			δ_{10}			
$3/2^+$	$\Delta(1232)$	$\Sigma(1385)$	$\Xi(1530)$	Ω^-	1.0-1.5	$\chi^2/\text{DF}=58/3$	
$7/2^+$	$\Delta(1950)$	$\Sigma(2030)$					

^a Masses in parentheses are the nominal masses used in the Baryon Table. The Ξ members have masses as calculated by using Eqs. (1) and (2) with the mixing angle θ derived from the decay widths.

^b See text for a discussion of the $1/2^-$ mixing angle.

^c The first values of θ and α are obtained by using a plus sign for the amplitudes of both $N(1535) \rightarrow N\eta$ and $\Lambda(1670) \rightarrow \Lambda\eta$. The second values use a minus sign for the second amplitude. Both fits, however, have a bad χ^2 , mostly due to the two baryon- η amplitudes.

^d Coupling constants updated from Ref. 14, using new $\Xi(1530)$ data.

vector current remains an octet in the presence of symmetry breaking) and it was advocated by Graham.²¹ For the $1/2^-$ nonet it was used in this form first by Gell-Mann.²² The couplings for the $A(1405)$ to the $N\bar{K}$ and $\Sigma\pi$ channels are the same as those used in Ref. 14.

$3/2^+$ Decuplet

The agreement among the coupling constants obtained for the four rates in this decuplet is very bad. The fit made using form (a) for B_c has $\chi^2=58$ for 3 degrees of freedom; the one made with form (b) for B_c has $\chi^2/DF=13/3$. The broken SU(3) relation (11), however, is very well satisfied.

B. SU(3) CLASSIFICATION OF MESON RESONANCES

All of the discussion above applies, except that for bosons the GMO formula is usually applied to the *square* of the masses, as opposed to the first power for fermions. Thus for example, Eq. (2) becomes

$$4\hat{K} = 3\hat{\eta} + \hat{\pi}. \quad (2')$$

The symbol \hat{K} was introduced by Glashow and Socolow⁴ for the *square* of the K mass, etc.

Because of the difference between Eqs. (5') and (5''), there is also an extra factor of (M_N/M_R) in Eqs. (6) and (7). The three established nonets ($0^-, 1^-, 2^+$) and their mixing angles are listed at the bottom of the Meson Table.

C. FLAVOR SYMMETRY BREAKING IN A UNITARY MIXING SCHEME

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Because of unitarity, the GMO formula is an approximation, which breaks down especially for broad resonances. The imaginary part of the mass matrix M^2 is related to the widths, and for mesons is of the form²³⁻²⁵

$$\text{Im } M_{ij}^2(s) = \text{Im } \Pi_{ij}(s) = \sum_{ab} g_{iab} g_{jab} k_a^{2\ell+1} s^{-1/2} F(k_a) \quad (12)$$

where i and j are resonance indices, Π the hadronic self-energy diagram, g_{iab} the coupling constant for the channel $i \rightarrow ab$, and F a hadronic form factor. Only the real part of M^2 approximately satisfies the GMO mass formula. More generally, one can assume $\text{Re } M^2$ to be given by quark masses in a bare-mass term plus a hadronic mass shift $\text{Re } \Pi_{ij}(s)$ satisfying analyticity constraints. In general, $\Pi_{ij}(s)$ is non-diagonal and breaks flavor symmetry through the different threshold positions. It is diagonalized by a complex orthogonal matrix which determines the generally complex mixing angles. The imaginary part of a mixing angle is related to the overlap between the resonance.²³

The masses and widths are given respectively by the real and imaginary parts of the eigenvalues of M^2 .

In this way one can, using both masses and widths as data, within the same framework, determine²⁶ both the deviation δ from ideal mixing (i.e., $\theta = 35.3^\circ + \delta$) and the Q_A-Q_B mixing angle ϕ listed at the end of the Meson Table.

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Appendix III

GROWTH OF INFORMATION

From time to time we have presented figures demonstrating the amount of experimental work which has gone into spectroscopy, and the amount of new information available as a result. The 1982 versions of these figures are shown as Figs. 1 and 2.

Figure 1 is a simple count of the number of meson resonances listed in the Tables, categorized as those "understood" -- i.e., all quantum numbers are believed known -- and those simply "listed". A rapid recent increase in both of these categories occurred because of the discovery of the J/ψ and related particles.

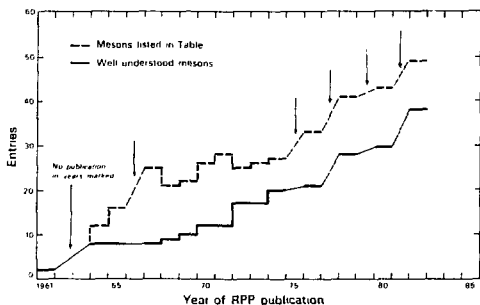


Fig. 1. Number of meson resonances listed in the Tables (dashed line) and those for which all quantum numbers are known (solid line), as a function of year of publication of the Review of Particle Properties.

In Figure 2 we present similar information for the baryon resonances, but concentrate here on the "growth of understanding". That is, the number of known baryons (we include for this figure only those with known J^P) has grown only very slowly with time (dashed line); the real progress has been in the measurement of the properties of those baryons. Therefore we show as the solid line a count of the number of baryonic properties -- mass, width, and branching ratios. Most of these results are from partial-wave analyses.

A history of the values of some of the constants in the Review of Particle Properties is presented in Figs. 3-7. It may be said that one can estimate the age of a high energy physicist by asking him or her the mass of the Λ . If the answer is 1115.44 MeV, he probably was deep into his graduate training in 1965.

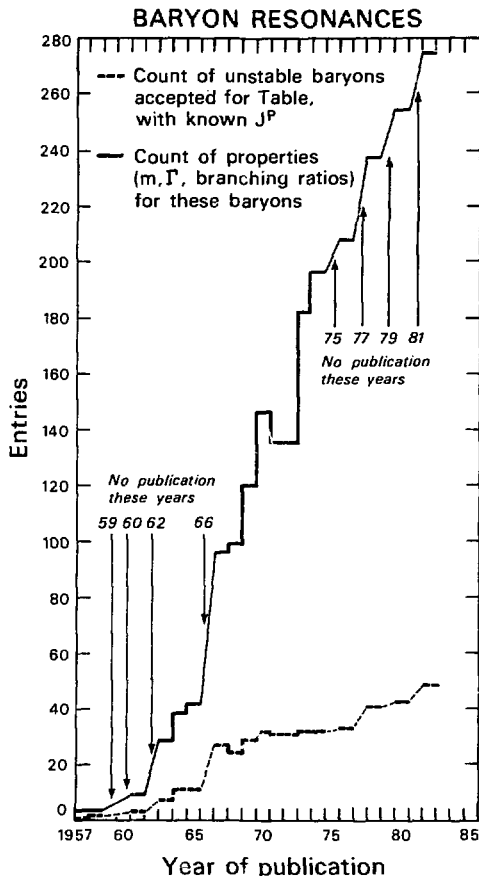


Fig. 2. Total amount of information (mass + width + branching ratios) on baryon resonances listed in the Tables, restricted to those with well-established J^P (solid line). Dashed line shows numbers of such resonances listed. Abscissa shows year of publication of Review of Particle Properties.

A history of this sort has more than whimsical value. We may use it as a guide to develop a "feel" for the reliability of current values. In Fig. 3 we show how the generally accepted values for the speed of light and a couple of other constants have changed with time. The "generally accepted value" is usually an average over several experiments, performed by a compiler (in Fig. 3, the compiler is other

than the Particle Data Group in all cases, although we do quote the compiled results). The abscissa on all these figures is the date of publication of the value shown. Clearly there is a general progression toward better understanding -- at least as measured by the size of the error bars. However, the size of the error bars does not tell the full story, as we can see by the frequency with which the "best" value has changed by more than one standard deviation. Changes in these values can come from several sources: a new experimental measurement, re-evaluation of an old measurement (which can come about if a previously unrecognized source of bias is discovered and corrected, or if a new value for one of the input constants, e.g. the electric charge, is available), or a change in the averaging procedure.

In Fig. 4 we show the history of some masses (including the Λ , for radioactive Λ dating of your colleagues), based on averages which we ourselves performed. These are adapted from those originally presented by Rosenfeld¹ in 1975. The publication date refers to the publication of the Review of Particle Properties.

In Fig. 5 we show the best estimates for the lifetimes of some of the particles stable against strong decay. These and subsequent figures have been compiled since publication of the Rosenfeld article.¹ In Fig. 6 we show the widths of some of the resonances, and in Fig. 7, the values of some of the branching fractions. All values are taken from the Tables. Before 1964, very few branching fractions were listed in the Tables. In all cases, a representative sample is chosen, usually from those with a lot of activity (a limited number of special requests for a more complete set of such figures may be honored, for those seriously interested in the history of the "best" values of physical constants). In each figure, the heavy inner error bar represents the statistical error computed in the averaging procedure, and the thin outer error bars, when present, indicate the increase in the error due to the "scale factor". The scale factor is described in the introductory text, Sec. VII. It represents an attempt to quantify the increase in the uncertainty which is present in the case of experiments which disagree by more than a certain amount. In the case where the error represents an "educated guess," rather than a calculation, the inner error bar is absent.

On the whole, the number of times the values have changed by more than one standard deviation over the years is remarkably few.

Even those branching fractions which involve rare decays and which are therefore presumably difficult to measure (Fig. 7) are, for the most part, within one or two standard deviations in 1978 of their value in any year since 1960. This is in spite of the vast amount of new experimental input, and indicates the general reliability of the results.

Of course, the data points for a given quantity are hardly independent of each other, but those differing by several years frequently have quite different experimental input. The relative lack of change is a comment both on the experiments and on the averaging procedures. We, of course, are responsible only for the averages (except Fig. 3). These averages entail considerable exercise of judgment: there are conflicting experiments, experiments with impossibly small errors, "preliminary" results, and so forth. Statistical procedures will tell us that two experiments do not agree; they do not give a clue as to which (if either) is a good representation of the truth. Major decisions, and their motivations, are usually discussed on a case-by-case basis in the Data Card Listings; general comments may be found in Sec. II of the text and in Rosenfeld¹. Note that, occasionally, the error bars increase from one publication to the next. This is usually the result of decision making by the compiler, e.g., to cease using a particular result, or because of new results in poor agreement with the old results.

We show these figures not only to demonstrate that there is not much change in these averages in the usual case, but also to show that there exist cases with relatively large changes. There is a psychological danger in preparing tables of "right" answers. The old joke about the experimenter who fights the systematics until he or she gets the "right" answer (read "agrees with previous experiments"), and then publishes, contains a germ of truth (presumably, those who compile and average experimental results are also not immune to this disease). A result can disagree with the average of all previous experiments by five standard deviations, and still be right. Hence, perhaps it is of value to show that large changes can (and do) sometimes occur.

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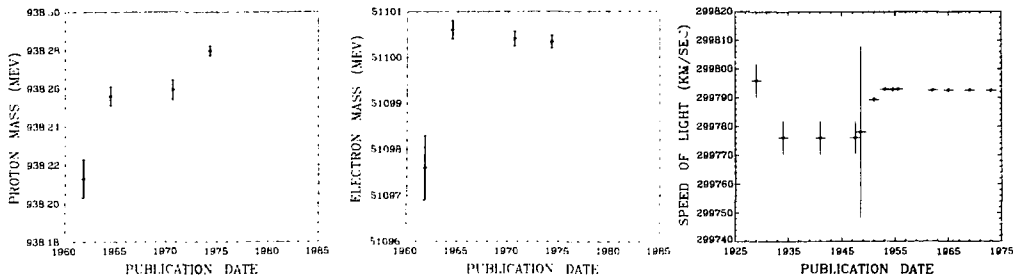


Fig. 3. The "generally accepted values" of the proton mass, the electron mass, and the speed of light, as a function of the publication date of the compilation used (not done by the Particle Data Group). Data for the speed of light plot courtesy of E. R. Cohen, Rockwell International Science Center. See the Stable Particle Data Card Listings for references on proton and electron masses.

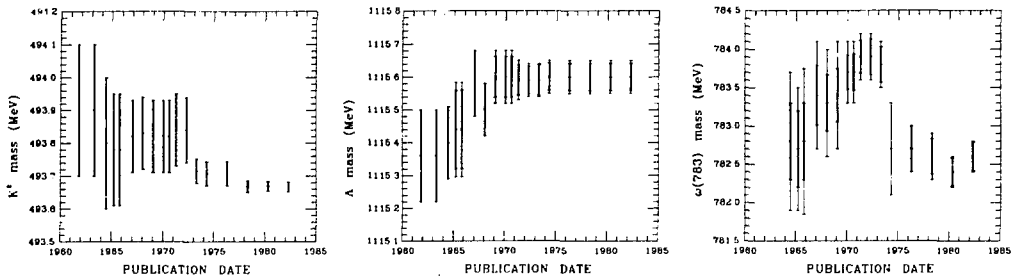


Fig. 4. Particle Data Group averages of the masses of various particles, as a function of date of publication of Review of Particle Properties (Adapted, with permission, from *Annual Review of Nuclear Science*, Volume 25. Copyright 1975 by Annual Reviews, Inc. All rights reserved). Full error bar indicates quoted error; thick-lined portion indicates quoted error with "scale factor" removed (see Sec. VII of introductory text).

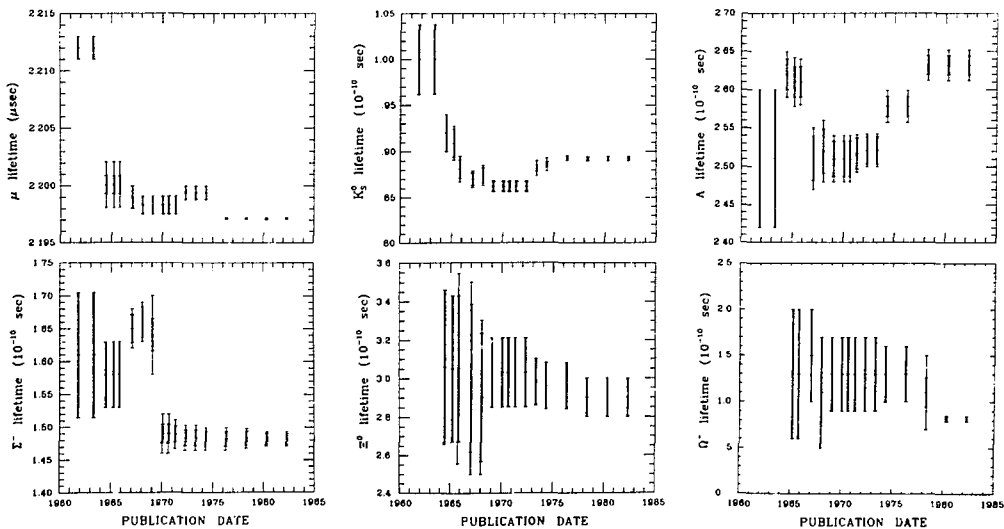


Fig. 5. Particle Data Group averages of the lifetimes of various particles, as a function of publication date of RPP.

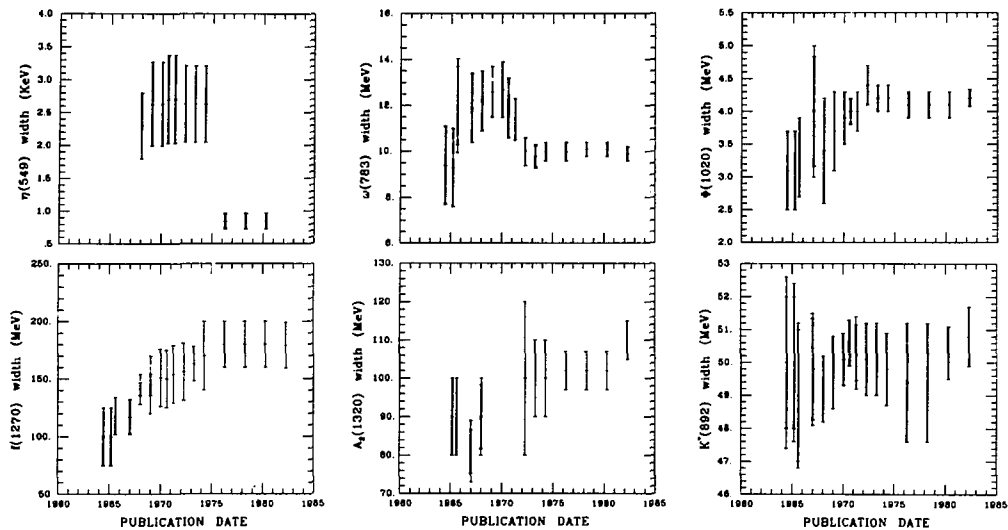


Fig. 6. Particle Data Group averages of the widths of various resonances, as a function of date of publication of RPP. The gap in the A_2 data indicates the years when the A_2 was thought to be split.

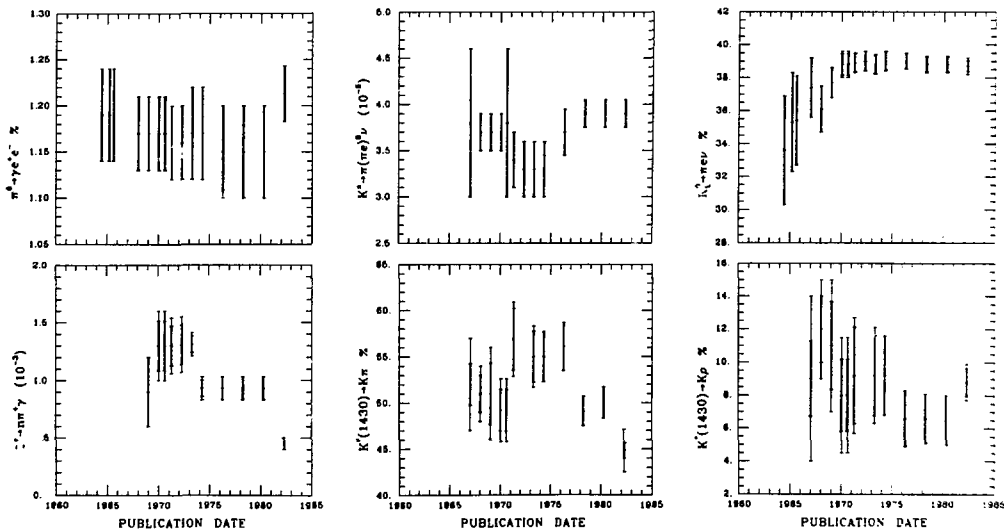


Fig. 7. Particle Data Group averages of various branching fractions, as a function of date of publication of RPP.