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ACDOS1: A COMPUTER CODE TO CALCULATE DOSE RATES
FROM NEUTRON ACTIVATION OF NEUTRAL BEAMLINES AND
OTHER FUSION-REACTOR COMPONENTS

Gregory S. Keney
(M.S. thesis)

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**ACDOS1: A Computer Code to Calculate Dose
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By Gregory S. Keney

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**Accelerator and Fusion Research Division
Lawrence Berkeley Laboratory
University of California
Berkeley, California, 94720**

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ABSTRACT

A computer code has been written to calculate neutron induced activation of neutral-beam injector components and the corresponding dose rates as a function of geometry, component composition, and time after shutdown. The code, ACDOS1, was written in Fortran IV to calculate both activity and dose rates for up to 30 target nuclides and 50 neutron groups. Sufficient versatility has also been incorporated into the code to make it applicable to a variety of general activation problems due to neutrons of energy less than 20 MeV.

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Introduction

This project has originated from a need in the fusion community to calculate dose rates from neutron-induced activation of fusion components. In particular, the subject matter of this work revolves around the calculations required to determine the dose rates from neutral-beam injector components as a function of geometry, component composition, and time after shutdown.

The methodology for solving the overall problem is broken down into essentially five major steps. The first is that of determining a mathematical model for the production of neutrons by the neutral-beam injector. In practice, neutron production is not uniform but occurs during a succession of evenly spaced pulses of short time duration. During the pauses, no neutrons are produced. An effective steady neutron source term is generated by utilizing an appropriate duty factor over the time duration of the pulses. The neutron production rate is then calculated on the basis of the energy and current of the neutral beam. The second step required writing and solving the pertinent ordinary differential equation that describes the system undergoing activation and then algebraically modifying the standard solution to make it amenable to the arbitrary injector test schedule and adapting the results to multigroup calculations. With the unity-normalized

average group fluxes assumed to be available from a previous calculation, the third step involves calculating the flux weighted group cross sections by making use of a library containing microscopic neutron cross sections. Prerequisite to this calculation is the determination of a weighting flux for use in averaging cross section data over the appropriate energy groups. This was accomplished by assuming a room-temperature Maxwell-Boltzmann distribution for the thermal region, a $1/E$ function for intermediate energies, and an exponential function for the fast neutron groups. The forms of the equations are all known. Arbitrary constants in the weighting fluxes are determined by integral and/or boundary conditions on the known flux for the particular energy group in question. With the weighting flux a determinable function of energy, the required flux weighted group cross sections can be evaluated. The fourth step involves a calculation of the resultant activities given the neutron production rate, the flux weighted group cross sections, target nuclide mass and type, and the specific times after shutdown at which induced activities are desired. The routine which calculates activity interrogates the cross section library to find the atomic weight of the specific target nuclide and in addition, the product nuclide and its half-life. The fifth and final step is a calculation of the dose rates as a function of time and geometry given the

previously calculated activities. This last routine interrogates a decay library and extracts the necessary gamma ray energy and intensity information. The program incorporates a gamma-flux-to-dose-rate function for several simple geometries. The user specifies the type of geometry to be used in the actual calculations-point, spherical, or cylindrical-on-axis, in addition to the distance from the activated component at which the dose rate is desired.

The result of the five steps is a code called ACDOS1 which performs the activation and dose rate calculations given the required input information. The code can handle from one to fifty neutron groups with up to thirty target nuclides and has an option which allows the user to substitute a neutron source term other than that from a neutral-beam injector, thereby increasing the versatility of the code. In addition, the entire code is written in ANSI FORTRAN IV which should minimize compiler compatibility problems at other user sites. Interested users should submit a written request for additional information. Address requests to: Head, Neutral Beam Development Group, Bldg. 5, Lawrence Berkeley Laboratory, Berkeley, California, 94720.

The end result of this project, then, is a code that is tailored to solve injector activation problems, but contains enough versatility to be useful in solving a variety of general activation problems produced by neutrons of energy less than 20 MeV.

Neutron Production Model

The first major step in solving the overall activation problem is to determine a mathematical model for the production of neutrons by the injector assembly. As previously mentioned, neutron production is not uniform but occurs during an arbitrary succession of evenly spaced injector pulses of short time duration. Neutron production coincides with the injector pulses. Figure 1 depicts schematically how this pulse-pause nature of neutron production might be viewed. INSNPS is the instantaneous number of neutrons produced per second. T1 is the length of time the short pulses and pauses are occurring and T2 is the length of time in which no pulsing-pausing occurs. It is assumed throughout the activation calculations that T1 and T2 do not vary during the operation of the injector.

The neutrons originate from two sources within the injector system. Gaseous deuterium that leaks from the neutralizer and ion-source adsorbs on the surface of the copper target located in the ion-dump and neutral-target sections of the injector. A (d,n) reaction occurs at the far end of the injector as D⁺ impinges upon the deuterium. The reaction is:

In addition, a similar reaction occurs as D^+ impinges upon deuterium adsorbed on the copper target located in the ion-dump section.

The calculation done to determine the instantaneous number of neutrons produced per second is empirical in nature and taken under accelerator test conditions.⁽¹⁾ The expression used is:

$$INSNPS = 8.64 \times 10^4 \frac{\text{neutrons}}{\mu\text{coulomb}} \times A \times 1.0 \times 10^6 \frac{\mu\text{coulomb}}{\text{sec}} \times F \times CF$$

where the variables are defined as below:

INSNPS: the instantaneous number of neutrons produced per second

A: current in amperes

F: fraction of the beam that is monatomic

CF: yield correction factor for voltages different from 150 keV

The yield correction factor for voltages different than 150 keV is taken from the calculation⁽²⁾ reproduced in Figure 2. To facilitate the calculation of INSNPS during execution of the program, a power fit of the form $y=a^x$ was applied to the points read from the curve to determine an analytical expression for CF. Three different fits were used to insure a high coefficient of determination. These expressions are incorporated into subroutine SOURCE which

Figure 1

Figure 2

is the subprogram which calculates INSNPS.

In modeling the beam pulses, INSNPS is multiplied by the fraction of the time the beam is on (the duty factor DF) to obtain the average number of neutrons produced per second during time T1, AVENPS. The short successive pulses are now represented by a single pulse (now referred to as a "test") with the same time length T1 but reduced in magnitude as shown in Figure 3. This type of pulse modeling is used throughout the project in connection with activation calculations. Typically, T1 is several hours. (T2, which represents an outage period as required by maintenance or overnight shutdowns, is also several hours in duration). The quantities, T1 and T2, and N, the total number of tests, are variables and must be specified by the user prior to execution of the program.

Figure 3

Derivation & Modification of the Activation Equation

The balance equation for the activation of target nuclides in the injector system is of the following form:

$$\text{rate of change of a product nuclide} = \text{rate of production} - \text{rate of decay}$$

In differential form, this equation becomes

$$\frac{dB(t)}{dt} = \sigma_A \phi A - B(t) \lambda_B$$

where the variables are defined as below:

$B(t)$: number of atoms of the product nuclide B

λ_B : decay constant for the product nuclide B

σ_A : group activation cross section for target nuclide A

A: number of atoms in the target nuclide A

ϕ : group neutron flux

The solution is:

$$B(t) = \frac{\sigma_A \phi A}{\lambda_B} \left[1 - e^{-\lambda_B t} \right]$$

This is the basic differential equation that is used in

subsequent activation calculations. This form, however, must be algebraically modified to take into account pauses of length T_2 in the injector test schedule as shown in Figure 4. T_1 is the length of the test, T_2 is the length of the pause, $t=0$ is the time at which shutdown occurs, and t_1 and t_2 are arbitrarily selected times after shutdown where an activity is required. N is the total number of tests. It is desired to calculate the amount of an activated nuclide at $t=0$ due to an arbitrary neutron testing history. The method of approach is to calculate the amount at $t=0$ due to each test prior to $t=0$. The total amount at $t=0$ will then be the sum of the contributions from each test. For some time t after shutdown, the amount of nuclide B will be given by:

$$B(t) = \frac{A_A \phi}{\lambda_B} \left[1 - e^{-\lambda_B \cdot T_1} \right] \cdot \left[\sum_{i=1}^N e^{-\left\{ \left[\begin{smallmatrix} N-i \\ T_1+T_2 \end{smallmatrix} \right] \lambda_B \right\}} \right] e^{-\lambda_B t}$$

This expression, however, represents only 1 group of neutrons. To generalize the relation for up to n neutron groups, the calculation for $B(t)$ above must be done for each group, that is:

Figure 4

$$B(t) = \sum_{K=1}^n \frac{A \sigma_{AK} \phi}{\lambda_B} \left[1 - e^{-\lambda_B \cdot T1} \right] e^{-\lambda_B t}$$

$$\left[\sum_{i=1}^N e^{-\{[N-i][T1+T2]\} \lambda_B} \right]$$

For the general case of the reaction $A(n,b)B$, the activity of B in Bq, t seconds after shutdown will be:

$$B(t) \lambda_B = \underset{B}{\text{Activity of Product Nuclide}} =$$

$$\sum_{K=1}^n A \sigma_{AK} \phi \left[1 - e^{-\lambda_B \cdot T1} \right] e^{-\lambda_B t}$$

$$\left[\sum_{i=1}^N e^{-\{[N-i][T1+T2]\} \lambda_B} \right] \quad (1)$$

Since the amount of a target nuclide (A in this case) is entered into computer memory in terms of kilograms, the number of target atoms "A" is computed by multiplying the input mass by the quantity:

$$\frac{(1000\text{g/kg}) (6.023 \times 10^{23} \text{ atoms/mol})}{\text{atomic weight in g/mol}}$$

Equation 1 is the activation equation which is used in ACDOS1 for calculating total activities. Note that the equation is applicable to situations where the flux is constant as opposed to a "pulsing" nature by the disappearance of the second bracketed term when $N=1$.

Determination of the Weighting Flux & Flux Weighted Group Cross Sections

In order to average cross section data over the appropriate energy groups to acquire flux weighted group cross sections, a weighting flux is required. The flux weighted group cross section $\bar{\sigma}$, can then be evaluated from the expression:

$$\bar{\sigma} = \frac{\int_{\Delta E} \sigma(E) \phi(E) dE}{\int_{\Delta E} \phi(E) dE}$$

The problem was solved by assuming that the thermal region could be represented by a classical Maxwell-Boltzmann distribution with $kT=0.025\text{eV}$, intermediate energies by a $1/E$ function and fast groups by an exponential function, with all equations known except for arbitrary constants.

The decision to use this approach came in part from results of slowing-down theory, and, in addition, from results of the Monte Carlo calculations which produced the data for the sample problem. (The latter involves a neutral beam injector which is surrounded by thick concrete walls.) The group fluxes calculated⁽³⁾ for this case are given in Table 1. To verify the above assumptions, column 3 (neutron/cm².eV.source neutron) was plotted as a function of the arithmetic average of the group boundaries found in column 2 for groups 1 through 19. See Figure 5.

NOTE: The convention will be adopted that Group 1
represents the highest energy group.

GP	Energy Interval(eV)	Neutrons/ cm ² .eV. Source Neutron	Neutrons/ cm ² . Source Neutron
1	2.385E+6 to 2.307E+6	1.013E-10	7.901E-6
2	2.307E+6 to 1.827E+6	2.879E-13	1.382E-7
3	1.827E+6 to 1.108E+6	1.959E-13	8.833E-8
4	1.108E+6 to 5.502E+5	1.896E-13	1.058E-7
5	5.502E+5 to 1.576E+5	3.510E-13	1.378E-7
6	1.576E+5 to 1.111E+5	6.111E-13	2.842E-8
7	1.111E+5 to 5.248E+4	8.172E-13	4.790E-8
8	5.248E+4 to 2.479E+4	1.394E-12	3.860E-8
9	2.479E+4 to 2.188E+4	2.550E-12	7.420E-9
10	2.188E+4 to 1.033E+4	3.049E-12	3.522E-8
11	1.033E+4 to 3.355E+3	6.552E-12	4.570E-8
12	3.355E+3 to 1.234E+3	1.896E-11	4.021E-8
13	1.234E+3 to 5.829E+2	4.673E-11	3.043E-8
14	5.829E+2 to 1.013E+2	1.502E-10	7.234E-8
15	1.013E+2 to 2.902E+1	6.481E-10	4.684E-8
16	2.902E+1 to 1.068E+1	1.642E-9	3.011E-8
17	1.068E+1 to 3.059E+0	5.184E-9	3.951E-8
18	3.059E+0 to 1.125E+0	1.787E-8	3.456E-8
19	1.125E+0 to 4.140E-1	4.327E-8	3.076E-8
20	4.140E-1 to 1.000E-5	2.002E-6	8.288E-7

Table 1

Figure 5

Neutrons/cm².ev
vs
Neutron Energy

Figure 5

Figure 5

Inspection of the curve does suggest a $1/E$ type of behavior for groups 4 through 19 since a quantity that varies as a constant/ E will plot as a straight line on log-log paper and will have a slope of -1 . Calculation of the slope at various points along the line did reveal a slope close to -1 . At higher energies (above $2 \times 10^4 \text{ eV}$) the curve departs somewhat more from $1/E$ behavior and finally reaches a region just below fast energies where the curve is neither $1/E$ nor exponential. Even though this region just below the exponential part of the curve deviates from $1/E$ behavior, it is assumed, for ease of calculations, that it does behave as $1/E$. Since neutrons are "born" with a narrow group of discrete energies in the injector as contrasted with a fission spectrum in a reactor, one would intuitively expect the fast region to appear more as a "spike" than a "spread out" distribution. The graph definitely illustrates this behavior and for this reason the assumption of an exponential form for the fast groups appears justified. As stated previously, the Maxwell-Boltzmann distribution is used for describing the thermal region. At this point, the forms of the equations for describing the weighting flux are known except for arbitrary constants. To evaluate the constants, each analytic expression for the weighting flux is integrated over the appropriate energy group and set equal to the numerical value of the corresponding group quantities in column 4 of Table 1, which are normalized

total group fluxes. For example, for the thermal region, group 20:

$$\int_{1.0 \times 10^{-5}}^{4.14 \times 10^{-1}} C \sqrt{E} e^{-\frac{E}{0.025}} dE = 8.288 \times 10^{-7}$$

therefore,

$$C = 2.36 \times 10^{-4}$$

For an intermediate neutron group, group 17

$$\int_{3.059}^{10.68} C/E dE = 3.951 \times 10^{-8}$$

therefore,

$$C = 3.160 \times 10^{-8}$$

The determination of the constants associated with fast neutron groups, however, is somewhat more involved. Since the assumed form for the fast exponential is

$$\phi(E) = A e^{\frac{E_1 - E_2}{B}} \quad \text{where } E_1 \text{ and } E_2 \text{ are knowns and}$$

E_2 is the highest energy of that group, two equations are required to evaluate the two constants. One equation results from the usual integral condition that

$$(2) \int_{E_1}^{E_2} A e^{\frac{E_1 - E_2}{B}} dE = \text{total flux for that group}$$

while the other condition results from a continuity of flux requirement at the boundary of the two groups. Equation 2 already represents one equation. A second equation resulting from the continuity requirement is:

$$C/E_1 = A e^{\frac{E_1 - E_2}{B}}$$

where C is a known constant from the integral condition placed on the last 1/E group. (Analysis is always done starting with the lowest energy group). See Figure 6. With two equations and two unknowns, the constants can be determined by solving the resultant transcendental equations. If there is another fast group, the same method is applied. One equation results from the required integral condition on that group, and a second equation results from a continuity of flux requirement at the boundary with the preceding group. The result of this method is that the first exponential group is always matched to the last 1/E group, and any additional fast groups are always matched to the preceding fast groups. In the case of only two groups, one thermal and one fast, a match is required at the boundary between the groups to allow evaluation of the two constants in the assumed exponential function following the Maxwellian. For 1/E groups, only one equation is re-

Figure 6

A continuity of flux requirement is used at energy E_1 to obtain another equation required to evaluate the two unknown constants in the assumed exponential function for group 2. This equation is:

$$C/E_1 = A e^{\frac{E_1 - E_2}{B}}$$

See appendix A for the forms and solutions to these transcendental equations.

quired as there is only one constant to be evaluated. Since $1/E$ constants are calculated solely from the integral condition placed on the equation for that group, discontinuities sometimes result in the flux at the group boundary between two adjacent $1/E$ groups or between the thermal and the first $1/E$ group. This is not a serious problem because the discontinuities are not significantly large. The important consideration is that the integral of the weighting flux function over the group be equal to the total flux for that group so that the correct averaging of the flux over that same interval is maintained. Table 2 shows the results of this approach to the weighting flux function determination. The integral of these functions over their proper energy range always results in the numerical value of the total flux for that group as it should. A plot of the data in Table 2 (Groups 1-19) is shown in Figure 7. As previously mentioned, slight flux discontinuities at boundaries between $1/E$ groups exists. Note that the discontinuities are greater in the region just below the two fast groups. This is not surprising since that region deviates somewhat from $1/E$ but was assumed to be $1/E$ in weighting flux calculations. Comparison of Figure 5 with Figure 7 shows a close resemblance as one would definitely want and expect.

During the execution of ACDO51, subroutine WTFLUX will determine the values of the constants associated with these

Table 2
Flux Densities for the Sample Problem

Group	Group Boundaries (eV)	Equation $\phi(E)$
20	1.0×10^{-5} to 0.414	$2.36 \times 10^{-4} \text{Eexp}$ $(-\frac{E}{0.025})$
19	0.414 to 1.125	$3.077 \times 10^{-8} \text{E}^{-1}$
18	1.125 to 3.059	$3.455 \times 10^{-8} \text{E}^{-1}$
17	3.059 to 10.68	$3.160 \times 10^{-8} \text{E}^{-1}$
16	10.68 to 29.02	$3.012 \times 10^{-8} \text{E}^{-1}$
15	29.02 to 101.3	$3.747 \times 10^{-8} \text{E}^{-1}$
14	101.3 to 582.9	$4.134 \times 10^{-8} \text{E}^{-1}$
13	582.9 to 1.234×10^3	$4.057 \times 10^{-8} \text{E}^{-1}$
12	1.234×10^3 to 3.355×10^3	$4.020 \times 10^{-8} \text{E}^{-1}$
11	3.355×10^3 to 1.033×10^4	$4.064 \times 10^{-8} \text{E}^{-1}$
10	1.033×10^4 to 2.188×10^4	$4.693 \times 10^{-8} \text{E}^{-1}$
9	2.188×10^4 to 2.479×10^4	$5.942 \times 10^{-8} \text{E}^{-1}$
8	2.479×10^4 to 5.248×10^4	$5.147 \times 10^{-8} \text{E}^{-1}$
7	5.248×10^4 to 1.111×10^5	$6.387 \times 10^{-8} \text{E}^{-1}$
6	1.111×10^5 to 1.576×10^5	$8.129 \times 10^{-8} \text{E}^{-1}$
5	1.576×10^5 to 5.502×10^5	$1.102 \times 10^{-7} \text{E}^{-1}$
4	5.502×10^5 to 1.108×10^6	$1.511 \times 10^{-7} \text{E}^{-1}$
3	1.108×10^6 to 1.827×10^6	$1.666 \times 10^{-7} \text{E}^{-1}$
2	1.827×10^6 to 2.307×10^6	$6.62 \times 10^{-13} \text{exp}$ $(\text{E}-2.307 \times 10^6 /$ $2.42 \times 10^5)$
1	2.307×10^6 to 2.385×10^6	$7.07 \times 10^{-10} \text{exp}$ $(\text{E}-2.385 \times 10^6 / 11,184)$

Figure 7

Figure 7

assumed functions. The particular forms of the weighting functions chosen depend upon the number of neutron groups in any given activation problem as follows:

<u>Number of Neutron Groups</u>	<u>Group Energy Designation</u>	<u>Form(s) of Weight- ing Function(s)</u>
1	Thermal	Maxwellian
2	1 Thermal 1 Fast	Maxwellian Exponential
3	1 Thermal 1 Intermedi- ate 1 Fast	Maxwellian 1/E Exponential
4 to 50	1 Thermal 1 or more Intermedi- ate 1 or more Fast	Maxwellian 1/E Exponential

For four or more neutron groups, the user must specify the number of fast groups, NF, so that the correct apportionment of 1/E groups can be made.

With the weighting flux function known for each group, the flux weighted group cross section can be calculated from the expression:

$$\bar{\sigma} = \frac{\int_{\Delta E} \sigma(E) \phi(E) dE}{\int_{\Delta E} \phi(E) dE}$$

where $\phi(E)$ is the previously determined weighting flux function. During execution of the subroutine that calculates activation, ACTVAT, the cross section information for the target nuclide in question is read from the ACTLMFE library, which is a subset (without fission cross sections) of the ACTL library.⁽⁴⁾ The data consist of two numbers, an energy in MeV and an associated cross section for that energy in barns. To calculate $\bar{\sigma}$, a subroutine, SMOOTH, first merges the group boundary values into the energy-cross-section pair data and performs a linear interpolation to calculate the cross section at the group boundaries.⁽⁵⁾ The result is two arrays, one with energy values and the other with cross section values. There is always a one to one correspondence. See Figure 8 (less caption). GP(J) and GP(J+1) are arbitrary group boundaries which initially had no cross section value associated with them. The value is obtained by a linear interpolation using the first value on each side of a group boundary assuming that cross sections are linearly interpolable in energy. For two arbitrary cross sections, $\sigma(I)$ and $\sigma(I+1)$ whose values lie within the group boundaries GP(J) and GP(J+1), an analytical expression is found for the line connecting them, using the common two point formula:

$$Y-Y_1 = \frac{Y_2-Y_1}{X_2-X_1} (X-X_1)$$

Figure 8

Straight lines are connected between two cross section points using the two point formula. The resulting equation for the line is used as the analytical expression for $\sigma(E)$ between $E(I)$ and $E(I+1)$. $\sigma(E)$ is used in the evaluation of

$$\Delta \bar{\sigma} = \frac{\int_{E(I)}^{E(I+1)} \sigma(E) \phi(E) dE}{\int_{E(I)}^{E(I+1)} \phi(E) dE}$$

where $\Delta \bar{\sigma}$ is an arbitrary fraction of the total flux weighted group cross section, $\bar{\sigma}$, for the group in question.

In this case y is $\sigma(E)$ and x is energy E in MeV. The equation takes the form:

$$\sigma(E) - \sigma(I) = \frac{\sigma(I+1) - \sigma(I)}{E(I+1) - E(I)} [E - E(I)]$$

so that $\sigma(E)$ for $\sigma(I) \leq E \leq \sigma(I+1)$ is

$$\sigma(E) = \left\{ \frac{\sigma(I+1) - \sigma(I)}{E(I+1) - E(I)} [E - E(I)] + \sigma(I) \right\}.$$

See Figure 8.

Now that an analytical expression exists for $\sigma(E)$ and $\phi(E)$, the product of these two functions is numerically integrated between $E(I)$ and $E(I+1)$ and the value kept in a running sum. This process is continued, point by point, until the upper boundary is reached. The particular form of $\phi(E)$ used depends upon the energy range in which the specific group boundaries lie. For example, in the thermal region where $GP(J) = 1 \times 10^{-5}$ eV and $GP(J+1) = 0.414$ eV, $\phi(E) = 2.360 \times 10^{-4} \sqrt{E} \exp(-E/0.025)$ from Table 1, so that the expression for $\sigma(E) \phi(E)$ is:

$$\left\{ \frac{\sigma(I+1) - \sigma(I)}{E(I+1) - E(I)} [E - E(I)] + \sigma(I) \right\} \left\{ 2.360 \times 10^{-4} \sqrt{E} e^{-\frac{E}{0.025}} \right\}$$

It is this expression that is numerically integrated between each thermal $E(I)$ and $E(I+1)$. When all numerical

integrations are completed for a specific group, the running sum is divided by the total flux for that group. The quotient, $\bar{\sigma}$, is the desired quantity--the flux weighted group cross section for that group. This entire process is repeated until all energy groups have been addressed. Several test cases have been run to acquire results for comparison purposes with another averaging program which originated at LLNL.⁽⁵⁾ The values obtained by this method are remarkably close to those obtained by the LLNL code for fast and intermediate neutron groups. For the thermal region, however, the above mentioned procedure produced better results than the LLNL code.

Two further comments are in order considering the method of averaging cross sections. First, for the case of nuclear reactions which are threshold oriented, a zero is substituted for each group cross section whose energy range is below the threshold energy. Calculations begin only when the energy at which a cross section value was measured equals or exceeds the threshold energy. Typically, that first cross section value is located somewhere within the group boundaries at which the calculations begin. The contribution to that particular group cross section comes solely from evaluation of the pertinent quantities for energies greater than or equal to the first cross section value and less than or equal to the first encountered group boundary. This is shown in Figure

Figure 9

In this example, the reaction has a threshold located between the arbitrary group boundaries, GP(J+3) and GP(J+4). The group cross section for group (J+3) results strictly from the evaluation of the pertinent quantities from the threshold value, T_v , to the first encountered group boundary, GP(J+4). Even if there is only one cross section value located in group (J+3), the calculation proceeds since there is a cross section value associated with GP(J+4) as a result of the initial merging of group boundary energies into the energy-cross-section data. In any case, however, zeros are substituted for group cross sections in group (J) through (J+2) in accordance with the above discussion.

The second comment concerns the calculation of the thermal group cross section. Since no provision is made for interpolating cross section values at the first group boundary, averaging calculations begin at the first encountered cross section value above the energy of the first group boundary as shown in Figure 10. As a result, a very small portion of the Maxwellian (generally below 1.15×10^{-10} eV) is not accounted for. However, this is unimportant since the contribution from this part of the Maxwellian is exceedingly small.

Figure 10

Program to Calculate Activities

The fourth major step in the solution of the problem required writing a standard FORTRAN program to calculate activities given the necessary information. The code is written in a structured format form using one executive program to call up a number of subroutines. Variably dimensioned arrays are used where possible to conserve on memory. The name and purpose of each subroutine is discussed below.

- 1) INPUT: for entering all "variable" input data into computer memory. Error traps are provided to check input data. All data is printed out for user verification.
- 2) ARAYIN: for entering all "array" data. Array data is printed out for user verification.
- 3) SOURCE: given the current, voltage, beam fraction, and duty factor, calculates the instantaneous and average number of source neutrons produced per second during a test. This subroutine is bypassed if the user wishes to use his own source term.
- 4) GPFLUX: uses the calculated values of the average number of source neutrons

produced per second during a test or the user supplied source term and calculates the average flux for each group by taking the product of the unity-normalized total group flux and source neutron term.

- 5) WTFLUX: determines the constants for the assumed weighting flux functions by applying integral and/or group boundary constraints to the pertinent equations. Also prints out the integrals of the weighting flux functions over their appropriate energy intervals for user verification.
- 6) AVRAGE: determines an analytical expression for $\sigma(E)$ over a specific ΔE and numerically integrates the product of $\sigma(E)$ and $\phi(E)$, the weighting function, over the group energy interval. It then takes the sum of the integrated products over the unity normalized total flux for the same group to calculate the flux weighted group cross section, $\bar{\sigma}$.
- 7) SMOOTH: first determines if the energy associated with the first energy-cross-section pair read from the ACTLMFE library

is below the greatest group boundary. If so, the routine then merges the group boundary energies into the energy-cross-section data read from the ACTLMFE library and linearly interpolates to find the values of the cross sections at the group boundaries. If not, the particular reaction is skipped and the next one considered.

- 8) POSITN: for positioning the file marker in the ACTLMFE library.
- 9) POSIT2: same as POSITN.
- 10) ACTVAT: takes the following parameters and calculates activities according to Equation 1.
 - a) number of tests
 - b) length of the tests
 - c) length of the pauses
 - d) number of kilograms of a particular nuclide
 - e) specific times after shutdown
 - f) average group fluxes
 - g) flux weighted group cross sections

The result of each activation calculation is printed out. This allows the user to determine what reaction is the most significant for a particular target nuclide.

Moreover, two running sums are maintained in order to print out the total activity from the activation of a particular target nuclide which is subject to many different reactions and, in addition, the total system activity when all target nuclides have been addressed. If, in the case of threshold reactions, the energy read from the first energy-cross-section pair is greater than the highest group boundary, the reaction is bypassed and interrogation continues. Only one pass through the entire ACTLMFE library is required since the nuclide data is arranged in order of increasing ZA (1000Z+A) as is the user specified target nuclide information. The accuracy of the calculations, of course, depend directly on the reliability of the cross section data and the flux calculations. At the termination of subroutine ACTVAT, all activation results will have been stored in memory for future dose rate calculations.

Program To Calculate Dose Rates

The fifth major and final step required writing a program to calculate dose rates as a function of geometry and time after shutdown. This subprogram, DOSRTE, accomplishes the desired calculation by performing essentially four major tasks:

- 1) interrogation of the decay library, LEVDEC, which is a subset of ENSL⁽⁶⁾, for the energy and multiplicity associated with each applicable gamma ray produced by a specific nuclide,
- 2) determination of an effective particle flux that yields a unit absorbed dose in soft tissue,
- 3) calculation of the dose rate using the previously calculated activities, user specified geometry, applicable multiplicities, and appropriate effective particle fluxes, and
- 4) interrogation of the LEVDEC library, on a second pass, for radioactive daughters.

Provisions are made in the program for keeping two running sums of dose rate values, one for a particular target nuclide and one for the system as a whole. The first allows the user to determine which of the particular target nuclides is contributing to the highest dose rate values. The second, of course, gives the dose rate for the entire system. After all parent and daughter

nuclides have been addressed, the results are printed out and the ACDOS1 object program terminates execution.

Prior to performing the calculations, parent ID numbers are rearranged in order of increasing ZA, 1000 Z+A, to match the format of nuclides listed in the library. The associated target ID numbers and parent activities and half-lives are also rearranged at the same time to maintain the correct association of data. All rearrangement occurs automatically during the execution of DOSRTE and is a prerequisite for efficient interrogation of the library.

Since the same nuclides (parents or daughters) may result from two different target nuclides by different reactions, a temporary holding array is utilized to store the decay information when two or more successive parents have the same ID numbers. This allows repeated interrogation of the information stored in the holding array when the routine finds parents with the same ID number. The use of the array in conjunction with the initial rearrangement of parent ID numbers makes it possible to address all parents on one pass through the LEVDEC library. Dose rates due to each parent are calculated and then subsequently stored in memory with the correct target affiliation prior to addressing any potential radioactive daughters.

After calculating dose rates for applicable parents, radioactive daughters are considered. Since the

identity of any daughters is not known in advance, a list of such potential candidates (read from the LEVDEC library) is compiled and stored in memory with the proper parent ID number and parent activity affiliation. After all parents have been addressed on the first pass through LEVDEC, and the resulting dose rates calculated, a second pass is made to search the library for the previously compiled daughters. Since the same interrogation methodology is employed as with the parents, only one pass through the library is required.

After finding a particular nuclide in the decay library, parent or daughter, the first consideration prior to calculating any dose rates is the determination of an effective particle flux that yields a unit absorbed dose in soft tissue. For each applicable gamma ray energy read from the LEVDEC library, an effective particle flux is calculated. This calculation was accomplished by first fitting 13 curves to the data⁽⁷⁾ shown in Table 3. Twelve of the curves are of the form ax^b and one of the form ae^{bx} . A large number of curves were chosen so as to accurately reproduce the data over such a wide energy range. (As a result, all coefficients of determination arising from the curve fitting process are in excess of 0.99.) Table 4 shows the results of the curve fittings. The appropriate equation for calculating the effective particle flux, FLUXE, is chosen according to the value

<u>E (MeV)</u>	<u>Dose Rate of 2.5 mrad/h corresponds to ϕ_0 QUANTA/cm².S</u>
0.01	956
0.015	2,310
0.020	4,320
0.030	9,980
0.040	17,400
0.050	23,100
0.060	25,200
0.080	23,200
0.100	18,600
0.150	10,800
0.300	5,020
0.400	3,660
0.500	2,920
0.600	2,440
0.800	1,880
1.00	1,550
1.25	1,330
1.50	1,130
2.0	912
3.0	686
4.0	559
5.0	480
6.0	420
8.0	339
10.0	284

Table 3

<u>Energy Range (MeV)</u>	<u>Particle Flux Equation</u>
0.010 to 0.020	FLUXE=21495330.64*E** 2.1759447
0.020 to 0.040	FLUXE=11460233.22*E** 2.0136022
0.040 to 0.050	FLUXE=1036918.975*E** 1.2698653
0.050 to 0.060	FLUXE=96497.38032*E** 0.4772399
0.060 to 0.080	FLUXE=11225.12221*E** (-0.2874410)
0.080 to 0.150	FLUXE=55482.78326*EXP (E*(-10.9132069))
0.150 to 0.400	FLUXE=1330.569205*E** (-1.1036085)
0.400 to 0.600	FLUXE=1462.198966*E** (-1.0004513)
0.600 to 1.000	FLUXE=1546.997798*E** (-0.8891058)
1.000 to 1.500	FLUXE=1549.091320*E** (-0.7797676)
1.500 to 3.000	FLUXE=1508.205757*E** (-0.7188784)
3.000 to 6.000	FLUXE=1488.362190*E** (-0.7051622)
6.000 to 10.000	FLUXE=1656.852874*E** (-0.7649757)

Table 4

of E read from the decay library for a particular gamma ray.

The second consideration is that of geometry. Three choices are possible. The first is a point source approximation given by the following equation for the dose rate:

$$\text{DOSRAT} = \frac{S_0}{4\pi D^2} \quad .$$

S_0 is the source strength divided by the flux-per-unit-dose-rate as given by the formula

$$S_0 = \frac{\text{MULT} \times 2.5 \times \text{ACT}}{\text{FLUXE}}$$

where MULT is the multiplicity of the gamma ray, and FLUXE is the effective photon flux corresponding to a dose rate of 2.5 mrem/h in soft tissue. D is the radial distance in cm to the point where the dose rate is desired. The second choice is that of cylindrical-on-axis geometry subject to the condition that the cylinder be non-absorbing. The equation⁹ that is used to calculate the dose rate as a function of the distance from one end is:

$$\text{DOSRAT} = \frac{S_v}{4} \left\{ (HT+D) \left[\ln \left(1 + \frac{R^2}{(HT+D)^2} \right) + \right. \right. \\ \left. \frac{2R}{HT+D} \text{ARCTAN} \frac{HT+D}{R} \right] - D \left[\ln \left(1 + \frac{R^2}{D^2} + \right. \right. \\ \left. \left. \frac{2R}{D} \text{ARCTAN} \frac{D}{R} \right) \right] \left. \right\}$$

where S_v is the ratio of S_o to the volume in cm^3 , D is the distance in cm, and R and HT are the radius and height in cm. The third and final choice, is that of spherical geometry also subject to the condition that the sphere be non-absorbing. The equation⁹ is:

$$\text{DOSRAT} = \frac{S_v}{4d} \left[2Rd - (d^2 - R^2) \ln \left(\frac{R+d}{d-R} \right) \right]$$

where S_v is defined as above, R is the radius in cm, and $d=D+R$ is the distance from the center of the sphere in cm.

In addition to the constraint that the cylinder and sphere be non-absorbing, the following assumptions apply to all dose rate calculations:

- 1) any potential radioactive daughters result only from β - decay,
- 2) the daughter decays from the ground state,
- 3) the daughter radiations appear coincidentally with those of the parents, i.e. have zero half lives,
- 4) the daughter of the daughter is stable, and
- 5) dose rates due to parents and daughters with half-lives less than one second can be neglected.

After all parents and radioactive daughters have been located and their contributions determined, the results are printed out starting with the lowest target ID number and progressing upward in ascending order. For each target nuclide, the dose rate is printed out for each reaction that occurs. and then, for the target as a whole. This is, of course, just the sum of the dose rates due to all of the reactions generated by that target nuclide. Finally, after all target nuclides are addressed, the dose rates for the entire system are printed out and the ACDOS1 object program terminates execution.

Code Description

1. ACTLMFE - Description
2. LEVDEC - Description
3. Code Input Description for
ACDOS1 Data Deck
4. Code Output Description for
ACDOS1
5. Program Flowchart
6. Subroutine Level Chart
7. Description of ACDOS1 Tape
Structure

ACTLMFE Description

The ACTLMFE library contains 20,177 lines. The format of the file is:

<u>Record</u>	<u>Column</u>	<u>Format</u>
1	1-6	ZA (1000Z+A)
	7-13	Ignore
	14-24	Target Mass (amu)
	25-35	Ignore
	36-46	Level of Target (MeV)
	47	Ignore
	48-58	Target Half-Life(s)
2	1-2	Reaction ID Number
	3-8	Ignore
	9-20	Q Value (MeV)
	21-32	ZA of Product
	33-44	Level of Product (MeV)
	45-56	Product Half-Life(s)
3	1-3	Number of Energy-Cross-Section Pairs (NP)
3 et seq	1-66	(EN(K), CS(K), K=1, NP) where EN is energy in MeV and CS is cross section in barns
After the last line of EN, CS pairs	72	Reaction Separator Sentinel

71X, 11

The above pattern is repeated for each reaction. The last reaction in ACTLMFE is for U^{240} , (N, γ).

LEVDEC Description

The LEVDEC library contains 62,221 lines. The format for the first record of each set is:

<u>Parameter</u>	<u>Format</u>
ZA(1000Z+A)	I6
Level (MeV)	E11.4
Parity	F4.1
Spin	F5.1
Half-Life(s)	E11.4
Number of Decay Modes to Follow (NDMODE)	I3

2nd et seq. records to NDMODE

<u>Parameter</u>	<u>Format</u>
Blank	38X
Mode of Decay	I2
ZA of the Daughter	I7
Level of Daughter	E11.4
Probability of Decay to that Level	E12.4

Two comments are in order concerning the LEVDEC library. First, of the twelve possible modes of decay shown on Table 5, only three are considered in dose rate calculations, γ , β^+ , and β^- .

<u>Decay Identifier</u>	<u>Mode of Decay</u>
1	neutron
2	proton
3	deuteron
4	triton
5	He ³
6	α
7	γ
8	β^+
9	β^-
10	EC
18	unresolved EC+ β^+
99	no decay- stable ground state

Table 5

During the interrogation of the LEVDEC library for a particular parent or radioactive daughter nuclide, subroutine DOSRTE will check the decay identifier. If the identifier is a 7, 8, or 9, the reaction is considered. If not, the reaction is skipped and interrogation continues. Secondly, when the identifier is 8, (which is β^+ decay) two annihilation γ 's of 0.511 MeV each will automatically be included in the dose rate calculations with the appropriate multiplicities.

Code Input Description for ACDOS1 Data Deck

<u>Card No.</u>	<u>Variable Description</u>	<u>Columns</u>	<u>Format</u>
1	Any information to be printed out regarding problem	1-80	10(A8)
2	A, V, F, DF, T1, T2, N: current (amperes), voltage (KV), fraction of beam that is monatomic, duty factor, length of test (h), length of pause (h), number of tests. NOTE: If the user specifies his own neutron source term, A, V, F, and DF will not be needed. In this case:	1-75	(6(E10.3, 2X), I3)
2	T1, T2, N: length of test (h), length of pause (h), number of tests.	1-27	(2(E10.3, 2X), I3)
3	NOEGPS-number of energy groups	1-2	I2
4	NOPAS-number of points after shutdown	1-2	I2
5	R, D, HT, IGEO: radius of sphere or cylinder (m), distance from point source or surface of sphere or cylinder (m), height of cylinder (m), geometry designation: 1-point source, 2-spherical, 3-cylindrical-on-axis.	1-37	(3(E10.3, 2X, I1)
6	NONUCL-number of nuclides, 30 maximum	1-2	I2
7	NF-number of fast neutron groups; use 0 if there are none	1-2	I2

8	(BFLUX(I), I=1, NOEGPS: unit normal- ized input fluxes	1-72	(6(E11.4, 1X))
9	(STAFS (I), I=1, NOPAS): specific time after shut-down (h)	1-72	(6(E10.3, 2X))
10 to as many as needed	ZNAME (I), IDNO (I), MASS (I): nuclide name, ID number-1000Z+A, mass of nuclide (kg). Maxi- mum of 30 target nuclides per run	1-25	(A8, 1X, I5, 1X, E10.3)
from above to as many as needed	(GP (I), I=1, NG): group boundaries in MeV in ascending order	1-72	(6(E10.3, 2X))

Code Output Description for ACDOS1

Prior to the execution of activation calculations, all input information is printed out for user verification under the following heading: "THE FOLLOWING DATA HAS BEEN ENTERED INTO MEMORY." This information is:

AMPERES	KILO-VOLTS	BEAM FRACTION	DUTY FACTOR
T1	T2	N	
NUMBER OF ENERGY GROUPS-NOEGPS			
NUMBER OF TARGET NUCLIDES-NONUCL			
NUMBER OF POINTS AFTER SHUTDOWN-NOPAS			
RADIUS R OF SPHERE OR CYLINDER (M)			
HEIGHT H OF CYLINDER (M)			
DISTANCE D FROM POINT SOURCE OR SURFACE OF SPHERE OR CYLINDER (M)			
NUMBER OF GROUP BOUNDARIES-NG			
NUMBER OF FAST NEUTRON GROUPS-NF			
SPECIFIC TIMES AFTER SHUTDOWN (H)			
UNIT NORMALIZED FLUXES			
GROUP BOUNDARIES (MEV)			
NAME	ID-NUMBER	MASS (KG)	

Quantities calculated by ACDOS1 are printed out under the heading: "THE FOLLOWING HAS BEEN CALCULATED BY ACDOS1." The quantities are:

INSTANTANEOUS NUMBER OF NEUTRONS PRODUCED PER SECOND
AVERAGE NUMBER OF NEUTRONS PRODUCED PER SECOND

INTEGRALS OF THE FITTING FUNCTION OVER THE GROUP INTERVALS

Calculated activities are printed out under the general heading: "THE FOLLOWING OUTPUT DATA ARE CALCULATED ACTIVITIES IN BQ." Three sub-headings are used to display the activation results. The first is: "TARGET PRODUCT ACT (T1) ACT(T2)...ACT(T12)." The second is: "TARGET ACT (T1) ACT(T2)...ACT(T12)." The third is: "SYSTEM ACT (T1) ACT(T2)...ACT(T12)." Above each of these three sub-headings, the times after shutdown in hours, are printed out to expedite review of the activation results. Calculated dose rates are printed out in exactly the same format as activation results. They appear under the general heading: "THE FOLLOWING OUTPUT DATA ARE CALCULATED DOSE RATES IN MRAD/H." The corresponding three sub-headings are:

- 1) "TARGET PRODUCT DSR(T1) DSR(T2)...DSR(T12)"
- 2) "TARGET DSR(T1) DSR(T2) DSR(T3)...DSR(T12)"
- 3) "SYSTEM DSR(T1) DSR(T2) DSR(T3)...DSR(T12)"

General Program Flow Chart*

*Not in detail--only to show the general flow of logic.

Subroutine Level Chart

Description of the ACDOS1 Tape Structure

Tape ACDOS1 has 5 files:

<u>File #</u>	<u>Name</u>	<u>Physical Records</u>	<u>Logical Records</u>
1	ACDOS1 Source Listing (ANSI FORTRAN IV)		.
2	ACTLMFE Library	449	20,177
3	LEVDEC 1	551	24,772
4	LEVDEC 2	544	24,476
5	LEVDEC 3	289	12,973

Logical records are 80 column card images.

Tape Characteristics

Type: 9 track, 1600 bpi

Blocking Factor: 45

Format: EBSCDIC

Sample Problem

Determine the dose rates resulting from activation of the copper target in the target section of the injector.

Composition-820lbs Cu

Chosen Geometry-Spherical

From the "Chart of the Nuclides", 9th Edition, 1966

<u>Element</u>	<u>lbs</u>	<u>Isotope</u>	<u>Wt%</u>	<u># of kg's</u>
Cu	820	Cu ⁶³	69.1	257.55
		Cu ⁶⁵	30.9	115.17

Volume of a sphere = $\frac{4}{3}\pi R^3$. Therefore, the volume of

copper is: $\frac{3.727 \times 10^5 \text{ gm}}{8.96 \text{ gm/cm}^3} = 4.160 \times 10^4 \text{ cm}^3$

$$R = \sqrt[3]{\frac{3 \cdot \text{VOL}}{4\pi}}$$

$$R = \sqrt[3]{\frac{3 \cdot 4.160 \times 10^4 \text{ cm}^3}{4\pi}}$$

$$R = \sqrt[3]{9931.27 \text{ cm}^3}$$

$$R = 0.215\text{m}$$

The following sample problem parameters apply:

A	current (amperes)	65
V	voltage (kilovolts)	170
F	beam fraction	1.0
DF	duty factor	0.10
T1	length of test (h)	8
T2	length of pause (h)	16
N	number of tests	7
NOEGPS	number of energy groups	20
NOPAS	number of points after shutdown	12
R	radius of sphere	0.215
D	distance from surface of sphere (m)	2
IGEOM	geometry designator (sphere)	2
NONUCL	number of nuclides	2
NF	number of fast neutron groups	2

GP(I), I=1,
NG (MeV)

BFLUX (I), I=1, NOEGPS
(neutrons/cm².sec/Source
Neutron)

1X10 ⁻¹¹	8.288X10 ⁻⁷
4.14X10 ⁻⁷	3.076X10 ⁻⁸
1.125X10 ⁻⁶	3.456X10 ⁻⁸
3.059X10 ⁻⁶	3.951X10 ⁻⁸
1.068X10 ⁻⁵	3.011X10 ⁻⁸
2.902X10 ⁻⁵	4.684X10 ⁻⁸
1.013X10 ⁻⁴	7.234X10 ⁻⁸
5.829X10 ⁻⁴	3.043X10 ⁻⁸
1.234X10 ⁻³	4.021X10 ⁻⁸
3.355X10 ⁻³	4.570X10 ⁻⁸
1.033X10 ⁻²	3.522X10 ⁻⁸
2.188X10 ⁻²	7.420X10 ⁻⁹
2.479X10 ⁻²	3.860X10 ⁻⁸
5.248X10 ⁻²	4.790X10 ⁻⁸
1.111X10 ⁻¹	2.842X10 ⁻⁸
1.576X10 ⁻¹	1.373X10 ⁻⁷
5.502X10 ⁻¹	1.058X10 ⁻⁷
1.108	8.333X10 ⁻⁸
1.827	1.382X10 ⁻⁷
2.307	7.901X10 ⁻⁶
2.385	

Results of Sample Problem

Inspection of the printout shows a dose rate of 2.67 mrad/h at time $t=0$ after shutdown. The distance from the surface of the sphere is 2m. Of this system dose rate, the majority of the contribution is coming from Copper-64 which was produced by the (n,γ) reaction on Copper-63. Due to the 12.9 hour half-life of Copper-64, one would need to exercise some caution for a period of several hours after shutdown if maintenance were to be performed near the beam dumps. (At 10 cm from the spherically modeled source, the corresponding dose rate is about 129 mrad/h due to Copper-64 alone.) It is interesting to note that the total activities of Copper-64 and Copper-66 only differ by a Factor 3, while the corresponding dose rates differ by a Factor of 7. This difference, of course, is due to the energy and decay probability considerations specific to the two different nuclides.

Although an activity was calculated for Nickel-63, there was no contribution to the dose rate. The long-lived state of Nickel-63 decays to Copper-63 by β^- emission. Since Copper-63 is formed in the ground state, no dose rate results.

Project Conclusions

The objective of this project was to determine the dose rates associated with neutron activation of neutral beam injector components. This objective has been accomplished by creating a Fortran IV program to calculate dose rates as a function of geometry, component composition and amount, and time after shutdown.

The work was successfully completed by performing five major tasks:

- 1) determination of a suitable mathematical model for the production of neutrons by the injector,
- 2) modification of the differential equation that describes the system undergoing activation (to reflect the pulse-pause nature of neutron production) and adapting the results to multigroup calculations,
- 3) development of a program to average microscopic cross-section data over the appropriate energy range (to calculate flux weighted group cross sections),
- 4) creation of a program to calculate activation given the injector testing history, pertinent beam parameters, component composition and amount, and desired times after shutdown, and finally,
- 5) development of a program to calculate dose rates given the previously calculated activities, user

specified geometry, and distance from the activated component.

The five tasks manifest themselves through 14 subroutine subprograms and one executive calling program. Together, they constitute the program, ACDOS1, which performs the desired dose rate calculations given the necessary information.

Although ACDOS1 is a complete program for calculating dose rates due to neutron activation, there are some areas where additional work could improve the code from the standpoint of completeness and accuracy. Coupling of a neutron transport code to the "front end" of the program would eliminate the need for the user to supply the input fluxes. With ACDOS1, the user must obtain the input fluxes from a separate calculation. This addition would make the program more self-contained computationally and could also reduce the input data required to run the program. Also, other geometry options could be added to increase the flexibility in modeling unusual source distributions that might arise in non-injector applications. Addition of these options could be made by incorporating the appropriate equation into subroutine CALC which is the subprogram that calculates the dose rates. (Any additional arguments would have to be passed in the calling statement.)

To produce more accurate dose rate calculations, three immediate improvements could be made to the program:

1) a more realistic treatment of radioactive daughters, 2) inclusion of self absorption in the geometrically modeled sources, and 3) a within-cavity geometry option. Currently, radioactive daughters are assumed to have a half life short compared to that of the parent. Although computationally expedient, this is not always realistic. The improvement would involve incorporating a more complex activation equation into the program; and in addition, some major changes in software logic. Inclusion of self absorption in the dose rate calculations would remove a major conservatism inherent in the present treatment of the calculations. Although modification of the dose rate equation would be a relatively simple matter, there would have to be some provision made for generating or acquiring γ ray attenuation coefficients since they vary strongly with composition and energy. An additional library containing such data would probably be the best solution. Another routine would also be necessary for interpolating γ attenuation coefficients between data points. Finally, a within-cavity geometry option could be utilized to include dose rates from activation of the concrete walls and/or floor surrounding the injector. Presently, there is no way to include any contribution to the dose rate from the walls of the injector housing.

At this time a second version of ACDOS1, ACDOS2, is being written and will incorporate many of the above men-

tioned improvements. These modifications will remove many of the approximations used in ACDOS1, and should yield more accurate dose rate calculations.

A final comment involves a paper¹⁰ that was written at ORNL concerning the calculation of dose rates in the Tokamak Fusion Test Reactor Test Cell. Basically, it is a comparison of resultant dose rates in the test cell with and without a neutral beam injector present--the objective being to determine the contribution to the test cell dose rates due to the presence of the injector and its penetration. Although no code details are presented in the article, a comparison of the similarities, differences, and general methodologies used would offer a perspective from which to evaluate or "bench mark" the ACDOS1 code work since both projects have similar objectives but different origins.

A tabular comparison of the two projects is shown in Table 6.

Table 6

<u>ORNL Project</u>	<u>ACDOS1 Project</u>
Photon transport calculations are carried out using Monte Carlo methods.	No photon transport calculations are done. Photon flux at a dose point is determined from standard geometric equations. no self attenuation or buildup is taken into account in geometrically modeled sources.

ORNL Project

Concrete walls and ceiling of the test cell were considered.

Uses cylindrical geometry to calculate spatially averaged dose rates.

Considers contributions from parents, daughters, and grand-daughters.

Allows for identification of those nuclides that are contributing to significant dose rates.

Arbitrary pulse sequences are allowed.

Calculates dose rates for a variety of times after shutdown.

Neglects burnout of activated nuclei.

Activation photons produced inside the igloo which stream into the injector through the penetration were considered.

Appears to be specific to neutral beam activation and dose rate calculations.

ACDOS1 Project

Doesn't consider effects due to activation of the concrete. Considers only activation of injector components inside the injector shield.

Choice of three different geometries. Dose rates are point values rather than spatial averages.

Considers contributions from parents and daughters. Grand-daughters are assumed to be stable.

Allows for identification of those nuclides that are contributing to significant dose rates.

Arbitrary pulse sequences are allowed.

Calculates dose rates for a variety of times after shutdown.

Neglects burnout of activated nuclei.

No provisions for including this contribution to dose rates.

Can be used for non-injector activation and dose rate calculations.

Review of Table 6 illuminates the similarities and differences. Some of these differences will most likely be eliminated with the second version of ACDOS1, ACDOS2.

Inclusion of self attenuation and buildup in dose rate calculations along with dose contributions from concrete injector walls will probably be addressed. When these modifications are made to ACDOS1, calculated dose rates will probably be more in line with those resulting from a transport treatment of the γ photons. The ACDOS1 code was never intended to consider activation photons produced in an igloo and which stream into the injector through the penetration. This results from the nature of the design work being done at LBL and the forthcoming upgrade program which will require prolonged testing periods at higher beam currents and duty factors.

In general terms, the ORNL calculation is probably more accurate due to the photon transport approach to dose rate calculation. ACDOS1, however, appears to be more versatile. With upcoming improvements to ACDOS1 as mentioned above, improved accuracy coupled with the code's versatility should make it a useful tool for injector design studies and non-injector applications.

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Appendix A

Solutions to Solving Transcendental Equations

For problems involving two or more neutron groups, transcendental equations will result due to the assumed form of the fast group exponential used for representing the weighting flux function at higher energies.

The exponential is:

$$Ae^{\frac{E-E_1}{B}}$$

where A and B are constants to be determined by applying the appropriate boundary and integral constraints on the pertinent equations, and E_1 is the highest energy in the particular fast group being evaluated--the upper boundary for that group. There are three cases where transcendental equations will result:

- 1) Matching the thermal flux to a fast flux
- 2) Matching a $1/E$ flux to a fast flux, and
- 3) Matching a fast flux to a fast flux

The match is always required at the boundary between the two groups.

Consider Case 1. It is required to match the Maxwellian to the fast group at the boundary between the two groups as shown in Figure 11. $\phi(E)$ is the fast weighting flux function and $GP(J)$, $J=1, 2, 3$ are the group boundaries. In addition, it is required that the integral of $\phi(E)$ over group 2 be equal to the total flux as

Figure 11

discussed on page . Since the constant in the Maxwellian ($C\sqrt{E} \exp(-(E/0.025))$) has already been evaluated, the flux at the boundary (FLUXBD) is known. It is:

$$FLUXBD = CONST(1) * \sqrt{GP(2)} * \exp(-(GP(2)/0.025))$$

where CONST(1) is the constant associated with the Maxwellian and GP is in eV. The first equation is:

$$FLUXBD = A e^{\frac{GP(2) - GP(3)}{B}} \quad (1)$$

in accordance that the flux match at the group boundary GP(2). The second equation is:

$$\int_{GP(2)}^{GP(3)} A e^{\frac{E - GP(3)}{B}} dE = \text{total flux for that group} = BFLUX(NG)$$

where NG=2 is the total number of groups and BFLUX is an array holding the total flux values. Evaluating this integral:

$$\begin{aligned} ABe^{\frac{(E - GP(3))}{B}} \Big|_{GP(2)}^{GP(3)} &= BFLUX(NG) \\ AB - ABe^{\frac{GP(2) - GP(3)}{B}} &= BFLUX(NG) \\ AB(1 - e^{\frac{GP(2) - GP(3)}{B}}) &= BFLUX(NG) \end{aligned} \quad (2)$$

From equation (1), $A = \text{FLUXBD} \cdot e^{-\left(\frac{\text{GP}(2) - \text{GP}(3)}{B}\right)}$

Equating this to equation (2), we get

$$\text{FLUXBD} \cdot e^{-\left(\frac{\text{GP}(2) - \text{GP}(3)}{B}\right)} \cdot B(1 - e^{\left(\frac{\text{GP}(2) - \text{GP}(3)}{B}\right)}) = \text{BFLUX(NG)}$$

As a result,

$$B = \frac{\text{BFLUX(NG)}}{\text{FLUXBD}} \left\{ e^{\frac{\text{GP}(3) - \text{GP}(2)}{B}} - 1 \right\}^{-1} \quad (3)$$

This is the above mentioned transcendental equation. Once B is found, A can be determined by substituting B into equation (2). Subroutine WTFLUX, which evaluates the constants associated with the assumed weighting functions, used the Newton-Raphson Method to solve equation (3). The numerator in equation (3), $\text{BFLUX(NG)}/\text{FLUXBD}$, is always a known quantity and is of the same order of magnitude as the numerator in the exponential, $\text{GP}(3) - \text{GP}(2)$. Therefore, $\text{BFLUX(NG)}/\text{FLUXBD}$ is used as an initial guess to start the iterative routine which solves the equation.

The second case involves matching a 1/E flux to a

fast flux as illustrated in Figure 12. The same type of method is used as before. There are two requirements: (1) continuity of flux at the boundary, $GP(NG)$, and (2) an integral constraint on the fast group. From the first requirement,

$$FLUXBD = \frac{CONST(ICOUNT)}{GP(NG)} = Ae^{\frac{GP(NG) - GP(NG+1)}{B}} \quad (4)$$

where $CONST(ICOUNT)$ is the known constant associated with the previous $1/E$ group. $ICOUNT$ is a neutron group counter. From equation (4), A can be solved in terms of B :

$$A = \frac{CONST(ICOUNT)}{GP(NG)} e^{-\left(\frac{GP(NG) - GP(NG+1)}{B}\right)} \quad (5)$$

From the second requirement:

$$\int_{GP(NG)}^{GP(NG+1)} Ae^{\frac{E - GP(NG+1)}{B}} dE = \text{total flux for the fast group} = BFLUX(NG)$$

$$AB - ABe^{\frac{GP(NG) - GP(NG+1)}{B}} = BFLUX(NG)$$

Figure 12

$$AB(1-e^{\frac{GP(NG)-GP(NG+1)}{B}}) = BFLUX(NG)$$

$$A = \frac{BFLUX(NG)}{B} \left\{ 1 - e^{\frac{GP(NG)-GP(NG+1)}{B}} \right\}^{-1} \quad (6)$$

equating equations (5) and (6)

$$B = \frac{GP(NG) \cdot BFLUX(NG)}{CONST(ICOUNT)} \left\{ e^{\frac{GP(NG+1)-GP(NG)}{B}} - 1 \right\}^{-1} \quad (7)$$

As before, the numerator in this equation is always known and is of the same order of magnitude as the exponential. Therefore, $GP(NG) \cdot BFLUX(NG) / CONST(ICOUNT)$ is used as an initial guess to start the Newton-Raphson iterative routine.

The third case involves matching a fast flux to a fast flux as shown in Figure 13. Again, the same method applies. The requirements yield two equations. From the continuity of flux requirement at $GP(ICOUNT+1)$:

$$FLUXBD = CONST(ICOUNT) = Ae^{\frac{GP(ICP1)-GP(ICP2)}{B}}$$

Figure 13

where for brevity, $ICP1 = ICOUNT+1$ and $ICP2 = ICOUNT+2$. Note that the exponential associated with $CONST(ICOUNT)$ equals 1 when evaluated at $GP(ICP1)$. Therefore,

$$A = CONST(ICOUNT) e^{-\left(\frac{GP(ICP1) - GP(ICP2)}{B}\right)} \quad (8)$$

The integral constraint requires that

$$\int_{GP(ICP2)}^{GP(ICP1)} A e^{\frac{E - GP(ICP2)}{B}} dE = \text{total flux for the fast group} = BFLUX(ICP1)$$

$$ABe^{\frac{E - GP(ICP2)}{B}} \Big|_{GP(ICP1)}^{GP(ICP2)} = BFLUX(ICP1)$$

$$AB(1 - e^{\frac{GP(ICP1) - GP(ICP2)}{B}}) = BFLUX(ICP1)$$

$$A = \frac{BFLUX(ICP1)}{B} \left\{ 1 - e^{\frac{GP(ICP1) - GP(ICP2)}{B}} \right\}^{-1} \quad (9)$$

Equating equation (8) and (9),

$$B = \frac{BFLUX(ICP1)}{CONST(ICOUNT)} \left\{ e^{\frac{GP(ICP2) - GP(ICP1)}{B}} - 1 \right\}^{-1}$$

As before, the numerator (BFLUX(ICP1)/CONST(ICOUNT)) is used as an initial guess. A is then found from equation (8).

In all three cases, the methodology employed to solve the equations is the same. Since the numerators in the pertinent equations are all known quantities when they are required, the user does not have to enter an initial guess. Subroutine WTFLUX performs all of the calculations required to determine the constants associated with the assumed weighting functions. Furthermore, all of the calculated constants are stored in memory for future use in subroutine AVRAVE which calculates the flux weighted group cross sections. The A's are stored in array CONST(I) and the B's in array FCONST(I).

One word of caution is appropriate. It is entirely possible that converging problems may occur during the execution of the routine that solves the transcendental equations.

The Newton-Raphson Method

This method is used to solve the transcendental equations that result when two or more neutron groups are used. The iterative expression is:

$$X_{i+1} = X_i - f(x_i)/f'(x_i)$$

where x_i is the initial guess.

The parameter B (see Appendix A) is always given by the following form:

$$B = \text{KNOWN QUANTITY} \left\{ e^{\frac{GP(N+1) - GP(N)}{B}} - 1 \right\}^{-1}$$

where GP(N+1) is a group boundary higher in energy than GP(N). Therefore,

$$f(B) = \text{KNOWN QUANTITY} \left\{ e^{\frac{GP(N+1) - GP(N)}{B}} - 1 \right\}^{-1}$$

and,

$$f'(B) = \frac{\text{KNOWN QUANTITY} \left(\frac{GP(N+1) - GP(N)}{B} \right) e^{\frac{GP(N+1) - GP(N)}{B}}}{B^2 \left[e^{\frac{GP(N+1) - GP(N)}{B}} - 1 \right]^2} - 1$$

These two expressions are incorporated into subroutine WTFLEX for evaluating the constants associated with the exponential weighting functions. When the difference between B_{i+1} and B is less than 0.1, the iteration terminates.

Appendix C

PROGRAM ACDO51 (INPUT,OUTPUT,TAPE5=INPUT,TAPE7,TAPE6=OUTPUT)

PROGRAM ACDO51 (INPUT,OUTPUT,TAPE5=INPUT,TAPE7,TAPE6=OUTPUT)

ACDO51 IS A STANDARD FORTRAN 4 CODE DEVELOPED FOR CALCULATING DOSE RATE FROM NEUTRON ACTIVATION OF NEUTRAL-BEAM INJECTORS. SUFFICIENT VERSATILITY HAS ALSO BEEN INCORPORATED INTO THE CODE TO MAKE IT APPLICABLE TO A WIDE VARIETY OF GENERAL ACTIVATION PROBLEMS DUE TO NEUTRONS OF ENERGY LESS THAN 20 MEV. FOR FURTHER INFORMATION CONCERNING THIS PROGRAM, SEE CLRL REPORT AC-12244 BY GREGORY S KENEY

VARIABLES AND ARRAYS ARE DEFINED AS BELOW

A	CURRENT IN AMPERES
AC(I)	PARENT ACTIVITY
ACT	ACTIVITY DUE TO ACTIVATION OF A PARTICULAR TARGET NUCLIDE BY A SPECIFIC NEUTRON REACTION
AVEHPS	AVERAGE NUMBER OF NEUTRONS PRODUCED PER SECOND
BFLUX(I)	UNIT NORMALIZED FLUX
CF	NEUTRON YIELD CORRECTION FACTOR FOR VOLTAGES DIFFERENT FROM 150 KV
CONST(I)	CONSTANTS FOR WEIGHTING FUNCTIONS
CS(I)	CROSS SECTION VALUE
CSII(I)	CROSS SECTION VALUE
C	DISTANCE FROM POINT SOURCE OR SURFACE OF SPHERE OR CYLINDER
DF	DLTY-FACTOR
DIFF	DIFFERENCE BETWEEN THE LATEST REFINED GUESS AND THE PREVIOUS GUESS
DLEVL	LEVEL OF THE PARENT OR DAUGHTER
DLEVL(I)	LEVEL OF THE PARENT OR DAUGHTER
DCSDAU	DOSE RATE DUE TO A RADIOACTIVE DAUGHTER
DCSPRC(I)	DOSE RATE DUE TO A PARENT
DOSRAT	CALCULATED DOSE RATE
DCSSUN(I)	DOSE RATE DUE TO A PARTICULAR TARGET NUCLIDE FOR UP TO 12 DIFFERENT TIMES
DPVAR	VALUE OF DEPENDENT VARIABLE IN NUMERICAL INTEGRATIONS
E	ENERGY OF GAMMA RAY
EN(I)	ENERGY VALUE ASSOCIATED WITH A CROSS SECTION
EI(I)	ENERGY VALUE ASSOCIATED WITH A CROSS SECTION
F	BEAM FRACTION
FCNST(I)	EXPONENTIAL CONSTANTS FOR FAST NEUTRON GROUPS
FLUXBD	WEIGHTING FLUX AT A GROUP BOUNDARY
FPRIME	EVALUATED VALUE OF THE DERIVATIVE OF THE TRANSCENDENTAL EQUATION FOR B
FXSUBI	EVALUATED VALUE OF THE TRANSCENDENTAL EQUATION FOR B
GP	GROUP BOUNDARIES
GRFLX(I)	FEW GROUP FLUXES-UP TO 50
GPXSEC(I)	FEW GROUP CROSS SECTIONS-UP TO 50
GUESS1	INITIAL GUESS FOR SOLVING TRANSCENDENTAL EQUATIONS
H(I)	PARENT HALF LIFE
HT	HEIGHT OF CYLINDER(M)

••PROGRAM ACCOS1 (INPUT,CUTPUT,TAPES=INPUT,TAPE7,TAPE6=OUTPUT)••

C	IBASE	MARKS THE BEGINNING OF THE PARENT NUCLIDE LIST
C	IBEGIN	MARKS THE BEGINNING OF THE LIKE PARENTS NUCLIDE LIST
C	ICOUNT	NEUTRON GROUP COUNTER
C	ICP1	ICOUNT+1
C	ICP2	ICOUNT+2
C	ICAUG(I)	CAUGHTER ID NUMBER
C	IDAUGH	DAUGHTER ID READ FROM LEVDEC
C	ICEKMO	DECAY MODE IDENTIFIER
C	IDCKM(I)	DECAY MODE IDENTIFIER
C	IDMCI(I)	5 DIGIT ID NUMBER OF TARGET NUCLIDE
C	IDUMB	DUMMY VARIABLE FOR ADVANCING DISK FILE RECORDS
C	IEND	MARKS THE END OF THE LIKE PARENT NUCLIDES IN THE LIST
C	IETPT	COUNTER FOR INTERMEDIATE NEUTRON GROUPS
C	IGEOM	DESIGNATES PROBLEM GEOMETRY
C	IMODE	DECAY MODE COUNTER
C	INSNPS	INSTANTANEOUS NUMBER OF NEUTRONS PRODUCED PER SECOND
C	IDCDEI(I)	HOLDING ARRAY FOR PRODUCT NUCLIDE ID NUMBERS
C	IPROD	PRODUCT NUCLIDE COUNTER
C	IITCP	MARKS THE END OF THE LIST OF ARRANGED PARENT NUCLIDES
C	ITAPE	READ VARIABLE FOR DISK FILE DATA
C	ITER	ITERATION COUNTER
C	MASS(I)	NUMBER OF KILOGRAMS OF A PARTICULAR NUCLIDE
C	MULT	MULTIPLICITY OF GAMMA RAY
C	N	NUMBER OF TESTS
C	NEPTH	NUMBER OF INTERMEDIATE NEUTRON GROUPS
C	NEI(I)	NUMBER OF E1 C1 PAIRS
C	NP	NUMBER OF FAST NEUTRON GROUPS
C	NG	NUMBER OF GROUP BOUNDARIES-UP TO 51
C	NCEGFS	NUMBER OF ENERGY GROUPS-UP TO 50
C	NGNUCL	NUMBER OF TARGET NUCLIDES-UP TO 30 PER RUN
C	NOPAS	NUMBER OF POINTS AFTER SHUTDOWN-UP TO 12
C	NCTCAC(I)	HOLDING ARRAY FOR PRODUCT NUCLIDE NUMBERS THAT HAVE NOT BEEN RE-ARRANGED IN ASCENDING ORDER
C	IND(I)	ID NUMBER OF DAUGHTER
C	NP	NUMBER OF ENERGY-CROSS-SECTION PAIRS
C	NP(I)	ID NUMBER OF PARENT
C	NTMP1	TEMPORARY STORAGE FOR PARENT ID NUMBERS
C	NTMP2	TEMPORARY STORAGE FOR TARGET NUCLIDE ID NUMBERS
C	NT(I)	ID NUMBER OF TARGET
C	OPTION	VARIABLE THAT DETERMINES WHETHER OR NOT AVENPS WILL BE CALCULATED BY ACCOS OR READ FROM A DATA CARD AS PRE-DETERMINED INPUT
C	FARSUM	PARTIAL SUM IN NUMERICAL INTEGRATIONS
C	PRCB(I)	DECAY PROBABILITY
C	PRCDHL(I,J)	PRODUCT NUCLIDE HALF-LIFE (S)
C	PRCDNU(I,J)	5 DIGIT ID NUMBER OF PRODUCT NUCLIDE
C	R	RADIUS OF HOMOGENEOUS NON-ABSORBING SPHERE OR CYLINDER (M)
C	REALCUM	REAL DUMMY VARIABLE FOR ADVANCING DISK FILE RECORDS
C	STAFS(I)	SPECIFIC TIMES AFTER SHUTDOWN (M)
C	STCPHL(I)	HOLDING ARRAY FOR PRODUCT NUCLIDE ID NUMBERS
C	STCTAC(I,J)	HOLDING ARRAY FOR PRODUCT NUCLIDE ACTIVITIES
C	SUMACT(I)	SUM OF ACTIVITIES PRODUCED FROM A PARTICULAR

```

**PROGRAM ACDCS1 (INPUT,CUTPUT,TAPE5=INPUT,TAPE7,TAPE6=OUTPUT)**
C
C      TARGET NUCLEIDE FOR UP TO 12 DIFFERENT TIMES
C      EXPONENTIAL EFFECT OF TESTS AND PAUSES ON THE
C      SHUTDOWN TIME ACTIVITY OF A SPECIFIC PRODUCT
C      NUCLEIDE
C      SUMDEL      RUNNING SUM OF AVERAGE CROSS SECTIONS
C      SUMINT      RESULT OF NUMERICAL INTEGRATIONS
C      SYSACT(I)    SYSTEM ACTIVITY FOR UP TO 12 DIFFERENT TIMES
C      SYSDESE(I)   SYSTEM DESE RATE FOR UP TO 12 DIFFERENT TIMES
C      TARMAS      ATOMIC WEIGHT OF TARGET NUCLEIDE (AMU)
C      TEMPS       TEMPORARY STORAGE FOR DO LOOP SUMS
C      TIMACT(I,J,K) INDUCED ACTIVITY AS A FUNCTION OF TARGET
C      NUCLEIDE, REACTION, AND TIME
C      T1          LENGTH OF TEST (H)
C      T2          LENGTH OF PAUSE (H)
C      V           VOLTAGE (KV)
C      WIDTH       DELTA E USED IN NUMERICAL INTEGRATIONS
C      XSBIPL      LATEST REFINED GUESS FOR B
C      XSUBI       INDEPENDENT VARIABLE OF THE TRANSCENDENTAL
C      EQUATION FOR B
C      ZINCR       INCREMENT IN ENERGY USED IN NUMERICAL
C      INTEGRATIONS
C      ZNAME       VARIABLE FOR NUCLEIDE NAMES
C      ZNUMER      NUMERATOR IN THE TRANSCENDENTAL EQUATION FOR B
C      ZTEMP3      TEMPORARY STORAGE FOR PARENT ACTIVITIES
C      ZTEMP4      TEMPORARY STORAGE FOR PARENT HALF LIVES
C
C
C
C
C      DIMENSION TIMACT(30,13,12),SYSACT(12),SUMACT(12),SYSDOS(12),PRODNL
C      *(30,13),PRCDNU(30,13),DCSSUM(12),ZNAME(30),STAFS(12),BFLUX(50),IDN
C      *(120),GPXSEC(50),PASS(30),GPFUX(50),GP(51),CCNST(50),FCONST(50),D
C      *OSNUC(13)
C      COMMON A,V,F,DF,T1,T2,A,R,D,MT,CPTIGN,IGEDM
C      REAL INSNPS,MASS
C      CALL INPUT(ICNUCL,NDEGPS,NOPAS,NG,NF)
C      CALL ARRAYIN(BFLUX,ZNAME,IGN,MASS,STAFS,GP,NDEGPS,NONUCL,NOPAS,NG)
C      IF(DFLTION.EQ.1,GOCCGOO) GO TO 1
C      AVEFPS=CPTIGN
C      GO TO 2
C      1 CALL SOURCE(AVEFPS)
C      2 CALL GPFUX(AVEFPS,NDEGPS,GPFUX,BFLUX)
C      CALL WFLUX(BFLUX,GP,CCNST,FCONST,NDEGPS,NF,NG)
C      CALL ACTVAT(TIMACT,SYSACT,SUMACT,GPFUX,PRODNU,PRCDNU,GPXSEC,IDNO,
C      *PASS,STAFS,GP,CCNST,FCONST,BFLUX,NONUCL,NDEGPS,NOPAS,NG,NF)
C      CALL DESPT(TIMACT,PRCDNU,ICNUCL,NGPAS,DCSSUM,SYSDOS,IDNO,STAFS,PR
C      *CDHL,DCOSNUC)
C      STOP
C      END

```

SLURCUTINE INPLT(ACAUCL,NCEGPS,NOPAS,AG,AF)

ARGUMENTS IN-NONE
ARGUMENTS CUT-ACNUCL,NOEGPS,NOPAS,NG,NE

GC TC 95


```

**SUBROUTINE INPUT(INCL,NOEGPS,NOPAS,NG,NF)**
9 WRITE(6,12)
12 FORMAT(1H,11H-PROGRAM 'ACDOS' ABORTED-NUMBER OF ENERGY GROUPS MUST
+BE INTEGER AND LESS THAN OR EQUAL TO 50-RECHECK THIS DATA)
STOP
13 WRITE(6,14)
14 FORMAT(1H,11H-PROGRAM 'ACDOS' ABORTED-NUMBER OF POINTS AFTER SHUT
+DOWN MUST BE INTEGER AND LESS THAN OR EQUAL TO 12-RECHECK THIS DATA)
+A)
STOP
15 WRITE(6,17)
17 FORMAT(1H,15H-PROGRAM 'ACDOS' ABORTED-NUMBER OF NUCLIDES IN SYSTEM
+MUST BE INTEGER AND LESS THAN OR EQUAL TO 30)
STOP
20 WRITE(6,52)
52 FORMAT(1H,18H-PROGRAM 'ACDOS' ABORTED-CURRENT MUST BE GREATER THAN
+0 AMPERES-RECHECK THIS DATA)
STOP
37 WRITE(6,53)
53 FORMAT(1H,19H-PROGRAM 'ACDOS' ABORTED-VOLTAGE MUST BE BETWEEN 38 A
+ND 300 KV INCLUSIVE-RECHECK THIS DATA)
STOP
38 WRITE(6,54)
54 FORMAT(1H,19H-PROGRAM 'ACDOS' ABORTED-BEAM FRACTION MUST BE GREATER
+R THAN 0 AND LESS THAN OR EQUAL TO 1.0-RECHECK THIS DATA)
STOP
39 WRITE(6,55)
55 FORMAT(1H,19H-PROGRAM 'ACDOS' ABORTED-DUTY FACTOR MUST BE GREATER
+THAN 0.0 AND LESS THAN OR EQUAL TO 1.0-RECHECK THIS DATA)
STOP
40 WRITE(6,56)
56 FORMAT(1H,19H-PROGRAM 'ACDOS' ABORTED-PULSE LENGTH MUST BE GREATER
+THAN 0.0-RECHECK THIS DATA)
STOP
41 WRITE(6,57)
57 FORMAT(1H,19H-PROGRAM 'ACDOS' ABORTED-PAUSE LENGTH MUST BE GREATER
+THAN OR EQUAL TO 0.0-RECHECK THIS DATA)
STOP
44 WRITE(6,58)
58 FORMAT(1H,19H-PROGRAM 'ACDOS' ABORTED-NUMBER OF PULSES MUST BE IN
+TEGER AND LESS THAN OR EQUAL TO 999-RECHECK THIS DATA)
STOP
49 WRITE(6,60)
60 FORMAT(1H,19H-PROGRAM 'ACDOS' ABORTED-R MUST BE GREATER THAN OR
+EQUAL TO 0.0 AND 0 GREATER THAN 0.0-RECHECK THIS DATA)
STOP
77 WRITE(6,78)
78 FORMAT(1H,18H-PROGRAM 'ACDOS' ABORTED-HT MUST BE GREATER THAN OR E
+QUAL TO 0.0-RECHECK THIS DATA)
STOP
67 WRITE(6,68)
68 FORMAT(1H,16H-PROGRAM 'ACDOS' ABORTED-ICECH MUST BE 1,2,OR 3-RECHE
+CK THIS DATA)
STOP
19 WRITE(6,81) Z1,Z2,Z3,Z4,Z5,Z6,Z7,Z8,Z9,Z10
81 FORMAT(1H,10(A6))
WRITE(6,20)
20 FORMAT(///,1H,47H-THE FOLLOWING DATA HAS BEEN ENTERED INTO MEMORY.
+///)

```

SUBROUTINE 7APLT(ACAUCL,NCEGPS,NCPAS,NG,NF)

```

WRITE(6,21)
21 FORMAT(1H0,85H#PREFS      KILO-VOLTS      REACTION      OUTVEA
+CTCP      T1      T2      N1
IF(CPTION.EQ.1,CCCCCG) WRITE(6,13) A,V,F,DF,T1,T2,N
13 FORMAT(1H0,E8.3,4X,E8.3,8X,E8.3,11X,E8.3,4X,E8.3,3X,E8.3,3X,I3,///
+1
IF(CPTION.EQ.1,CCCCCG) WRITE(6,65) T1,T2,N
45 FORMAT(1H0,55H#*****
+***.6X,21E8.3,2X,I3,///)
WRITE(6,22) NCEGPS
22 FORMAT(1H ,39HNUMBER OF ENERGY GROUPS=NCEGPS      ,12,/)
WRITE(6,75) NCAUCL
75 FORMAT(1H ,35HNUMBER OF TARGET NUCLIDES=NCAUCL      ,12,/)
WRITE(6,23) NOPAS
23 FORMAT(1H ,39HNUMBER OF POINTS AFTER SHUTDOWN=NOPAS ,12,/)
WRITE(6,25) R
25 FORMAT(1H ,35HRADIALS R OF SPHERE OR CYLINDER (M)      ,E10.3,/)
WRITE(6,79) HT
79 FORMAT(1H ,35HHEIGHT HT OF CYLINDER (M)      ,E10.3,/)
WRITE(6,26) D
26 FORMAT(1H ,68HDISTANCE D FROM POINT SOURCE OR SURFACE OF SPHERE OR
+ CYLINDER (M)      ,E10.3,/)
NG=NCEGPS+1
WRITE(6,72) NG
72 FORMAT(1H ,35HNUMBER OF GROUPS BOUNDARIES=NG      ,12,/)
WRITE(6,76) NF
76 FORMAT(1H ,39HNUMBER OF FAST NEUTRON GROUPS=NF      ,12,/)
IF(CPTION.EQ.1,CCCCCG) WRITE(6,66) OPTCN
66 FORMAT(1H ,46HAVERAGE NUMBER OF NEUTRONS PRODUCED PER SECCND,4X,E1
+ C,4,/)
RETURN
END

```

```

**SUBROUTINE ARAYIN(BFLUX,ZNAME,IDNO,MASS,STAFS,GP,NCEGPS,NCONUCL,NOPAS,NG)
SUBROUTINE ARAYIN(BFLUX,ZNAME,IDNO,MASS,STAFS,GP,NCEGPS,NCONUCL,NOPAS,NG)
C
C THIS PROGRAM IS USED FOR ENTERING 'ARRAY' DATA INTO MEMORY
C
C ARGUMENTS IN-STAFS,NCEGPS,NCONUCL,NOPAS,NG
C ARGUMENTS OUT-BFLUX,ZNAME,IDNO,MASS,GP
C
C
C DIMENSION BFLUX(NCEGPS),ZNAME(NCONUCL),IDNO(NCONUCL),MASS(NCONUCL),STAFS(NOPAS),GPING)
REAL MASS
READ(5,4)(BFLUX(I),I=1,NCEGPS)
READ(5,4)(STAFS(I),I=1,NOPAS)
4 FORMAT(6(E10.3,2X))
DO 7 I=1,NCONUCL
  READ(5,8) ZNAME(I),IDNO(I),MASS(I)
  8 FORMAT(A8,1X,1E,1X,1E10,3)
  7 CONTINUE
  READ(5,4)(GP(I),I=1,NG)
  WRITE(6,24)
  24 FORMAT(/,1H,33,SPECIFIC TIMES AFTER SHUTDOWN (H),/)
  DO 34 I=1,NOPAS
    WRITE(6,35) I,STAFS(I)
  35 FORMAT(1H,4HTIME,1X,12,10X,E10.3)
  34 CONTINUE
  WRITE(6,27)
  27 FORMAT(/,1H,22HUNIT NORMALIZED FLUXES,/)
  DO 32 I=1,NCEGPS
    I=NCEGPS+1
    WRITE(6,28) I,BFLUX(I)
  28 FORMAT(1H,5HGROUP,1X,12,10X,E11.4)
  32 CONTINUE
  WRITE(6,73)
  73 FORMAT(/,1H,22HGROUP BOUNDARIES (MEV),/)
  WRITE(6,74)(I,GPING(I)-1,I=1,NG)
  74 FORMAT(1H,8HBOUNDARY,1X,13,5X,E11.4)
  WRITE(6,29)
  29 FORMAT(1H,/,1H,2X,4HNAME,4X,9HID-NUMBER,3X,8HMASS(KG),/)
  DO 30 I=1,NCONUCL
    WRITE(6,31) ZNAME(I),IDNO(I),MASS(I)
  31 FORMAT(1X,A8,5X,15,5X,E8.3)
  30 CONTINUE
  RETURN
END

```

SUBROUTINE SOURCE(AVENPS)

SUBROUTINE SOURCE(AVENPS)

C

CTHIS PROGRAM CALCULATES THE SOURCE NEUTRON PRODUCTION RATE

C

C ARGUMENTS IN=NONE

C

C ARGUMENTS OUT=AVENPS

C

C

COMMON A,V,F,CF,T1,T2,N,R,D,HT,CPTION,IGECN

REAL INSNPS

IF(V.GE.38.0.AND.V.LE.80.0) GO TO 3

IF(V.GT.80.0.AND.V.LE.180.0) GO TO 4

IF(V.GT.180.0.AND.V.LE.300.0) GO TO 5

3 CF=(3.25917*V**2.83194)/38.5

GO TO 7

4 CF=(1.51725*V**2.46206)/38.5

GO TO 7

5 CF=(9.8466)*V**2.14122/38.5

7 INSNPS=(8.64E+4)*(A)*(1.0E+6)*F*(CF*1.0E-4)

AVENPS=INSNPS*CF

WRITE(6,8)

8 FORMAT(///,1H,43HTHE FOLLOWING HAS BEEN CALCULATED BY ACQCS1,///)

WRITE(6,1) INSNPS

1 FORMAT(///,1H,52HINSTANTANEOUS NUMBER OF NEUTRONS PRODUCED PER SEC

+CAND,4X,E11.4)

WRITE(6,2) AVENPS

2 FORMAT(///,1H,46HAVERAGE NUMBER OF NEUTRONS PRODUCED PER SECOND,4X,

+E11.4)

RETURN

END

```

      **SUBROUTINE GFLLX(AVENPS,NCEGPS,GPFLUX,BFLUX)**
      SUBROUTINE GRFLUX(AVENPS,NCEGPS,GPFLUX,BFLUX)
C-----
C THIS PROGRAM CALCULATES UP TO 50 GROUP FLUXES
C-----
C ARGUMENTS IN-AVENPS,NCEGPS,BFLUX
C ARGUMENTS OUT-GPFLUX
C-----
C
C DIMENSION GPFLUX(NCEGPS),BFLUX(NCEGPS)
C DO 1 I = 1,NCEGPS
C   GPFLUX(I) = BFLUX(I)*AVENPS
C 1 CONTINUE
C RETURN
C END

```

```

**SUBROUTINE WFLUX(BFLUX,GP,CONST,FCONST,NCEGPS,NF,NG)**
SLBRCLTAE WFLUX(BFLUX,GP,CONST,FCONST,NCEGPS,NF,NG)
C
C THIS PROGRAM DETERMINES THE CONSTANTS ASSOCIATED WITH THE ASSUMED
C WEIGHTING FUNCTIONS
C
C ARGUMENTS IN=BFLUX,GP,NCEGPS,NF,NG
C ARGUMENTS OUT=CONST,FCONST
C
C
C DIMENSION BFLUX(NCEGPS),GP(NG),CONST(NCEGPS),FCONST(NCEGPS)
C CO 106 JC=1,NG
C GP(JC)=GP(JC)+1.00000
106 CONTINUE
C CALCULATE THE THERMAL GROUP CONSTANT
ICOUNT=1
SUMINT=0.0
WIDTH=(GP(2)-GP(1))/200.0
ZINCR=WIDTH/2.0
DPVAR=GP(1)+ZINCR
DO 1 I=1,200
FARSUM=SQRT(DPVAR)*EXP(-(DPVAR/0.0251)*WIDTH)
SUMINT=SUMINT+FARSUM
DPVAR=DPVAR+WIDTH
1 CONTINUE
CONST(ICOUNT)=BFLUX(ICOUNT)/SUMINT
IF(NCEGPS.EQ.1) GO TO 1000
IF(NCEGPS.EQ.2) GO TO 100
C CALCULATE THE 1/E GROUP CONSTANTS
ICOUNT=ICOUNT+1
IEPTH=C
2 SUMINT=0.0
WIDTH=(GP(ICOUNT+1)-GP(ICOUNT))/200.0
ZINCR=WIDTH/2.0
DPVAR=GP(ICOUNT)+ZINCR
DO 3 I=1,200
PARSUM=(1.0/DPVAR)*WIDTH
SUMINT=SUMINT+PARSUM
DPVAR=DPVAR+WIDTH
3 CONTINUE
CONST(ICOUNT)=BFLUX(ICOUNT)/SUMINT
IEPTH=IEPTH+1
NEPT=NCEGPS-1-NF
IF(IEPTH.EQ.NEPTH) GO TO 150
ICOUNT=ICOUNT+1
GO TO 2
100 ICOUNT=ICOUNT+1
C THE 1 GROUP FAST SPECTRUM WILL BE MATCHED TO THE MAXWELLIAN AT THE
C GROUP BOUNDARY
FLLXBD=CONST(1)*SQRT(GP(2))*EXP(-(GP(2)/0.0251))
ZNUMBER=GUESS=BFLUX(NCEGPS)/FLLXBD
ITER=1
6 FXSUBI=2*NUMBER/(EXP(GP(3))-GP(2))/GUESS)-1-GUESS
XSUBI=GUESS
FPRIME=2*NUMBER*(GP(3)-GP(2))*EXP(1/2*(GP(3)-GP(2))/GUESS)/(GUESS**2-1
+EXP((GP(3)-GP(2))/GUESS)-1)**2)-1

```

```

-----
**SUERCUT INE WTFLUX(BFLUX,GP,CONST,FCONST,NOEGPS,NF,NG)**
-----
XSLBI=XSLBI-FXSLBI/FPRIME
C1FF=ABS(XSLBI-GUESS1)
IF(DIFF.LT.0.1) GC TC 5
ITER=ITER+1
IF(ITER.GT.100) GO TO 75
GLESSI=XSLBI
GC TC 6
75 WRITE(6,76)
76 FORMAT(1H,114****PROGRAM ABORTED****THE TRANSCENDENTAL EQUATION
+USED TO MATCH THE THERMAL GROUP TO FAST GROUP IS NOT CONVERGING)
WRITE(6,77)
STOP
5 E=XSLBI
A=BFLUX(NCEGPS)/(B*(1-EXP((GP(2)-GP(3))/B)))
CONST(ICOUNT)=A
FCONST(ICOUNT)=B
GC TC 1000
150 IF(NF.GT.1) GC TO 160
C THE 1 GROUP FAST SPECTRUM WILL BE MATCHED TO THE LAST 1/E GROUP
ITER=1
NGP1=NCEGPS+1
ZNUMBER=GUESS1*(BFLUX(NOEGPS)*GP(NOEGPS)/CONST(ICOUNT))
51 FXSLBI=ZNUMBER/(EXP((GP(NGP1)-GP(NOEGPS))/GUESS1)-1)-GUESS1
XSLBI=GUESS1
FPRIME=ZNUMBER*(GP(NGP1)-GP(NOEGPS))*EXP((GP(NGP1)-GP(NOEGPS))/GUESS1)
+SIGMA(GUESS1**2*(EXP((GP(NGP1)-GP(NOEGPS))/GUESS1)-1)**2)-1
XSLBI=XSLBI-FXSLBI/FPRIME
C1FF=ABS(XSLBI-GUESS1)
IF(DIFF.LT.0.1) GC TC 92
ITER=ITER+1
IF(ITER.GT.100) GC TC 25
GLESSI=XSLBI
GC TC 91
52 B=XSLBI
A=BFLUX(NCEGPS)/(B*(1-EXP((GP(NOEGPS)-GP(NGP1))/B)))
CONST(NCEGPS)=A
FCONST(NCEGPS)=B
GO TO 1000
160 ICP1=ICOUNT+1
C THE FIRST OF 2 OR MORE FAST GROUPS WILL BE MATCHED TO THE LAST 1/E
CEFCUP
ICP2=ICOUNT+2
ITER=1
ZNUMBER=GUESS1*(BFLUX(ICP1)*GP(ICP1)/CONST(ICOUNT))
20 FXSLBI=ZNUMBER/(EXP((GP(ICP2)-GP(ICP1))/GUESS1)-1)-GUESS1
XSLBI=GUESS1
FPRIME=ZNUMBER*(GP(ICP2)-GP(ICP1))*EXP((GP(ICP2)-GP(ICP1))/GUESS1)
+SIGMA(GUESS1**2*(EXP((GP(ICP2)-GP(ICP1))/GUESS1)-1)**2)-1
XSLBI=XSLBI-FXSLBI/FPRIME
C1FF=ABS(XSLBI-GUESS1)
IF(DIFF.LT.0.1) GC TC 30
ITER=ITER+1
IF(ITER.GT.100) GO TO 25
GLESSI=XSLBI
GO TO 20
25 WRITE(6,26)
26 FORMAT(1H,116****PROGRAM ABORTED****THE TRANSCENDENTAL EQUATIONS
+WHICH MATCH 1/E TO FAST OR FAST TO FAST GROUPS ARE NOT CONVERGING)

```

```

**SLBRCUTLINE B7FLUX(BFLUX,GP,CCNST,FCCNST,NCEGPS,NF,NG)**

WRITE(6,27)
27 FORMAT(1H,75MCST LIKELY PROBLEM IS UNREALISTIC OR UNUSAL FLUX
+AND GROUP BOUNDARY DATA)
STOP
30 E=XSBIPI
A=BFLUX(ICP1)/(B*(1-EXP((GP(ICP1)-GP(ICP2))/B)))
ICCOUNT=ICCOUNT+1
ICPL=ICCOUNT+1
ICP2=ICCOUNT+2
CCAST(ICCOUNT)=A
FCCAST(ICCOUNT)=B
IF(ICCOUNT.EQ.NCEGPS) GO TO 1000
ZNUMBER=BFLUX(ICP1)/CCNST(ICCOUNT)
ITER=1
GO TO 20
1000 WRITE(6,11)
11 FORMAT(//,1H,63H INTEGRALS OF THE FITTING FUNCTION OVER THE G
+ROUP INTERVAL(S))
WRITE(6,12)
12 FORMAT(1H,74HGROUP LOWER BOUNDARY UPPER BOUNDARY REQUIRED INTEGR
+AL CALCULATED INTEGRAL/)
CCALCULATE THE THERMAL INTEGRAL
ICCOUNT=1
SUMINT=0.0
WIDTH=(GP(21)-GP(1))/200.0
ZINCR=WIDTH/2.0
DPVAR=GP(1)+ZINCR
DO 21 I=1,200
PARSUM=CCNST(I)*SQRT(DPVAR)*EXP(-(DPVAR/0.025))*WIDTH
SUMINT=SUMINT+PARSUM
DPVAR=DPVAR+WIDTH
21 CONTINUE
IF(ICCOUNT.EQ.NCEGPS) GC TC 109
IF(ICCOUNT.EQ.2) GC TC 46
JJ=NG-ICCOUNT
WRITE(6,22) JJ,GP(ICCOUNT),GP(ICCOUNT+1),BFLUX(ICCOUNT),SUMINT
22 FORMAT(1H,2X,12,5X,E10.3,4X,E10.3,5X,E10.3,9X,E10.3)
ICCOUNT=ICCOUNT+1
CCALCULATE THE 1/E INTEGRAL(S)
45 SUMINT=0.0
WIDTH=(GP(ICCOUNT+1)-GP(ICCOUNT))/200.0
ZINCR=WIDTH/2.0
DPVAR=GP(ICCOUNT)+ZINCR
DO 23 I=1,200
PARSUM=(CCNST(ICCOUNT)/DPVAR)*WIDTH
SUMINT=SUMINT+PARSUM
DPVAR=DPVAR+WIDTH
23 CONTINUE
JJ=NG-ICCOUNT
WRITE(6,22) JJ,GP(ICCOUNT),GP(ICCOUNT+1),BFLUX(ICCOUNT),SUMINT
ICCOUNT=ICCOUNT+1
IF(ICCOUNT.LE.NDEPTH+1) GO TO 45
GC TC 16
46 JJ=2
WRITE(6,22) JJ,GP(ICCOUNT),GP(ICCOUNT+1),BFLUX(ICCOUNT),SUMINT
IF(NG.EQ.2) ICCOUNT=ICCOUNT+1
16 SUMINT=C.C
CCALCULATE THE FAST INTEGRAL(S)

```



```

**SUBROUTINE WTELL(BFLUX,GP,CCNST,FCCAST,NCEGPS,NF,NG)**
-----
WIDTH=(GP(ICCUNT+1)-GP(ICOUNT))/200.0
ZINCR=WIDTH/2.0
DPVAR=GP(1CCUNT)+ZINCR
CC 24 I=1,200
PARSUM=CCAST(1CCUNT)*EXP((DPVAR-GP(1CCUNT+1))/FCCNST(1CCUNT))*WIDT
+
SLPIAT=SUMINT+PARSUM
DPVAR=DPVAR+WIDTH
24 CONTINUE
JJ=NG-1CCUNT
WRITE(6,22) JJ,GP(1CCUNT),GP(1CCUNT+1),BFLUX(1CCUNT),SUMINT
1CCUNT=1CCUNT+1
IF(1CCUNT.LE.NCEGPS) GC TC 16
GC TC 15
109 JJ=1
WRITE(6,22) JJ,GP(1CCUNT),GP(1CCUNT+1),BFLUX(1CCUNT),SUMINT
15 DO 105 JC=1,NG
GP(JC)=GP(1CC)/1.000E6
105 CONTINUE
RETURN
END
-----

```

```

**SUBROUTINE ACTVAT(TIMACT,SYSACT,SUMACT,GPFLUX,PRCDNU,PRCDHL,GPXSEC**
SUBROUTINE ACTVAT(TIMACT,SYSACT,SUMACT,GPFLUX,PRCDNU,PRCDHL,GPXSEC
*,IDNO,MASS,STAFS,GP,CONST,ECONST,BFLUX,NONUCL,NCEGPS,NCPAS,NG,NFI
C
C
CTHIS PROGRAM CALCULATES ACTIVITIES DUE TO NEUTRON INDUCED REACTIONS
C
C
C ARGUMENTS IN=GPFLUX,GPXSEC,IDNO,MASS,STAFS,GP,CONST,ECONST,BFLUX
C ARGUMENTS I=-ACNUCL,NCEGPS,NCPAS,NG,NFI
C ARGUMENTS OUT=TIMACT,SYSACT,PRCDNU,PRCDHL
C
C
C DIMENSION TIMACT(NONUCL,13,NCPAS),SYSACT(NCPAS),SUMACT(NCPAS),GPFL
+UX(NCEGPS),PRCDNU(NONUCL,13),PRCDHL(NONUCL,13),GP(NG),GPXSEC(NCEGP
+S),CONST(NCEGPS),ECONST(NCEGPS),BFLUX(NCEGPS),IDNO(NONUCL),MASS(INO
+NUCL),STAFS(NCPAS)
C CPXPCN A,V,F,CF,T1,T2,N,R,D,HT,OPTION,IGEDM
C REAL MASS
C WRITE(6,102)
102 FORMAT(///,1H ,57H-THE FOLLOWING OUTPUT DATA ARE CALCULATED ACTIVIT
+IES IN BQ)
C I=1
C J=0
C INITIALIZE FLAGS AND ARRAYS
C DO 40 P=1,NONUCL
C DC 40 MM=1,13
C PRCDNU(M,MM)=0.0
C PRCDHL(M,MM)=0.0
C 40 CONTINUE
C DC 41 M=1,NONUCL
C CC 41 MM=1,13
C DC 41 MM=1,NCPAS
C TIMACT(M,MM,MM)=0.0
C 41 CONTINUE
C DC 22 JN=1,NCPAS
C SYSACT(JN)=SUMACT(JN)=0.0
C 22 CONTINUE
C KFLAG=LELAG=C
C PRINT HEADINGS
C WRITE(6,37)(STAFS(K),K=1,NCPAS)
37 FORMAT(1H ,4X,7HTIME(H),4X,10E8.3,1X),2X,2(E8.3,2X)
C WRITE(6,1)
1 FORMAT(1H ,125H-TARGET PRODUCT ACT(T1) ACT(T2) ACT(T3) ACT(T4)
+ ACT(T5) ACT(T6) ACT(T7) ACT(T8) ACT(T9) ACT(T10) ACT(T11)
+ACT(T12))
C 24 READ(7,25) ITAFE,TARMAS
C 25 FORMAT(16,7X,E11.4)
C IF(ITAFE.EQ.IDNO(1)) GO TO 7
C IF(ITAFE.GT.IDNO(1)) GO TO 44
C CALL POSITCN
C GO TO 24
C THE ABOVE STATEMENTS SEARCH THE TAPE FOR A MATCHING TARGET ID NUMBER
C 44 WRITE(6,45) IDNO(1)
C 45 FORMAT(///,1H ,31HATTENTION-TARGET NUCLIDE NUMBER,16,88HDOES NOT
+EXIST IN DATA LIBRARY-THEREFORE-NO CONTRIBUTION FROM THIS NUCLIDE

```

```

**SUBROUTINE ACTIVAT(IMPACT, SYSACT, SUMACT, GPFLUX, PRODN, PRODH, GPXSEC)
*IN DESE CALCS,////)
IF(ITAPE.EQ.(IDNO(I+1))) GO TO 46
I=I+1
CALL PCSITCN
CPRINT HEADINGS
WRITE(6,37)(STAFS(K),K=1,NCPAS)
WRITE(6,1)
GO TO 24
46 I=I+1
CPRINT HEADINGS
WRITE(6,37)(STAFS(K),K=1,NCPAS)
WRITE(6,1)
7 J=J+1
READ(7,8) PRODN(I,J), PRODH(I,J)
READ(7,9) NP
59 FORMATT(13)
IF(ITAPE.EQ.(IDNO(I+1))) LFLAG=1
IF(PRODH(I,J).GT.9.1000E+50) GO TO 50
8 FORMAT (20X,E12.4,12X,E12.4)
CCALL AVERAGE TO CALCULATE GROUP CROSS SECTIONS
CALL AVERAGE(GP)SEC,CCAST,FCCNST,BFLUX,GP,NCEGPS,NG,NF,NP)
IF(GPXSEC(I).GT.C.500E+49) GO TO 103
CINITIALIZE SUM1A AND SUM2A
SUM1A=SUM2A=C.C
CBEGIN ACTIVATION CALCULATIONS
DO 11 J=1,N
TEMPS=EXP(-(C.G-3))*(1-J)*((T1*3600.0)+(T2*3600.0))/PRODH(I,J)
+1)
SUM1A=SUM1A+TEMPS
11 CONTINUE
DO 12 L=1,NDEGPS
TEMPS=SUM1A*(1.-Q-2-1)*GPXSEC(L)*GPFLUX(L)
SUM2A=SUM2A+TEMPS
12 CONTINUE
CCALCULATE ACTIVITY AT T=0
ACT=(MASS(1)*ICCD.C/TARMA5)*6.023E+23*SUM2A*(1-EXP(-0.693*3600.0*T
+1/PRODH(I,J)))
CCALCULATE ACTIVITY AS A FUNCTION OF TIME
DO 13 K=1,NCPAS
TIMACT(I,J,K)=ACT*EXP(-(STAFS(K)*Q.693*3600.0/PRODH(I,J)))
CSUM ACTIVITIES
SUMACT(K)=SUMACT(K)+TIMACT(I,J,K)
13 CONTINUE
CPRINT ACTIVITIES
103 WRITE(6,20) IDNO(I), IFIX(PRODN(I,J)), (TIMACT(I,J,K),K=1,NCPAS)
20 FORMAT(1X,1X,15,2X,15,2X,10(EB,3,1X),2X,2(EB,3,2X))
CBEGIN INTERREGATION OF THE ACTLME LIBRARY
READ(7,25) ITAPE,TARMA5
IF(IECF(1).NE.0) KFLAG=1
IF(KFLAG.EQ.1) GO TO 27
IF(LFLAG.EQ.1.AND.ITAPE.NE.(IDNO(I+1))) GO TO 27
IF(ITAPE.EQ.(IDNO(I+1))) LFLAG=1
IF(ITAPE.NE.(IDNO(I))) GO TO 27
GO TO 7
CPRINT HEADINGS
27 WRITE(6,37)(STAFS(K),K=1,NCPAS)

```

```

**SUBROUTINE ACTVAT(IT,IMACT,SYSACT,SUMACT,GPFLUX,PROONU,PRODHL,GPXSEC**
WRITE(6,14)
14 FORMAT(1H,125H TARGET ACT(T1) ACT(T2) ACT(T3) ACT(T4)
+ ACT(T5) ACT(T6) ACT(T7) ACT(T8) ACT(T9) ACT(T10) ACT(T11)
+ACT(T12))
CPRINT SUM OF ACTIVITIES DUE TO A PARTICULAR TARGET NUCLIDE
WRITE(6,15) IDNC(I), (SUMACT(K), K=1, NCPAS)
15 FORMAT(1X,5X,15,5X,10(E8.3,1X),2X,2(E8.3,2X))
DC 10 K=1, NCPAS
SYSACT(K)=SYSACT(K)+SUMACT(K)
10 CONTINUE
IF(ITAPE.EQ.IDNC(I+1)) GO TO 32
IF(FLAG.EC.1) GO TO 16
32 DO 23 JN=1, NCPAS
SUMACT(JN)=2.0
23 CONTINUE
IF(ITAPE.EQ.IDNC(I+1)) GO TO 31
IF(ITAPE.GT.IDNC(I+1)) GO TO 47
I=I+1
CALL PCSITCN
CPRINT HEADINGS
WRITE(6,37) (STAFS(K), K=1, NCPAS)
WRITE(6,11)
J=0
GO TO 24
47 WRITE(6,45) IDNC(I+1)
CALL PCSITCN
I=I+2
IF(ITAPE.GT.IDNC(I)) WRITE(6,45) IDNC(I)
IF(ITAPE.GT.IDNC(I)) I=I+1
CPRINT HEADINGS
WRITE(6,37) (STAFS(K), K=1, NCPAS)
WRITE(6,11)
J=C
GO TO 24
CPRINT HEADINGS
31 WRITE(6,37) (STAFS(K), K=1, NCPAS)
WRITE(6,11)
I=I+1
J=0
GO TO 7
CPRINT HEADINGS
16 WRITE(6,37) (STAFS(K), K=1, NCPAS)
WRITE(6,21)
21 FORMAT(1H,125H SYSTEM ACT(T1) ACT(T2) ACT(T3) ACT(T4)
+ ACT(T5) ACT(T6) ACT(T7) ACT(T8) ACT(T9) ACT(T10) ACT(T11)
+ACT(T12))
CPRINT SYSTEM ACTIVITY
WRITE(6,51) (SYSACT(K), K=1, NCPAS)
51 FORMAT(1X,15X,10(E8.3,1X),2X,2(E8.3,2X))
IF(XFLAG.EC.0) GO TO 42
RETURN
42 CALL PCSITCN
43 READ(7,25) ITAPE,TARMAS
IF(EOF(7).NE.C) RETURN

```

****SUBROUTINE ACTVAT IT IMACT, SYS ACT, SUMACT, GPFLUX, PRODN, PRODM, GPXSEC****

CALL POSITON

GC TC 43

50 CALL PCSITZ(NP)

GO TO 103

END

```

**SUBROUTINE AVERAGE(GPXSEC,CONST,FCONST,BFLUX,GP,NCEGPS,NG,NF,NP)**
SUBROUTINE AVERAGE(GPXSEC,CONST,FCONST,BFLUX,GP,NCEGPS,NG,NF,NP)
C
C THIS PROGRAM CALCULATES FLUX WEIGHTED GROUP CROSS SECTIONS BY AVERAGING
C MICROSCOPIC CROSS SECTION DATA
C
C
C ARGUMENTS IN=CONST,FCONST,BFLUX,GP,NCEGPS,NG,NF,NP
C ARGUMENTS OUT=GPXSEC
C
C
C DIMENSION GPXSEC(NCEGPS),CONST(NCEGPS),FCONST(NCEGPS),BFLUX(NCEGPS
+1,GP(NG),CS(150),EA(150),CS(201),E(201)
C INITIALIZE ARRAY GPXSEC
DO 41 JK=1,NCEGPS
GPXSEC(JK)=0.0
41 CONTINUE
CLCAD ENERGY AND CROSS SECTION DATA FROM THE ACTLME LIBRARY
READ(7,3)(EN(I),CS(I),I=1,NP)
3 FORMAT(6(E11.4))
CDETERMINE IF THE ENERGY OF THE FIRST ENERGY-CROSS SECTION PAIR READ
CFROM THE ACTLME LIBRARY IS BELOW THE GREATEST GROUP BOUNDARY
IF(EN(1)-CPING) 105,102,102
105 CALL SMGCTH(GP,AG,EA,CS,NP,E(1),CS(1),NEI)
I=1
J=1
SLMINT=0.0
CDETERMINE IF THE ENERGY AT WHICH THE CROSS SECTION WAS MEASURED IS
CBELW THE NEXT GROUP BOUNDARY
50 IF(E(1)-LT,GP(J+1)) GO TO 51
GPXSEC(J)=0.0
J=J+1
GO TO 50
52 IF(E(1)-LT,GP(J+1)) GO TO 51
CCALCULATE THE GROUP CROSS SECTION
GPXSEC(J)=SUMINT/BFLUX
IF(J.EQ.NCEGPS) GO TO 106
J=J+1
SUMINT=0.0
GO TO 52
51 IF(J.EQ.1) GO TO 1000
IF(NCEGPS.GE.3.AND.J.GT.1.AND.J.LE.(NCEGPS-NF)) GO TO 2000
GO TO 3000
1000 SUMDEL=0.0
CAVERAGE MICROSCOPIC CROSS SECTION DATA OVER THERMAL ENERGIES
WIDTH=(E(1+1)-E(1))/200.0
ZINCR=WIDTH/2.0
CPVAR=E(1)+ZINCR
DO 1001 KK=1,2/C
PARSUM=(CONST(J)*SORT(CPVAR*1.0E6)*EXP(-(CPVAR*1.0E6/0.025)))+(CS
+1(I)-CS(1))+(CPVAR-E(1))/(E(1+1)-E(1))+(CS(1)-E(1)))*WIDTH*1.0E6
SLMDEL=SUMDEL+PARSUM
CPVAR=CPVAR+WIDTH
1001 CONTINUE
SUMINT=SUMINT+SLMDEL
I=I+1
GO TO 52

```

```

**SUBROUTINE AVERAGE(GPXSEC,CONST,FCONST,BFLUX,GP,NCEGPS,NG,NF,NPI)**
2000 SLMDL=0.0
CAVERAGE MICROSCOPIC CROSS SECTION DATA OVER INTERMEDIATE ENERGIES
  WIDTH=(E1(I+1)-E1(I))/200.0
  ZINCR=WIDTH/2.0
  DPVAR=E1(I)+ZINCR
  CC 2001 KK=1,200
  PAR SUM=(CCNST(I))/(DPVAR*1.0E6)*WIDTH*1.0E6*(CS1(I+1)-CS1(I))*(DP
+VAR-E1(I))/(E1(I+1)-E1(I))+CS1(I))
  SLMCEL=SLMDL+FAR SUM
  DPVAR=DPVAR+WIDTH
2001 CONTINUE
  SUM1(I)=SUM1(I)+SUM1(I)
  -S.13
  GC TC 32
3000 SLMDL=0.0
CAVERAGE MICROSCOPIC CROSS SECTION DATA OVER FAST ENERGIES
  WIDTH=(E1(I+1)-E1(I))/200.0
  ZINCR=WIDTH/2.0
  DPVAR=E1(I)+ZINCR
  CC 3001 KK=1,200
  PAR SUM=CONST(I)*EXP((DPVAR-GP(I+1))*1.0E6/FCONST(I))*WIDTH*1.0E6*(
+CS1(I+1)-CS1(I))*(DPVAR-E1(I))/(E1(I+1)-E1(I))+CS1(I))
  SLMDL=SLMDL+FAR SUM
  DPVAR=DPVAR+WIDTH
3001 CONTINUE
  SUMINT=SUMINT+SUMCEL
  I=I+1
  GC TC 32
102 GPXSEC(1)=1.0E+50
106 READ(7,51)IDUMB
5 FCFMAT(7IX,11)
  RETURN
END

```

```

**SUBROUTINE SMOOTH(X1,NX1,X2,Y2,NX2,X3,Y3,NX3)**
SUBROUTINE SMOOTH(X1,NX1,X2,Y2,NX2,X3,Y3,NX3)
C
C THIS PROGRAM CALCULATES CROSS SECTIONS AT GROUP BOUNDARIES BY LINEAR
C INTERPOLATION
C
C ARGUMENTS IN=GE,NG,EN,CS,DP
C ARGUMENTS OUT=E1,CS1,NE1
C
C
C DIMENSION X1(1),X2(1),Y2(1),X3(1),Y3(1)
C K=1
12 IF(X2(K)-X1(1)) 20,11,13
20 K=K+1
GC TO 12
11 X3(1)=X2(K)
Y3(1)=Y2(K)
K=K+1
L=2
11=2
GC TO 14
13 IF(X2(K)-X1(1)) GO TO 11
24 11=11+1
IF(X2(K)-X1(1)) GC TO 25
GO TO 24
25 X3(1)=X2(K)
Y3(1)=Y2(K)
K=K+1
L=2
14 DO 1 I=1,NX1
5 IF(X2(K)-X1(I)) 2,3,4
2 X3(I)=X2(K)
Y3(I)=Y2(K)
L=L+1
K=K+1
IF(K-NX2) 5,5,10
3 X3(L)=X2(K)
Y3(L)=Y2(K)
L=L+1
K=K+1
IF(K-NX2) 1,1,10
4 X3(L)=X1(1)
Y3(L)=Y3(L-1)+((Y2(K)-Y3(L-1))/(X2(K)-X3(L-1)))*(X3(L)-X3(L-1))
L=L+1
1 CONTINUE
10 NX3=L-1
RETURN
END

```



```

**SUBROUTINE ADVAN(NDMODE)**
SUBROUTINE ADVAN(ACPCDEI)
C
C THIS PROGRAM IS USED FOR POSITIONING THE FILE MARKER IN THE LEVDEC
C LIBRARY
C
C
C ARGUMENTS IN-NDMODE
C ARGUMENTS CUT-ACNE
C
C
C DC 305 JB=1,NDMODE
C READ(7,202) IDLHB
202 FORMAT(38X,12)
305 CCATINUE
C RETURN
C END

```

```

**SUBROUTINE PSITCA**
SUBROUTINE PSITCA
C
C THIS PROGRAM IS USED FOR POSITIONING THE FILE MARKER IN THE ACTLEME
C LIBRARY
C
C
C ARGUMENTS IN-NONE
C ARGUMENTS CUT-ACNE
C
C
C READ(7,301) NP
301 FORMAT(7,13)
NP=INT(FLCAT(NP)/3.000+0.800)
C DC 52 ICCUT=1,API
C READ(7,531) REALDUM
531 FORMAT(E11.6)
52 CCATINUE
C READ(7,5) IDLHB
5 FORMAT(71X,11)
C RETURN
C END

```

```

**SUBROUTINE PCSIT2(NP)**
SUBROUTINE PCSIT2(NP)
C
C
C THIS PROGRAM IS USED FOR POSITIONING THE FILE MARKER IN THE ACTLENF
CLIBRARY
C
C
C   ARGUMENTS IN=NP
C   ARGUMENTS OUT=ACNE
C
C
C   NP1=INT(FLCAT(NP)/3.000+0.800)
C   DO 52 I=1,NP1
C     REAC(7,53) REALDUM
53  FORMAT(11.4)
52  CONTINUE
C   REAC(7,51) IDUME
5  FORMAT(71X,11)
C   RETURN
C   END

```

```

**SLERCUTINE CCSRTE(TIMACT, PRODN, NONUCL, NOPAS, DOSSUM, SYSDCS, IDNO, ST**
SUBROUTINE DOSRTE(TIMACT, PRODN, NONUCL, NOPAS, DOSSUM, SYSDCS, IDNO, ST
+AFS, PRODHL, DCSAUC)
C
C
C THIS PROGRAM CALCULATES DOSE RATES AS A FUNCTION OF COMPONENT COMPOSIT-
C ION, GEOMETRY, AND TIME AFTER SHUTDOWN
C
C
C ARGUMENTS IN-TIMACT, PRODN, NONUCL, NOPAS, DOSSUM, SYSDCS, IDNO, ST AFS
C ARGUMENTS I6-PRODHL, DCSNUC
C ARGUMENTS CL7-ACNE
C
C
C DIMENSION TIMACT(ACNUCL,13,ACPAS), PRODN(ACNUCL,13), DOSSUM(NOPAS,
+SYSDCS(NOPAS), IDNO(NONUCL), STAFS(NOPAS), PRODHL(NONUCL,13), DCSAUC(N
+CPAS), DCSAUC(390), NT(390), NP(390), AC(390), H(390), NDI(390), DOSEPR(39
+C), IDEM(30), IDALG(30), DLEV(30), PRCB(30)
C CMPCA, A, V, F, CE, T1, T2, N, R, D, HT, CPTION, IGEDM
C REAL MULT
C WRITE(6,100)
100 FORMAT(///,1H ,61H THE FOLLOWING OUTPUT DATA ARE CALCULATED DOSE RA
+TES IN MRAD/H,/)
C INITIALIZE ARRAYS
C DC 1 I=1,ACPAS
C SYSDCS(I)=0.0
C DCSNUC(I)=0.0
C DOSSUM(I)=0.0
1 CONTINUE
C INITIALIZE ARRAYS
C DC 101 I=1,390
C NT(I)=NP(I)=NDI(I)=0
C AC(I)=DCSPRC(I)=H(I)=DCSAUC(I)=0.0
101 CONTINUE
C INITIALIZE ARRAYS
C DC 2 J=1,13
C K=(I-1)*13+J
C NP(K)=I*13+PRCAUC(I,J)
C HK=PRODHL(I,J)
C AC(K)=TIMACT(I,J,1)
C AT(K)=IDNC(I)
2 CONTINUE
C ARRANGE TARGET ID NUMBERS, PARENT ACTIVITIES, AND PARENT HALF-LIVES IN
C ORDER OF INCREASING PARENT ID NUMBER PRIOR TO SEARCHING LEYDEC FOR
C PARENT DECAY INFORMATION
C AC=ACNUCL*13-1
C DC 3 J=1,NO
C NTEMP1=NP(J+1)
C NTEMP2=NT(J+1)
C ZTEMP3=AC(J+1)
C ZTEMP4=H(J+1)
C DC 4 K=1,J
C I=J+1-K
C IF (ATEMP1,GE,NF(I)) GO TO 5
C AP(I+1)=AP(I)
C IT(I+1)=IT(I)
C AC(I+1)=AC(I)

```

```

**SUBROUTINE DOSRTEITIMACT,PRODN,NUCL,NOPAS,DOSSUM,SYSDCS, IDNO,ST**
      F(I+1)=H(I)
4 CONTINUE
      I=C
      5 AP(I+1)=ATEMP1
      NT(I+1)=NTEMP2
      AC(I+1)=ZTEMP3
      F(I+1)=ZTEMP4
      3 CONTINUE
CSEARCH THE NP ARRAY FOR THE FIRST NON-ZERO PARENT ID NUMBER
      L=1
      7 IF(NP(L).NE.C) GO TO 205
      L=L+1
      GC TC 7
CHARK THE BEGINNING OF THE PARENT NUCLIDE LIST-THIS LIST IS IN ASCENDING
CG ORDER
      205 I=BASE=L
CBEGIN INTERROGATION OF THE LEVDEC LIBRARY FOR A PARTICULAR PARENT ID
NUMBER
      200 READ(7,16) ITAPE,HALFL,NDMGDE
      16 FORMAT(16,20X,E11.4,I3)
      IF(IECF(7).NE.0) GC TC 200
      IF(ITAPE.EQ.AP(I)) GC TC 17
      IF(ITAPE.GT.NP(L)) GO TO 280
      CALL ADVAN(INCMCDE)
      GC TC 200
      280 CALL ADVAN(INCMCDE)
      28 WRITE(6,29) AP(I)
      29 FORMAT(////,1H,12HATTENTION-PRODUCT NUCLIDE NUMBER,16,90HDOES NO
      +T EXIST IN DECAY LIBRARY-THEREFORE-NO CONTRIBUTION FROM THIS NUCL
      +IDE IN DOSE CALC)
      WRITE(6,37) NT(I)
      37 FORMAT(1H,19HTHE TARGET ATCM WAS,15,////)
      L=L+1
CDISCONTINUE IF THE LAST PARENT NUCLIDE HAS BEEN ADDRESSED
      IF(NP(L).EQ.0) GC TC 813
      IF(ITAPE.GT.NP(L)) GC TC 28
      GC TC 200
CCHECK TO SEE IF NP(L) EQUALS NP(L+1)-THIS OCCURS WHEN THE SAME PARENTS
CAFE PRODUCED FROM DIFFERENT TARGET NUCLIDES. IN THIS CASE, THE DECAY
CINFORMATION IS STORED IN VARIOUS HOLDING ARRAYS AND IS INTERROGATED AS
CMANY TIMES AS THERE ARE ADDITIONAL LIKE PARENTS. THIS NEGATES HAVING
CTC RE-READ LEVDEC FOR EACH LIKE PARENT.
      17 IF(NP(L).EQ.NP(L+1)) GC TC 800
CCOMPARE HALF LIVES
      313 IF(F(L)/HALFL.LT.0.8.OR.H(L)/HALFL.GT.1.2) GO TO 19
      IF(HALFL.EG.1.00E+30.CR.AC(L).EQ.0.0) GO TO 19
      IMDE=C
      302 IMCDE=IMCDE+1
CBEGIN INTERROGATION OF THE LEVDEC LIBRARY
      REAC(7,3CG) IDEKMD,IDAUGH,DLEVEL,MULT
      3CC *CPAT(38X,12,17,E11.4,E12.4)
CIF DECAY IS B) BETA MINUS,STORE DAUGHTER ID NUMBER
      IF(IDEKMD.EQ.9) NDIL=IDAUGH
      IF(IDEKMD.EQ.7.CP.ICEKMD.EQ.9) GO TO 301
      IF(IDEKMD.EQ.8) GC TC 310
      303 IF(IMCDE.EC.ACNCDE) GO TO 304
      GC TC 302
      201 IF(DLEVEL.EQ.C.C) GO TO 303

```

```

**SUBROUTINE DCSRTX(TIMACT, PRCONV, NONUCL, NOPAS, DOSSUM, SYSDOS, IDNO, ST**)
  CALL CURVE(OLEVL, FLUXE)
  CALL CALC(AC(IL), DCSPRD(IL), MULT, FLUXE)
  GC TO 303
210 IF(OLEVL.EQ.0.0) GO TO 311
  CALL CURVE(OLEVL, FLUXE)
  CALL CALC(AC(IL), DCSPRD(IL), MULT, FLUXE)
CIN ANY CASE ADD IN TWO 0.511 MEV ANNIHILATION GAMMAS
211 GC 312 IBETA=1.2
  CALL CURVE(C.511, FLUXE)
  CALL CALC(AC(IL), DCSPRD(IL), MULT, FLUXE)
212 CCATIALE
  GC TO 303
  19 CALL ADVAN(NCMODE)
CRRAD THE NEXT LEVCEC NUCLEIDE
204 READ(7,16) ITAPE,HALFLF,NMODE
  IF(CEC(7).NE.0) GO TO 304
  IF(ITAPE.EC.NPIL) GC TO 313
  IF(ITAPE.EQ.NP(L+1)) GO TO 21
CISCONTINUE IF THE LAST NUCLEIDE HAS BEEN ADDRESSED
  IF(INP(L+1).EC.C) GC TO 813
  CALL ADVAN(NCMODE)
  L=L+1
  GC TO 204
21 L=L+1
  GC TO 17
CLOAD DECAY INFORMATION INTO THE HOLDING ARRAYS
800 GC 801 IL=1, NCMODE
  READ(7,30C) IDEKM(IL), IDAUG(IL), OLEV(IL), PROB(IL)
801 CONTINUE
CPARK THE BEGINNING OF LIKE PARENT NUCLEIDES IN THE LIST
  IBEGIN=L
  EC2 L=L+1
  IF(INP(L).EC.NPIL+1) GO TO 802
CMARK THE END OF LIKE PARENT NUCLEIDES IN THE LIST
  IEND=L
810 GC 805 IP=IBEGIN, IEND
CDETERMINE IF A PARENT NUCLEIDE WILL BE CONSIDERED FOR DOSE RATE CALCS
  IF((IP/HALFLF.GT.C.8.AND.H(IP)/HALFLF.LT.1.2.AND.HALFLF.NE.1.0E+
    +50.AND.AC(IP).NE.0.0) GC TO 804
  GC TO 805
804 GC 820 IL=1, NCMODE
CBEGIN INTERCAGATION OF THE HOLDING ARRAYS
  IF(IDEKM(IL).EQ.9) ND(IP)=IDAUG(IL)
  IF(IDEKM(IL).EC.7.GR.IDEKM(IL).EQ.9) GO TO 501
  IF(IDEKM(IL).EC.8) GC TO 510
  GO TO 820
501 IF(OLEVL(IL).EC.0.0) GC TO 820
  CALL CURVE(OLEVL(IL), FLUXE)
  CALL CALC(AC(IP), DCSPRD(IP), PROB(IL), FLUXE)
  GC TO 820
510 IF(OLEVL(IL).EC.0.0) GC TO 511
  CALL CURVE(OLEVL(IL), FLUXE)
  CALL CALC(AC(IP), DCSPRD(IP), PROB(IL), FLUXE)
511 GC 821 IBETA=1.2
  CALL CURVE(C.511, FLUXE)
  CALL CALC(AC(IP), DCSPRD(IP), PROB(IL), FLUXE)
821 CONTINUE
820 CONTINUE

```

SUBROUTINE DOSRTEITIMACT,PRODNU,NONUCL,NOPAS,DOSSUM,SYSOCS,IGNO,ST

```

      EC5 CONTINUE
      CCETERMINE IF THE NEXT LEVDEC NUCLIDE IS THE SAME AS THE PREVIOUS ONE
      806 READ(7,16) ITAPE,HALFLF,NMODE
      IF(ECF(7).NE.C) GC TC 806
      IF(ITAPE.EQ.NP(1BEGIN)) GO TO 808
      GC TC 807
      EC8 DO 809 IL=1,ADPCOE
      REAC(7,300),IDEXN(11),DAUG(11),OLEVL(11),PROB(11)
      809 CONTINUE
      GC TC 810
      CCETERMINE IF THE NEXT LEVDEC NUCLIDE IS EQUAL TO THE NEXT PARENT ID
      CALMPER FOLLOWING THE LIST OF LIKE PARENT NUCLIDES
      EC7 IF(ITAPE.EQ.NP(IEND+1)) GO TO 811
      CALL ADVAN(NCMCEE)
      L=IEND+1
      IF(ITAPE.GT.NP(IEND+1)) GO TO 812
      GC TC 200
      811 L=IEND+1
      GO TO 17
      812 WRITE(6,29) AP(1)
      WRITE(6,37) AT(1)
      L=L+1
      IF(ITAPE.GT.NP(L)) GC TC 812
      GO TO 200
      813 PRINT 7
      CMCVE THE FILE MARKER THROUGH THE ACTLNEE LIBRARY AND POSITION AT THE
      CBEGINNING OF LEVDEC
      49 READ(7,47) ITAPE
      47 FORMAT(I6)
      IF(EDF(7).NE.0) GO TO 700
      CALL PCSITCN
      GO TO 45
      CMZK THE END OF THE LIST OF ARRANGED PARENT NUCLIDES
      700 ISTOP=L
      L=IBASE
      CBEGIN INTERROGATION OF THE LEVDEC LIBRARY FOR POSSIBLE RADIOACTIVE
      CDAUGHTERS-READ UNTIL THE ID NUMBER IS MATCHED
      701 IF(ND(L).NE.0) GO TO 702
      L=L+1
      IF(L.EQ.1STOP+1) GC TC 913
      GO TO 701
      702 READ(7,16) ITAPE,HALFLF,NMODE
      IF(ECF(7).NE.C) GC TC 702
      IF(ITAPE.EQ.ND(L)) GO TO 703
      IF(ITAPE.GT.ND(L)) GC TC 926
      CALL ADVAN(NCMCEE)
      GO TO 702
      526 L=L+1
      IF(L.EQ.1STOP+1) GC TC 913
      IF(ITAPE.EQ.ND(L)) GO TO 703
      CALL ADVAN(NCMCEE)
      GO TO 701
      CCHECK THE HALF LIFE OF THE DAUGHTER
      703 IF(HALFLF.GT.1.EE+48) GC TC 704
      GO TO 705
      704 CALL ADVAN(NCMCEE)
      604 L=L+1
      IF(L.EQ.1STOP+1) GO TO 913

```

```

**SUBRELTIME CCSRTEITIMACT,PRCDNU,NCNUCL,NCPAS,DCSSUM,SYSDCS,IDNO,ST**

GC TO 701
705 IMODE=0
502 IMODE=IMODE+1
C BEGIN INTERPGATION OF THE LEVDEC LIBRARY FOR DAUGHTER DECAY INFO
      REAC(7,707) IDEKMC,DLEVEL,MULT
707 FORMAT(38),12,7X,E11.4,E12.4)
      IF(IDEKMC.EQ.7.CR.IDEKM0.EQ.9) GO TO 901
      IF(IDEKMC.EQ.8) GC TC 910
503 IF(IMODE.EQ.NMODE) GO TO 904
      GC TC 902
501 IF(DLEVEL.EQ.C.C) GC TC 903
      CALL CURVE(DLEVEL,FLUXE)
      CALL CALC(AC(L),DCSDAU(L),MULT,FLUXE)
      GC TC 903
910 IF(DLEVEL.EQ.0.0) GO TO 911
      CALL CURVE(DLEVEL,FLUXE)
      CALL CALC(AC(L),DCSDAU(L),MULT,FLUXE)
C1A ANY CASE ACC. IN TWO C.S. ALL MEV ANNILATION GAMMAS
911 DC 912 IBETA=1.2
      CALL CURVE(0.511,FLUXE)
      CALL CALC(AC(L),DCSDAU(L),MULT,FLUXE)
512 CONTINUE
      GC TO 903
513 REWIND 7
      NC=NCNUCL+13
      DO 915 I=1,NONUCL
C PRINT HEADINGS
      WRITE(6,81) STAFS(K),K=1,NCPAS
8  FORMAT(1,1H,4X,7H TIME(H),4X,10(E8.3,1X),2X,2(E8.3,2X))
      WRITE(6,400)
400 FORMAT(1H,125# TARGET PRODUCT CSR(T1) DSR(T2) DSR(T3) DSR(T4)
+ DSR(T5) DSR(T6) DSR(T7) DSR(T8) DSR(T9) DSR(T10) DSR(T11)
+ CSR(T12))
      DC 818 L=IBASE,AC
      IF(INT(L).EQ.IDNO(I)) GO TO 817
      GC TC 818
CSUM THE DECAY RATES DUE TO PARENTS AND DAUGHTERS
817 DO 819 K=1,NCPAS
      DCSSUM(K)=DCSFR(L)*EXP(-(STAFS(K)*0.693*3600.0/H(L))+DCSDAU(L)*E
+XP(-(STAFS(K)*0.693*3600.0/H(L)))
      CCSSUM(K)=DCSSUM(K)+DCSSUM(K)
      SYSDCS(K)=SYSDCS(K)+DCSSUM(K)
819 CONTINUE
C PRINT DECAY RATES
      WRITE(6,201) IDAC(I),NP(L),IDCSNUC(K),K=1,NCPAS
201 FORMAT(1X,1X,15,2X,15,2X,10(E8.3,1X),2X,2(E8.3,2X))
818 CONTINUE
C PRINT HEADINGS
      WRITE(6,81) STAFS(K),K=1,NCPAS
      WRITE(6,101)
101 FORMAT(1H,125# TARGET DSR(T1) DSR(T2) DSR(T3) DSR(T4)
+ DSR(T5) DSR(T6) DSR(T7) DSR(T8) DSR(T9) DSR(T10) DSR(T11)
+ CSR(T12))
      WRITE(6,231) IDAC(I),IDCSSUM(K),K=1,NCPAS

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```

**SUBROUTINE CCSRTE(ITIMACT,PROONU,NONUCL,NOPAS,DOSSUM,SYSDDS,IONO,ST**
23 FORMAT(1X,5X,15,5X,10(E8.3,1X),2X,2(E8.3,2X))
DC B22 K=1,NOPAS
CCSSUM(K)=0.0
B22 CONTINUE
B15 CONTINUE
WRITE(6,8)ISTAFS(K),K=1,NOPAS)
CPRINT SYSTEM HEADINGS
WRITE(6,2E)
26 FORMAT(1H,125H SYSTEM DSR(T1) DSR(T2) DSR(T3) DSR(T4)
+ DSR(T5) DSR(T6) DSR(T7) DSR(T8) DSR(T9) DSR(T10) DSR(T11)
+ DSP(I,12))
CPRINT SYSTEM CASE
WRITE(6,51)(SYSDCS(K),K=1,NOPAS)
51 FORMAT(1X,15X,10(E8.3,1X),2X,2(E8.3,2X))
RETURN
END

```



```

**SUBROUTINE CALC(ACT,DSPRD,MULT,FLUXE)**
SLBRETIME CALC(ACT,DSPRD,MULT,FLUXE)
C
C THIS PROGRAM CALCULATES DCSE RATE FOR 3 DIFFERENT GEOMETRIES
C
C ARGUMENTS IN-ACT,MULT,FLUXE
C ARGUMENTS OUT-USPRD
C
C CMPCN A,V,F,CF,T1,T2,N,R,D,HT,OPTION,IGECM
REAL MULT
GO TO(31,32,33) IGECM
31 DCSEAT=(ACT*MULT*2.5/(4.0*3.1416*FLUXE*(D*100.0)**2))
DSFRD=CSPRD+DCSEAT
RETURN
C THE ABOVE DCSE RATE CALCULATION IS FOR A POINT SOURCE
32 VOL=4*3.1416*((R*100.0)**3)/3
S=(ACT*MULT*2.5)/(VOL*FLUXE*4.0)
DCSEAT=S*(2*R*(R*D*10000.0-D*(2*R*D)*10000.0*(ALOG((2*R*D)/D)))/(
+R*D)*100.0)
DSFRD=CSPRD+DCSEAT
RETURN
C THE ABOVE DCSE RATE CALCULATION IS FOR A HOMOGENEOUS NON-ABSORBING
CSPHERE
33 VOL=3.1416*((R*100.0)**2)*HT*100.0
S=(ACT*MULT*2.5)/(VOL*FLUXE*4.0)
DCSEAT=S*((HT+C*100.0*(ALOG(1+(R**2)/((HT*D)**2))+2*R*ATAN((HT*D)
+R)/(HT*D))-C*100.0*(ALOG(1+(R**2)/(D**2))+2*R*ATAN(D/R)/D))
DSFRD=CSPRD+DCSEAT
RETURN
C THE ABOVE CALCULATION IS FOR A HOMOGENEOUS NON-ABSORBING CYLINDER WHERE
C THE OBSERVER FACES THE AXIS
END

```

```

**SUBROUTINE CURVE(E,FLUXE)**
SUBROUTINE CURVE(E,FLUXE)
C
C
CGIVEN THE ENERGY OF THE GAMMA RAY,THIS PROGRAM CALCULATES THE PARTICLE
CFLUX(QUANTA/CM**2*SEC) THAT CORRESPONDS TO A DCE RATE OF 2.5 NR/HR --
CFLUXE
C
C
C ARGUMENTS IN-CLEVEL
C ARGUMENTS OUT-FLUXE
C
C
IF(E.GE.C.C1C.AND.E.LE.0.020) GO TO 1
IF(E.GT.0.020.AND.E.LE.0.040) GO TO 2
IF(E.GT.0.040.AND.E.LE.0.050) GO TO 3
IF(E.GT.0.050.AND.E.LE.0.060) GO TO 4
IF(E.GT.0.060.AND.E.LE.0.080) GO TO 5
IF(E.GT.0.080.AND.E.LE.0.150) GO TO 6
IF(E.GT.0.150.AND.E.LE.0.400) GO TO 7
IF(E.GT.0.400.AND.E.LE.0.600) GO TO 8
IF(E.GT.0.600.AND.E.LE.1.000) GO TO 9
IF(E.GT.1.000.AND.E.LE.1.500) GO TO 10
IF(E.GT.1.500.AND.E.LE.3.000) GO TO 11
IF(E.GT.3.000.AND.E.LE.6.000) GO TO 12
IF(E.GT.6.000.AND.E.LE.10.000) GO TO 13
FLUXE=1.0E+99
RETURN
1 FLUXE=21495330.64*E**2.1759447
RETURN
2 FLUXE=11460233.22*E**2.0136022
RETURN
3 FLUXE=1036515.975*E**1.2698653
RETURN
4 FLUXE=96497.38032*E**0.4772399
RETURN
5 FLUXE=11225.12221*E**(-0.2874410)
RETURN
6 FLUXE=554E2.78326*EXP(E*(-10.9132065))
RETURN
7 FLUXE=1330.569205*E**(-1.1036085)
RETURN
8 FLUXE=1462.198546*E**(-1.0004513)
RETURN
9 FLUXE=1546.557758*E**(-0.8891058)
RETURN
10 FLUXE=1545.05132*E**(-0.7797676)
RETURN
11 FLUXE=1508.205757*E**(-0.7188784)
RETURN
12 FLUXE=1468.362150*E**(-0.7051622)
RETURN
13 FLUXE=1656.852E74*E**(-0.7649757)
RETURN
END

```

DETERMINATION OF THE DOSE RATE FROM NEUTRON ACTIVATION OF THE COPPER BEAM DUMP

THE FOLLOWING DATA HAS BEEN ENTERED INTO MEMORY

AMPERES	KILO-VOLTS	BEAM FRACTION	DUTY FACTOR	T1	T2	N
.650E+02	.170E+03	.100E+01	.100E+00	.800E+01	.160E+02	7

NUMBER OF ENERGY GROUPS-NEGRPS	20
NUMBER OF TARGET NUCLIDES-NNUCL	2
NUMBER OF POINTS AFTER SHUTDOWN-NGPAS	12
RADIUS R OF SPHERE OR CYLINDER (IN)	.215E+00
HEIGHT HT OF CYLINDER (IN)	0.
DISTANCE D FROM POINT SOURCE OR SURFACE OF SPHERE OR CYLINDER (IN)	.200E+01
NUMBER OF GROUPS BOUNDARIES-NG	21
NUMBER OF FAST NEUTRON GROUPS-NF	2

SPECIFIC TIMES AFTER SHUTDOWN (H)

TIME 1	0.
TIME 2	.100E+C1
TIME 3	.200E+C1
TIME 4	.300E+C1
TIME 5	.400E+C1
TIME 6	.500E+C1
TIME 7	.600E+C1
TIME 8	.700E+C1
TIME 9	.800E+C1
TIME 10	.900E+C1
TIME 11	.100E+C2
TIME 12	.200E+C2

UNIT NORMALIZED FLUXES

GROUP 1	.7531E-C5
GROUP 2	.1382E-C6
GROUP 3	.8333E-C7
GROUP 4	.1058E-C6
GROUP 5	.1378E-C6
GROUP 6	.2842E-C7
GROUP 7	.4790E-C7
GROUP 8	.3860E-C7
GROUP 9	.7420E-C8
GROUP 10	.3522E-C7
GROUP 11	.4570E-C7
GROUP 12	.4021E-C7
GROUP 13	.3043E-C7
GROUP 14	.7234E-C7
GROUP 15	.4684E-C7
GROUP 16	.3011E-C7
GROUP 17	.3551E-C7
GROUP 18	.3456E-C7
GROUP 19	.3076E-C7
GROUP 20	.8288E-C6

GROUP BOUNDARIES (MEV)

BOUNDARY	1	.2385E+01
BOUNDARY	2	.2307E+01
BOUNDARY	3	.1827E+01
BOUNDARY	4	.1108E+01
BOUNDARY	5	.5502E+00
BOUNDARY	6	.1576E+00
BOUNDARY	7	.1111E+00
BOUNDARY	8	.1248E-01
BOUNDARY	9	.2479E-01
BOUNDARY	10	.2198E-01
BOUNDARY	11	.1033E-01
BOUNDARY	12	.3355E-02
BOUNDARY	13	.1234E-02
BOUNDARY	14	.5629E-02
BOUNDARY	15	.1013E-03
BOUNDARY	16	.2902E-04
BOUNDARY	17	.1068E-04
BOUNDARY	18	.3059E-05
BOUNDARY	19	.1125E-05
BOUNDARY	20	.4140E-06
BOUNDARY	21	.1000E-10

NAME ID-NUMBER MASS(KG)

CCPFR63	29063	.258E+03
CCPFR65	29065	.115E+03

THE FOLLOWING HAS BEEN CALCULATED BY ACDD51

INSTANTANEOUS NUMBER OF NEUTRONS PRODUCED PER SECOND .7645E+13

AVERAGE NUMBER OF NEUTRONS PRODUCED PER SECOND .7645E+12

INTEGRALS OF THE FITTING FUNCTION OVER THE GROUP INTERVALS
GROUP LOWER BOUNDARY UPPER BOUNDARY REQUIRED INTEGRAL CALCULATED INTEGRAL

20	.100E-04	.414E+00	.829E-06	.829E-06
19	.414E+00	.113E+01	.308E-07	.308E-07
18	.113E+01	.306E+01	.344E-07	.344E-07
17	.306E+01	.107E+02	.395E-07	.395E-07
16	.107E+02	.250E+02	.301E-07	.301E-07
15	.290E+02	.101E+03	.468E-07	.468E-07

14	.1C1E+C3	.583E+C3	.723E-07	.723E-07
13	.583E+03	.123E+C4	.304E-C7	.304E-07
12	.123E+04	.336E+04	.402E-07	.402E-07
11	.336E+C4	.1C3E+C5	.457E-C7	.457E-07
10	.1C3E+05	.219E+C5	.352E-C7	.352E-07
9	.219E+05	.248E+C5	.742E-08	.742E-08
8	.248E+C5	.525E+C5	.3E4E-C7	.386E-07
7	.525E+05	.111E+06	.479E-C7	.479E-07
6	.111E+C6	.158E+06	.284E-07	.284E-07
5	.158E+C6	.55CE+C6	.138E-06	.138E-06
4	.55CE+06	.111E+07	.1C6E-C6	.1C6E-06
3	.111E+C7	.183E+07	.833E-07	.833E-07
2	.183E+07	.231E+07	.126E-C6	.138E-06
1	.231E+07	.239E+07	.790E-05	.790E-05

THE FOLLOWING OUTPUT DATA ARE CALCULATED ACTIVITIES IN HQ

TIME(H)	PRODUCT	ACT(T1)	ACT(T2)	ACT(T3)	ACT(T4)	ACT(T5)	ACT(T6)	ACT(T7)	ACT(T8)
29063	29062	.0	.0	.0	.0	.0	.0	.0	.0
29063	28062	.0	.0	.0	.0	.0	.0	.0	.0
29063	27062	.0	.0	.0	.0	.0	.0	.0	.0
29063	28063	.176E+04	.176E+04	.176E+04	.176E+04	.176E+04	.176E+04	.176E+04	.176E+04
29063	27060	.0	.0	.0	.0	.0	.0	.0	.0
29063	27060	.0	.0	.0	.0	.0	.0	.0	.0
29063	29064	.431E+10	.409E+10	.387E+10	.367E+10	.347E+10	.329E+10	.312E+10	.295E+10

TIME(H)	TARGET	ACT(T1)	ACT(T2)	ACT(T3)	ACT(T4)	ACT(T5)	ACT(T6)	ACT(T7)	ACT(T8)
29063		.431E+10	.409E+10	.387E+10	.367E+10	.347E+10	.329E+10	.312E+10	.295E+10

ACT(T9)	ACT(T10)	ACT(T11)	ACT(T12)
.0	.0	.0	.0
.0	.0	.0	.0
.0	.0	.0	.0
.176E+04	.176E+04	.176E+04	.176E+04
.0	.0	.0	.0
.0	.0	.0	.0
.280E+10	.265E+10	.251E+10	.146E+10

ACT(T9)	ACT(T10)	ACT(T11)	ACT(T12)
.280E+10	.265E+10	.251E+10	.146E+10

TIME(H)	.0	.100E+01	.200E+01	.300E+01	.400E+01	.500E+01	.600E+01	.700E+01
TARGET PRODUCT	ACT(T1)	ACT(T2)	ACT(T3)	ACT(T4)	ACT(T5)	ACT(T6)	ACT(T7)	ACT(T8)
29065 29064	.0	.0	.0	.0	.0	.0	.0	.0
29065 29063	.0	.0	.0	.0	.0	.0	.0	.0
29065 28064	.0	.0	.0	.0	.0	.0	.0	.0
29065 27061	.0	.0	.0	.0	.0	.0	.0	.0
29065 28065	.0	.0	.0	.0	.0	.0	.0	.0
29065 27062	.0	.0	.0	.0	.0	.0	.0	.0
29065 27062	.0	.0	.0	.0	.0	.0	.0	.0
29065 25066	.148E+10	.426E+06	.123E+03	.353E-01	.102E-04	.293E-08	.842E-12	.242E-15

TIME(H)	.0	.100E+01	.200E+01	.300E+01	.400E+01	.500E+01	.600E+01	.700E+01
TARGET	ACT(T1)	ACT(T2)	ACT(T3)	ACT(T4)	ACT(T5)	ACT(T6)	ACT(T7)	ACT(T8)
29065	.148E+10	.426E+06	.123E+03	.353E-01	.102E-04	.293E-08	.842E-12	.242E-15

TIME(H)	.0	.100E+01	.200E+01	.300E+01	.400E+01	.500E+01	.600E+01	.700E+01
SYSTEM	ACT(T1)	ACT(T2)	ACT(T3)	ACT(T4)	ACT(T5)	ACT(T6)	ACT(T7)	ACT(T8)
	.576E+10	.405E+10	.387E+10	.367E+10	.347E+10	.329E+10	.312E+10	.295E+10

.800E+01	.900E+01	.100E+02	.200E+02
ACT(T9)	ACT(T10)	ACT(T11)	ACT(T12)
.0	.0	.0	.0
.0	.0	.0	.0
.0	.0	.0	.0
.0	.0	.0	.0
.0	.0	.0	.0
.0	.0	.0	.0
.698E-19	.201E-22	.578E-26	.226E-61

.800E+01	.900E+01	.100E+02	.200E+02
ACT(T9)	ACT(T10)	ACT(T11)	ACT(T12)
.698E-19	.201E-22	.578E-26	.226E-61

.800E+01	.900E+01	.100E+02	.200E+02
ACT(T9)	ACT(T10)	ACT(T11)	ACT(T12)
.280E+10	.265E+10	.251E+10	.146E+10

THE FOLLOWING OUTPUT DATA ARE CALCULATED DOSE RATES IN MRAD/H

TIME(H)	.0	.100E+01	.200E+01	.300E+01	.400E+01	.500E+01	.600E+01	.700E+01
TARGET	PRDCLCT	DSR(T1)	DSR(T2)	DSR(T3)	DSR(T4)	DSR(T5)	DSR(T6)	DSR(T7)
29063	27C59	.0	.0	.0	.0	.0	.0	.0
29063	27060	.0	.0	.0	.0	.0	.0	.0
29063	27C60	.0	.0	.0	.0	.0	.0	.0
29063	28062	.0	.0	.0	.0	.0	.0	.0
29063	28063	.0	.0	.0	.0	.0	.0	.0
29063	29C62	.0	.0	.0	.0	.0	.0	.0
29063	29C64	.234E+01	.221E+01	.210E+01	.199E+01	.188E+01	.178E+01	.169E+01

TIME(H)	.0	.100E+01	.200E+01	.300E+01	.400E+01	.500E+01	.600E+01	.700E+01
TARGET	DSR(T1)	DSR(T2)	DSR(T3)	DSR(T4)	DSR(T5)	DSR(T6)	DSR(T7)	DSR(T8)
29C63	.234E+01	.221E+01	.210E+01	.199E+01	.188E+01	.178E+01	.169E+01	.160E+01

.800E+01	.900E+01	.100E+02	.200E+02
DSR(T9)	DSR(T10)	DSR(T11)	DSR(T12)
.0	.0	.0	.0
.0	.0	.0	.0
.0	.0	.0	.0
.0	.0	.0	.0
.0	.0	.0	.0
.0	.0	.0	.0
.151E+01	.143E+01	.136E+01	.791E+00

.800E+01	.900E+01	.100E+02	.200E+02
DSR(T9)	DSR(T10)	DSR(T11)	DSR(T12)
.151E+01	.143E+01	.136E+01	.791E+00

TIME(H)	.0	.100E+01	.200E+01	.300E+01	.400E+01	.500E+01	.600E+01	.700E+01
TARGET PRODUCT	OSR(T1)	OSR(T2)	OSR(T3)	OSR(T4)	OSR(T5)	OSR(T6)	OSR(T7)	OSR(T8)
29065 27061	.0	.0	.0	.0	.0	.0	.0	.0
29065 27062	.0	.0	.0	.0	.0	.0	.0	.0
29065 27062	.0	.0	.0	.0	.0	.0	.0	.0
29065 28064	.0	.0	.0	.0	.0	.0	.0	.0
29065 28065	.0	.0	.0	.0	.0	.0	.0	.0
29065 29063	.0	.0	.0	.0	.0	.0	.0	.0
29065 29064	.0	.0	.0	.0	.0	.0	.0	.0
29065 29066	.332E+00	.556E-04	.275E-07	.792E-11	.228E-14	.656E-18	.189E-21	.544E-25

TIME(H)	.0	.100E+01	.200E+01	.300E+01	.400E+01	.500E+01	.600E+01	.700E+01
TARGET	OSR(T1)	OSR(T2)	OSR(T3)	OSR(T4)	OSR(T5)	OSR(T6)	OSR(T7)	OSR(T8)
29065	.332E+00	.556E-04	.275E-07	.792E-11	.228E-14	.656E-18	.189E-21	.544E-25

TIME(H)	.0	.100E+01	.200E+01	.300E+01	.400E+01	.500E+01	.600E+01	.700E+01
SYSTEM	OSR(T1)	OSR(T2)	OSR(T3)	OSR(T4)	OSR(T5)	OSR(T6)	OSR(T7)	OSR(T8)
	.267E+01	.221E+01	.210E+01	.199E+01	.188E+01	.178E+01	.165E+01	.160E+01

.800E+01	.900E+01	.100E+02	.200E+02
OSR(T9)	OSR(T10)	OSR(T11)	OSR(T12)
.0	.0	.0	.0
.0	.0	.0	.0
.0	.0	.0	.0
.0	.0	.0	.0
.0	.0	.0	.0
.0	.0	.0	.0
.0	.0	.0	.0
.157E-28	.451E-32	.130E-35	.508E-71

.800E+01	.900E+01	.100E+02	.200E+02
OSR(T9)	OSR(T10)	OSR(T11)	OSR(T12)
.157E-28	.451E-32	.130E-35	.508E-71

.800E+01	.900E+01	.100E+02	.200E+02
OSR(T9)	OSR(T10)	OSR(T11)	OSR(T12)
.151E+01	.143E+01	.136E+01	.791E+00

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