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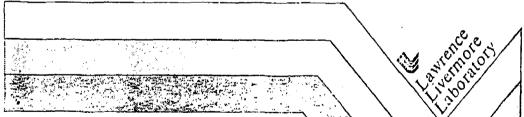
RADIONUCLIDE TRANSPORT IN SANDSTONES WITH WIPP BRINE

H. C. Weed, F. Bazan, J. Fontanilla J. Garrison, J. Rego, and A. M. Winslow

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ABSTRACT

Retardation factors (R) have been measured for the transport of 3H, 95mTc, and 85Sr in WIPP brine using St. Peter, Berea, Kayenta, and San Felipe sandstone cores. If tritium is assumed to have R=1, 95mTc has R=1.0 to 1.3 and therefore is essentially not retarded. Strontium-85 has R = 1.0 to 1.3 on St. Peter, Berca, and Kayenta, but R=3 on San Felipe. This is attributed to sorption on the matrix material of San Felipe, which has 45 volume Z matrix compared with 1-10 volume % for the others. Retardation factors (R_s) for ^{85}Sr calculated from static sorption measurements are unity for all the sandstones. Therefore, the static and transport results for 85Sr disagree in the case of San Felipe, but agree for St. Peter, Berea, and Kayenta.

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INTRODUCTION

The work described in this report was performed under Task 3 of the Waste/Rock Interaction Technology (WRIT) program during FY 80. An important goal of the program is to provide defensible models of radionuclide transport behavior in the vicinity of waste repositories. Since an experimental data base is required for these models, this study is designed to directly measure radionuclide retardation factors for the data base and to compare them with those calculated from static sorption ratio measurements on the same rock/solution systems. Sandstones and (WIPP) Waste Isolation Pilot Plant Brine B were chosen because salt deposits suitable for repositories usually have sandstone aquifers associated with them. Other recent core transport measurements have been reported on dolomite (Lynch, 1979) granite, and tuff (Daniels et al., 1980).

EXPERIMENTAL WORK

Starting Materials

The solution used in the transport measurements was WIPP Brine B diluted with deionized water to 80% of the original concentration. Table 1 shows the composition of the diluted brine. Before use in the core transport experiments it was filtered through a 0.2µ m pore size Nuclepore filter and then stored under laboratory atmosphere. No attempt was made to control Eh or pH.

TABLE 1

Coppound	<u>C, g/1</u>	Compound	<u>C, g/1</u>
KaCl	229.6	Na2B407 10H20	0.013
Na 504	5.0	NaHCO3	0.011
CaCl ₂	2.07	KI	0.010
NaBr	0.42	FeCl 6H20	0.008
KC1	0.023	RbCl	0.00053
SrCl ₂ 6H ₂ 0	0.036		

COMPOSITION OF BO% WIPP BRINE B SOLUTION FOR CORE TRANSPORT STUDIES

The sandstones used were St. Peter, Berea, and Kay-nta quartz arenites, and San Felipe subarkose. They are described briefly elsewhere (Coles et al. 1980) on the basis of X-ray diffraction data on powders. Core samples and the powdered samples used for static sorption measurements were prepared from the same blocks of material in order to minimize the effects of rock variability. They were right circular cylinders 25 mm diameter x 25.5 mm length. Spike solutions contained 85 Sr, 95m Tc, and 3 H. They were prepared by pipetting appropriate amounts of the individual radionuclide stock solutions into a volumetric flask, evaporating them to dryness under an IR lamp with a gentle stream of filtered, compressed air, and then redissolving them in the 80% WIPP Brine B. Tritium was added as a water solution before final dilution to volume with the brine. An aliquot of the solution was evaporated to dryness on a filter paper and then X - counted in the LLNL Nuclear Chemistry Division's X - analysis facility in order to determine the 95m Tc and 85 Sr concentrations; 3 H concentrations were determined by scintillation counting on the distillate from another aliquot of the spike solution.

Apparatus

The basic design of the transport apparatus has been described elsewhere (Coles et al. 1980). This design was used for the present study with only minor changes. A Camberra 30 pulse height analyzer and Nal (T1) well counter were used for 95mTc and 85Sr analyses, and a Packard Tri-Carb liquid scintillation counter for 3H analyses.

Procedure

The following procedure was used for experimental runs: confining pressure on the sample was set at 10 MPa, steady state flow was established, 0.5 cm³ of spike solution was injected into the inlet stream, the effluent radioauclide concentration was monitored as a function of time, and the run was continued until it appeared that at least one of the radioauclides (usually, 95mTc) had been completely transported through the core. After the run the core was blown dry with compressed air, and the dry core was removed from the sample holder and jacket and impregnated with EPON/TEFA resin. The core was sectioned perpendicular to the core axis, and δ - counted. Also, the

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parts of the disassembled system were rinsed with $6\underline{M}$ HNO₃ and the rinse liquid \eth - counted. Unfortunately, it was later found that the rinse liquids and all the cores except San Felipe were not counted long enough to get good statistics. The results cannot be used to determine the concentration profile or the amount of activity left in the system. The San Felipe core sections were counted in the LLNL Nuclear Chemistry Division's \eth - counting facility so that the total amount of $\85 Sr activity in them is known within 1.3% and the very small amount of 95Tc activity within 9%. Effluent samples containing the 95mTc peak were distilled to remove the brine salts and analyzed for $\3 H by scintillation counting of the distillate. Runs were performed at 25°C.

DISCUSSION AND CONCLUSIONS

Experimental results for the effluent solutions are expressed in terms of C/v_0 vs. N_e, where C/C_0 is the ratio of radionuclide concentrations in the effluent solution and the original spike solution, and N_e is the number of static pore volumes of effluent. In order to calculate the retardation factors, we have used the analytic solution of the one-dimensional advection-dispersion equation (Scheidegger 1974) for constant pore velocity v, an injected pulse of negligible width compared with the transport path, and an infinite homogeneous medium. This is a moving Gaussian of the form

$$C(x,t') = \frac{(const.)}{\sqrt{Dt'}} \exp \left[\frac{-(x-vt')^2}{4Dt'} \right]$$
(1)

where x is the axial coordinate, t' = t/R is the time scale for a nuclide with retardation factor R, t is the time, and D is the dispersion coefficient. We normalize C(x,t') so that the integral over volume is equal to the number of injected tracer atoms. Rewriting Equation (1) in terms of the number of pore volumes and applying it at the space position x = L, the exit face of the sample, we get an asymmetric function of pore volume N':

$$\begin{pmatrix} C \\ C \\ O \end{pmatrix} = \underbrace{B}_{\sqrt{DN'}} \exp \left[- \left(\underbrace{Lv}_{4D} \frac{(1-N')^2}{N'} \right) \right]$$
(2)

where

B = a normalization factor with a theoretical value of B = $(V_0/V_g)\sqrt{Lv/4\Pi}$, V_0 = the initial spike volume, V_g = the static pore volume.

Equation (2) has a single maximum with a steep rise and a more gadual fall according to the value of the Peclet number (Lv/D).

Differentiating Equation (2) with respect to N' in order to locate the maximum in (C/C_{n}) and rearranging:

$$N'_{max} = \sqrt{(D/Lv)^2 + 1} - (D/Lv).$$
(3)

For two nuclides, one unretarded and the other retarded, the ratio of pore volumes is

$$\frac{N}{N'} = \frac{vt/L}{vt'/L} = \frac{t}{t'} = R.$$
(4)

N is related to experimental quantities by:

where

 $N_e = V/V_s$, V = the total volume of effluent, $N_r =$ the reference value of N_e which corresponds to the origin of the theoretical curve at N' = 0 in Equation (2).

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Substituting (3) and (5) into (4),

$$R = \frac{(N_{e max} - N_{r})}{\sqrt{(D/Lv)^{2} + 1 - (D/Lv)}}.$$
 (6)

In this treatment it is assumed that the dynamic pore volume is equal to the static pore volume. Although the flow inside the core is not strictly one-dimensional, Equation (2) can be fitted to the data, and R calculated from Equation (6). Equation (2) was fitted to the data by means of the routine CURVE on the Livermore Time-Sharing System. This routine displays data points superimposed on a graph of Equation (2) and the quantities \mathbf{B}_{p} , \mathbf{N}_{p} and \mathbf{D} are adjusted to fit the graph to the data by inspection. Assuming that R = 1 for a reference nuclide, (usually 3 H), R for other nuclides can be calculated from Equation (6) using values of N e max, N and D from the curvefitting routine.

The estimated error in the retardation factors is 20%. The estimated experimental error in the (C/C_p) data for Sr and Tc due to weighing of samples, & - counting, and the uncertainty in the value of C is between 8% and 11% in the peak regions. The estimated experimental error for 3 H data is 10% to 20%. No 3 H data are reported for San Felipe sandstone since they are too irregular to serve as a reference standard for ^{95m} Tc and ⁸⁵Sr. Although some problems exist in the $\frac{3}{4}$ analysis of the brine solutions, we are using $\frac{3}{11}$ as a reference nuclide because it is widely accepted as a standard (Payme 1972). We believe that the experimental errors affect the peak heights but not the peak locations (N). Similar considerations apply to Tc and Sr; though some of the experiments show low fractional recovery of the nuclides, the main effect is on the peak heights rather that the peak locations (N $_{\rm e}$ max). Values of the retardation factors and other parameters are shown in Table 2.

TABLE 2.

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VALUES OF FITTING PARAMETERS AND RETARDATION FACTORS FOR ³H. ^{95m}Tc, AND ⁸⁵Sr IN SANDSTONES WITH WIPP BRINE B

Sandsten¢	Nuclide	Time, mín	10 ⁴ v cn/s	16 ⁴ D cm ² /s	10 ³ Be	Nr	N _{e max}	R(a)	F.E.
St. Peter 10 3H Ø = .139 95mTc p = 2.24 85Sr		10107	4.69	2.6	2.2	3.92	4.73	1.0	.7
	95aTc				2.3	3.86	4.71	1.1	1
	85 _{Sr}				2.0	3.74	4.79	1.3	1
St. Peter 4	Peter 4 3 _H	7215	24,3	30.	5.3	4.42	5.05	1.0	.9
Ø = .139	95mTc			50.	4.7	4.65	5.16	1.1	.9
	85 _{Sr}			50.	3.5	4.70	5.26	1.2	.7
Kayenta	3 _H	2859	3.39	.80	1.2	5.35	6.25	1.0	.8
Ø = .201	95¤Tc				1.2	5.27	6.22	1.0	.7
p = 2.13	85 _{Sr}				1.2	5.20	6.25	1.2	.8
Berea	3 _H	2634	1.08	.16	.64	. 58	1.52	1.0	.5
£ = .198	95arc				.76	.46	1.42	1.0	.8
	85 _{Sr}				.80	. 3/.	1.45	1.2	.9
San Felipe	= .118 85 _{Sr}	3938	1.43	0.6	.74	.95	1.93	1.2	.5
g = .118 $\rho = 2.42$.20	.45	3.07	3	.5

Ø = fractional porosity based on static pore volume
 P = bulk density, g/cm³
 F.E. = fraction of spike activity recovered in effluent.
 (a) Estimated error 20%

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In Figures 1-5, the curves are calculated from Equation (2) and the values of R, B_{e} , N_{r} , and D are in Table 2. The vertical arrows indicate the time of spike injection. The delay between injection and appearance of the leading edge of the peak is due to dead volume effects in the system. The retardation factors (R) for 95mTc and 85Sr vary from 1.0 to 1.3, so that they are essentially unretarded relative to 3 H in St. Peter, Kayenta, and Borea sandstones. In the case of St. Peter 10 sandstone the curves for all three nuclides lie close together (Figure 1). The ³H peak was not sampled completely which caused the low fractional recovery in the effluent. In Figure 2, the effluent curves for St. Peter 4 sandstone have a steeper rise and a more gradual fall than those for St. Peter 10. This is due to the higher pore velocity and higher dispersion coefficient required to fit the theoretical curves to the data. In Figure 3, the effluent curves for Kayenta sandstone agree very closely with each other and are more symmetrical than those for the St. Peter sandstones. The curves for Berea sendstone in Figure 4 are fitted to the data at the rise and fall of the peak since the data at N = 1.5 appear anomalo:sly low compared with the rest. The apparent delay of the ³H peak relative to the others is .07 to 0.1 pore volume, which is considered to be within the limits of experimental error for N.

There is no obvious correlation between the retardation factors and mineralogical characteristics of St. Peter, Kayenta and Berea sandstones. St. Peter has I volume % matrix material, while Kayenta and Berea have 10 volume %, the rest being mostly quartz. Technetium-95m is considered to be unretarded because of the scarcity of anion adsorbers in rock materials. Strontium-85 may be unretarded because of occupation of Sr sorption sites by the Sr in the WIPP Brine B, or because the matrix material has a low selectivity ratio K^{Sr}₁₀.

In the case of San Felipe sandstone, 95mTc shows little retardation (R = 1.2) but 85Sr behaves very differently. It shows definite retardation (R = 3) and the theoretical curve fits the data quite will below about 5.5 pore volumes, but thereafter the data and the curve diverge. About 11% of the originally injected 85Sr is distributed along the core with an essentially uniform profile, but there is less that 1% 95mTc. Since the core has 45 volume % matrix matrial, the observed retardation may be due to irreversible sorption

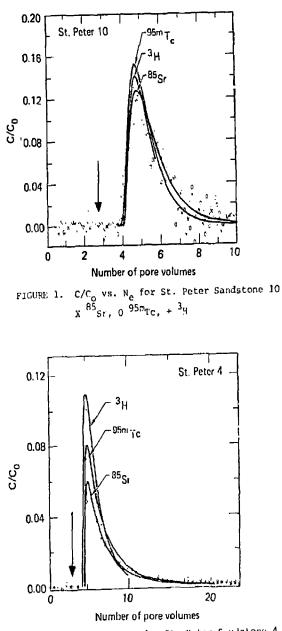


FIGURE 2. C/C_o vs. N_e for St. Peter Sumdstone 4

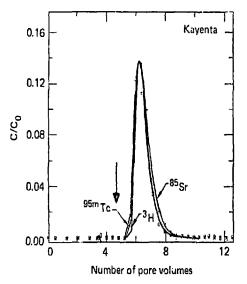


FIGURE 3. C/C_0 vs. N_e for Kayenta Sandstone

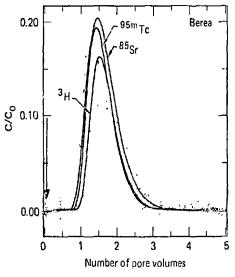


FIGURE 4. C/C_0 vs. W_p for Berea Sandstone

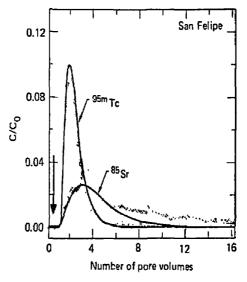


FIGURE 5. C/C_0 vs. N_e for San Felipe Sandstone

of the 85 Sr by the clay in the matrix. Also, if small amounts of colloidal clay particles transport 85 Sr in the solution, this may account for the additional 85 Sr appearing in the effluent after the initial peak. We have no definite explanation for the low fractional recovery of 95m Tc and 85 Sr. Some possible causes are incomplete filling of the spike loop, or leakage at the spike value or high pressure inlet fittings.

Retardation factors for 85 Sr calculated from static sorption measurements in Coles et al. (1980) on the sandstones in this study all have the value R_s = 1. Agreement between core transport and static sorption values is good for all the sandstones except for San Felipe, where the static sorption value is R_e = 1 compared with R = 3 from core transport measurements. Also, the static surption experiments give no indication of the complexities in the behavior of 85 Sr shown in Figure 5.

ACKNOWLEDCMENTS

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