

THERMAL DIFFUSIVITY OF MTR-ETR TYPE FUEL PLATES

FINAL REPORT

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Purpose of Study

The purpose of this study was to determine the thermal diffusivity of aluminum clad MTR-ETR type fuel plates* from room temperature to 600°C.

Description of Fuel Plates

The fuel plates were approximately 0.13 cm (0.05 inch) thick with a 0.051 cm (0.020 inch) fuel core clad on both sides with 0.038 cm (0.015 inch) aluminum. A description of the fuel plates and core material are described in Table I.

TABLE I

Sample Number	Thickness		Fuel Core
	cm	in.	
P-1-1048	0.130	(0.051)	52.6 wt% UAl_3 + 0.19 wt% B_4C + X8001 Al
P-2-555	0.130	(0.051)	41.6 wt% U_3O_8 + 0.19 wt% B_4C + X8001 Al
O-4-833	0.128	(0.0505)	63.15 wt% Nb coated UO_2-ZrO_2 + X8001 Al
P-5-576	0.130	(0.051)	57.7 wt% UAl_3 + 0.19 wt% B_4C + X8001 Al
P-1-1047	0.132	(0.052)	52.6 wt% UAl_3 + 0.19 wt% B_4C + X8001 Al
79-11	0.082	(0.032)	50 wt% U + 3 wt% Sn + Al P.A APM-257

Samples for thermal diffusivity measurements were 0.635 cm (0.25 inch) diameter. No changes were made in the fabricated thickness of plates. Some samples were coated with water suspension of graphite to improve transient temperature measurements.

*Measurements were made for and samples were furnished by Phillips Petroleum Company, NRTS, Idaho Falls, Idaho, under related services Purchase Order AEC-74953.

Thermal Diffusivity Apparatus

A single pulse, transient method⁽¹⁾ was used to measure the thermal diffusivity. The technique was similar to the "flash" method originally described by Parker, et al.⁽²⁾ A laser produces the pulse of radiant energy to the specimen and the temperature transient on the back surface of the specimen was measured with liquid N₂ cooled indium antimonide detector. The signal from the detector was displayed on a cathode ray oscilloscope and recorded on Polaroid film.

The pulse was supplied by a ruby crystal [1.27 cm (0.5 inch) diameter and 15.25 cm (6.0 inch) long with TIR end] laser. Measurements were made using pulse energies between 2 and 5 joules and wavelength of 6943° A. The duration of the laser pulse was 0.54 x 10⁻³ seconds.

The wafer specimens were inserted into a lavite holder having a tapered center hole into which the specimens are lightly pressed. The holder was positioned in a tantalum tube which was heated by electrical resistance. Temperatures were measured by Pt versus Pt-Ph and chromel-alumel thermocouples inserted into the specimen holder through a hole in the tantalum tube.

The atmosphere during the measurements was purified argon containing < 1 ppm O₂ and <10 ppm H₂O.

Correction for Heat Losses and Finite Pulse Time

Corrections to thermal diffusivity data for radiation losses and finite pulse time were made using methods of Cape and Lehman⁽³⁾ and Taylor, et al.⁽⁴⁾

Because of the low temperatures, <600°C, no radiation loss corrections were required. However, corrections were made for finite pulse time effects.

When the heat pulse (τ) is of infinitesimal duration, the thermal diffusivity (α) is given by

$$\alpha = 1.37 a^2 / \pi^2 t_{1/2};$$

where $t_{1/2}$ is the time required for the back surface to reach half the maximum temperature, a is the sample thickness and

$$t_c = \frac{a^2}{\pi^2 \alpha}$$

where t_c is the characteristic rise time of the back surface of the sample. When τ is much smaller than t_c , $t_{1/2} / t_c = 1.37$. However, when $\tau \approx t_c$, it is

apparent that $t_{1/2} / t_c > 1.37$ and α should be corrected for the finite pulse time effect.

The shape of the pulse is important. Since the laser pulse more closely approximates a square wave, Taylor's method for finite pulse time correction was used. Estimated error for the thicker specimens (0.130 cm thick) was approximately 5% where $t_{1/2}/t_c = 1.44$.

Results

Data are given in Table II.

An estimate of the accuracy of the thermal diffusivity data was obtained using published thermal conductivity data for wrought aluminum alloys.

The thermal diffusivity is defined as

$$\alpha = k C_p^{-1} \rho^{-1},$$

where k is thermal conductivity, C_p is heat capacity, and ρ is the density.

Assuming nominal values of 0.223 cal/g for C_p , 2.7 g/cm³ for ρ , and 1.2 watt/cm °C for k , at room temperature,

$$\alpha_{\text{cal}} = \frac{1.2 \text{ watt/cm } ^\circ\text{C}}{0.223 \text{ cal/g/}^\circ\text{C} \times 4.186 \text{ watt/cal/sec} \times 2.7 \text{ g/cm}^3}$$

$$\alpha_{\text{cal}} = 0.475 \text{ cm}^2/\text{sec}.$$

Although higher than the experimental values, differences in density and heat capacity could explain these variations. It is interesting to note that the samples cores which contained ceramic, (P-2-555 and P-4-833) exhibited higher thermal diffusivity values than the all metallic cores.

TABLE II

<u>P-1-1048</u> (Coated)		<u>P-1-1048</u> (Noncoated)	
<u>Temperature</u> °C	$\frac{\alpha}{\text{cm}^2/\text{sec}}$	<u>Temperature</u> °C	$\frac{\alpha}{\text{cm}^2/\text{sec}}$
25	0.383	105	0.390
97	0.366	105	0.376
100	0.366	202	0.372
201	0.358	200	0.376
200	0.372	301	0.348
301	0.356	308	0.342
300	0.358	400	0.339
400	0.348	400	0.342
500	0.336	500	0.311
500	0.324	502	0.328

<u>P-2-555</u> (Coated)		<u>O-4-833</u> (Noncoated)	
<u>Temperature</u> °C	$\frac{\alpha}{\text{cm}^2/\text{sec}}$	<u>Temperature</u> °C	$\frac{\alpha}{\text{cm}^2/\text{sec}}$
20	0.462	105	0.469
20	0.487	102	0.484
105	0.449	195	0.442
106	0.449	190	0.445
195	0.437	298	0.433
200	0.442	298	0.433
300	0.417	395	0.390
302	0.425	395	0.406
402	0.417	503	0.440
400	0.417	503	0.402
496	0.404	105	0.495
494	0.383	195	0.490
		305	0.474
		303	0.433
		402	0.433
		500	0.464

TABLE II (Cont)

<u>P-5-576</u> (Coated)		<u>P-1-1047</u> (Coated)	
<u>Temperature</u> <u>°C</u>	<u>α</u> <u>cm²/sec</u>	<u>Temperature</u> <u>°C</u>	<u>α</u> <u>cm²/sec</u>
25	0.331	25	0.311
25	0.336	25	0.335
92	0.306	105	0.298
107	0.311	100	0.294
210	0.278	195	0.290
202	0.282	200	0.298
300	0.274	295	0.278
297	0.269	297	0.288
305	0.278	404	0.265
402	0.274	398	0.268
400	0.247	500	0.265
500	0.251	503	0.250
500	0.254	551	0.233

<u>79-11</u> (Coated)	
<u>Temperature</u> <u>°C</u>	<u>α</u> <u>cm²/sec</u>
25	0.423
25	0.415
100	0.372
102	0.378
195	0.366
200	0.360
300	0.337
295	0.347
400	0.332
400	0.327
500	0.317
500	0.320
560	0.300

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